

New physics @ horizon

Mind the Cap !

Iosif Bena

IPhT, CEA Saclay

with

Nick Warner, Emil Martinec, Jan deBoer, Micha Berkooz, Simon Ross, Stefano Giusto, Rodolfo Russo, Guillaume Bossard, Masaki Shigemori, Monica Guică, Nikolay Bobev, Bert Vercnocke, Andrea Puhm, David Turton, **Stefanos Katmadas**, **Ruben Monten**, **Daniel Mayerson**, **Johan Blåbäck**, **Pierre Heidmann**



FQXi

JOHN TEMPLETON
FOUNDATION

Agence Nationale de la Recherche
ANR

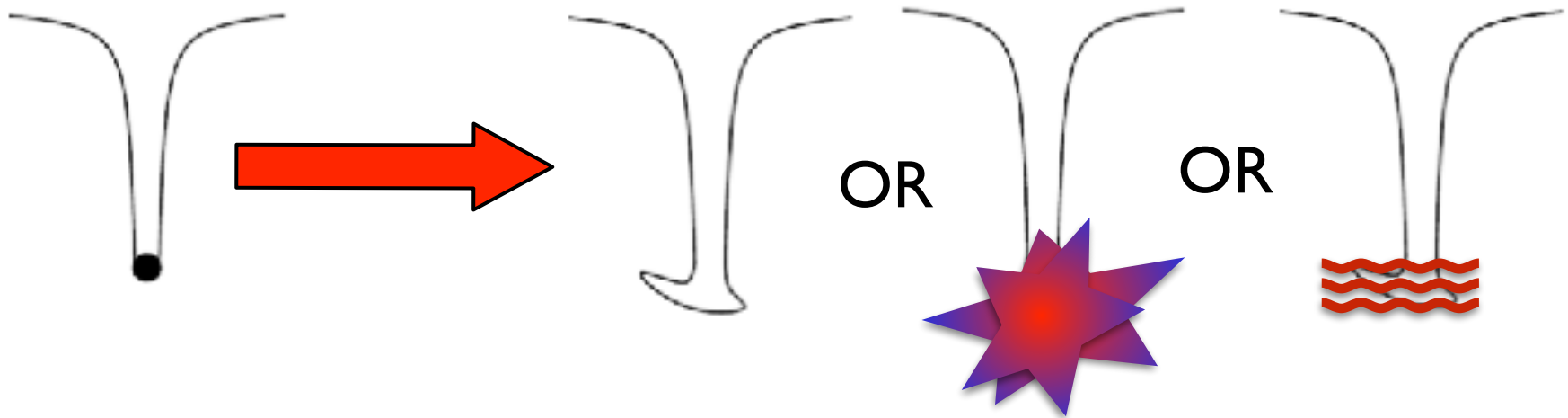
DE LA RECHERCHE À L'INDUSTRIE
cea
SACLAY

BLACK HOLES: **CONFLICT** between **General Relativity** and **Quantum Mechanics**

There must exist **STRUCTURE @ horizon** scale

Mathur 2009, Almheiri, Marolf, Polchinski, Sully 2012

Fuzzball / Firewall / etc.



How does **STRUCTURE @ horizon** look like ?

our work over the past 15 years

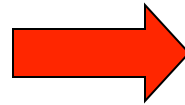
Analogy with ideal gas

Thermodynamics

(Air = ideal fluid)

$$P V = n R T$$

$$dE = T dS + P dV$$



Statistical Physics

(Air -- molecules)

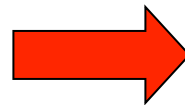
e^S microstates

typical

atypical

Thermodynamics

Black Hole Solution



Statistical Physics

Microstate geometries

Long distance physics

Gravitational lensing

Physics at horizon

Information loss

Gravity waves ?

Word of caution

- Everybody&their brother - replace classical BH by BH-sized object
 - Gravastar, quark-star, boson-star
 - Infinite density firewall hovering just above horizon
 - Gas of wormholes
 - Bose-Einstein condensate of gravitons
 - LQG configuration...

3 very stringent tests:

1. Same growth with G_N !!!

Horowitz

BH size **grows** with G_N ; “normal objects” **shrink**

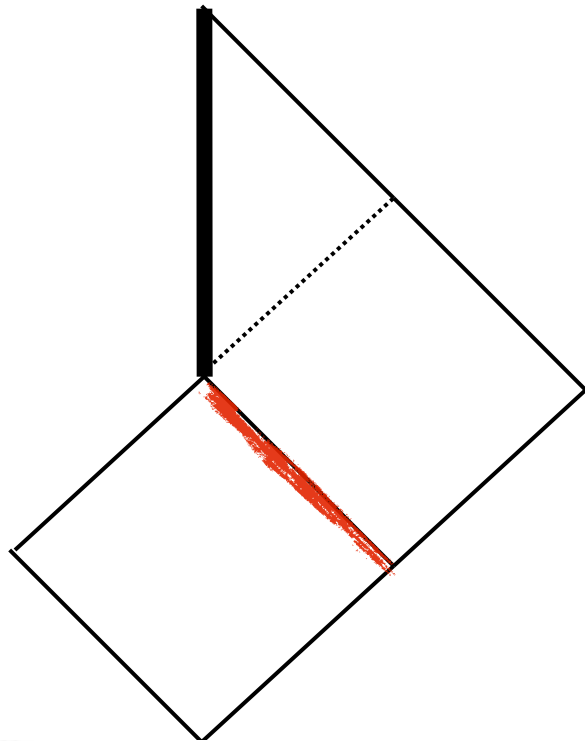
- BH **microstate** geometries **pass** this test
- **Highly nontrivial** mechanism: $G_N = g_s^2$
- D-branes = solitons, **tension** $\sim 1/g_s \rightarrow$ lighter as G_N increases



To build structure@horizon, non-perturbative degrees of freedom you must use !

2. Mechanism not to fall into BH

Very difficult !!!



GR Dogma:

**Thou shalt not put anything
at the horizon !!!**

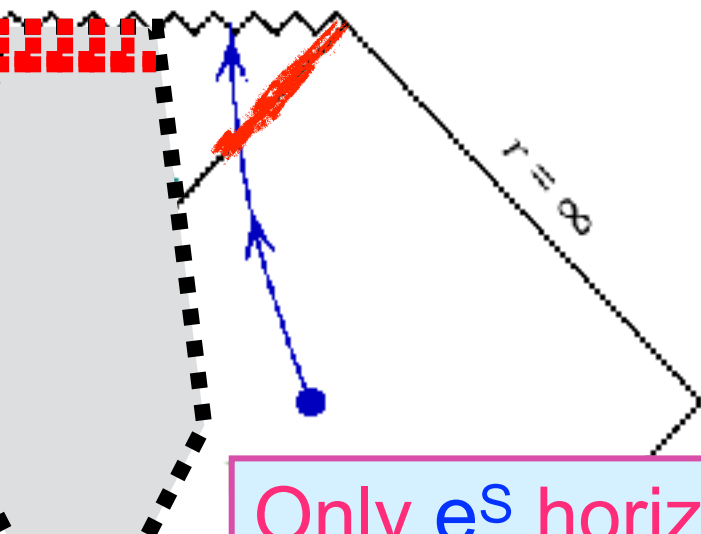
- Null \rightarrow speed of light.
- If massive: ∞ boost \rightarrow ∞ energy
- If massless: dilutes with time
- Nothing can live there !
(or carry degrees of freedom)
- No membrane, no spins, no “quantum stuff”
- No (fire)wall

*If support mechanism have you not,
b.s. you are doing*



3. Avoid forming a horizon

- Collapsing shell forms horizon Oppenheimer and Snyder (1939)
- If curvature is low, no reason not to trust classical GR
- By the time shell becomes **curved-enough for quantum effects to become important**, horizon in causal past (60 hours for NGC 4889 BH)



Backwards in time - **illegal** !

BH has e^S microstates with no horizon

Small tunneling probability = e^{-S}

Will tunnel with probability **ONE** !!!

Kraus, Mathur; Bena, Mayerson, Puhm, Vercnocke

Only e^S horizon-sized microstates can do it !

Black hole entropy the structure must have



Microstates geometries

- Where is the BH charge ?

$$L = q A_0$$

magnetic

$$L = \dots + A_0 F_{12} F_{34} + \dots$$

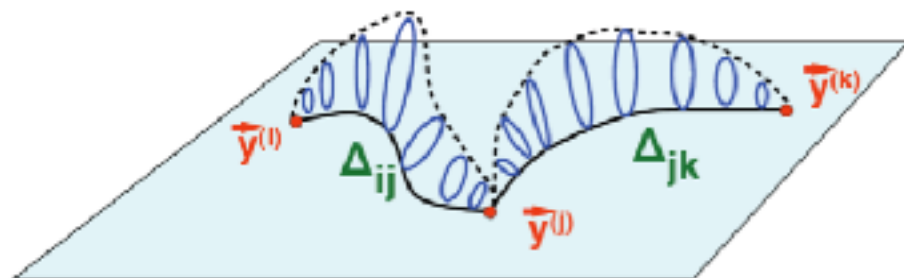
- Where is the BH mass ?

$$E = \dots + F_{12} F^{12} + \dots$$

- BH angular momentum

$$J = E \times B = \dots + F_{01} F_{12} + \dots$$

2-cycles + magnetic flux



Bubbling Geometries

The charge is dissolved in magnetic fluxes. No singular sources.

Largest family of solutions known to mankind

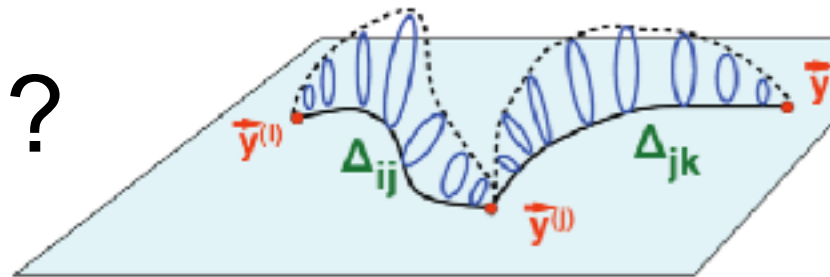
Arbitrary functions of **two** variables: $\infty \times \infty$ parameters
Bena, Giusto, Russo, Shigemori, Warner

$$\begin{aligned}
 ds_{10}^2 &= \frac{1}{\sqrt{\alpha}} ds_6^2 + \sqrt{\frac{Z_1}{Z_2}} ds_4^2, \\
 ds_6^2 &= -\frac{2}{\sqrt{\mathcal{P}}} (dv + \beta) \left[du + \omega + \frac{\mathcal{F}}{2} (dv + \beta) \right] + \sqrt{\mathcal{P}} ds_4^2, \\
 e^{2\phi} &= \frac{Z_1^2}{\mathcal{P}}, \\
 B &= \frac{Z_4}{\mathcal{P}} (du + \omega) \wedge (dv + \beta) + a_4 \wedge (dv + \beta) + \delta_2, \\
 C_0 &= \frac{Z_4}{Z_1}, \\
 C_2 &= \frac{Z_2}{\mathcal{P}} (du + \omega) \wedge (dv + \beta) + a_1 \wedge (dv + \beta) + \gamma_2, \\
 C_4 &= \frac{Z_4}{Z_2} \widehat{\text{vol}}_4 - \frac{Z_4}{\mathcal{P}} \gamma_2 \wedge (du + \omega) \wedge (dv + \beta) + x_3 \wedge (dv + \beta) + \mathcal{C}, \\
 C_6 &= \widehat{\text{vol}}_4 \wedge \left[-\frac{Z_1}{\mathcal{P}} (du + \omega) \wedge (dv + \beta) + a_2 \wedge (dv + \beta) + \gamma_1 \right] \\
 &\quad + \frac{Z_4}{\mathcal{P}} \mathcal{C} \wedge (du + \omega) \wedge (dv + \beta), \\
 \alpha &\equiv \frac{Z_1 Z_2}{Z_1 Z_2 - Z_4^2}, \quad \mathcal{P} \equiv Z_1 Z_2 - Z_4^2.
 \end{aligned}$$

$$\begin{aligned}
 & - \frac{Rr}{\sqrt{2} k_2 (m_1^2 - 1)} \frac{m_1 (k_2 + m_1 + 1) \Delta_{k_2+m_1-1, m_1-1} + (k_2 + m_1 - 1) \Delta_{k_2+m_1}}{(r^2 + a^2)^2} \\
 & - \frac{R}{\sqrt{2} k_2 (m_1^2 - 1) a^2 \sin \theta \cos \theta} \left[2(m_1 - 1) \Delta_{k_2+m_1-3, m_1-1} \right. \\
 & \quad + (m_1 - 1)(m_1 - 2) \Delta_{k_2+m_1-1, m_1-1} + m_1 (k_2 - 2) \Delta_{k_2+m_1-1, m_1+1} \\
 & \quad \left. - m_1 (m_1 - 1) \Delta_{k_2+m_1+1, m_1-1} + (m_1^2 (k_2 - 1) + 1) \Delta_{k_2+m_1+1, m_1+1} \right], \\
 & - \frac{R}{\sqrt{2}} \frac{\Delta_{k_2+m_1+1, m_1+1}}{\Sigma} \sin^2 \theta - \frac{R}{\sqrt{2} k_2 (m_1^2 - 1) a^2} \left[2(m_1 - 1) \Delta_{k_2+m_1-3, m_1-1} \right. \\
 & \quad + (m_1^2 - 2m_1 + k_2 - 1) \Delta_{k_2+m_1-1, m_1-1} + m_1 (k_2 - 2) \Delta_{k_2+m_1-1, m_1+1} \\
 & \quad + m_1 (k_2 - m_1 - 1) \Delta_{k_2+m_1+1, m_1-1} + (k_2 (m_1^2 + m_1 - 1) - m_1 (m_1 + 1)) \Delta_{k_2+m_1+1, m_1+1} \\
 & \quad \left. - \frac{R}{\sqrt{2}} \frac{\Delta_{k_2+m_1+1, m_1+1}}{\Sigma} \cos^2 \theta - \frac{R}{\sqrt{2} k_2 (m_1^2 - 1) a^2} \left[(k_2 - 1)(m_1 - 1) \Delta_{k_2+m_1-1, m_1-1} \right. \right. \\
 & \quad - 2(m_1 - 1) \Delta_{k_2+m_1-3, m_1-1} - (m_1 - 1)(m_1 - 2) \Delta_{k_2+m_1-1, m_1-1} \\
 & \quad + (m_1 - 1)(k_2 - 3) \Delta_{k_2+m_1-1, m_1+1} + m_1 (m_1 - 1) \Delta_{k_2+m_1+1, m_1-1} \\
 & \quad \left. \left. + (m_1 - 1)(m_1 (k_2 - 1) + 1) \Delta_{k_2+m_1+1, m_1+1} \right] \right].
 \end{aligned}$$

Habemus Superstratum !!!

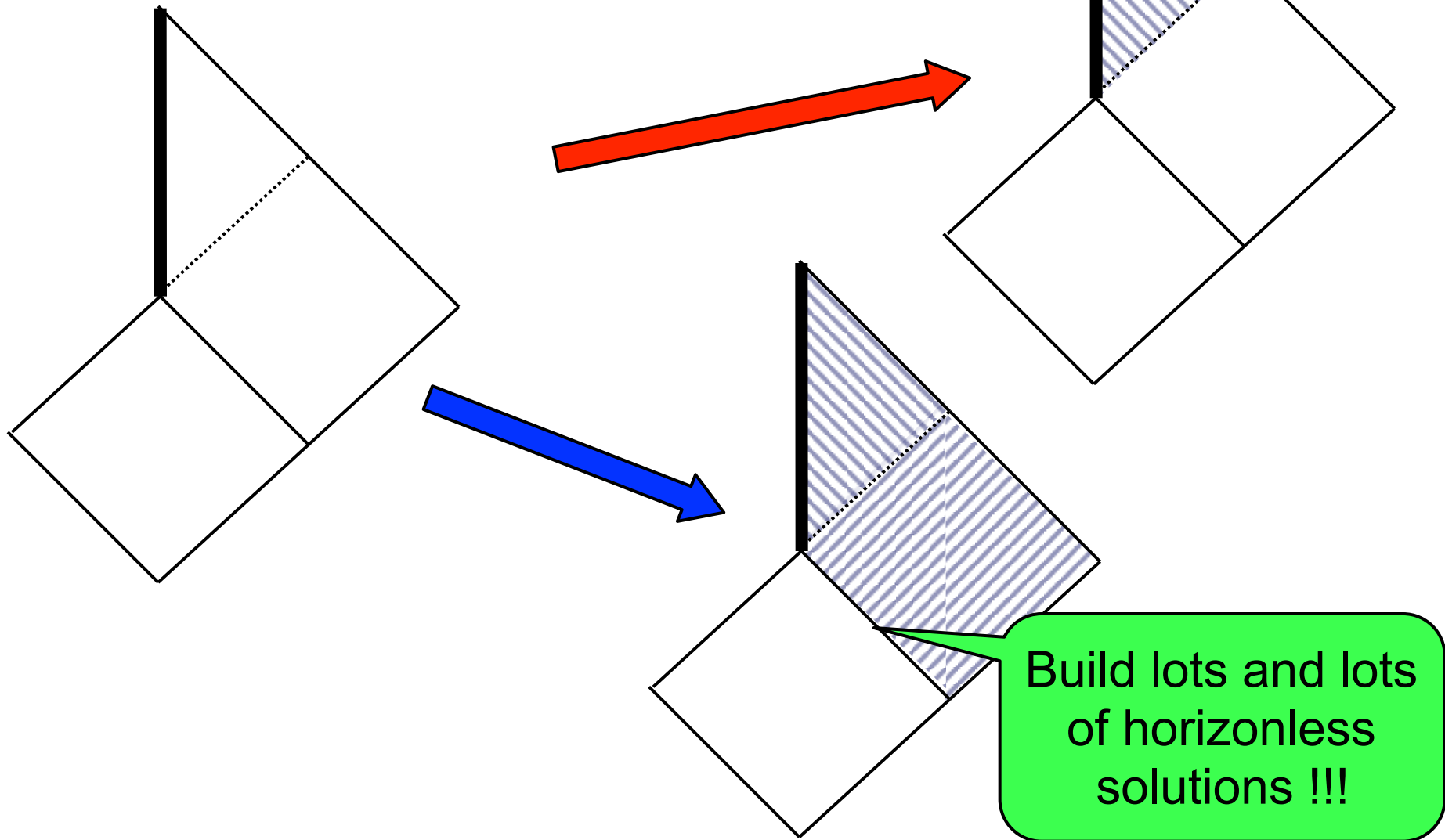
Why not collapsing ?



- 5(+6)d : smooth solutions + **quantized** magnetic flux on topologically-nontrivial **2-cycles**
 - cycles smaller \rightarrow increases energy
 - bubbling = **only** mechanism to avoid collapse in semiclassical limit Gibbons, Warner
 - If **any** state in the **e^S -dimensional** BH Hilbert space has a semiclassical limit, it **must** be a microstate geometry !
- 4(+6)d : multicenter solutions Denef
 - smooth GH centers with negative charge \rightarrow centers with **negative D6 charge** and **negative mass**
 - common in String Theory (e.g. orientifolds); **nowhere else**
 - **Highly unusual** matter from a 4d perspective
 - Usual matter does not hang around, just falls in BH

$Q = M$: extremal BH

Need non-extremal



Very few known. Extremely hard to build...

– Coupled nonlinear 2nd order PDE's do not factorize

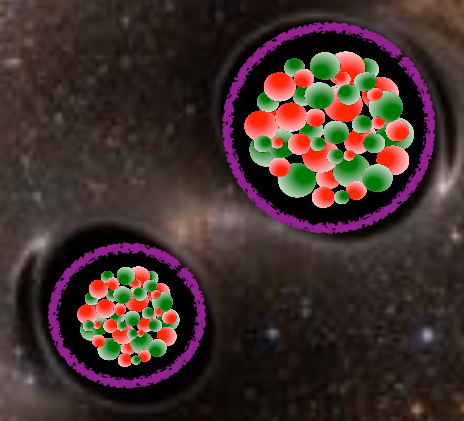
- Near-extremal Bena, Puhm, Verhocke
- A few solutions far from extremality with large angular momenta Bena, Bossard, Katmadas, Turton
- Extremal Kerr (NHEK) Heidmann (upcoming paper)

Universal feature:

Yuge™ amount of new degrees of freedom @ horizon

- Effect on gravity waves ?
- Supermassive BH formation easier ?

What difference do they make?



Horizon viscosity

Microstate mountains

Distortion of Love numbers / Kerr multipoles

Horizon viscosity

more d.o.f - more dissipative effects (electron moving parallel to a conducting plate)

reflection

lots of d.o.f. - expect **complete absorption**
(no reflection *Ashfordi & friends*)



Already taken into account by GR BH ?

Microstate mountains

microstates non-spherical (mountains on neutron stars). Jitters on infalling stuff



Typical microstates may be very spherical ?

Planck-size bumps ?

Distortion of Kerr multipole moments :

spinning ball of dust or liquid

spinning solid shell

spinning BH with vacuum at horizon

spinning **STRUCTURE @ horizon** scale

effect on collisions with non-aligned spins

Love numbers : response to tidal stress Cardoso&al
similar story



Fuzzball complementarity: unusual matter forming **STRUCTURE @ horizon** gives “classical horizon experience” to infalling observer. Same as vacuum.