

# Investigation of Gaseous Detectors with Laser Induced Electrons

Lothar Naumann



**hzdr**

HELMHOLTZ  
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ROSSENDORF

# Introduction

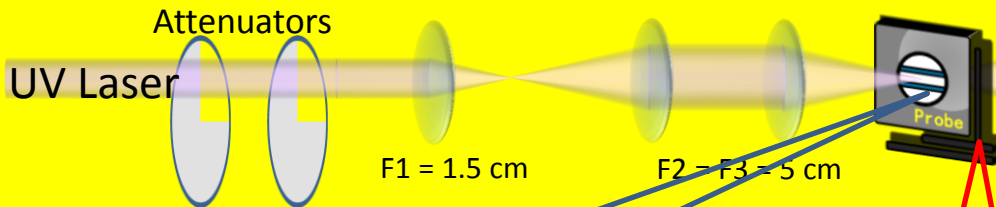
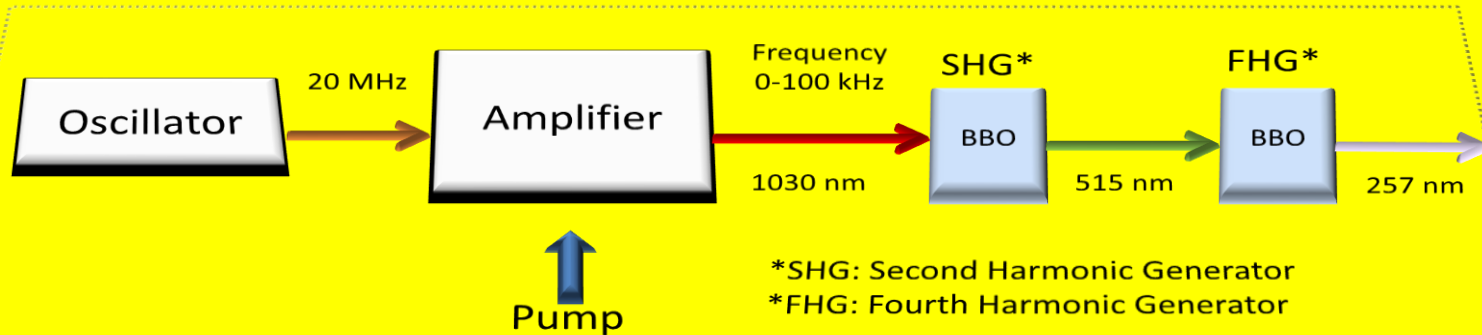
- Application of UV laser beams for calibration and surveying of Wire Chambers since 40 years *[M. Anderhub et al.; NIM 166 (1979)]*  
Resistive Plate Chambers since 30 years *[E. Gorini et al.; NIMA 425 (1999)]*
- RPC operating in strong and homogeneous electric fields at atmospheric pressure → relevant gas parameters obtained in reduced electric field
- Drift Detectors operating in inhomogeneous electric field-topologies → relevant gas parameters obtained in simplified electric field topologies

# Introduction

**Micro-plasma** creation and precise **micro-positioning** inside the active volume of gaseous detectors allows to improve the detector tests:

- **Resistive Plate Chambers:** timing and trigger RPC samples
- **Mini Drift Cells:** HADES-like MDC topology
- **Drift Tube:** for laser facility calibration purposes

# HZDR Laser Facility



## UV Laser Properties:

wavelength: 257nm (4.8 eV)  
frequency : 0 -100 kHz  
pulse width :  $\leq 2$  ps  
beam waist:  $\geq 5 \mu\text{m}$   
Debye length:  $\pm 100 \mu\text{m}$   
Intensity:  $\leq 10^{12} \text{ W/cm}^2$

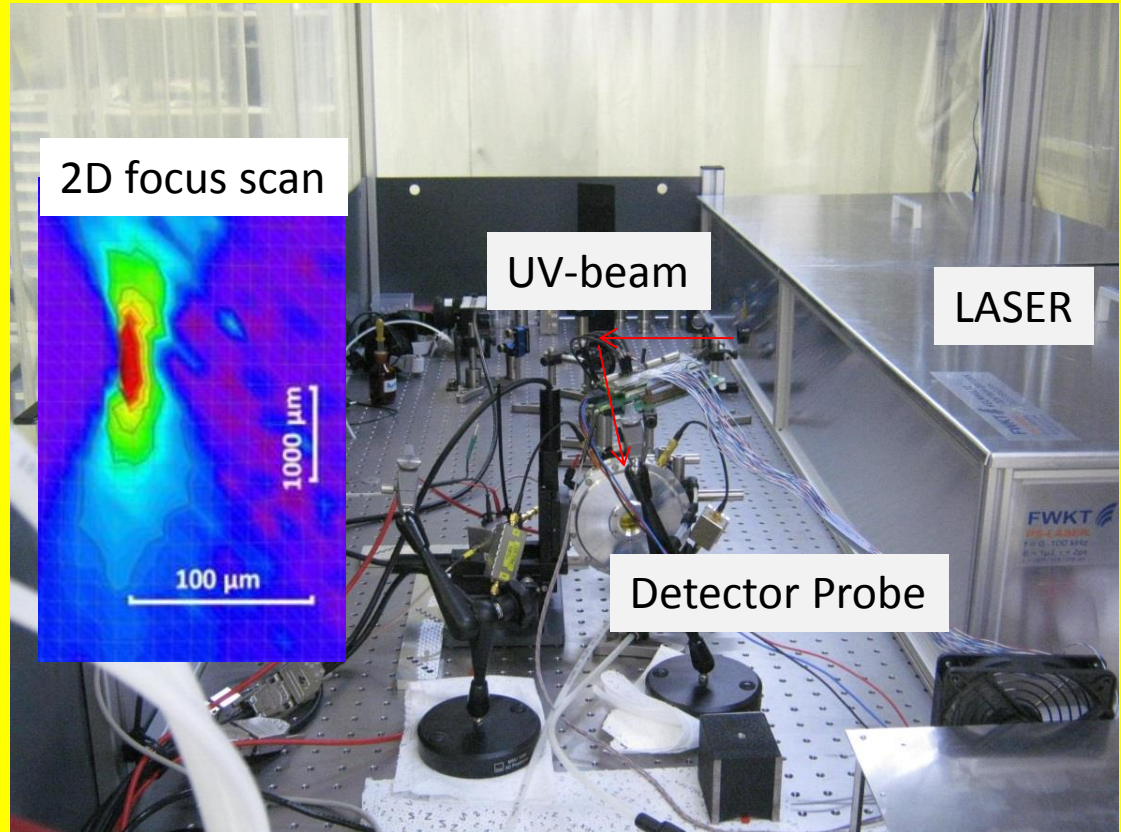
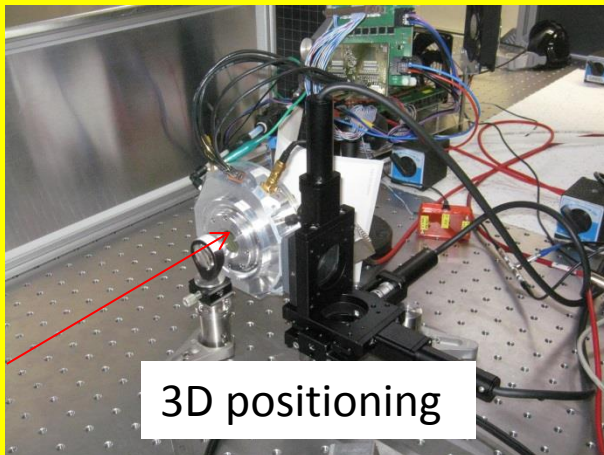
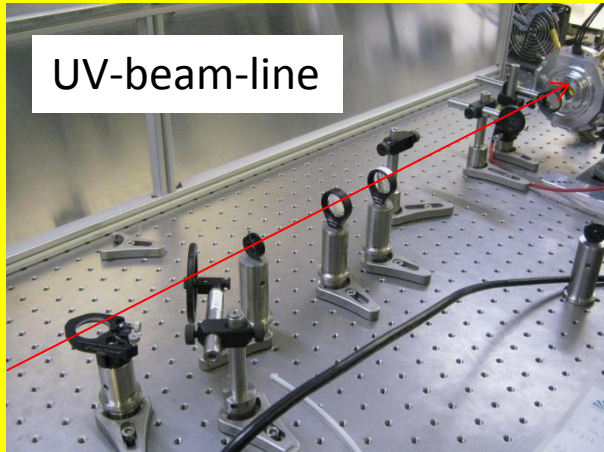
## Dedicated detector probes:

Resistive Plate Chamber  
Mini Drift Cell  
Drift Tube

## 3D Step Driver :

range: 8 mm  
resolution:  $\leq 1 \mu\text{m}$

# HZDR Laser Facility

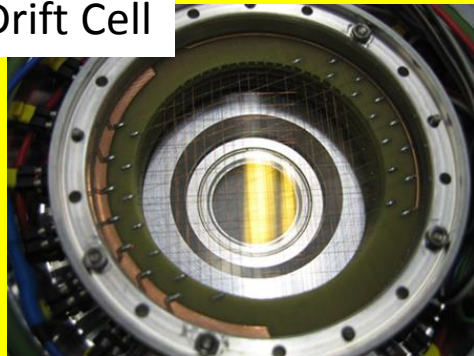


# HZDR Laser Facility

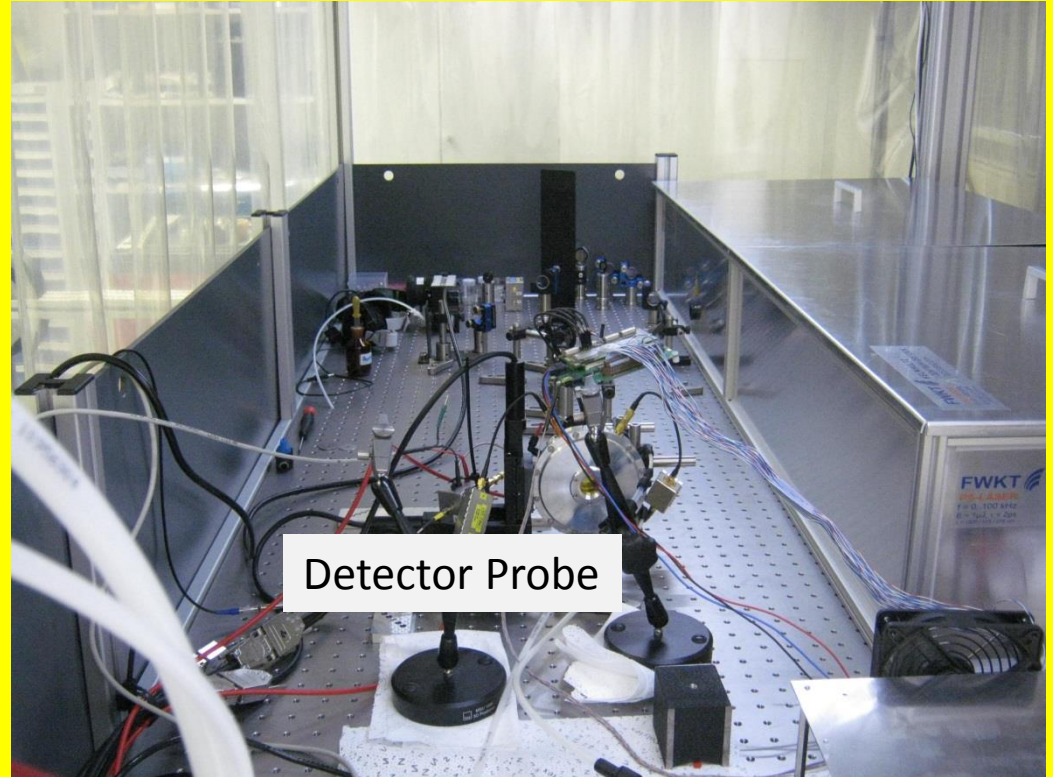
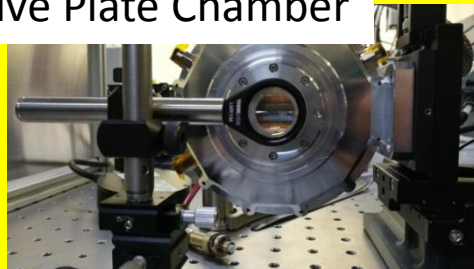
Drift Tube



Mini Drift Cell



Resistive Plate Chamber

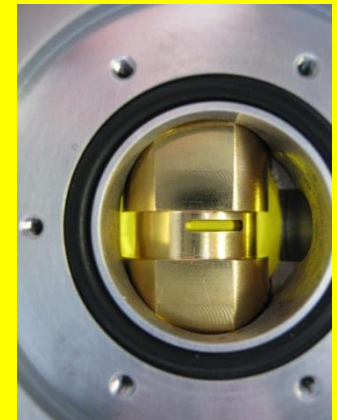
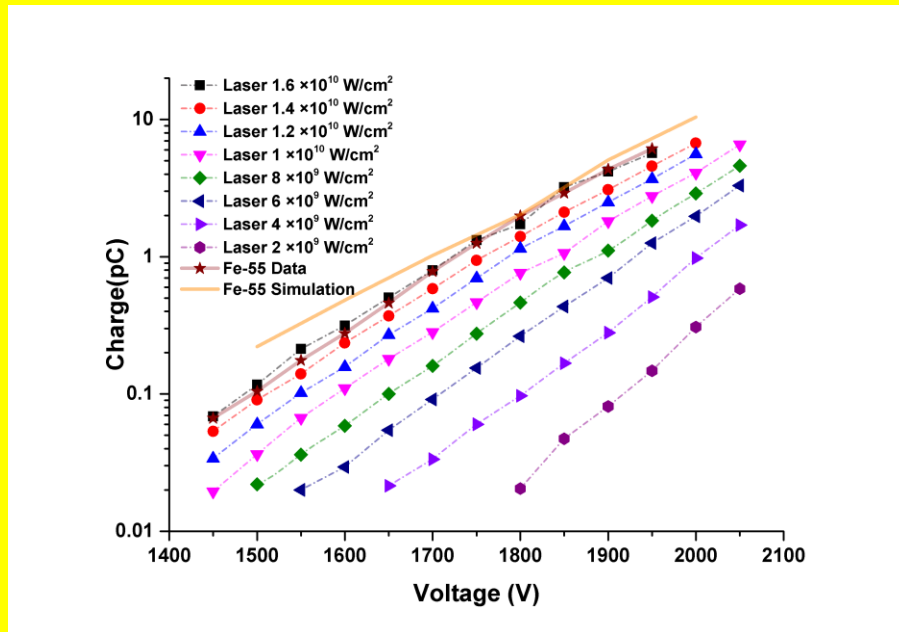
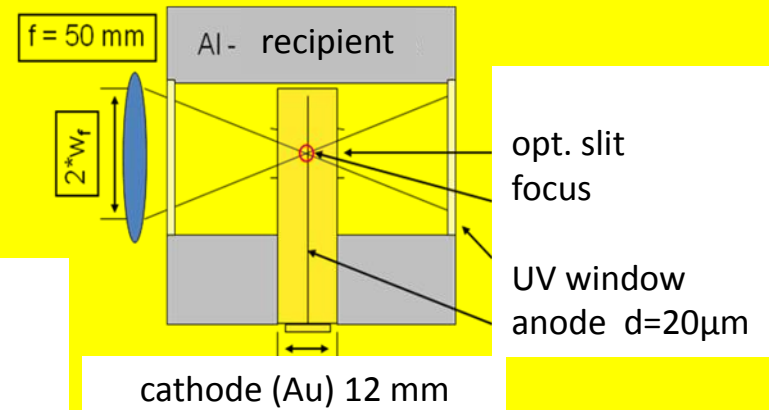


Detector Probe

# Drift Tube

$\lambda = 257 \text{ nm} \rightarrow \text{photon energy } 4,8 \text{ eV}$   
 $\rightarrow \text{multi-photon ionization } \geq 2$

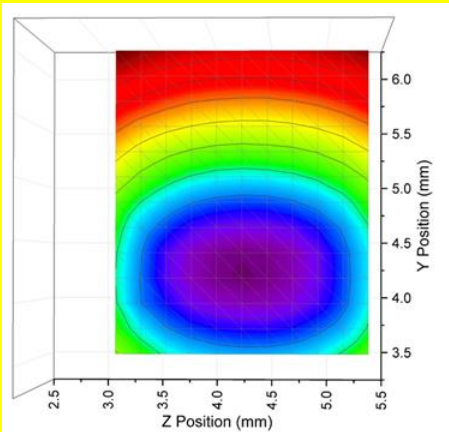
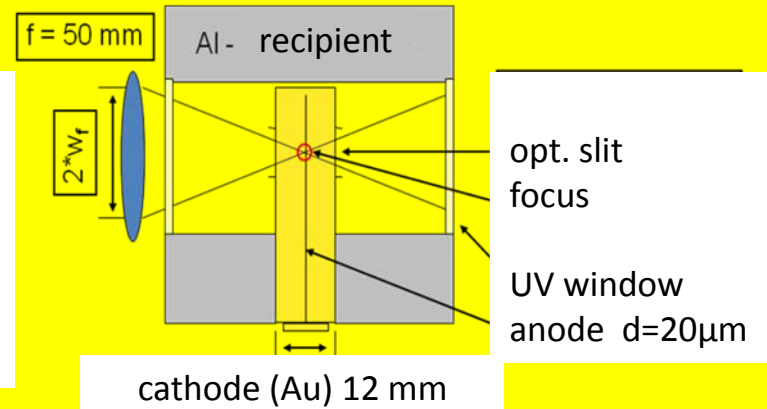
drift tube detector



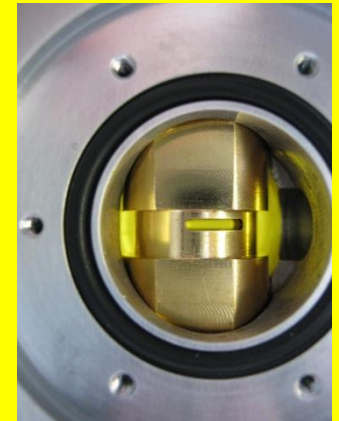
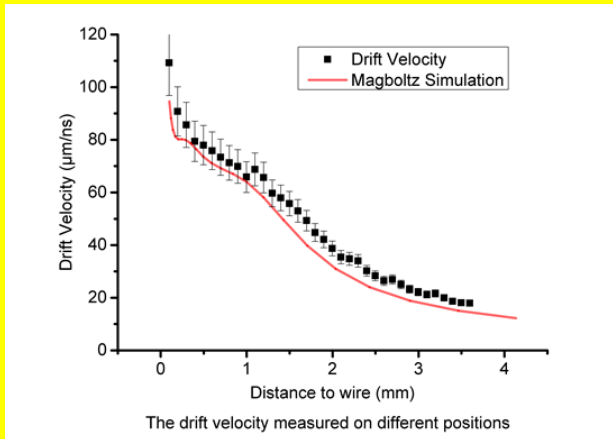
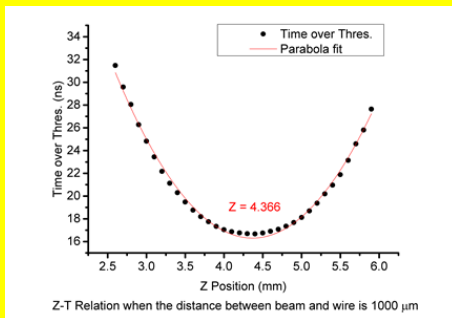
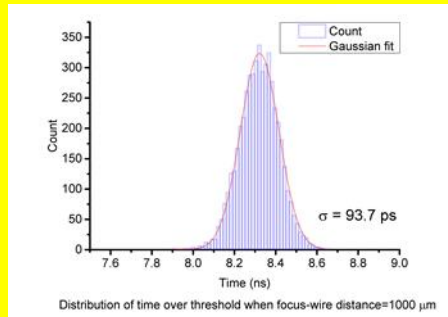
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drift tube detector

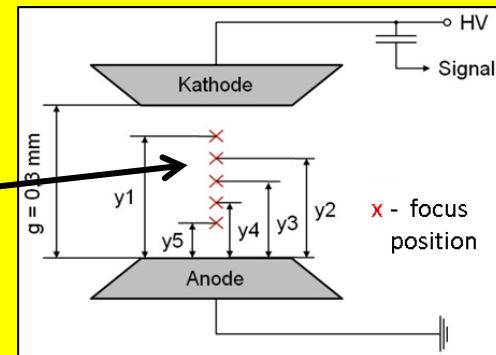
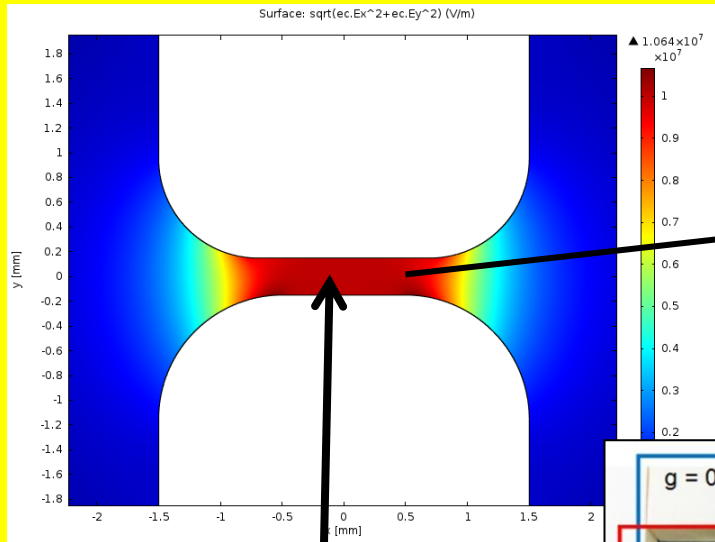


2D scan of time over threshold

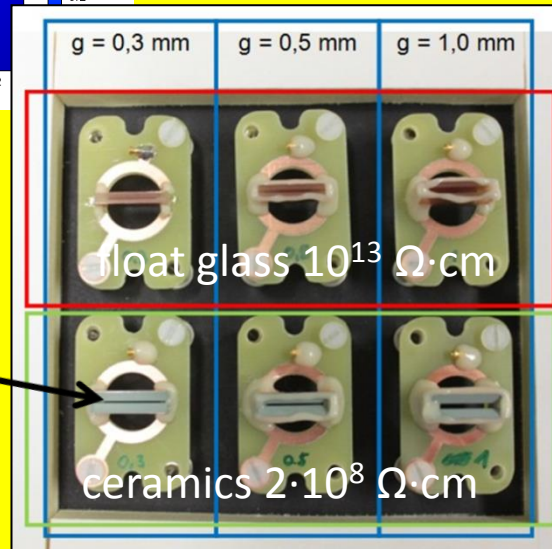




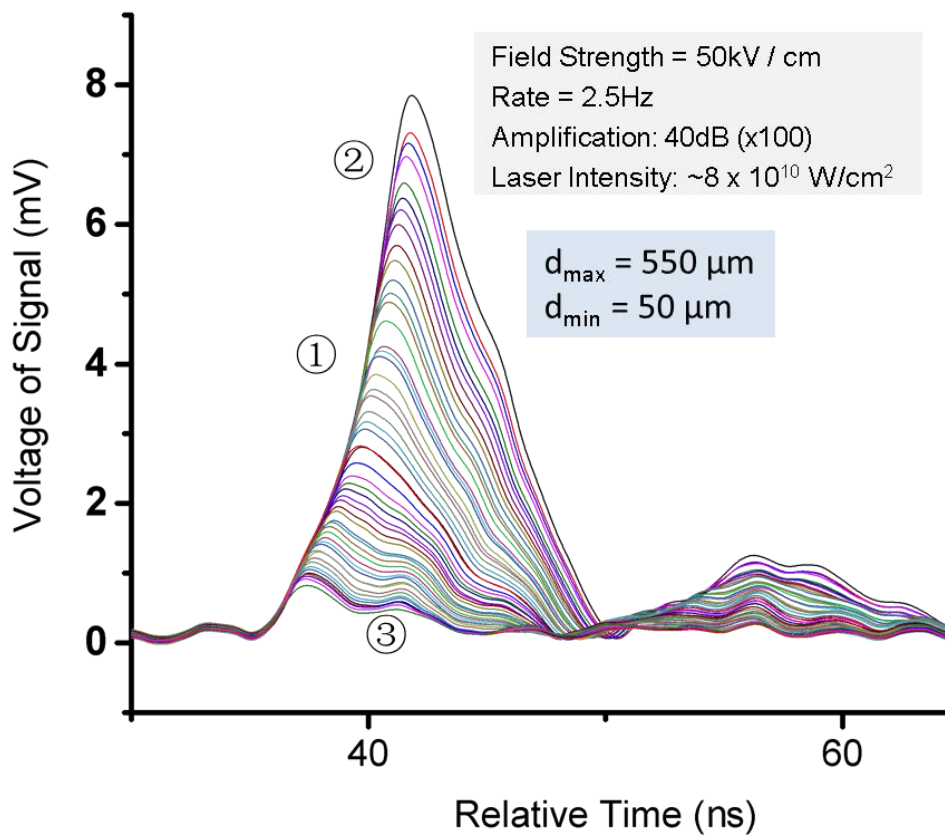
# RPC probe



gap size:  
width 300  $\mu\text{m}$   
length 1 mm  
radius 1 mm

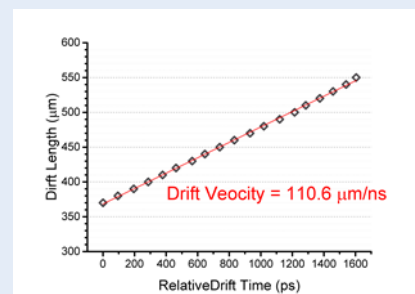


# Trigger-RPC: 50 kV/cm; 1.0 mm

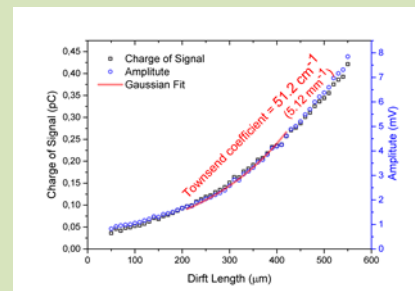


① Rising edge of the signal is overlapped, no matter where the avalanche is started.

② By setting a relative threshold on the waveform, we can measure the drift velocity:

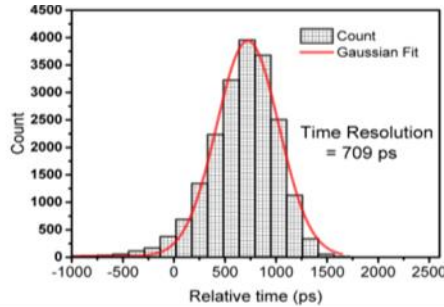


③ The area represents the charge of signal, amplification (eff. Townsend coefficient) can be measured:

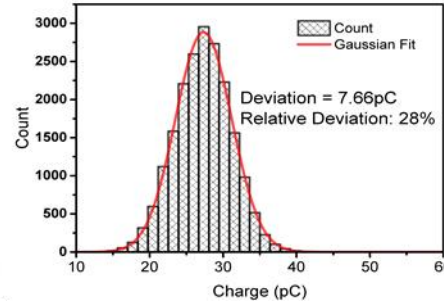


# Trigger-RPC: 50 kV/cm; 1.0 mm

time



charge

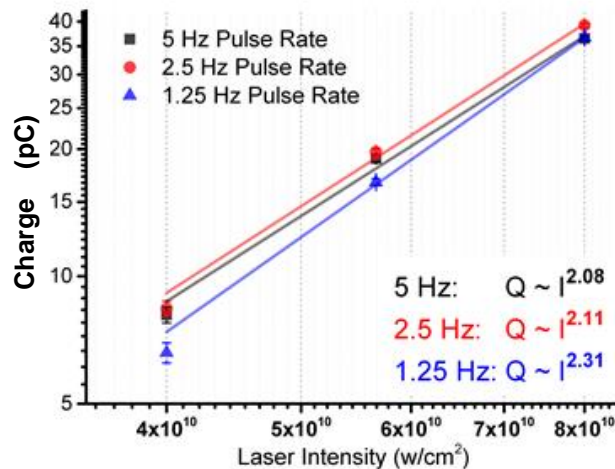


Efficiency: 100%

Time resolution: 700 ps

rel. charge deviation: 28%

charge vs. laser intensity



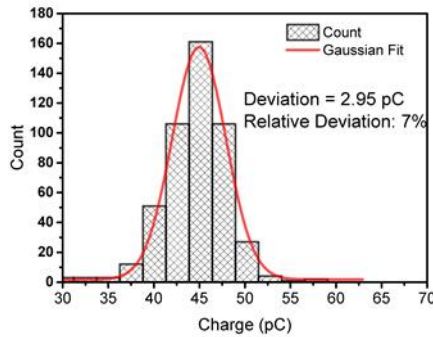
Proportionality of charge(Q), number of primary electrons( $N_e$ ) and intensity( $I^x$ ) :

$$Q \sim N_e \sim I^x$$
$$x \geq 2$$

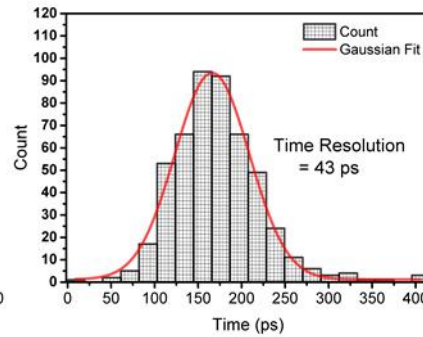
→ **double-photon-ionization**

# Timing-RPC: 100 kV/cm; 0.5 mm

charge



time

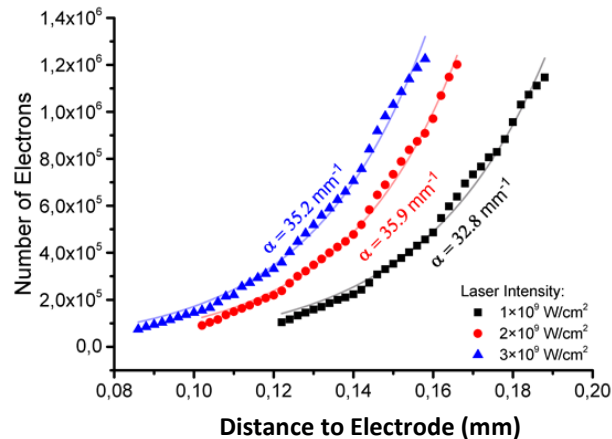


Efficiency: 100%

Time resolution: < 50 ps

rel. charge deviation: 7%

eff. Townsend coefficient



Rate = 0.5Hz

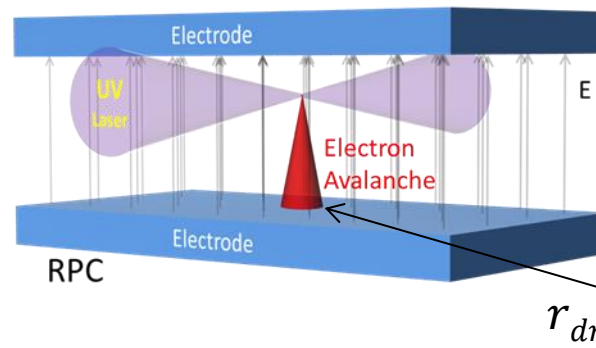
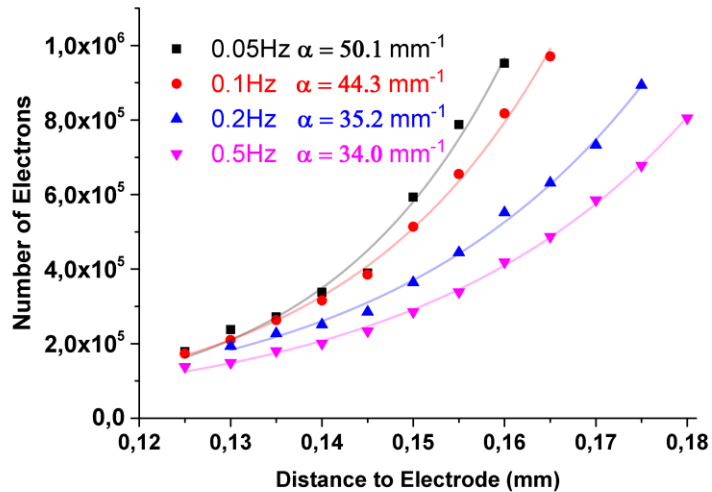
Increment : 2  $\mu\text{m}$

The effective Townsend coefficient is independent on the number of primary electrons ( $N_e$ )

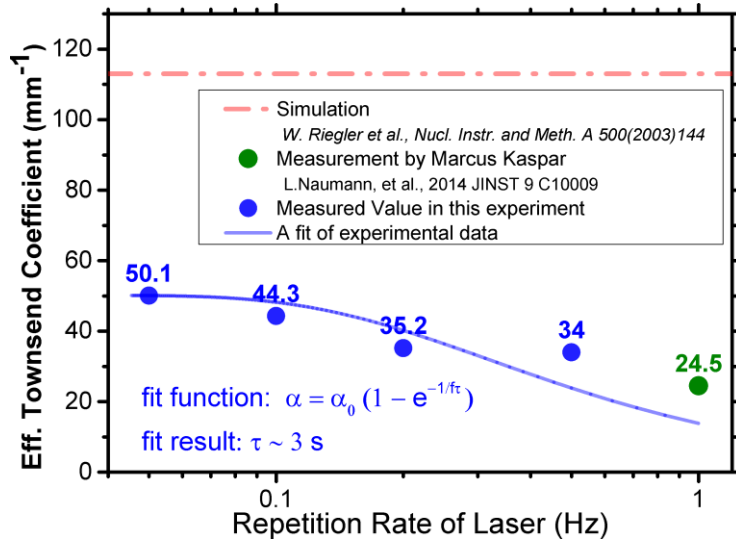
$$\alpha_{eff} = \text{const. for } Q \sim N_e \sim I^2$$

→ **no space-charge effect**

# $\alpha_{\text{eff}}$ Timing-RPC: 100 kV/cm; 0.5 mm

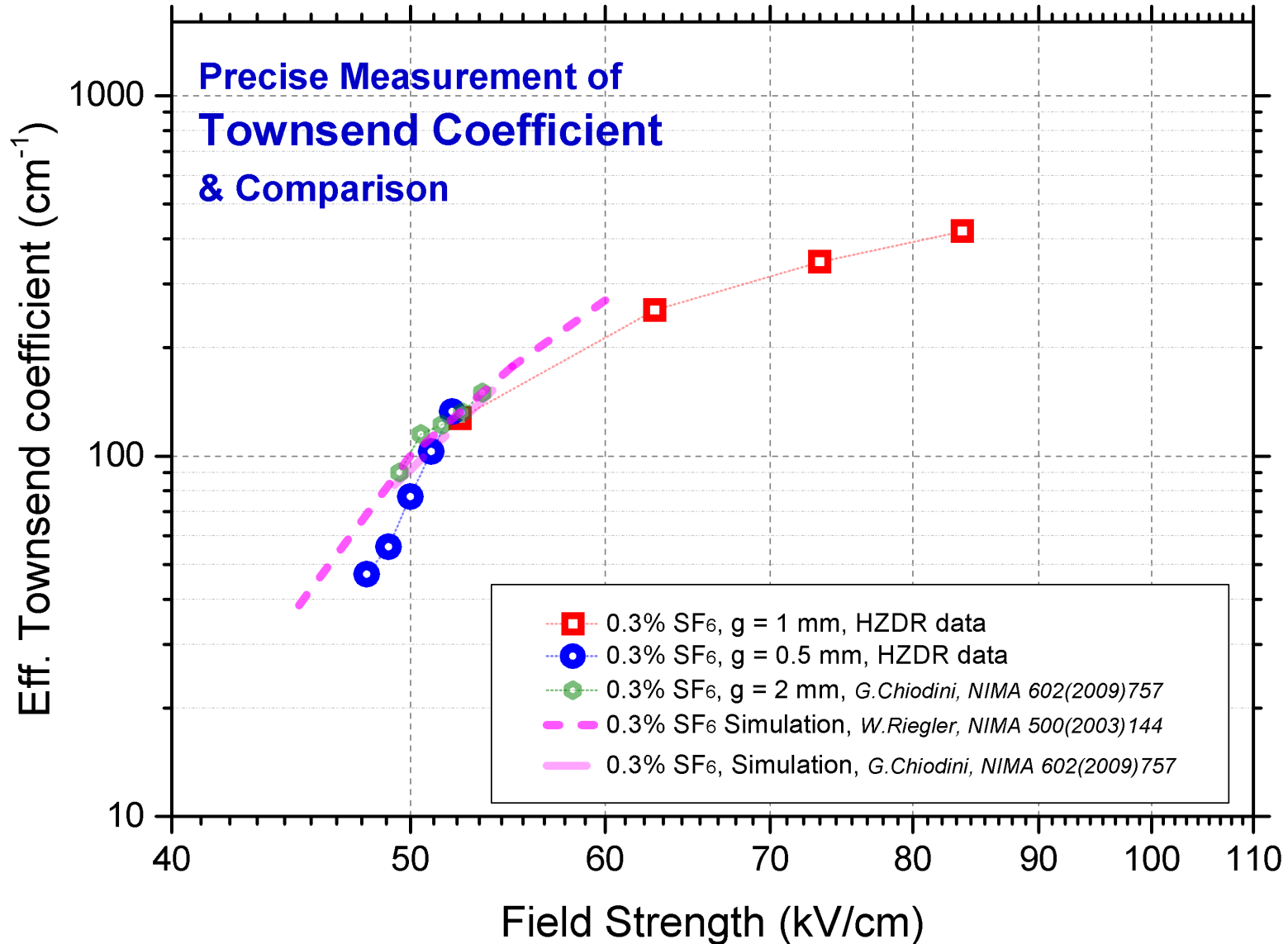


Float glass sample  
 $\tau = \epsilon_0 \epsilon \rho = (3 \pm 0.5) \text{ s}$

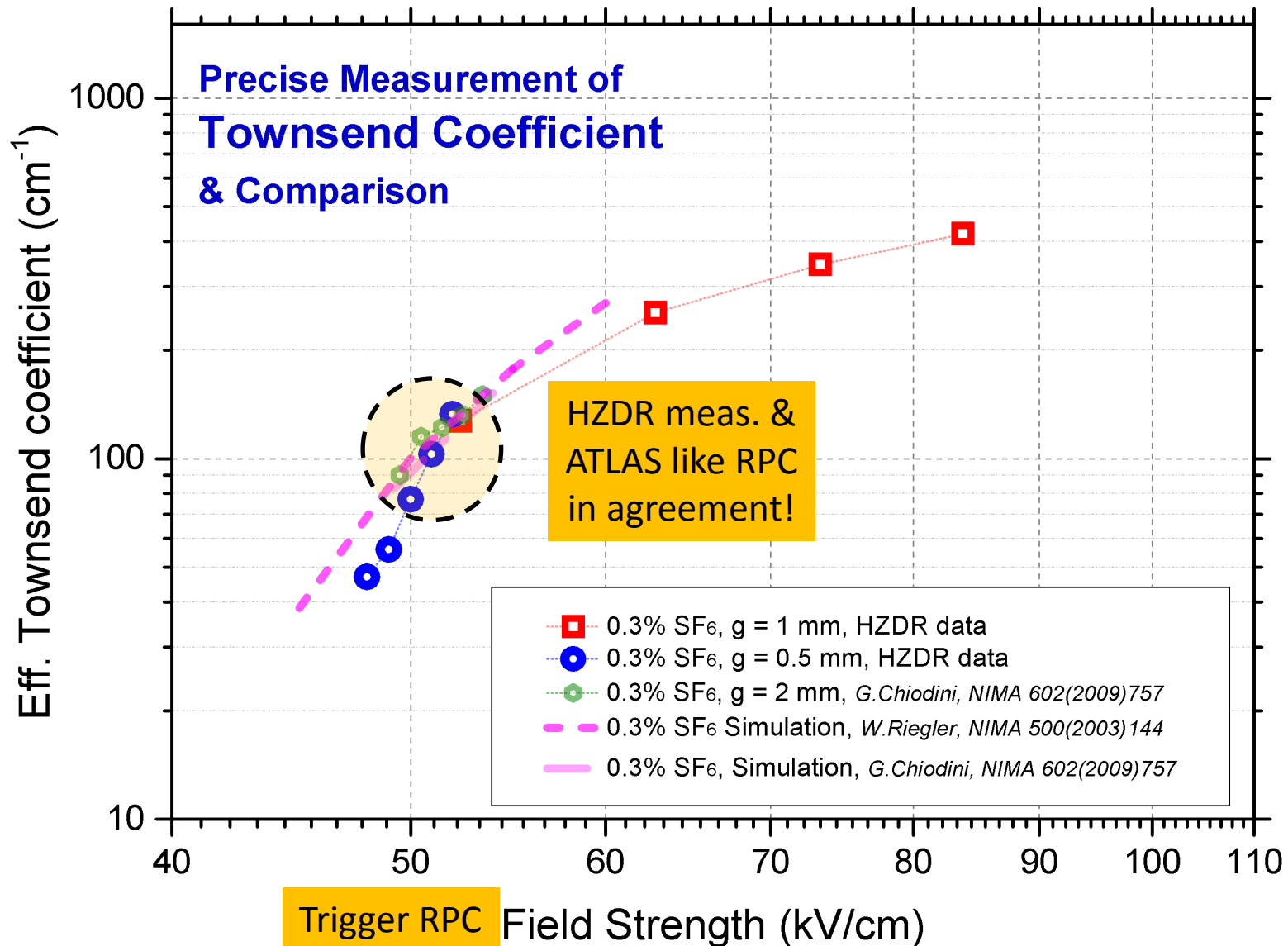


- The eff. Townsend coefficient depends on the laser repetition rate.
- The Time constant of the float glass sample is in agreement with the data.
- Data reach the horizontal asymptote at  $\leq 0.1 \text{ Hz}$ .
- The ionisation occurs always at the same micro-volume and the charges are accumulated on the same area of the electrode surfaces  $\rightarrow$   $0.1 \text{ Hz}/(\text{aval. area})$  is comparable to  $\geq 1 \text{ kHz}/\text{cm}^2$

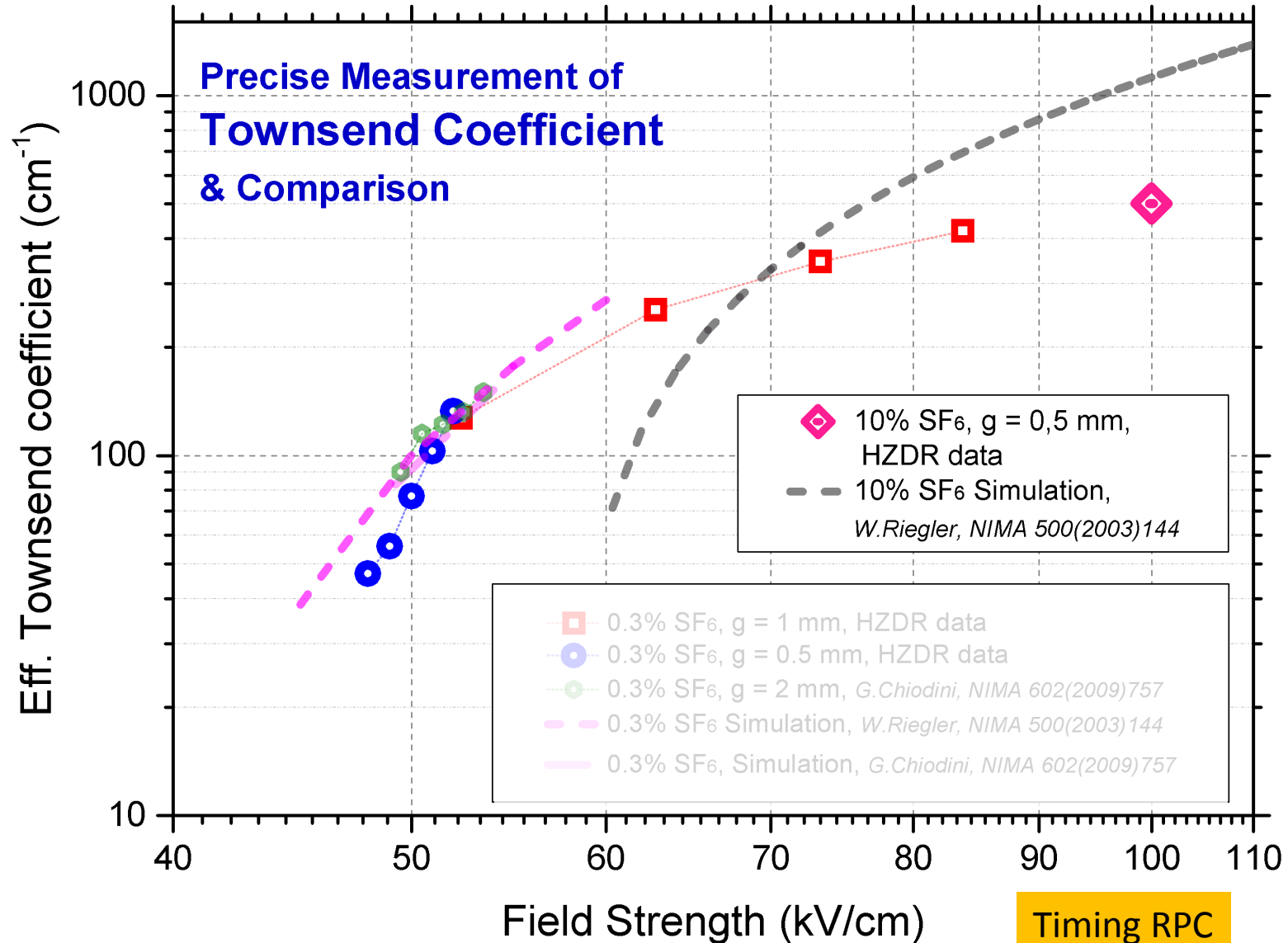
# Effective Townsend coefficient (1)



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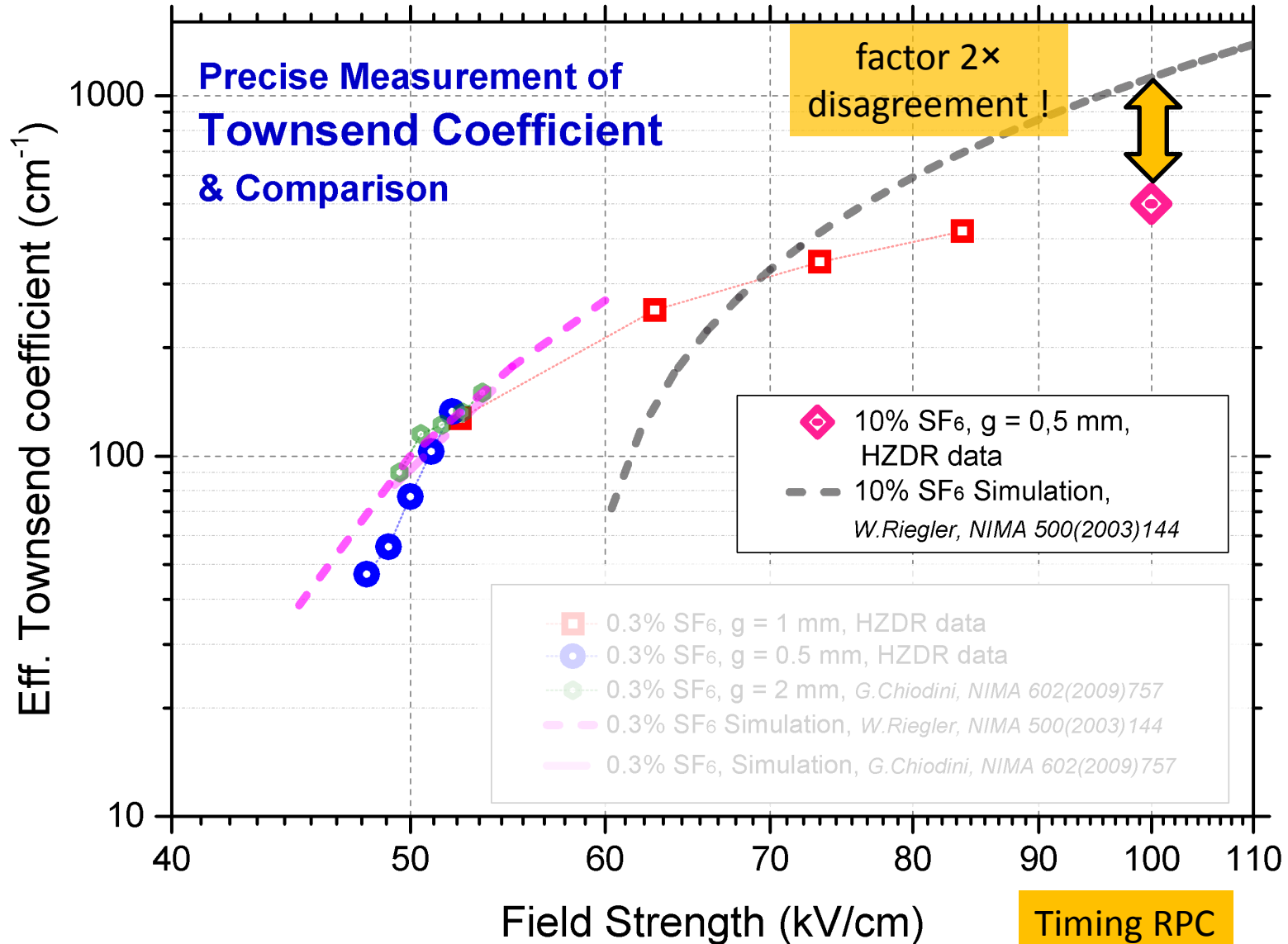


# Effective Townsend coefficient (2)

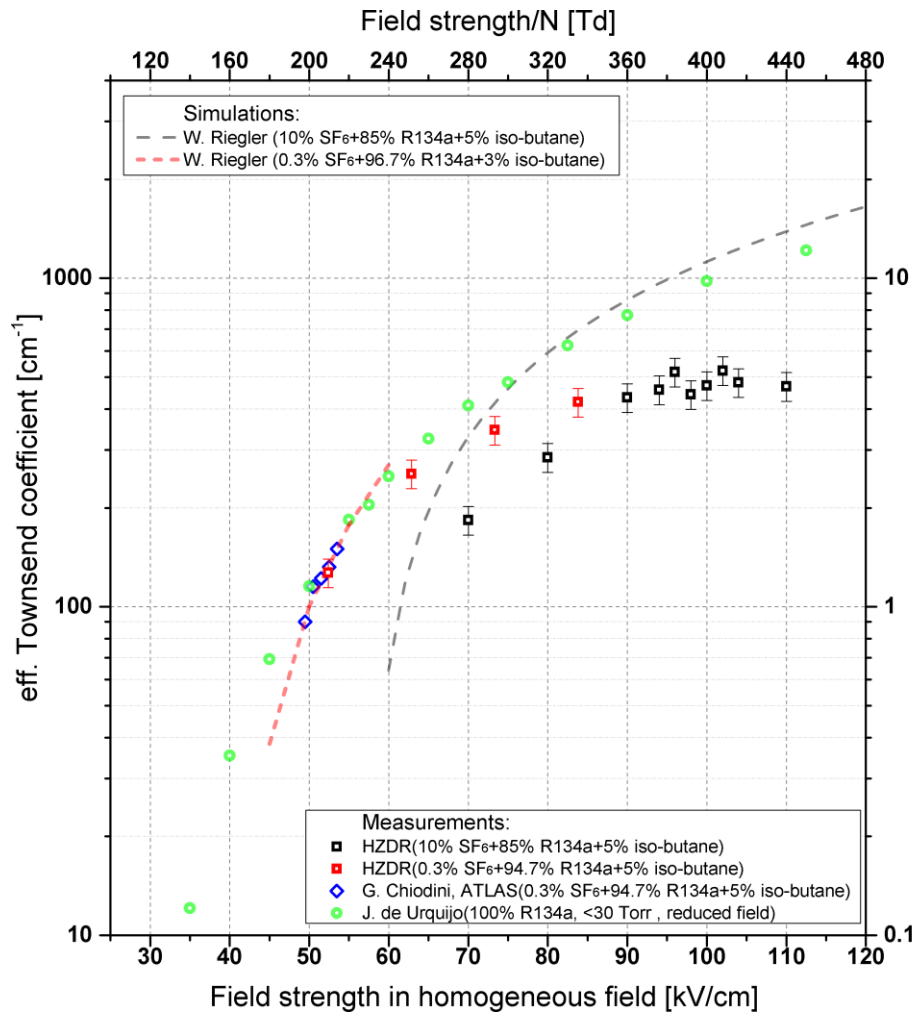




# Effective Townsend coefficient (2)

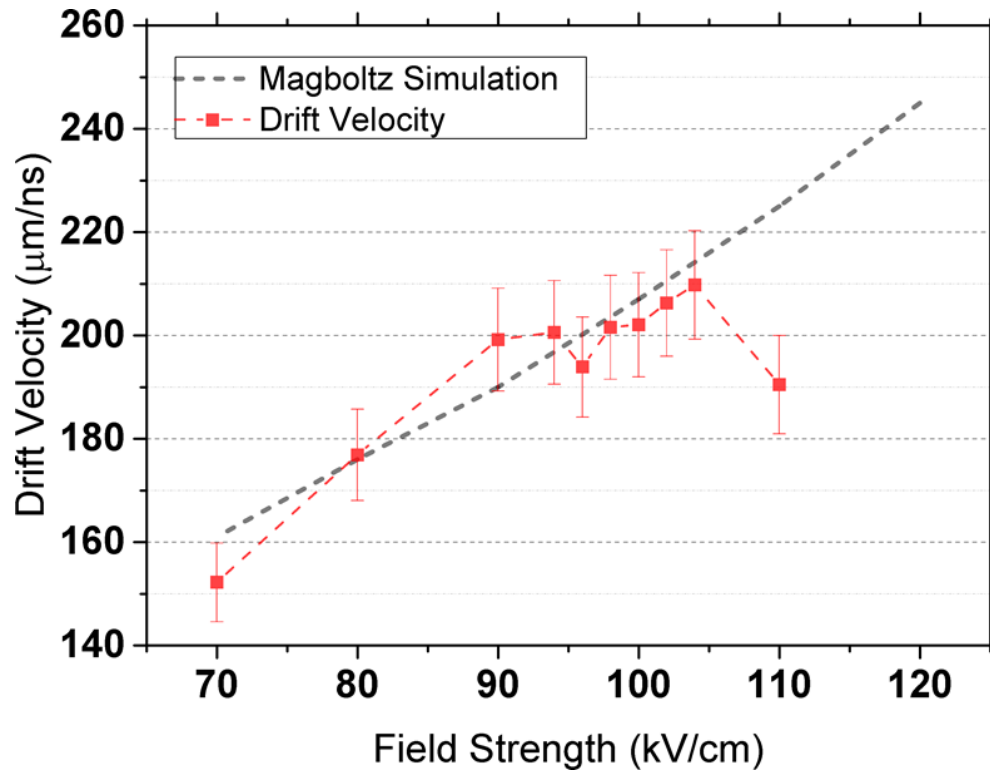


# Effective Townsend coefficient (3)



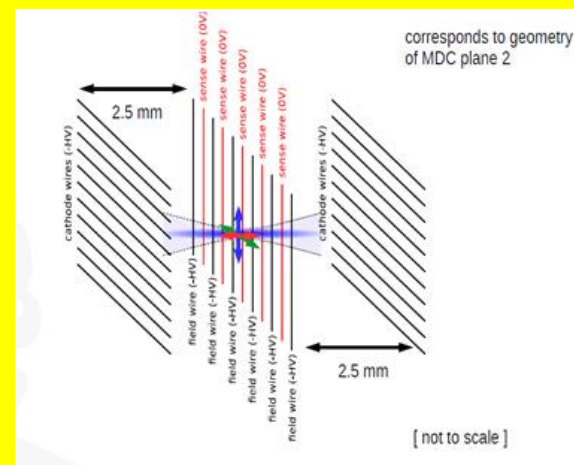
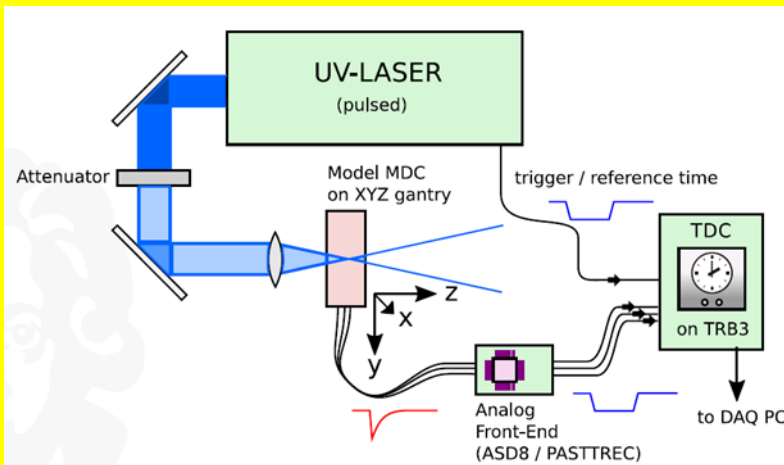
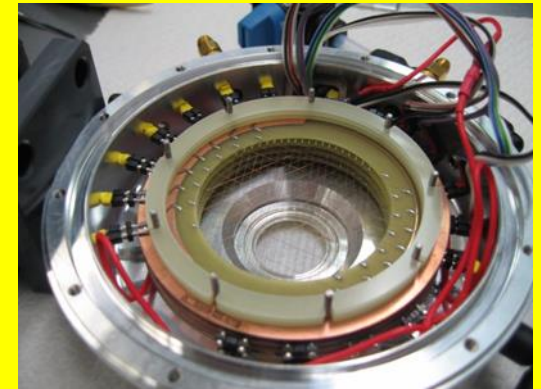
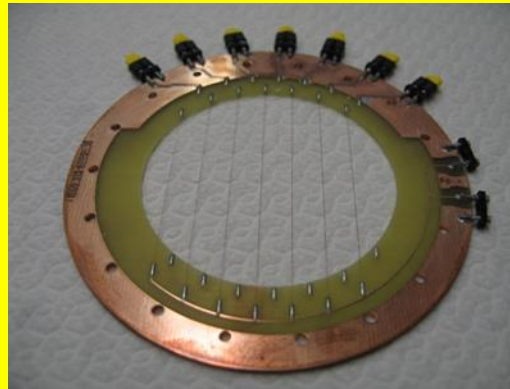
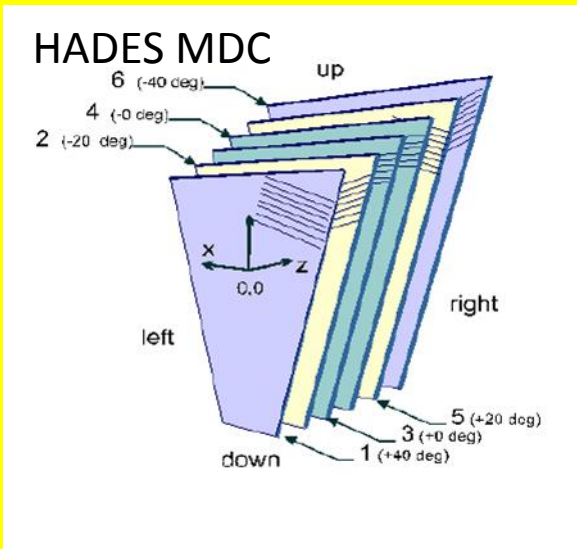
- reduced field data are in agreement with Magboltz simulation
- HZDR data shows a saturation above 90 kV/cm

# Electron drift velocity



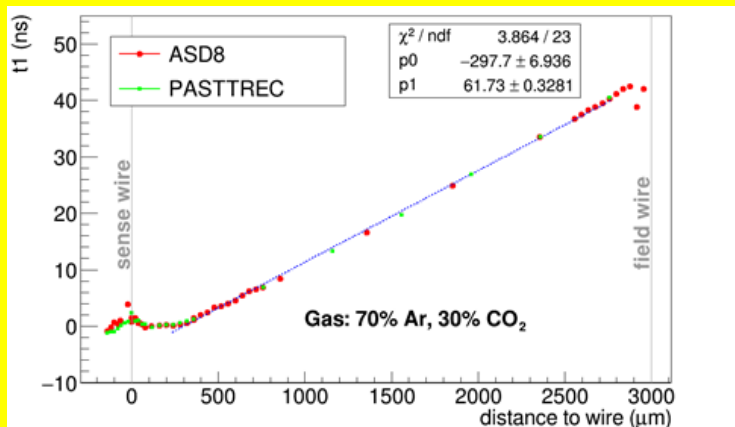
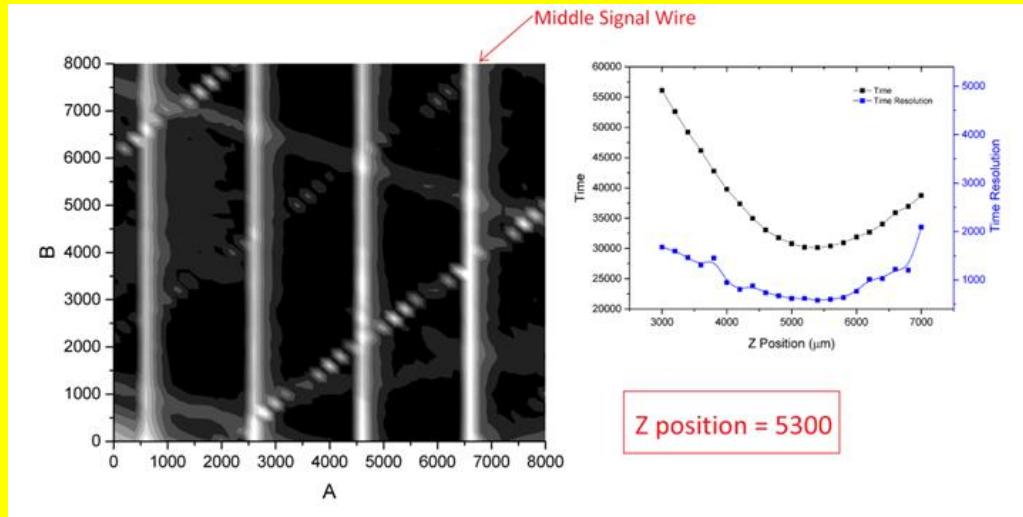
■ data are in agreement with Magboltz simulation

# Mini Drift Cell (HADES-like)

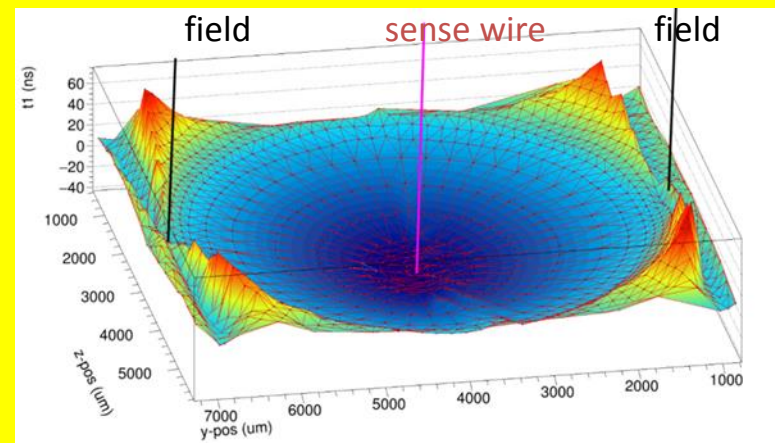


# Mini Drift Cell

## MDC tomography



Electron drift velocity  $v_e = 62 \mu\text{m}/\text{ns}$   
 Spatial resolution  $\leq 60 \mu\text{m}$



Drift time distribution: 1290 data points  
 with 8000 ev. /point

# Summary

The UV-Laser driven test facility for gaseous detectors at HZDR is works very stable in an automatically regime to provide detector tests with micro-positioning of the generated micro-plasma

For HADES-like Mini Drift Cells operating in inhomogeneous electric fields has been obtained:

- a deeper understanding of the field topology
- spatial resolution better than 60  $\mu\text{m}$

# Summary

For RPC operating in strong and homogeneous electric fields at atmospheric pressure has been shown:

- **Agreement** of the eff. Townsend coefficient for Freon(94.7%)+IB(5%)+SF<sub>6</sub>(0.3%) **at 50 kV/cm** for ATLAS-like and HZDR RPC prototype measurements and MAGBOLTZ simulation.
- **Disagreement (factor 2)** of the eff. Townsend coefficient measurement for Freon(85%)+IB(5%)+SF<sub>6</sub>(10%) **at 100 kV/cm** and the MAGBOLTZ simulation.
- **Agreement** of the drift velocity with simulation.
- **Agreement** of RPC parameter measurements of rate capability, time and energy resolution with model predictions.

# Next Tests

- Investigation of RPC rate capability test with low resistive RPC
- Investigation of RPC double hit behavior
- Evaluation of environmentally friendly gas mixtures for RPC application
- Investigation of the Townsend puzzle



## Acknowledgement

HZDR: **X. Fan**, B. Kämpfer, M. Siebold, M. Sobiella, D. Stach  
GSI: C. Wendisch, **M. Wiebusch**



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