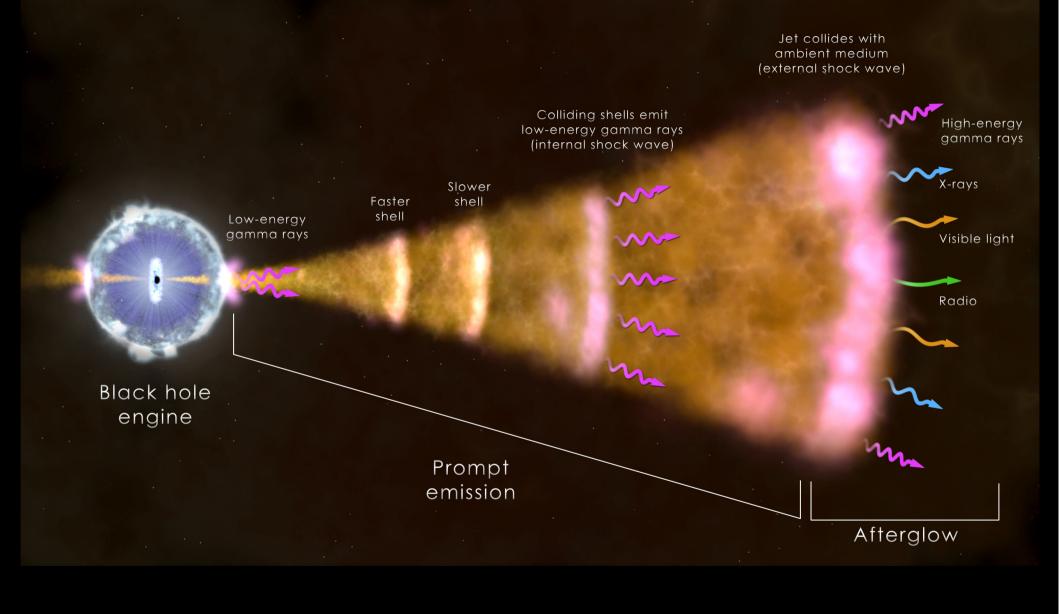
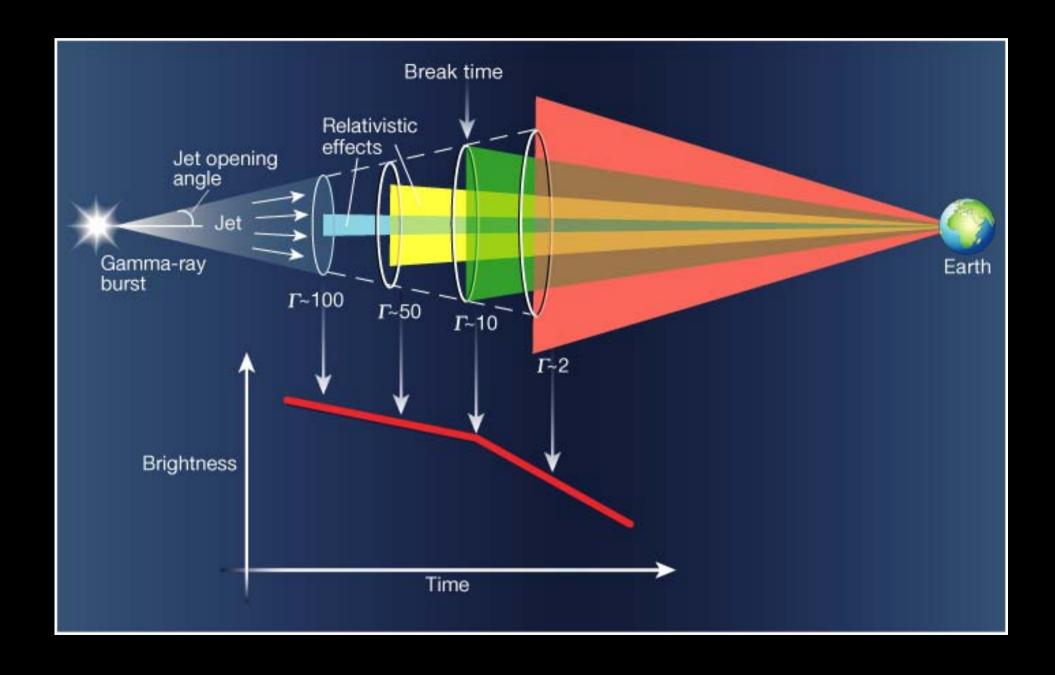
Afterglow emission from GRBs

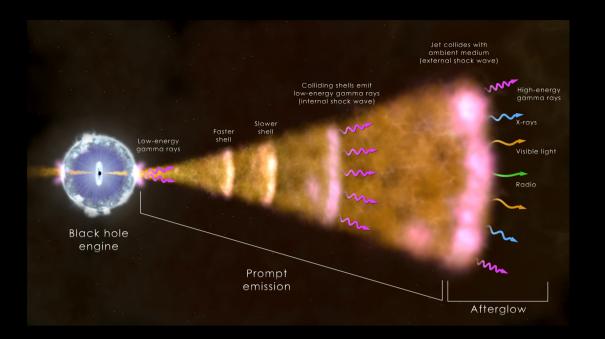
Fabio De Colle

Instituto de Ciencias Nucleares, UNAM



Jet structure



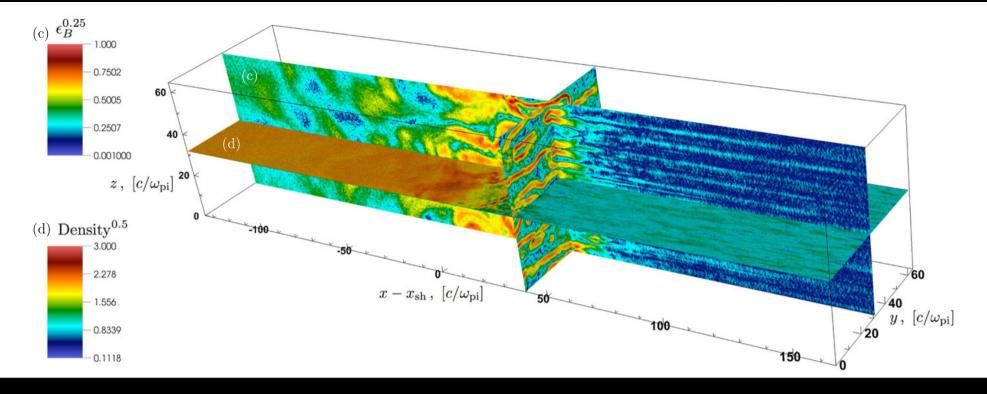


$$\frac{\partial \gamma \rho}{\partial t} + \nabla \cdot (\gamma \rho \vec{v}) = 0$$

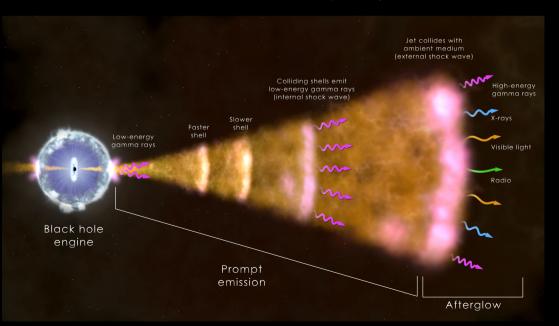
$$\frac{\partial \vec{\rho} \vec{v} \gamma^2 h}{\partial t} + \nabla \cdot (\vec{\rho} \vec{v} \vec{v} \gamma^2 h + p \mathcal{I}) = 0$$

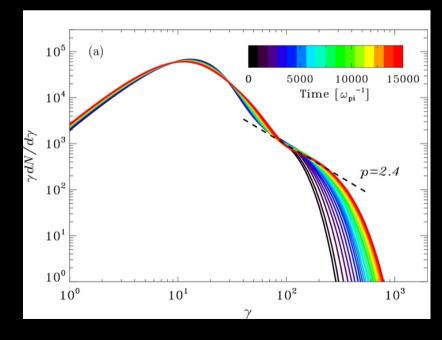
$$\frac{\partial e}{\partial t} + \nabla \cdot ((e + p) \vec{v}) = 0$$

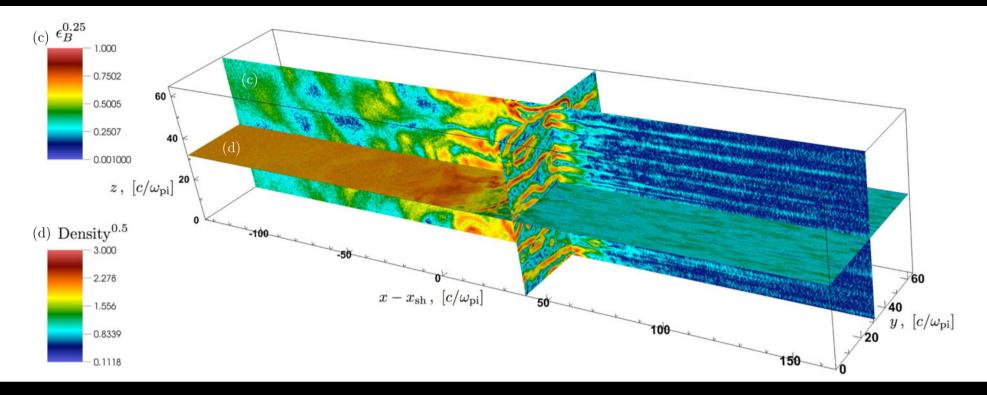




Sironi et al. (2013)



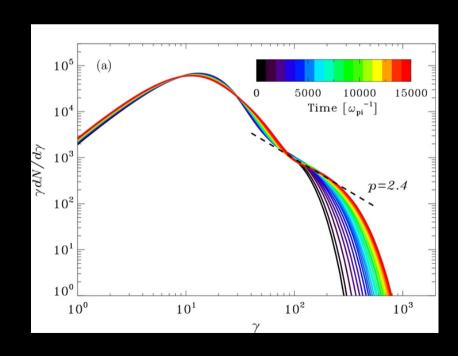


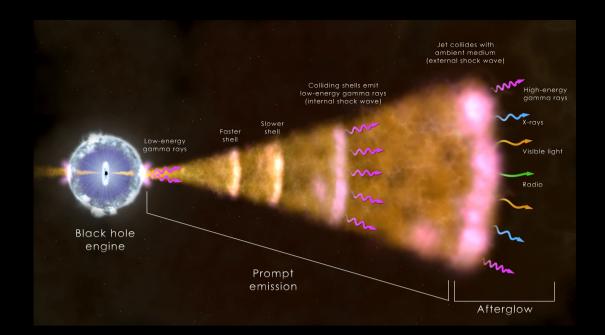


Sironi et al. (2013)

Parametrization of our "ignorance"...

Magnetic field is amplified by the shock: $\epsilon_{_B}$ Electrons are accelerated by the shock: $\epsilon_{_e}$ Electron number density is a power-law: $n \propto \gamma^{-p}$

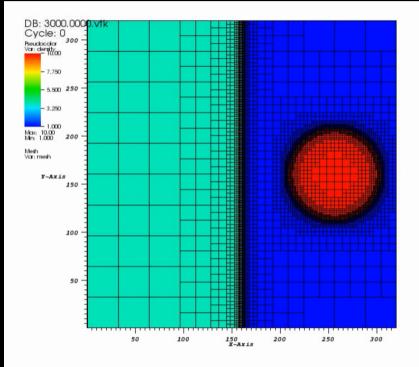




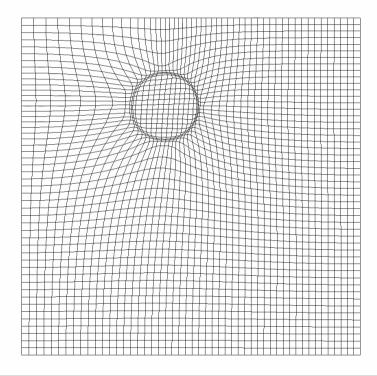
Main problem

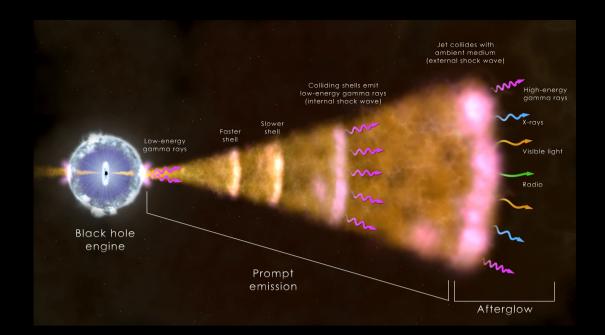
• • •

Huge spatial range (10⁸-10¹⁸ cm)





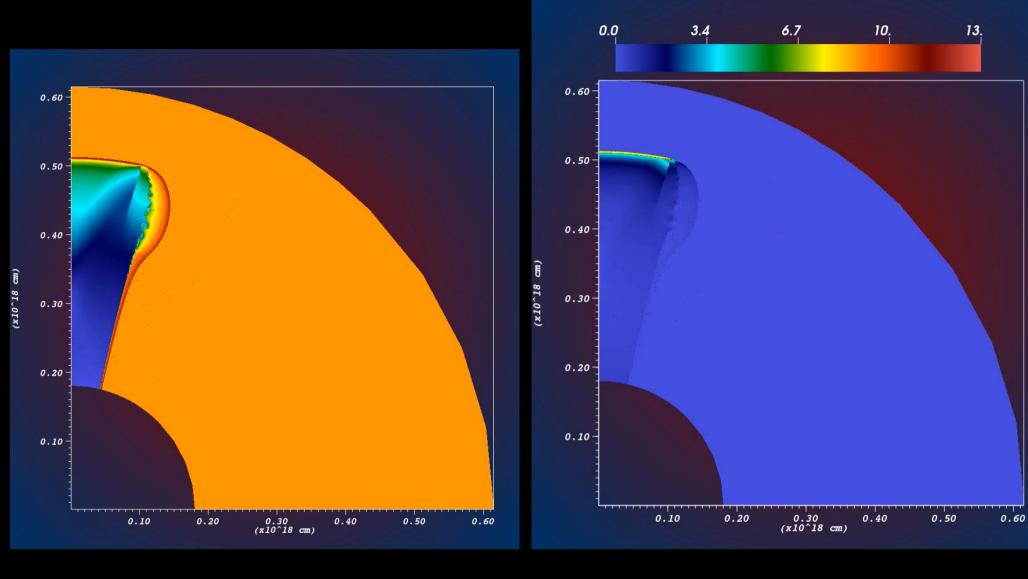




Main problem

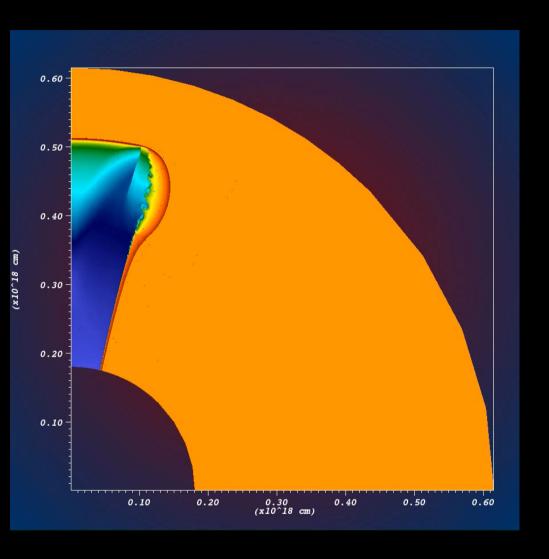
• • •

Huge spatial range (10⁸-10¹⁸ cm)



De Colle et al. (2012)

The dynamics can be rescaled...



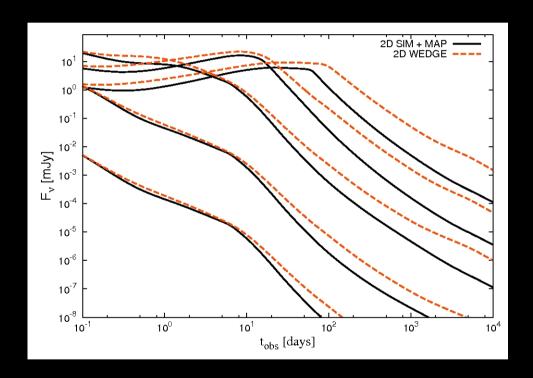
$$E = \int_0^\infty e dV = \int_0^\infty 4\pi r^2 e dr$$

$$E' = kE$$

$$n' = \lambda n$$

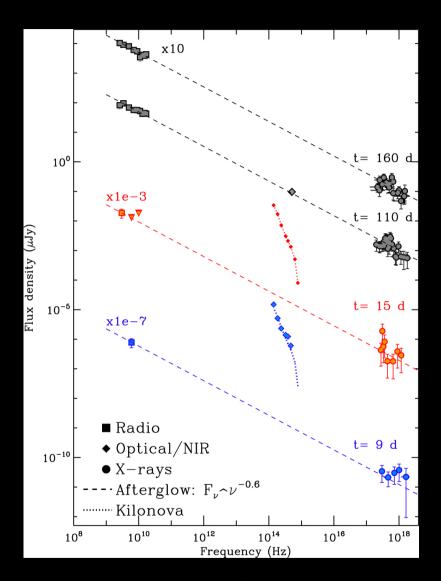
$$\Rightarrow \frac{r'}{r} = \frac{t'}{t} = \left(\frac{k}{\lambda}\right)^{1/3}$$

And the radiation too... (see VanEerten 2012, Granot 2012)

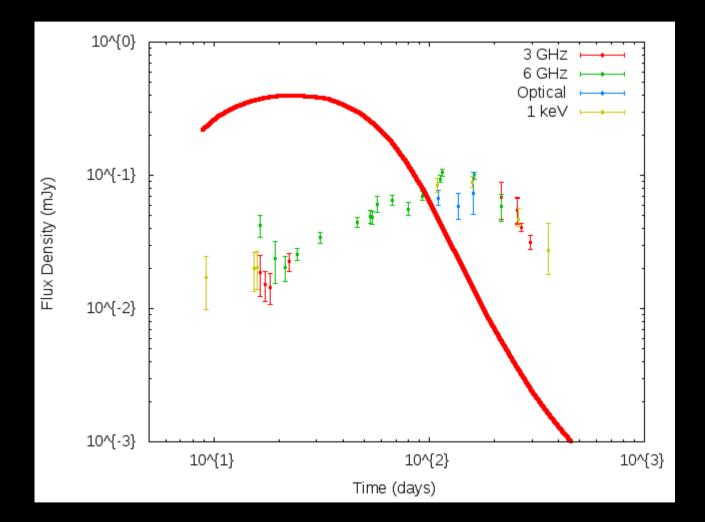


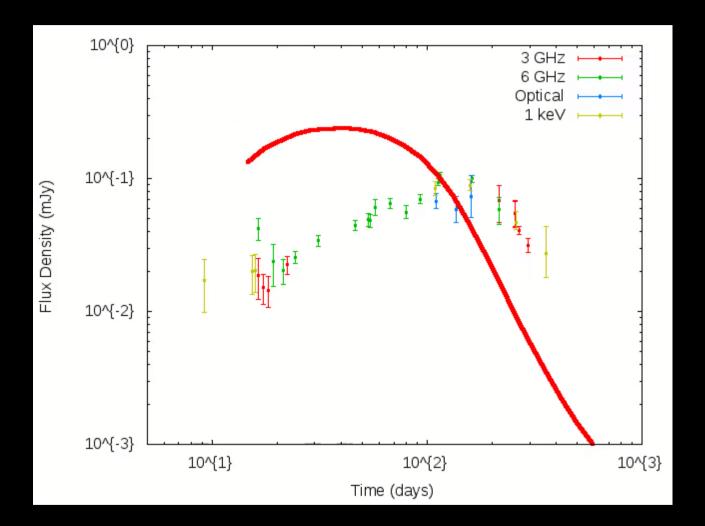
$$\frac{F'}{F} = k\lambda^{(1+p)/4} \left(\frac{\epsilon'_e}{\epsilon_e}\right)^{p-1} \left(\frac{\epsilon'_B}{\epsilon_B}\right)^{(1+p)/4}$$

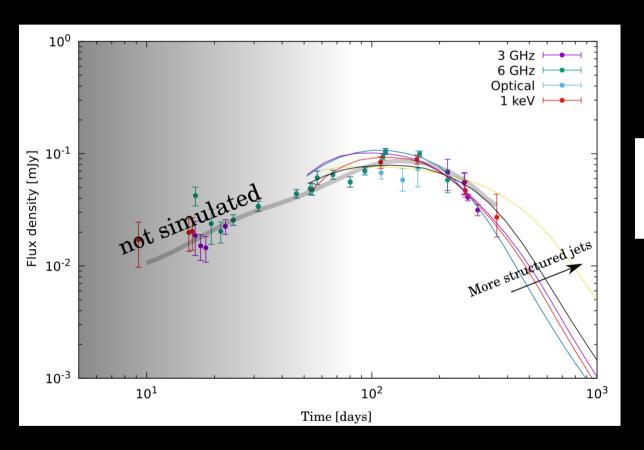
GRB 170817A



Margutti et al. (2018)



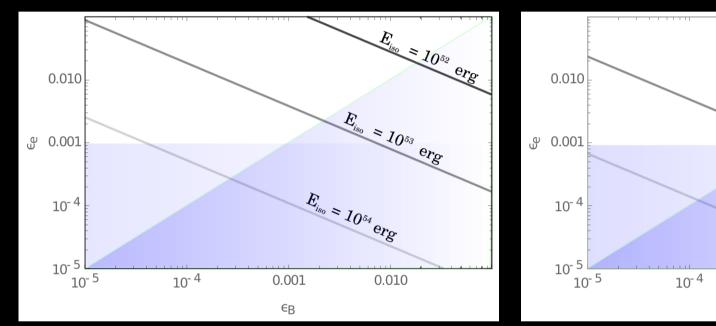


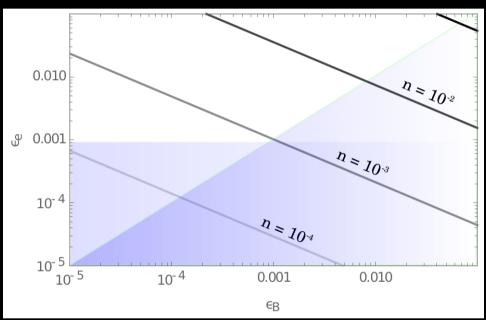


Urrutia, FDC et al. (in preparation)

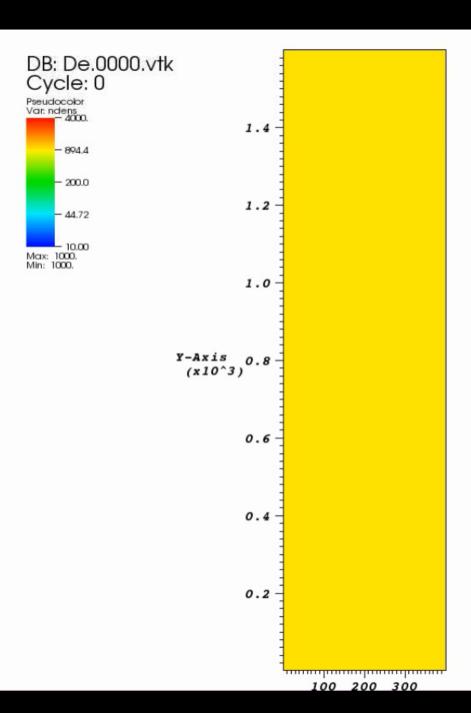
$$\frac{F'}{F} = k\lambda^{(1+p)/4} \left(\frac{\epsilon'_e}{\epsilon_e}\right)^{p-1} \left(\frac{\epsilon'_B}{\epsilon_B}\right)^{(1+p)/4}$$

$$\frac{t'}{t} = \left(\frac{k}{\lambda}\right)^{1/3}$$



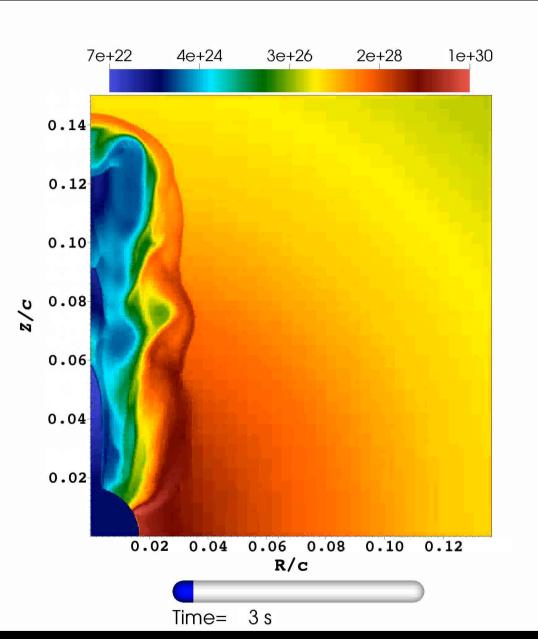


Urrutia, FDC et al. (in preparation)

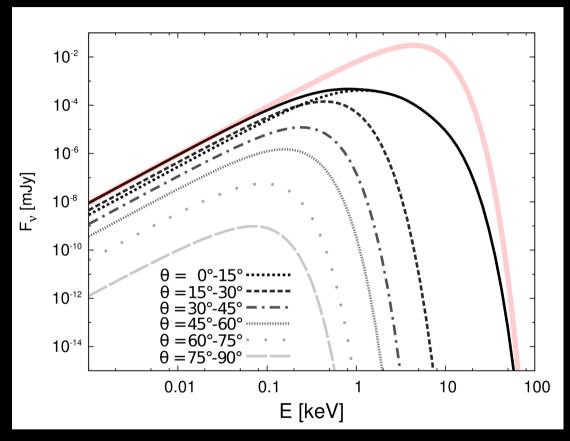


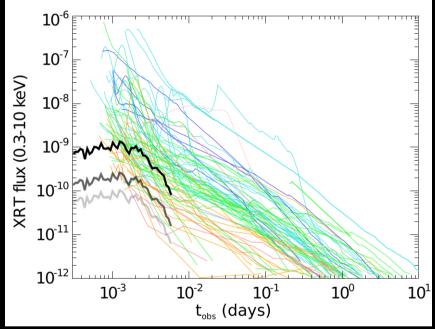
user: fdc Thu Jan 31 12:51:35 2008





De Colle et al. (2018a, 2018b)





De Colle et al. (2018a)

LETTER

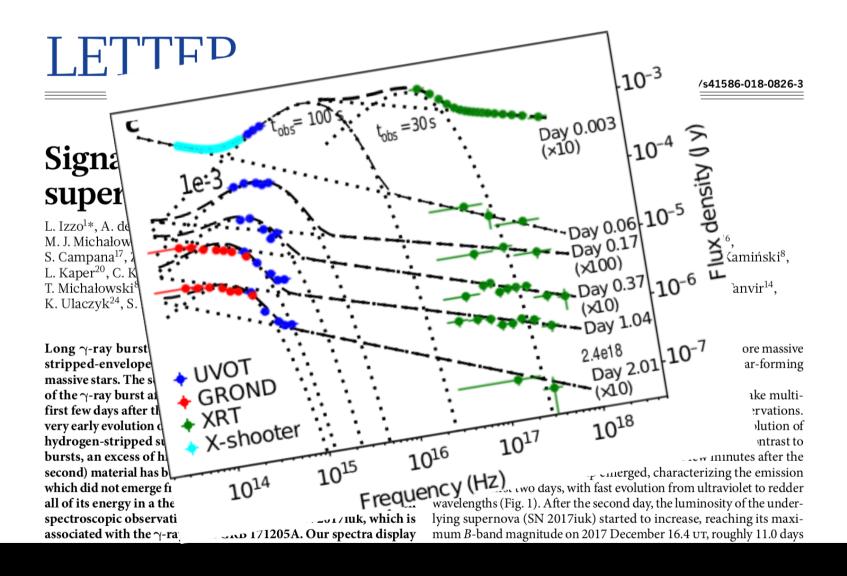
Signatures of a jet cocoon in early spectra of a supernova associated with a γ -ray burst

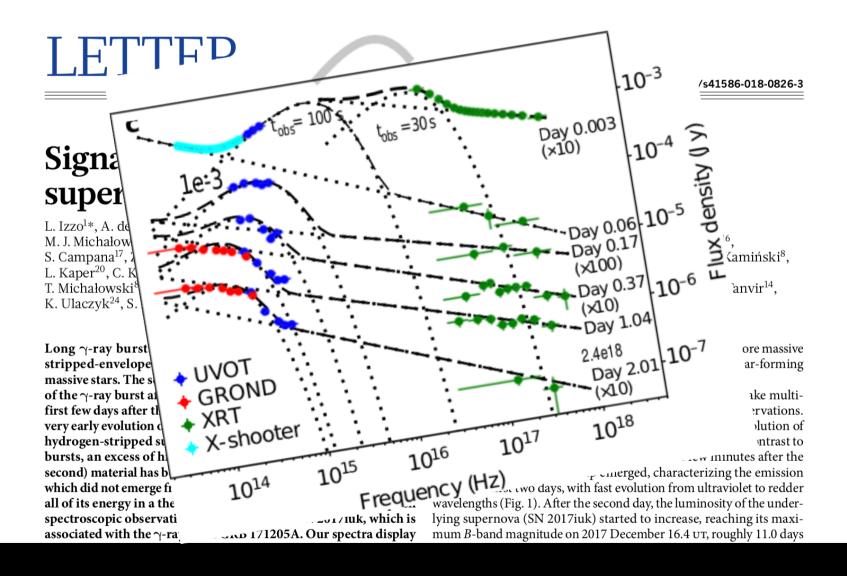
L. Izzo¹*, A. de Ugarte Postigo¹,², K. Maeda³, C. C. Thöne¹, D. A. Kann¹, M. Della Valle¹,⁴,⁵,⁶, A. Sagues Carracedo⁻, M. J. Michałowski³, P. Schady⁵,¹0, S. Schmidl¹¹, J. Selsing²,¹2,¹3, R. L. C. Starling¹⁴, A. Suzuki¹⁵, K. Bensch¹, J. Bolmer⁵,¹6, S. Campana¹⁻, Z. Cano¹, S. Covino¹⁻, J. P. U. Fynbo¹²,¹3, D. H. Hartmann¹³, K. E. Heintz¹²,¹3,¹9, J. Hjorth², J. Japelj²⁰, K. Kamiński³, L. Kaper²⁰, C. Kouveliotou²¹,²², M. Krużyński³, T. Kwiatkowski³, G. Leloudas²,²³, A. J. Levan²⁴, D. B. Malesani²,¹²,¹³, T. Michałowski³, S. Piranomonte²⁵, G. Pugliese²⁰, A. Rossi²⁶, R. Sánchez-Ramírez²⁻, S. Schulze²³, D. Steeghs²⁴, N. R. Tanvir¹⁴, K. Ulaczyk²⁴, S. D. Vergani²ց & K. Wiersema¹⁴,²⁴

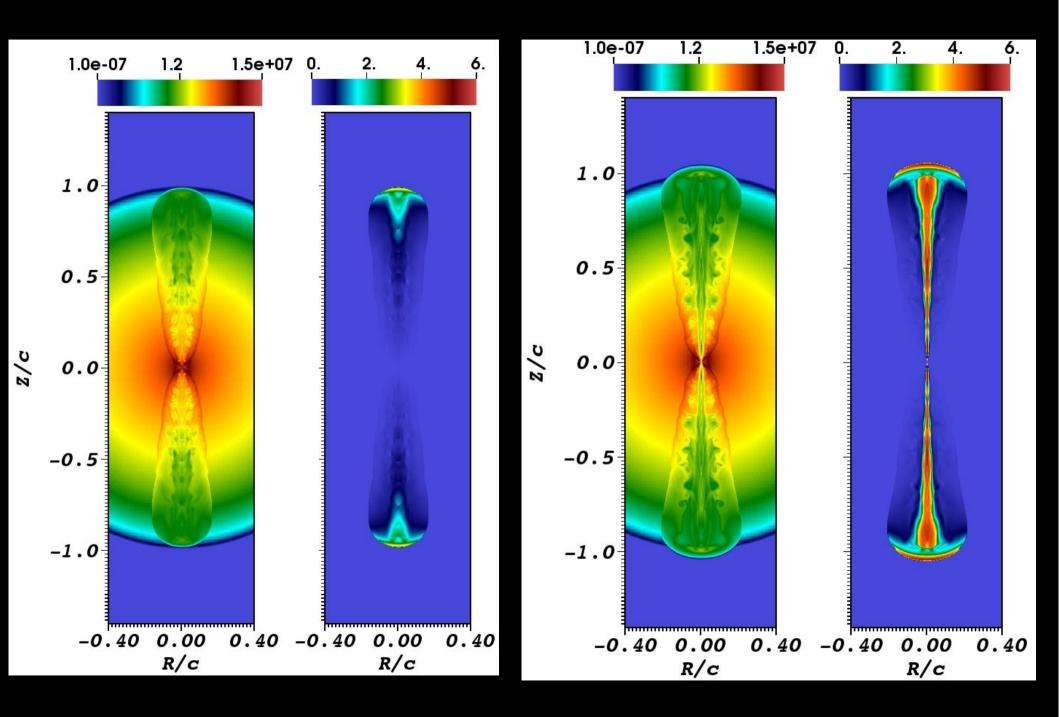
Long γ -ray bursts are associated with energetic, broad-lined, stripped-envelope supernovae^{1,2} and as such mark the death of massive stars. The scarcity of such events nearby and the brightness of the γ -ray burst afterglow, which dominates the emission in the first few days after the burst, have so far prevented the study of the very early evolution of supernovae associated with γ -ray bursts³. In hydrogen-stripped supernovae that are not associated with γ -ray bursts, an excess of high-velocity (roughly 30,000 kilometres per second) material has been interpreted as a signature of a choked jet, which did not emerge from the progenitor star and instead deposited all of its energy in a thermal cocoon⁴. Here we report multi-epoch spectroscopic observations of the supernova SN 2017iuk, which is associated with the γ -ray burst GRB 171205A. Our spectra display

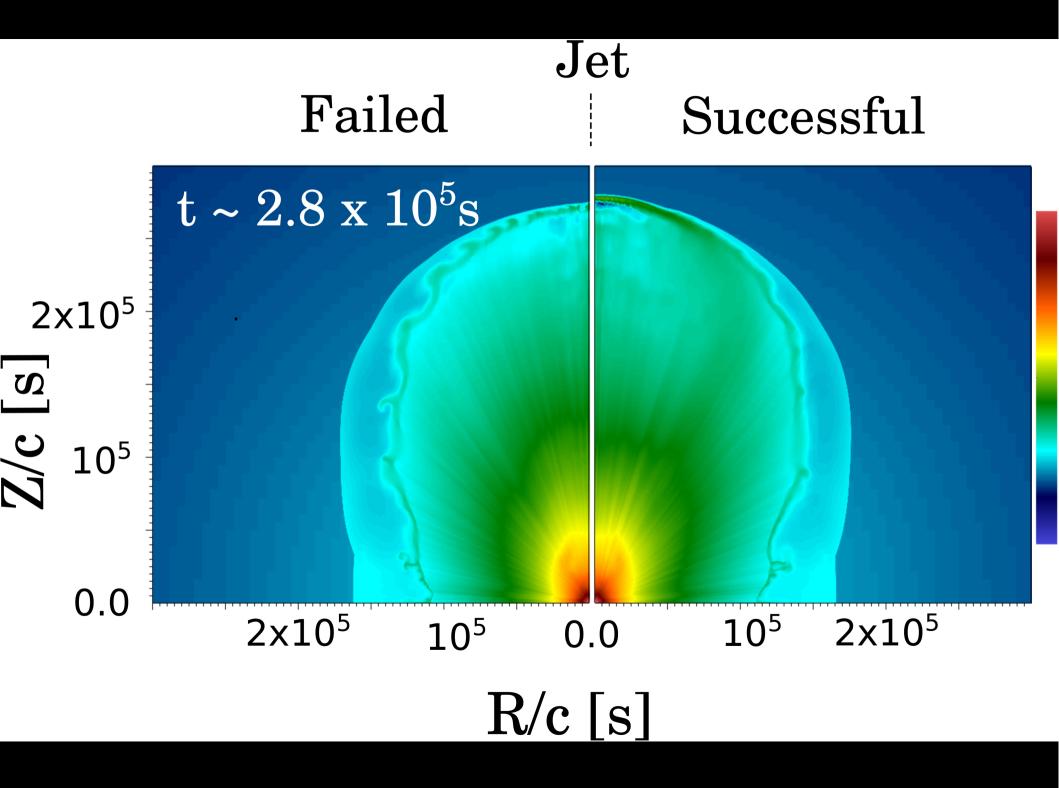
metallicity ($12 + \log(O/H) = 8.41$; Methods). It is much more massive than typical GRB hosts, which are normally metal-poor, star-forming dwarf galaxies, particularly at low redshift¹².

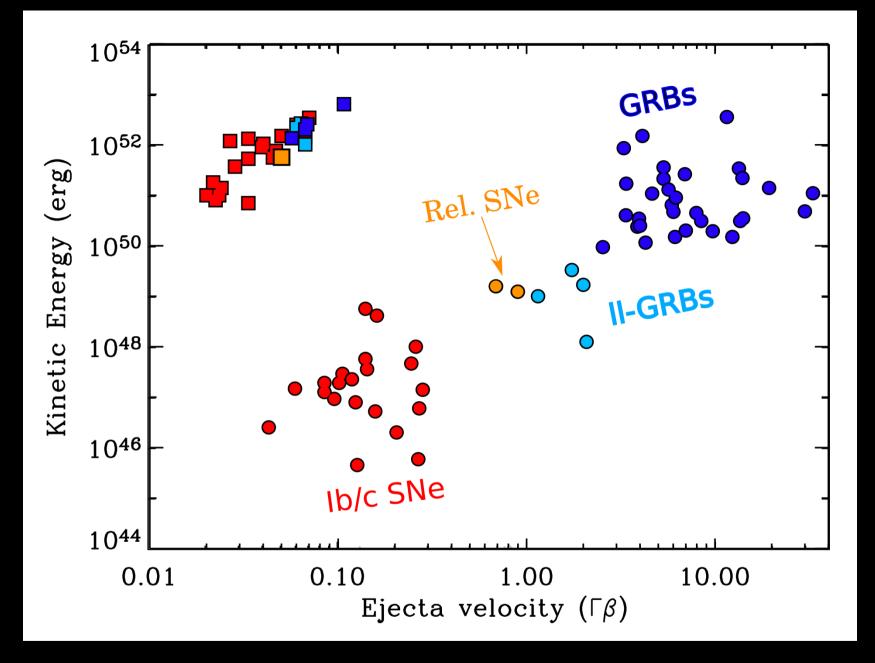
The proximity of GRB 171205A motivated us to undertake multi-wavelength photometric and spectroscopic follow-up observations. The light curve exhibits unusual behaviour, with colour evolution of the optical and ultraviolet emission at very early phases, in contrast to the rapid decay observed in X-ray emission. A few minutes after the burst, a first light-curve bump emerged, characterizing the emission during the first two days, with fast evolution from ultraviolet to redder wavelengths (Fig. 1). After the second day, the luminosity of the underlying supernova (SN 2017iuk) started to increase, reaching its maximum *B*-band magnitude on 2017 December 16.4 ut, roughly 11.0 days



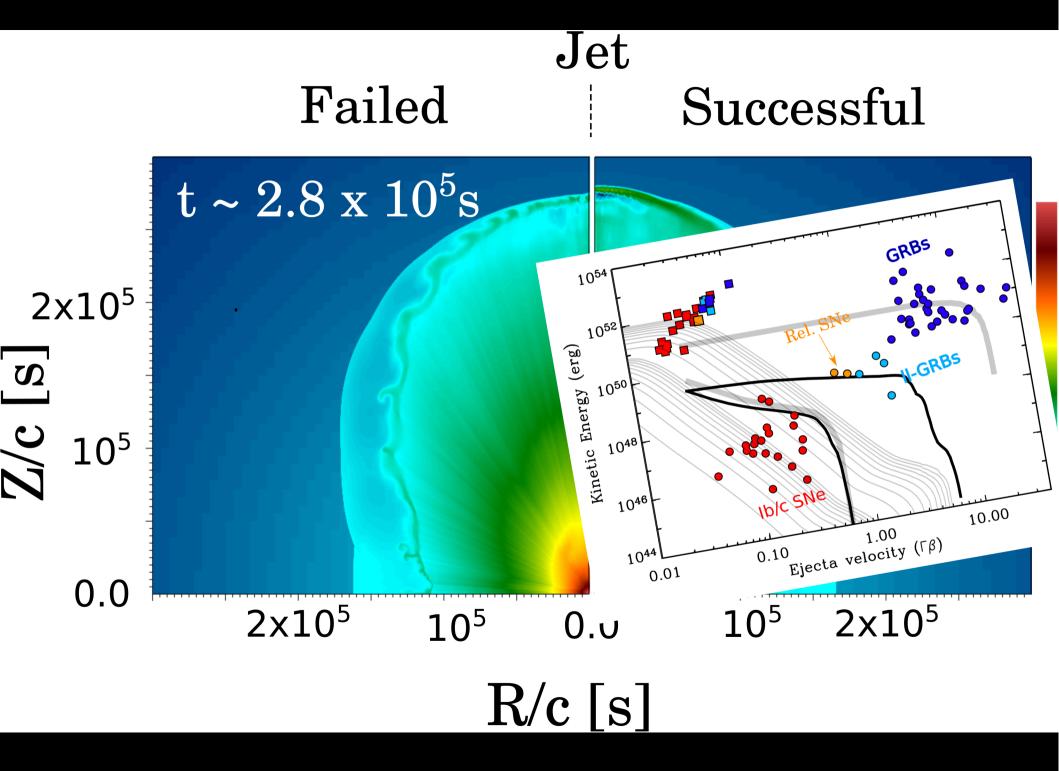


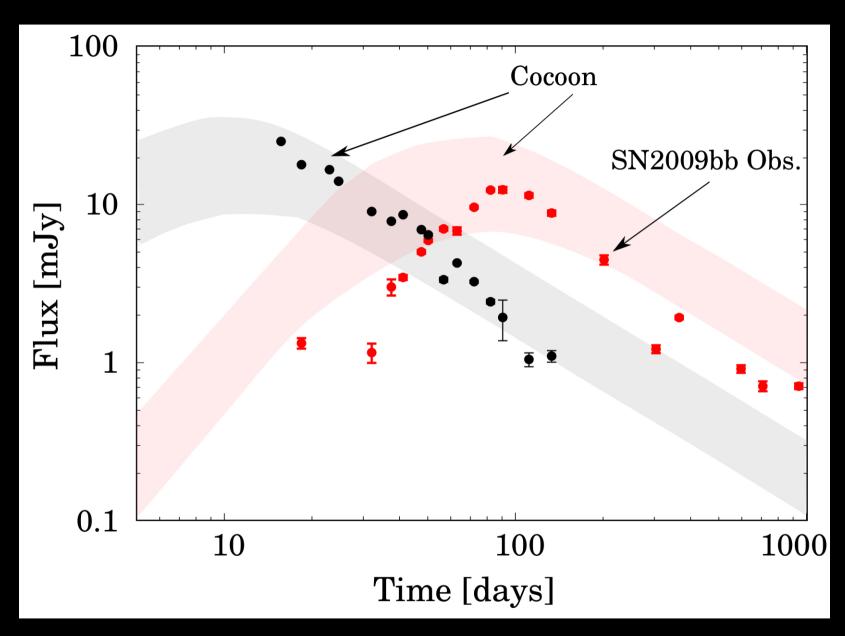






Soderberg et al. (2010), Margutti et al. (2014)





De Colle et al. (2018b)

Conclusions

- Determination of the physical parameters from the observations is complicated
- Cocoon emission is important!
 - Thermal emission can help us to understand the progenitor structure
 - Non-thermal emission contributes and in some cases dominates the emission...