New modelling techniques for Generalised Parton Distributions

Cédric Mezrag

INFN Roma1

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In collaboration with: N. Chouika, H. Moutarde and J. Rodriguez-Quintero

> Eur.Phys.J. C77 (2017) no.12, 906 Phys.Lett. B780 (2018) 287-293

Modelling GPDs

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• Generalized Parton Distributions (GPDs):

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Definition of GPDs



- Generalized Parton Distributions (GPDs):
 - are defined according to a non-local matrix element,

$$\begin{split} &\frac{1}{2}\int \frac{e^{ixP^+z^-}}{2\pi} \langle P + \frac{\Delta}{2} |\bar{\psi}^q(-\frac{z}{2})\gamma^+\psi^q(\frac{z}{2})|P - \frac{\Delta}{2}\rangle \mathrm{d}z^-|_{z^+=0,z=0} \\ &= \frac{1}{2P^+} \bigg[H^q(x,\xi,t)\bar{u}\gamma^+u + E^q(x,\xi,t)\bar{u}\frac{i\sigma^{+\alpha}\Delta_{\alpha}}{2M}u \bigg]. \end{split}$$

$$\begin{split} &\frac{1}{2}\int \frac{e^{ixP^+z^-}}{2\pi} \langle P + \frac{\Delta}{2} |\bar{\psi}^q(-\frac{z}{2})\gamma^+\gamma_5\psi^q(\frac{z}{2})|P - \frac{\Delta}{2}\rangle \mathrm{d}z^-|_{z^+=0,z=0} \\ &= \frac{1}{2P^+} \bigg[\tilde{H}^q(x,\xi,t)\bar{u}\gamma^+\gamma_5u + \tilde{E}^q(x,\xi,t)\bar{u}\frac{\gamma_5\Delta^+}{2M}u \bigg]. \end{split}$$

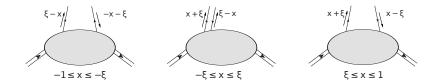
D. Müller et al., Fortsch. Phy. 42 101 (1994)
 X. Ji, Phys. Rev. Lett. 78, 610 (1997)
 A. Radyushkin, Phys. Lett. B380, 417 (1996)

4 GPDs without helicity transfer + 4 helicity flip GPDs

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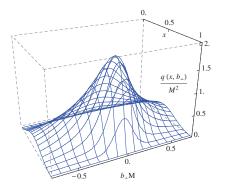
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M. Burkardt, Phys. Rev. D62, 071503 (2000)



Pion GPD in Impact parameter space from: C. Mezrag *et al.*, Phys. Lett. **B741**, 190-196 (2015)

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 - can be related to the Energy-Momentum tensor (GFF) through their n = 1 Mellin moments

X. Ji, PRL 78, 610-613 (1997)

X. Ji, J. Phys. G24, 1181-1205 (1998)

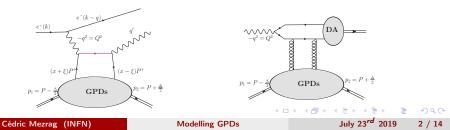
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 - can be related to the Energy-Momentum tensor (GFF) through their n = 1 Mellin moments
 - are univeral, *i.e.* are related to the Compton Form Factors (CFFs) of various exclusive processes through convolutions

$$\mathfrak{H}(\xi,t) = \int \mathrm{d}x \ C(x,\xi) H(x,\xi,t)$$



GPDs theoretical constraints



• Polynomiality:

$$\int_{-1}^{1} \mathrm{d}x \; x^{m} H^{q}(x,\xi,t) = \sum_{j=0}^{\left[\frac{m}{2}\right]} \xi^{2j} C_{2j}^{q}(t) + mod(m,2)\xi^{m+1} C_{m+1}^{q}(t)$$

Lorentz Covariance

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GPDs theoretical constraints

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• Polynomiality:

Lorentz Covariance

• Positivity:

$$\left|H^q(x,\xi,t)-\frac{\xi^2}{1-\xi^2}E^q(x,\xi,t)\right|\leq \sqrt{\frac{q\left(\frac{x+\xi}{1+\xi}\right)q\left(\frac{x-\xi}{1-\xi}\right)}{1-\xi^2}}$$

A. Radysuhkin, Phys. Rev. **D59**, 014030 (1999)
B. Pire *et al.*, Eur. Phys. J. **C8**, 103 (1999)
M. Diehl *et al.*, Nucl. Phys. **B596**, 33 (2001)
P.V. Pobilitsa, Phys. Rev. **D65**, 114015 (2002)

Positivity of Hilbert space norms

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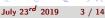
• Positivity:

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Positivity of Hilbert space norms

• Support:

 $x \in [-1; 1]$ M. Diehl and T. Gousset, Phys. Lett. **B428**, 359 (1998) Relativistic Quantum Mechanic





- Polynomiality:
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Soft Pion theorem

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Positivity of Hilbert space norms

Relativistic Quantum Mechanic

M.V. Polyakov, Nucl. Phys. **B555**, 231 (1999) C. Mezrag *et al.*, Phys. Lett. **B741**, 190 (2015)

PCAC and Axial-Vector WTI

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- Polynomiality:
- Positivity:
- Support:
- Soft Pion theorem

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Positivity of Hilbert space norms

Relativistic Quantum Mechanic

PCAC and Axial-Vector WTI

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Problems

There is no model (until now) fulfilling a priori all these constraints





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• GPDs are related to Double Distributions (DDs) through:

$$H(x,\xi,t) = \int_{\Omega} d\beta d\alpha \left(F(\beta,\alpha,t) + \xi G(\beta,\alpha,t) \right) \delta \left(x - \beta - \xi \alpha \right)$$

The Dirac δ insures that the polynomiality is fulfilled, independently of our choice of F and G



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- They also appear naturally in covariant modelling attempts

Positivity property is not guaranteed, and may be violated.

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• On the light front, hadronic states can be expanded on a Fock basis:

$$|P,\pi
angle \propto \sum_{eta} \Phi_{eta}^{qar{q}} |qar{q}
angle + \sum_{eta} \Phi_{eta}^{qar{q},qar{q}} |qar{q},qar{q}
angle + \dots$$

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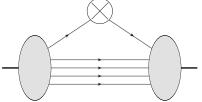
- Non-perturbative physics is contained within the N-particle LFWFs Φ^N
- This formalism allows to recover the probabilistic picture of non-relativistic quantum mechanics

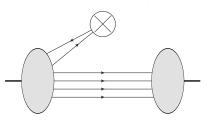
- Same N I FWFs
- No ambiguity

Classical modelling techniques II LFWFs approach to GPDs

On the light front, hadronic states can be expanded on a Fock basis

DGLAP: $|x| > |\xi|$





ERBL: $|x| < |\xi|$

- N and N + 2 partons LFWFs
- Ambiguity

M. Diehl et al., Nucl. Phys. B596 (2001) 33-65



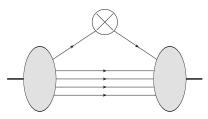
Modelling GPDs



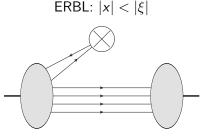
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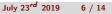
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LFWFs formalism has the positivity property inbuilt but polynomiality is lost by truncating both in DGLAP and ERBL sectors.

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• Mathematical properties of GPDs:

$$H(x,\xi,t) - D(x/\xi,t) = \int_{-1}^{1} \mathrm{d}\beta \int_{-1-|\beta|}^{1-|\beta|} \mathrm{d}\alpha F(\beta,\alpha,t) \delta(x-\beta-\alpha\xi)$$

D. Müller *et al.*, Fortsch. Phys. 1994, 42, 101
 A. Radyushkin, Phys.Rev., 1997, D56, 5524-5557
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• *H* in the DGLAP region only allows to obtain *F* up to *D*-term-like contributions (which remains of physical interest)

J. Boman and E. T. Quinto, Duke Math. J. 55, 943 (1987). N. Chouika, C. Mezrag, *et al.*, Eur.Phys.J. C77 (2017) no.12, 906



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 Relation between DGLAP and ERBL region allowing the fulfilment of polynomiality



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- Relation between DGLAP and ERBL region allowing the fulfilment of polynomiality
- Additional constraints are needed to fix the D-term

N. Chouika, C. Mezrag et al., Phys.Lett. B780 (2018) 287-293

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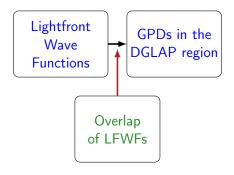
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Lightfront Wave Functions

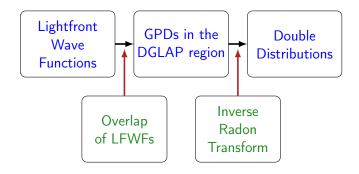






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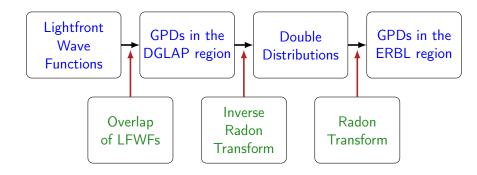




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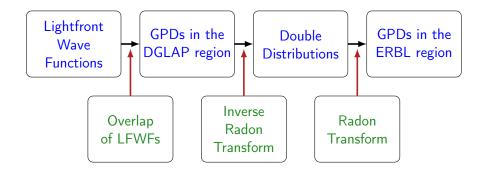
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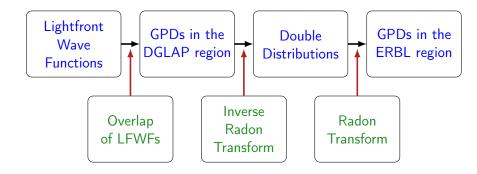
Reshuffling of the series in the ERBL region \rightarrow Polynomiality is fulfilled at every order in *N*.

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Applied and tested in various cases

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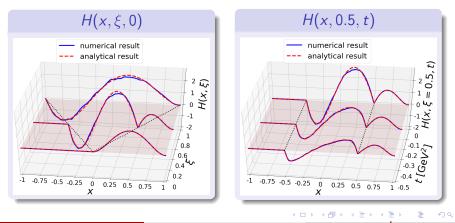
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Numerical Benchmarking Nabil Chouika Ph.D. Thesis



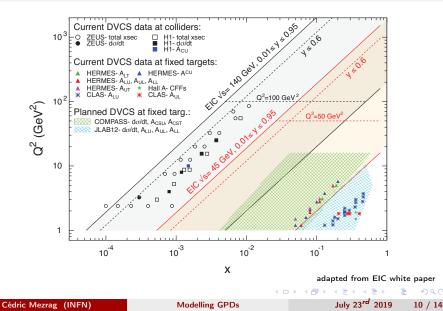
- The inverse Radon transform is an ill-posed problem
- Numerical implementation can be challenging due to noise



DVCS at EIC

Kinematical range

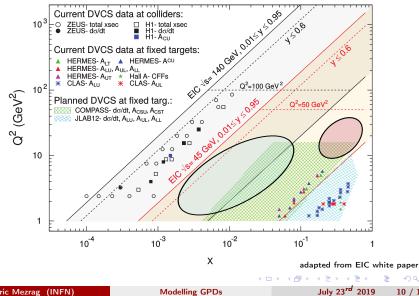




DVCS at EIC

Kinematical range





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Motivations

- Small range of Q^2 available today in the valence region
- LO/LT not a good approximation at JLab Kinematics

M. Defurne et al., Phys.Rev. C92 (2015) no.5, 055202 M. Defurne et al., Nature Commun. 8 (2017) no.1, 1408





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Challenges

• Unclear whether the EIC will be able to measure high- x_B DVCS





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Challenges

• Unclear whether the EIC will be able to measure high- x_B DVCS

Work in Progress

- Recently started working with Jinlong Zhang (Stony Brook U.)
- Use PARTONS together with Stony Brook's EIC DVCS event generator
- Under which conditions could we obtain relevant valence-x DVCS data?

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Conclusion



Modelling GPDs

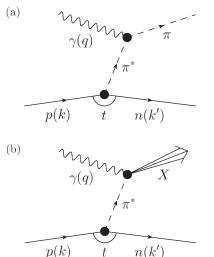
- New formalism to model GPDs which fulfils by construction all theoretical constraints for the first time
- Based on Lightfront, from whichever non-perturbative approach you prefer (Bethe-Salpeter, Holographic QCD, Hamiltonian formulation...)
- Use for phenomenological applications?

EIC valence-x DVCS measurements

- Valence-x DVCS measurement may be possible (need to be checked)
- A welcome high- Q^2 complementary measurement to fix target experiments

Addendum: Mesons Structure at EIC





 "White paper" for Pion and Kaon physics at an EIC

arXiv:1907.08218

- Two key point to study the internal structure of pions :
 - internal structure of Goldstone bosons, and the dynamic behind it
 - meson distribution amplitude play a significant role in DVMP

figure from arXiv:1907.08218

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Thank you for your attention

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