

Minimal field theories for active systems

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Phénomènes hors d'équilibre
IPhT
02 October 2018



3 ways into non-equilibrium

Group

Active Matter

Tracers in a
bacterial bath

Incomplete
Phase
Separation

- **Very slow relaxation towards equilibrium**

So slow that the transient can be considered a bona fide state of matter

Ex: glasses, ...

- **Global external forcing of an otherwise equilibrium system**

imposing T , μ gradients, mechanical stirring ...

Ex: turbulence, ...

- **Out of equilibrium in the bulk**

Each microscopic component is out of equilibrium

Ex: Active matter,

Group

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H. Chaté



I. Dornic



C. Nardini



A. Patelli
(post-doc, '18-'19)



G. Fausti
(Ph.D., '18-'21)



B. Ventejou
(Ph.D., '18-'21)

Collaborations

Theory/Modeling

- M. Cates (Cambridge, phenomenology of active systems)
- A. Morozov (Edinburgh, hydrodynamic interactions)
- J. Tailleur, F. van Wijland (Paris 6, minimal models)
- B. Delamotte (Paris 7, universality)
- X. Shi (Soochow, large scale numerics)
- E. Bertin (Grenoble, continuum theories)
- K. Mallick (IPhT, Palm funding)
- ...

Experimentalists

- M. Sano (Tokyo, bacterial thin films)
- H.P Zhang (Shanghai, bacterial thin films)
- W. Poon (Edinburgh, 3d bacterial suspensions)
- G. Duménil (Pasteur, in vivo bacterial behavior)
- G. Theraulaz (Toulouse, self-organization in fish groups)
- L. Blanchoin, M. Théry (Grenoble, cytoskeleton, Impulsion funding)

Active Matter: a brief introduction

Group

Active Matter

Tracers in a bacterial bath

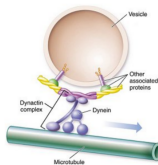
Incomplete Phase Separation

Any active unit is able to extract energy from the environment and dissipate it to self-propel

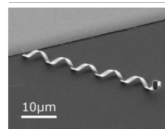
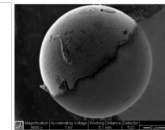
Human scale



Micrometer-scale (biological)



Micrometer-scale (synthetic)



Active matter: a brief introduction

Novel collective phenomena

[Chen et al., Nature, 2017]

- Most of the work on systems where units are μm -sized (bacteria, microtubules, cells – tissues, synthetic active colloids, ...)
- Self-organized structures, collective intelligence, smart materials, ...

Active matter: a brief introduction

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Understanding requires theoretical physics tools

(not only, of course!)

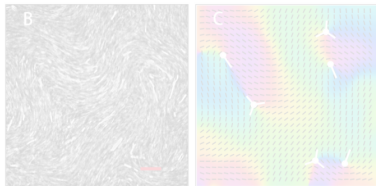
And gives back to theoretical physicists new – intriguing – problems to solve



- In our group [Li et al., PNAS 2018]: quantitative prediction of defects speed and shape; comparison with experiments on 2d bacterial suspensions

- 2-d active nematics: topological defects [Dogic et al., Nature 2013, ...]
- $+1/2$ defects are motile

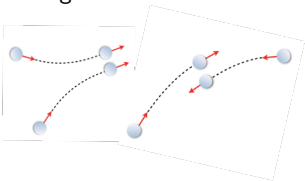
2d swarming filamentous bacteria



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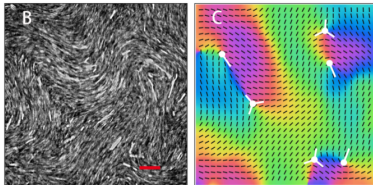
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Passive tracer in a bacterial suspension

Group

Active Matter

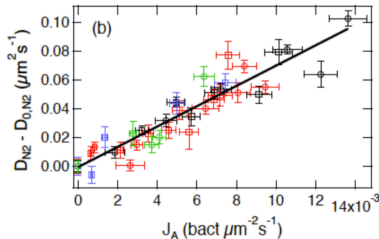
Tracers in a bacterial bath

Incomplete Phase Separation

Tracer diffusivity in 3-d bacterial suspension?

Extensive experimental work

(Wu et al. PRL, 2000; Goldstein et al. PRL, 2009; Poon et al., PRE(R) 2013; Polin et al. Nat. Phys. 2015, ...)



Very low densities (bacteria packing fraction $\phi \sim 0.01 - 0.1\%$)
 \sim bacteria don't interact with each other

Passive tracer in a bacterial suspension

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Active Matter

Tracers in a bacterial bath

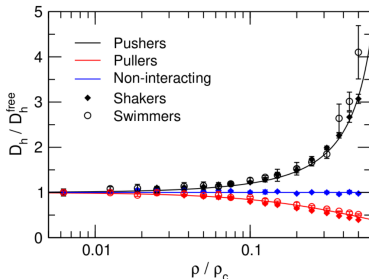
Incomplete Phase Separation

Can we describe tracer diffusivity at higher densities (but still low $\phi \sim 1\%$)?

- Explicit coarse graining (hydrodynamics \sim long-range)
- Stochastic continuum theory for $f(\mathbf{r}, \mathbf{p}, t) \sim$ Kinetic theory for plasma

$$\partial_t f + \nabla_{\mathbf{r}}(\dot{\mathbf{r}}f) + \partial_{\mathbf{p}}(\dot{\mathbf{p}}f) = \text{noise}$$

- $\dot{\mathbf{r}} = v_s \mathbf{p} + U_{\text{mean-field}}, \dot{\mathbf{p}} = \dots$



- So far: analytical and numerical predictions
- $\rho_c \sim \phi \sim 1\%$
- Experiments in progress in W. Poon group (Edinburgh)

Minimal field theory for phase separation



M.E. Cates



E. Tjhung



F. Caballero

- Tjhung et al., PRX 031080, 2018
- F. Caballero et al., PRL 020601, 2018
- F. Caballero et al., ArXiv:1809.10433, 2018

Simplest active particles

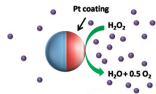
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Active Matter

Tracers in a
bacterial bathIncomplete
Phase
Separation

- Self-propelled spherical particles
- Two-body steric interactions
- No alignment

Ex: Janus colloids



- microscale: persistent Brownian motion (speed v , rotational diffusivity $1/\tau$)
- large-scales: Brownian motion with diffusivity $D \sim v^2\tau$



Can many (so-simple!) particles do anything interesting?

Simplest active particles

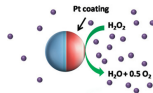
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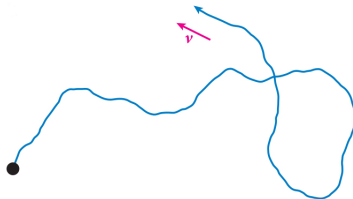
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Interactions: TRS breakdown

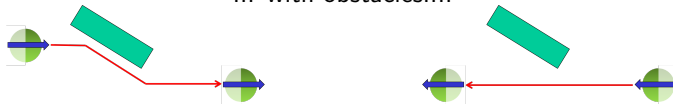
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Active Matter

Tracers in a bacterial bath

Incomplete Phase Separation

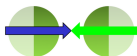
... with obstacles...



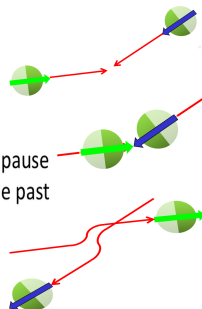
... or among particles...



long pause
to rotate



short pause
to slide past

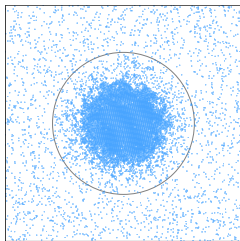


Measuring breakdown of TRS

Entropy production rate: $S = \lim_{T \rightarrow \infty} \frac{1}{T} \log \left(\mathcal{P}_F / \mathcal{P}_B \right)$

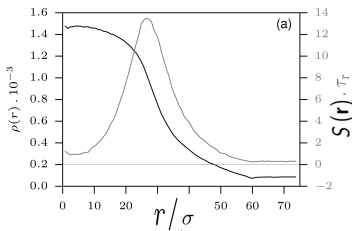
$$S = \frac{\sqrt{\tau}}{2D} \sum_{ij} \left\langle (p_i - p_j)^3 : \nabla^3 u(r_i - r_j) \right\rangle \quad U = \frac{1}{2} \sum_{ij} u(r_i - r_j)$$

$S(r)$: contribution to S given by particles in a mesoscopic box centered in r



$$D = 1, \tau = 10, N = 10^4$$

$S(|r|)$ for different interface widths



Part I: Phase separation in active systems

Phase separation of purely repulsive particles

At large scales

~

liquid-gas phase separation

[Tailleur et al. PRL 2008; Ann. Rev. Cond. Mat. 2015]

Effective equilibrium idea
thrived huge number of works

[Marchetti, Speck, Joanny, Di Leonardo, Brader, Redner,

Szamel, ...]

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Clusters and bubbly phase separation

Cluster Phases

Bubbly phase separation

Experiments – Janus particles

[J. Palacci et al., *Science*, 2013; Buttinoni et al, PRL 2013; J Schwarz-Linek et al, PNAS 2012; I Theurkauff et al, PRL 2012]

Simulations – purely repulsive
particles

[Stenhammar et al., *Soft Matter*, 2014]

Cluster phases / bubbly phase
separation generic?

Are they related?

Large-scale properties?

[PRX 031080, 2018]

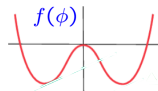
Minimal Field Theory

Minimal modification of Model B to break Time-Reversal Symmetry

$$\dot{\phi} = -\nabla \cdot (\mathbf{J} + \eta) \quad \mathbf{J} = -\nabla\mu + \zeta(\nabla^2\phi)\nabla\phi$$

$$\mu(\mathbf{r}) = \frac{\delta\mathcal{F}}{\delta\phi(\mathbf{r})} + \lambda|\nabla\phi|^2 \quad \mathcal{F} = \int d\mathbf{r} [f(\phi) + K|\nabla\phi|^2]$$

$$\langle \eta(\mathbf{r}, t)\eta(\mathbf{r}', t') \rangle = 2D \mathbf{1} \delta(\mathbf{r} - \mathbf{r}') \delta(t - t')$$



- J : current
- μ : chemical potential
- ~~μ : chemical potential~~

Obtained from

- symmetries and gradient expansion
- explicit coarse graining of particles models

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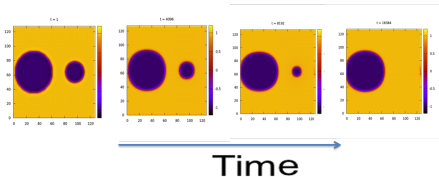
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Ostwald Ripening in equilibrium

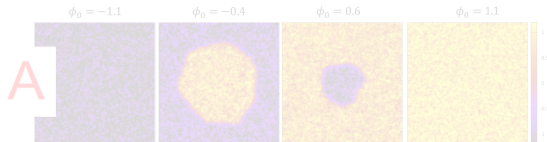
Model B: Ostwald Ripening



Reason of:

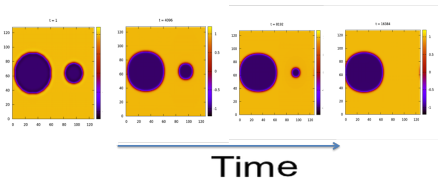
- Pastis + water \implies opaque mixture
- Re-crystallization of water in ice-cream giving crunchy texture

Pressure difference inside two bubbles/clusters of different radii
(and no momentum conservation)



Ostwald Ripening in equilibrium

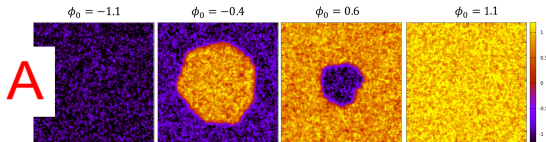
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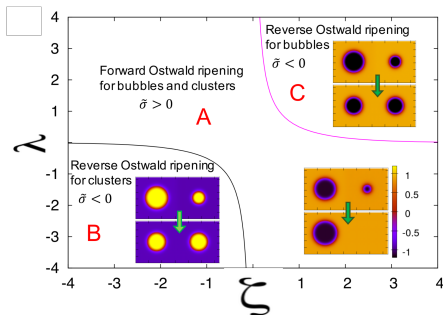
Reason of:

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Mean-field analysis (analytical)



$$\dot{R} \propto \frac{\sigma}{R} \left(\frac{1}{R_c} - \frac{1}{R} \right) + \mathcal{O} \left(\frac{1}{R^3} \right)$$

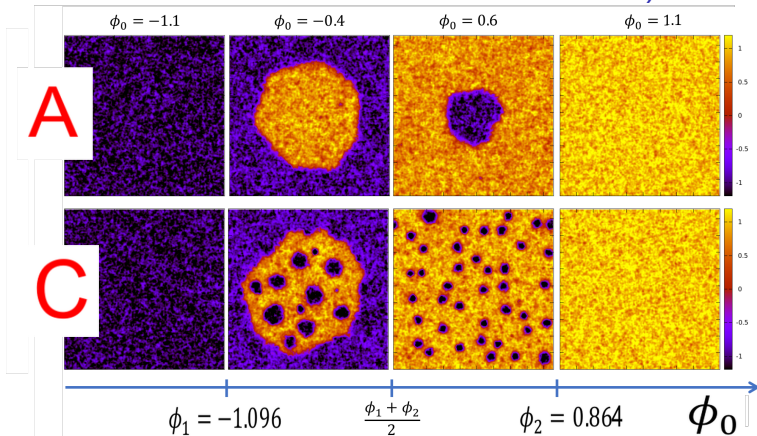
Formally (but not physically!) equivalent to
Ostwald Ripening in equilibrium

$$\sigma = \frac{K}{\zeta - 2\lambda} \left[\zeta \sigma_A e^{\frac{\zeta - 2\lambda}{K} \phi_+} - 2\lambda \sigma_B \right]$$

$$\sigma_A = \int_0^\infty \phi'(y)^2 dy$$

$$\sigma_B = \int_0^\infty \phi'^2(y) e^{\frac{\zeta - 2\lambda}{K} y} dy.$$

Phase diagram (with noise, numerical)



B Symmetry $(\phi, \lambda, \zeta) \rightarrow -(\phi, \lambda, \zeta)$

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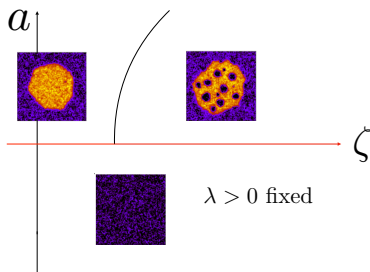
Active Matter

Tracers in a
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Incomplete
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Manifest TRS breaking

Does TRS-breaking survives at criticality?



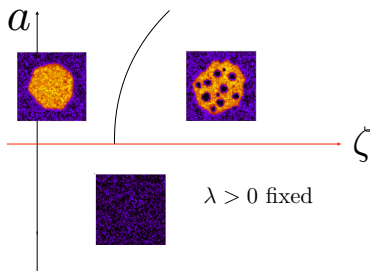
**Critical transition to phase
separation**
(densities of the two phases almost
identical)

$$\dot{\phi} = \nabla^2 [-a\phi + u\phi^3 - K\nabla^2\phi + \lambda|\nabla\phi|^2] - \zeta\nabla \cdot [\nabla^2\phi\nabla\phi] + \nabla \cdot \eta$$

Renormalisation Group analysis

Dimensional analysis suggest: TRS-breaking terms (λ, ζ) irrelevant

Does TRS-breaking survives at criticality?



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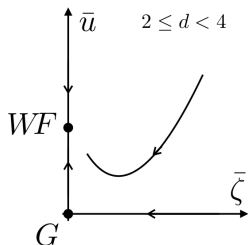
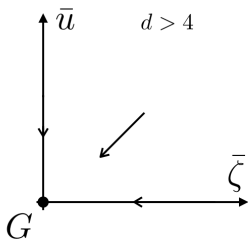
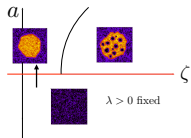
Renormalisation Group analysis

Dimensional analysis suggest: TRS-breaking terms (λ, ζ) irrelevant

1-loop results: bulk phase-separation

Dimensional regularization close to
 $d = 4$

Bare parameters: region A
(bulk phase separation)

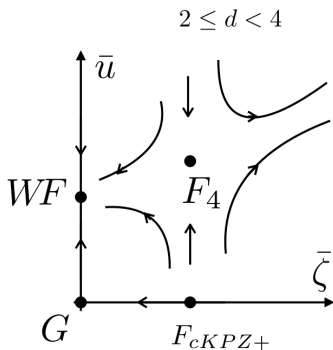
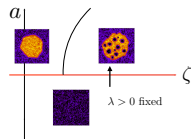


Equilibrium fixed points of the RG flow

- G : Gaussian
- WF : Wilson-Fisher

1-loop results: bubbly phase-separation

Bare parameters: region B,C
(micro-/bubbly- phase separation)



- F_{cKPZ+} : transition mounding phase in a surface growth model (PRL 020601 2018)

(Probably) new non-equilibrium universality class describes the transition to bubbly/microphase separation

[ArXiv:1809.10433 (2018)]

Conclusions

Take home message:

- ~~Working in Active matter allows to give talks containing nice movies~~
- Theoretical physics tools are crucial for understanding Active matter
- Even the simplest active systems (spherical particles, no alignment): intriguing large scale properties

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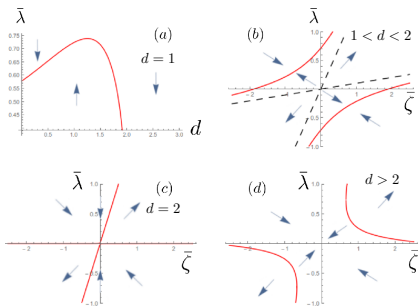
1-loop results

Local free energy absent: $f = 0$: cKPZ+ equation

$$\dot{\phi} = \nabla^2(-K\nabla^2\phi + \lambda|\nabla\phi|^2) - \zeta\nabla \cdot (\nabla^2\phi\nabla\phi) + \nabla \cdot \eta$$

model for growing surfaces with mass conservation

(in the comoving frame)



RG and numerics suggest the presence of a **strong coupling regime** akin to KPZ equation (rough surfaces)