

# GPDs and gravitational form factors of the proton

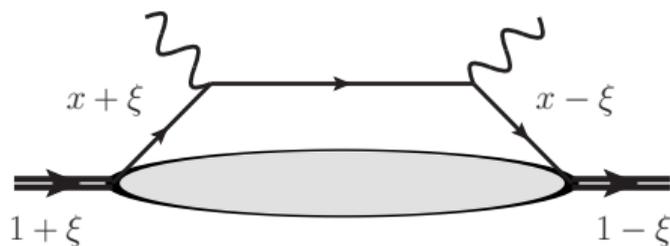
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## Why GPDs at an EIC?

- GPDs are the soft (non-perturbative) part of DVCS, DVMP, *etc.*



- **Hard exclusive reactions** are easier at an EIC!
  - Recoil proton/ion more easily measured than in fixed target experiments.
  - High luminosity helpful for events with low cross sections.
  - Hermetic coverage (central timing-optimized silicon, RICH for forward scattering) great for exclusive reactions.
- GPDs are related to literal spatial distributions by 2D Fourier transforms

$$\int \frac{d^2\mathbf{k}_\perp}{(2\pi)^2} H^q(x, \xi = 0, t = -\mathbf{k}_\perp^2) e^{-i(\mathbf{k}_\perp \cdot \mathbf{b}_\perp)} = \rho^q(\mathbf{b}_\perp)$$

- GPDs encode information about decomposition & distribution of **mass**, **angular momentum**, and **forces** through polynomiality relations.

## Gravitational form factors and the EMT

- **Gravitational form factors** encode information in the energy-momentum tensor (EMT):

$$\langle p'\lambda' | T_{\mu\nu}^q(0) | p\lambda \rangle = \bar{u}^{\lambda'}(p') \left[ \gamma_{\{\mu} P_{\nu\}} A^q(t) + \frac{iP_{\{\mu} \sigma_{\nu\}} \Delta}{2m_N} B^q(t) + \frac{\Delta_\mu \Delta_\nu - \Delta^2 g_{\mu\nu}}{4m_N} C^q(t) \right. \\ \left. + m_N g_{\mu\nu} \bar{c}^q(t) + \frac{iP_{[\mu} \sigma_{\nu]} \Delta}{2m_N} D^q(t) \right] u^\lambda(p)$$

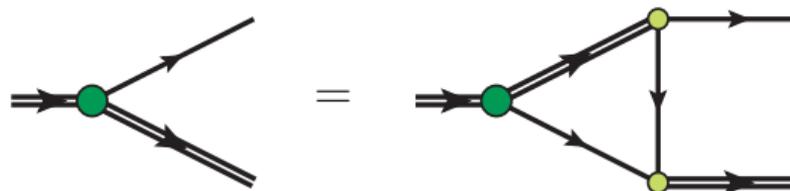
- Three of the GFFs accessible through GPDs:

$$\int dx x H^q(x, \xi, t) = A^q(t) + \xi^2 C^q(t), \quad \int dx x E^q(x, \xi, t) = B^q(t) - \xi^2 C^q(t)$$

- GFFs encode information about distribution of **energy**, **angular momentum**, and **forces**.
- **Cannot do graviton-exchange experiments**—but GPDs allow GFFs to be experimentally accessed! (DVCS, DVMP, ...)

# Quark-diquark model

- **Quark-diquark model:** proton is dominated by configurations where two quarks are in a **diquark correlation**.
- Faddeev equation replaced by two-body equation with quark exchange kernel.



- Two species of diquarks dominate.
  - ① Scalar, isoscalar ( $S = 0, T = 0$ )
  - ② Axial vector, isovector ( $S = 1, T = 1$ )
- Evidence for diquarks in flavor-separated form factors

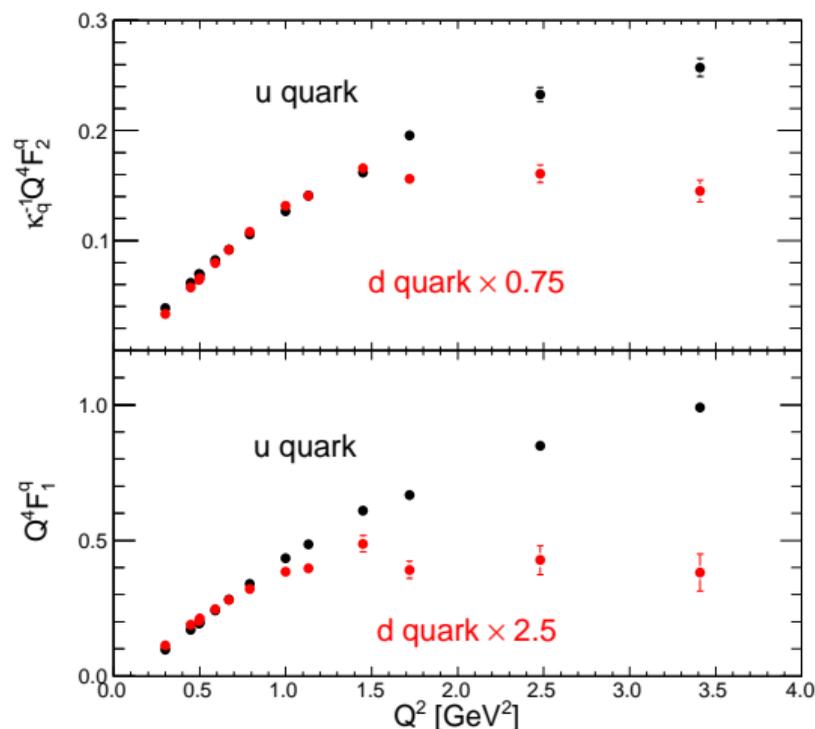
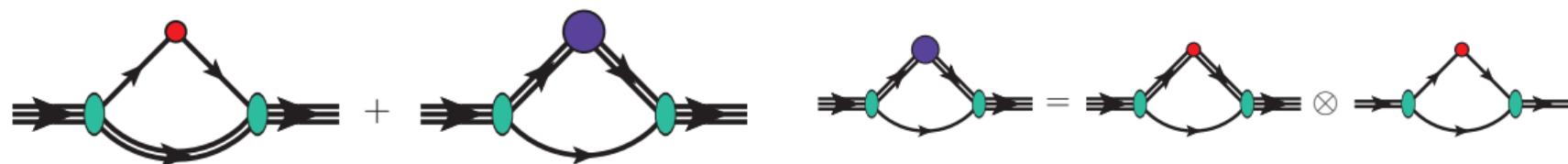


Figure from Cates *et al.*, PRL106 (2011) 252003

# Diagrams and convolution



## Proton GPD diagrams

- GPD gets contributions from the lonely quark and from the diquark.
- Diquark contribution given by a convolution formula

$$H_{X,i}(x, \xi, t) = \sum_j \int \frac{dy}{|y|} h_{Y/X,ij}(y, \xi, t) H_{Y,j} \left( \frac{x}{y}, \frac{\xi}{y}, t \right)$$

- Diquark GPDs include scalar and axial species, along with a transition GPD
- The dressed quark will also have a GPD, to be folded in through a second layer of convolution
- Practically, do model calculation in NJL model

# Dressed quark GPDs

- GPD defined through **current quark fields**; proton made of **three dressed quarks**.
- Bethe-Salpeter equation gives GPD of dressed quarks in terms of current quarks:

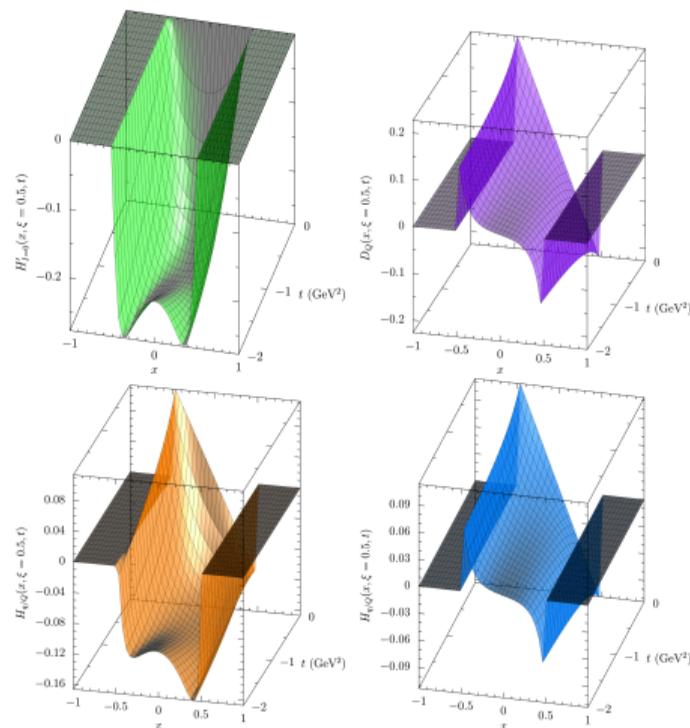


- Fold quark GPDs into proton GPD using another layer of convolution.
- Isospin mixture solutions:

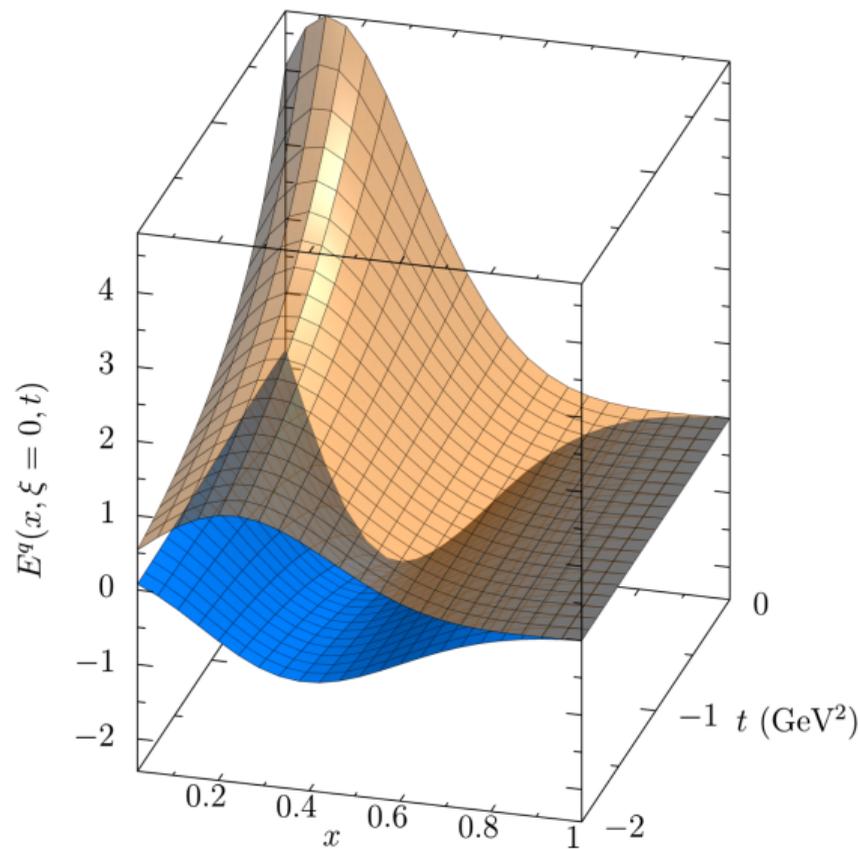
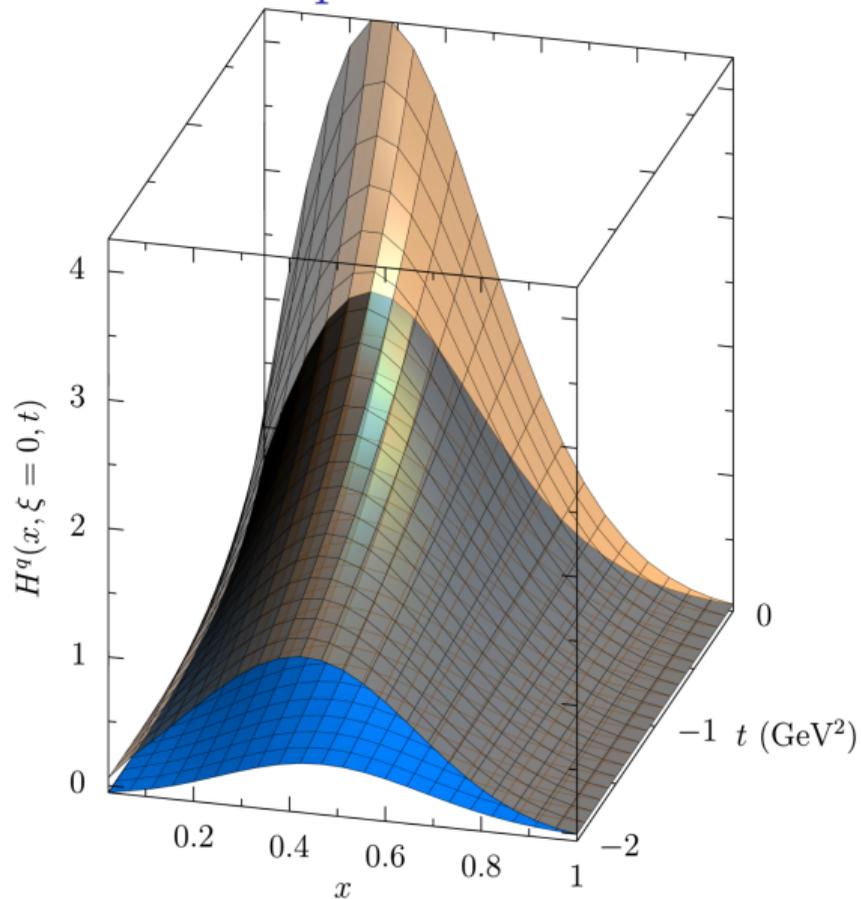
$$H_I(x, \xi, t) = \delta(1-x) + H'_I(x, \xi, t) + \delta_{I,0} D_Q(x, \xi, t)$$

$$E_I(x, \xi, t) = -\delta_{I,0} D_Q(x, \xi, t)$$

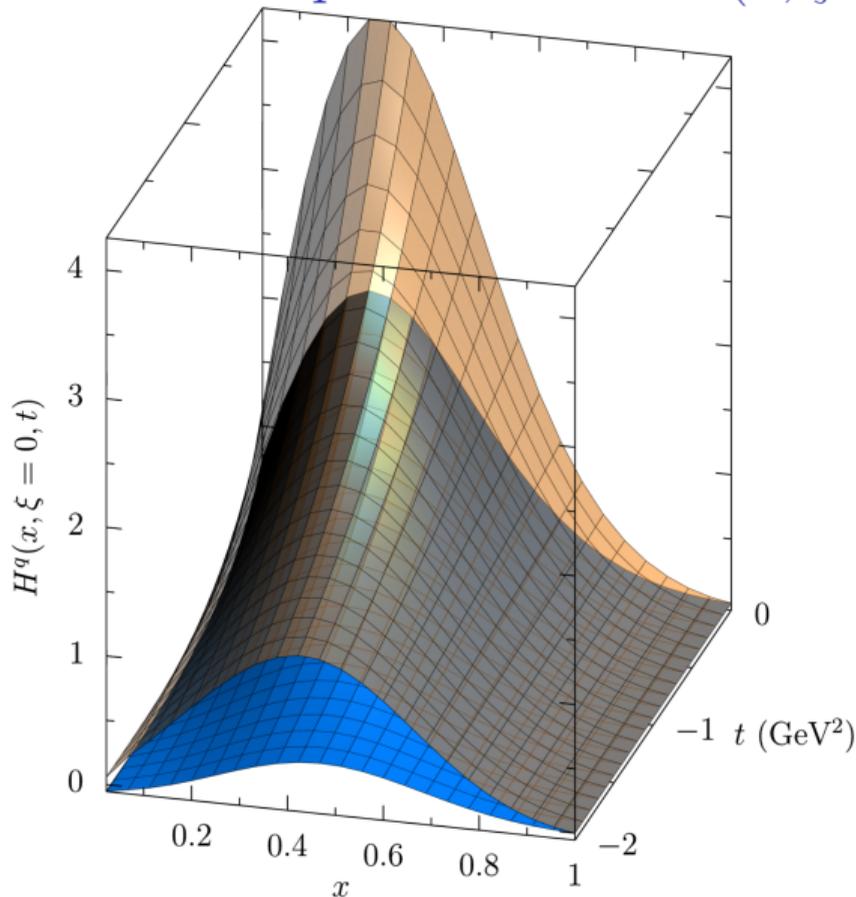
- Orange for  $u/U$  or  $d/D$ ; blue for  $d/U$  or  $u/D$ .



# Non-skewed proton GPD results



## Non-skewed proton GPD: $H^q(x, \xi = 0, t)$

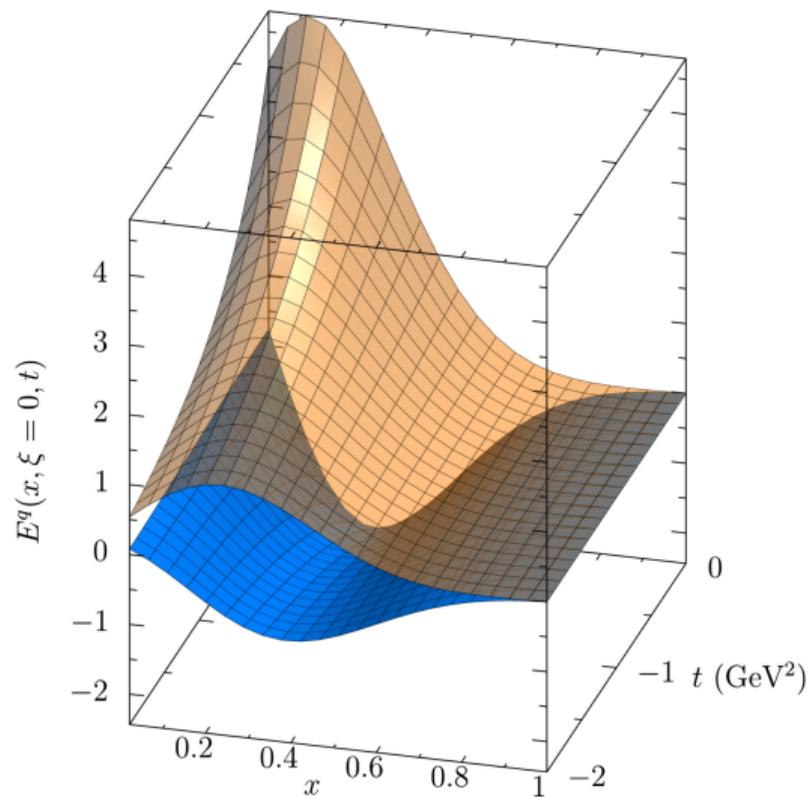


Orange is up; blue is down.

- GPD not dressed in DGLAP region
- **Forward limit:**  $H^q(x, 0, 0) = q(x)$
- up-down differences due to diquark correlations
- **Scalar diquarks** dominate at low  $-t$ ;  
up quarks carry more  $x$  than down
  - At  $t = 0$ ,  $\langle\langle x_u \rangle\rangle = 0.34$
  - At  $t = 0$ ,  $\langle\langle x_d \rangle\rangle = 0.32$
- **Axial diquarks** dominate at high  $-t$ ;  
down quark carries more  $x$ 
  - At  $t = -2$  GeV<sup>2</sup>,  $\langle\langle x_u \rangle\rangle = 0.45$
  - At  $t = -2$  GeV<sup>2</sup>,  $\langle\langle x_d \rangle\rangle = 0.50$

# Non-skewed proton GPD: $E^q(x, \xi = 0, t)$

Orange is up; blue is down.



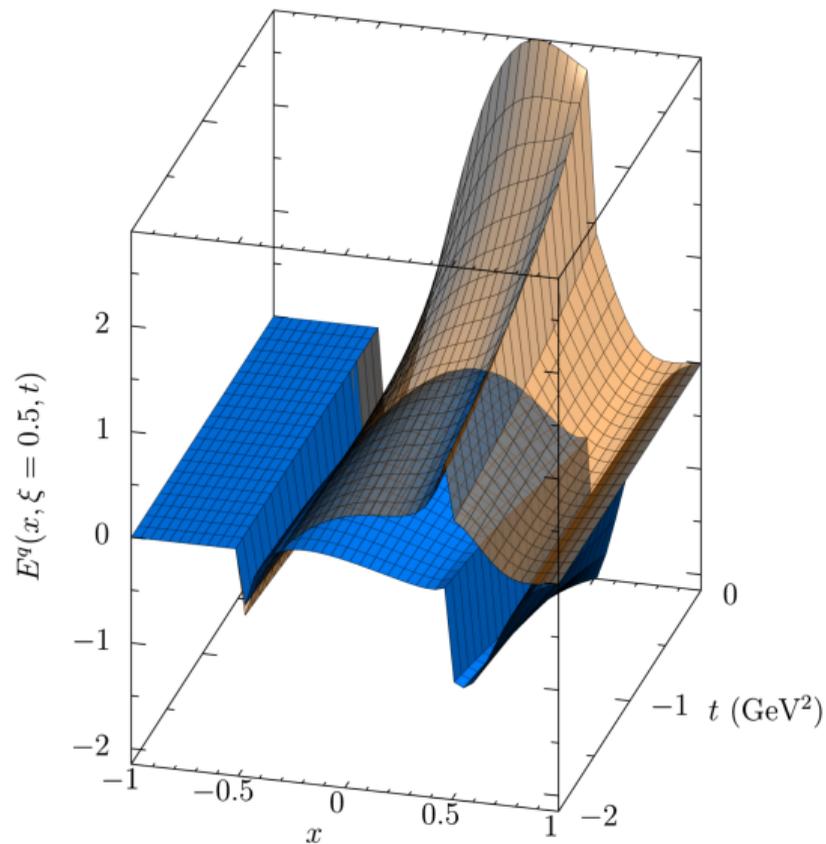
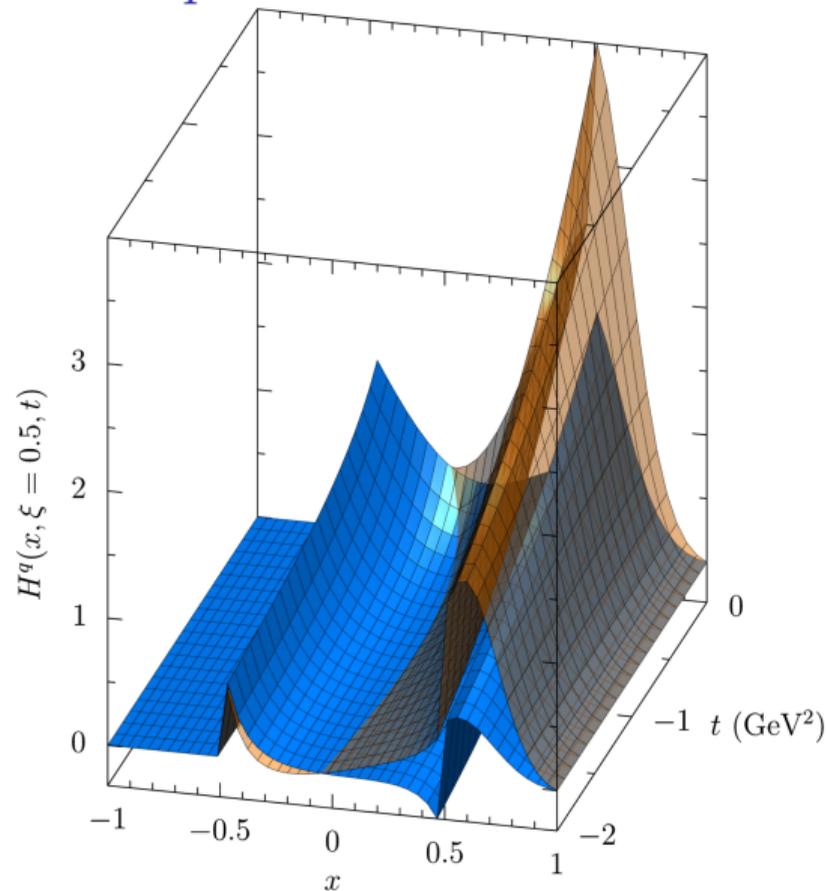
- **Axial diquarks** dominate at all  $-t$ ;  
down quarks carry more  $x$
- No forward limit; understand by moments

$$\int dx E^q(x, 0, t) = F_2^q(t)$$

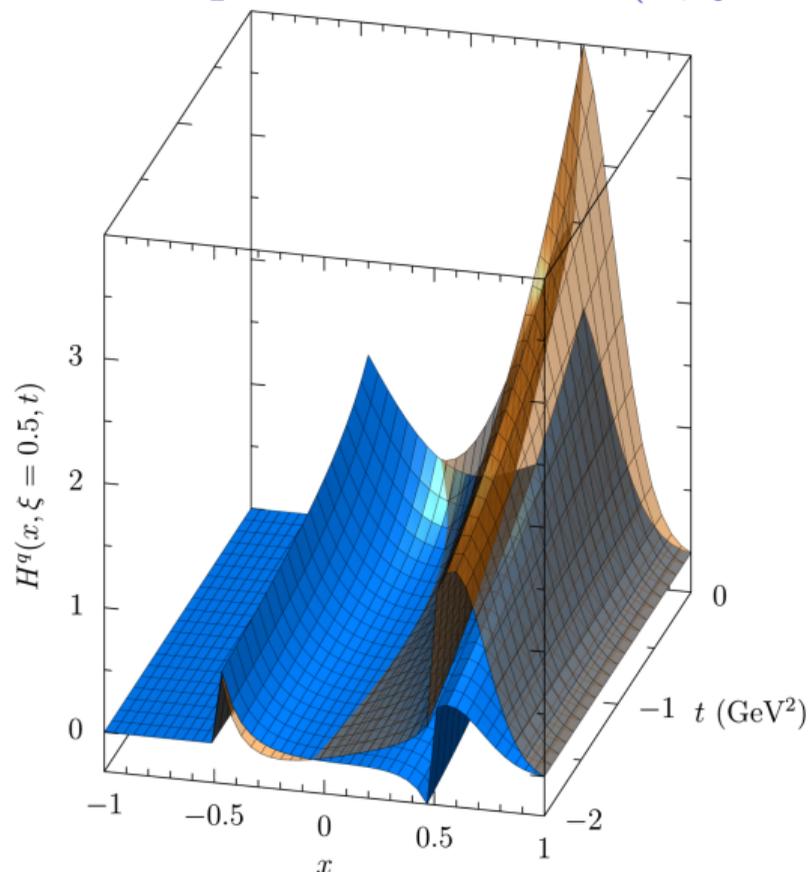
$$\int dx x E^q(x, 0, t) = B^q(t) = 2J^q(t) - A^q(t)$$

- $J^u$  and  $J^d$  **have opposite signs**
  - $J^u(0) = 0.52$  (close to total)
  - $J^d(0) = -0.02$
- $\kappa^u$  and  $\kappa^d$  **have same sign**
  - $e_u F_2^u(0) = 1.15$
  - $e_d F_2^d(0) = 0.34$

# Skewed proton GPD results



## Skewed proton GPD: $H^q(x, \xi = 0.5, t)$



Orange is up; blue is down.

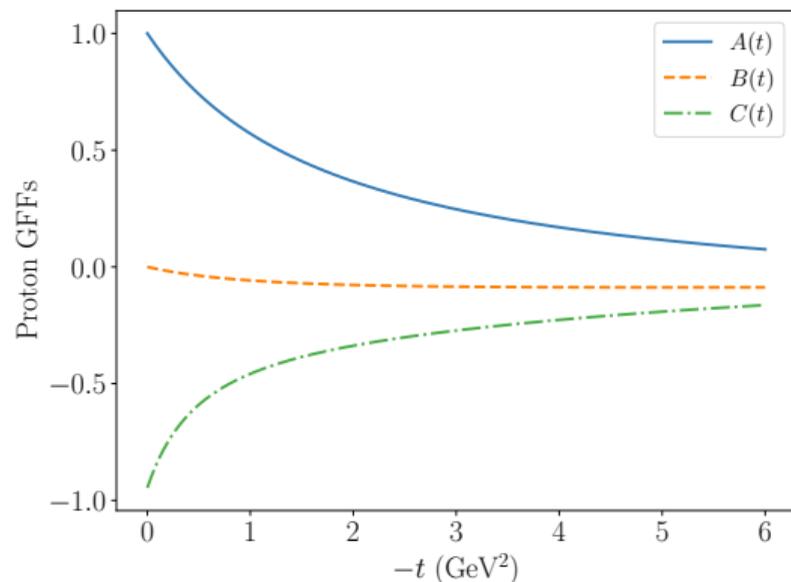
- GPD dressed in ERBL region:  $|x| < |\xi|$
- up-down difference in ERBL region due to isospin-dependent dressing
- Near  $\rho$ - $\omega$  mass degeneracy:  $H_{I=0} \approx H_{I=1}$

$$H_{u/U}(x, \xi, t) \approx H_{I=0}(x, \xi, t) + \frac{1}{2}D_Q(x, \xi, t)$$

$$H_{d/U}(x, \xi, t) \approx \frac{1}{2}D_Q(x, \xi, t)$$

- Down quark GPD dominated by  $D_Q(x, \xi, t)$  in ERBL region

# Gravitational form factors of the proton



Can extract **gravitational form factors** from GPDs; model is Lorentz covariant

$$\int dx x H^q(x, \xi, t) = A^q(t) + \xi^2 C^q(t)$$

$$\int dx x E^q(x, \xi, t) = B^q(t) - \xi^2 C^q(t)$$

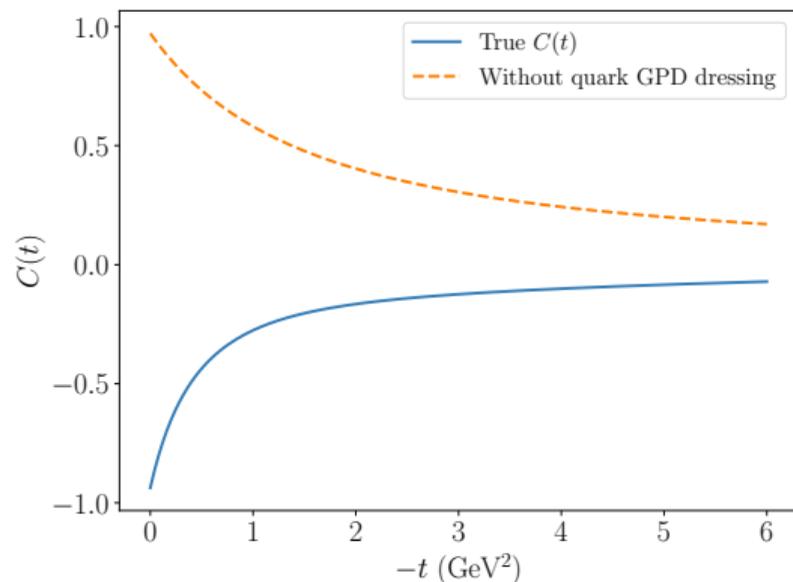
Sum rules:  $A(0) = 1$ ,  $B(0) = 0$

- $A(t)$  gives a mass form factor
- $J(t) = \frac{1}{2} (A(t) + B(t))$  gives an angular momentum form factor; *n.b.*,  $J(0) = \frac{1}{2}$
- The sum rules are guaranteed by Poincaré invariance of the model

## D-term of the proton

$C(t)$  encodes **pressure distribution** (see Peter Schweitzer's talk)

- $C(0)$  not constrained by conservation laws
- $C(0) < 0$  is a stability condition
- $C(t)$  is affected by quark GPD dressing, but not  $A(t)$  or  $B(t)$



- $C(0) = -0.94 < 0$  in the NJL model
- Ignoring quark GPD dressing gives  $C(0) = 0.97 > 0$
- Proton stability appears to **require** dressing the quark GPD!

# Conclusions & outlook

## Conclusions

- We've calculated leading-twist, helicity-independent proton GPDs in the NJL model.
- The GPDs show manifestations of diquark correlations.
- Can extract predictions for gravitational form factors, since model is covariant.
- Dressing of light cone correlator needed for correct description of ERBL region.
- Dressing required for proton stability.

## Outlook

- Calculations for helicity-dependent and helicity-flip GPDs also possible.
- Plans to extend model to include intrinsic (non-perturbative) glue.

## See also

- Talks by Peter Schweitzer (parallel session), Cédric Lorcé (Thursday)
- Posters by Arkadiusz Trawinski, David Arturo Amor Quiroz, & Peter Lowdon
- arXiv:1907.08256