

Discussion on global fits

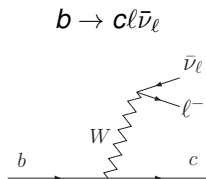
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GdR-InF annual meeting, Arles 07/11/18



Two transitions of interest



tree (charged) ($V - A$)

$$\bar{B} \rightarrow D \ell \bar{\nu}_\ell$$

$$\bar{B} \rightarrow D^* \ell \bar{\nu}_\ell$$

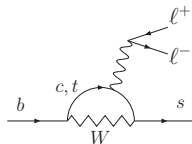
Total Br

$$\ell = \tau, \mu, e$$

$$R_{D^{(*)}} = \frac{\text{Br}(B \rightarrow D^{(*)} \tau \nu)}{\text{Br}(B \rightarrow D^{(*)} \ell \bar{\nu}_\ell)}$$

$(\ell = e, \mu)$

$$b \rightarrow s \ell^+ \ell^-$$



loop (neutral)

$$B \rightarrow K \ell \ell$$

$$B \rightarrow K^* \ell \ell, B_s \rightarrow \phi \ell \ell$$

$d\Gamma/dq^2 + \text{Angular obs}$

$$\ell = \mu, e$$

$$R_{K^{(*)}} = \frac{\text{Br}(B \rightarrow K^{(*)} \mu \mu)}{\text{Br}(B \rightarrow K^{(*)} e e)}$$

$\text{Br}(K, K^*, \phi + \mu \mu)$
angular obs (e.g., P'_5)

SM
Spin 0
Spin 1
Observables
with
Tensions

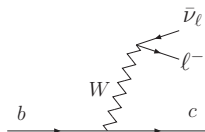
Two transitions exhibiting interesting patterns of deviations from SM

A major tool: effective Hamiltonian

Starting from SM (or extensions)
and integrating out heavy/energetic
degrees of freedom

$$\mathcal{H}^{\text{eff}} = CKM \times \mathcal{C}_i \times \mathcal{O}_i$$

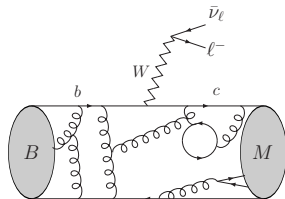
$$\langle M | \mathcal{H}^{\text{eff}} | B \rangle = CKM \times \mathcal{C}_i \times \langle M | \mathcal{O}_i | B \rangle$$



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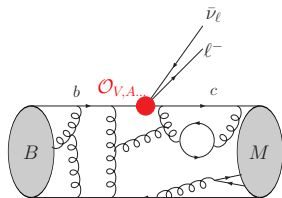
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involving hadronic quantities such as **form factors** (and others)

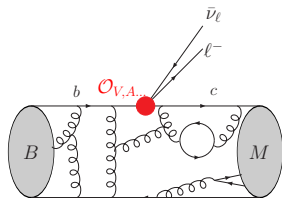
selecting processes for accurate predictions:

- semileptonic decays (form factors, not more complicated objects)
- ratios of branching ratios with different leptons
- ratios of observables with similar dependence on form factors
 \implies observables with limited sensitivity to (ratio of form) factors

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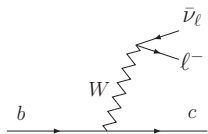
Two possible uses of effective approaches

- fixing \mathcal{C}_i , computing SM and comparing with the data
- determining \mathcal{C}_i from the data and compare with SM or NP models

How various analyses can differ ?

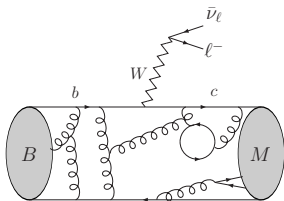
- Statistical treatment (Bayesian vs frequentist)
- Form factors
 - low q^2 /large-meson recoil: light-cone sum rules (LCSR)
 - large q^2 /low-meson recoil: lattice QCD (LQCD)
- Other hadronic inputs (intermediate resonances...)
- Help (or not) from effective theories ($m_b \rightarrow \infty$)
 - low q^2 /large-meson recoil: Soft-Collinear Effective Theory (SCET)
 - large q^2 /low-meson recoil: Heavy Quark Effective Theory (HQET)
- Sticking to one particular quark decay (e.g., $b \rightarrow s\mu\mu$ only) or trying to connect several of them through assumptions/symmetries
- NP scenarios considered
 - in all \mathcal{C}_i or only some of them ?
 - correlating NP contributions among \mathcal{C}_i ?
 - including imaginary parts or not ?
 - only violating lepton universality or also lepton universal ?

$b \rightarrow cl\bar{\nu}_\ell$ effective Hamiltonian



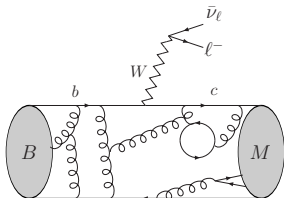
$$\mathcal{H}^{\text{eff}}(b \rightarrow cl\nu) \propto G_F V_{cb} \sum C_i \mathcal{O}_i$$

$b \rightarrow cl\bar{\nu}_\ell$ effective Hamiltonian



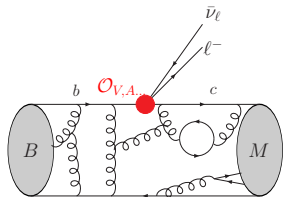
$$\mathcal{H}^{\text{eff}}(b \rightarrow cl\nu) \propto G_F V_{cb} \sum C_i \mathcal{O}_i$$

$b \rightarrow c\ell\bar{\nu}_\ell$ effective Hamiltonian

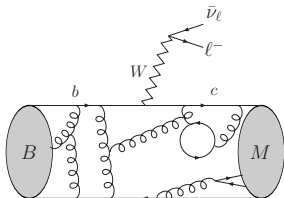


$$\mathcal{H}^{\text{eff}}(b \rightarrow c\ell\nu) \propto G_F V_{cb} \sum C_i \mathcal{O}_i$$

- In the SM
 - $\mathcal{O}_{V_L} = (\bar{c}\gamma^\mu P_L b)(\bar{\ell}\gamma_\mu P_L \nu_\ell)$ [W exchange]
 - $C_{V_L} = 1$ and universal for all three leptons
- Hadronic uncertainties all summarised in form factors defined from $\langle M | \mathcal{O}_i | B \rangle$

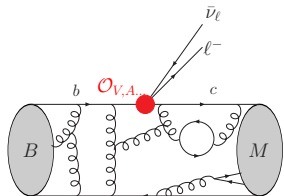


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 - $C_{V_L} = 1$ and universal for all three leptons
- Hadronic uncertainties all summarised in form factors defined from $\langle M | \mathcal{O}_i | B \rangle$
- NP changes short-distance C_i for SM or new long-distance ops \mathcal{O}_i



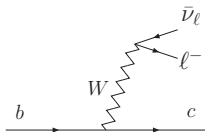
- Chirally flipped ($W \rightarrow W_R$)
- (Pseudo)scalar ($W \rightarrow H^+$)
- Tensor operators ($W \rightarrow T$)

$$\mathcal{O}_{V_L} \rightarrow \mathcal{O}_{V_R} \propto (\bar{c}\gamma^\mu P_R b)(\bar{\ell}\gamma_\mu P_L \nu_\ell)$$

$$\mathcal{O}_{V_L} \rightarrow \mathcal{O}_{S_L} \propto (\bar{c}P_L b)(\bar{\ell}P_L \nu_\ell), \mathcal{O}_{S_R}$$

$$\mathcal{O}_{V_L} \rightarrow \mathcal{O}_{T_L} \propto (\bar{c}\sigma^{\mu\nu} P_L b)(\bar{\ell}\sigma_{\mu\nu} P_L \nu_\ell)$$

$b \rightarrow c \ell \bar{\nu}_\ell$ typical observables



$$B \rightarrow D^* \ell \nu$$

$$B \rightarrow D \ell \nu$$

$$B \rightarrow J/\psi \ell \nu$$

$$B_c \rightarrow \tau \nu$$

$$R_D = \frac{Br(B \rightarrow D \tau \nu)}{Br(B \rightarrow D \ell \nu)}$$

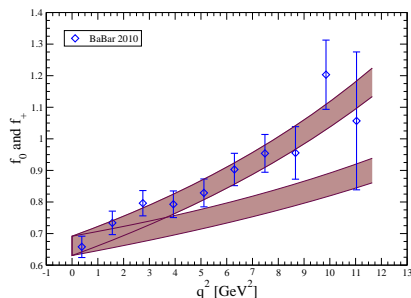
$$R_{D^*} = \frac{Br(B \rightarrow D^* \tau \nu)}{Br(B \rightarrow D^* \ell \nu)}, P_\tau, f_L(D^*)$$

$$R_{J/\psi} = \frac{Br(B_c \rightarrow J/\psi \tau \nu)}{Br(B_c \rightarrow J/\psi \ell \nu)}$$

$$Br(B_c \rightarrow \tau \nu) \text{ [bound from } \Gamma(B_c)\text{]}$$

$\ell = e, \mu$ (integrated over all phase space)

Hadronic uncertainties

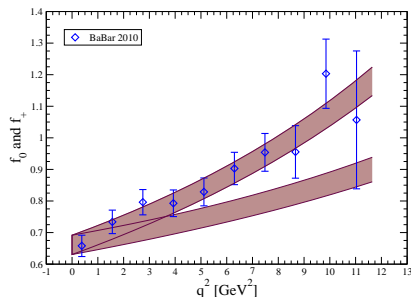


$B \rightarrow D \ell \bar{\nu}_\ell$ branching ratios

- SM: 2 form factors f_+ (vector) and f_0 (scalar)
- NP: 1 more f_T (tensor)
- From lattice QCD, extrapolated over whole kinematic range

[HPQCD, FNAL/MILC collaborations]

Hadronic uncertainties



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[HPQCD, FNAL/MILC collaborations]

$B \rightarrow D^*\ell\bar{\nu}_\ell$ branching ratios

- 4 form factors $V, A_{0,1,2}$ (vector/axial) + NP: 3 more $T_{1,2,3}$ (tensor)
- No complete lattice determination [Fajfer, Kamenik, Nisandzic]
 - Belle: $B \rightarrow D^*\ell\bar{\nu}_\ell$ ($\ell = e, \mu$) form factors, assuming no NP [Jung, Straub]
 - Supplemented by HQET considerations
 - relations in the limit $m_b \rightarrow \infty$, normalisation in the no-recoil limit
 - corrections to be estimated with theoretical prejudices (CLN param)
 - Or fit using a generic polynomial z -expansion (BGL param)
 - Ongoing discussion, potential impact on $|V_{cb}|$ but not much on R_{D^*}

[Bigi, Gambino, Schacht; Grinstein, Koback; Bernlochner, Ligeti, Papucci, Robinson]

Further data and questions

$R_{J/\psi}(B_c \rightarrow J/\psi \ell \bar{\nu}_\ell)$

- LHCb measurement low compared to theoretical estimates, and rather interestingly

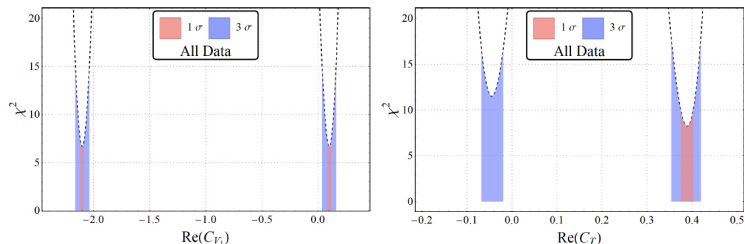
$$\frac{R_D}{R_{D;SM}} \simeq \frac{R_{D^*}}{R_{D^*;SM}} \simeq \frac{R_{J/\psi}}{R_{J/\psi;SM}}$$

- But current estimates of the form factors mainly based on models with uncertainties difficult to assess (compared to R_D and R_{D^*})

Potential issues

- Cross-checks of the form factors would be very welcome !
- Radiative corrs 3-4% for $E_\gamma^{\text{cut}} = 20 - 40$ MeV [De Boer, Kitahira, Nisandzic]
- Width of the D^* in $D\pi$ [Chavez-Saab, Toledo; Le Yaouanc, Leroy, Roudeau]
 - Contribution from longitudinal polarisation to R_{D^*} ?
 \implies Possibly shift $\simeq 9\%$ in SM direction,
potentially reducing discrepancy from 3.4σ to 2.2σ
 - Tail of D^* at high $D\pi$ mass contributing to $B \rightarrow [D\pi]_{\text{broad}} \ell \nu_\ell$?

Global fits in $b \rightarrow c\ell\bar{\nu}_\ell$

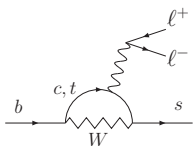


[Bhattacharyya, Nandi, Patra; Alok, Kumar, Kumar, Kumbhakar, Uma Sankar; Kumar, London, Watanabe; Freytsis, Ligeti, Ruderman]

- Often NP only for $\ell = \tau$, assuming real Wilson coefficients (no CP violation), but some studies with also imaginary contributions
- Right-handed and (pseudo)scalar couplings disfavoured by B_c width (bound on $B_c \rightarrow \tau\nu$) and shape of $d\Gamma(B \rightarrow D^*\tau\nu)/dq^2$
- Tensor could describe the data, strong impact on P_τ and f_L
- Most simple explanation: NP in $C_{VL\tau}$, change of G_F for $b \rightarrow c\tau\bar{\nu}_\tau$

R_D and R_{D^*} not enough: more observables (angular analysis !)

$b \rightarrow sll$ effective Hamiltonian



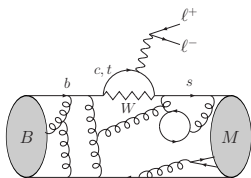
$$\mathcal{H}(b \rightarrow s \gamma^{(*)}) \propto G_F V_{ts}^* V_{tb} \sim C_i \mathcal{O}_i$$

to separate short and long distances ($\mu_b = m_b$)

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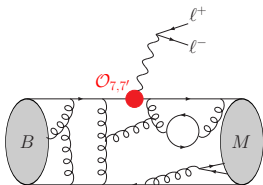
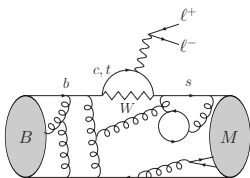


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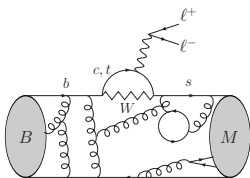
- $\mathcal{O}_7 = \frac{e}{g^2} m_b \bar{s} \sigma^{\mu\nu} (1 + \gamma_5) F_{\mu\nu} b$ [real or soft photon]



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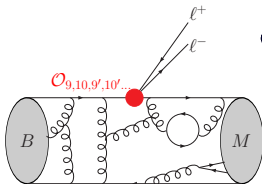
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- $\mathcal{O}_9 = \frac{e^2}{g^2} \bar{s} \gamma_\mu (1 - \gamma_5) b \bar{l} \gamma^\mu l$ [$b \rightarrow s\mu\mu$ via Z /hard γ ...]

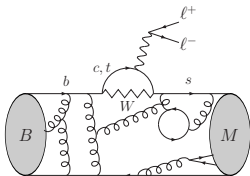
- $\mathcal{O}_{10} = \frac{e^2}{g^2} \bar{s} \gamma_\mu (1 - \gamma_5) b \bar{l} \gamma^\mu \gamma_5 l$ [$b \rightarrow s\mu\mu$ via Z]



$b \rightarrow sll$ effective Hamiltonian

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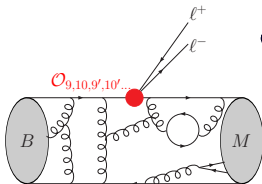
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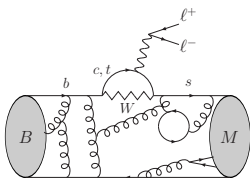


$$C_7^{\text{SM}} = -0.29, \quad C_9^{\text{SM}} = 4.1, \quad C_{10}^{\text{SM}} = -4.3$$

$b \rightarrow sll$ effective Hamiltonian

$$\mathcal{H}(b \rightarrow s\gamma^{(*)}) \propto G_F V_{ts}^* V_{tb} \sim \mathcal{C}_i \mathcal{O}_i$$

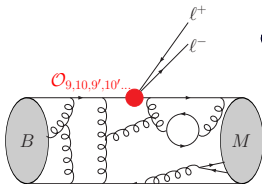
to separate short and long distances ($\mu_b = m_b$)



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$$\mathcal{C}_7^{\text{SM}} = -0.29, \quad \mathcal{C}_9^{\text{SM}} = 4.1, \quad \mathcal{C}_{10}^{\text{SM}} = -4.3$$

NP changes short-distance \mathcal{C}_i or add new operators \mathcal{O}_i

- Chirally flipped ($W \rightarrow W_R$)

$$\mathcal{O}_7 \rightarrow \mathcal{O}_{7'} \propto \bar{s} \sigma^{\mu\nu} (1 - \gamma_5) F_{\mu\nu} b$$

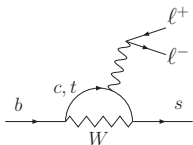
- (Pseudo)scalar ($W \rightarrow H^+$)

$$\mathcal{O}_9, \mathcal{O}_{10} \rightarrow \mathcal{O}_S \propto \bar{s} (1 + \gamma_5) b \bar{l} l, \quad \mathcal{O}_P$$

- Tensor operators ($\gamma \rightarrow T$)

$$\mathcal{O}_9 \rightarrow \mathcal{O}_T \propto \bar{s} \sigma_{\mu\nu} (1 - \gamma_5) b \bar{l} \sigma_{\mu\nu} l$$

$b \rightarrow sll$ typical observables

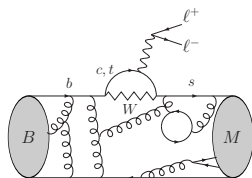


$B \rightarrow K^* \mu\mu$	$\text{Br}, P_{1,2}, P'_{4,5,6,8}, F_L$
$B \rightarrow K^* ee$	$P_{1,2,3}, P'_{4,5}, F_L$
$B \rightarrow K^* ll$	$R_K, R_{K^*}, Q_{4,5}$
$B_s \rightarrow \phi\mu\mu$	$\text{Br}, P_1, P'_{4,6}, F_L$
$B \rightarrow K\mu\mu$	Br
$B \rightarrow X_S\gamma$	Br
$B \rightarrow X_S\mu\mu$	Br
$B_s \rightarrow \mu\mu$	Br
$B_s \rightarrow \phi\gamma$	$\text{Br}, S_{\phi\gamma}, A_{\Delta\Gamma;\phi\gamma}$
$B \rightarrow K^*\gamma$	$\text{Br}, S_{K^*\gamma}$

binned in most cases, sometimes with CP asymmetries

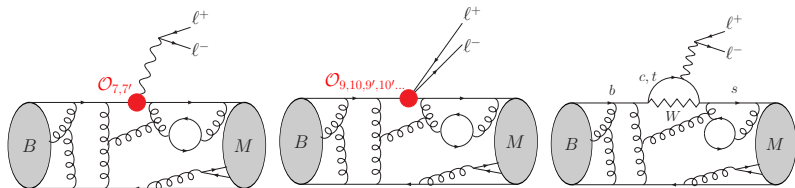
Two sources of hadronic uncertainties

$$A(B \rightarrow M \ell \ell) = \frac{G_F \alpha}{\sqrt{2} \pi} V_{tb} V_{ts}^* [(A_\mu + T_\mu) \bar{u}_\ell \gamma^\mu v_\ell + B_\mu \bar{u}_\ell \gamma^\mu \gamma_5 v_\ell]$$



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Form factors (local)

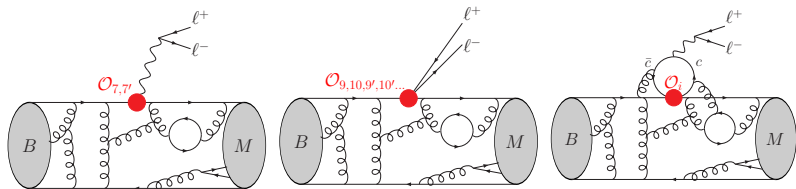
- Local contributions (more terms if NP in non-SM C_i): **form factors**

$$A_\mu = -\frac{2m_b q^\nu}{q^2} C_7 \langle M | \bar{s} \sigma_{\mu\nu} P_R b | B \rangle + C_9 \langle M | \bar{s} \gamma_\mu P_L b | B \rangle$$

$$B_\mu = C_{10} \langle M | \bar{s} \gamma_\mu P_L b | B \rangle$$

Two sources of hadronic uncertainties

$$A(B \rightarrow M \ell \ell) = \frac{G_F \alpha}{\sqrt{2} \pi} V_{tb} V_{ts}^* [(A_\mu + T_\mu) \bar{u}_e \gamma^\mu v_\ell + B_\mu \bar{u}_e \gamma^\mu \gamma_5 v_\ell]$$



Form factors (local)

Charm loop (non-local)

- Local contributions (more terms if NP in non-SM C_i): **form factors**

$$A_\mu = -\frac{2m_b q^\nu}{q^2} C_7 \langle M | \bar{s} \sigma_{\mu\nu} P_R b | B \rangle + C_9 \langle M | \bar{s} \gamma_\mu P_L b | B \rangle$$

$$B_\mu = C_{10} \langle M | \bar{s} \gamma_\mu P_L b | B \rangle$$

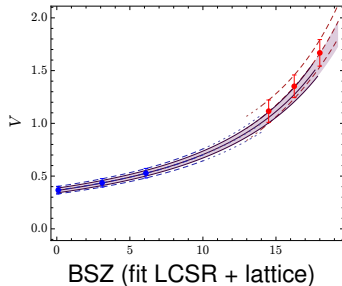
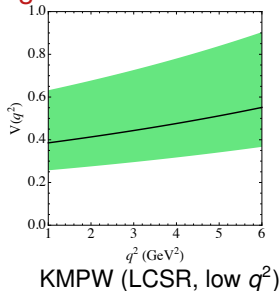
- Non-local contributions (charm loops): **hadronic contribs.**

T_μ contributes like $O_{7,9}$, but depends on q^2 and external states

Hadronic uncertainties: form factors

3 form factors for K , 7 form factors for K^* and ϕ

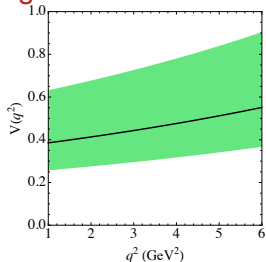
- low recoil: **lattice**, with correlations [Horgan, Liu, Meinel, Wingate; HPQCD collab]
- large recoil: **B-meson Light-Cone Sum Rule**, large error bars and no correlations [Khodjamirian, Mannel, Pivovarov, Wang]
recently reanalysed with correlations [Gubernari, Kokulu, van Dyk]
- all: fit **light-meson LCSR** + lattice, small errs, correls [Bharucha, Straub, Zwicky]



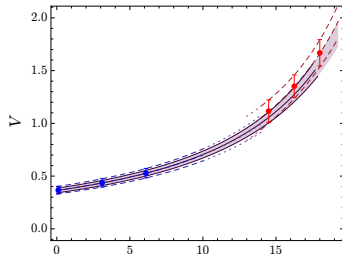
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KMPW (LCSR, low q^2)



BSZ (fit LCSR + lattice)

- former controversies about EFT to obtain/restore correlations for form factors discussed and all approaches in good agreement

[Jäger, Camalich; Capdevila, SDG, Hofer, Matias; Straub, Altmannshofer; Hurth, Mahmoudi]

Hadronic uncertainties: charm loops

Charm loops

- important for resonance regions (charmonia)
- SM effect contributing to C_9
- expected to depend on q^2
- . . . but lepton universal (little effect on R_K , even with NP)

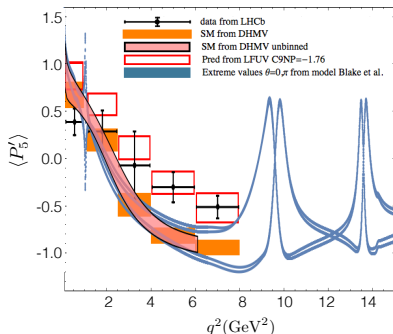
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Several theo/pheno approaches

- LCSR estimate [Khodjamirian, Mannel, Pivovarov, Wang]
- order of magnitude estimate for the fits (LCSR or Λ/m_b), check with bin-by-bin fits [Crivellin, Capdevila, SDG, Hofer, Matias; Straub, Altmannshoffer; Hurth, Mahmoudi]
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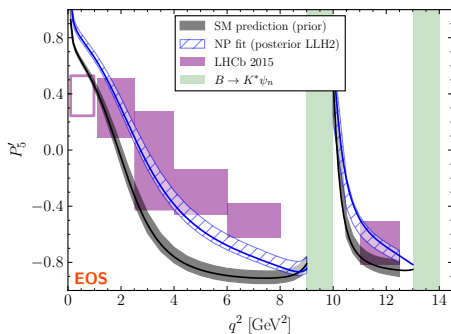
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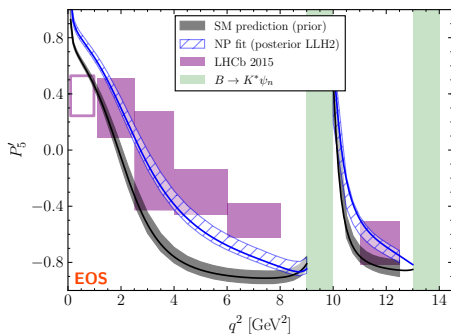
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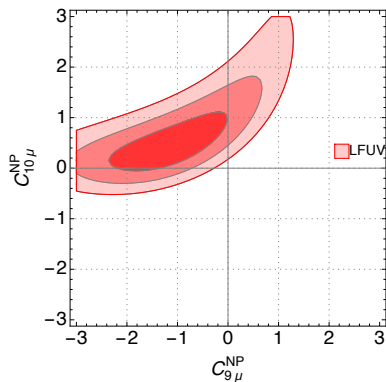
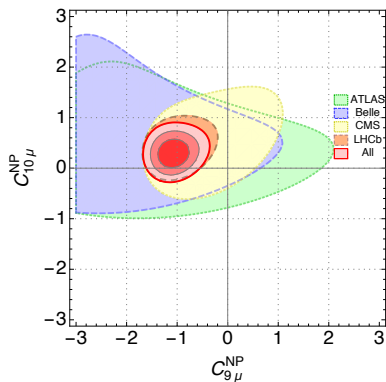
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No hint in the fits of missing large q^2 -dependent contribution

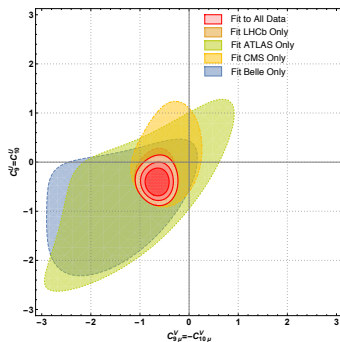
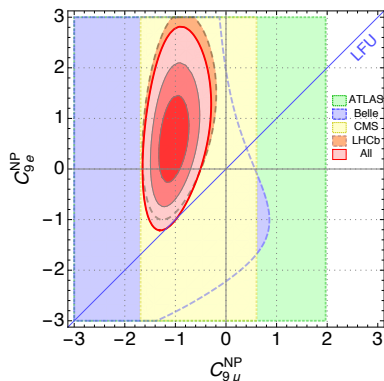
Global fits in $b \rightarrow sll$

- 175 obs in [Capdevila, Crivellin, SDG, Matias, Virto], agree well with other global fits
[Straub, Stangl, Altmannshoffer; Hurth, Mahmoudi, Neshatpour; Geng, Grinstein, Jäger, Camalich, Ren, Shi]
- Real contributions, based on absence of CP violation
- Favoured $C_{9\mu}^{NP} \simeq O(-1)$ + smaller corrections to other C_j
- $C_{9\mu}^{NP} = -C_{10\mu}^{NP}$ also good scenario (NP models with $SU(2)_L$)
- Overall consistency (All vs LFUV obs, channels, Brs vs angular)



Are we general enough ?

C_{il}^{NP} $\ell = e, \mu, \tau$ not related, to be fixed by data (scarce for e , none for τ)

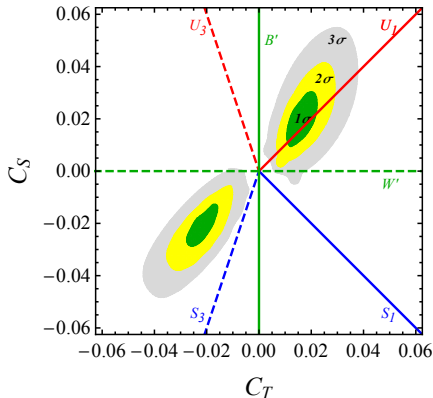


- Often no NP assumed in $b \rightarrow see$, but still room (eg NP in $C_{9\mu}, C_{9e}$)
- 4-par best-fit pt $(C_{9e}, C_{9\mu}, C_{10e}, C_{10\mu})^{NP} = -1.26, -1.18, 1.14, 0.23$
- Possibility of large LFU NP contribs, with preference for $V - A$ LFUV but $V + A$ for LFU

[Algueró, Capdevila, SDG, Masjuan, Matias]

$b \rightarrow cl\bar{\nu}_\ell$ and $b \rightarrow sl\ell$ combined analysis

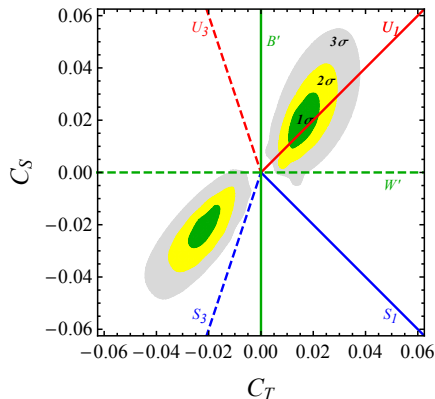
- SMEFT: NP higher-dim ops obeying $SU_C(3) \times SU_L(2) \times U_Y(1)$
- Restrictive (but reasonable) assumptions [Butazzo, Greljo, Isidroi, Marzocca]
 - Only left-handed fields
 - No lepton-flavour violating contributions
 - $U(2)_q \times U_\ell(2)$ symmetry for couplings with same structure



$$\lambda_{ij}^q \lambda_{ab}^\ell \left[\frac{C_S}{\Lambda^2} (\bar{Q}_L^i \gamma_\mu Q_L^j) (\bar{L}_L^a \gamma^\mu L_L^b) + \frac{C_T}{\Lambda^2} (\bar{Q}_L^i \gamma_\mu \sigma^\alpha Q_L^j) (\bar{L}_L^a \gamma^\mu \sigma^\alpha L_L^b) \right]$$
$$Q_L^i = \begin{pmatrix} V_{ji}^* u_L^j \\ d_L^i \end{pmatrix} \quad L_L^a = \begin{pmatrix} \nu_L^a \\ \ell_L^a \end{pmatrix}$$

Single-mediator models and more

[Butazzo, Greljo, Isidroi, Marzocca]



- Disfavours colourless vectors (W', Z' , green) and coloured scalars (S_1, S_3 leptoquarks, blue) + high p_T constraints
- Favours U_1 vector leptoquark ($3, 1, 2/3$), which also passes direct LHC production limits
- Same conclusions taking a general structure of the couplings

[Kumar, London, Watanabe]

Possible to consider models with more than 1 mediator

- Two scalar leptoquarks
- Three-generation Pati-Salam
- Composite models
- ...

[See P. Stangl's talk]

$b \rightarrow s\ell^+\ell^-$ and $b \rightarrow c\ell\bar{\nu}_\ell$

- Many observables, more or less sensitive to hadronic unc.
- Global fit to $b \rightarrow c\ell\bar{\nu}$ still relying on limited amount of information, with questions on hadronic uncertainties (form factors, D^* width)
- Global fit to $b \rightarrow s\ell^+\ell^-$ in favour of large deviation for C_9 in $b \rightarrow s\mu\mu$ and does not seem to favour hadronic explanations
- Global fits of both sets using SMEFT, with many models proposed for either or both sets of deviations
- Several tools to perform some fits (flavio, EOS, Hammer. . .)

Where to go ?

- Better measurements of q^2 and angular dependence
- More info on processes with e and/or τ
- Other LFU violating observables
- Provide lattice form factors over larger range
- Further constraints on $c\bar{c}$ loops
- New observables (CP-violation, time-dep, LFUV and LFV obs. . .)