

Supergravity phenomenology : minimal VS no-scale

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JRJC 09

Outline

- 1 A world of symmetries
- 2 Why Supersymmetry ?
- 3 Supergravity Basics
- 4 Where is the phenomenology?
- 5 Msugra VS no-scale

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It's (almost) all about symmetry

Building Standard model Lagrangian

- Symmetries (Space-time, interactions)
- Renormalizability

$$\begin{cases} \phi \rightarrow U\phi \\ \delta S = 0 \end{cases}$$

Doesn't change the physics

It's (almost) all about symmetry

Building Standard model Lagrangian

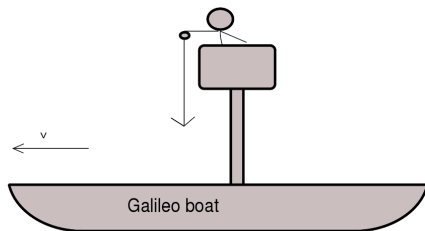
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Doesn't change the physics

Space-time symmetries

It started with a boat



Physics is the same in a all inertial frames

Poincaré group

- **Lorentz group** : Antisymmetric operator $M_{\mu\nu}$
 - ▶ Rotation in space-coordinates
 - ▶ Lorentz boost : rotation in space and time coordinates
- **Translation**

Global and local symmetry

Example of gauge interactions in Standard Model (QED)

Global phase transformation

$$\begin{cases} \phi \rightarrow e^{i\theta} \phi \\ \bar{\phi} \rightarrow e^{-i\theta} \bar{\phi} \end{cases} \Rightarrow \delta(\bar{\phi} \partial_{\mu} \phi) = 0$$

Covariant derivative

$$\begin{aligned} \partial_{\mu} &\rightarrow D_{\mu} = \partial_{\mu} + iA_{\mu} \\ \begin{cases} \phi \rightarrow e^{i\theta(x)} \phi \\ D_{\mu} \bar{\phi} \rightarrow e^{-i\theta(x)} D_{\mu} \bar{\phi} \end{cases} \end{aligned}$$

Restore local symmetry \Rightarrow introduce new field and interaction

Global and local symmetry

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Local phase transformation

$$\begin{cases} \phi \rightarrow e^{i\theta(x)}\phi \\ \bar{\phi} \rightarrow e^{-i\theta(x)}\bar{\phi} \end{cases} \Rightarrow \delta(\bar{\phi}\partial_\mu\phi) \sim i\partial_\mu\theta(x)$$

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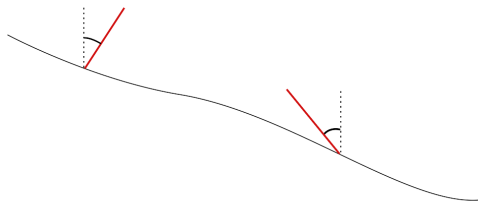
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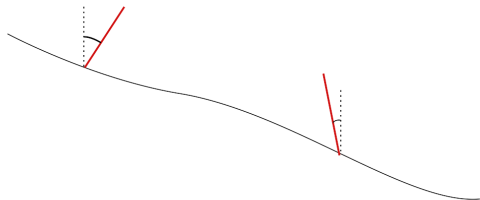
What really happens ?

Global : same phases everywhere



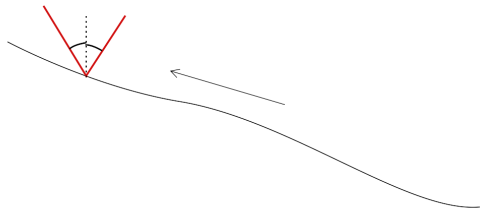
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Local : not the same everywhere



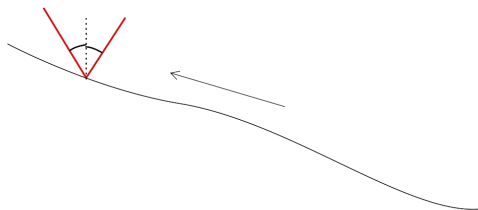
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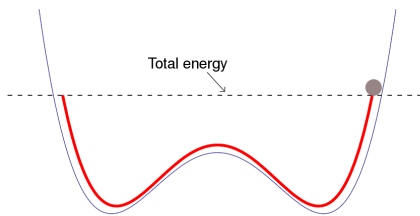
Local : not the same everywhere



Transport to the same place

Symmetries may be broken : $SU(2)_L \times U(1)_Y \rightarrow U(1)_{EM}$

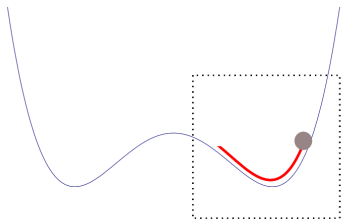
Higgs potential : Before



$\phi\phi WW$

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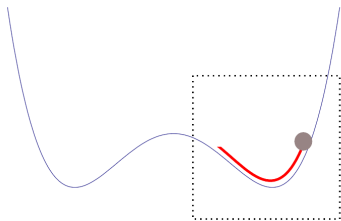
Higgs potential : After



$$v^2 WW + 2vhWW + hhWW$$

Symmetries may be broken: $SU(2)_L \times U(1)_Y \rightarrow U(1)_{EM}$

Higgs potential : After



$$v^2 WW + 2vhWW + hhWW$$

It's a bird, it's a plane, no it's a mass term

Important ideas

- Symmetries are everywhere
- Global \rightarrow Local : new field
- Scalar field at *low energy* :
 - ▶ Field \rightarrow vev + small fluctuation
 - ▶ Effective interactions : mass terms appears

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The question

Looking for operator Q such that

$$Q \phi(j) = \phi'(j' \neq j)$$

Changing the spin $\Rightarrow [Q_i, \text{rotation operator}] \neq 0$

Only supersymmetry :
anti-commuting algebra

$$\begin{cases} Q_i Q_j + Q_j Q_i = 0 \\ Q|j\rangle = |j \pm \frac{1}{2}\rangle \end{cases}$$

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Symmetry between bosons fermions

- each SM fermion : scalar partner
- each gauge boson : fermion partner
- And the Higgs ? scalar : fermion partner

partners have same masses

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partners have **same masses**

The Minimal Supersymmetric Standard Model

Minimal extension of standard model

- SM \rightarrow super-partner
 - ▶ fermions \rightarrow sfermions
 - ▶ gauge bosons \rightarrow gauginos
- Higgs $\times 2$
 - ▶ $H \rightarrow H_u, H_d$
- additional terms for masses and interactions (broken susy) **lifting degeneracy**
(For instance : $m_{sfermion} = m_{fermion} + \Delta m$)

How to generate such terms?

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Supergravity

Symmetry \rightarrow doesn't change the physics

$\epsilon \rightarrow \epsilon(x^\mu)$ space-time dependant

$$\delta S_0 \sim \int d^4x \partial_\mu \epsilon^a K_a^\mu$$

$$S = S_0 + S_G$$

$$\delta S_0 + \delta S_G = 0$$

Restore symmetry, introduce $S_G \sim \int d^4x K_a^\mu \Psi_\mu^a$

$$\Psi_\mu^a \rightarrow \Psi_\mu^a - \partial_\mu \epsilon^a$$

Introducing **gravitino** field (spin $\frac{3}{2}$)

Gravity multiplet

Completing supersymmetry :

- Introduce gravitino to restore local supersymmetry
- add kinetic term for gravitino
- add partner of gravitino

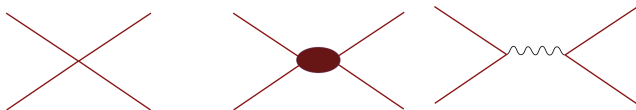
spin 2 graviton

Renormalizability

Non-renormalizable operators

$$L_{NR} \supset \frac{O_{grav}^{d+n}}{M_{Planck}^n}$$

Like Fermi weak theory



Effective gravity theory

$$L_{NR} \rightarrow 0 \text{ when } M_{Planck} \rightarrow \infty$$

Effective theory : link between high and low scales

- Lagrangian we construct start just below M_{Planck}
- Gravitational quantum effects are already integrated
- Real quantum gravity theory beyond M_{Planck} (string or ...)
- Next stop : GUT scale (grand unification)

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Scalar Potential

Phenomenology is encoded in :

$$V \sim e^G \left(G^i (G^{-1})_{\bar{i}} \bar{G}_{\bar{j}} - 3 \right) + \frac{1}{2} f_{ab} D_a D^a$$

- $W(\phi)$: Superpotential, analytic in ϕ , encode scalar masses and yukawas interactions
- $K(\phi^i, \bar{\phi}^{\bar{j}})$: Kähler function, appears in kinetic terms $K_{i\bar{j}} \partial_\mu \phi^i \partial^\mu \bar{\phi}^{\bar{j}}$ with $K_{i\bar{j}} = \partial_i \partial_{\bar{j}} K$
- $G(\phi, \bar{\phi}) = K(\phi, \bar{\phi}) + \ln |W(\phi)|^2$
- f_{ab} : gauge kinetic function

scalar masses shift (partner of fermions and higgs)

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gauginos masses shift (partner of gauge bosons)

Breaking and hidden sector

- Breaking $\rightarrow m_{\frac{3}{2}} \neq 0$
- need scalar vev : hidden field $\langle z \rangle$

Gravitino mass \rightarrow local supersymmetry breaking

Hidden field vev $\rightarrow m_{\frac{3}{2}}$

Hidden sector decouples at *low energy*

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Masses in minimal supergravity

$$\begin{cases} K = \phi^i \bar{\phi}^{\bar{j}} \\ f_{ab} = \delta_{ab} \end{cases}$$

$$M_{Planck} \rightarrow \infty$$

Effective potential at lower scale (usually **unification** scale)

$$V(\langle z \rangle, \phi^i) \sim \sum |\partial_i W|^2 + m_{\frac{3}{2}}^2 \sum |\phi^i|^2 + m_{\frac{3}{2}} \left(\sum \phi^i \partial_i W + (A - 3)W + h.c. \right)$$

soft mass term lifting degeneracy between fermion and scalar

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mSugra issues

Light version

- $V = 0$ put by hand
- Fine-tuning the cosmological constant
- necessary to identify $m_{\frac{3}{2}}$ as a gravitino mass

no-scale supergravity

Light version 2

- Non-minimal Kähler : $K = -3 \log(z + z^*)$
- Naturally $V = 0$
- Undetermined vacuum expectation value
- Flat potential at GUT scale
- no scalar masses, only gauginos

This kind of Kähler potential appears in some **string theories** after compactification

Renormalization group equations

High scale parameters (usually **unification scale** $\sim 10^{16}$ GeV)



RGE



Low scale parameters (**EW scale**)

Ordinary differential equations
Evolution of parameters with energy scale

Boundary conditions

mSugra, no-scale, or any other give us parameters at Grand Unification scale

Universality condition

- 1 m_0 : one mass for all scalar at GUT scale
- 2 $m_{\frac{1}{2}}$: one mass for all gauginos at GUT scale

ElectroWeak scale

- Spectrum at low scale
- Correct parameters to induce ElectroWeak Symmetry Breaking (EWSB)
- Impose all constraints

$$V = m_1^2 |H_1|^2 + m_2^2 |H_2|^2 + B\mu (H_1 H_2 + H_1^* H_2^2) + \frac{g^2 + g'^2}{8} (|H_1|^2 - |H_2|^2)^2$$

Parameters driven from high energy to low energy

SuSy breaking \Rightarrow EW breaking

My work

How can no-scale models survive ?

Main phenomenological issue

$m_0 = 0 \rightarrow$ charged LSP

(Many others, but not enough space)

Hard to find parameters such that good dark matter candidate

Outlook

Thunderbird

- Trying to enlarge Poincaré \rightarrow supersymmetry
- Local supersymmetry \rightarrow Supergravity
- Natural frame for Standard Model, EWSB, dark matter...
- Still don't know the correct breaking mechanism

Supergravity, what else?