

IFSC UNIVERSIDADE
DE SÃO PAULO
Instituto de Física de São Carlos



Testing fundamental physics with Astrophysical Sources: LIV

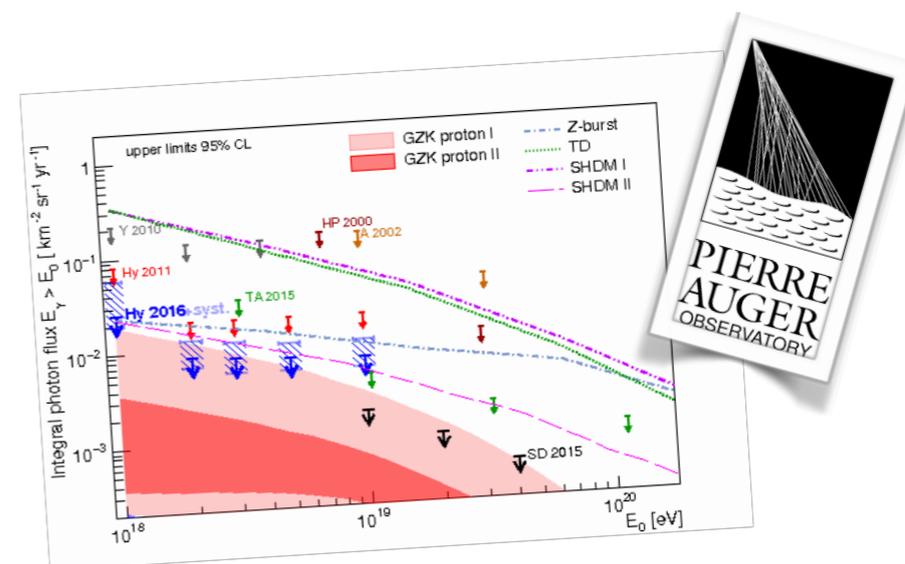
Humberto Martínez-Huerta,
IFSC-USP, Brazil

**Meeting of the Cosmic Rays Section of the Mexican
Physical Society
3-5 Oct 2018**

Index

- I. Lorentz invariance violation (LIV)
- II. Limits: TeV γ -rays
 - i. Time Energy Dependent delay
 - ii. Photon Decay
 - iii. Pair production threshold shifts
- III. UHECR

- i. GZK-photons + LIV
- ii. Limits: UHECR



I. Lorentz invariance violation (LIV)

II. Limits: TeV γ -rays

i. Time Energy Dependent delay

ii. Photon Decay

iii. Pair production threshold shifts

III. UHECR

i. GZK-photons + LIV

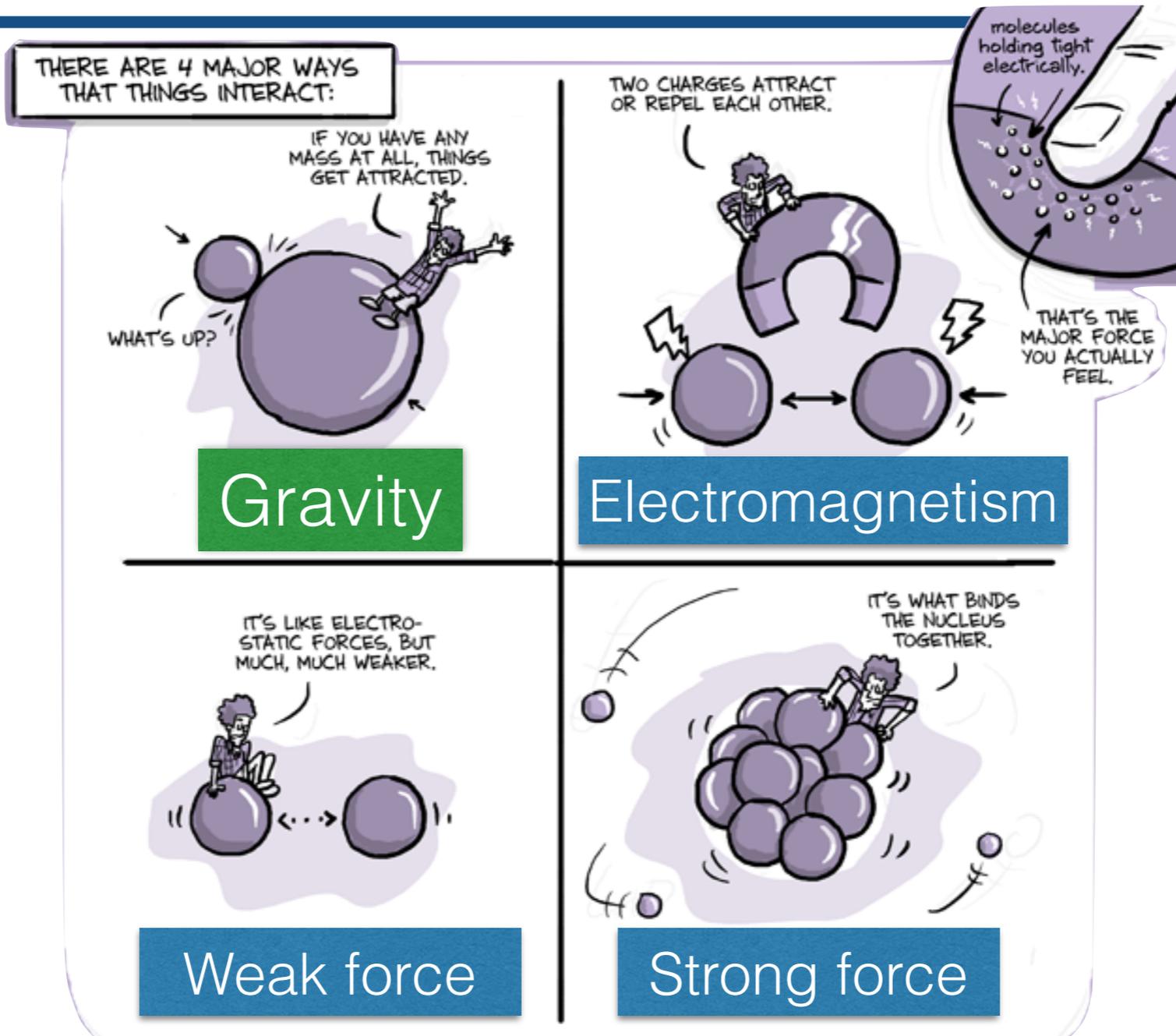
ii. Limits

Fundamental Forces of Nature

General Relativity



Geometrical Theory

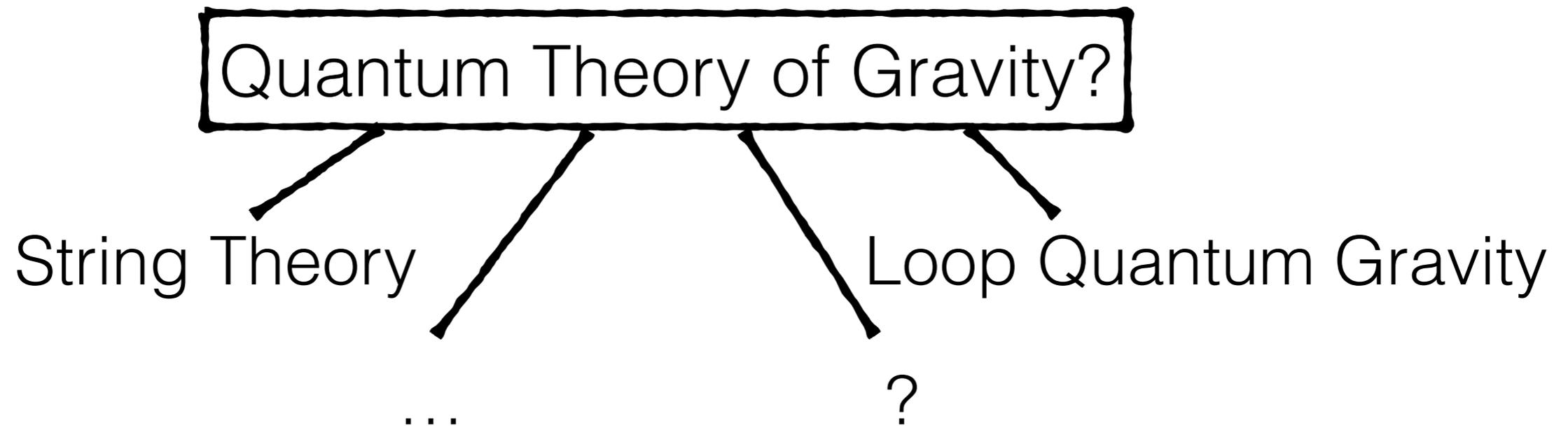


Standard Model



Quantum Field Theory

- SM & GR: the best theories describing the 4-fundamental Forces.
- No conflict with predictions from either of them.
- **They are fundamentally different.**



New Physics involves new features, such as:

- Higher Dimensions of s-t
- Brane World scenarios
- Noncomutative geometries
- ...
- The law of relativity might not hold exactly at all energy scales → Lorentz Invariance Violation (LIV)

... **LI may not be an exact symmetry of Nature**

Generic LIV dispersion relation

$$E^2 - p^2 \pm \epsilon A^2 = m^2,$$

$$E \gg m,$$

$$A = \{E, p\}$$

$$\epsilon \rightarrow \epsilon(A)$$

A general modification to the dispersion relation would rather involve a general function of energy and momentum

$$\epsilon(A)A^2 = \epsilon(0)A^2 + \epsilon'(0)A^{(2+1)} + \frac{\epsilon''(0)}{2!}A^{(2+2)} + \frac{\epsilon'''(0)}{3!}A^{(2+3)} + \dots$$

The dispersion relation:

$$E^2 - p^2 \pm \delta_n A^{n+2} = m^2,$$

$$\delta_n \stackrel{n \geq 1}{=} \epsilon^{(n)} / M^n = 1 / (E_{LIV}^{(n)})^n$$

it is not necessarily bound to a particular LIV-model, which allows to generalize to some point the search of LIV-signatures.

LIV negligible at the lower standard energies

Generic LIV dispersion relation

$$E^2 - p^2 \pm \epsilon A^2 = m^2,$$

$$E \gg m,$$

$$A = \{E, p\}$$

$$\epsilon \rightarrow \epsilon(A)$$

A general modification to the dispersion relation would rather involve a general function of energy and momentum

$$\epsilon(A)A^2 = \epsilon(0)A^2 + \underbrace{\epsilon'(0)A^{(2+1)}}_{n=1} + \frac{\epsilon''(0)}{2!}A^{(2+2)} + \frac{\epsilon'''(0)}{3!}A^{(2+3)} + \dots$$

The dispersion relation:

$$E^2 - p^2 \pm \delta_n A^{n+2} = m^2, \quad \delta_n \stackrel{n \geq 1}{=} \epsilon^{(n)} / M^n = 1 / (E_{LIV}^{(n)})^n$$

it is not necessarily bound to a particular LIV-model, which allows to generalize to some point the search of LIV-signatures.

LIV negligible at the lower standard energies

Generic LIV dispersion relation

$$E^2 - p^2 \pm \epsilon A^2 = m^2,$$

$$E \gg m,$$

$$A = \{E, p\}$$

$$\epsilon \rightarrow \epsilon(A)$$

A general modification to the dispersion relation would rather involve a general function of energy and momentum

$$\epsilon(A)A^2 = \epsilon(0)A^2 + \epsilon'(0)A^{(2+1)} + \frac{\epsilon''(0)}{2!}A^{(2+2)} + \frac{\epsilon'''(0)}{3!}A^{(2+3)} + \dots$$

$n=2$

The dispersion relation:

$$E^2 - p^2 \pm \delta_n A^{n+2} = m^2, \quad \delta_n \stackrel{n \geq 1}{=} \epsilon^{(n)} / M^n = 1 / (E_{LIV}^{(n)})^n$$

it is not necessarily bound to a particular LIV-model, which allows to generalize to some point the search of LIV-signatures.

LIV negligible at the lower standard energies

Generic LIV dispersion relation

$$E^2 - p^2 \pm \epsilon A^2 = m^2,$$

$$E \gg m,$$

$$A = \{E, p\}$$

$$\epsilon \rightarrow \epsilon(A)$$

A general modification to the dispersion relation would rather involve a general function of energy and momentum

$$\epsilon(A)A^2 = \epsilon(0)A^2 + \epsilon'(0)A^{(2+1)} + \frac{\epsilon''(0)}{2!}A^{(2+2)} + \frac{\epsilon'''(0)}{3!}A^{(2+3)} + \dots$$

The dispersion relation:

$$E^2 - p^2 \pm \delta_n A^{n+2} = m^2, \quad \delta_n \stackrel{n \geq 1}{=} \epsilon^{(n)} / M^n = 1 / (E_{LIV}^{(n)})^n$$

it is not necessarily bound to a particular LIV-model, which allows to generalize to some point the search of LIV-signatures.

LIV negligible at the lower standard energies

I. Lorentz invariance violation (LIV)

II. **Limits: TeV γ -rays**

i. Time Energy Dependent delay

ii. Photon Decay

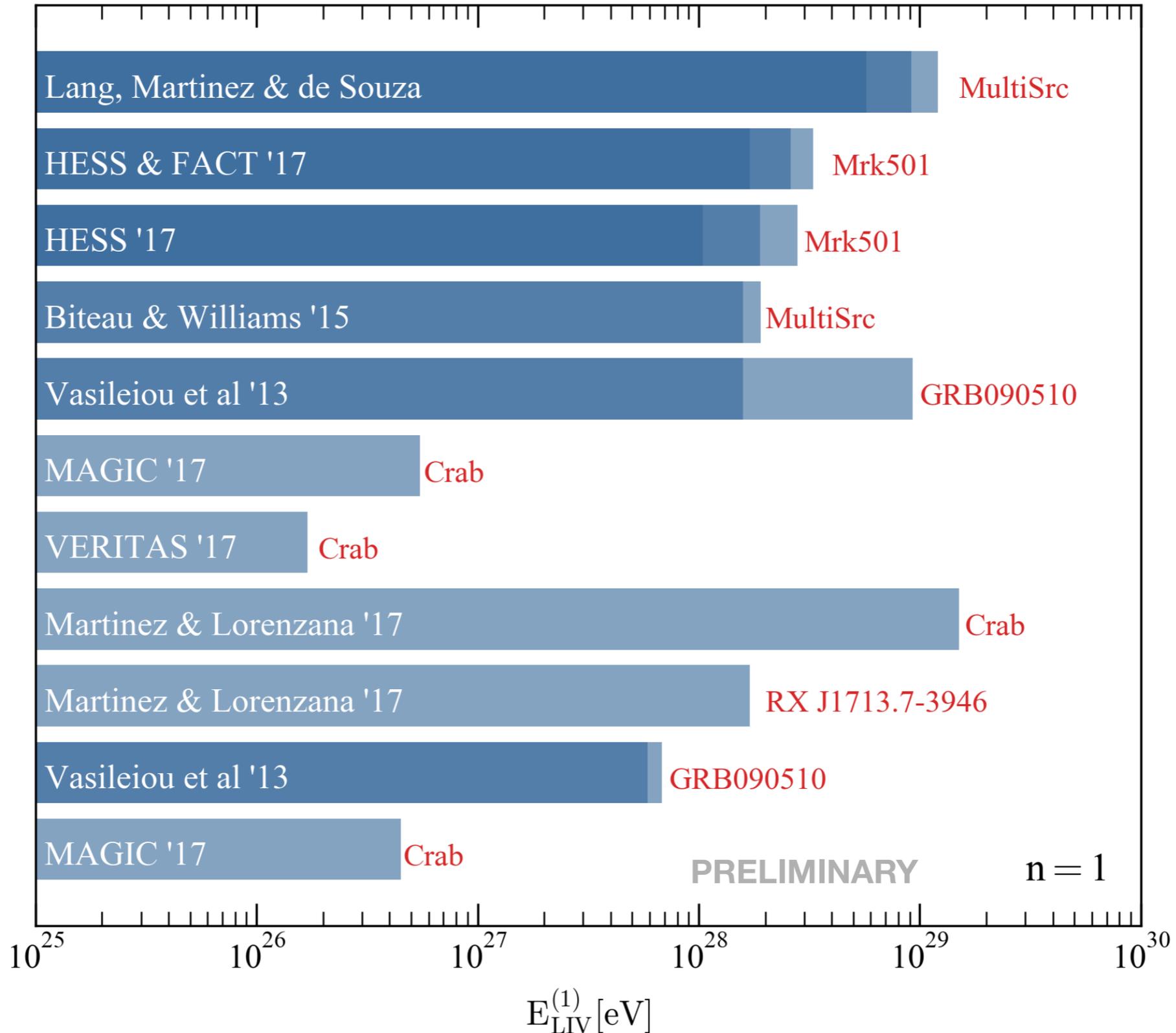
iii. Pair production threshold shifts

III. UHECR

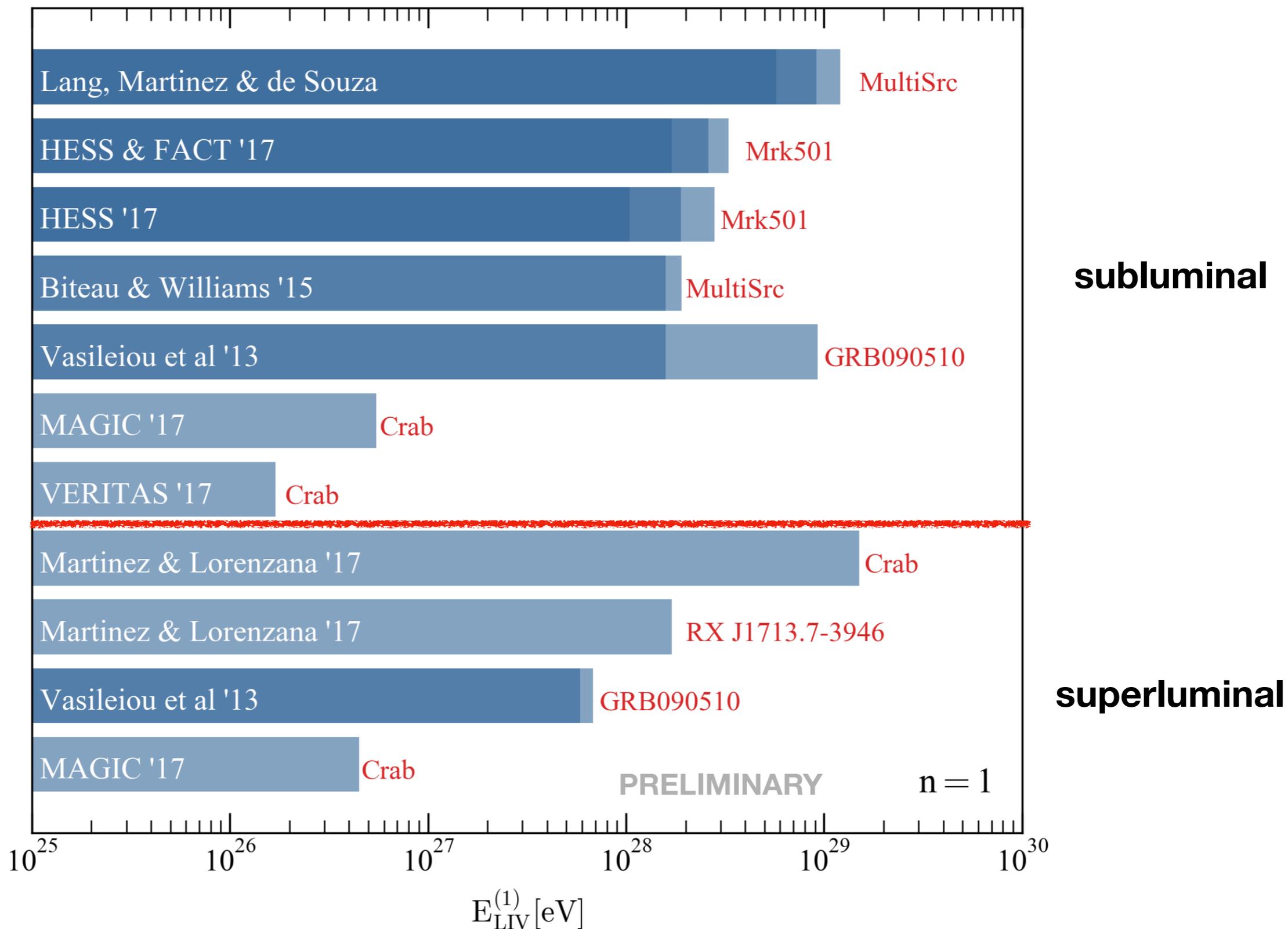
i. GZK-photons + LIV

ii. Limits

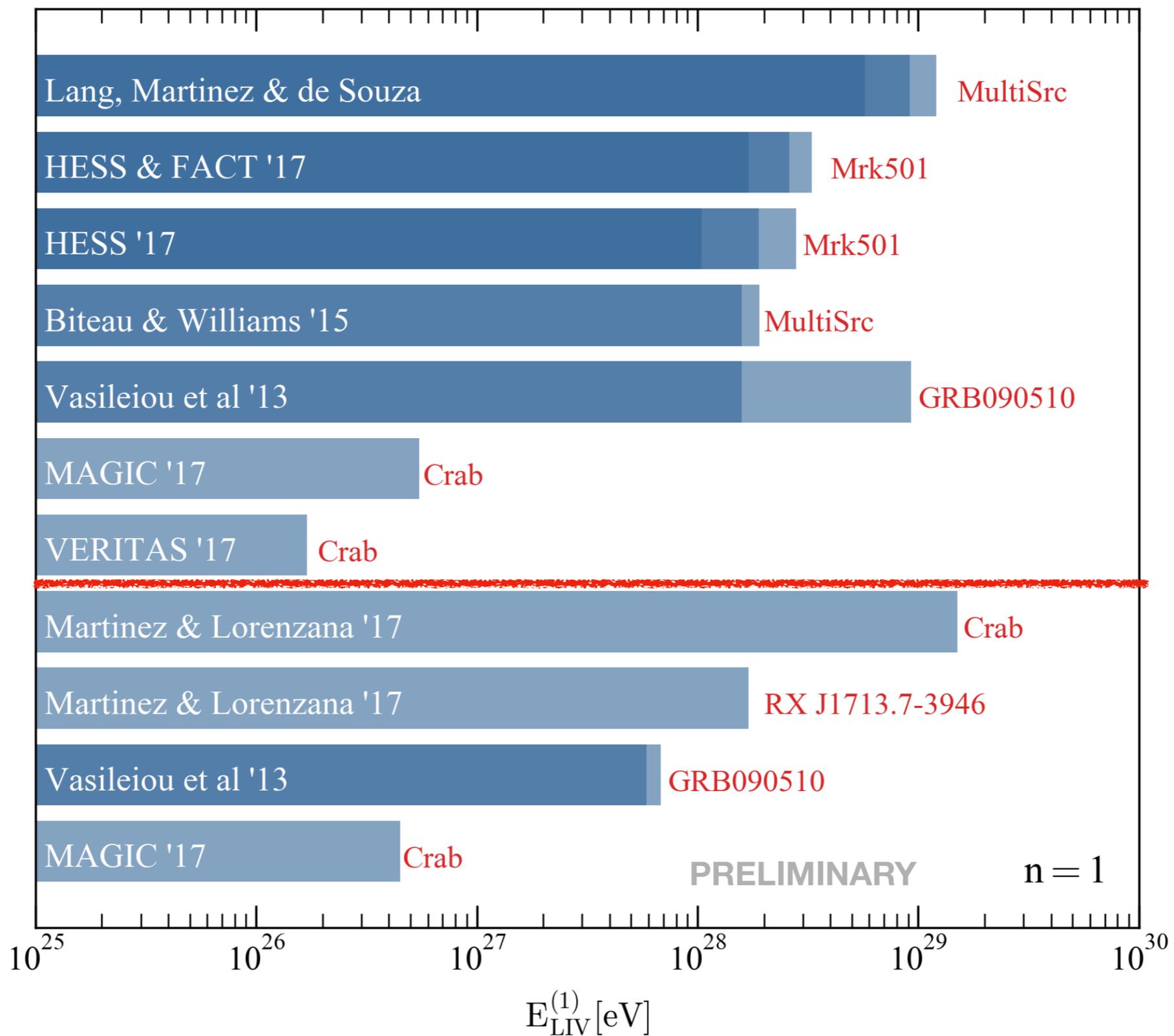
LIV limits by γ -rays



LIV limits by γ -rays



LIV limits by γ -rays



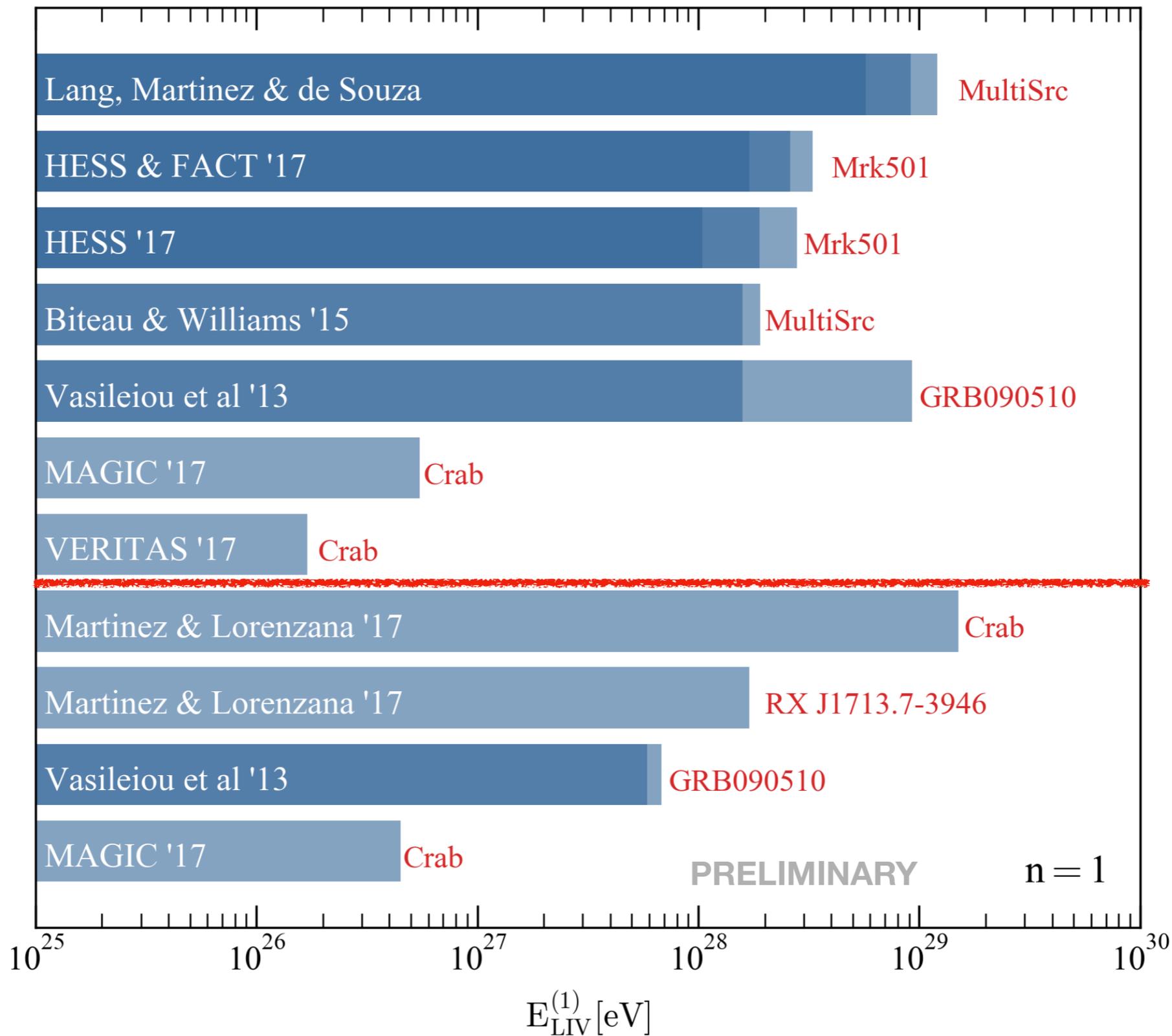
LIV limits by γ -rays

Pair production shift threshold

Time Energy Dependent -delay

Photon decay

Time Energy Dependent -delay



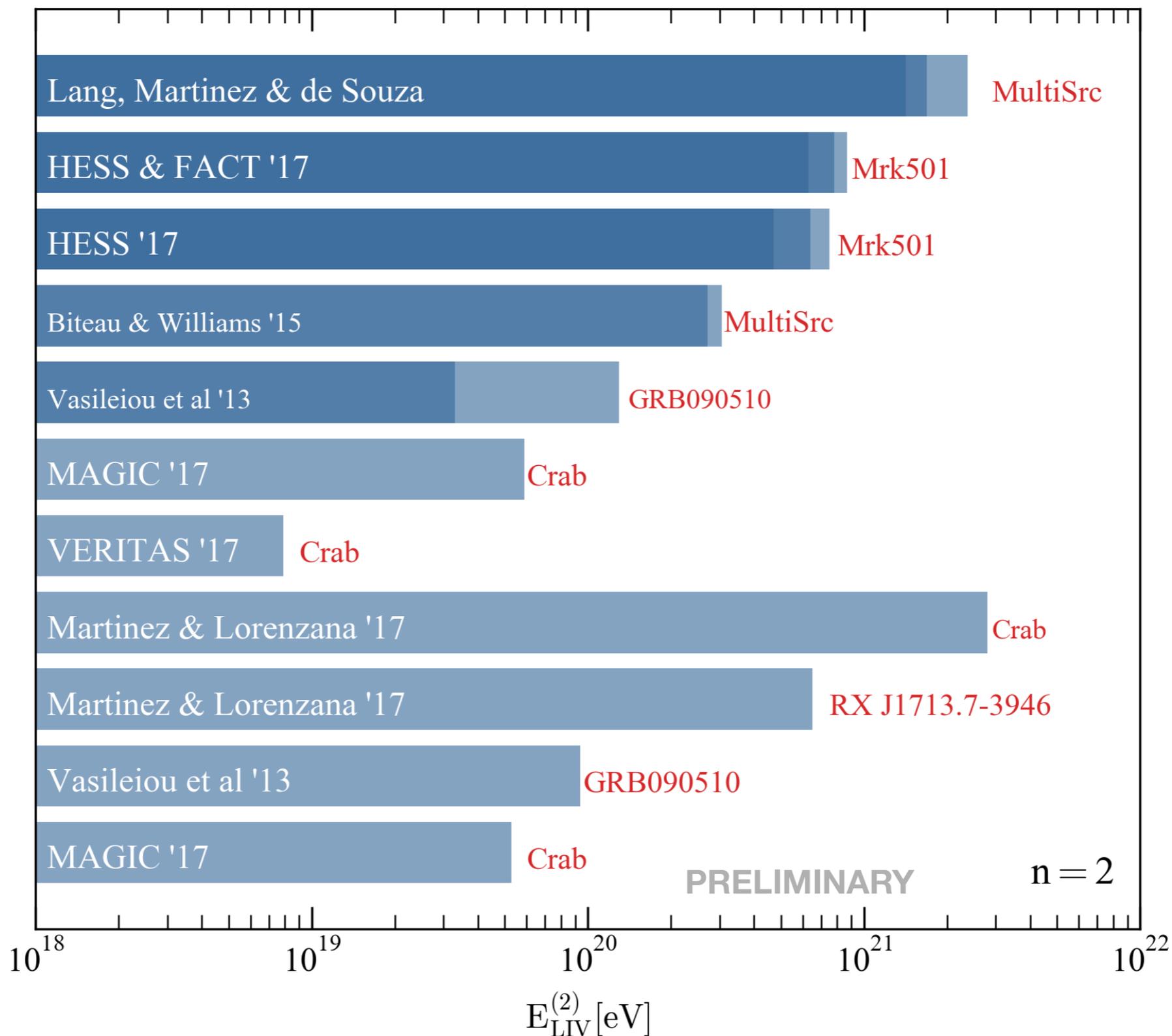
LIV limits by γ -rays

Pair production shift threshold

Time Energy Dependent -delay

Photon decay

Time Energy Dependent -delay



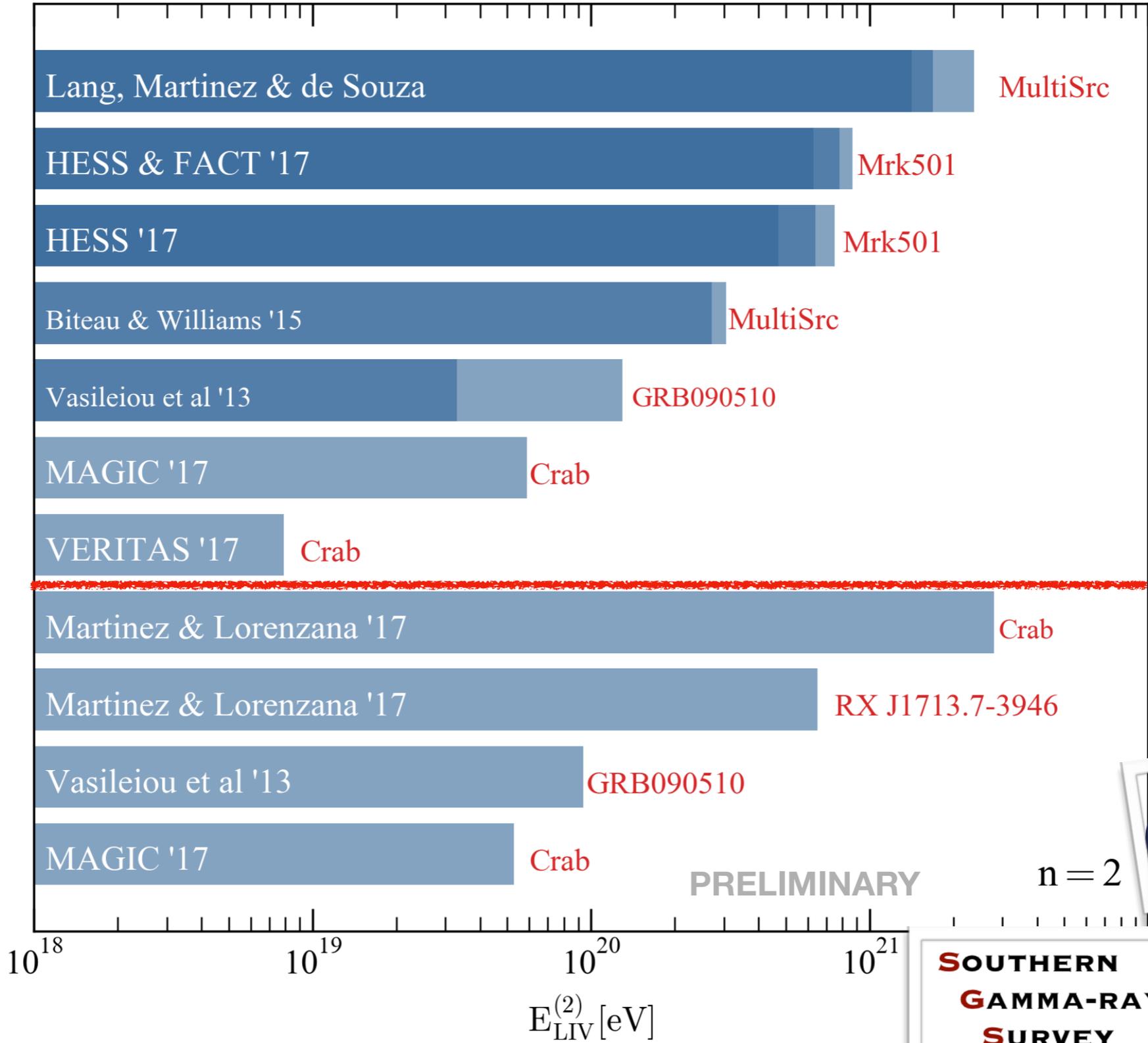
LIV limits by γ -rays

Pair production shift threshold

Time Energy Dependent -delay

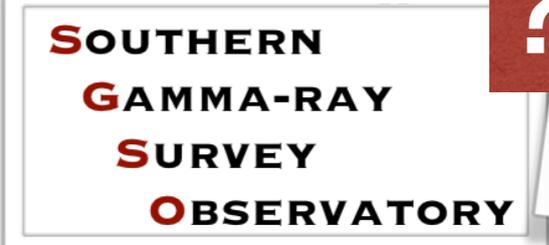
Photon decay

Time Energy Dependent -delay



PRELIMINARY

n = 2

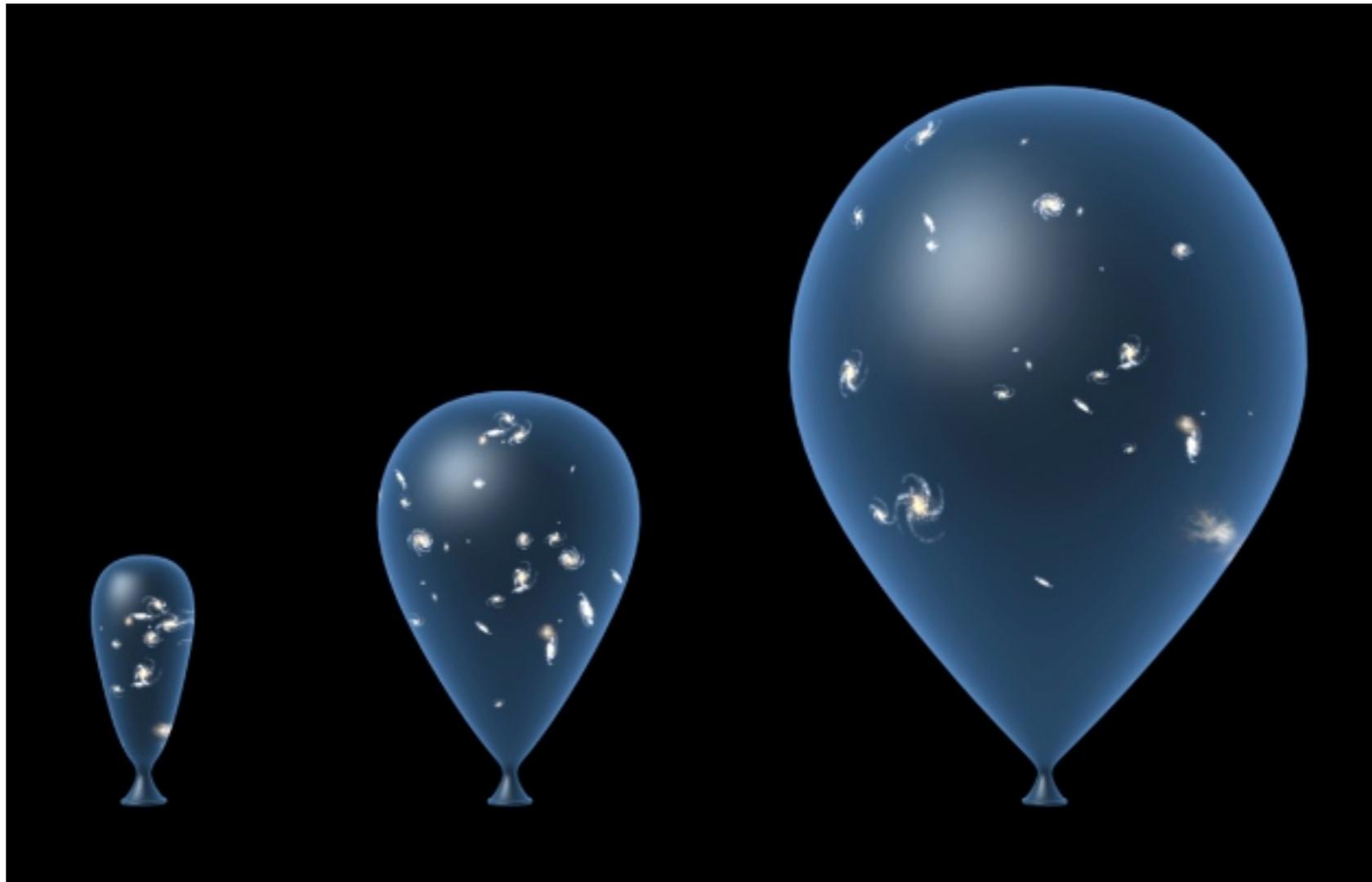


- I. Lorentz invariance violation (LIV)
- II. Limits: TeV γ -rays
 - i. Time Energy Dependent delay**
 - ii. Photon Decay
 - iii. Pair production threshold shifts
- III. UHECR
 - i. GZK-photons + LIV
 - ii. Limits

Time energy dependent delay

The corresponding differential relation between time and redshift is

$$dt = -H_0^{-1} \frac{dz}{(1+z)h(z)} \quad ; \quad h(z) = \sqrt{\Omega_\Lambda + \Omega_m(1+z)^3}$$



Time energy dependent delay

The corresponding differential relation between time and redshift is

$$dt = -H_0^{-1} \frac{dz}{(1+z)h(z)} \quad ; \quad h(z) = \sqrt{\Omega_\Lambda + \Omega_m(1+z)^3}$$

A particle with velocity u travels an elementary distance:

Two particles with velocities differing by Δu

$$u dt = -H_0^{-1} \frac{u dz}{(1+z)h(z)} \quad \longrightarrow \quad \Delta L = H_0^{-1} \int_0^z \frac{\Delta u dz}{(1+z)h(z)}$$

Time energy dependent delay

The corresponding differential relation between time and redshift is

$$dt = -H_0^{-1} \frac{dz}{(1+z)h(z)} \quad ; \quad h(z) = \sqrt{\Omega_\Lambda + \Omega_m(1+z)^3}$$

A particle with velocity u travels an elementary distance

Two particles with velocities differing by Δu

$$u dt = -H_0^{-1} \frac{u dz}{(1+z)h(z)} \quad \rightarrow \quad \Delta L = H_0^{-1} \int_0^z \frac{\Delta u dz}{(1+z)h(z)}$$

LIV - energy dependence group velocity

$$E_\gamma^2 = p^2 \pm \left(\frac{E_\gamma^2}{E_{LIV}^{(n)}} \right)^n \quad \rightarrow \quad u = \frac{dE}{dp} \approx 1 + \frac{n+1}{2} \left(\frac{E_\gamma}{E_{LIV}^{(n)}} \right)^n$$

Time energy dependent delay

The corresponding differential relation between time and redshift

$$dt = -H_0^{-1} \frac{dz}{(1+z)h(z)} \quad ; \quad h(z) = \sqrt{\Omega_\Lambda + \Omega_m(1+z)^3}$$

Let be u , a particle velocity

Δu : two particles with velocities differing by

$$u dt = -H_0^{-1} \frac{u dz}{(1+z)h(z)} \quad \rightarrow \quad \Delta L = H_0^{-1} \int_0^z \frac{\Delta u dz}{(1+z)h(z)}$$

LIV - energy dependence group velocity

$$E_\gamma^2 = p^2 \pm \left(\frac{E_\gamma^2}{E_{LIV}^{(n)}} \right)^n \quad \rightarrow \quad u = \frac{dE}{dp} \approx 1 + \frac{n+1}{2} \left(\frac{E_\gamma}{E_{LIV}^{(n)}} \right)^n$$

Time energy dependent delay

The corresponding differential relation between time and redshift

$$dt = -H_0^{-1} \frac{dz}{(1+z)h(z)} \quad ; \quad h(z) = \sqrt{\Omega_\Lambda + \Omega_m(1+z)^3}$$

Let be u , a particle velocity

Δu : two particles with velocities differing by

$$u dt = -H_0^{-1} \frac{u dz}{(1+z)h(z)} \quad \rightarrow \quad \Delta L = H_0^{-1} \int_0^z \frac{\Delta u dz}{(1+z)h(z)}$$

LIV - energy dependence group velocity

$$u = \frac{dE}{dp} \approx 1 + \frac{n+1}{2} \left(\frac{E_\gamma}{E_{LIV}^{(n)}} \right)^n$$

$$\Delta t = H_0^{-1} \left(\frac{\Delta E}{E_{LIV}^{(n)}} \right)^2 \int_0^z \frac{(1+z) dz}{h(z)}$$

Time energy dependent delay

Constraints on Lorentz Invariance Violation from *Fermi*-Large Area Telescope Observations of Gamma-Ray Bursts

V. Vasileiou,^{1,*} A. Jacholkowska,^{2,†} F. Piron,¹ J. Bolmont,² C. Couturier,²
 J. Granot,³ F. W. Stecker,^{4,5} J. Cohen-Tanugi,¹ and F. Longo^{6,7}

GRB Name	PairView		SMM		Likelihood ^a	
	95%	99.5%	$n=1, s_{\pm}=+1$ (E_{PI} units)		95%	99.5%
080916C	0.11	0.081	95%	99.5%	0.22	0.2
090510	7.6	1.3	0.09	0.067	5.2	4.2
090902B	0.17	0.13	5.9	1.2	0.12	0.074
090926A	–	0.55	0.15	0.11	1.2	0.45
			8	0.35		
			$n=1, s_{\pm}=-1$ (E_{PI} units)			
	95%	99.5%	95%	99.5%	95%	99.5%
080916C	18	0.33	5.4	0.31	0.2	0.18
090510	0.56	0.48	0.57	0.48	11	3.6
090902B	0.38	0.2	0.86	0.28	0.37	0.11
090926A	0.24	0.18	0.2	0.12	0.17	0.15
			$n=2, s_{\pm}=+1$ (10^{10} GeV units)			
	95%	99.5%	95%	99.5%	95%	99.5%
080916C	0.31	0.28	0.24	0.21	0.35	0.33
090510	6.7	3.3	13	3.3	8.6	6.4
090902B	0.8	0.72	0.73	0.64	0.64	0.49
090926A	0.67	0.48	9.1	1.6	0.48	0.47
			$n=2, s_{\pm}=-1$ (10^{10} GeV units)			
	95%	99.5%	95%	99.5%	95%	99.5%
080916C	–	0.69	–	5.2	0.34	0.32
090510	1.9	1.5	1.9	1.5	9.4	5.4
090902B	1.6	0.97	3.5	1.2	0.64	0.46
090926A	0.51	0.42	0.51	0.5	0.31	0.26

Time energy dependent delay

Constraints on Lorentz Invariance Violation from *Fermi*-Large Area Telescope Observations of Gamma-Ray Bursts

V. Vasileiou,^{1,*} A. Jacholkowska,^{2,†} F. Piron,¹ J. Bolmont,² C. Couturier,²
 J. Granot,³ F. W. Stecker,^{4,5} J. Cohen-Tanugi,¹ and F. Longo^{6,7}

GRB Name	PairView		SMM		Likelihood ^a	
	95%	99.5%	$n=1, s_{\pm}=+1$ (E_{PI} units)		95%	99.5%
080916C	0.11	0.081	0.09	0.067	0.22	0.2
090510	7.6	1.3	5.9	1.2	5.2	4.2
090902B	0.17	0.13	0.15	0.11	0.12	0.074
090926A	–	0.55	8	0.35	1.2	0.45
	95%	99.5%	$n=1, s_{\pm}=-1$ (E_{PI} units)		95%	99.5%
080916C	18	0.33	5.4	0.31	0.2	0.18
090510	0.56	0.48	0.57	0.48	11	3.6
090902B	0.38	0.2	0.86	0.28	0.37	0.11
090926A	0.24	0.18	0.2	0.12	0.17	0.15
	95%	99.5%	$n=2, s_{\pm}=+1$ (10^{10} GeV units)		95%	99.5%
080916C	0.31	0.28	0.24	0.21	0.35	0.33
090510	6.7	3.3	13	3.3	8.6	6.4
090902B	0.8	0.72	0.73	0.64	0.64	0.49
090926A	0.67	0.48	9.1	1.6	0.48	0.47
	95%	99.5%	$n=2, s_{\pm}=-1$ (10^{10} GeV units)		95%	99.5%
080916C	–	0.69	–	5.2	0.34	0.32
090510	1.9	1.5	1.9	1.5	9.4	5.4
090902B	1.6	0.97	3.5	1.2	0.64	0.46
090926A	0.51	0.42	0.51	0.5	0.31	0.26

Time energy dependent delay



$$\Delta t = H_0^{-1} \left(\frac{\Delta E}{E_{LIV}^{(n)}} \right)^2 \int_0^z \frac{(1+z)dz}{h(z)}$$

Spain
2 Telescopes

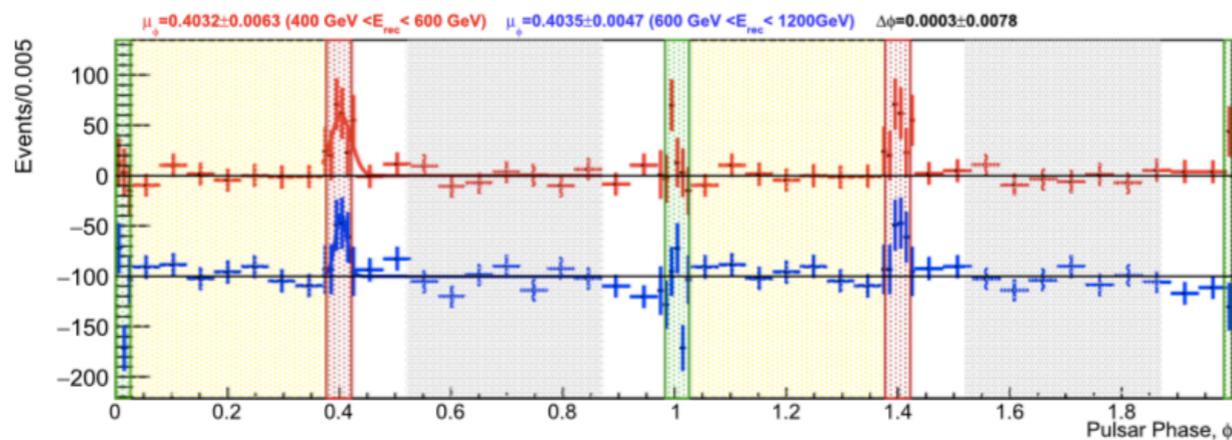
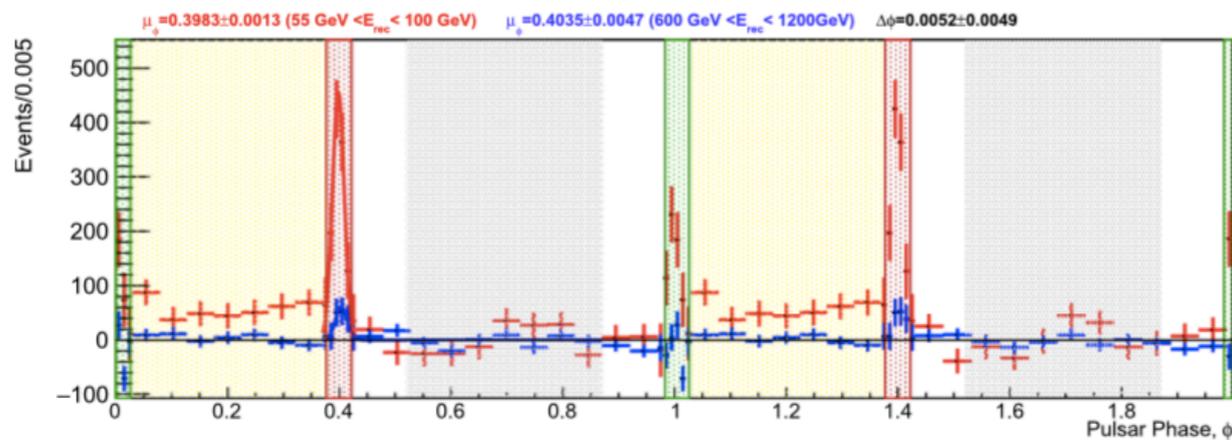


Table 2
Obtained 95%CL Limits from the Peak Comparison Method

Case	55–100 GeV versus 600–1200 GeV	400–600 GeV versus 600–1200 GeV
	E_{QG1} (GeV)	
$\xi_1 = +1$	2.5×10^{17}	1.1×10^{17}
$\xi_1 = -1$	6.7×10^{17}	1.1×10^{17}
	E_{QG2} (GeV)	
$\xi_2 = +1$	1.8×10^{10}	1.4×10^{10}
$\xi_2 = -1$	2.9×10^{10}	1.5×10^{10}

Time energy dependent delay



Source	Experiment	Limit on $E_{\text{QG}}^{(1)}$	Limit on $E_{\text{QG}}^{(2)}$	Distance	Δt	E_{max}	Ref.
HAWC Pulsar ref.	HAWC	10^{17} GeV	$9 \cdot 10^9$ GeV	2kpc	1 ms	500GeV	
HAWC GRB ref.	HAWC	$4.9 \cdot 10^{19}$ GeV	$1.1 \cdot 10^{11}$ GeV	$z = 1$	1 s	100GeV	

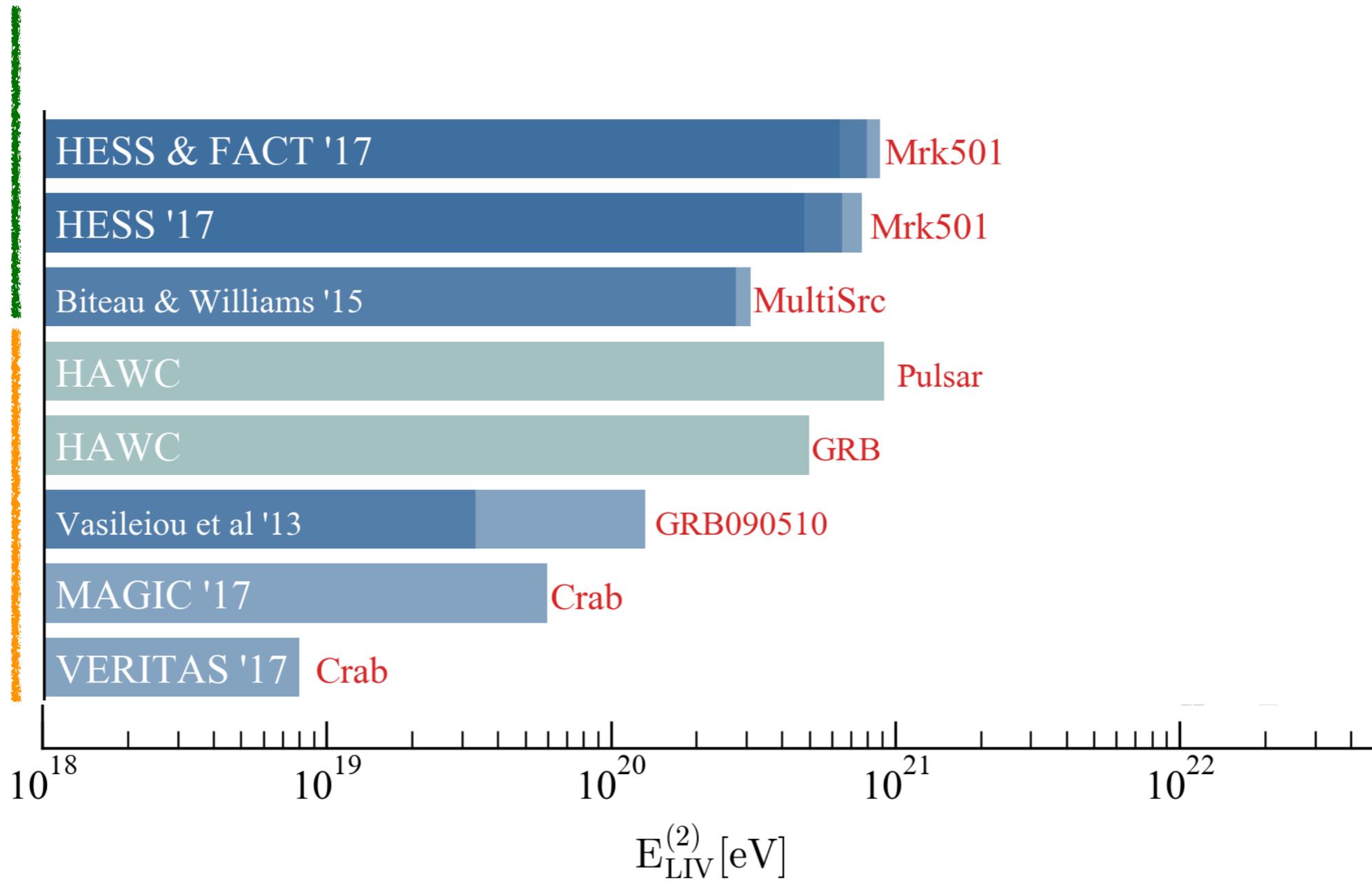
The potential of the HAWC observatory,
based on the reference scenarios

Time energy dependent delay



Pair production shift threshold

Time Energy Dependent



- I. Lorentz invariance violation (LIV)

- II. Limits: TeV γ -rays
 - i. Time Energy Dependent delay

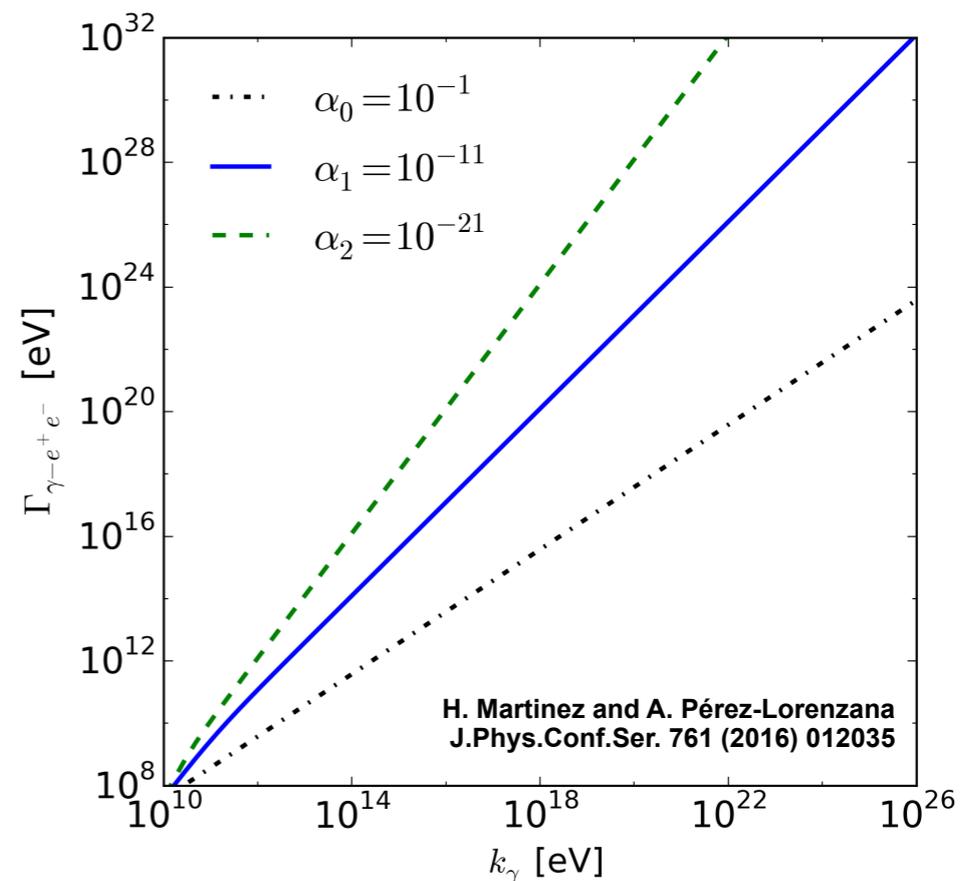
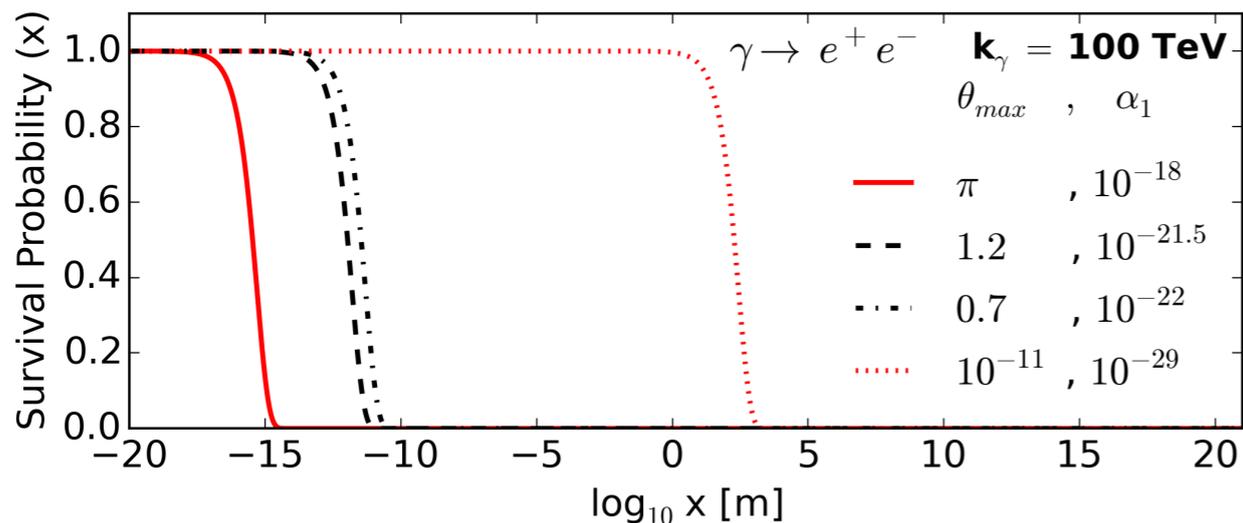
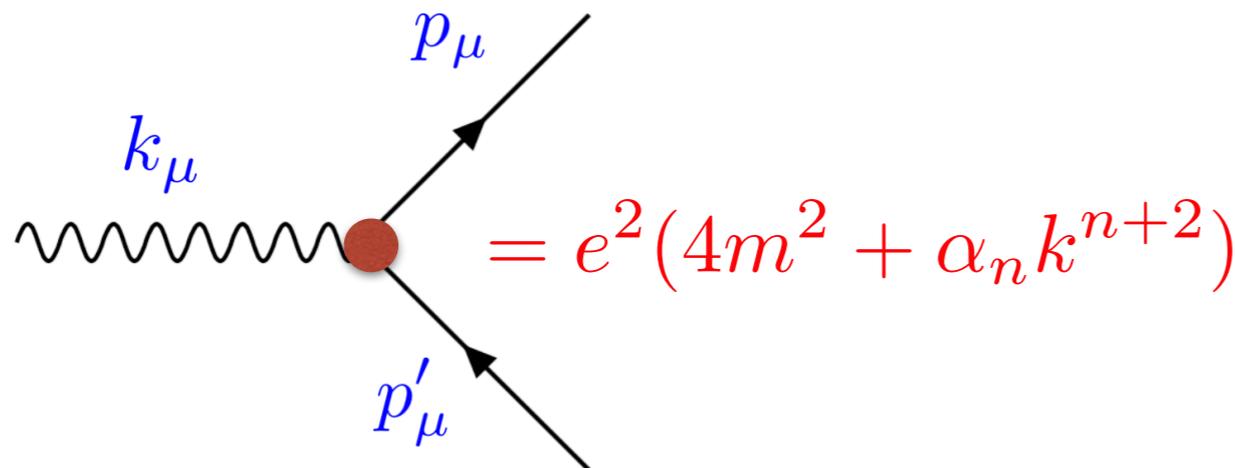
 - ii. Photon Decay**

 - iii. Pair production threshold shifts

- III. UHECR
 - i. GZK-photons + LIV

 - ii. Limits

Photon decay

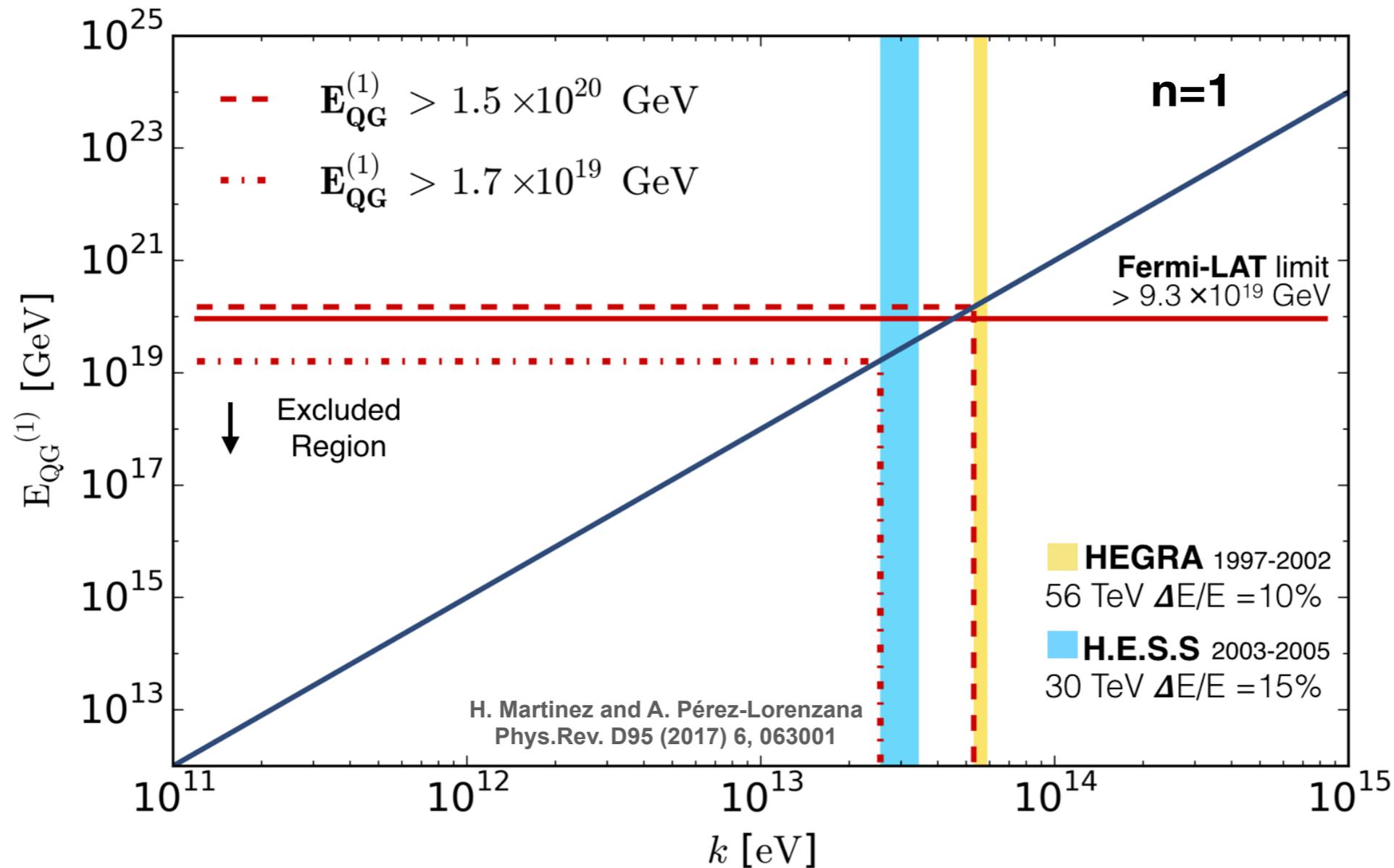


Above this energy threshold, the decay rate is quite efficient that photons should not arrive at Earth from astrophysical distances

$$E_{LIV}^{(n)} > E_\gamma \left[\frac{E_\gamma^2 - 4m_{e^-}^2}{4m_{e^-}^2} \right]^{1/n}$$

If you observe VHE gamma-rays, the LIV process is restricted!!

E_{LIV} excluded region due to $\gamma \rightarrow e^+e^-$



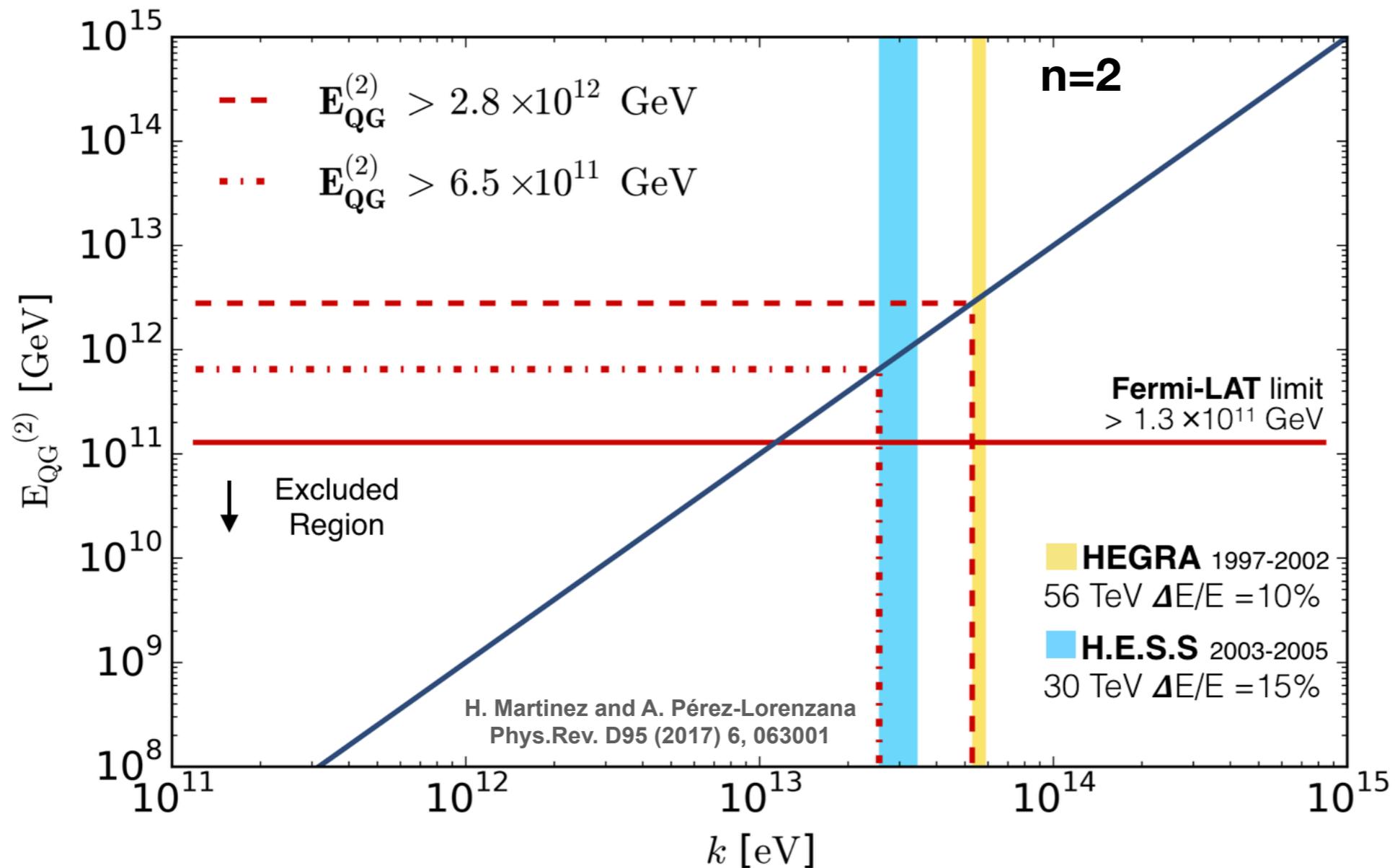
CRAB

HEGRA Collaboration,
 The Astrophysical Journal,
 614:897-913, 2004

SNR RX J1713.7.3946

HESS Collaboration,
 Astron. Astrophys, 449, 223
 (2007)

E_{LIV} excluded region due to $\gamma \rightarrow e^+e^-$

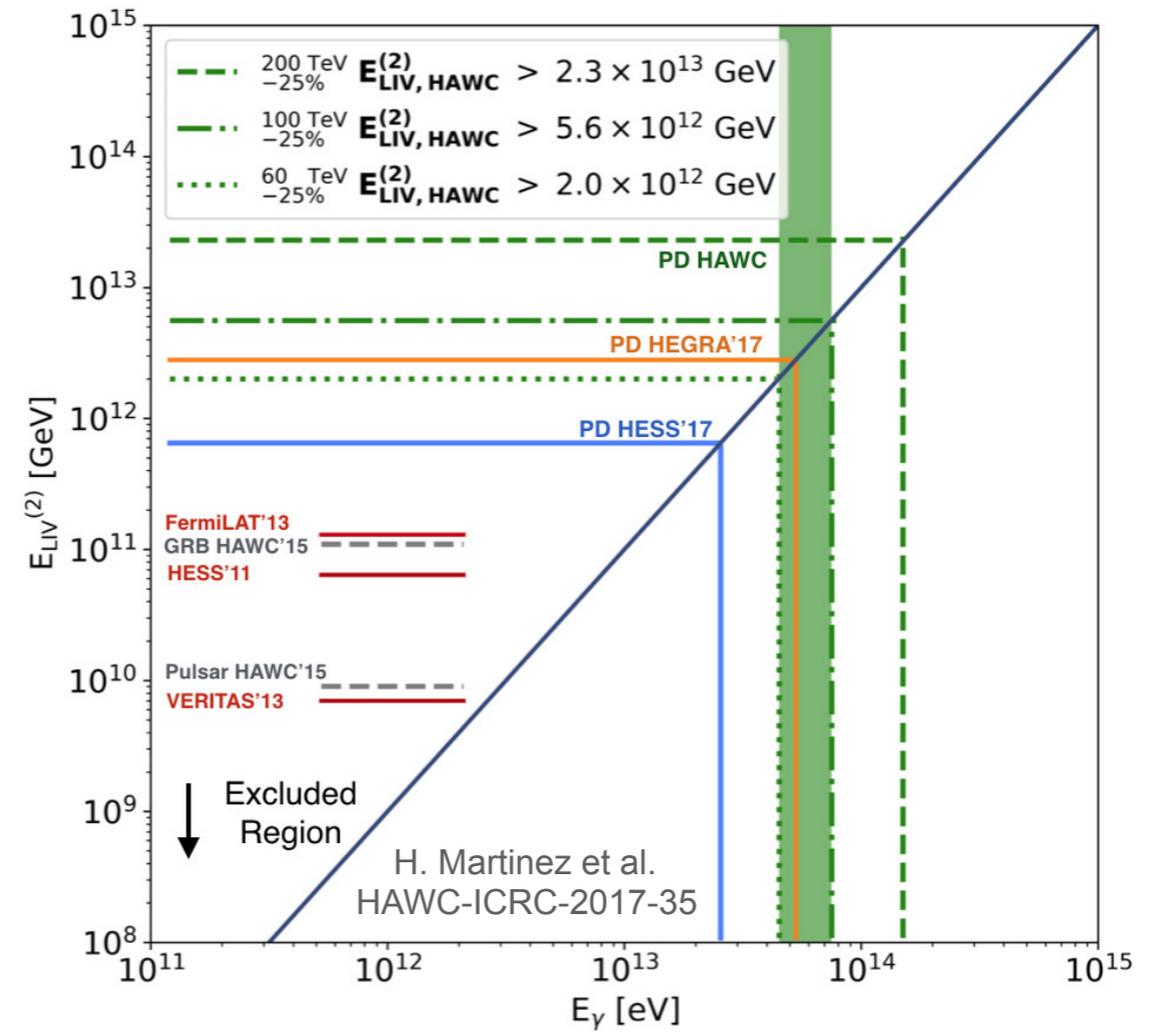
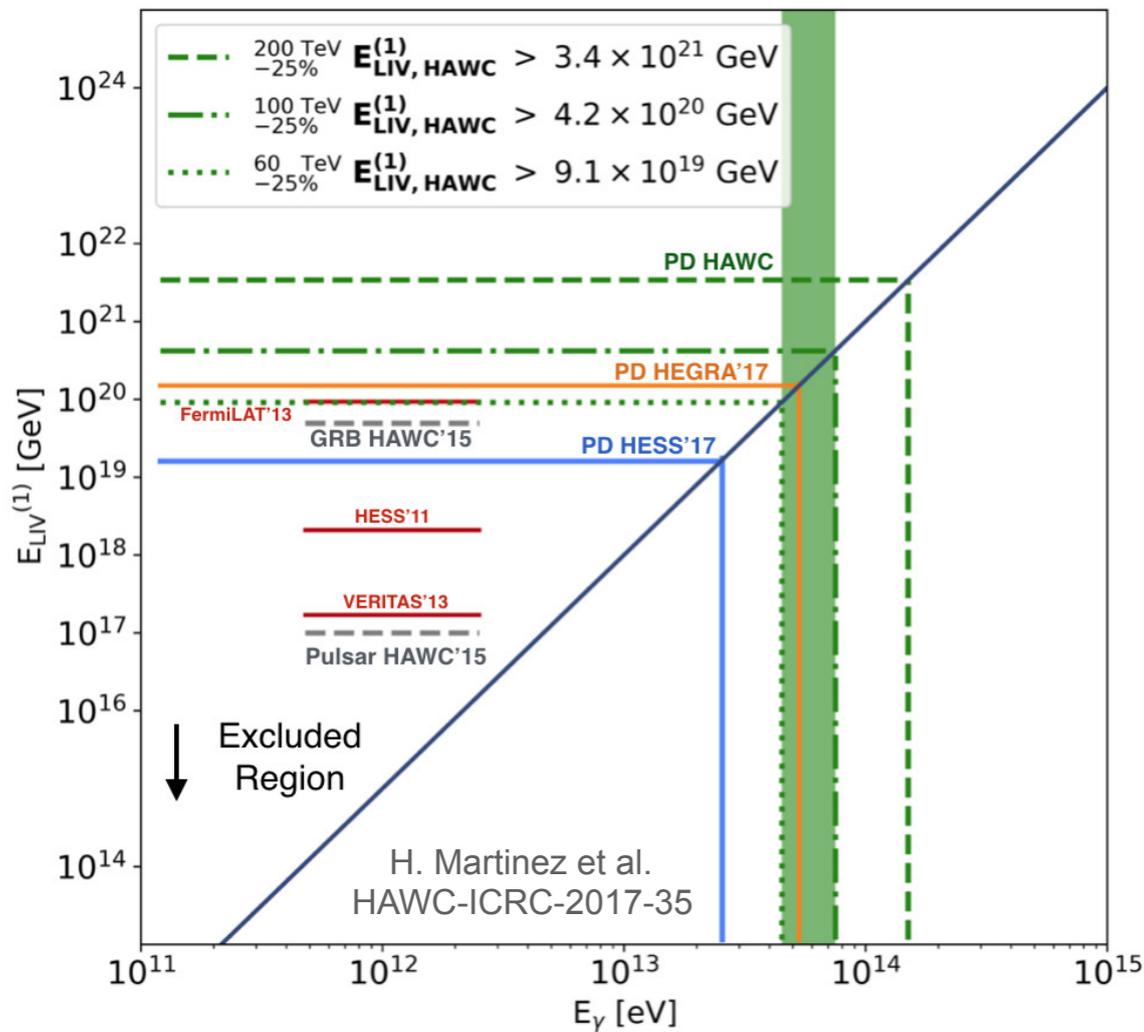


CRAB

HEGRA Collaboration,
 The Astrophysical Journal,
 614:897-913, 2004

SNR RX J1713.7-3946

HESS Collaboration,
 Astron. Astrophys, 449, 223
 (2007)

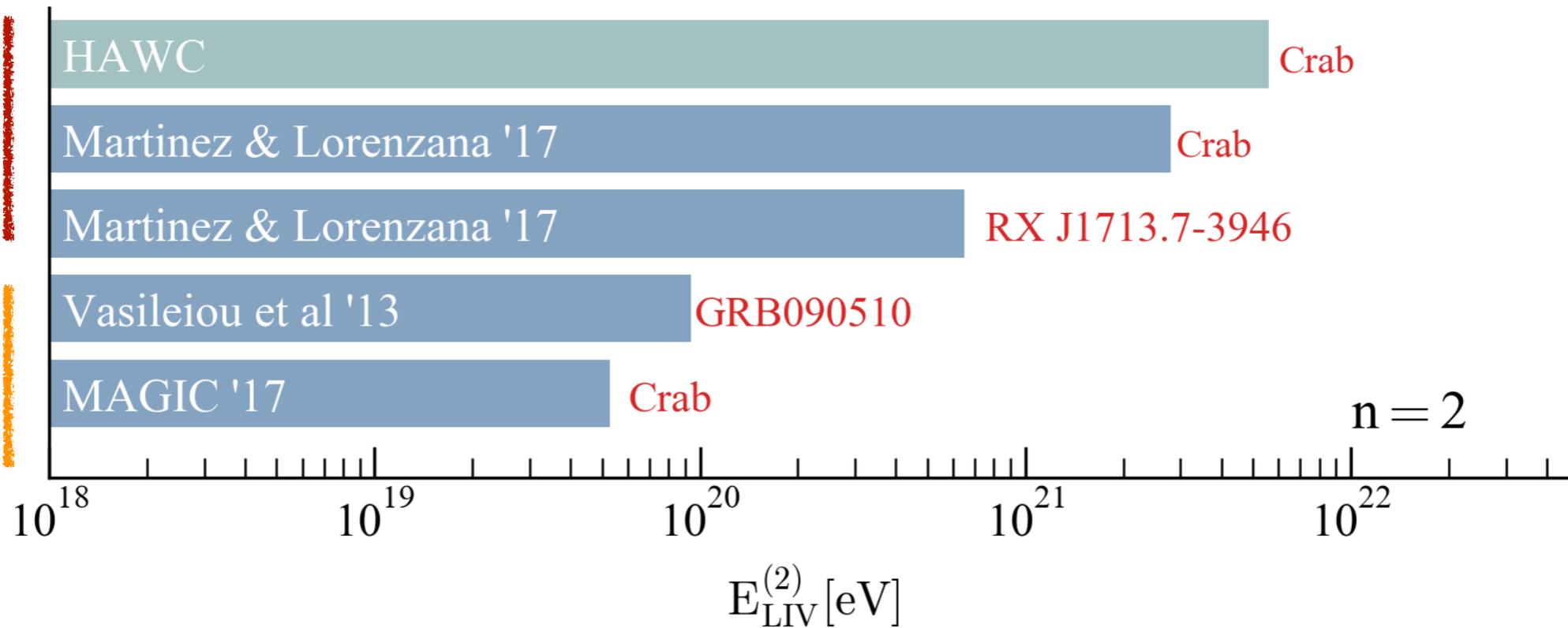


ELIV potential limits from HAWC
for $n=1$ (left), $n=2$ (right) and energy unc. of 25%.



Photon
decay

Time
Energy
Dependent



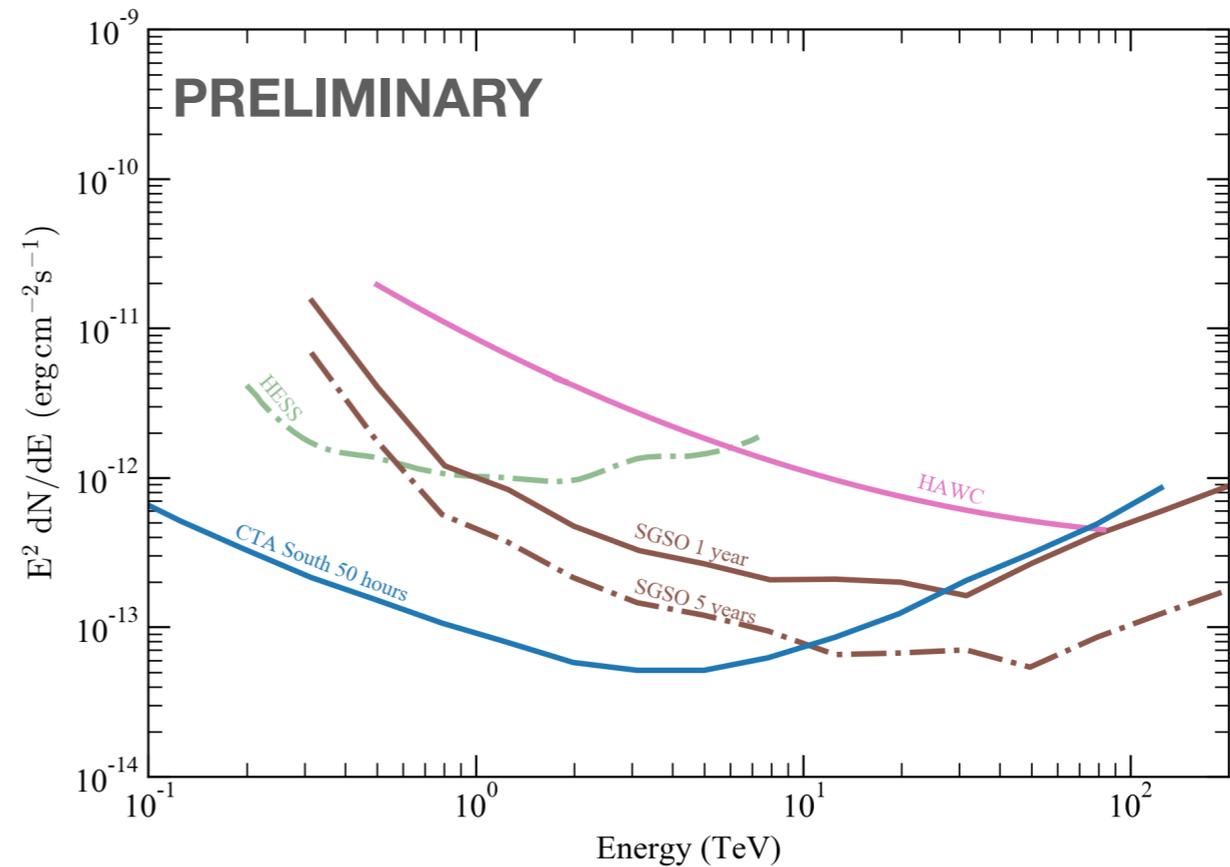
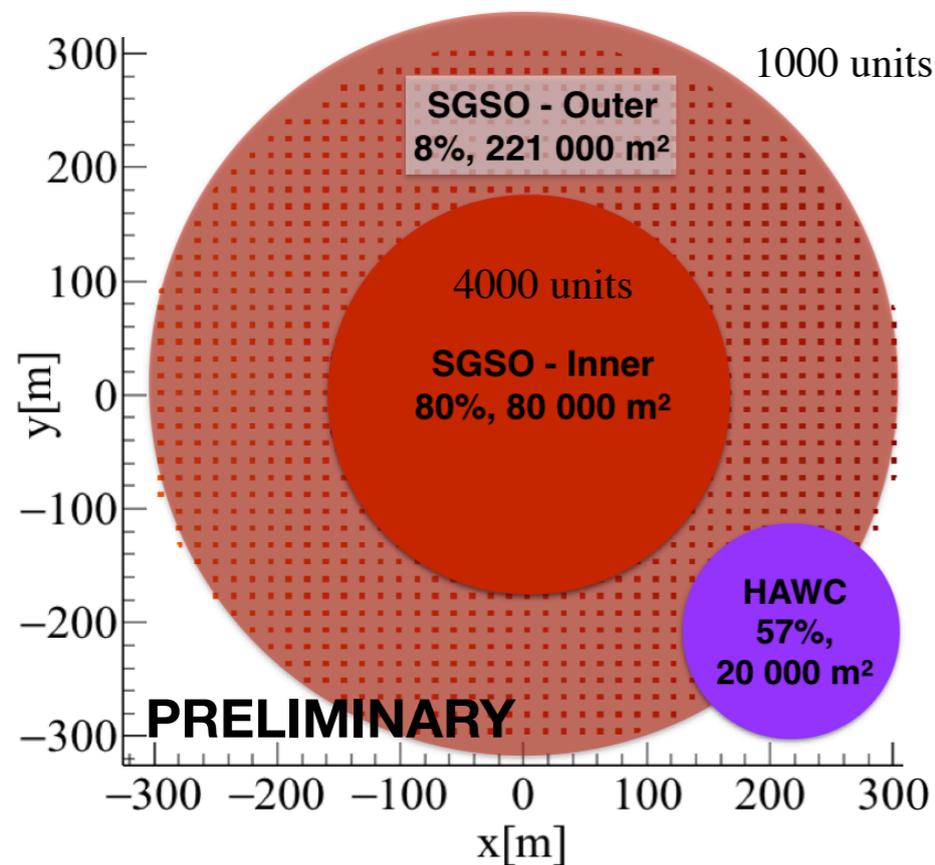
ELIV potential limits from HAWC

n=2 and energy unc. of 25%.

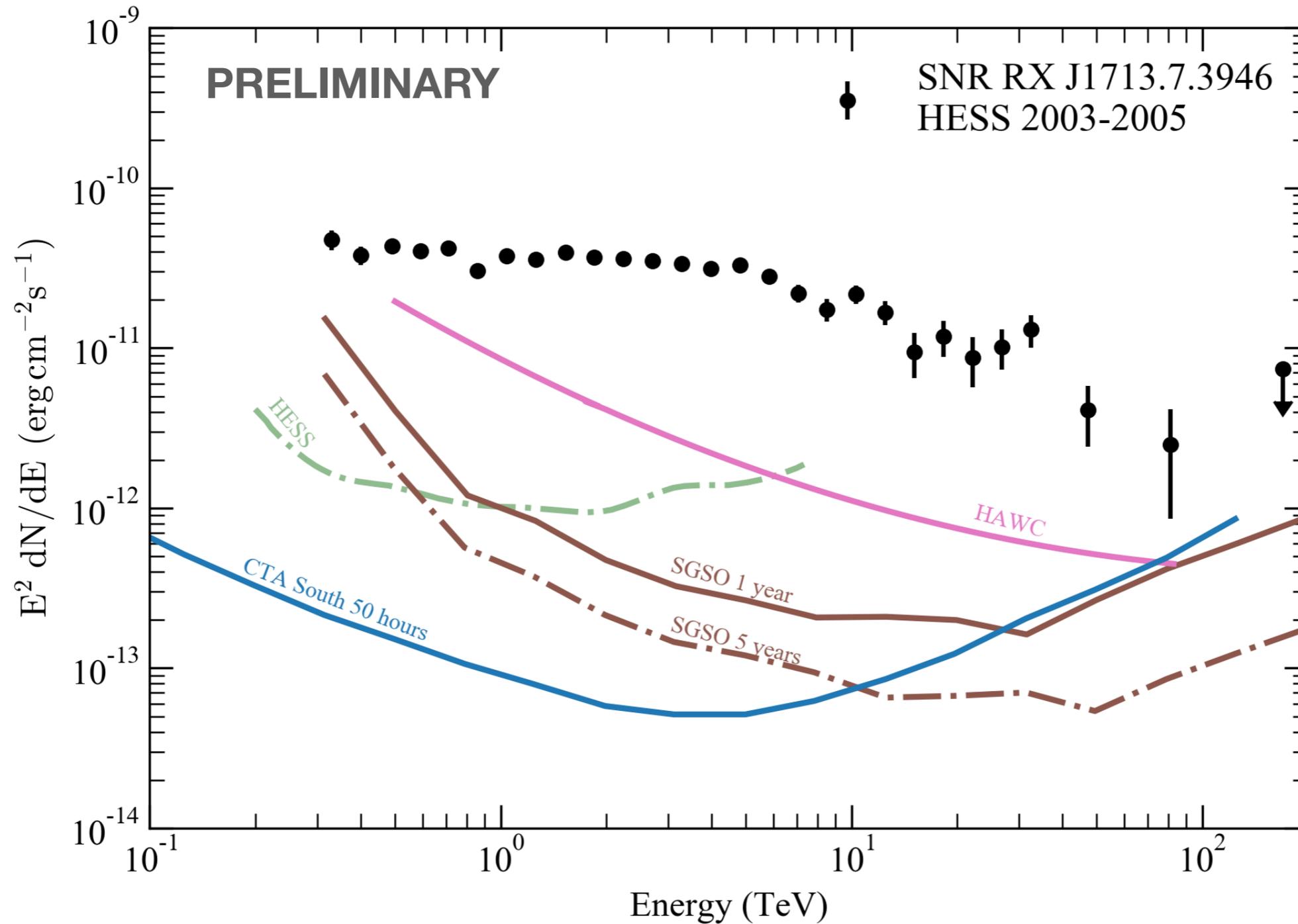
Searching LIV signatures with SGSO



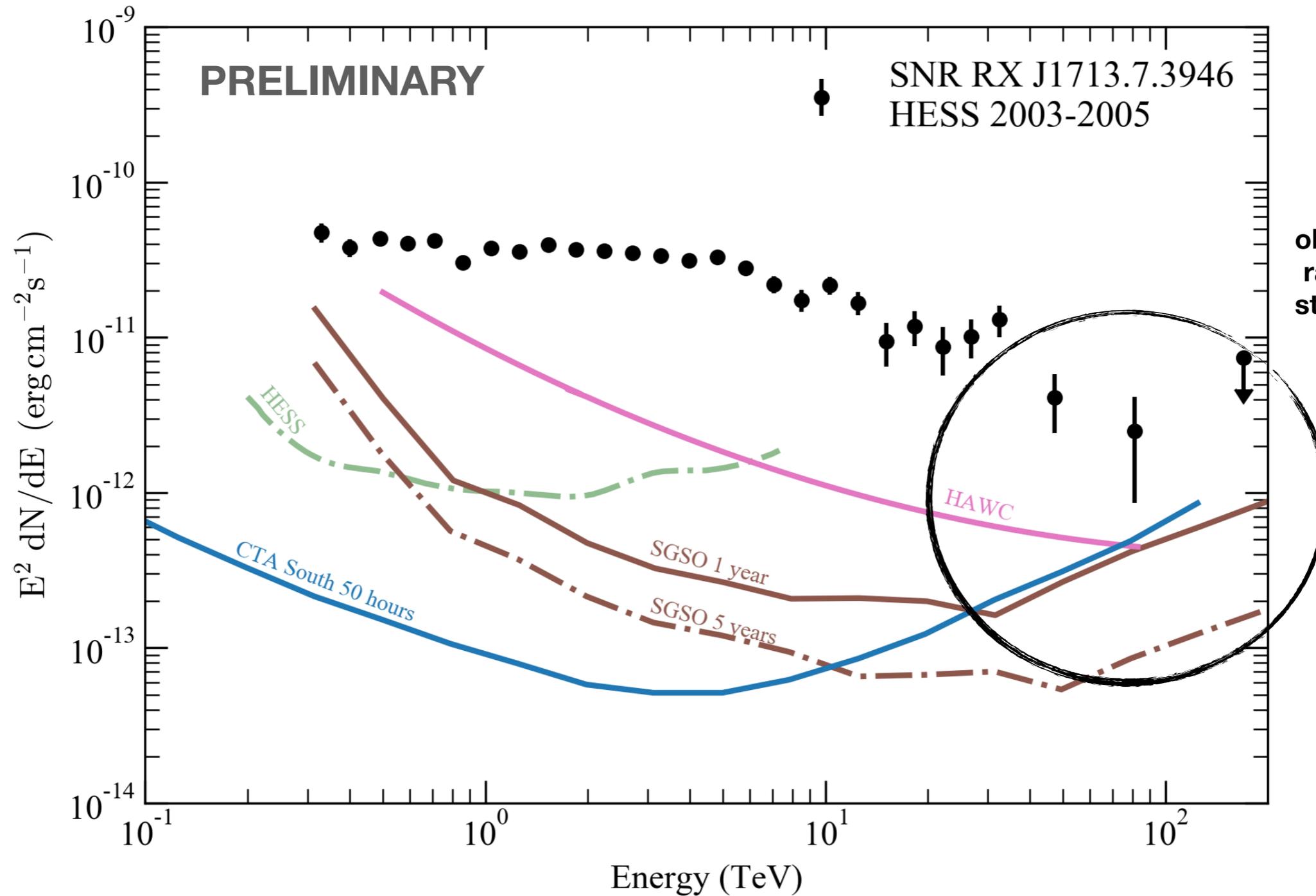
Cerro Vecar, province of Salta, Argentina.



Searching LIV signatures with SGSO



Searching LIV signatures with SGSO



The higher the observed gamma-ray energy is, the stringent the limit!

Searching LIV signatures with SGSO

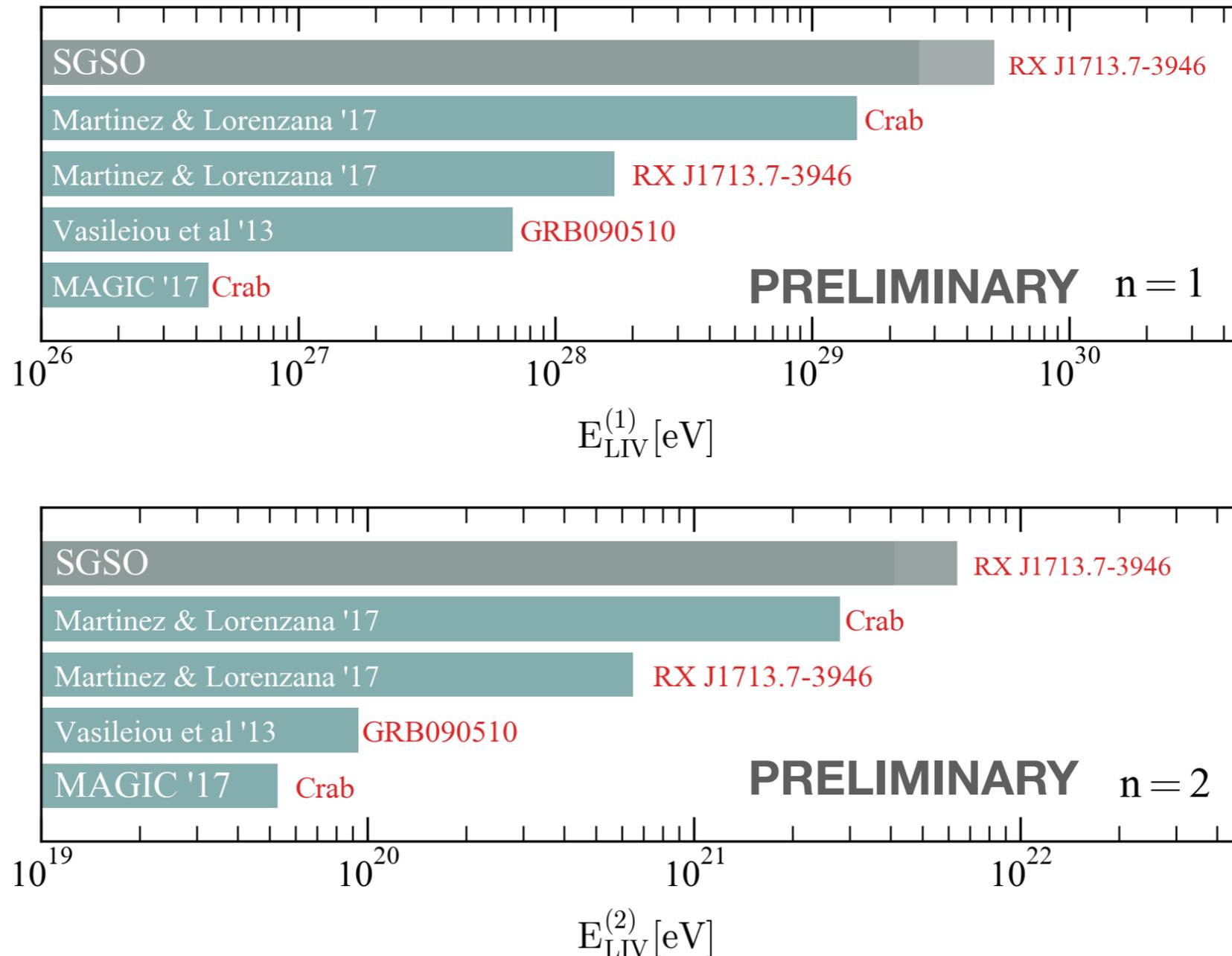


Fig. 1

LIV energy scale limits from superluminal searches including a potential reference of SGSO by measuring RXJ1713.7-3946 photons at **80 TeV** and **100 TeV** and the absence of photon decay into electron-positron pairs.

White paper in
preparation

- I. Lorentz invariance violation (LIV)

- II. Limits: TeV γ -rays
 - i. Time Energy Dependent delay

 - ii. Photon Decay

 - iii. Pair production threshold shifts**

- III. UHECR
 - i. GZK-photons + LIV

 - ii. Limits

Pair Production

$$\gamma_{VHE} \gamma_{BKG} \rightarrow e^+ e^-$$

$$\Lambda_{\gamma,n} x_{\gamma}^{n+2} + x_{\gamma} - 1 = 0$$

$$x_{\gamma} = \frac{E_{\gamma}}{E_{\gamma}^{LI}}, \quad \Lambda_{\gamma,n} = \frac{E_{\gamma}^{LI(n+1)}}{4\epsilon} \delta_{\gamma,n}$$

$$\Lambda_n < 0$$

Threshold-shifts

$$\Lambda_n = 0$$

LI scenario

$$\Lambda_n > 0$$

+2nd Threshold

The threshold equation

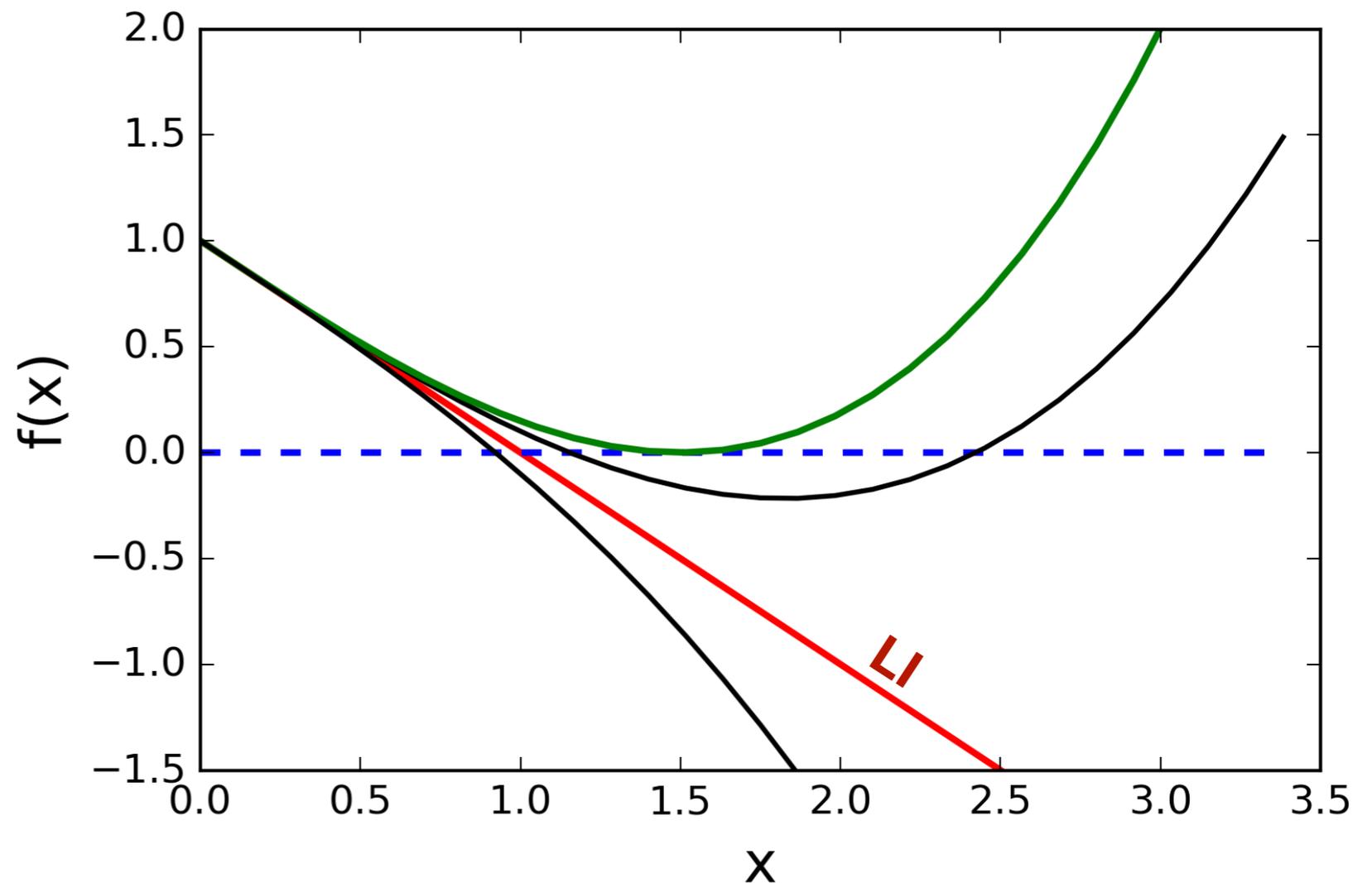
$$\delta_{\gamma,n} E_{\gamma}^{n+2} + 4E_{\gamma}\epsilon - m_e^2 \frac{1}{K(1-K)} = 0$$

Critical point

$$\delta_{\gamma,n}^{lim} = -4 \frac{\epsilon}{E_{\gamma}^{LI(n+1)}} \frac{(n+1)^{n+1}}{(n+2)^{n+2}}$$

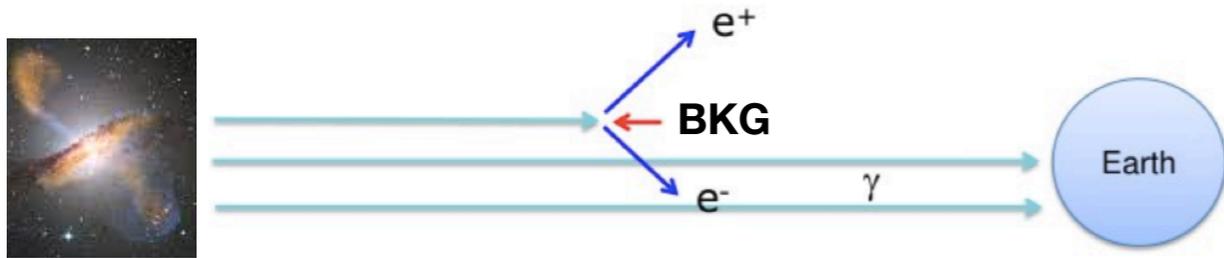
Background:

$$\epsilon_{th}^{LIV} = \frac{m_e^2}{4E_{\gamma}K(1-K)} - \frac{\delta_{\gamma,n}E_{\gamma}^{n+1}}{4}$$

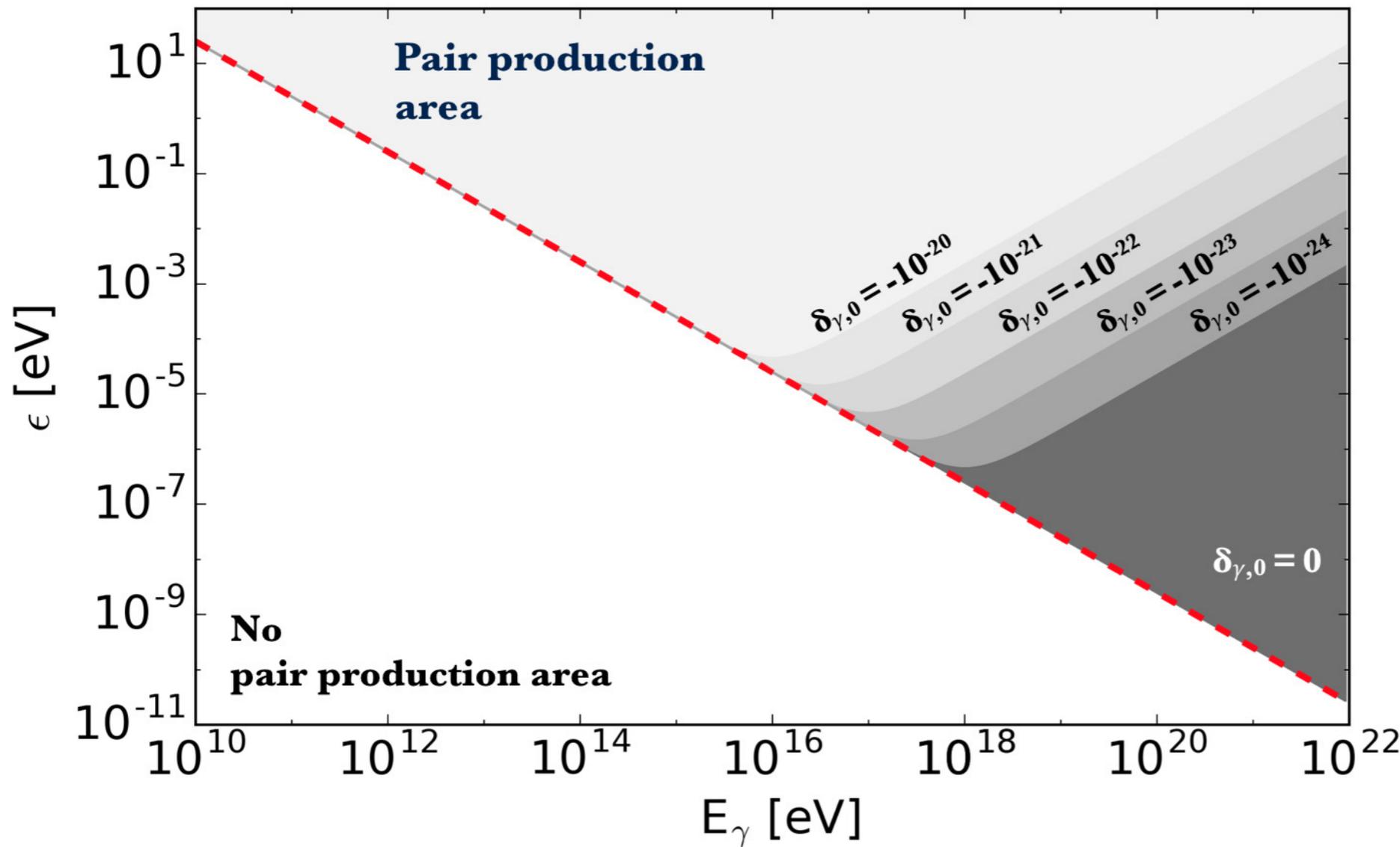


Pair Production

$$\gamma_{VHE} \gamma_{BKG} \rightarrow e^+ e^-$$



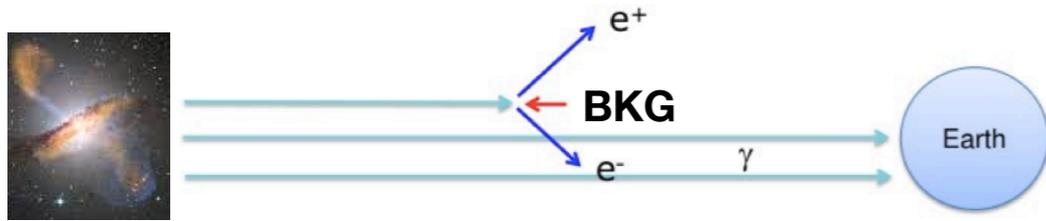
$$\epsilon_{th}^{LIV} = \frac{m_e^2}{4E_\gamma K(1-K)} - \frac{\delta_{\gamma,n} E_\gamma^{n+1}}{4}$$



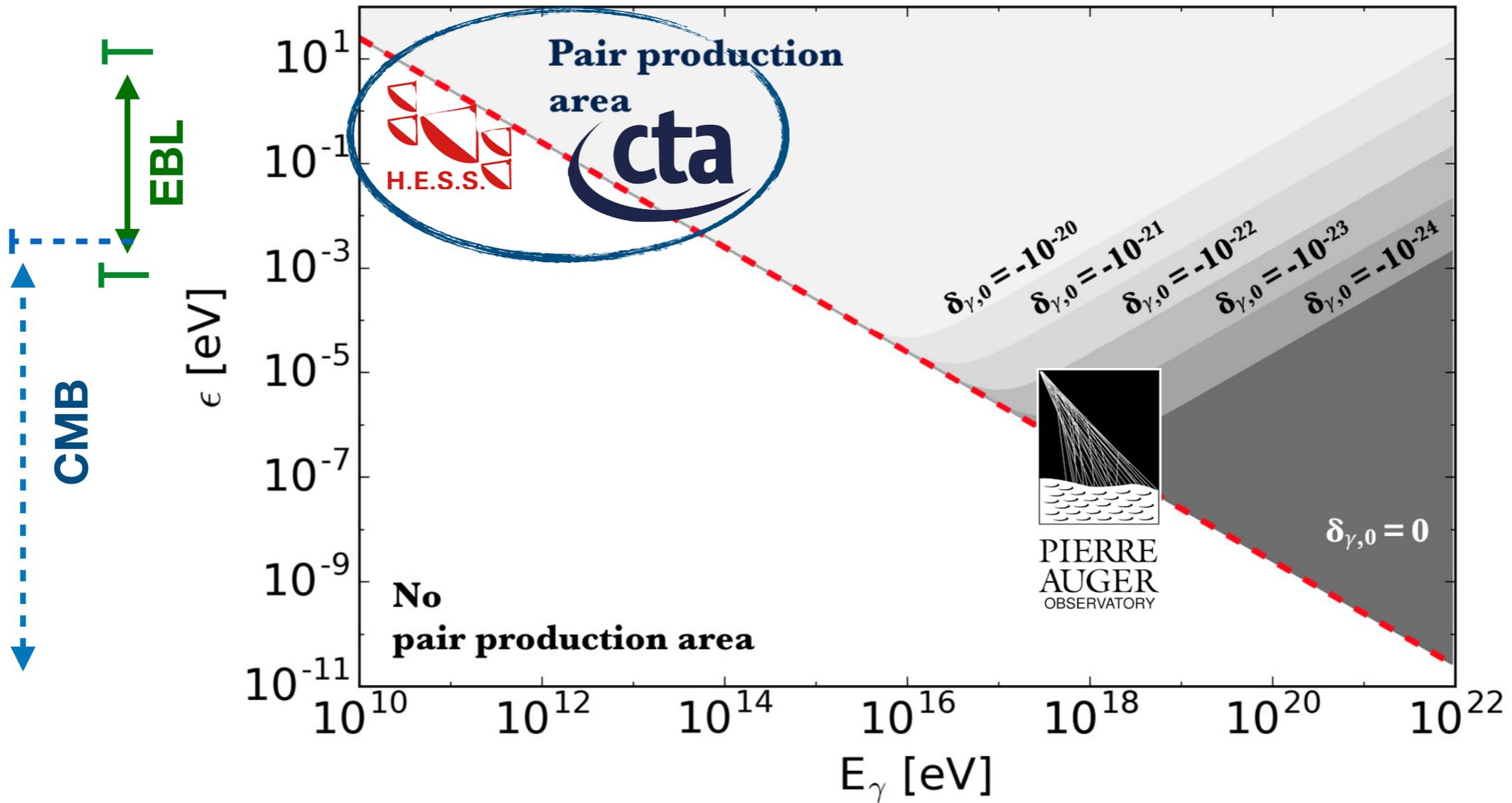
Allowed region change with the LIV parameter and the Energy

Pair Production

$$\gamma_{VHE} \gamma_{BKG} \rightarrow e^+ e^-$$

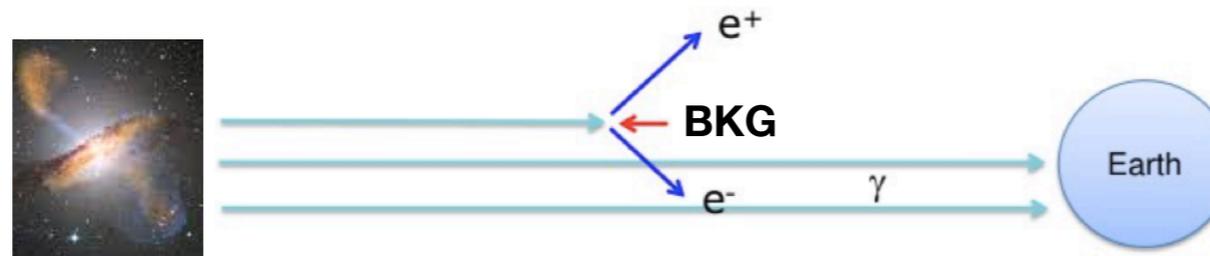


$$\epsilon_{th}^{LIV} = \frac{m_e^2}{4E_\gamma K(1-K)} - \frac{\delta_{\gamma,n} E_\gamma^{n+1}}{4}$$



... deeper LIV effects

Optical depth



$$\tau_{\gamma}(E_{\gamma}, z, n) = \int_0^z dz \frac{c}{H_0(1+z) \sqrt{\Omega_{\Lambda} + \Omega_M(1+z)^3}}$$

The distance element

$$\times \int_{\epsilon_{th}}^{\infty} d\epsilon n_{\gamma}(\epsilon, z)$$

Density of BKG photons

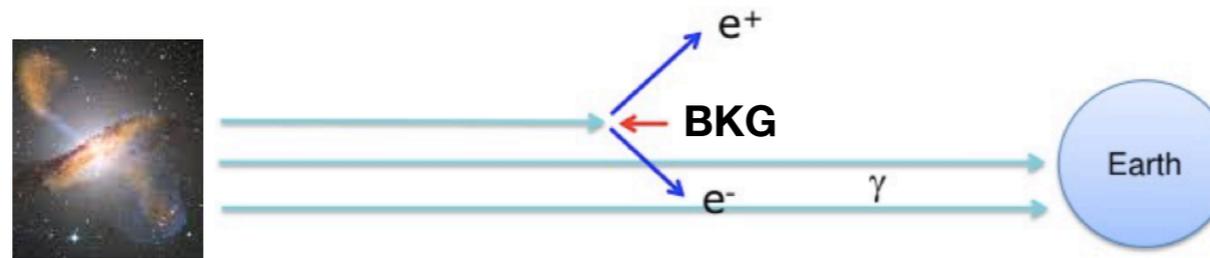
$$\times \int_{-1}^1 d(\cos \theta) \frac{1 - \cos \theta}{2} \sigma(E_{\gamma}, \epsilon, z, \cos \theta)$$

Pair Production cross section

Breit & Wheeler 1934; Heitler 1960

De Angelis, Alessandro et al.
Mon.Not.Roy.Astron.Soc.
432 (2013) 3245-3249

Optical depth + LIV



$$\tau_\gamma(E_\gamma, z, \eta, n, E_{LIV}^{(n)}) = \int_0^z dz \frac{c}{H_0(1+z) \sqrt{\Omega_\Lambda + \Omega_M(1+z)^3}}$$

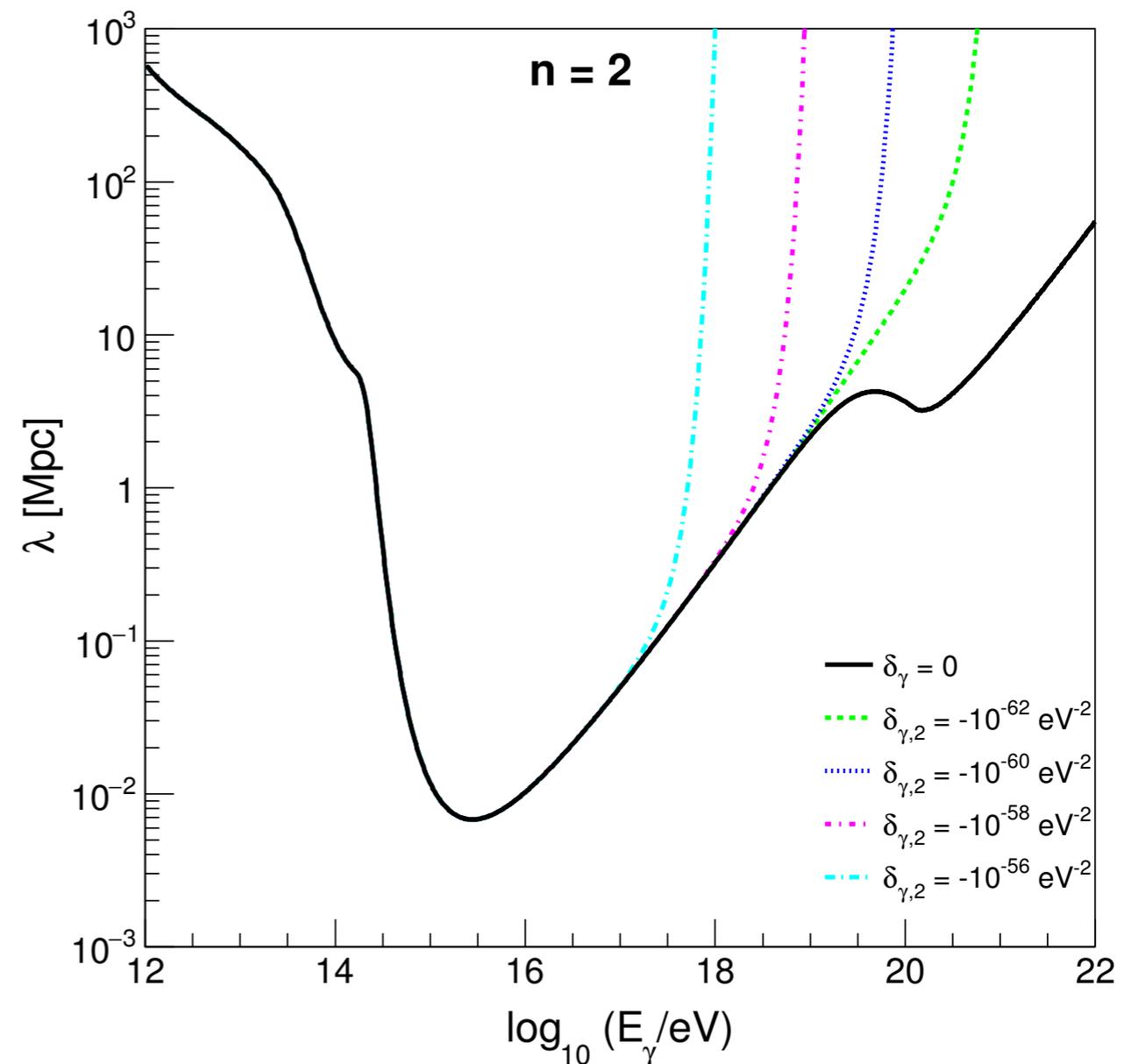
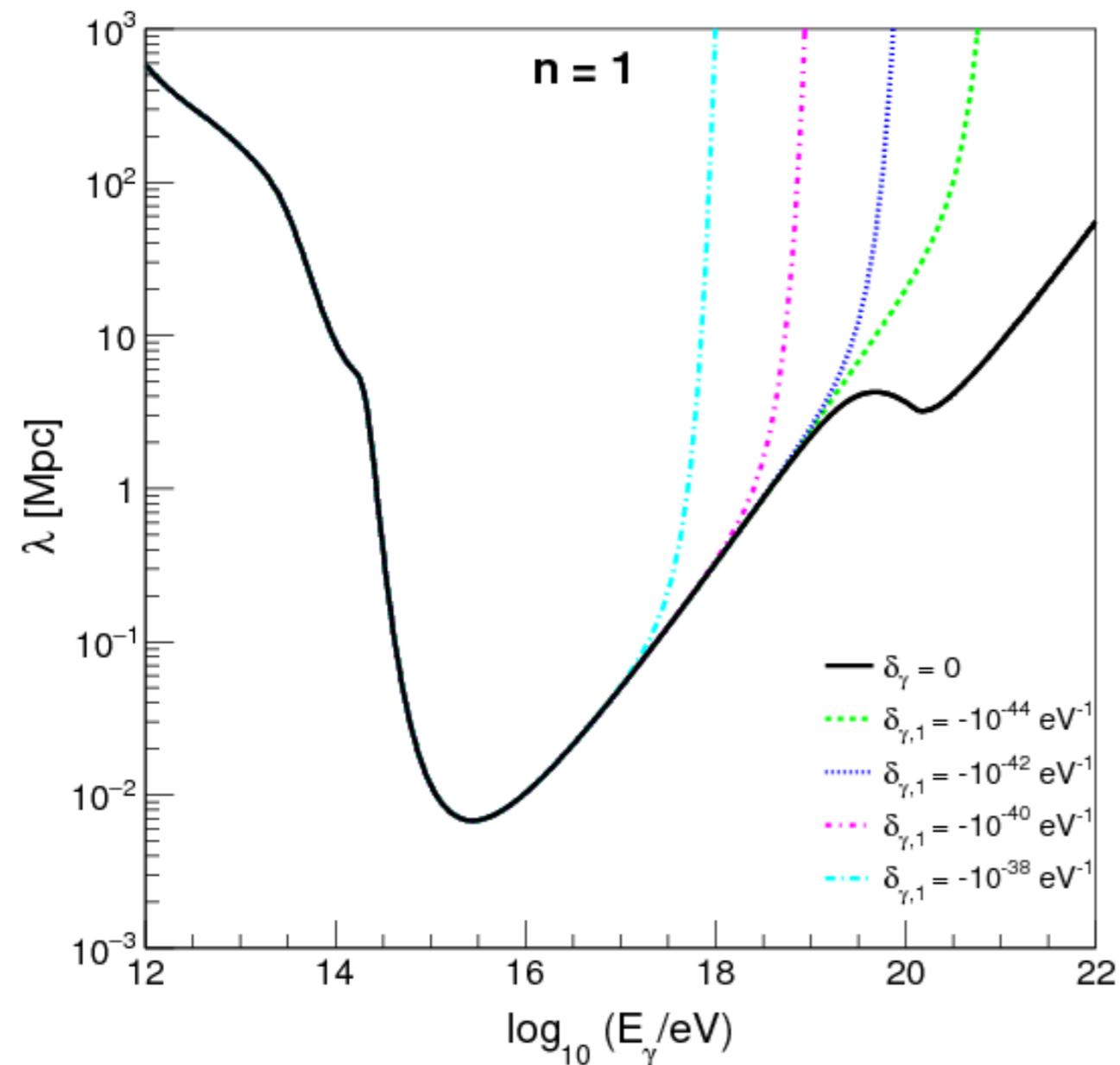
$$\times \int_{\epsilon_{th}^{LIV}}^\infty d\epsilon n_\gamma(\epsilon, z) \times \int_{-1}^1 d(\cos \theta) \frac{1 - \cos \theta}{2} \sigma(E_\gamma, \epsilon, z, \cos \theta)$$

↑
LIV

$$\epsilon_{th}^{LIV} = \frac{m_e^2}{4E_\gamma K(1-K)} - \frac{\delta_{\gamma,n} E_\gamma^{n+1}}{4}$$

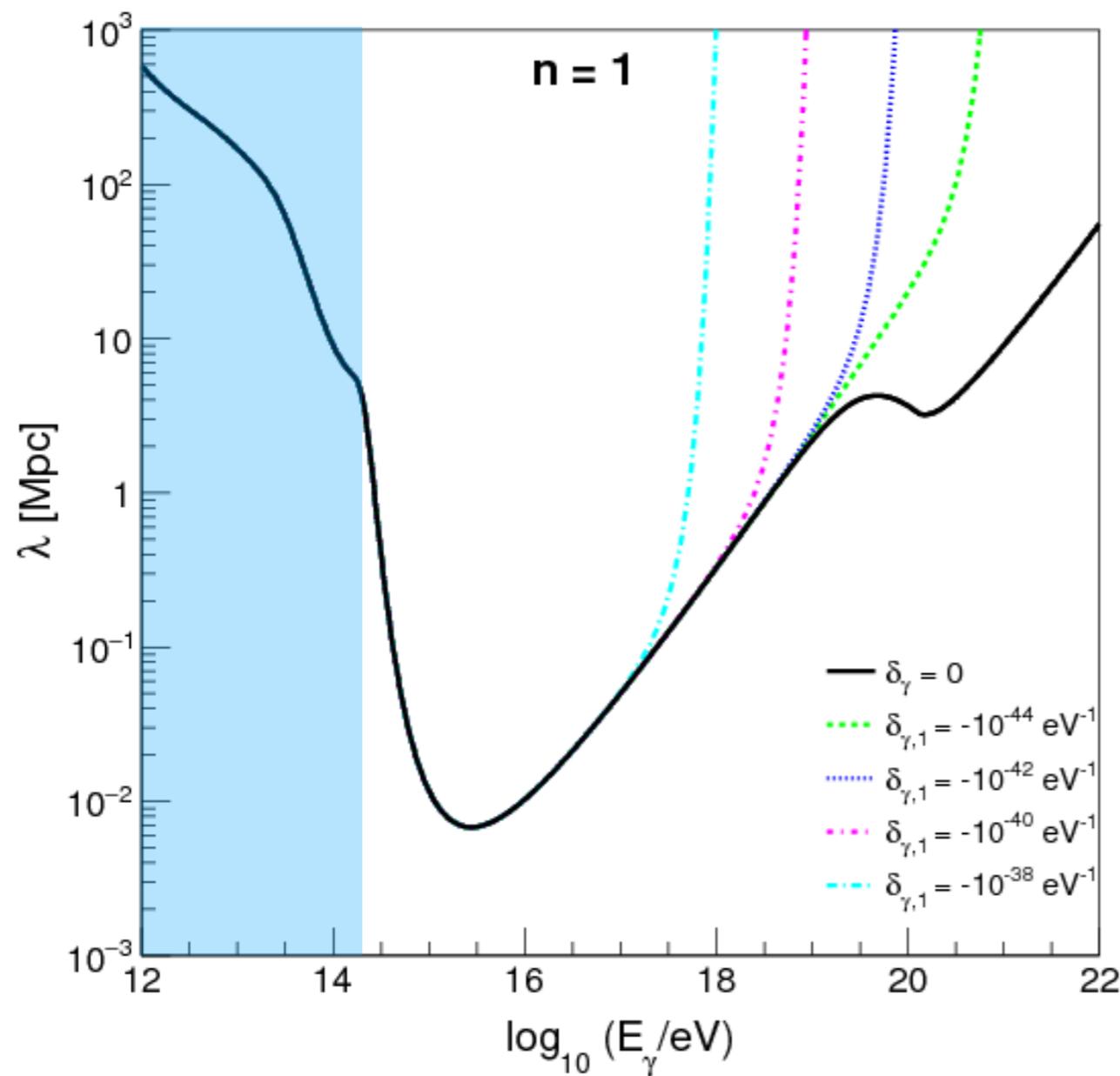
Optical depth + LIV

$$\tau_\gamma(E_\gamma, z, n, E_{LIV}^{(n)}) = \int_0^z dz \frac{c}{H_0(1+z)\sqrt{\Omega_\Lambda + \Omega_M(1+z)^3}} \int_{\epsilon_{th}^{LIV}}^\infty d\epsilon n_\gamma(\epsilon, z) \int_{-1}^1 d(\cos\theta) \frac{1 - \cos\theta}{2} \sigma(E_\gamma, \epsilon, z, \cos\theta)$$

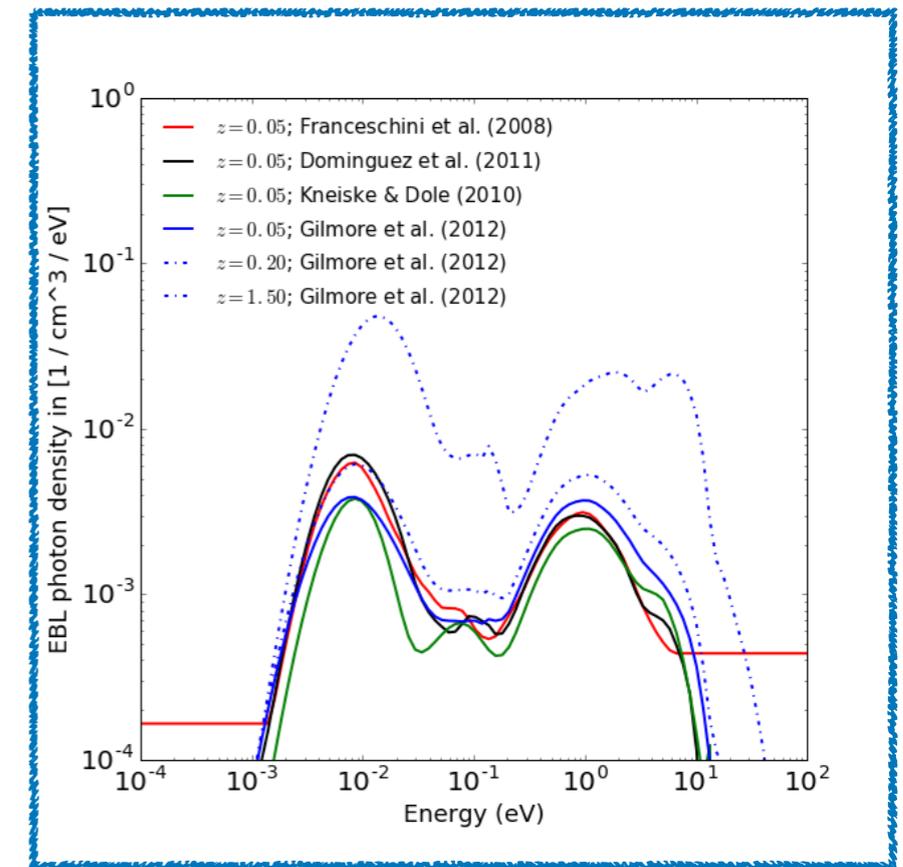


Optical depth + LIV

$$\tau_\gamma(E_\gamma, z, n, E_{LIV}^{(n)}) = \int_0^z dz \frac{c}{H_0(1+z)\sqrt{\Omega_\Lambda + \Omega_M(1+z)^3}} \int_{\epsilon_{th}^{LIV}}^\infty d\epsilon n_\gamma(\epsilon, z) \int_{-1}^1 d(\cos\theta) \frac{1 - \cos\theta}{2} \sigma(E_\gamma, \epsilon, z, \cos\theta)$$



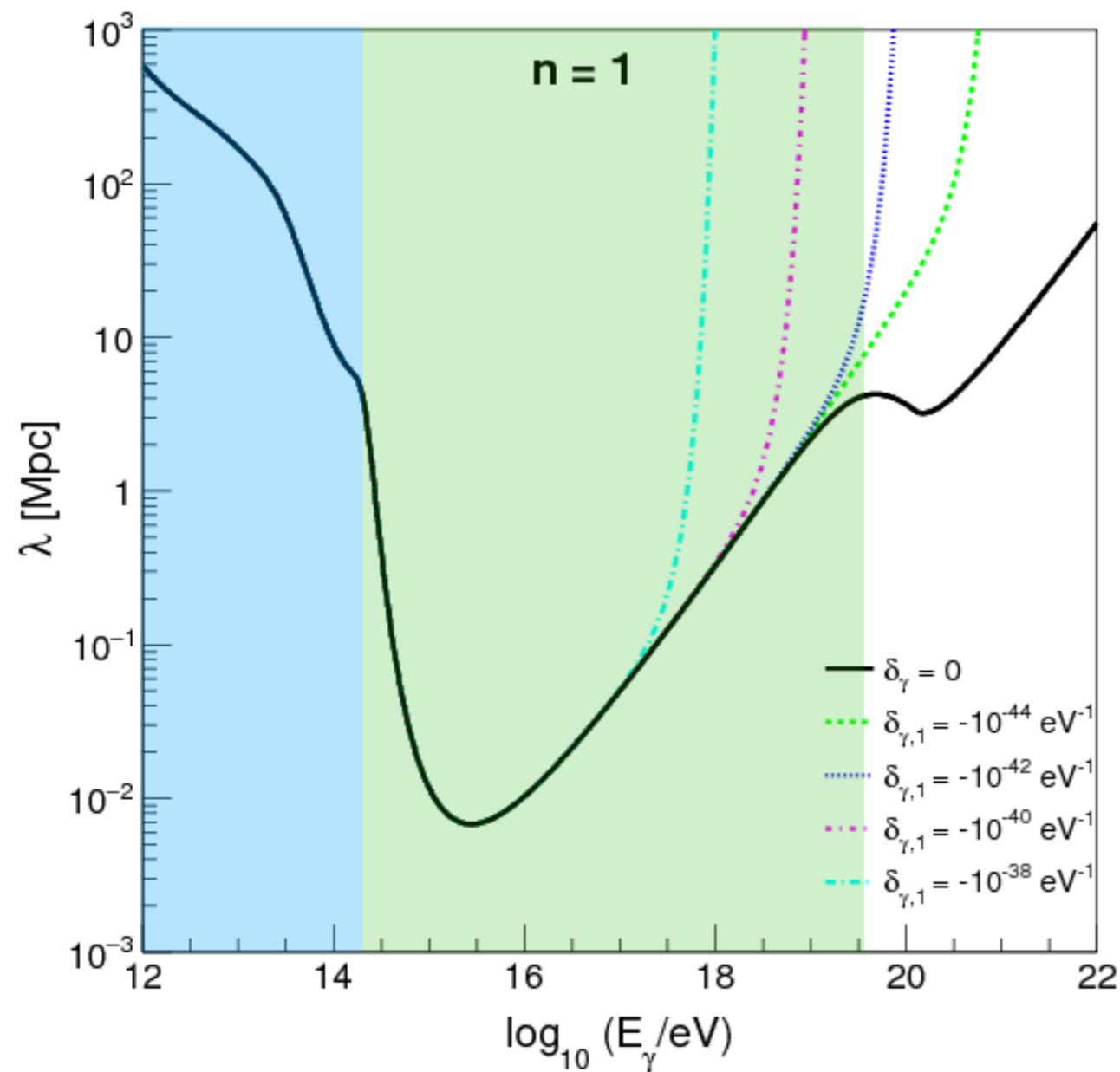
BKG density
EBL-photons



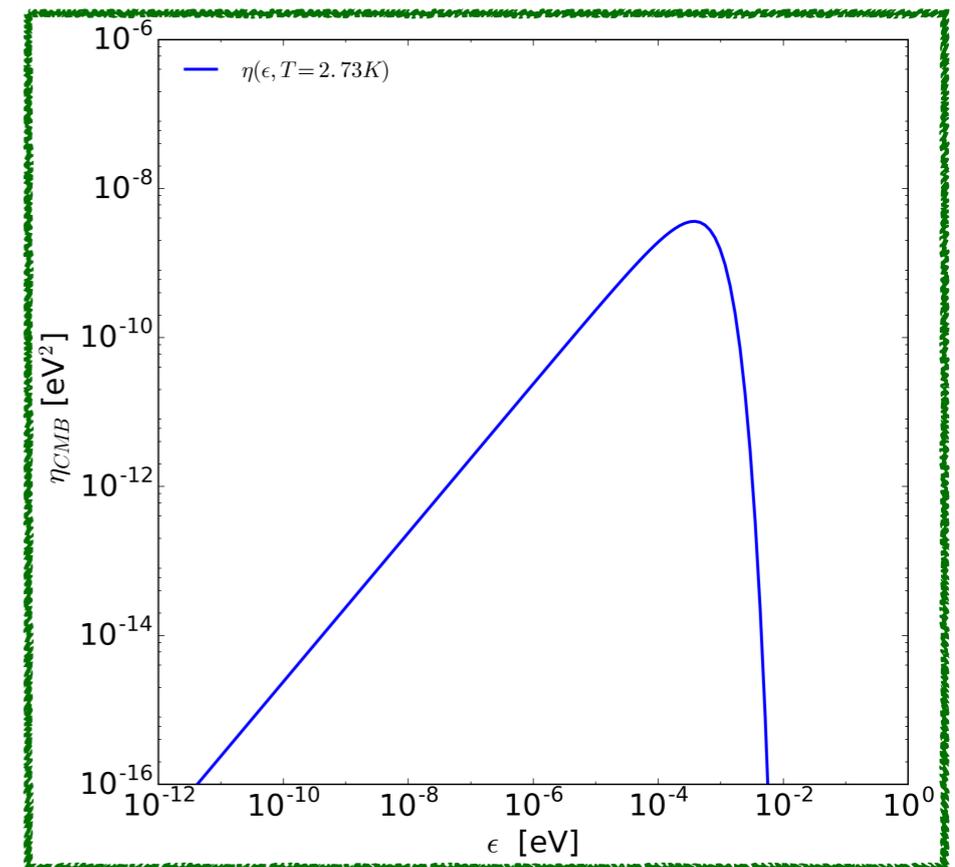
EBL: Gilmore & Ramirez-Ruiz (2010)

Optical depth + LIV

$$\tau_\gamma(E_\gamma, z, n, E_{LIV}^{(n)}) = \int_0^z dz \frac{c}{H_0(1+z)\sqrt{\Omega_\Lambda + \Omega_M(1+z)^3}} \int_{\epsilon_{th}^{LIV}}^\infty d\epsilon n_\gamma(\epsilon, z) \int_{-1}^1 d(\cos\theta) \frac{1 - \cos\theta}{2} \sigma(E_\gamma, \epsilon, z, \cos\theta)$$

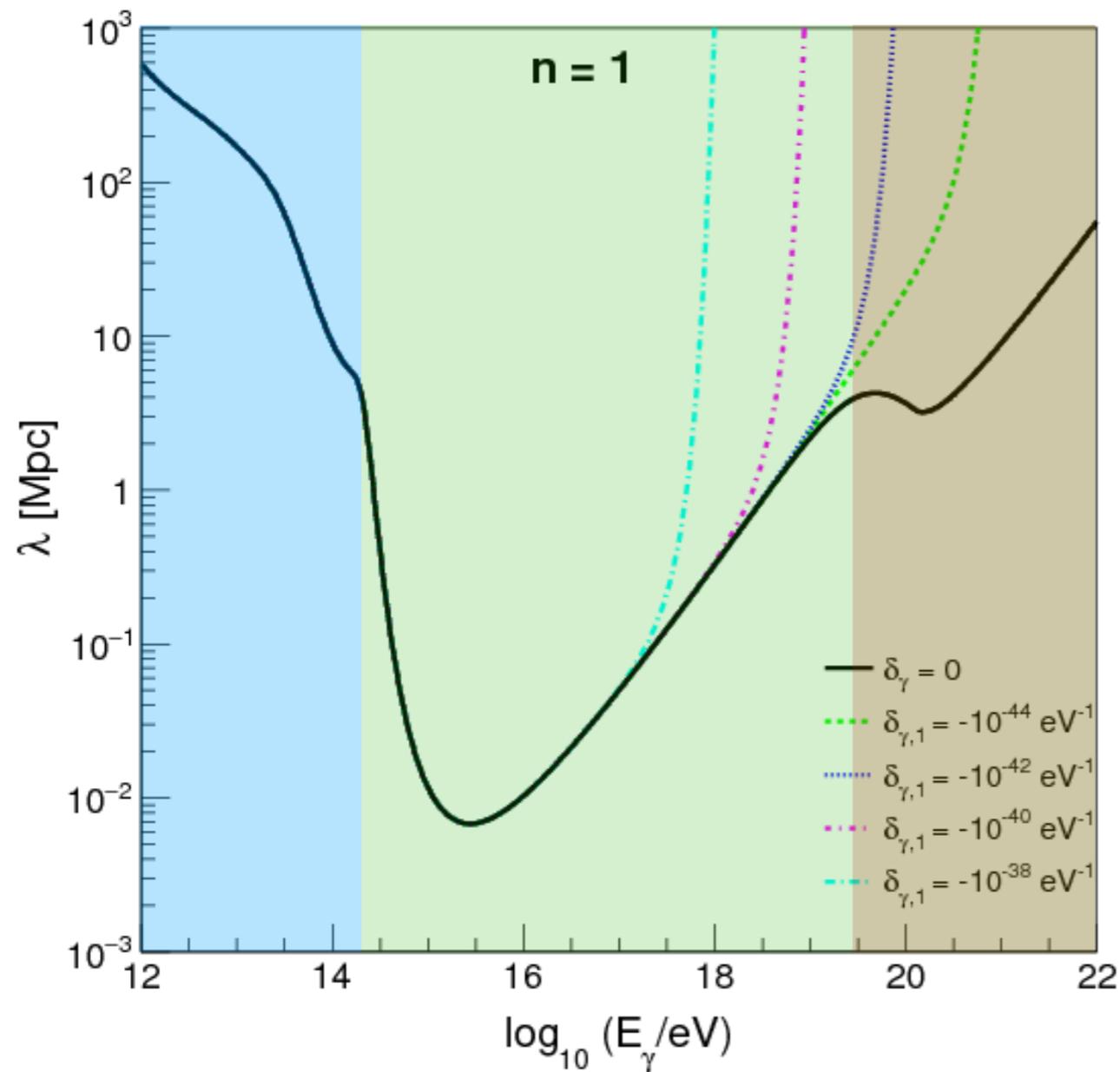


BKG density
CMB-photons



Optical depth + LIV

$$\tau_\gamma(E_\gamma, z, n, E_{LIV}^{(n)}) = \int_0^z dz \frac{c}{H_0(1+z)\sqrt{\Omega_\Lambda + \Omega_M(1+z)^3}} \int_{\epsilon_{th}^{LIV}}^\infty d\epsilon n_\gamma(\epsilon, z) \int_{-1}^1 d(\cos\theta) \frac{1 - \cos\theta}{2} \sigma(E_\gamma, \epsilon, z, \cos\theta)$$



BKG density
Radio-photons

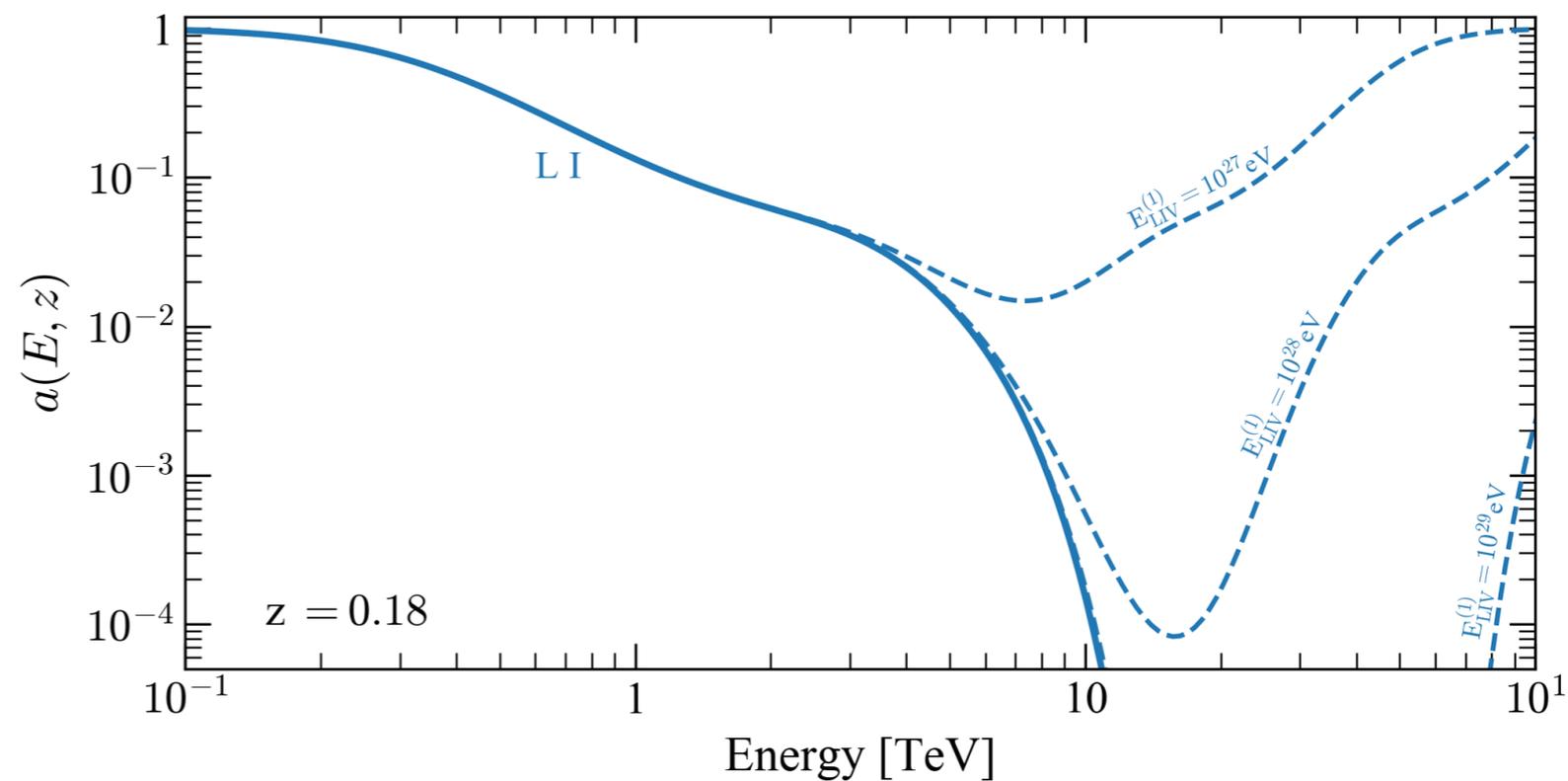
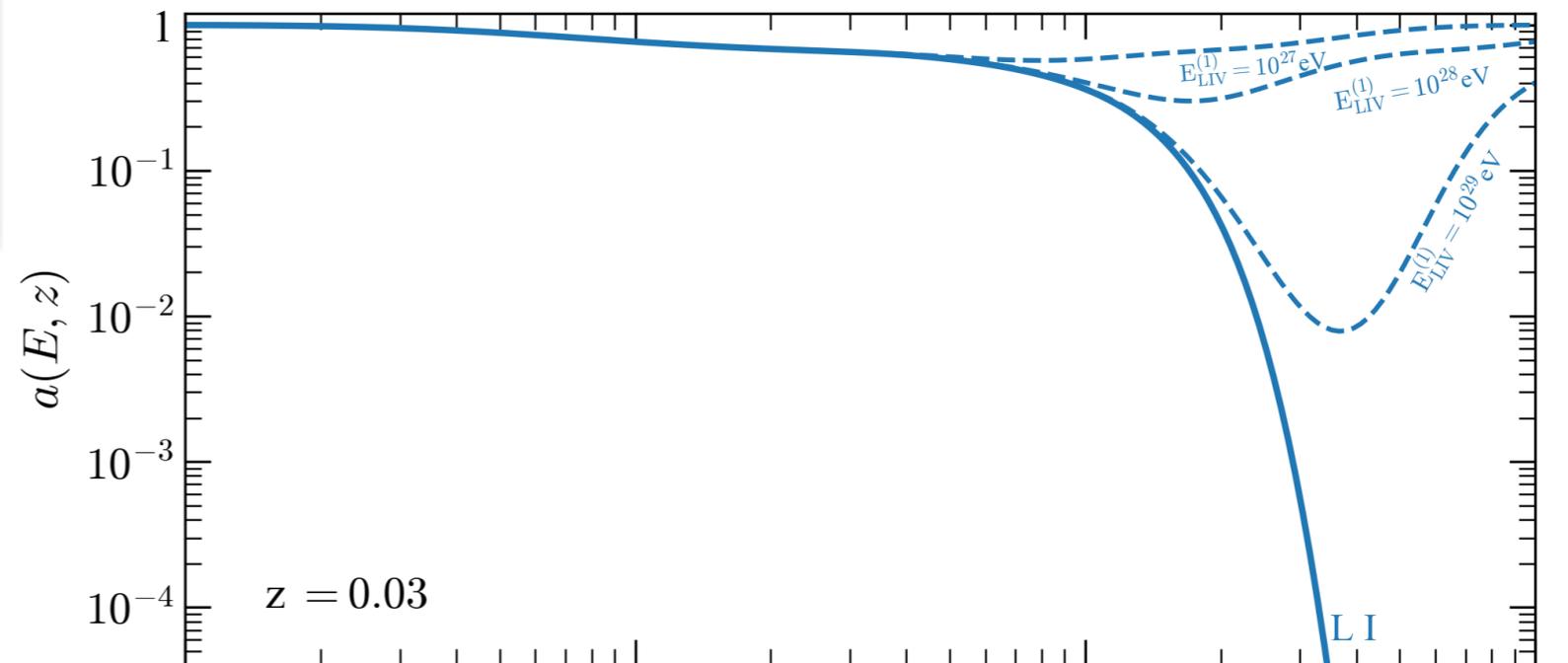
Data from Gervasi et al.
(2008)

EBL-Attenuation + LIV

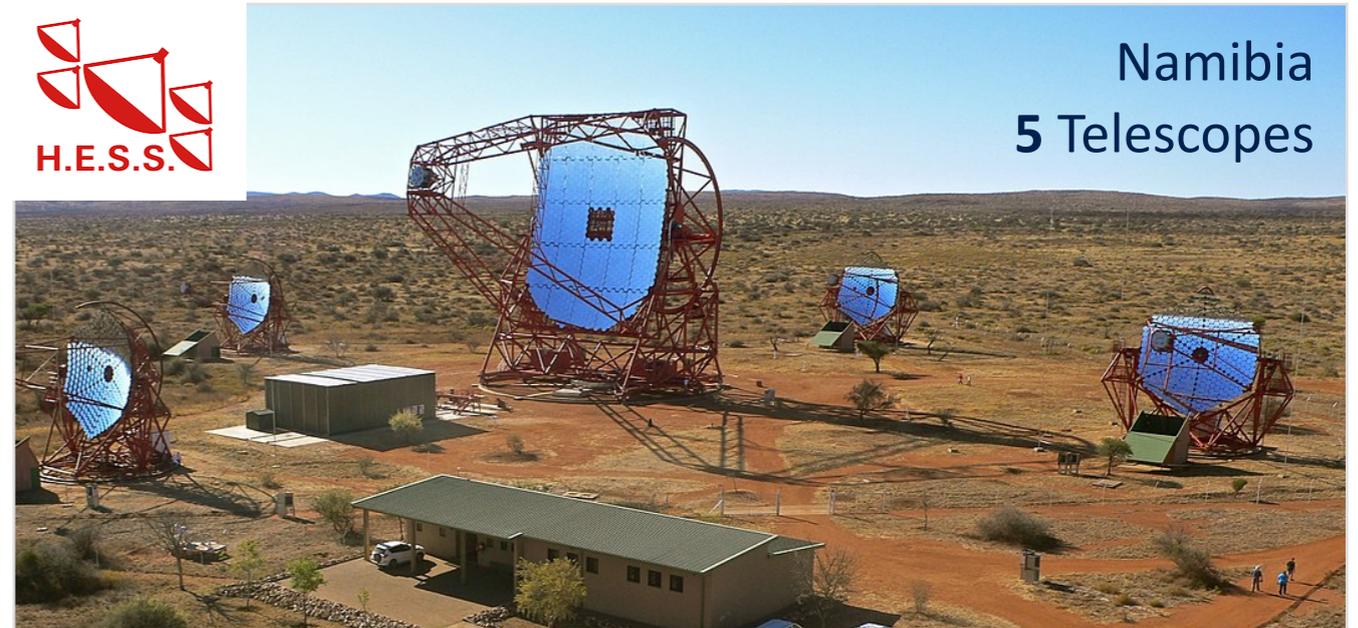
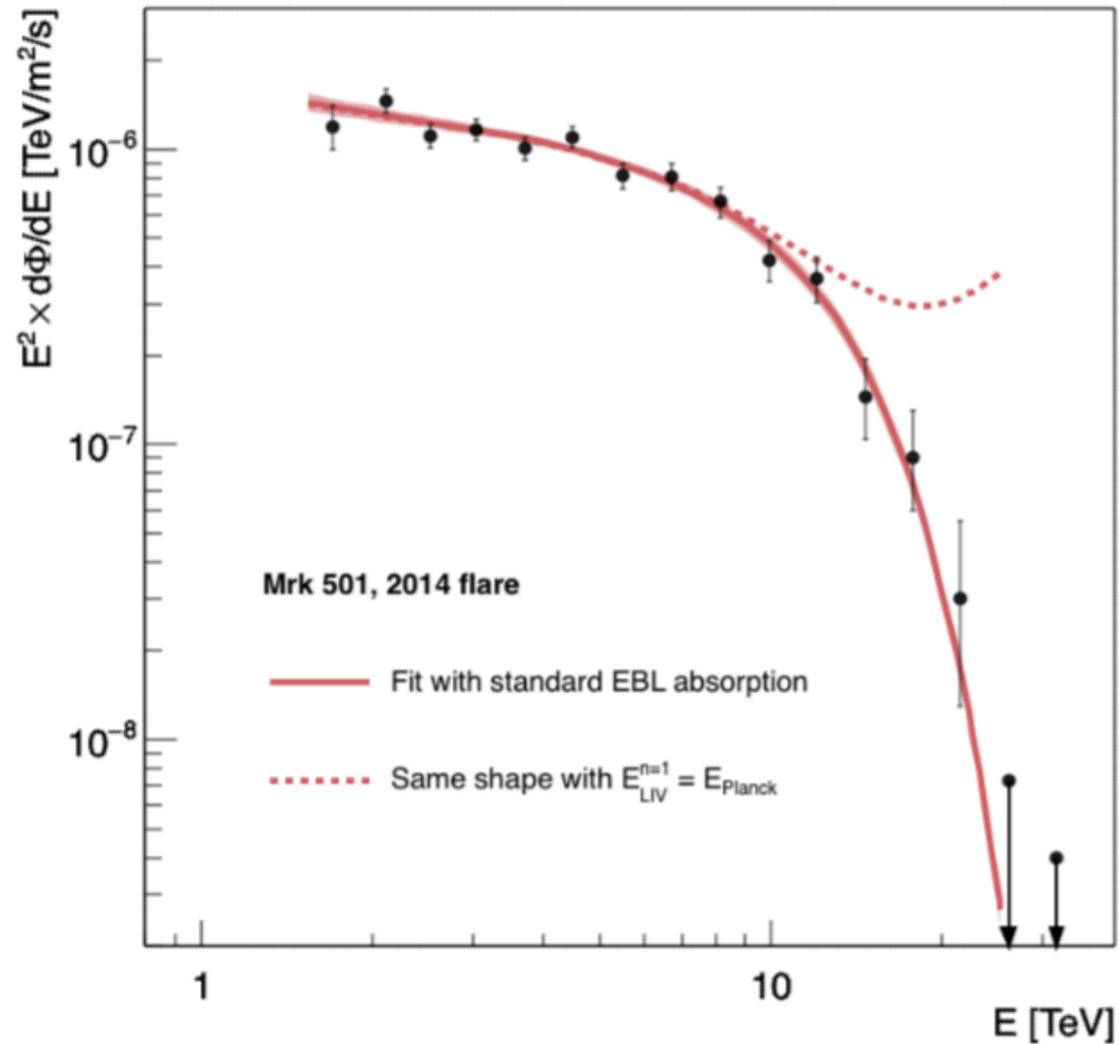
$$a(E, z) = e^{-\tau(E, z)}$$

The intensity of the LIV effect depends on

- ▶ E_γ :
The energy of the γ -ray
- ▶ E_{LIV} :
The LIV energy scale
- ▶ z :
The distance of the source.



EBL-Attenuation + LIV



	2σ	3σ	5σ
n=1	2.8×10^{28} eV ($2.29 \times E_{Planck}$)	1.9×10^{28} eV ($1.6 \times E_{Planck}$)	1.04×10^{28} eV ($0.86 \times E_{Planck}$)
n=2	7.5×10^{20} eV	6.4×10^{20} eV	4.7×10^{20} eV

Lorentz and Brun for the HESS collaboration, RICAP16, 2016.

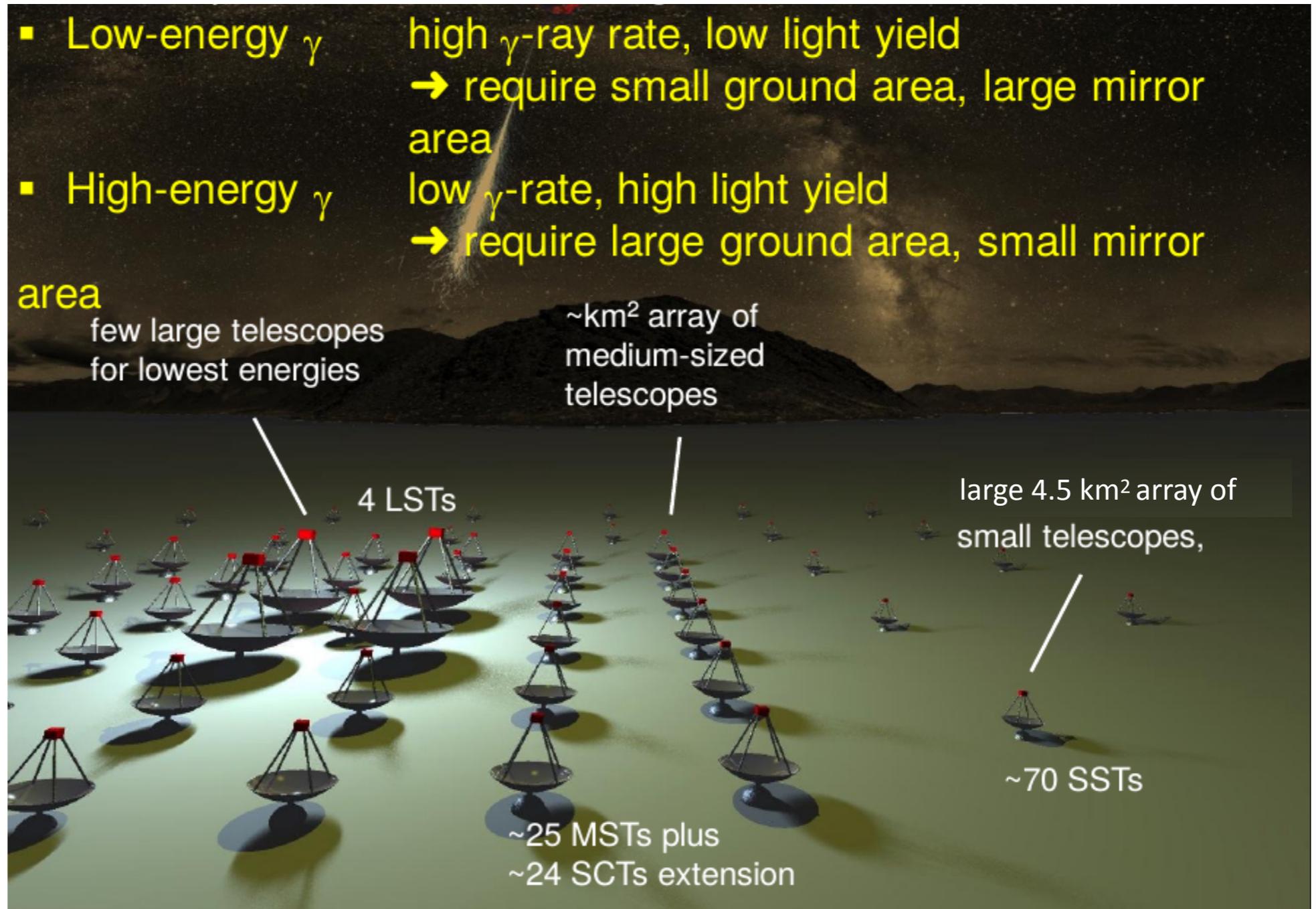
Cherenkov Telescope Array



99 Telescopes

- 4 LST
- 25 MST
- 70 SSTs

Energy range
20 GeV - 300 TeV





31 Countries
92 Parties
202 Institutes
1466 members (513 FTE)

UNAM



- 7 Scientist
 - Ruben Alfaro
 - Alejandro Lara
 - William Lee
 - Maria Magdalena
 - Lukas Nellen
 - Andrés Sandoval
 - Gagik Tovmassian
- 2 Engineers
 - Fernando Garfias
 - Arturo Iriarte



CTA Brazil

12 Institutions

25 Scientists

16 Students

5 Technicians

CTA - SP - MST

- **IFSC- USP**
 - Prof. Vitor de Souza
 - Profa. Manuela Vecchi
 - Profa. Cibelle Celestino
 - Dr. Humberto Huerta
 - Dr. Aion Viana
 - Edyvania Martins
 - Rodrigo Lang
 - Luan Arbeletche
 - Andres Delgado
 - Rodrigo Guedes Lang
 - Danielle Kaori
- **IF-USP**
 - Prof. Edivaldo Moura
 - Douglas Pimentel
- **UFABC**
 - Prof. Marcelo Leigui
 - Raquel de Almeida
- **UFSCar**
 - Dr. Gustavo Rojas
- **UFPR**
 - Prof. Rita de Cássia
- **EEL / USP**
 - Prof. Fernando Catalani
 - Prof. Carlos Todero
- **SAIFR / IFT - UNESP**
 - Dr. Fabio Iocco
 - Dr Ekaterina Karukes
 - Maria Benito

CTA - SP - SST

- **IAG – USP**
 - Profa. Elisabete dal Pino
 - Prof. Rodrigo Nemmen
 - Dr. Rafael Batistai
 - Dr. Chandra Singh
 - Dr. Grzegorz Kowal
 - Dr. Reinaldo Lima
 - Dr. Paramita Barai
 - Dr. Luis Kadowki
 - Dr. Claudio Melioli
 - Dr. Juan Ramirez
 - Tania Torrejon
 - Renato Gimenes
 - Pankaj Kushwaha
 - Saib Hussain
 - Carlos Fermino
 - Raniere Menezes
 - William Bohórquez
 - Lucas Santos
- **UNICSul**
 - Prof. Anderson Caproni
- **EACH / USP**
 - Prof. Diego Falceta-Gonçalves
 - Mohammad Ali

CTA - Rio

- **CBPF**
 - Prof. Ulisses de Almeida
 - Prof. Ronald Shellard
 - Bruno Arsioli
 - Bernardo Fraga
 - Rodrigo Cardoso
 - Amanda Carvalho



The Array Locations



Array Coordinates

Latitude: 24° 41' 0.34" South
Longitude: 70° 18' 58.84" West



CTA South
Chile, Paranal

~5 km²

area covered by the array of telescopes



CTA North
Spain, La Palma

~0.5 km²

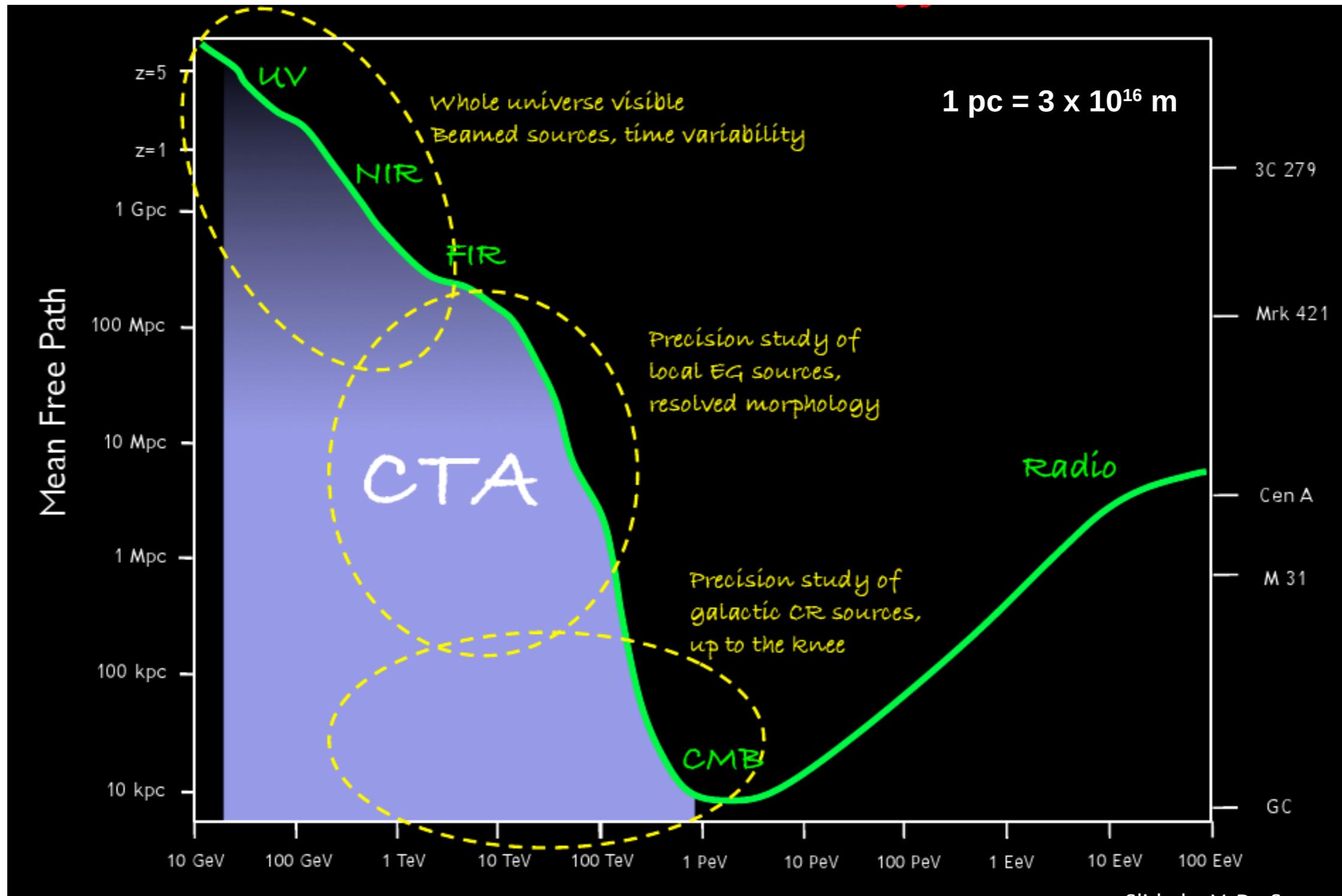
area covered by the array of telescopes



Array Coordinates

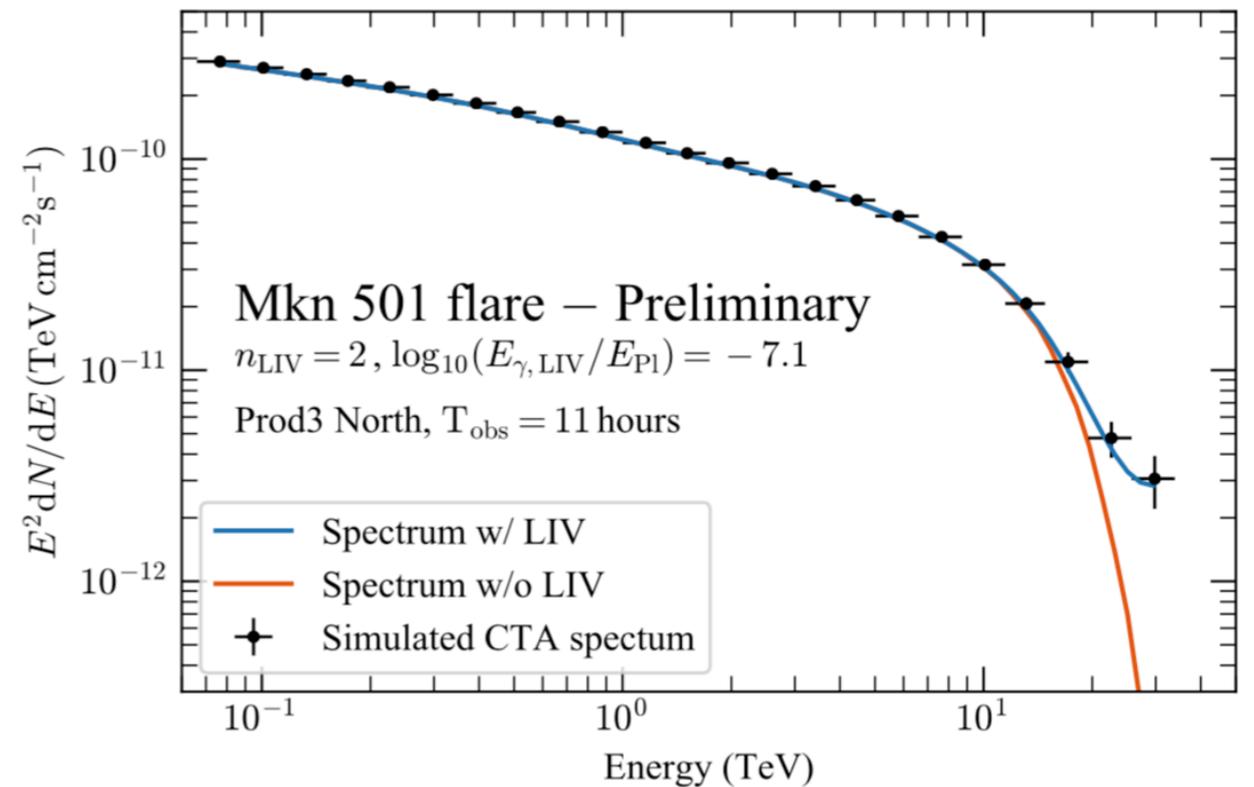
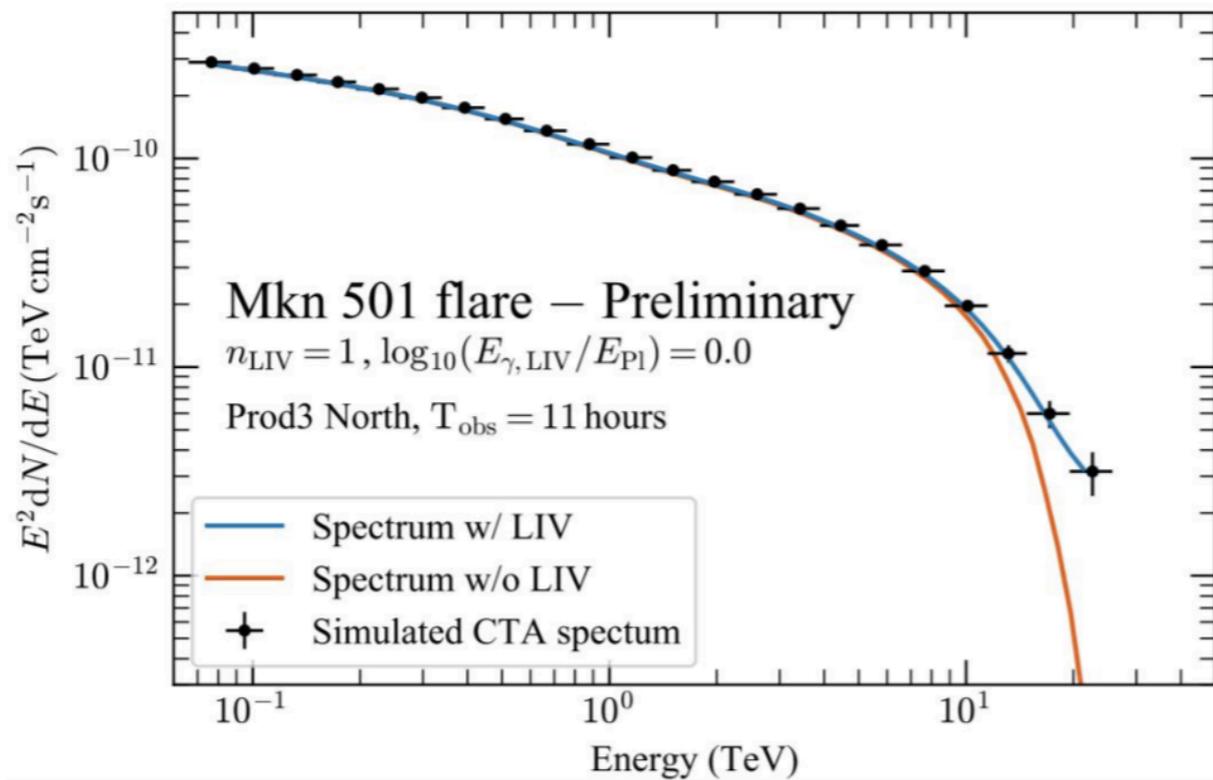
Longitude: 17° 53' 31.218" West
Latitude: 28° 45' 43.7904" North

The γ -ray horizon



Slide by V. De Souza

EBL-Attenuation + LIV



LIV TeV Horizon

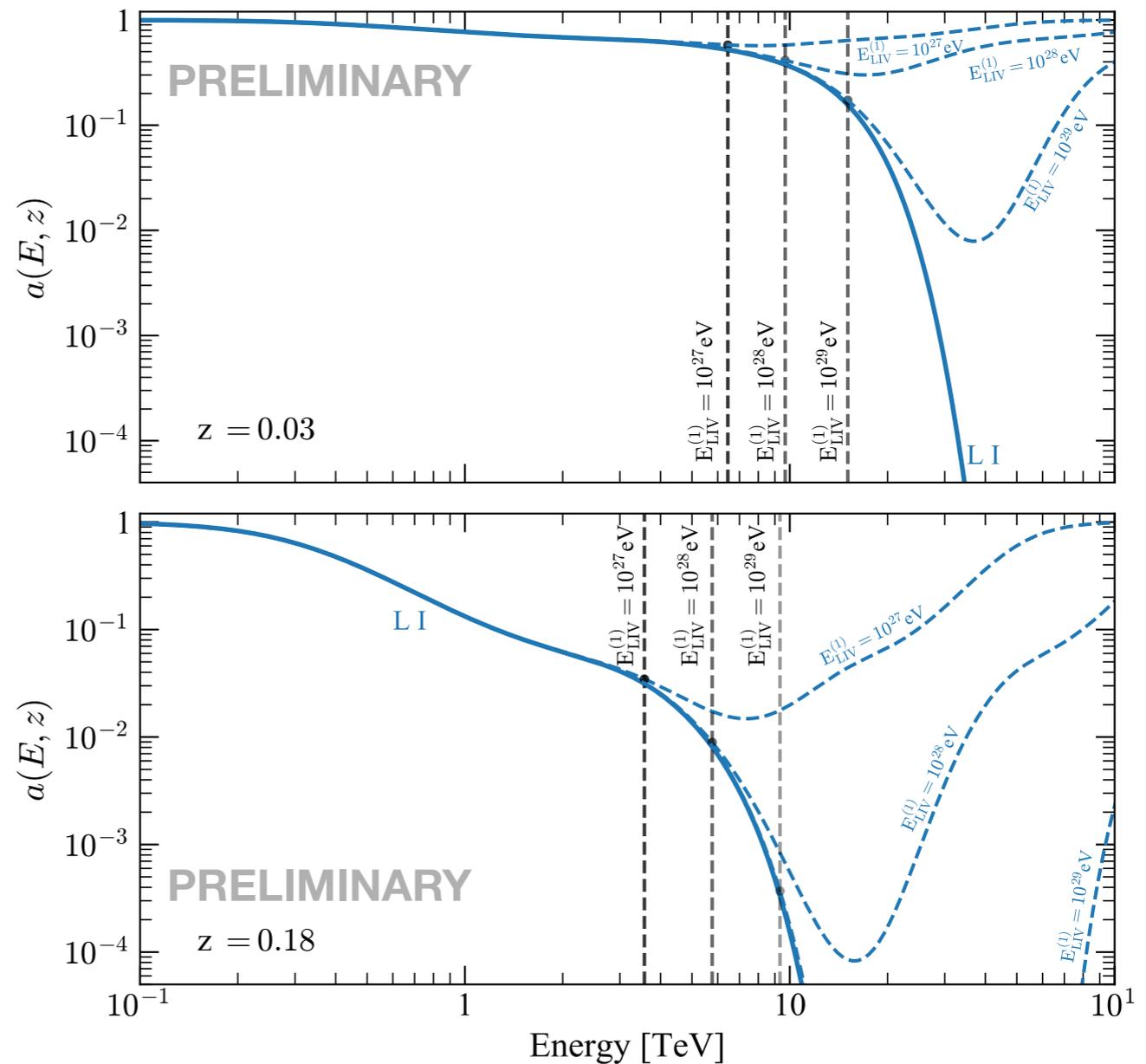
...Why use only one source?

**There are ~111
measured energy
spectra in the
TeVCat !**

LIV TeV Horizon

...Why use only one source?

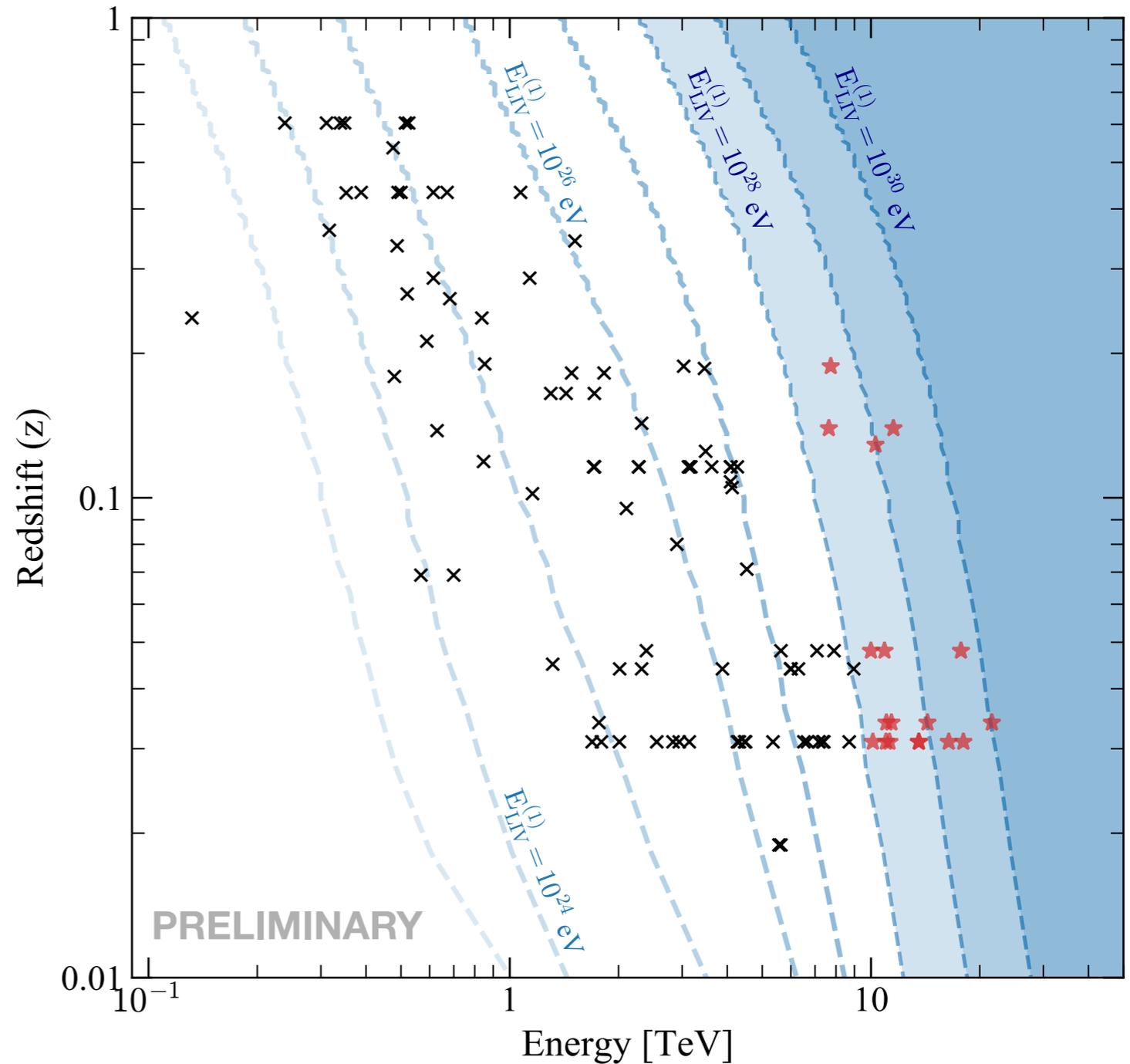
There are ~111
measured energy
spectra in the
TeVCat !



LIV TeV Horizon

...Why use only one source?

There are ~111
measured energy
spectra in the
TeVCat !



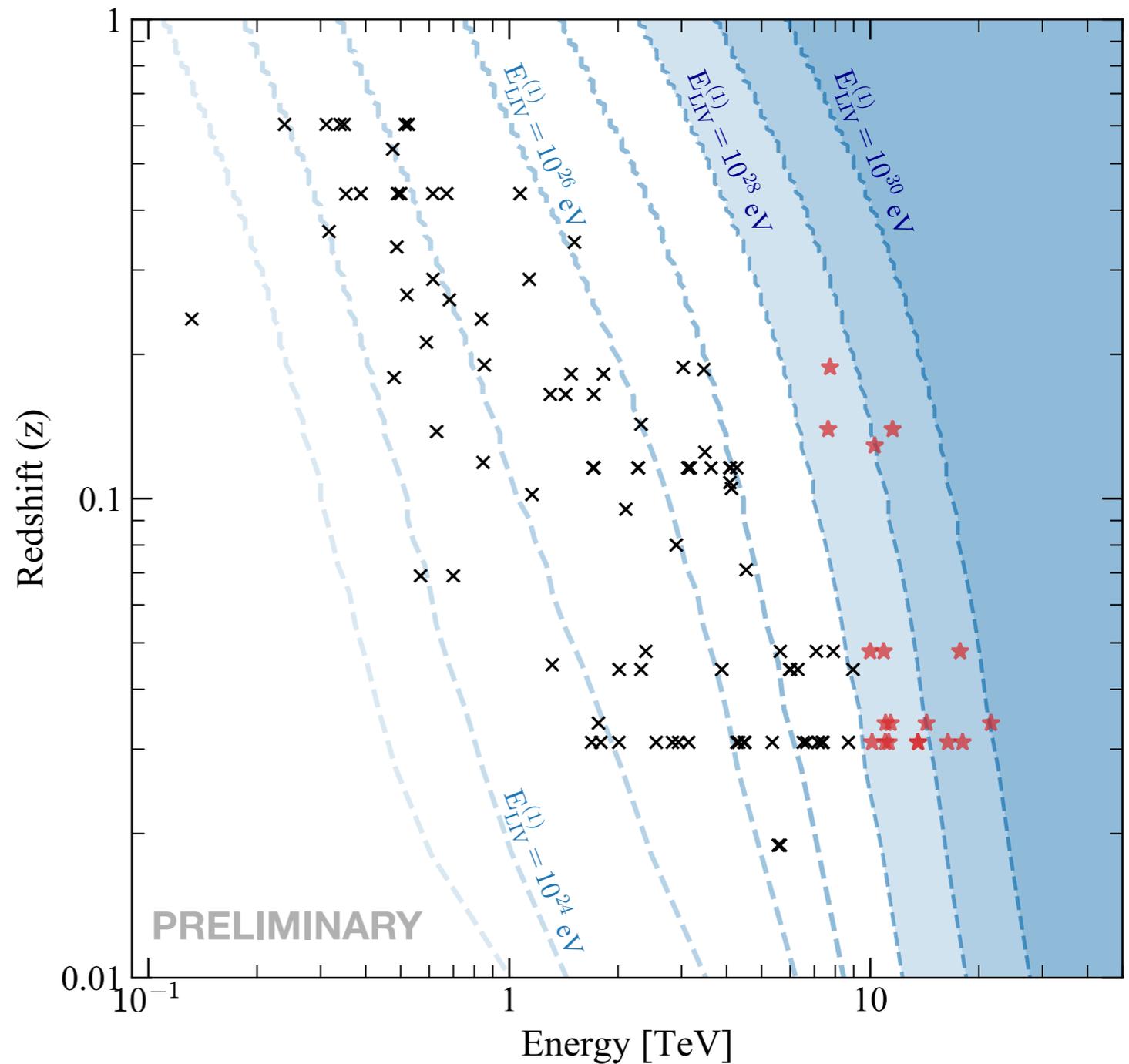
LIV TeV Horizon

...Why use only one source?

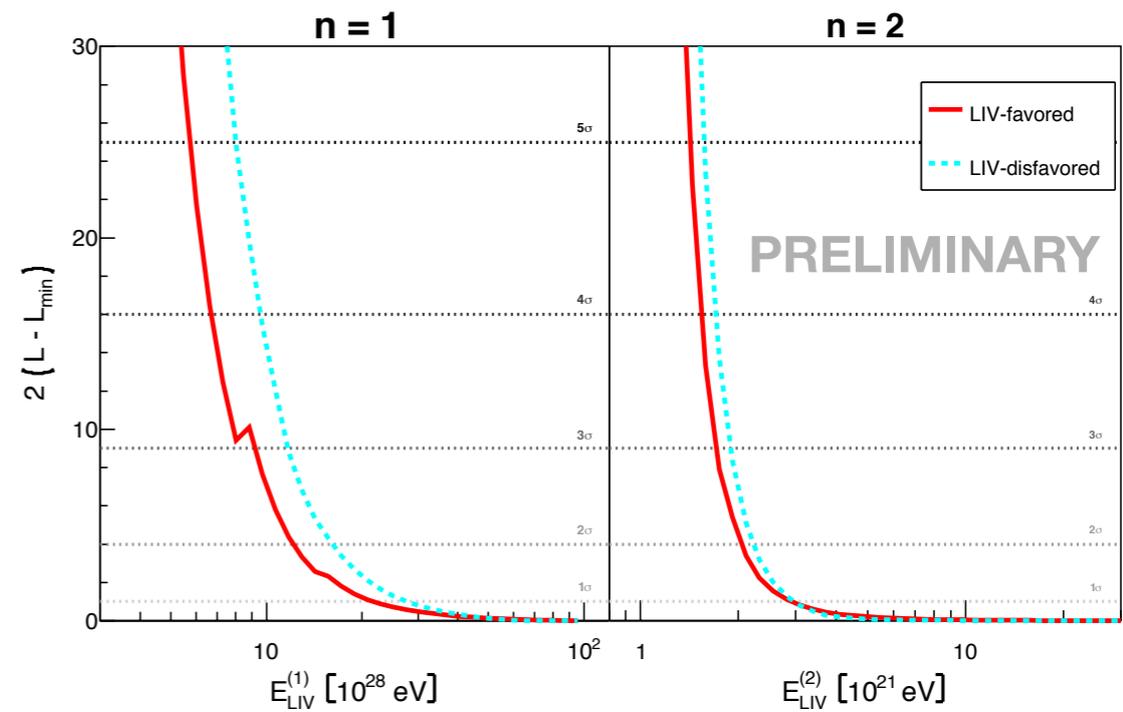
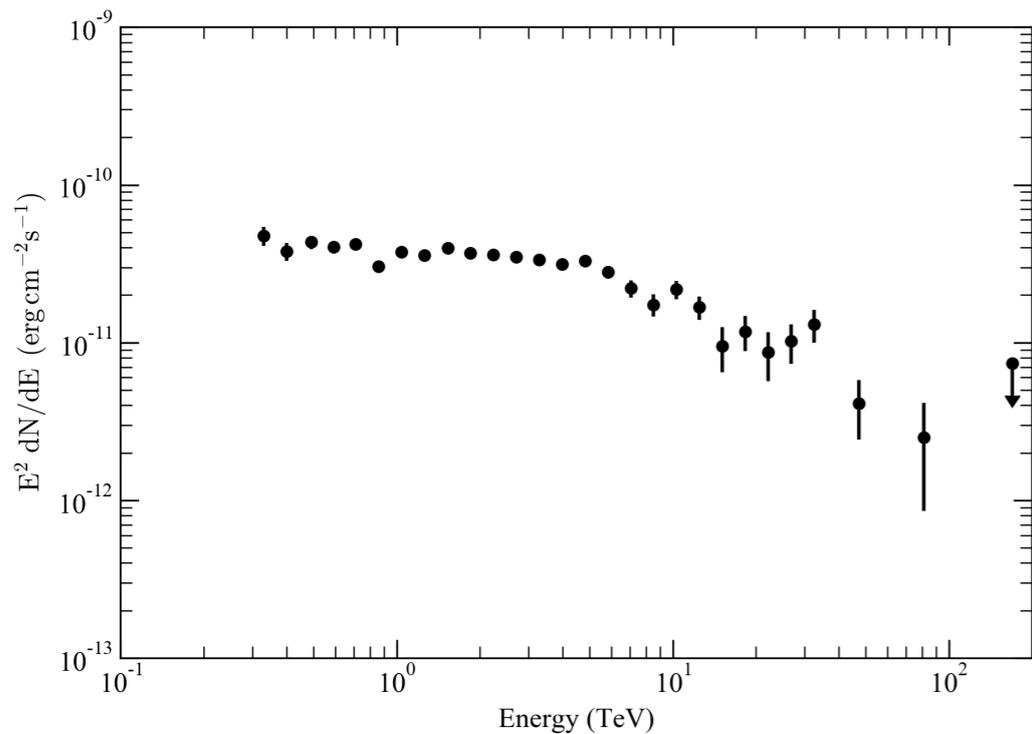
There are ~111
measured energy
spectra in the
TeVCat !



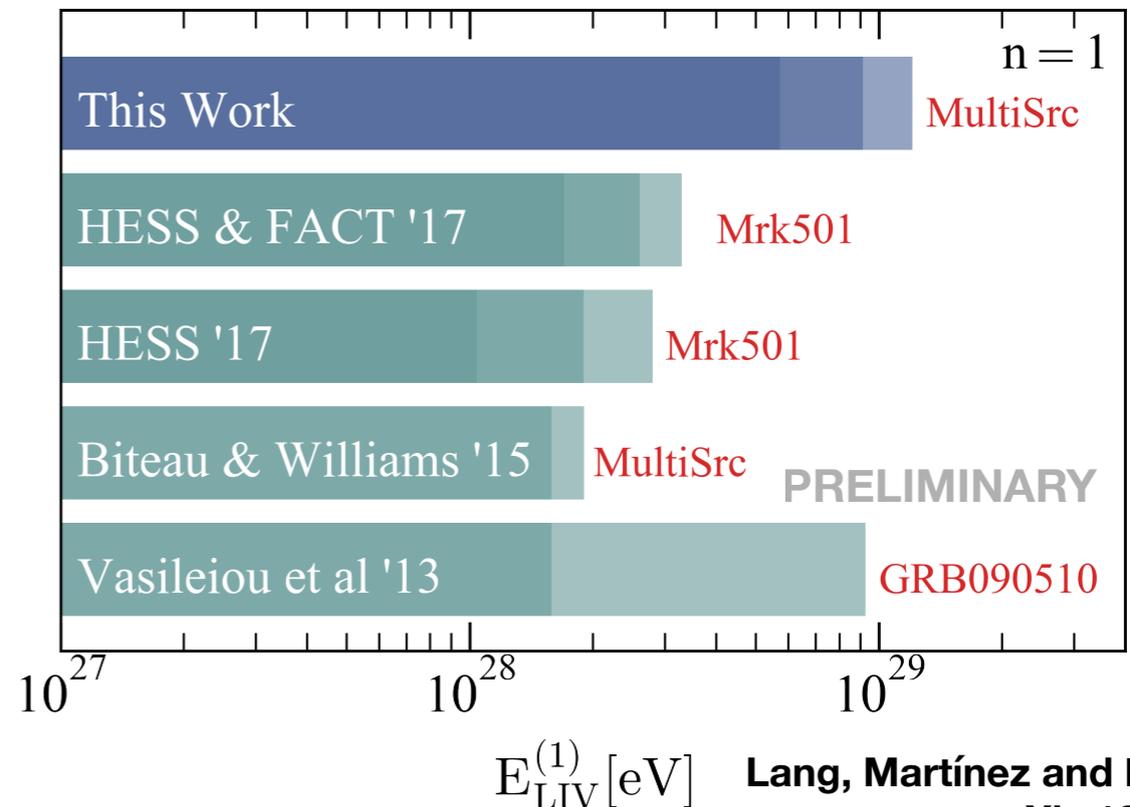
only 18 spectra from 6
sources
have the potential to
show LIV effects
(constrain LIV)



New best LIV limits!

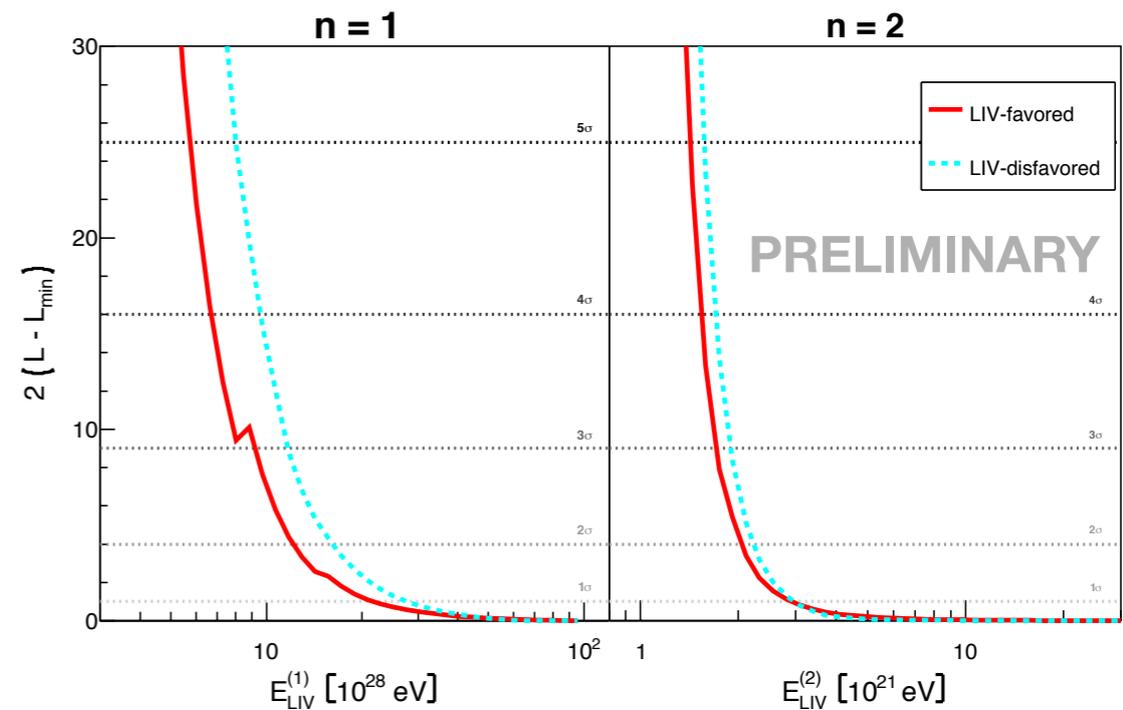
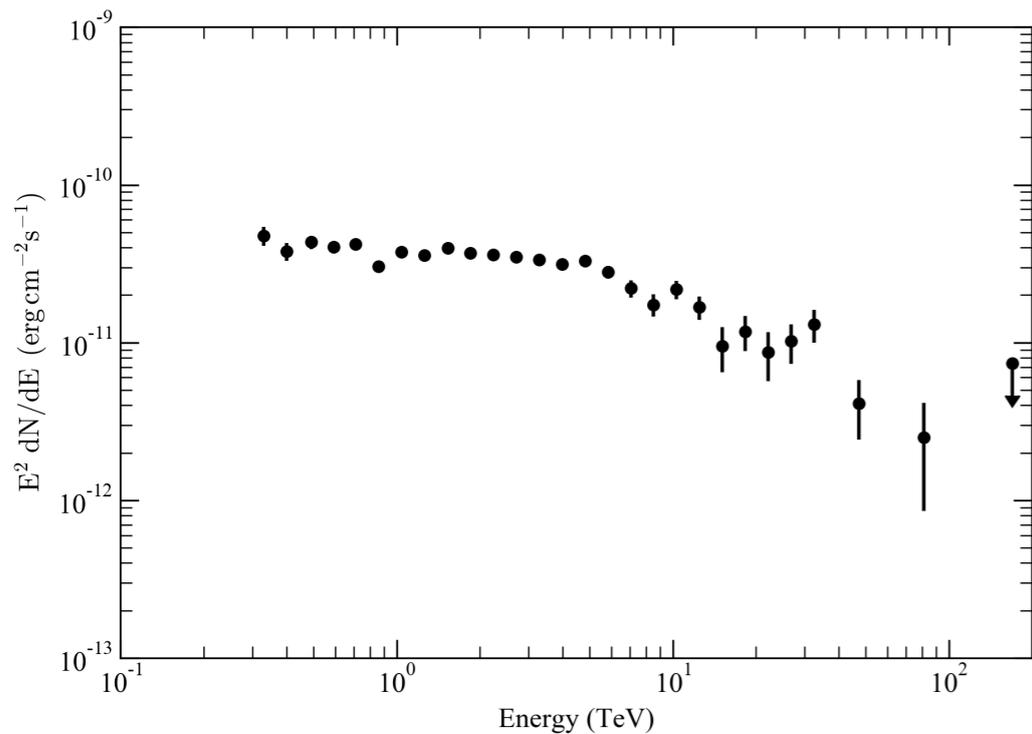


- ▶ Choices of the EBL models
- ▶ Model of the intrinsic spectrum
- ▶ Energy resolution
- ▶ Selection of spectra Selection of energy bins to be used in the calculation of the intrinsic energy spectra

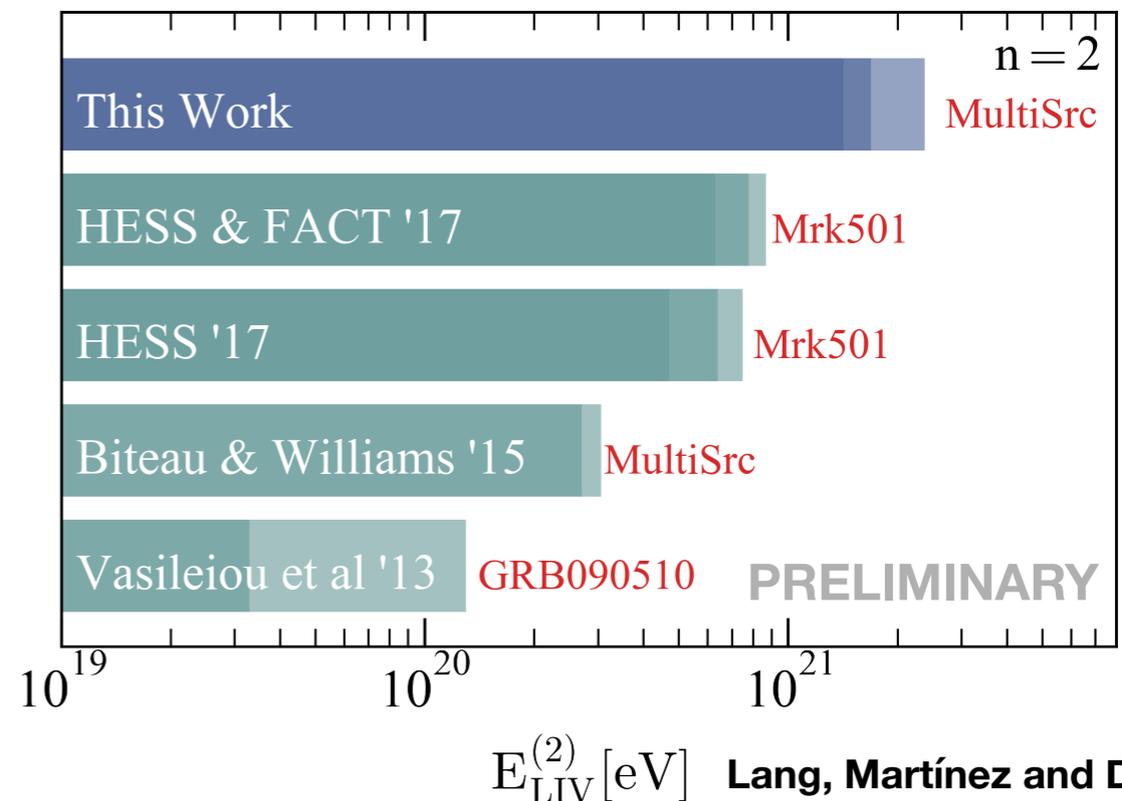


Lang, Martínez and De Souza
arXiv:1810.13215
Submitted

New best LIV limits!

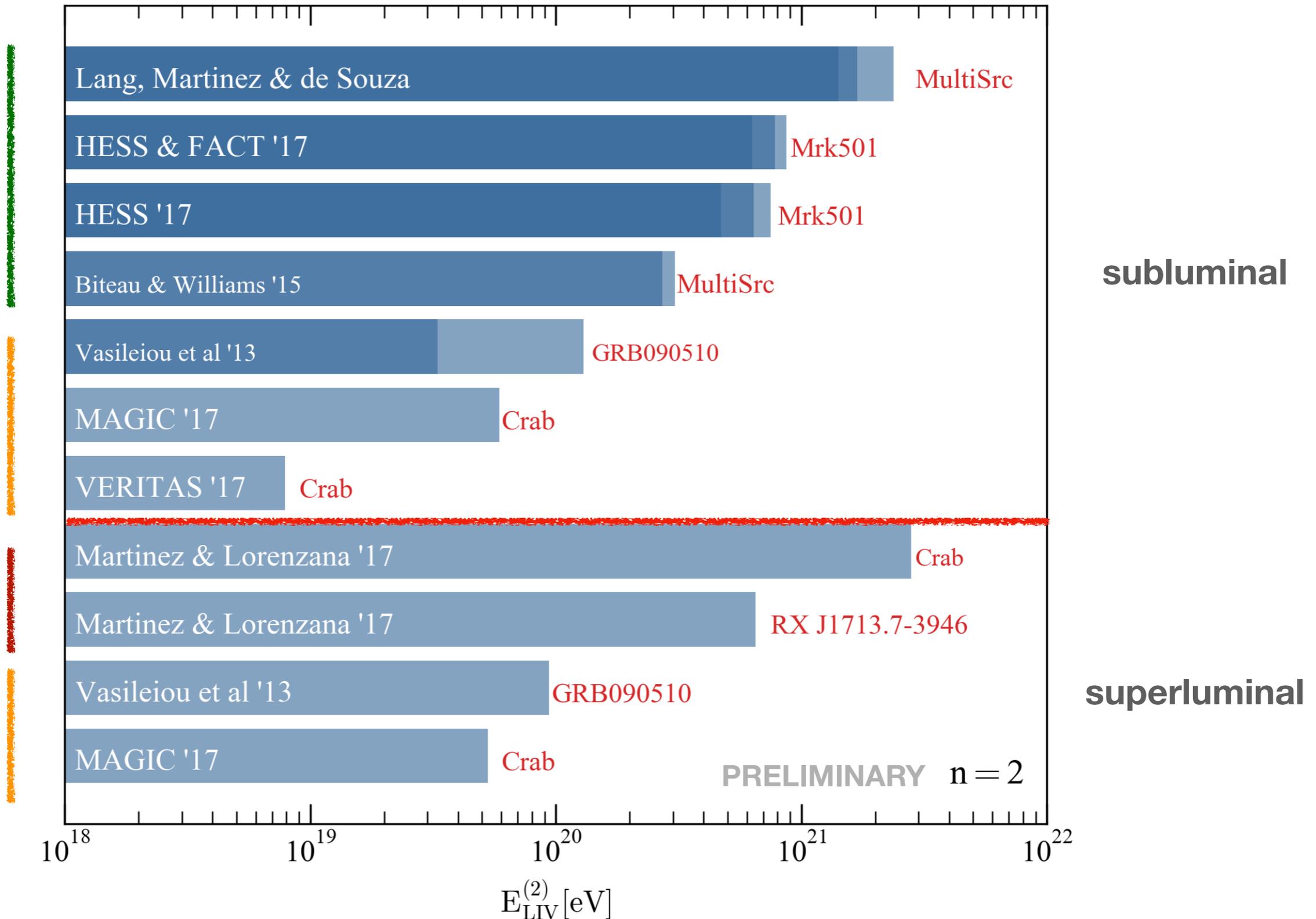


- ▶ Choices of the EBL models
- ▶ Model of the intrinsic spectrum
- ▶ Energy resolution
- ▶ Selection of spectra Selection of energy bins to be used in the calculation of the intrinsic energy spectra

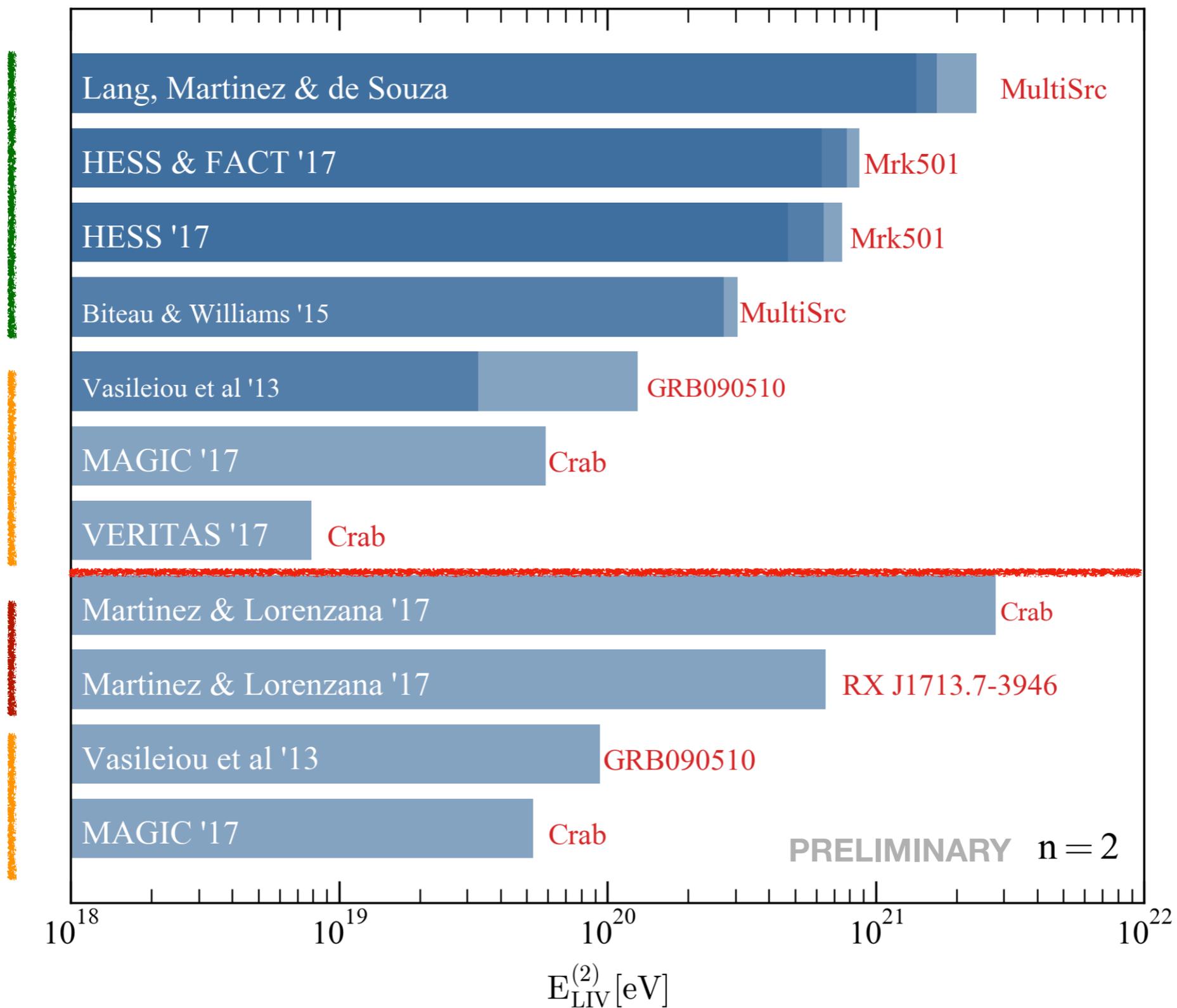


$E_{LIV}^{(2)}$ [eV] Lang, Martínez and De Souza
arXiv:1810.13215
Submitted

LIV limits : γ -rays



LIV limits : γ -rays



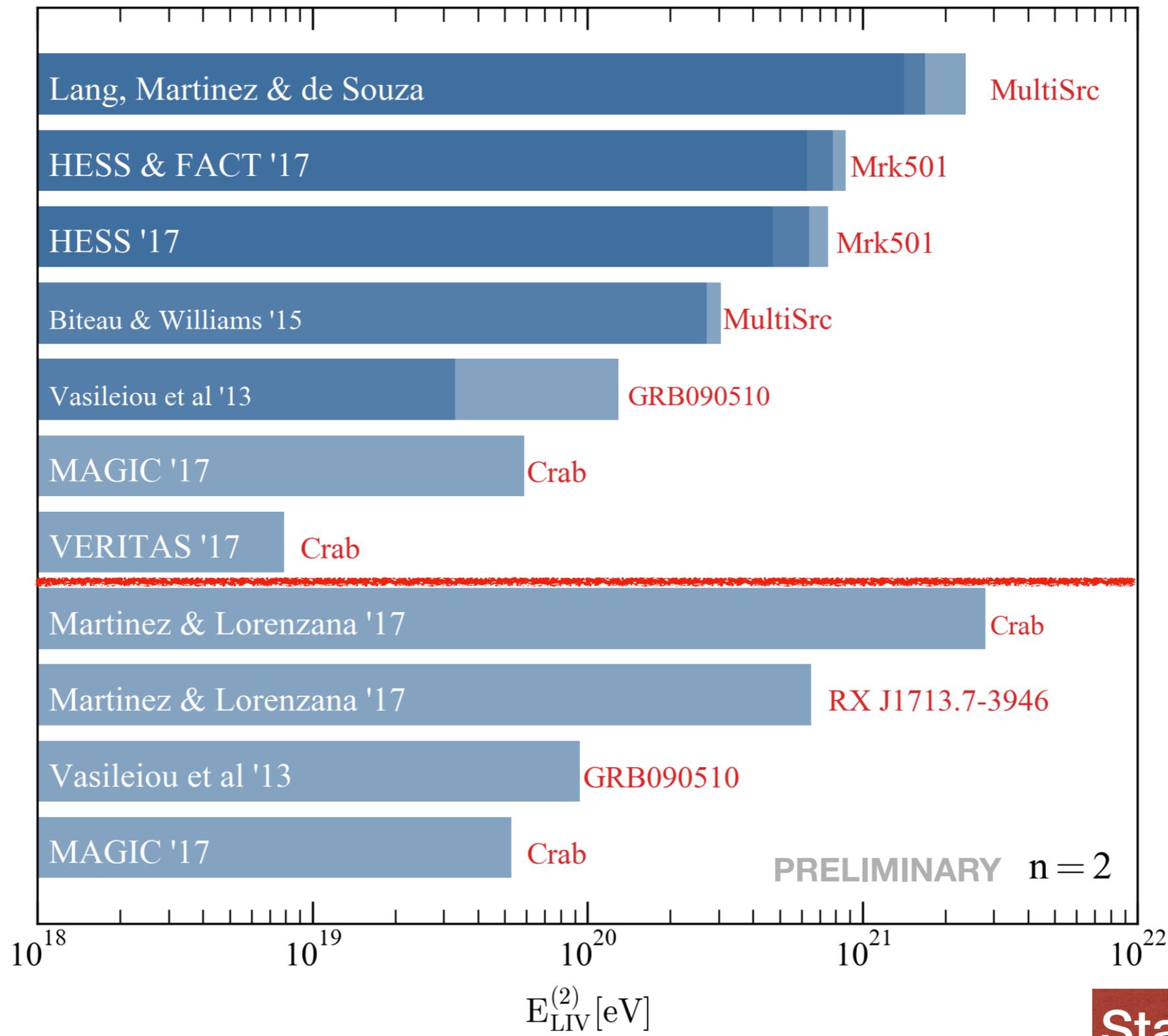
LIV limits : γ -rays

Pair production shift threshold

Time energy dependent delay

Photon decay

Time energy Dependent delay



Stay tuned...

Conclusions and remarks I

- ❖ **Astroparticle physics has recently reached the status of precision science** due to the construction of new observatories, operating innovative technologies and the detection of large numbers of events and sources.
 - ➔ The precise measurements of cosmic and gamma rays can be used as test for fundamental physics, such as the Lorentz invariance violation.
- ❖ We have established **the best limits** to the LIV energy scale to the superluminal and subluminal regime.
- ❖ We developed **a new analysis procedures for searching LIV signatures** using multiple TeV measured energy spectra.
- ❖ We are studying the potential to test / constrain LIV signatures with **HAWC, SGSO and CTA**

- I. Lorentz invariance violation (LIV)

- II. Limits: TeV γ -rays
 - i. Time Energy Dependent delay

 - ii. Photon Decay

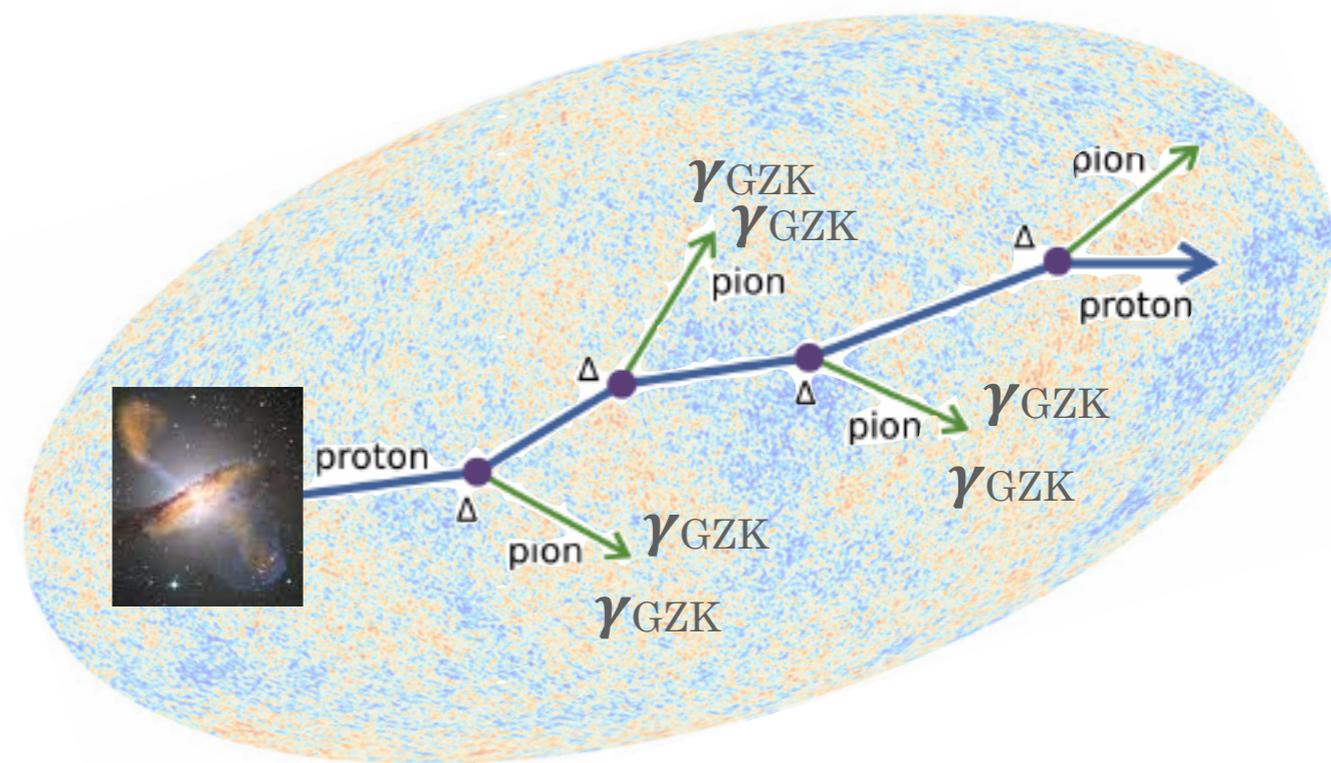
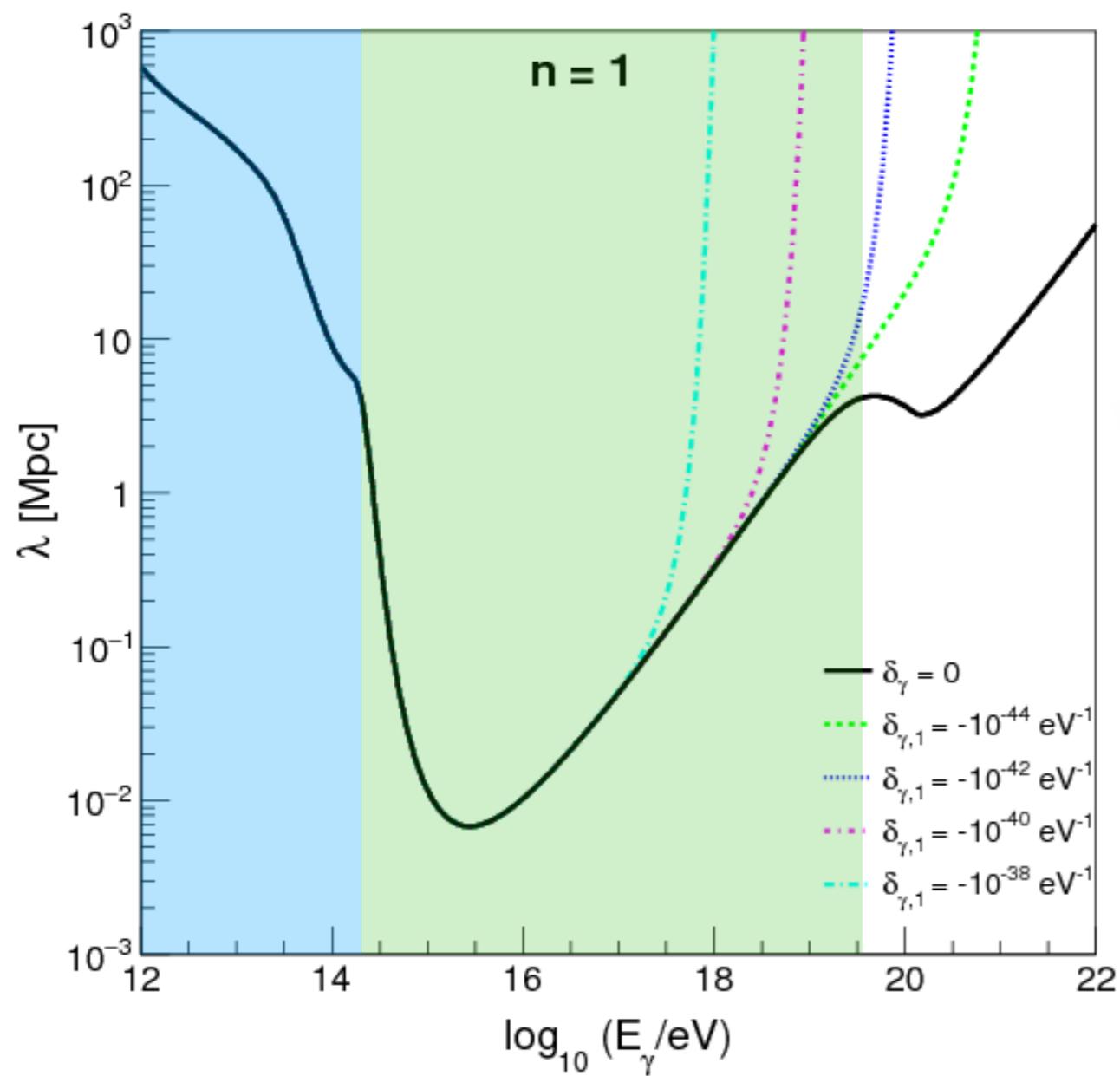
 - iii. Pair production threshold shifts

III. UHECR

- i. GZK-photons + LIV**

- ii. Limits

GZK photons + LIV



Model of UHECR Sources

$$\frac{dN}{dE_s} = \begin{cases} E_s^{-\Gamma}, & \text{for } R_s < R_{\text{cut}} \\ E_s^{-\Gamma} e^{1-R_s/R_{\text{cut}}}, & \text{for } R_s \geq R_{\text{cut}} \end{cases},$$

1. C_1 : Aloisio et al. (2014);
2. C_2 : Unger, Farrar, & Anchordoqui (2015)—Fiducial model (Unger et al. 2015);
3. C_3 : Unger et al. (2015) with the abundance of galactic nuclei from (Olive & Group 2014);
4. C_4 : Berezhinsky, Gazizov, & Grigorieva (2007)—Dip model (Berezhinsky et al. 2006).

Parameters of the Four Source Models Used in This Paper

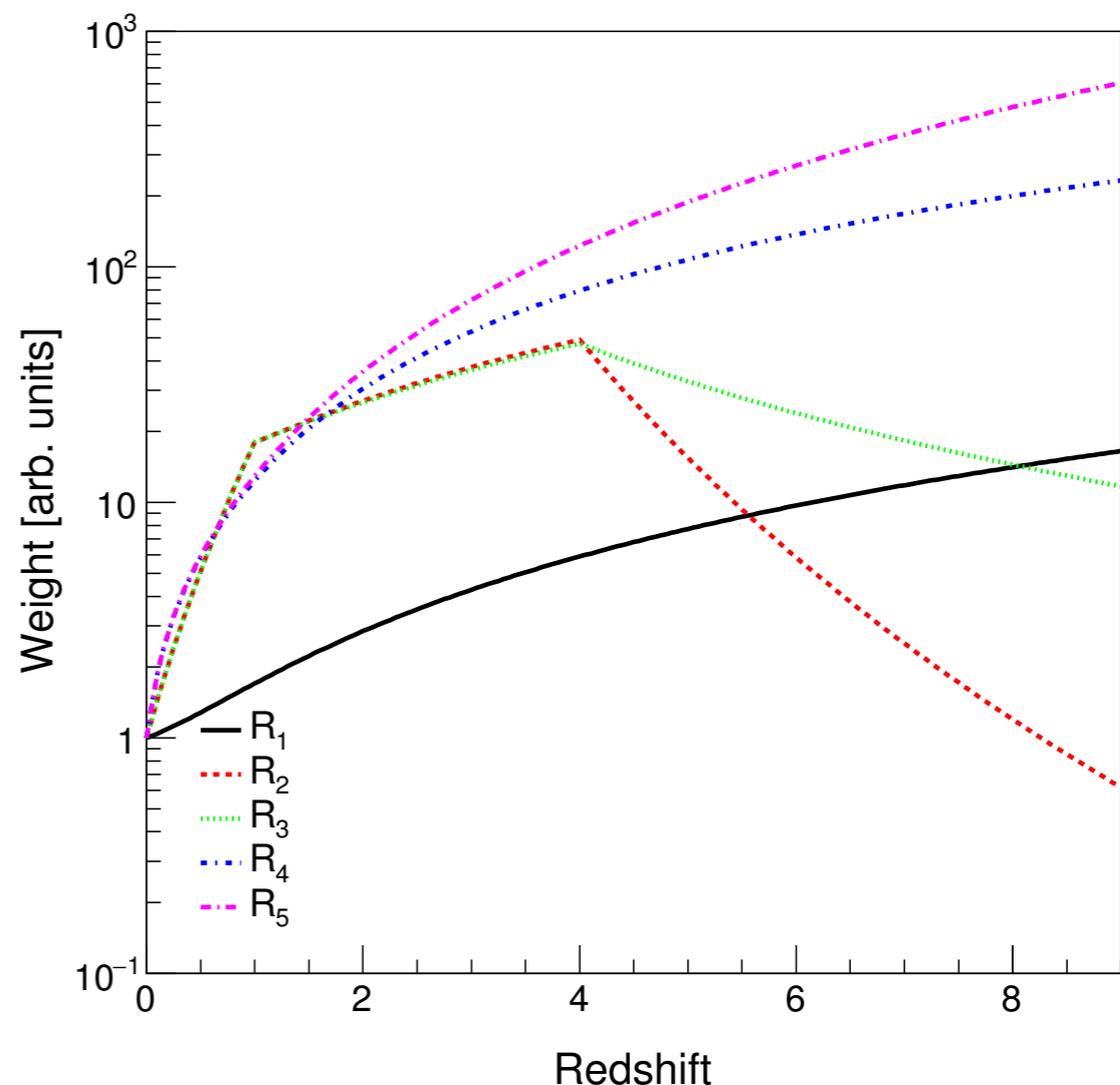
Model	Γ	$\log_{10}(R_{\text{cut}}/V)$	fH	fHe	fN	fSi	fFe
C_1	1	18.699	0.7692	0.1538	0.0461	0.0231	0.00759
C_2	1	18.5	0	0	0	1	0
C_3	1.25	18.5	0.365	0.309	0.121	0.1066	0.098
C_4	2.7	∞	1	0	0	0	0

Note. Γ is the spectral index, R_{cut} is the rigidity cutoff and fH , fHe , fN , fSi , and fFe are the fractions of each nuclei.

Model of UHECR Sources

Models of Source Distribution

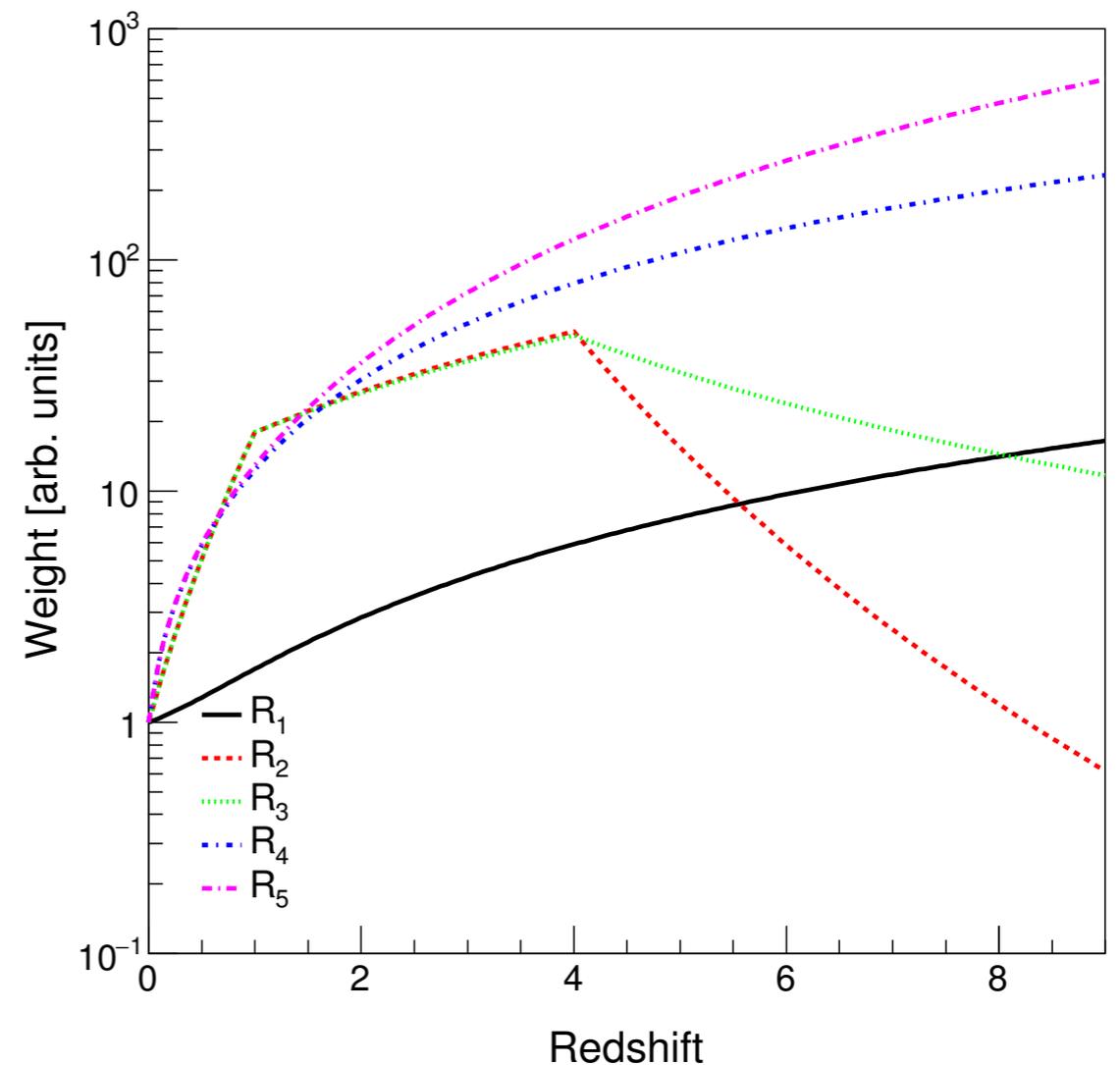
1. R_1 : sources are uniformly distributed in a comoving volume;
2. R_2 : sources follow the star formation distribution given in Hopkins & Beacom (2006). The evolution is proportional to $(1+z)^{3.4}$ for $z < 1$, to $(1+z)^{-0.26}$ for $1 \leq z < 4$ and to $(1+z)^{-7.8}$ for $z \geq 4$;
3. R_3 : sources follow the star formation distribution given in Yüksel et al. (2008). The evolution is proportional to $(1+z)^{3.4}$ for $z < 1$, to $(1+z)^{-0.3}$ for $1 \leq z < 4$ and to $(1+z)^{-3.5}$ for $z \geq 4$;
4. R_4 : sources follow the GRB rate evolution from Le & Dermer (2007). The evolution is proportional to $(1+8z)/[1+(z/3)^{1.3}]$;
5. R_5 : sources follow the GRB rate evolution from Le & Dermer (2007). The evolution is proportional to $(1+11z)/[1+(z/3)^{0.5}]$.



Model of UHECR Sources

Models of Source Distribution

1. R_1 : sources are uniformly distributed in a comoving volume;
2. R_2 : sources follow the star formation distribution given in Hopkins & Beacom (2006). The evolution is proportional to $(1+z)^{3.4}$ for $z < 1$, to $(1+z)^{-0.26}$ for $1 \leq z < 4$ and to $(1+z)^{-7.8}$ for $z \geq 4$;
3. R_3 : sources follow the star formation distribution given in Yüksel et al. (2008). The evolution is proportional to $(1+z)^{3.4}$ for $z < 1$, to $(1+z)^{-0.3}$ for $1 \leq z < 4$ and to $(1+z)^{-3.5}$ for $z \geq 4$;
4. R_4 : sources follow the GRB rate evolution from Le & Dermer (2007). The evolution is proportional to $(1+8z)/[1+(z/3)^{1.3}]$;
5. R_5 : sources follow the GRB rate evolution from Le & Dermer (2007). The evolution is proportional to $(1+11z)/[1+(z/3)^{0.5}]$.



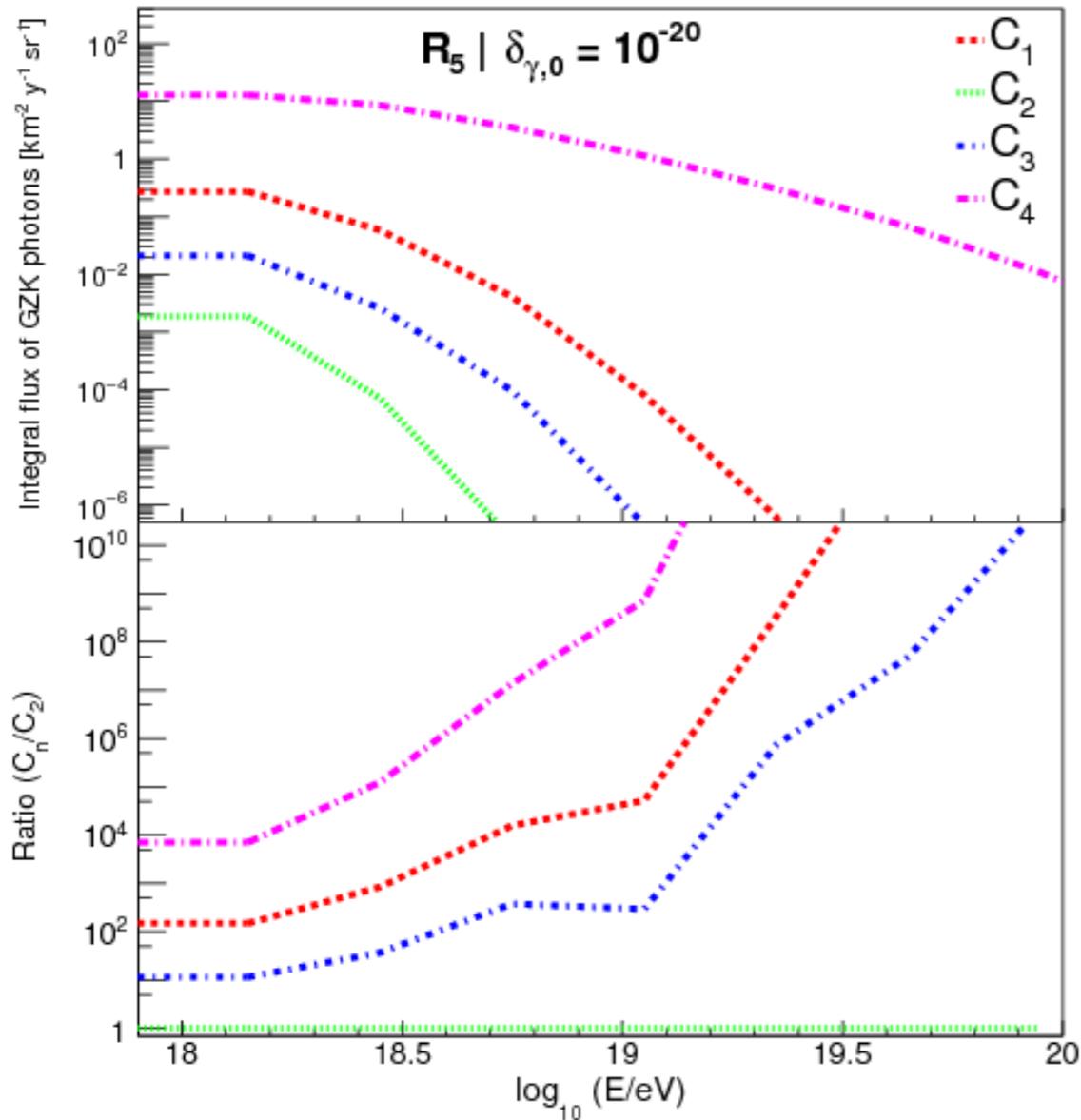
CRPropa3/EleCA

(Settimo & Domenico 2015; Batista et al. 2016)

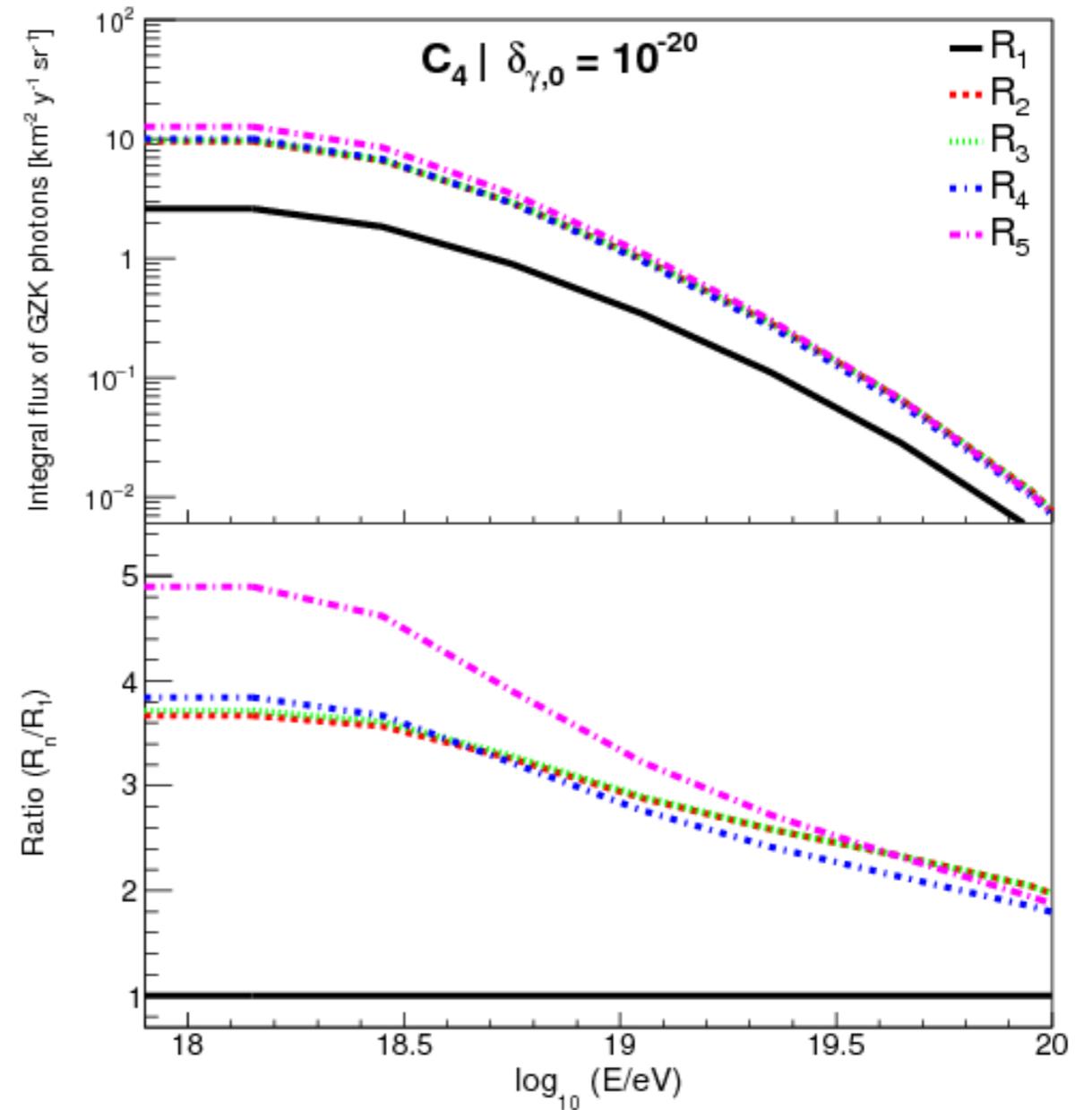
Lang, Martinez & De Souza
ApJ 853, no.1, 23 (2018)

Integral flux of GZK photons + LIV

... for each source model



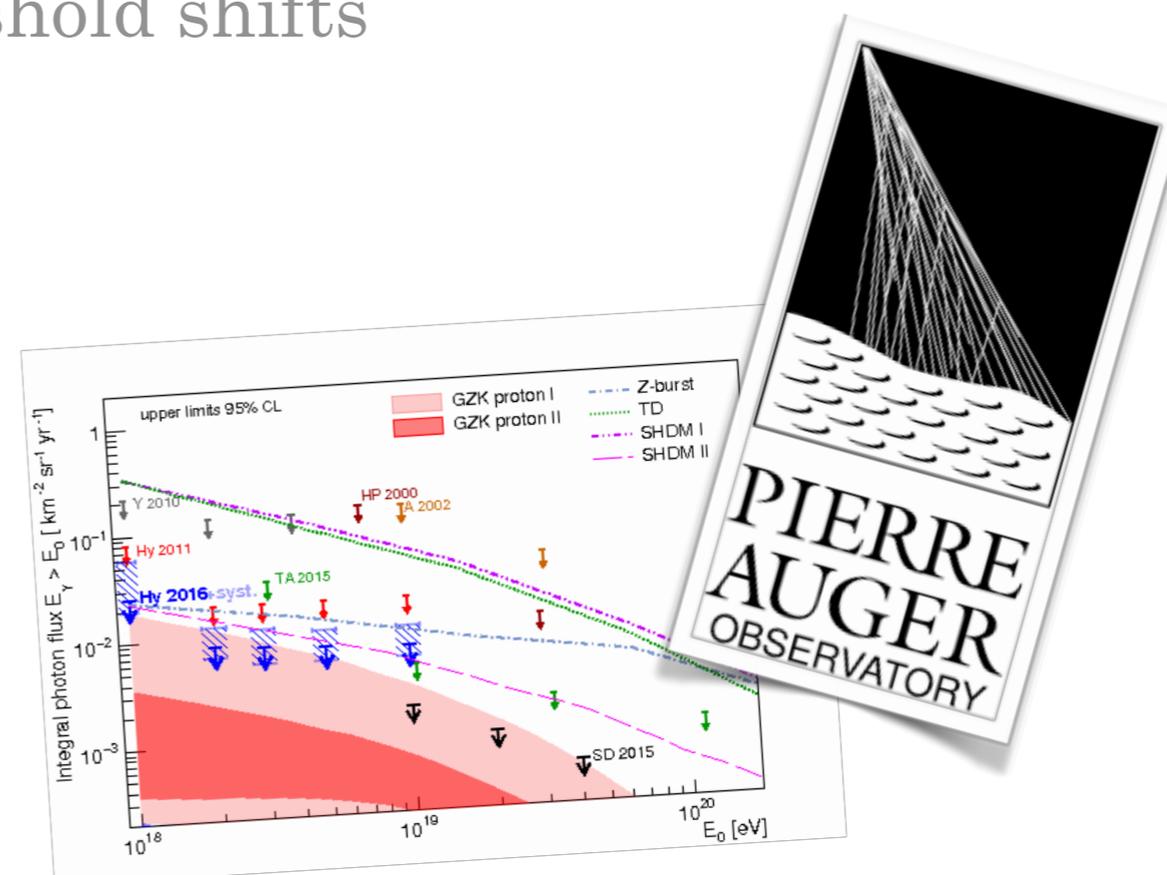
... for each source evolution model



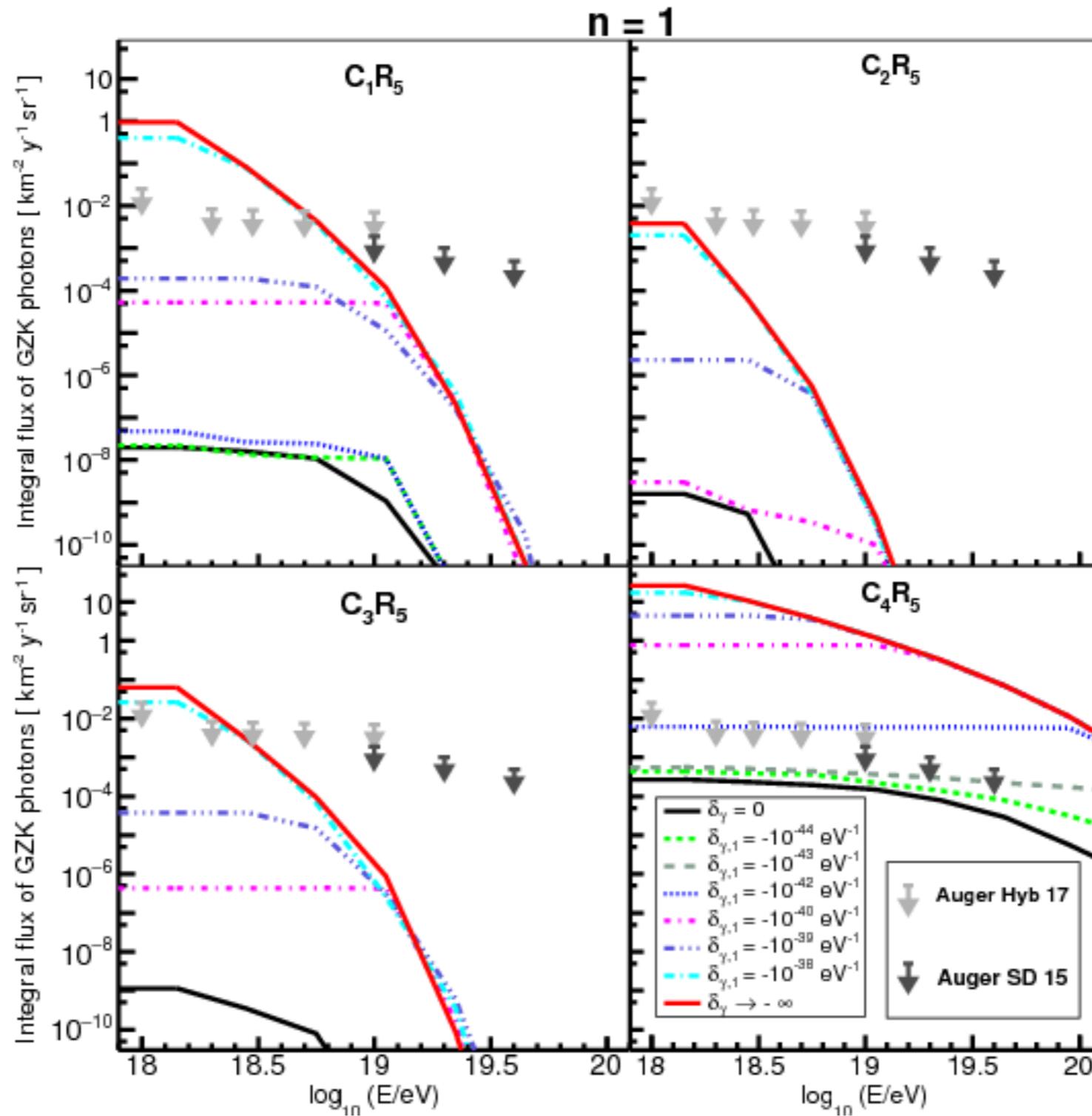
Different LIV coefficients result in a shift up an down

Index

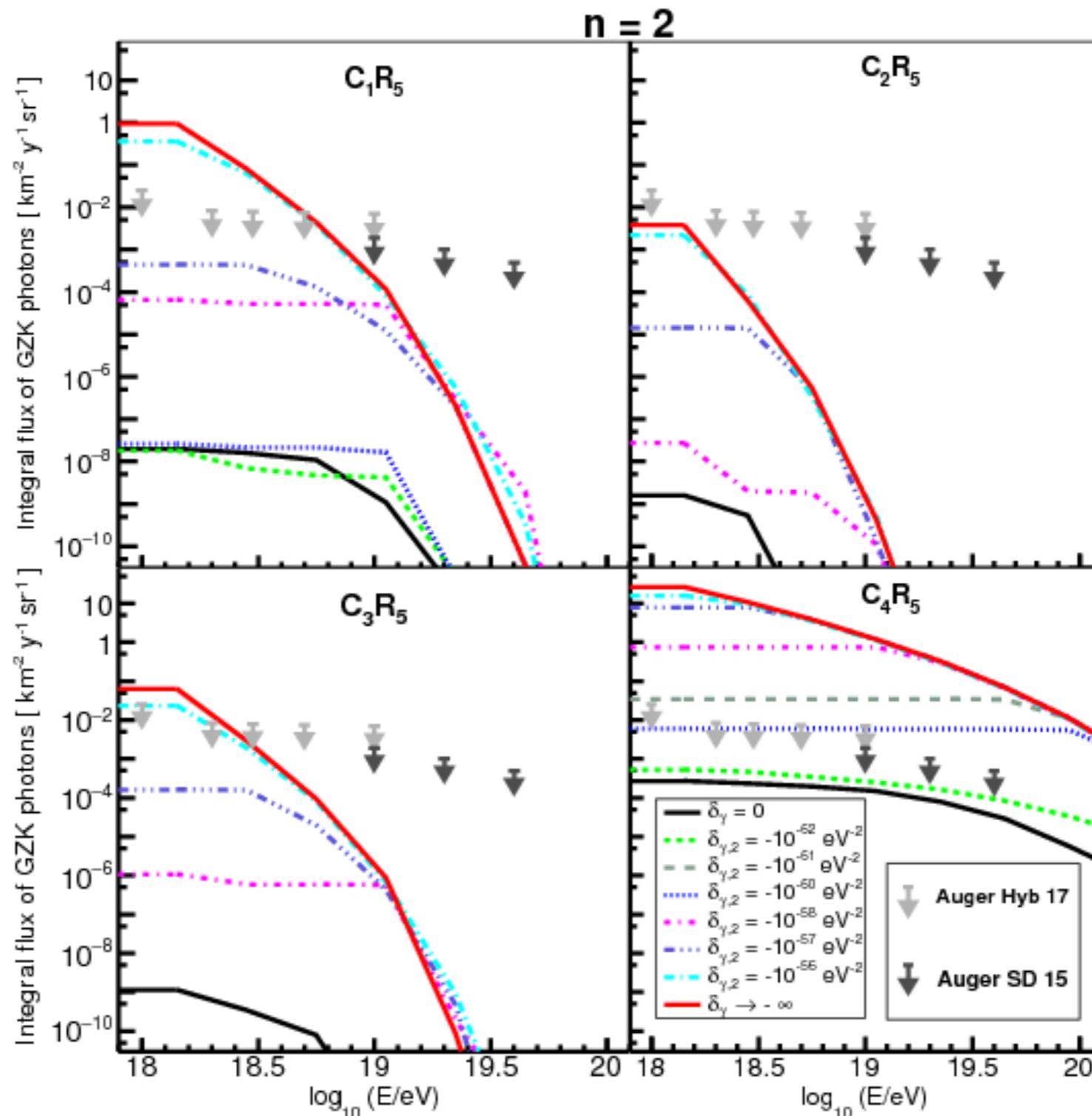
- I. Lorentz invariance violation (LIV)
- II. Limits: TeV γ -rays
 - i. Time Energy Dependent delay
 - ii. Photon Decay
 - iii. Pair production threshold shifts
- III. UHECR
 - i. GZK-photons + LIV
 - ii. **Limits**



GZK photon flux + LIV



GZK photon flux + LIV



Model C_3R_5 was shown to (best) describe the energy spectrum, composition, and arrival direction of UHECR*

*M. Unger et al 2015, Phys. Rev. D, 92, 123001

Limits on the LIV Coefficients Imposed by This Work for Each Source Model and LIV Order (n)

Model	$\delta_{\gamma,0}^{\text{limit}}$	$\delta_{\gamma,1}^{\text{limit}} (\text{eV}^{-1})$	$\delta_{\gamma,2}^{\text{limit}} (\text{eV}^{-2})$
C_1R_5	$\sim -10^{-20}$	$\sim -10^{-38}$	$\sim -10^{-56}$
C_2R_5
C_3R_5	$\sim -10^{-20}$	$\sim -10^{-38}$	$\sim -10^{-56}$
C_4R_5	$\sim -10^{-22}$	$\sim -10^{-42}$	$\sim -10^{-60}$

Limits on the LIV Coefficients Imposed by Other Works Based on Gamma-Ray Propagation

Model	$\delta_{\gamma,0}^{\text{limit}}$	$\delta_{\gamma,1}^{\text{limit}} (\text{eV}^{-1})$	$\delta_{\gamma,2}^{\text{limit}} (\text{eV}^{-2})$
Galaverni & Sigl (2008a)	...	-1.97×10^{-43}	-1.61×10^{-63}
H.E.S.S.—PKS 2155–304 (2011)	...	-4.76×10^{-28}	-2.44×10^{-40}
Fermi—GRB 090510 (2013)	...	-1.08×10^{-29}	-5.92×10^{-41}
H.E.S.S.—Mrk 501 (2017)	...	-9.62×10^{-29}	-4.53×10^{-42}

Limits on the LIV Coefficients Imposed by This Work for Each Source Model and LIV Order (n)

Model	$\delta_{\gamma,0}^{\text{limit}}$	$\delta_{\gamma,1}^{\text{limit}} (\text{eV}^{-1})$	$\delta_{\gamma,2}^{\text{limit}} (\text{eV}^{-2})$
C_1R_5	$\sim -10^{-20}$	$\sim -10^{-38}$	$\sim -10^{-56}$
C_2R_5
C_3R_5	$\sim -10^{-20}$	$\sim -10^{-38}$	$\sim -10^{-56}$
C_4R_5	$\sim -10^{-22}$	$\sim -10^{-42}$	$\sim -10^{-60}$

Limits on the LIV Coefficients Imposed by Other Works Based on Gamma-Ray Propagation

Model	$\delta_{\gamma,0}^{\text{limit}}$	$\delta_{\gamma,1}^{\text{limit}} (\text{eV}^{-1})$	$\delta_{\gamma,2}^{\text{limit}} (\text{eV}^{-2})$
Galaverni & Sigl (2008a)	...	-1.97×10^{-43}	-1.61×10^{-63}
H.E.S.S.—PKS 2155–304 (2011)	...	-4.76×10^{-28}	-2.44×10^{-40}
Fermi—GRB 090510 (2013)	...	-1.08×10^{-29}	-5.92×10^{-41}
H.E.S.S.—Mrk 501 (2017)	...	-9.62×10^{-29}	-4.53×10^{-42}

Conclusions and remarks II

- ❖ We studied the effect of possible LIV in the propagation of photons in the universe.
- ❖ The **mean-free path of the pair production** interaction was calculated **considering LIV effects**.
- ❖ We found that even moderate LIV coefficients introduce a significant change in the mean-free path of the interaction.
- ❖ **The GZK photon flux including LIV was obtained** for different source models and source distribution models.
- ❖ **Limits to the LIV** coefficient were established based on source models **compatible with the most updated data of UHECR**.
- ▶ The limits presented here are several orders of magnitude more restrictive than previous calculations based on the arrival time of TeV photons; however, the comparison is not straightforward due to different systematics of the measurements and energy of the photons.

Thanks!