

# A gauged horizontal $SU(2)$ symmetry and $R_k$

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GDR-InF annual workshop – 11/05/2018

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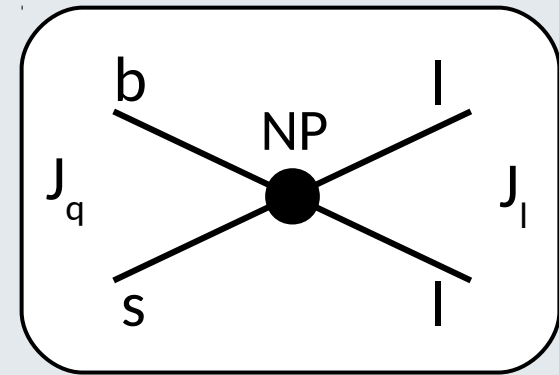
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Diego Guadagnoli & Olcyr Sumensari



# Introduction

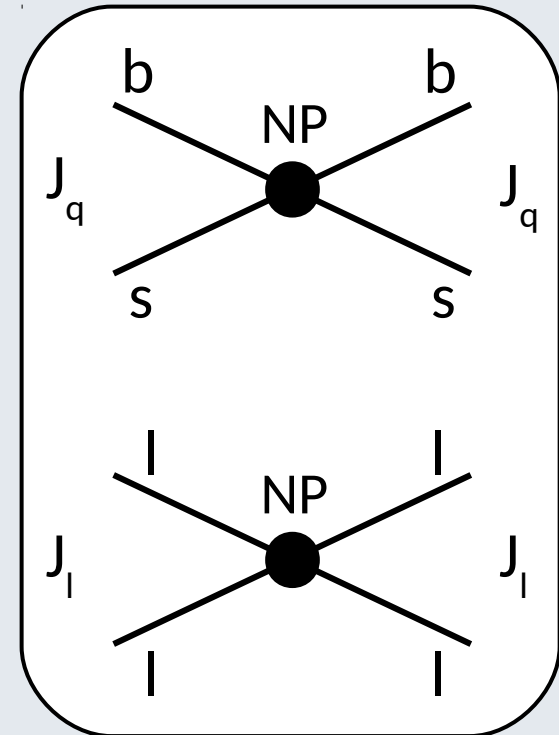
**B anomalies** can be explained by  
a **shift in  $C_9$  and  $C_{10}$**   
→ **15-20% correction to  $J_q \times J_l$**



**But:**

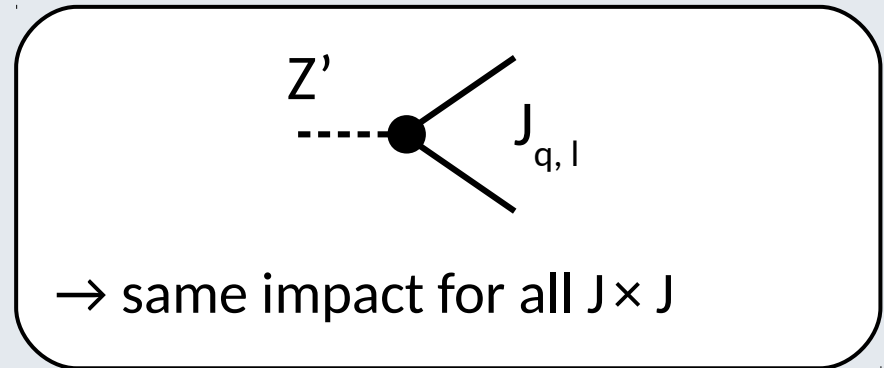
- $\Delta M_s \sim (\Delta M_s)_{SM}$   
→ **small correction to  $J_q \times J_q$**

- Lepton decays < experimental limit  
→ **small correction to  $J_l \times J_l$**

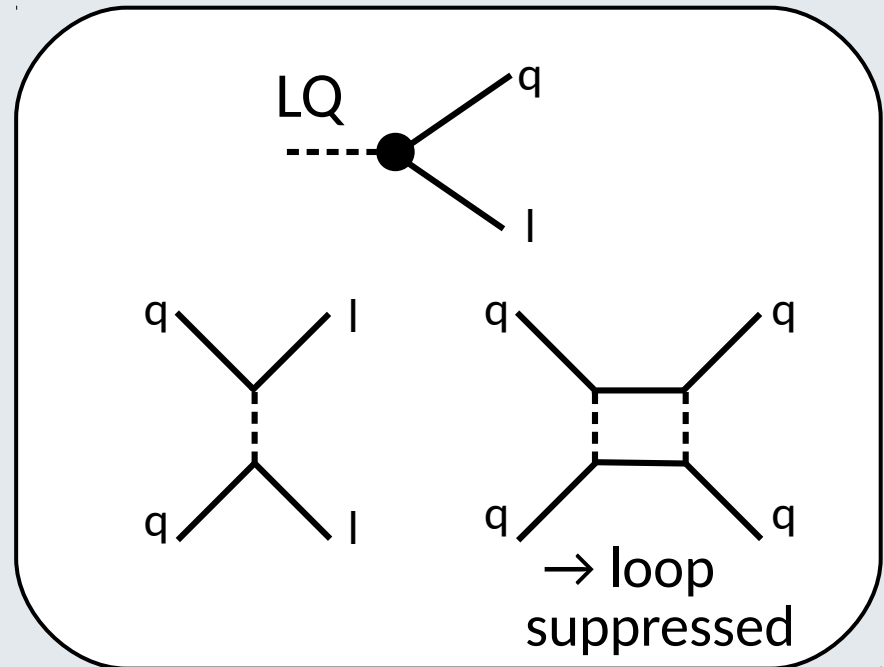


# Introduction – LQ

This can be **challenging** for  $Z'$  **models...**

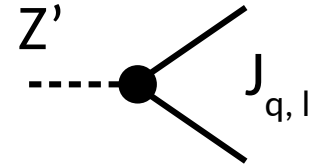


...but it makes **leptoquark** solutions very **appealing**



# Introduction – Question

This can be **challenging for  $Z'$  models...**



→ same impact for all  $J \times J$

... Can we get the same suppression **within gauge extension?**

# A gauged horizontal SU(2)

1) Add a **horizontal group**  $G_H$ , e.g. SU(2), acting on the **2<sup>nd</sup> and 3<sup>rd</sup> generations**

$$G_{SM} = SU(3)_c \times SU(2)_L \times U(1)_Y \times SU(2)_H$$

$$\delta L = \sum_{F,a} \bar{F} \gamma_\mu (g_L P_L G_L^{\mu a} + g_R P_R G_R^{\mu a}) \tau^a F$$

$$\text{with } F = \begin{pmatrix} f_2 \\ f_3 \end{pmatrix} \text{ i.e. } \begin{pmatrix} s_L \\ b_L \end{pmatrix}, \begin{pmatrix} u_L \\ \tau_L \end{pmatrix}, \begin{pmatrix} s_R \\ b_R \end{pmatrix} \dots$$

2) Integrate out heavy degrees of freedom

$$\delta L_{eff} = \sum_{F_1, F_2, a} \frac{g_L^2}{2M_{G_L a}^2} (\bar{F}_1 \gamma_\mu P_L \tau_a F_1) (\bar{F}_2 \gamma^\mu P_L \tau_a F_2)$$

# Basic argument

3) Rotate the field to the mass basis  $F = U_F \hat{F}$

$$\delta L_{\text{eff}} = \sum_{F_1, F_2, a} \frac{g_L^2}{2M_{G_L a}^2} (\overline{\hat{F}}_1 U_{F_1}^+ \gamma_{\mu L} \tau_a U_{F_1} \hat{F}_1) (\overline{\hat{F}}_2 U_{F_2}^+ \gamma_L^\mu \tau_a U_{F_2} \hat{F}_2)$$

For  $F_1 = F_2$ , if masses are degenerate,  $U_F$  can be absorbed in a **G basis redefinition**

→ **Only flavor diagonal terms!**

# Minimal setup

**But:**

Rotation also involve the **first generation**

→ Small **corrections** to  $K-\bar{K}$  mixing,  $\mu \rightarrow 3e\dots$  wich are **very constrained**

Generalize to 3 generations:  $F = \begin{pmatrix} u_L \\ s_L \\ b_L \end{pmatrix} \dots \quad F = U_F \hat{F}$

Apply all constraints:

- CKM =  $U_{uL}^+ U_{dL}$
- meson mixings
- lepton decays
- $R_K$

# Minimal setup

**But:**

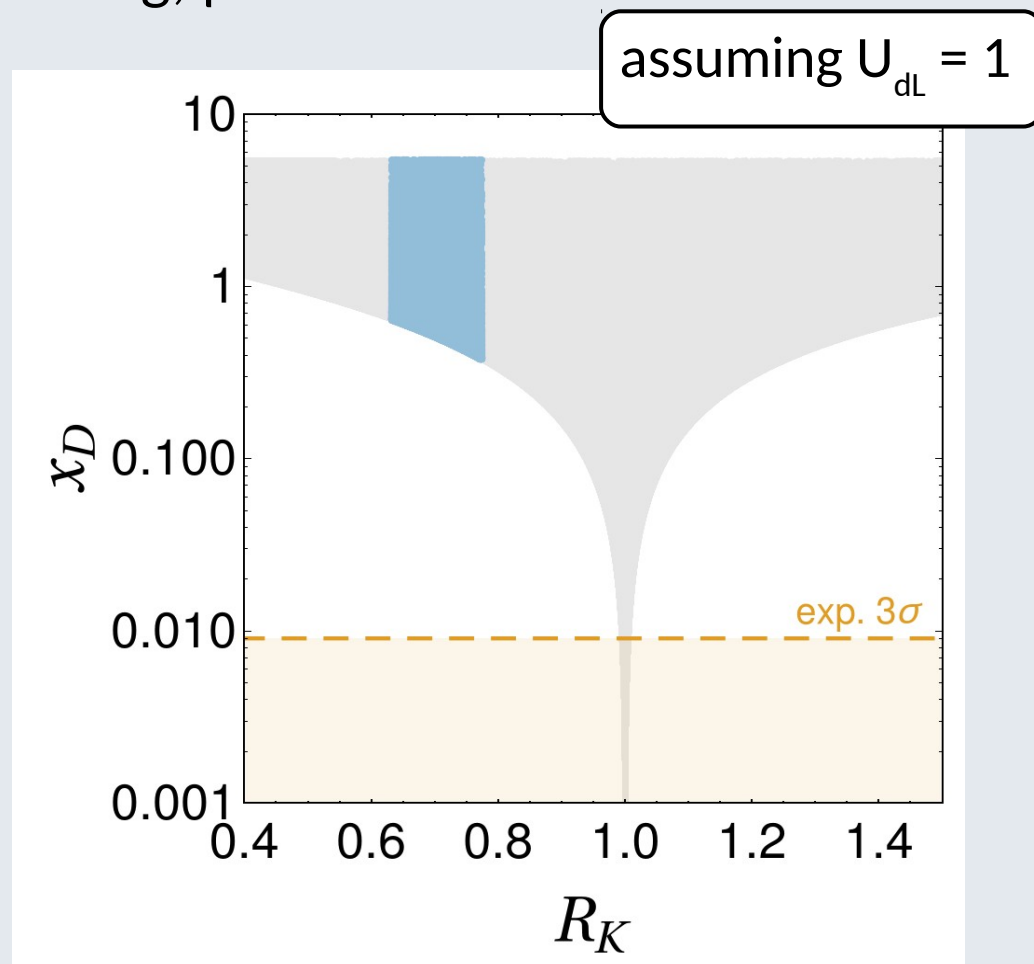
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Apply all constraints:

- CKM =  $U_{uL}^+ U_{dL}$
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# Non-degenerate masses

A direct generalization is obtained with **non-degenerate masses**:

$$M_{G_{L1}} = M_{G_{L2}} \ll M_{G_{L3}}$$

[-] This breaks the original SU(2) symmetry

[+] It passes all the experimental constraints:

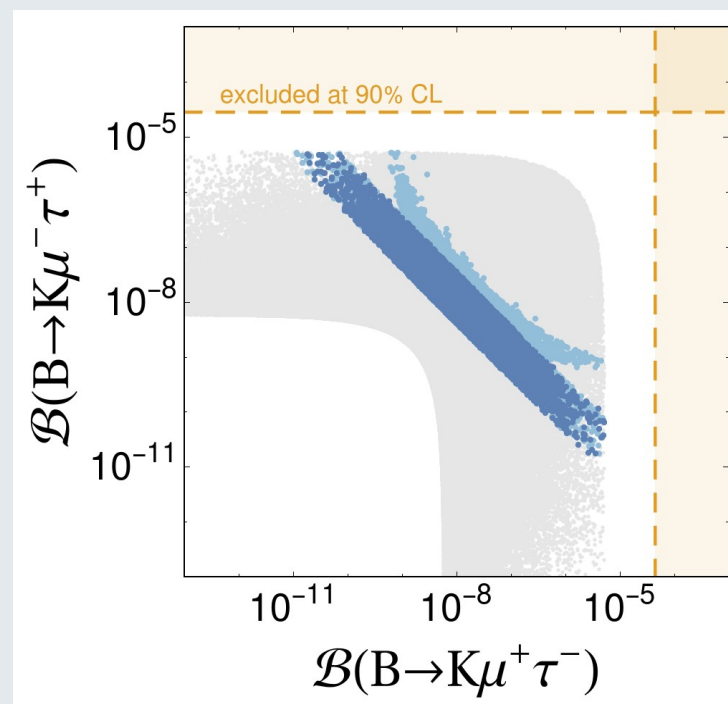
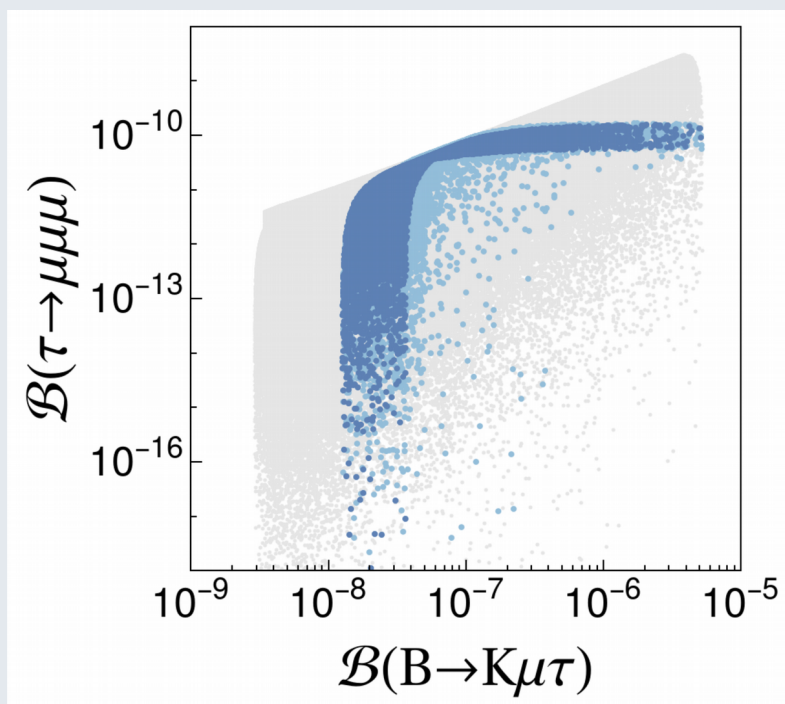
- $b \rightarrow sll$  measurements
- Meson mixings
- Small shift to  $B \rightarrow K\nu\nu$  (due to underlying SU(2) symmetry)
- Small shifts to  $\tau \rightarrow l\nu\nu$  and  $D^0 \rightarrow \mu\mu$

# Non-degenerate masses

A direct generalization is obtained with **non-degenerate masses**:

$$M_{G_{L1}} = M_{G_{L2}} \ll M_{G_{L3}}$$

- [-] This breaks the original SU(2) symmetry
- [+] It passes all the experimental constraints
- [+]  $\delta C_{9,10}^{\tau\tau} = -\delta C_{9,10}^{\mu\mu}$  and LFV predictions



Adding a **SU(2) horizontal symmetry** allows

- sizeable correction to  $J_q \times J_l$
- while keeping  $J_q \times J_q$  and  $J_l \times J_l$  small enough

The choice of SU(2) is motivated by the **absence of anomalies**, but one can consider a larger group

- Require **extra matter**
- May solve **other problems**

**Thank you!**

## Back-up

# Fermion masses

Accommodating quark masses and the SU(2) symmetry can be done by adding heavy fermions that mixes with the SM ones.

$$\begin{aligned}\psi_U^{L,R} &= (3, 1, 2/3; 2) & F_Q &= \begin{pmatrix} Q_L^2 \\ Q_L^3 \end{pmatrix} \\ \psi_D^{L,R} &= (3, 1, -1/3; 2) \\ \Phi_{1,2} &= (1, 1, 0; 2)\end{aligned}$$

$$L_{mass} = m_U \overline{\psi}_U^L \psi_U^R + m_D \overline{\psi}_D^L \psi_D^R$$

$$\begin{aligned}L_{mix} &= \sum_{a=1,2} [(\mathbf{y}_a^U)_i \overline{\psi}_U^L \Phi_a (u_R)_i + (u \Leftrightarrow d)] \\ &+ c_U \overline{F}_Q \tilde{H} \psi_U^R + c_D \overline{F}_Q H \psi_D^R + \text{h.c.}\end{aligned}$$