

EICUG
Electron-Ion Collider User Group Meeting

2019 JULY 22-26
PARIS
École Nationale Supérieure de Chimie

*The world's most powerful
microscope for studying the
"glue" that binds the building
blocks of visible matter*



Summary Talk of Parallel Sessions Theory Talks, Tuesday & Thursday

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FUNDAÇÃO DE AMPARO À PESQUISA
DO ESTADO DE SÃO PAULO

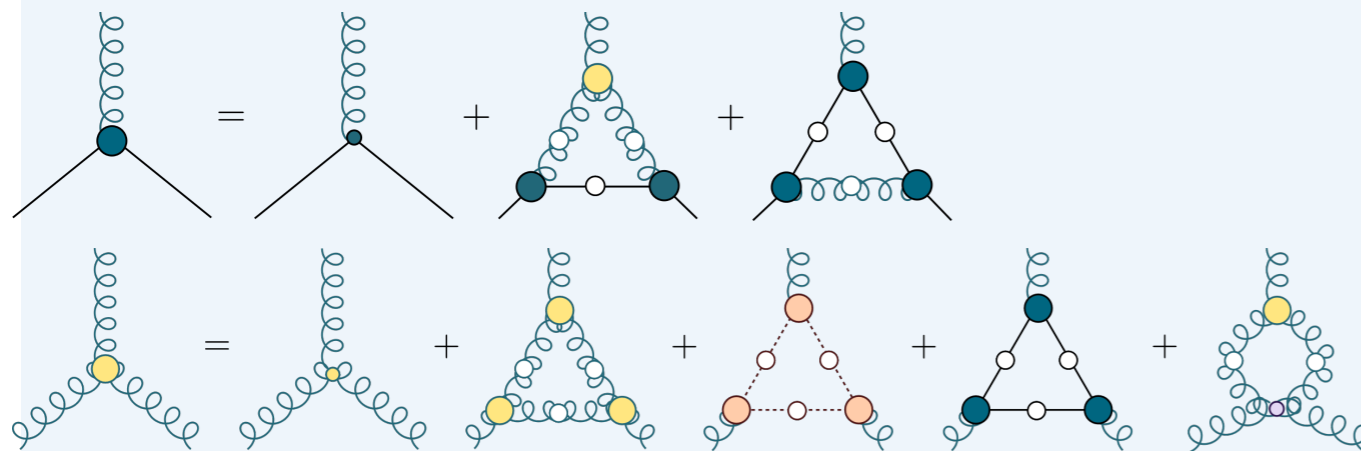
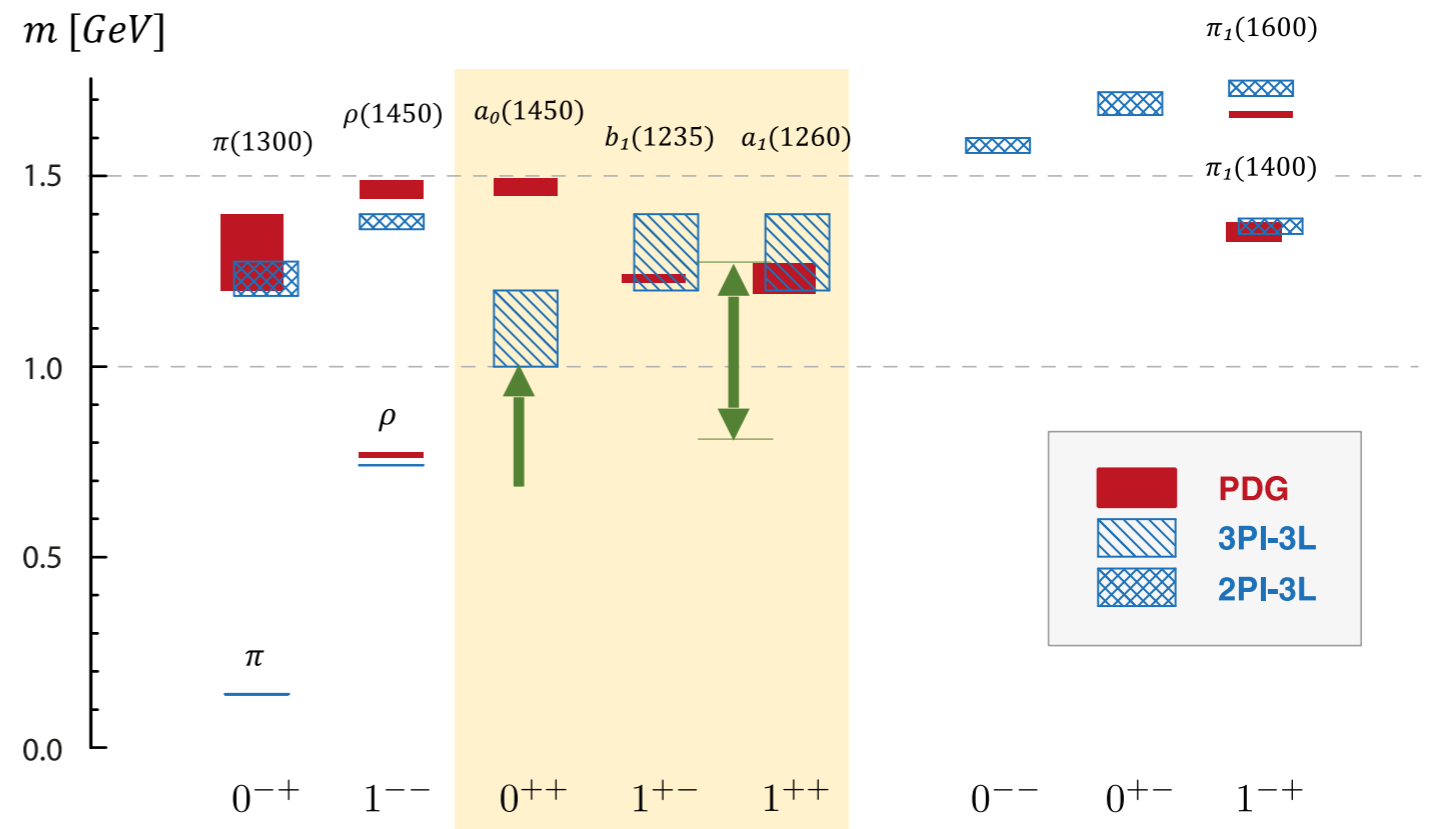


Conselho Nacional de Desenvolvimento
Científico e Tecnológico

Richard Williams — “Topics in QCD from Dyson-Schwinger Equations”

Switching from 2PI-2loop (Rainbow-Ladder) to 3PI-3loop

- Corrects $\rho - a_1$ splitting and degeneracy of axial-vectors.
- Pushes lightest $q\bar{q}$ scalar above 1 GeV.



Interaction vertices explicit d.o.f in 3PI approach.

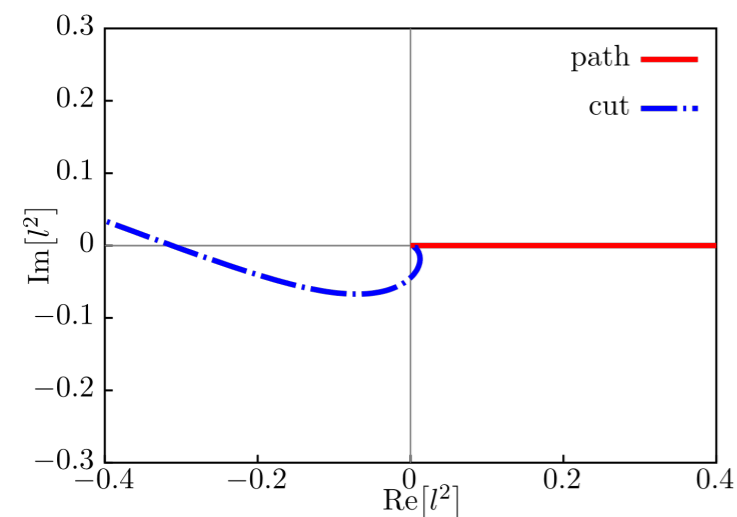
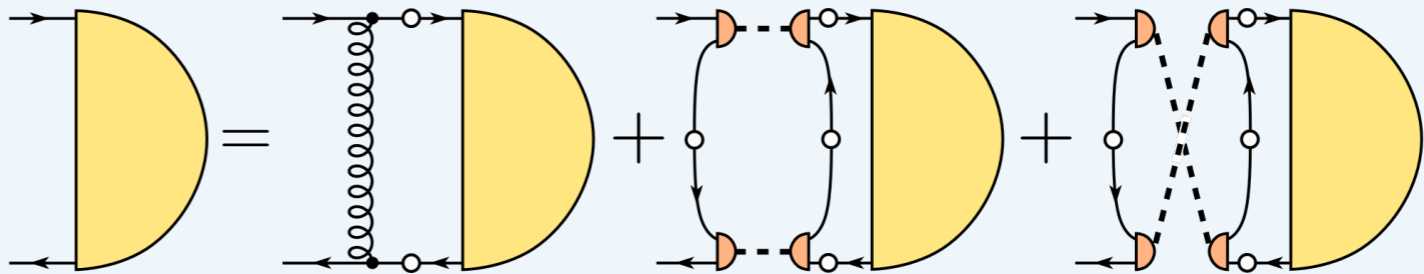
Less modelling.

Parameters are coupling g and the quark masses m_i .

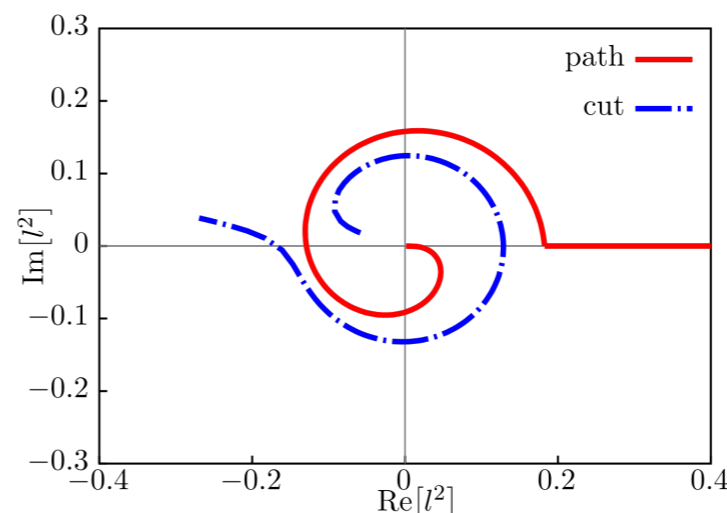
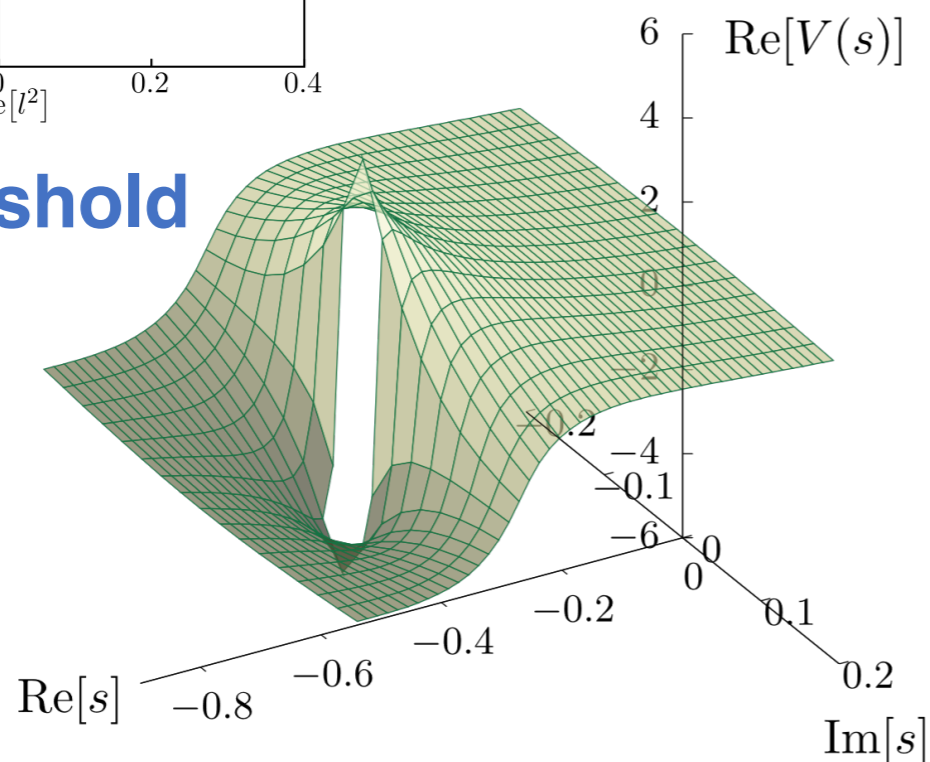
Richard Williams — “Topics in QCD from Dyson-Schwinger Equations”

Bound-state resonances:

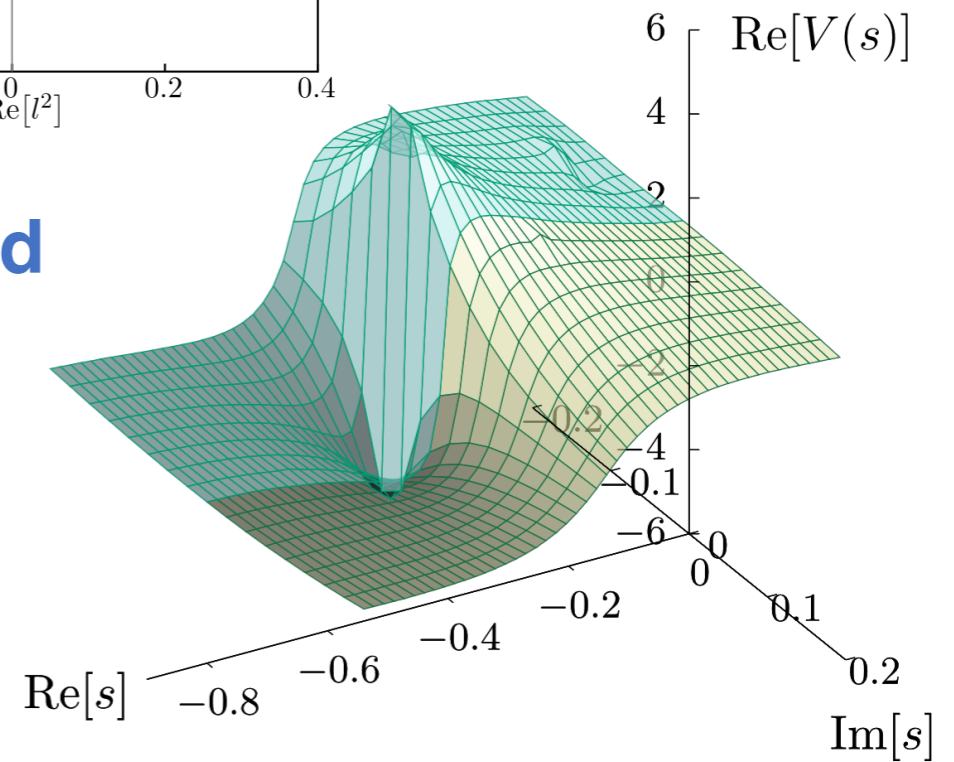
Including decay channels requires integration over intermediate particle poles



Below threshold

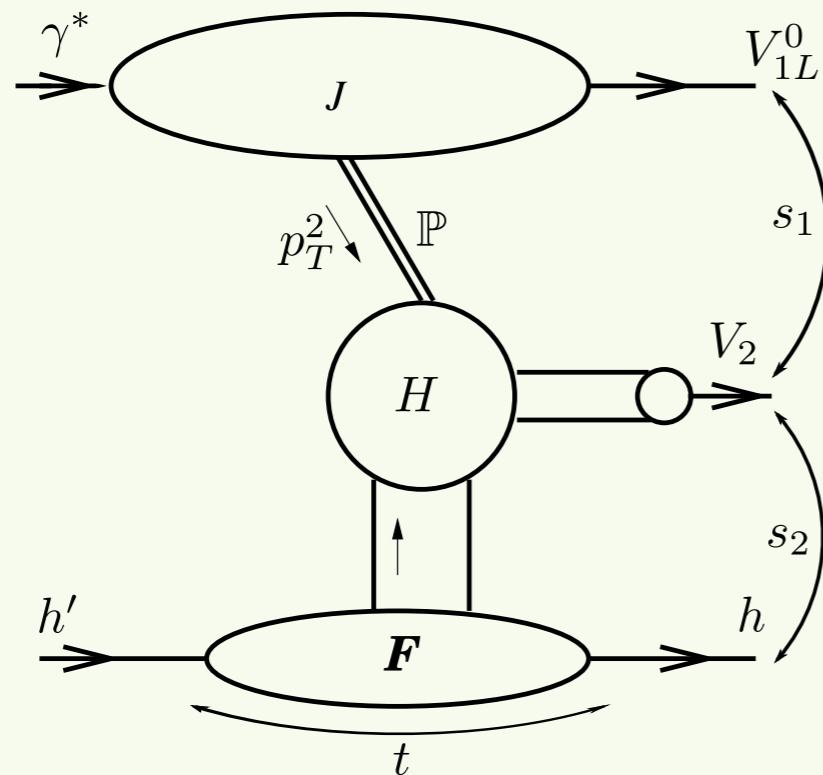


Above threshold



(Euclidean) bound-state pole found in 2nd Riemann sheet, as expected.

Diffractive electroproduction of two vector mesons



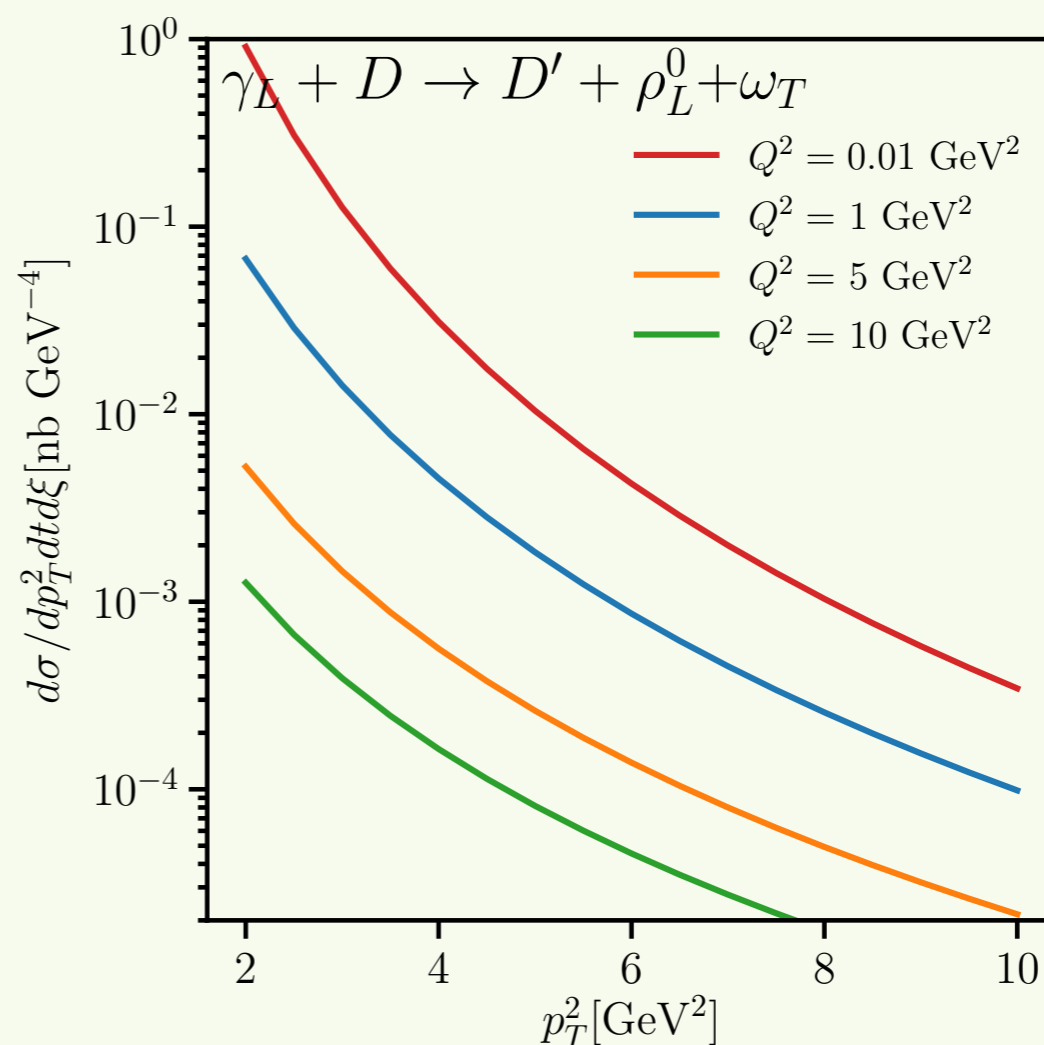
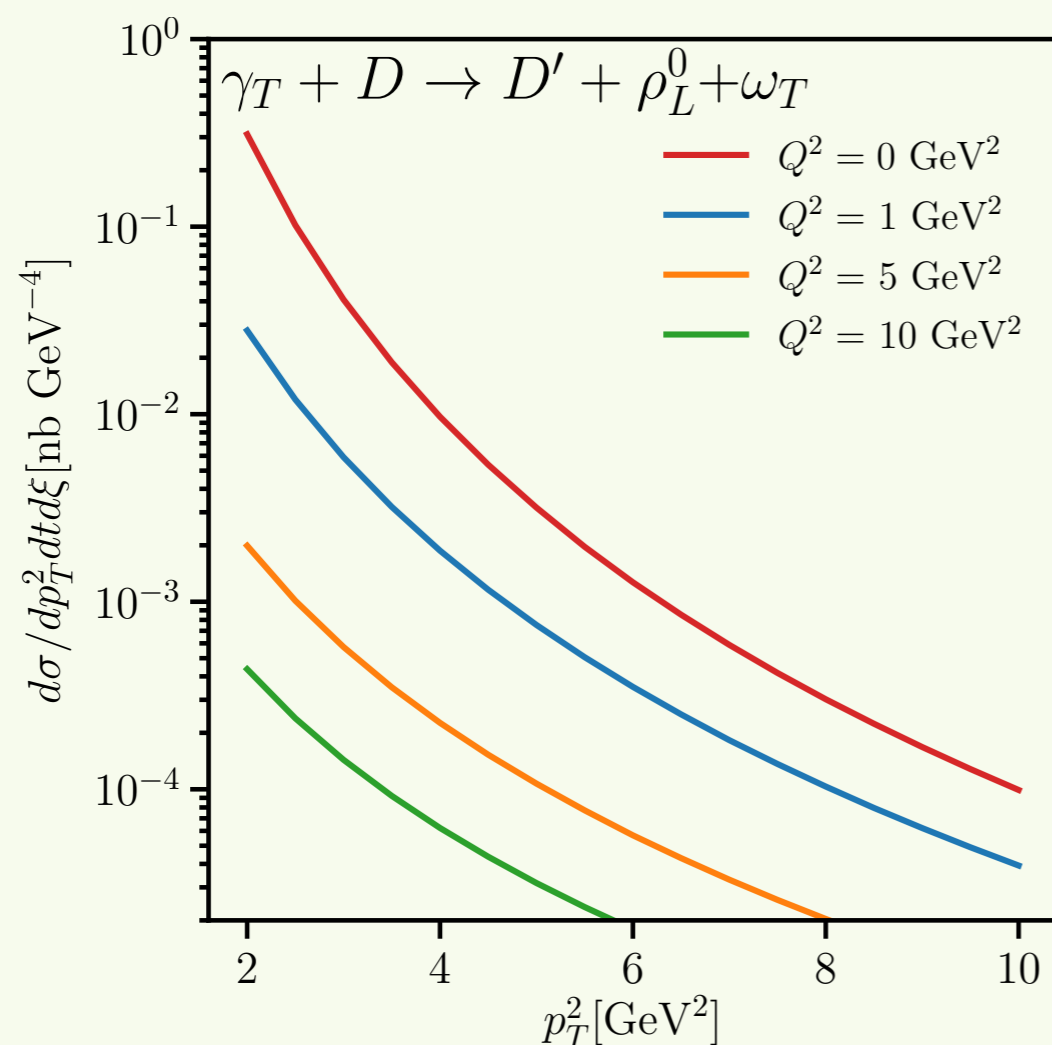
- Two vector mesons are separated by large rapidity gap
- Ordering of scales: invariant mass of two vector mesons $s_1 \gg \gg$ hadron- (ρ^+ / ω) invariant mass $s_2 \sim p^2 \gg \gg \Lambda_{\text{QCD}}$
- Pomeron \rightarrow two gluon exchange
- No gluon contribution as Pomeron is C-even
- Polarization V_2 determines chiral even or odd GPD entering

- All the hard work already done for the nucleon in
D.Yu. Ivanov, B. Pire, L. Szymanowski, PLB550 '02
R. Enberg, B. Pire, L. Szymanowski, EPJC47 '06
- Extension for deuteron straightforward: ω and ϕ ($s - \bar{s}$) production

Wim Cosyn — “Probing transversely GPDs in diffractive electroproduction on the proton and deuteron at an EIC”

$$\gamma_{L/T} d \rightarrow d \rho_L^0 \omega_T^0 \quad \text{PRELIMINARY}$$

- Chiral odd deuteron GPD
- Calculation at $\xi_D = 0.15$, $t \approx t_{\min} = -0.33 \text{ GeV}^2$



GPDs theoretical constraints



- Polynomiality:

Lorentz Covariance

- Positivity:

Positivity of Hilbert space norms

- Support:

Relativistic Quantum Mechanics

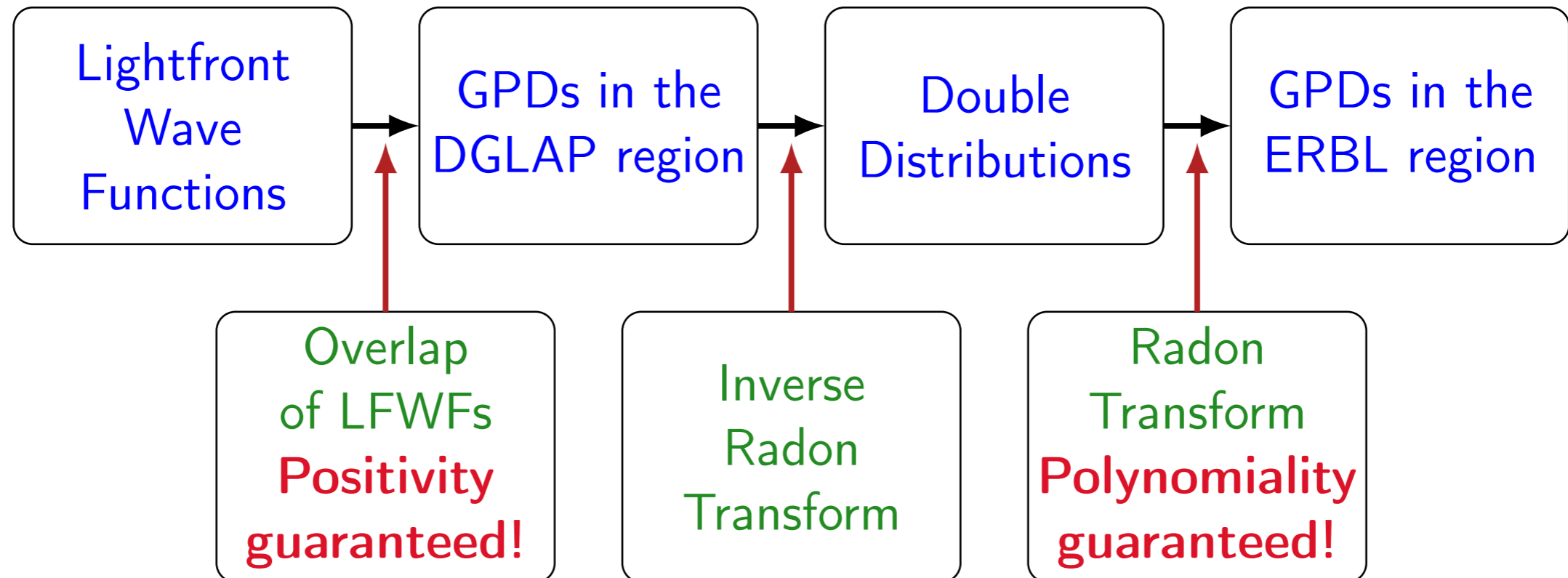
- Soft Pion theorem

PCAC and Axial-Vector WTI

Problems

There is no model (until now) fulfilling a priori all these constraints
People usually choose to fulfil either **polynomiality** or **positivity**
→ violation of either Lorentz covariance or Hilbert space properties

A new modelling procedure



Both positivity and polynomiality are fulfilled

Eur.Phys.J. C77 (2017) no.12, 906

Phys.Lett. B780 (2018) 287-293

Cédric Lorcé — “Mechanical Properties of Hadrons”

Bare QCD energy-momentum tensor

$$T^{\mu\nu} = \bar{\psi} \gamma^\mu \frac{i}{2} \overleftrightarrow{D}^\nu \psi - G^{a\mu\alpha} G^{a\nu}{}_\alpha + \frac{1}{4} \eta^{\mu\nu} G^2$$

Renormalized trace of the QCD EMT

Mass, spin and pressure all encoded in

$$T^\mu{}_\mu = \underbrace{\frac{\beta(g)}{2g} G^2}_{\text{Trace anomaly}} + (1 + \gamma_m) \bar{\psi} \underbrace{m}_{\text{Quark mass matrix}} \psi$$

$$T^{\mu\nu} = \begin{bmatrix} \text{Energy density} & & & \\ T^{00} & T^{01} & T^{02} & T^{03} \\ T^{10} & T^{11} & T^{12} & T^{13} \\ T^{20} & T^{21} & T^{22} & T^{23} \\ T^{30} & T^{31} & T^{32} & T^{33} \\ & \text{Energy flux} & \text{Momentum flux} & \end{bmatrix}$$

Shear stress
Normal stress (pressure)

- Key information about the structure encoded in EMT
- Breit frame defines the spatial distribution at rest
- Mechanical properties can be expressed in terms of spatial moments

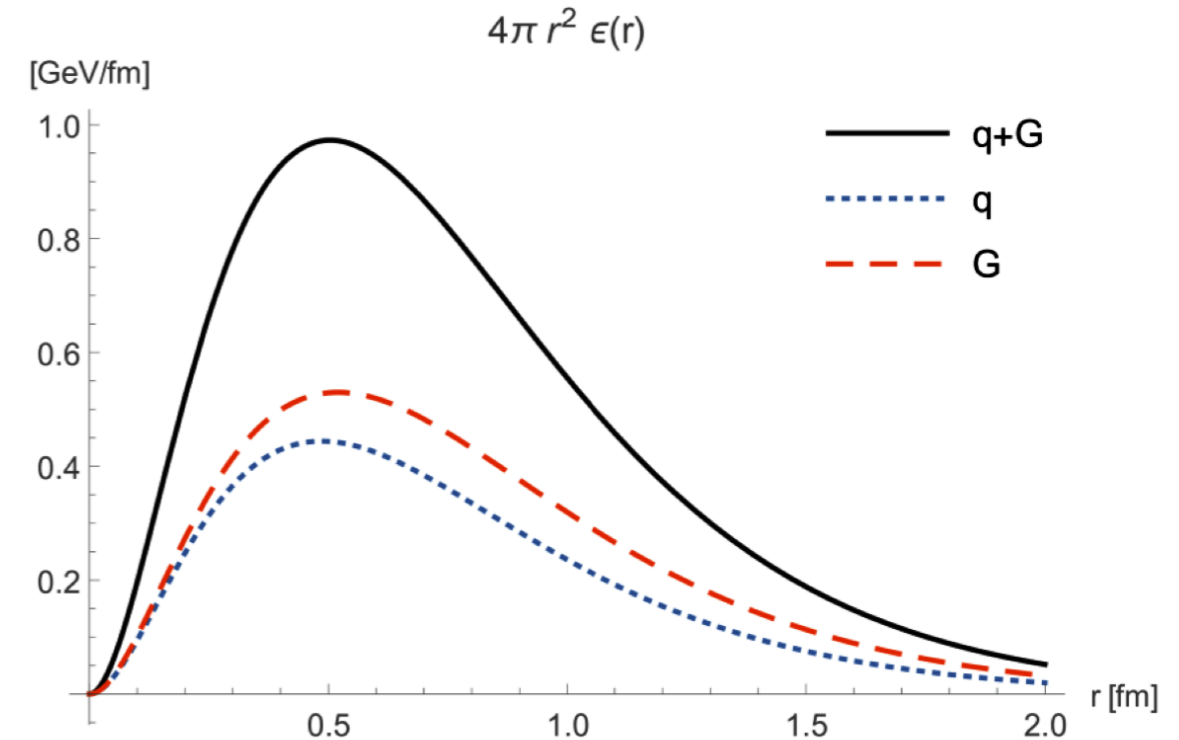
Mass, pressure, moment of inertia, ...

Cédric Lorcé — “Mechanical Properties of Hadrons”

Spin-1/2

$$\langle p', s' | T_a^{\mu\nu}(0) | p, s \rangle = \bar{u}(p', s') \Gamma_a^{\mu\nu}(P, \Delta) u(p, s)$$

$$\Gamma_a^{\mu\nu}(P, \Delta) = \frac{P^\mu P^\nu}{M} A_a(t) + \frac{\Delta^\mu \Delta^\nu - g^{\mu\nu} \Delta^2}{M} C_a(t) + M g^{\mu\nu} \bar{C}_a(t) \\ + \frac{P^{\{\mu} i\sigma^{\nu\}\lambda} \Delta_\lambda}{4M} J_a(t) + \frac{P^{[\mu} i\sigma^{\nu]\lambda} \Delta_\lambda}{4M} D_a(t)$$



$$\sqrt{\langle r^2 \rangle_M} = 0.91 \text{ fm}$$

$$\sqrt{\langle r^2 \rangle_Q} = 0.84 - 0.88 \text{ fm}$$

Energy density

$$\varepsilon_a = [A_a(0) + \bar{C}_a(0)] \frac{M}{V} \quad \Rightarrow \quad U_i = \varepsilon_i V$$

$$M = \underbrace{U_q}_{\sim 44\%} + \underbrace{U_g}_{\sim 56\%}$$

$$p_q = -p_g$$

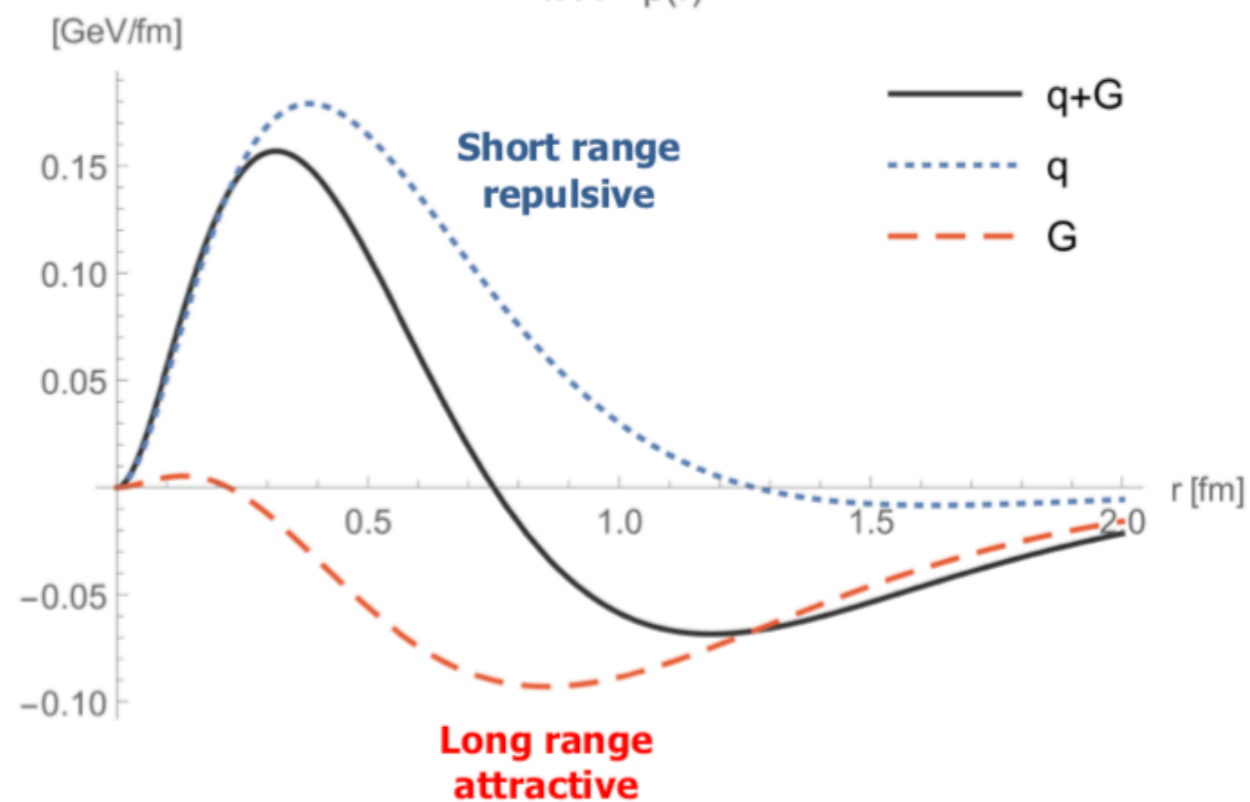
$$\mu = 2 \text{ GeV}$$

$$\sim 11\%$$

Cédric Lorcé — “Mechanical Properties of Hadrons”

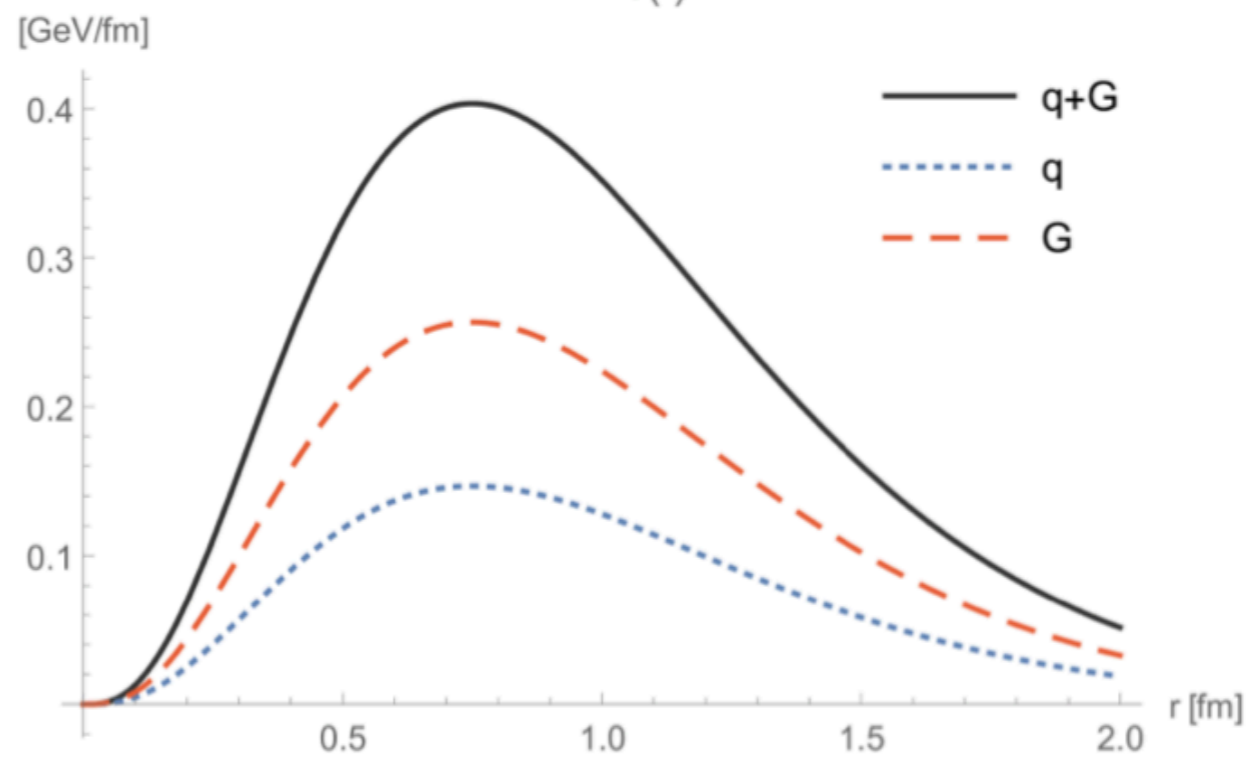
Radial pressure

$$4\pi r^2 \rho(r)$$



Pressure anisotropy

$$4\pi r^2 s(r)$$



Adam Freese — “Gravitational Form Factors of Mesons and Proton”

Gravitational form factors and the EMT

- **Gravitational form factors** encode information in the energy-momentum tensor (EMT):

$$\langle p' \lambda' | T_{\mu\nu}^q(0) | p \lambda \rangle = \bar{u}^{\lambda'}(p') \left[\gamma_{\{\mu} P_{\nu\}} A^q(t) + \frac{i P_{\{\mu} \sigma_{\nu\}} \Delta}{2m_N} B^q(t) + \frac{\Delta_\mu \Delta_\nu - \Delta^2 g_{\mu\nu}}{4m_N} C^q(t) \right. \\ \left. + m_N g_{\mu\nu} \bar{c}^q(t) + \frac{i P_{[\mu} \sigma_{\nu]} \Delta}{2m_N} D^q(t) \right] u^\lambda(p)$$

- Three of the GFFs accessible through GPDs:

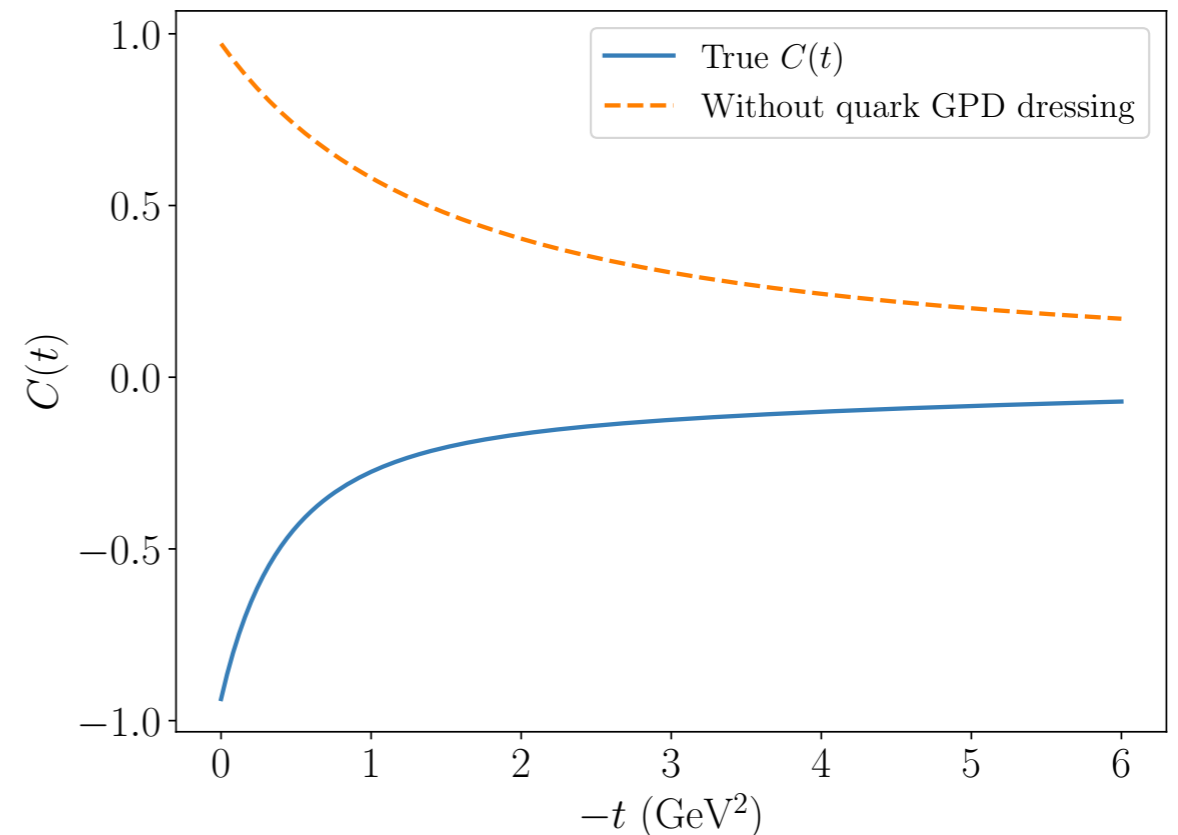
$$\int dx x H^q(x, \xi, t) = A^q(t) + \xi^2 C^q(t), \quad \int dx x E^q(x, \xi, t) = B^q(t) - \xi^2 C^q(t)$$

- GFFs encode information about distribution of **energy**, **angular momentum**, and **forces**.
- **Cannot do graviton-exchange experiments**—but GPDs allow GFFs to be experimentally accessed! (DVCS, DVMP, ...)

Adam Freese — “Gravitational Form Factors of Mesons and Proton”

$C(t)$ encodes **pressure distribution** (see Peter Schweitzer’s talk)

- $C(0)$ not constrained by conservation laws
- $C(0) < 0$ is a stability condition
- $C(t)$ is affected by quark GPD dressing, but not $A(t)$ or $B(t)$



- $C(0) = -0.94 < 0$ in the NJL model
- Ignoring quark GPD dressing gives $C(0) = 0.97 > 0$
- Proton stability appears to **require** dressing the quark GPD!

Peter Schweitzer — “D-Term of Nucleon”

D-term of nucleon

Peter Schweitzer, UConn

- energy-momentum tensor $T^{\mu\nu}$ fundamentally important

nucleon matrix elements of $T^{\mu\nu} \rightarrow$ 3 EMT form factors $A(t)$, $J(t)$, $D(t)$

constraints: **mass** $\Leftrightarrow A(0) = 1 \Leftrightarrow$ quarks + gluons carry 100% of nucleon momentum

spin $\Leftrightarrow J(0) = \frac{1}{2} \Leftrightarrow$ quarks + gluons carry 100% of nucleon spin

D-term $\Leftrightarrow D(0) \equiv D \rightarrow$ unconstrained! **Last global unknown!**

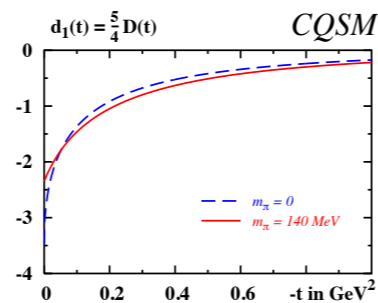
- **theory:** free fermion $D = 0$ i.e. D-term 100% dynamic effect!
Hudson, PS PRD97 (2018) 056003

- **interaction** $\rightarrow D \neq 0$

chiral quark soliton

Goeke et al,
PRD75 (2007) 094021

agrees with lattice & disp. relation



- **experiment**

$\pi^0 D^Q \approx -0.7$ from Belle data

on $\gamma\gamma^* \rightarrow \pi^0\pi^0$ in e^+e^-

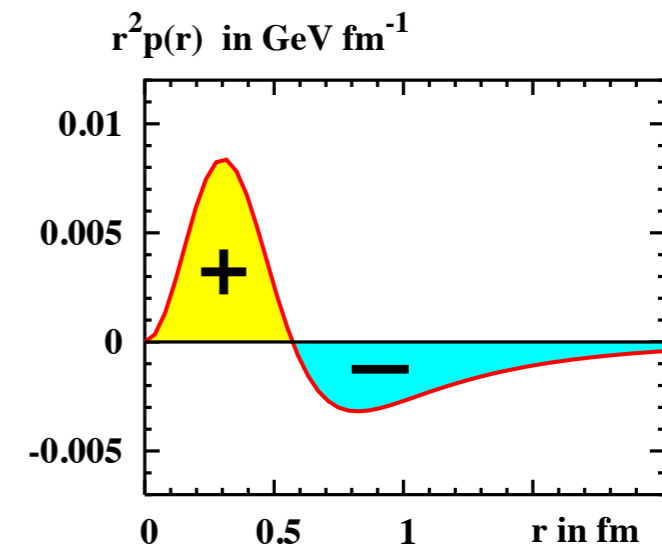
Kumano, Song, Teryaev, PRD97 (2018)

nucleon D from DVCS data from subtraction term in fixed- t dispersion relation of Re, Im part of DVCS amplitude (accessible *under assumptions* at this point)

Burkert, Elouadrhiri, Girod, **Nature** 557, 396 (2018)

K. Kumerički, **Nature** 570, 7759 (2019)

- **explore scale dependence at EIC**
extract $D(t)$ (!) (vs $A(t)$ and $J(t)$ much harder)
learn about forces inside hadrons!!



- **interpretation**

M.V.Polyakov (2003)

shows how internal forces balance related to D -term

stability $\int_0^\infty dr r^2 p(r) = 0$

$D = 4\pi m \int_0^\infty dr r^4 p(r) < 0$

review:

M.V.Polyakov, PS,

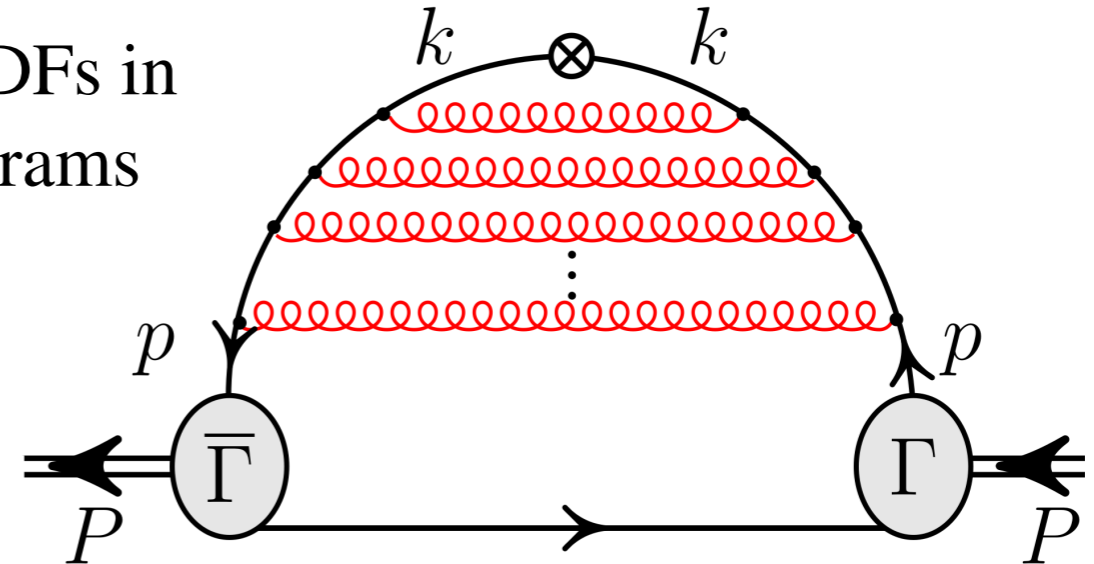
Int.J.Mod.Phys.A 33, 1830025 (2018)

[arXiv:1805.06596]

Pion PDFs – Self-Consistent DSE Calculations

- To self-consistently determine hadron PDFs in rainbow-ladder must sum all planar diagrams

$$q(x) \propto \text{Tr} \int \frac{d^4 p}{(2\pi)^4} \bar{\Gamma}_M(p, P) S(p) \times \Gamma_q(x, p, n) S(p) \Gamma_M(p, P) S(p - P)$$



- DSEs are formulated in Euclidean space – evaluate $q(x)$ by taking moments
- The *hadron dependent* vertex $\Gamma_q(x, p, n)$ satisfies an inhomogeneous BSE
- However can define a *hadron independent* vertex $\Lambda_q(x, p, n)$

$$\Gamma_q(x, p, n) = \iint dy dz \delta(x - yz) \delta\left(y - \frac{p \cdot n}{P \cdot n}\right) \Lambda_q(z, p, n)$$

- $\Lambda_q(x, p, n)$ satisfies the inhomogeneous BSE

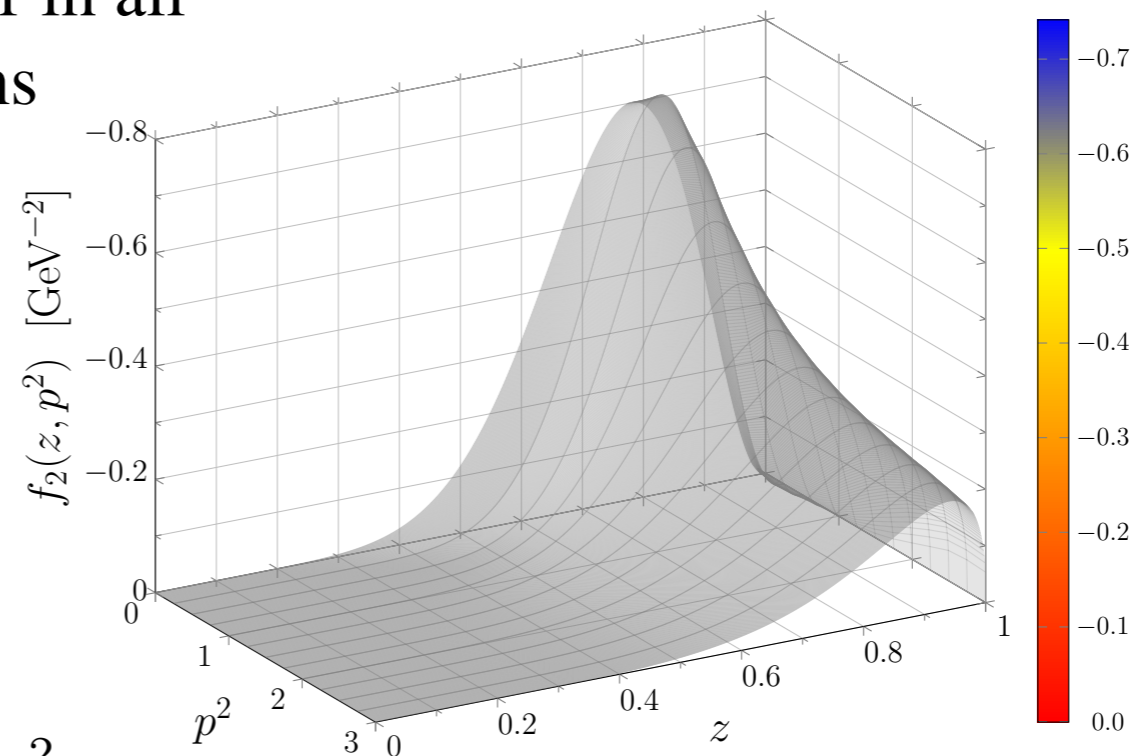
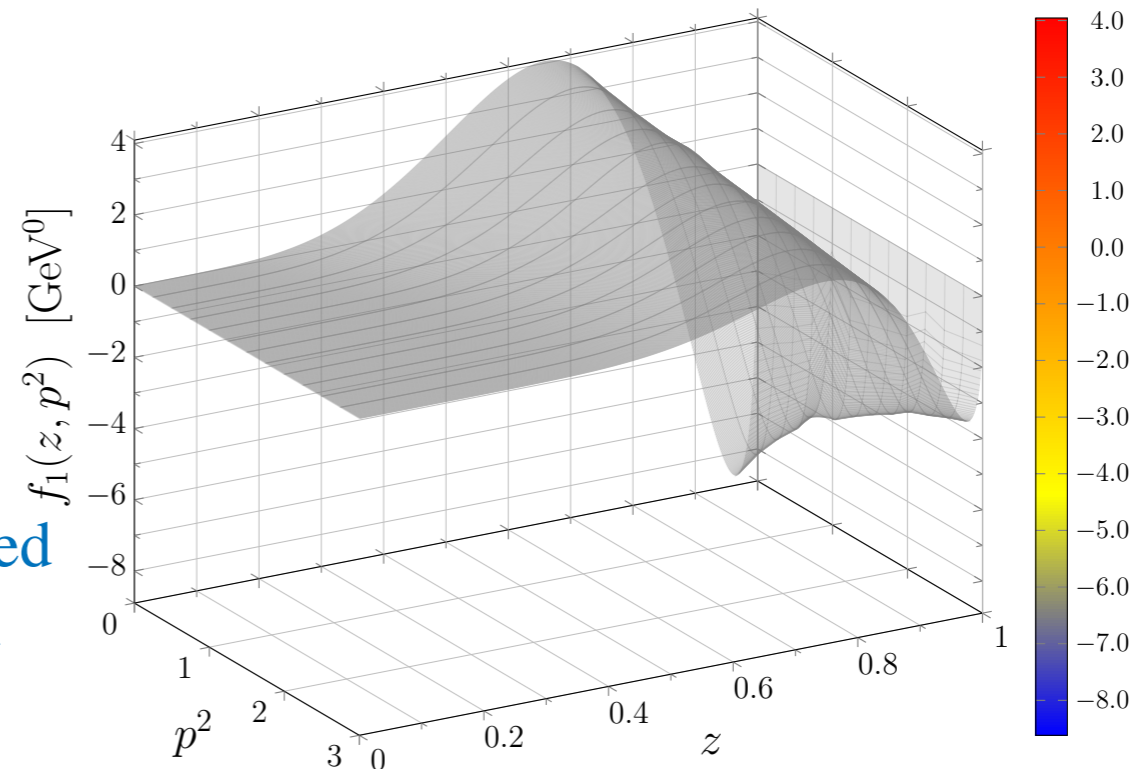
$$\Lambda_q(z, p, n) = iZ_2 \not{n} \delta(1 - z) - \iint du dw \delta(z - uw) \int \frac{d^4 \ell}{(2\pi)^4} \delta\left(w - \frac{\ell \cdot n}{p \cdot n}\right) \times \gamma_\mu \mathcal{K}_{\mu\nu}(p - \ell) S(\ell) \Lambda_q(u, \ell, n) S(\ell) \gamma_\nu$$

PDFs of a Dressed Quark

- *Hadron independent* vertex has form

$$\Lambda_q(z, p, n) = i\not{n} \delta(1 - z) + i\not{n} f_1^q(z, p^2) + n \cdot p [i\not{p} f_2^q(z, p^2) + f_3^q(z, p^2)]$$

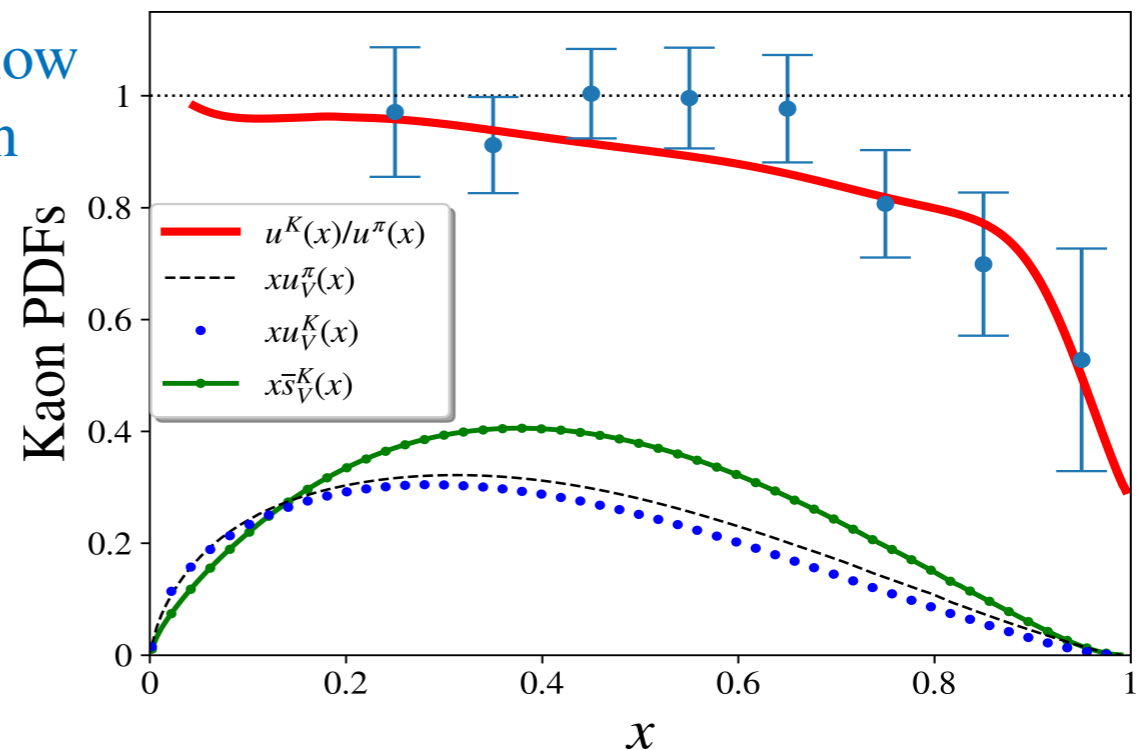
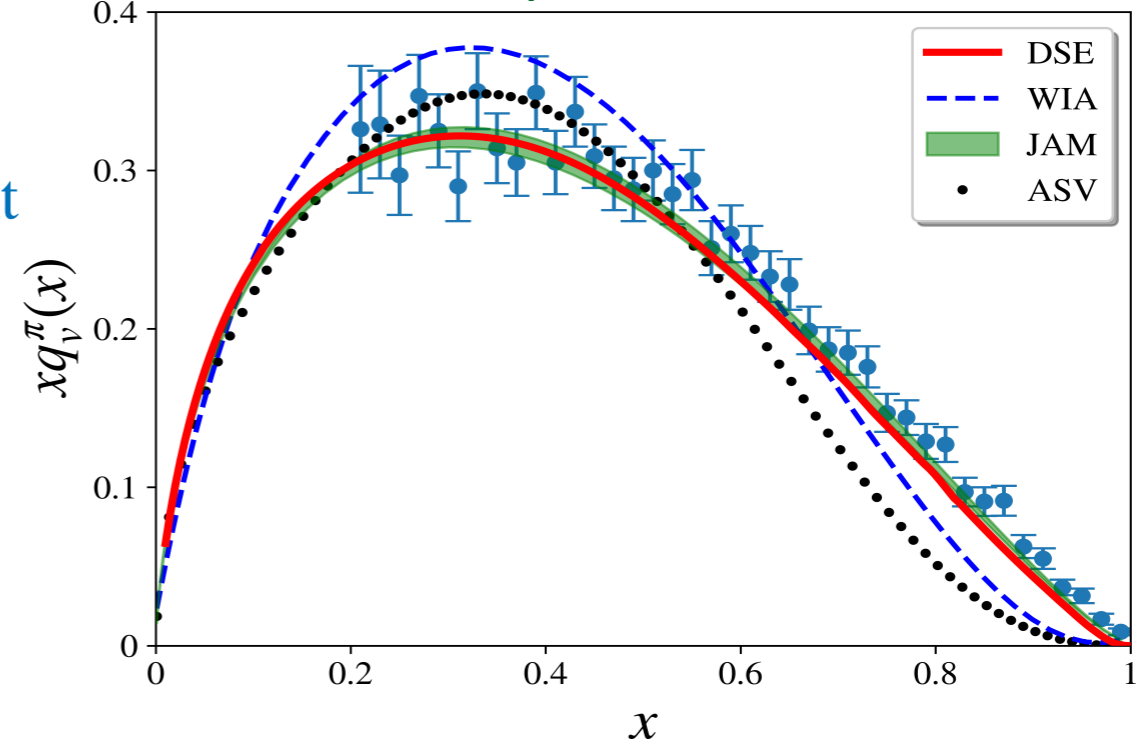
- the functions $f_i^q(z, p^2)$ can be interpreted as unpolarized PDFs in a dressed quark of virtuality p^2
- These functions are universal – appear in all RL-DSE unpolarized PDF calculations
- Distributed support in z is immediate indication gluons carry significant momentum
 - heavier s quark support nearer $z = 1$
 - WIA $\implies \Lambda_q(z, p, n) \propto \delta(1 - z)$
- Renormalization condition means dressing functions vanish when $p^2 = \mu^2$



Self-Consistent DSE Results

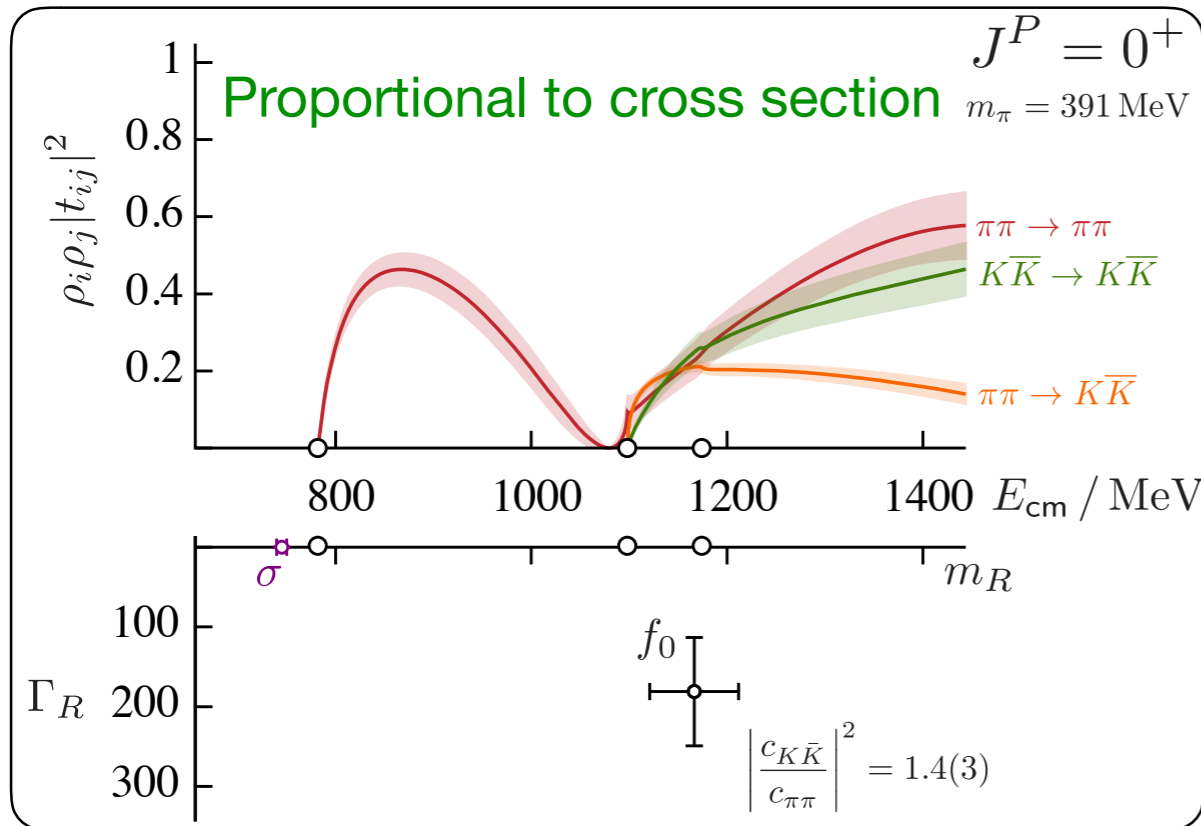
- For pion and kaon PDFs included for first time gluons self-consistently
 - correct RL-DSE pion PDFs in excellent agreement with Conway *et al.* data and recent JAM analysis
 - agrees with $x \rightarrow 1$ pQCD prediction
- Treating non-perturbative gluon contributions correctly pushes support of $q_\pi(x)$ to larger x
 - gluons remove strength from $q_\pi(x)$ at low to intermediate x – baryon number then demands increased support at large x
 - cannot be replicated by DGLAP – DSE splitting functions are dressed
- *Immediate consequence of gluon dressing is that gluons carry 35% of pion’s and 30% of kaon’s momentum*

[K. Bednar, ICC, P. Tandy, *et al.*, arXiv:1811.12310 [nucl-th]]



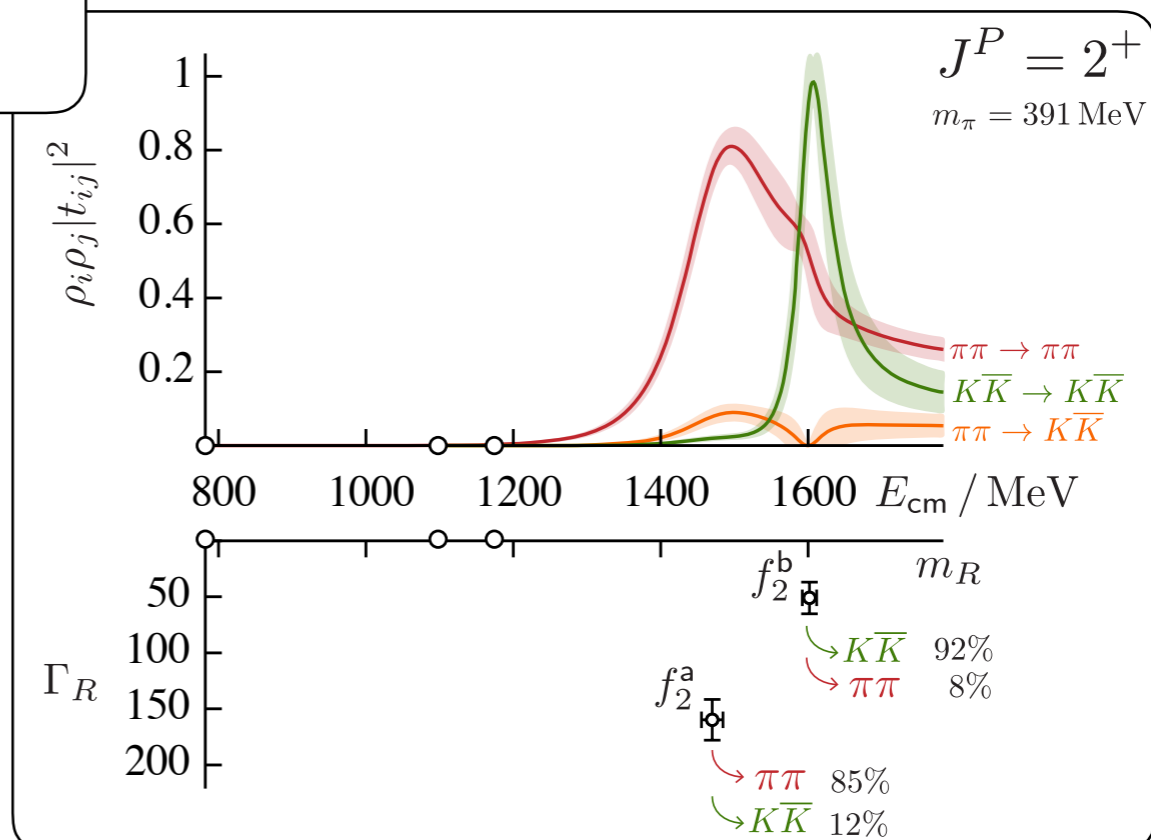
Raul Briceño — “QCD Spectroscopy on the Lattice”

Scalar $\pi\pi$ - KK



Isoscalar $\pi\pi, KK, \eta\eta$ scattering in S and D wave and the σ , $f_0(980)$, $f_2(1270)$ mesons from QCD.

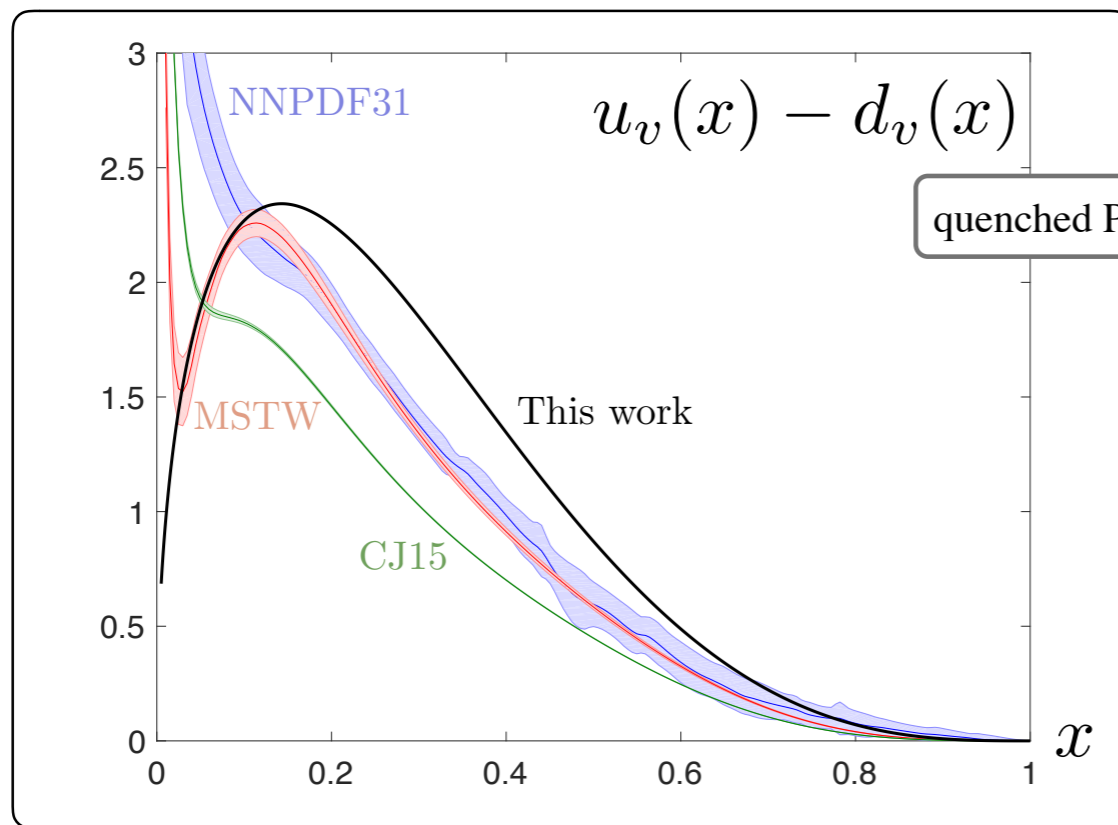
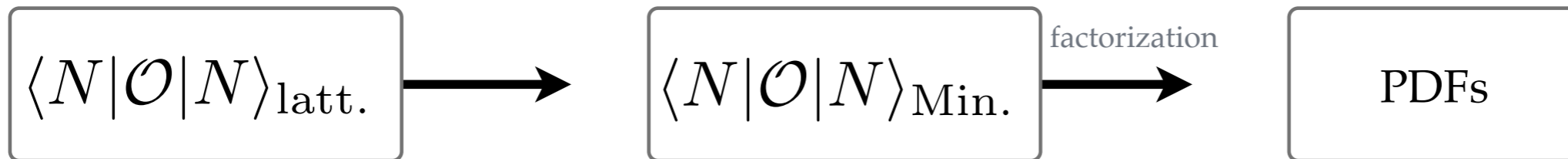
$f_0(980)$ manifests itself as a dip in the $\pi\pi$ cross section



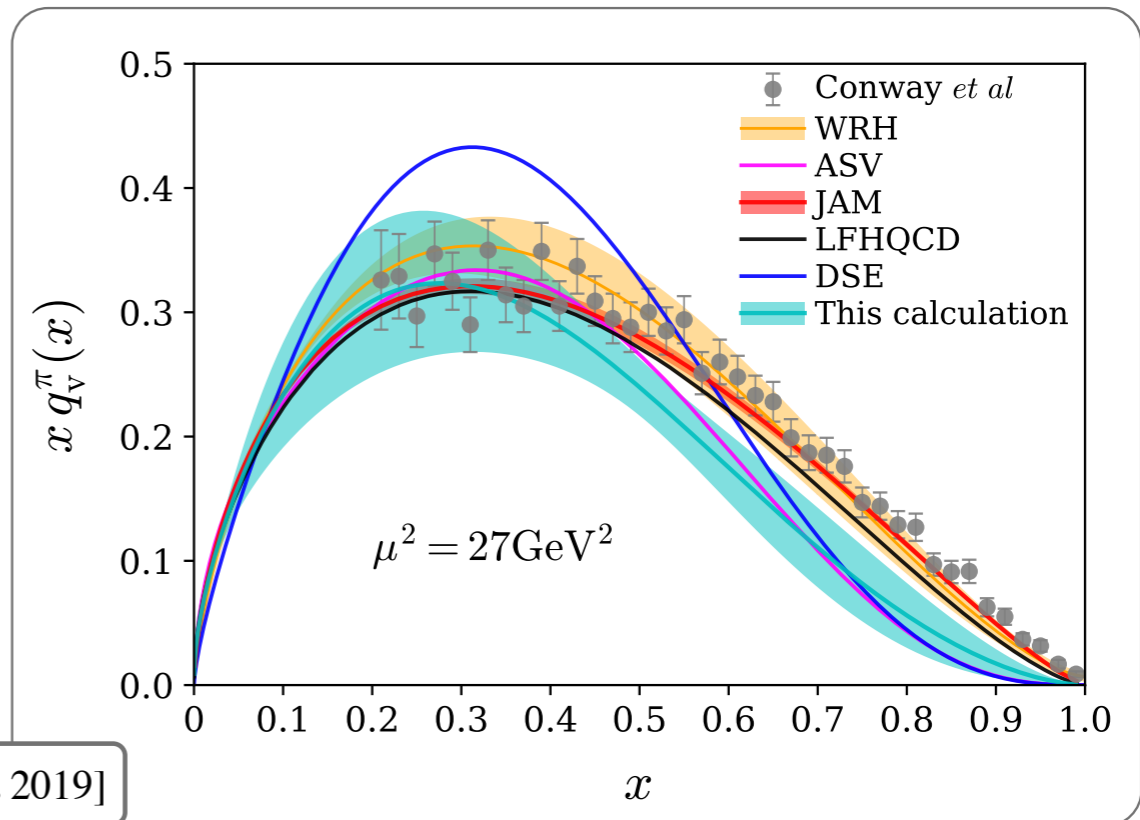
Raul Briceño — “QCD Spectroscopy on the Lattice”

PDFs in the lattice

Pseudo PDFs, not using moments



quenched PDFs [Orginos, et al. 2017]

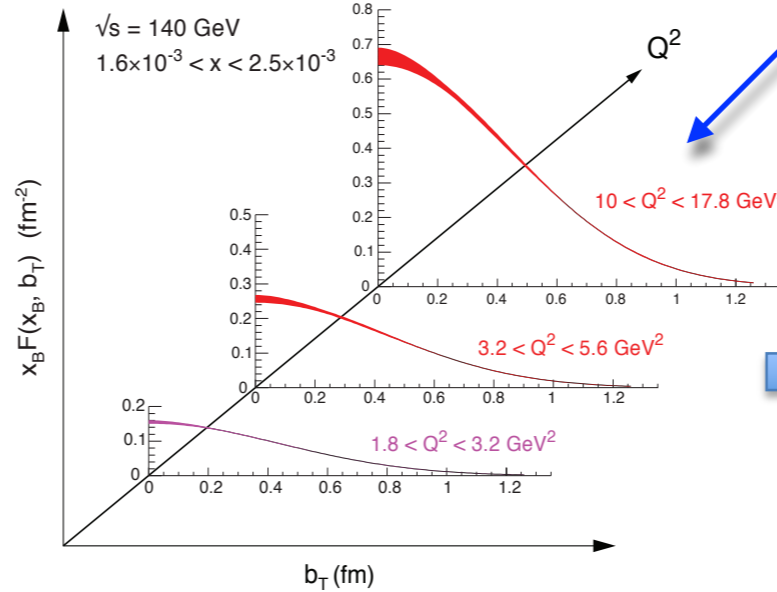
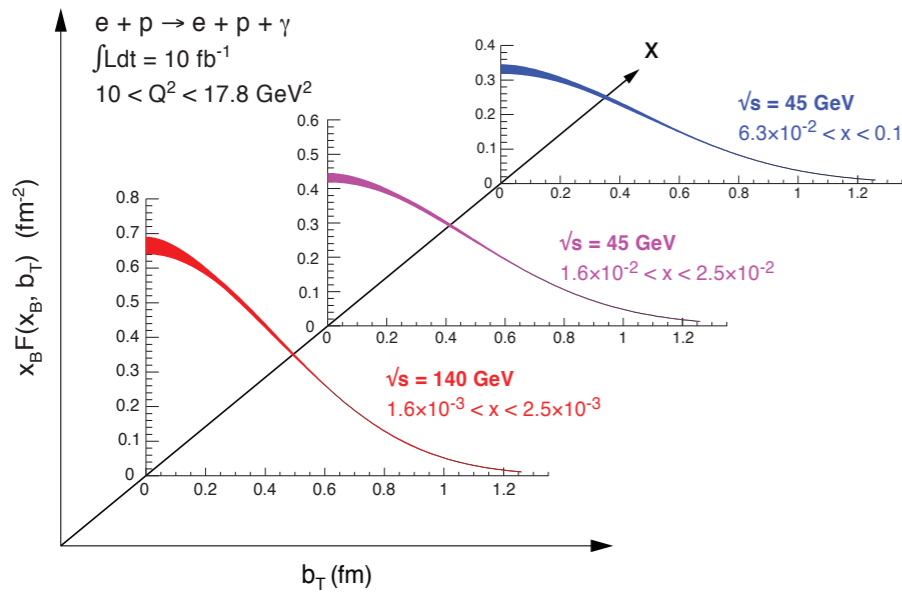


Pion Valence Quark Distribution @ $m_\pi = 400\text{MeV}$ [Sufian, et al. 2019]

Salvatore Fazio — “Support of Partonic Spatial Imaging at EIC”

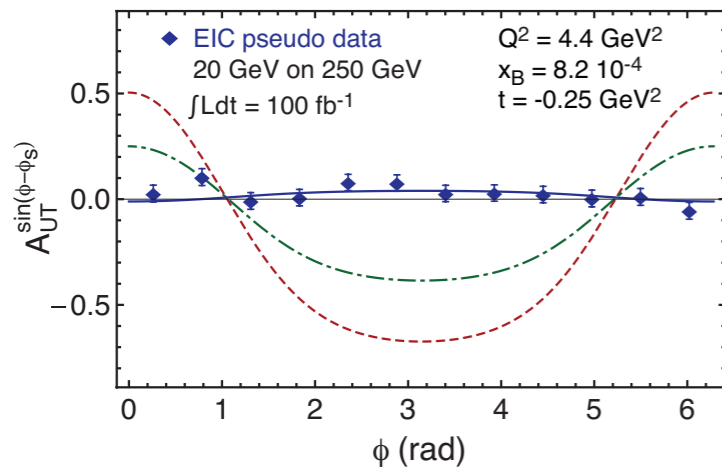
DVCS @ EIC

Average partonic densities obtained by Fourier transforming differential cross section measurements



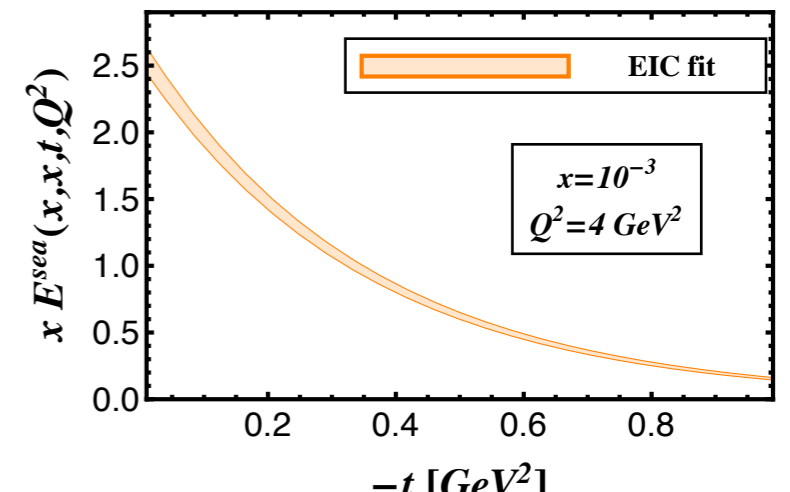
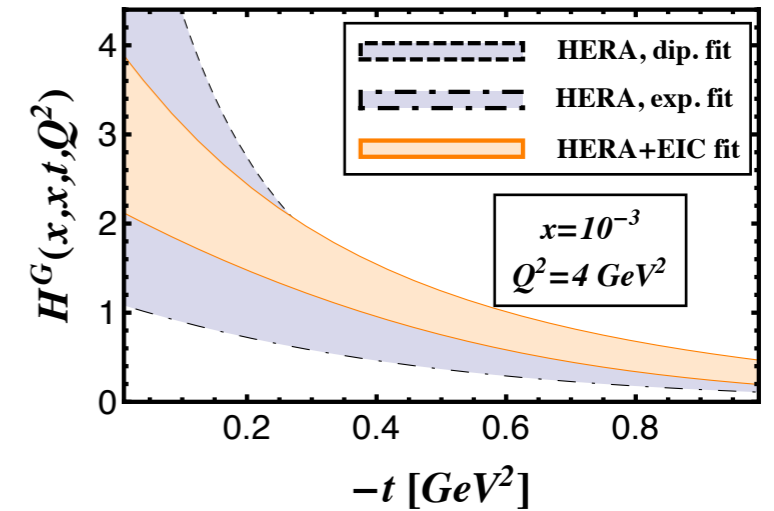
$$\int L = 10 \text{ fb}^{-1}$$

Constrains GPD H



$$\int L = 100 \text{ fb}^{-1}$$

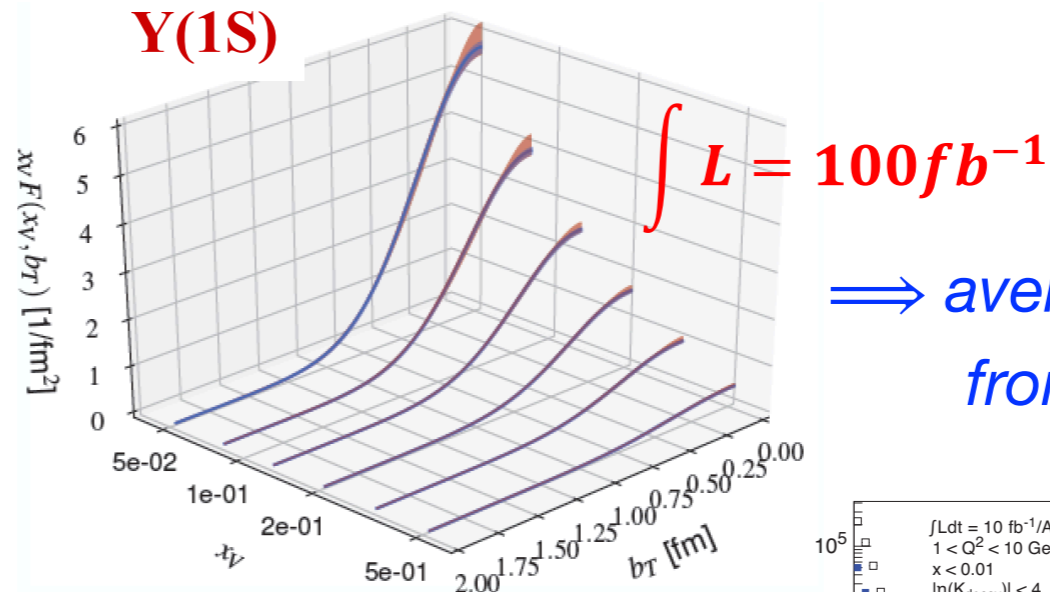
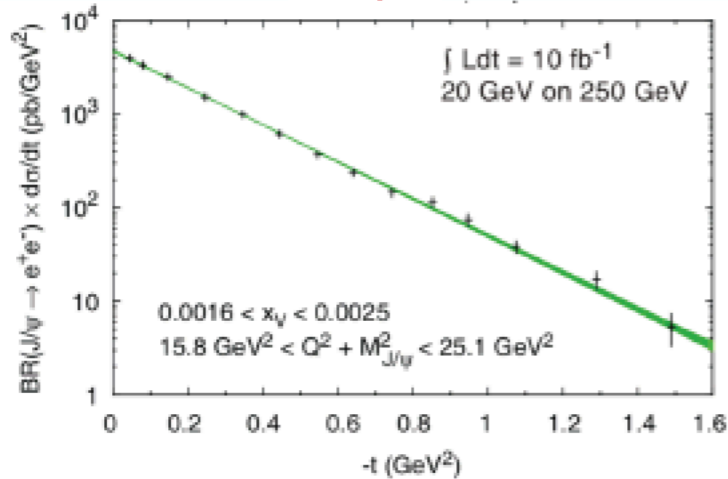
Sensitive to GPD E



- ❖ Studies performed at the times of the EIC WP era... (DVCS)
 - E.C.Aschenauer, S.Fazio, K.Kumerički, D. Müller, JHEP09(2013)093
- ❖ Global fits of DVCS world data and EIC pseudo-data
- ❖ Good constrain of H^{sea} , and H^g
- ❖ Opportunity for first precise extraction of sea-quarks GPD E

New excitement ahead

Imaging gluons with J/ψ and $Y(1S)$



⇒ average gluon density from t -spectrum

Goals (needed to fully exploit EIC data):

- ❖ Include mesons in global fits (flavor separation, precision in constraining gluons GPDs)
- ❖ Assess the sensitivity to extract “D-term”, related to distribution of forces inside the nucleon
- ❖ Extract neutron GPDs (via measuring DVCS on D/He3)
- ❖ Refine tools to study of GPDs in nuclei (and possible gluon saturation effects)
- ❖ Fully develop common software platforms (E.g. PARTONS)
- ❖ Go beyond: Can access gluon elliptic Wigner fcn. At EIC?
- Very engaged Community, stay tuned!

