

Turbulence of Gravitational Waves

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IAP, October 2018

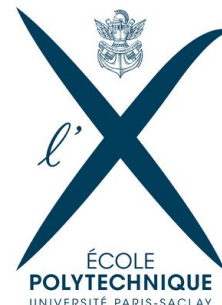


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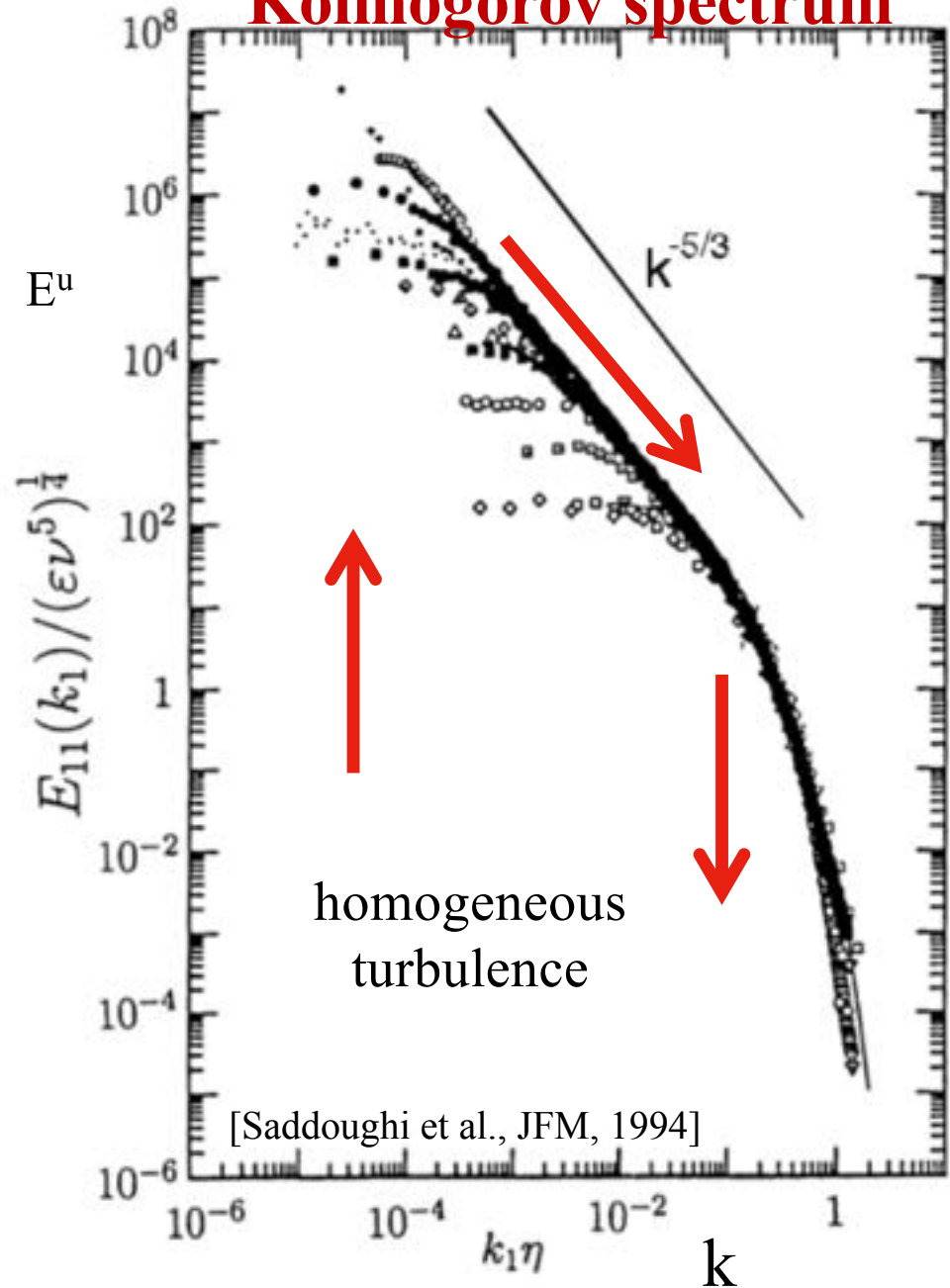
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Turbulence in 3D hydrodynamics



Kolmogorov spectrum



Turbulence in 3D hydrodynamics

Incompressible
Navier-Stokes equations

$$\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla P + \nu \Delta \mathbf{u} + \mathbf{f}$$

Fourier
transform

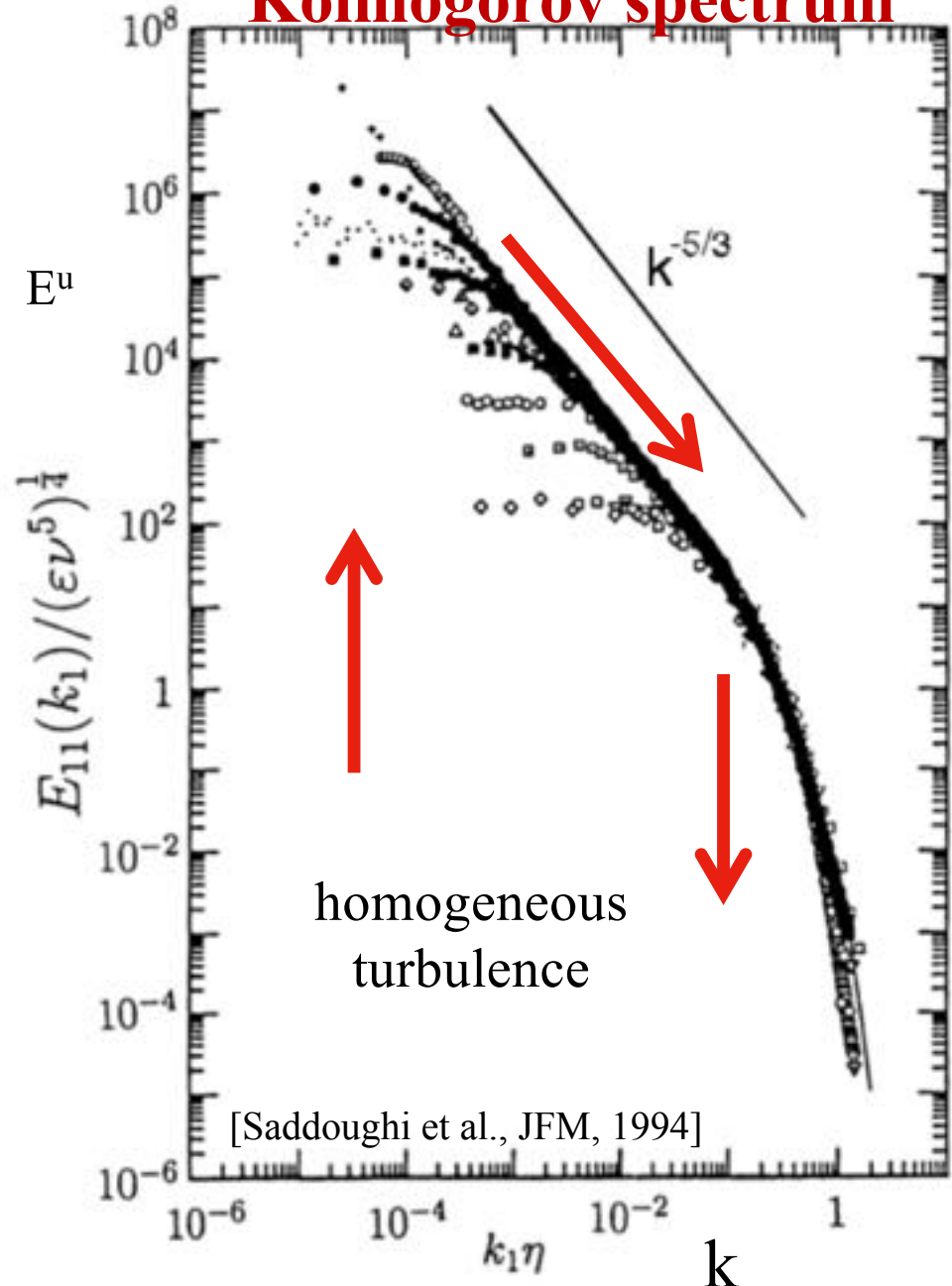
$$\partial_t \hat{u}_j(\mathbf{k}) = -i \int \left(\delta_{jn} - \frac{k_j k_n}{k^2} \right) q_m \hat{u}_m(\mathbf{p}) \hat{u}_n(\mathbf{q}) \delta(\mathbf{k} - \mathbf{p} - \mathbf{q}) d\mathbf{p} d\mathbf{q}$$

triadic interactions

nonlinear
analysis

$$\partial_t E(k) = NL - 2\nu k^2 E(k) + F(k)$$

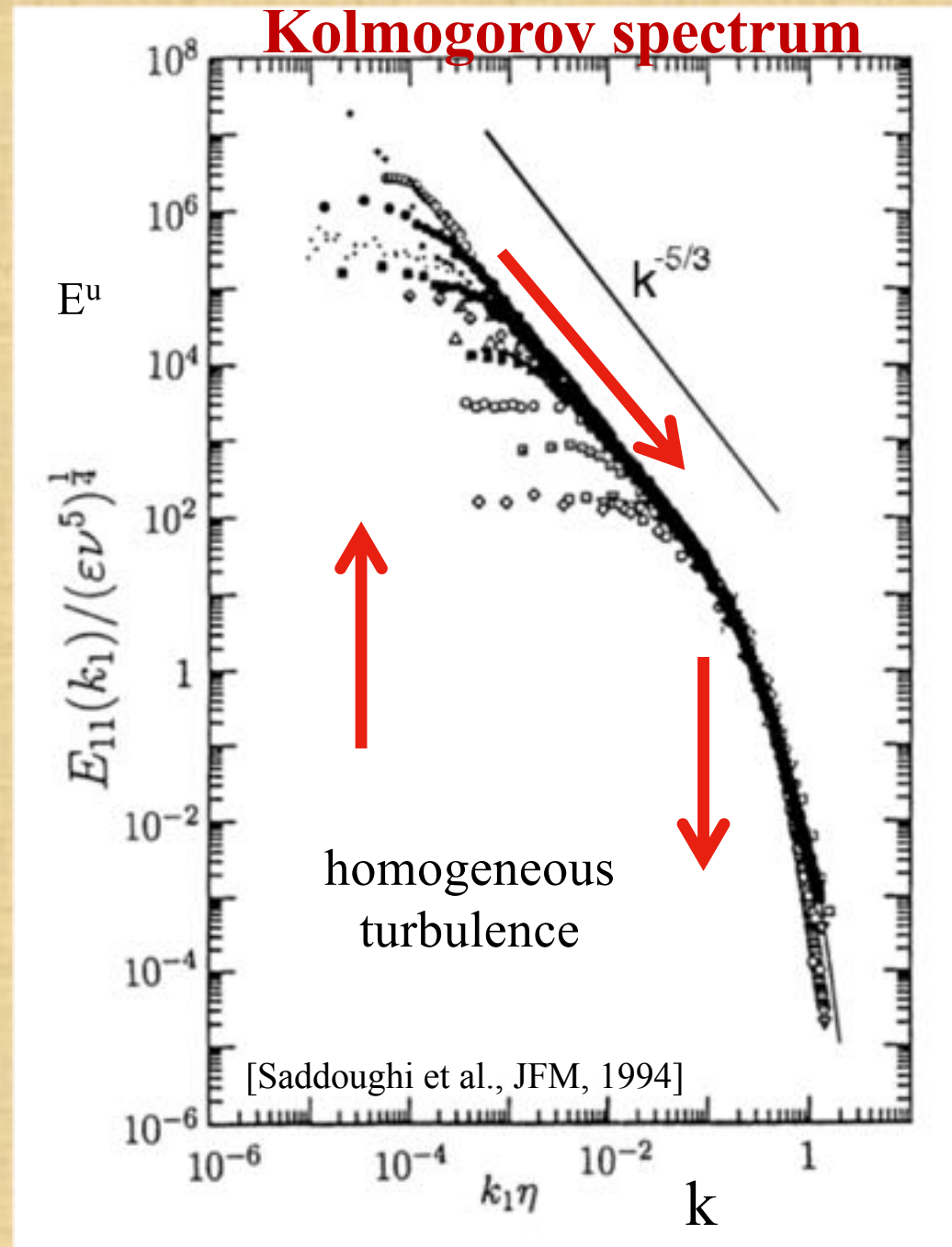
Kolmogorov spectrum



Turbulence in 3D hydrodynamics



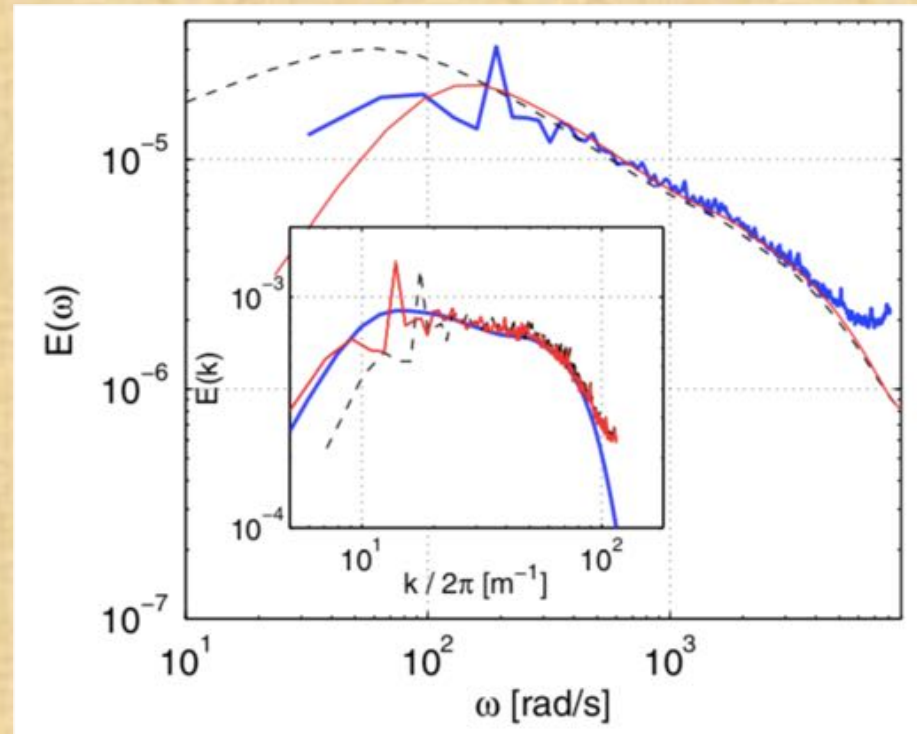
Wave turbulence



The sound of wave turbulence: the vibrating elastic plates

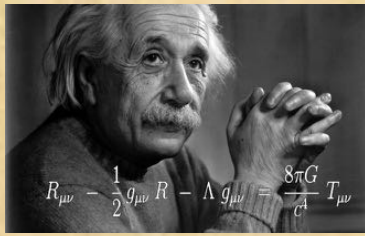


[During+, PRL, 2006; Cobelli+, PRL, 2009;
Mordant+, PRE, 2017; Haissaini+, 2018]



Analytical theory of weak wave turbulence

- + Natural **asymptotical closure** of the hierarchy of moment equations
[Benney & Saffman, PRSLA, 1966; Benney & Newell, JMP, 1967]
- + The kinetic equations admits **exact stationary finite flux** solutions
[Zakharov & Filonenko, SPD, 1967]
- Finite flux spectra not valid **over all k 's** \rightarrow strong turbulence
[SG+, JPP, 2000; Meyrand+, PRL, 2016]
- Experiments and dns show **some limitations** in the predictions
[Morize+, PoF, 2005; Nazarenko, NJP, 2007]



Gravitational waves

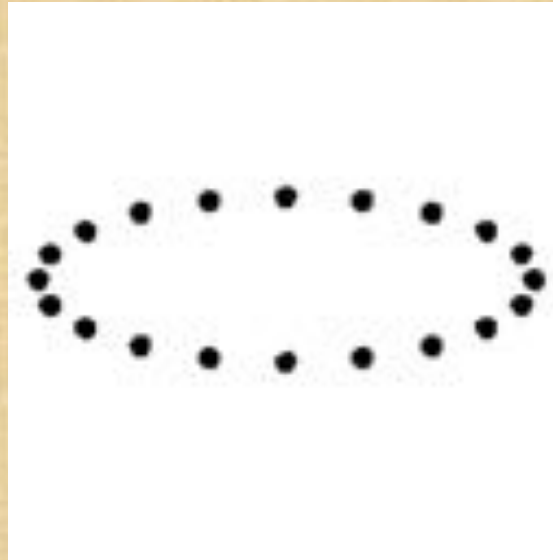
$$\Lambda=0$$

Exact linear solutions in an empty – flat – Universe:

$$R_{\mu\nu} = 0$$

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, \text{ where } h_{\mu\nu} \ll 1$$

Effect of a + gravitational wave
on a ring of particles
(h = 0.5)



$$\omega_{\mathbf{k}} = c|\mathbf{k}| = ck$$

$$h_{\mu\nu}^+ = a \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

Weakly nonlinear general relativity

$$\Lambda=0$$

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, \text{ where } h_{\mu\nu} \ll 1$$

$$R_{\mu\nu} = 0$$

$$R_{\mu\nu} = R_{\mu\nu}^{(1)} + R_{\mu\nu}^{(2)} + \dots$$

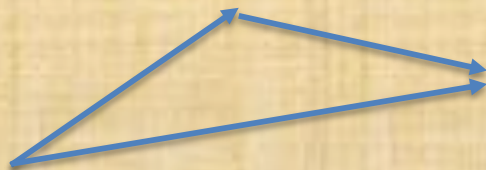
$$R_{\mu\nu}^{(1)} = -\frac{1}{2}\square h_{\mu\nu}$$

Triadic interactions:

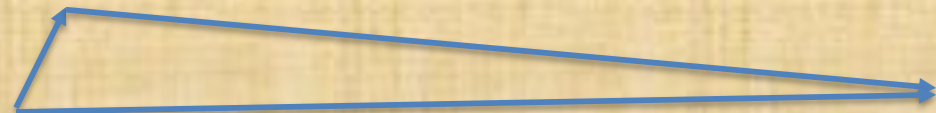
$$\mathbf{k} = \mathbf{k}_1 + \mathbf{k}_2 \text{ and } \omega_{\mathbf{k}} = \omega_{\mathbf{k}_1} + \omega_{\mathbf{k}_2}$$

$$\omega_{\mathbf{k}} = c|\mathbf{k}| = ck$$

$$\tau_{\text{lin}} \ll \tau_{\text{nonlin}}$$



local interactions



non-local interactions

We found that there is no contribution on the **resonant manifold**

Three-wave interactions in GW turbulence **are absent** !

Nonlinear general relativity equations

Einstein-Hilbert action:

$$S = \frac{1}{2} \int R \sqrt{-g} d^4x$$

g is the determinant of $g_{\mu\nu}$

R is the scalar curvature

Diagonal space-time metric:

$$\partial/\partial z = 0$$

[Hadad & Zakharov, JGP, 2014]

$$g_{\mu\nu} = \begin{pmatrix} -(H_0)^2 & 0 & 0 & 0 \\ 0 & (H_1)^2 & 0 & 0 \\ 0 & 0 & (H_2)^2 & 0 \\ 0 & 0 & 0 & (H_3)^2 \end{pmatrix}$$

$$H_0 = e^{-\lambda}\gamma, \quad H_1 = e^{-\lambda}\beta, \quad H_2 = e^{-\lambda}\alpha, \quad H_3 = e^{\lambda}$$

Spatial isotropy *is not* assumed

Lagrangian density

$$\Rightarrow \mathcal{L} = \frac{1}{2} \left[\underbrace{\frac{\alpha\beta}{\gamma} \dot{\lambda}^2 - \frac{\alpha\gamma}{\beta} (\partial_x \lambda)^2 - \frac{\beta\gamma}{\alpha} (\partial_y \lambda)^2}_{\text{Give the linear contribution}} - \frac{\dot{\alpha}\dot{\beta}}{\gamma} + \frac{(\partial_x \alpha)(\partial_x \gamma)}{\beta} + \frac{(\partial_y \beta)(\partial_y \gamma)}{\alpha} \right]$$

Give the linear contribution

$$\alpha = \beta = \gamma = 1 \quad \lambda \ll 1$$

$$\lambda = c_1 \exp(-i\omega_{\mathbf{k}} t + i\mathbf{k} \cdot \mathbf{x}) + c_2 \exp(i\omega_{\mathbf{k}} t + i\mathbf{k} \cdot \mathbf{x})$$

Hamiltonian formalism

Normal variables: $\lambda_{\mathbf{k}} = \frac{a_{\mathbf{k}} + a_{-\mathbf{k}}^*}{\sqrt{2k}}, \quad \dot{\lambda}_{\mathbf{k}} = \frac{\sqrt{k}(a_{\mathbf{k}} - a_{-\mathbf{k}}^*)}{i\sqrt{2}},$ Fourier space

Hamiltonian equation: $i\dot{a}_{\mathbf{k}} = \frac{\partial H}{\partial a_{\mathbf{k}}^*}$ where $H = H_{\text{free}} + H_{\text{int}}$

$$H_{\text{free}} = \sum_{\mathbf{k}} k |a_{\mathbf{k}}|^2$$

4 wave processes

$$\begin{aligned}
 H_{\text{int}} = & \frac{1}{4} \sum_{1,2,3,4,5} \frac{\delta_{123} \delta_{45}^1}{\sqrt{k_2 k_3 k_4 k_5}} \left\{ \left[\left(\frac{p_5}{p_1} + \frac{q_5}{q_1} \right) k_4 \left(-\frac{a_4 a_5 + a_{-4}^* a_{-5}^*}{k_5 + k_4} + \frac{a_{-4}^* a_5 + a_4 a_{-5}^*}{k_5 - k_4} \right) + \frac{p_4 q_5}{p_1 q_1} (a_4 + a_{-4}^*)(a_5 + a_{-5}^*) \right] \right. \\
 & k_2 k_3 (a_2 - a_{-2}^*)(a_3 - a_{-3}^*) + \left[- \left(\frac{p_5}{p_1} - \frac{q_5}{q_1} \right) (p_2 p_3 - q_2 q_3) k_4 \left(-\frac{a_4 a_5 + a_{-4}^* a_{-5}^*}{k_5 + k_4} + \frac{a_{-4}^* a_5 + a_4 a_{-5}^*}{k_5 - k_4} \right) \right. \\
 & \left. \left. + \frac{p_4 q_5}{p_1 q_1} (\mathbf{k}_2 \cdot \mathbf{k}_3) (a_4 + a_{-4}^*)(a_5 + a_{-5}^*) \right] (a_2 + a_{-2}^*)(a_3 + a_{-3}^*) \right\} \\
 & + \frac{1}{2} \sum_{\mathbf{k}, 1,2,3,4} \frac{\delta_{12}^{\mathbf{k}} \delta_{34}^{\mathbf{k}}}{\sqrt{k_1 k_2 k_3 k_4}} \left\{ \frac{(\mathbf{k} \cdot \mathbf{k}_2) k_1 p_3 q_4}{pq} \left(-\frac{a_1 a_2 + a_{-1}^* a_{-2}^*}{k_2 + k_1} + \frac{a_{-1}^* a_2 + a_1 a_{-2}^*}{k_2 - k_1} \right) (a_3^* + a_{-3})(a_4^* + a_{-4}) \right. \\
 & \left. + \frac{k_1 k_3 p_2 q_4}{pq} (a_1 a_2 + a_1 a_{-2}^* - a_{-1}^* a_2 - a_{-1}^* a_{-2}^*) (a_3^* a_4^* + a_3^* a_{-4} - a_{-3} a_4^* - a_{-3} a_{-4}) \right\}. \quad (6)
 \end{aligned}$$

Kinetic equation of GW turbulence

$$H_{3 \rightarrow 1} = 0$$

$$\dot{n}_{\mathbf{k}} = 4\pi \int |T_{\mathbf{k}_1 \mathbf{k}_2}^{\mathbf{k} \mathbf{k}_3}|^2 n_{\mathbf{k}_1} n_{\mathbf{k}_2} n_{\mathbf{k}_3} n_{\mathbf{k}} \left[\frac{1}{n_{\mathbf{k}}} + \frac{1}{n_{\mathbf{k}_3}} - \frac{1}{n_{\mathbf{k}_1}} - \frac{1}{n_{\mathbf{k}_2}} \right] \delta(\mathbf{k} + \mathbf{k}_3 - \mathbf{k}_1 - \mathbf{k}_2) \delta(\omega_{\mathbf{k}} + \omega_{\mathbf{k}_3} - \omega_{\mathbf{k}_1} - \omega_{\mathbf{k}_2}) d\mathbf{k}_1 d\mathbf{k}_2 d\mathbf{k}_3,$$

$$n_{\mathbf{k}} = \langle |a_{\mathbf{k}}|^2 \rangle$$

$$\text{with } T_{34}^{12} = \frac{1}{4}(W_{34}^{12} + W_{34}^{21} + W_{43}^{12} + W_{43}^{21}), \quad W_{34}^{12} = Q_{34}^{12} + Q_{12}^{34}$$

$$Q_{34}^{12} = \frac{1}{4\sqrt{k_1 k_2 k_3 k_4}} \left\{ 2 \left(\frac{p_4}{p_1 - p_3} - \frac{q_4}{q_1 - q_3} \right) \frac{k_2(p_1 p_3 - q_1 q_3)}{k_1 - k_3} - 2 \left(\frac{p_4}{p_1 - p_3} + \frac{q_4}{q_1 - q_3} \right) \frac{k_1 k_2 k_3}{k_1 - k_3} \right. \\ \left. + \left(\frac{p_2}{p_1 + p_2} - \frac{q_2}{q_1 + q_2} \right) \frac{k_1(p_3 p_4 - q_3 q_4)}{k_1 + k_2} - \left(\frac{p_2}{p_1 + p_2} + \frac{q_2}{q_1 + q_2} \right) \frac{k_1 k_3 k_4}{k_1 + k_2} + \frac{2k_1 k_3 p_2 q_4}{(p_1 + p_2)(q_1 + q_2)} + \frac{2k_1 p_3(q_2 k_4 + k_2 q_4)}{(p_1 - p_3)(q_1 - q_3)} \right\}. \quad (12)$$

[Zakharov & Filonenko, SPD, 1967]

[SG & Nazarenko, PRL, 2017]

Constant flux (stationary) solutions:

Energy

$$\mathcal{E} = \iint \omega_{\mathbf{k}} n_{\mathbf{k}} d\mathbf{k} = \text{const}$$

$$E_k^{(1D)} \sim \varepsilon^{1/3} k^0$$

Direct cascade

Wave action

$$\mathcal{N} = \iint n_{\mathbf{k}} d\mathbf{k} = \text{const}$$

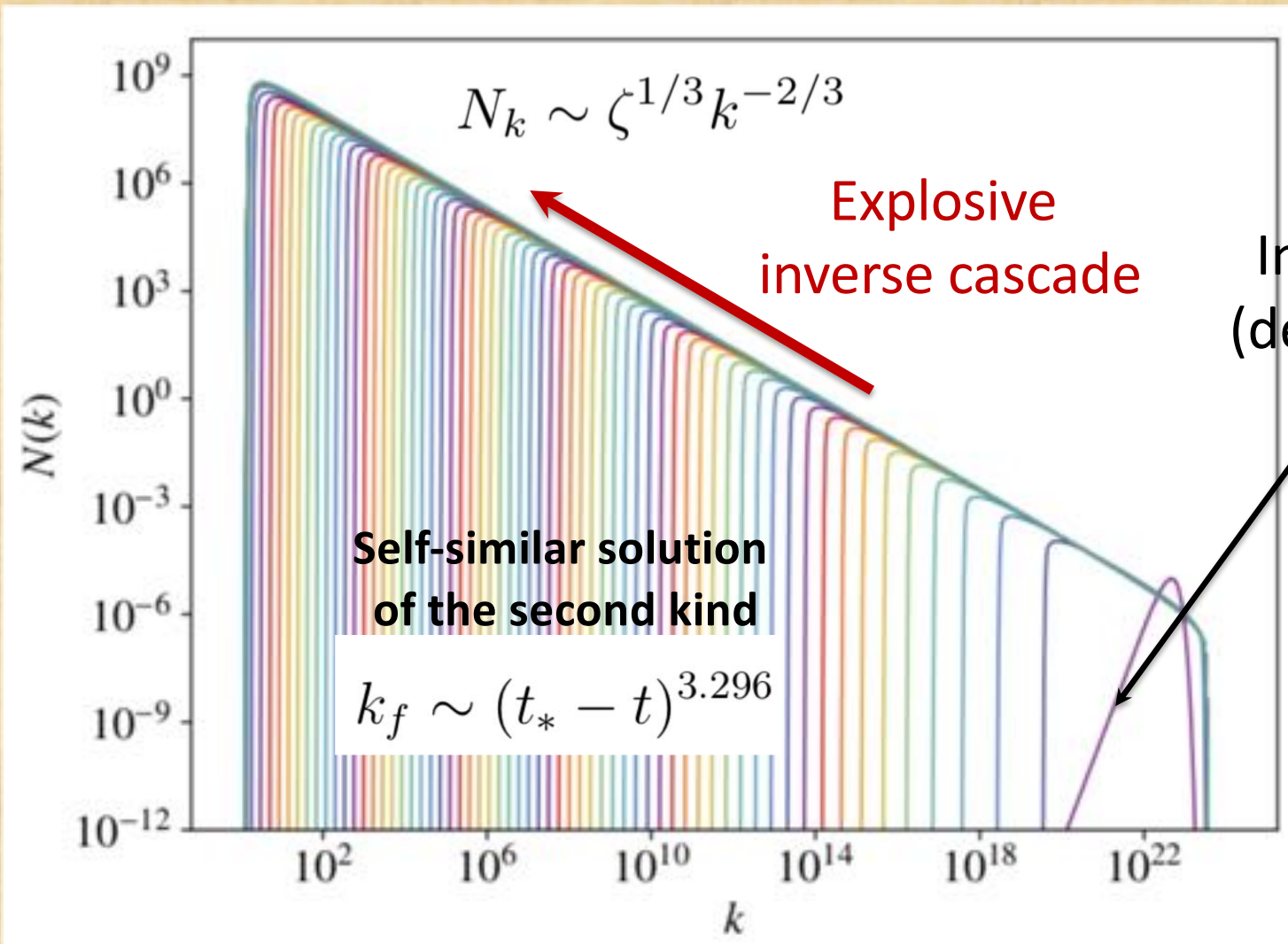
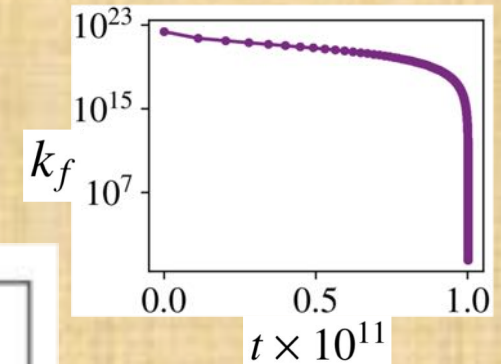
$$N_k^{(1D)} \sim \zeta^{1/3} k^{-2/3}$$

Inverse cascade

Isotropic
spectra

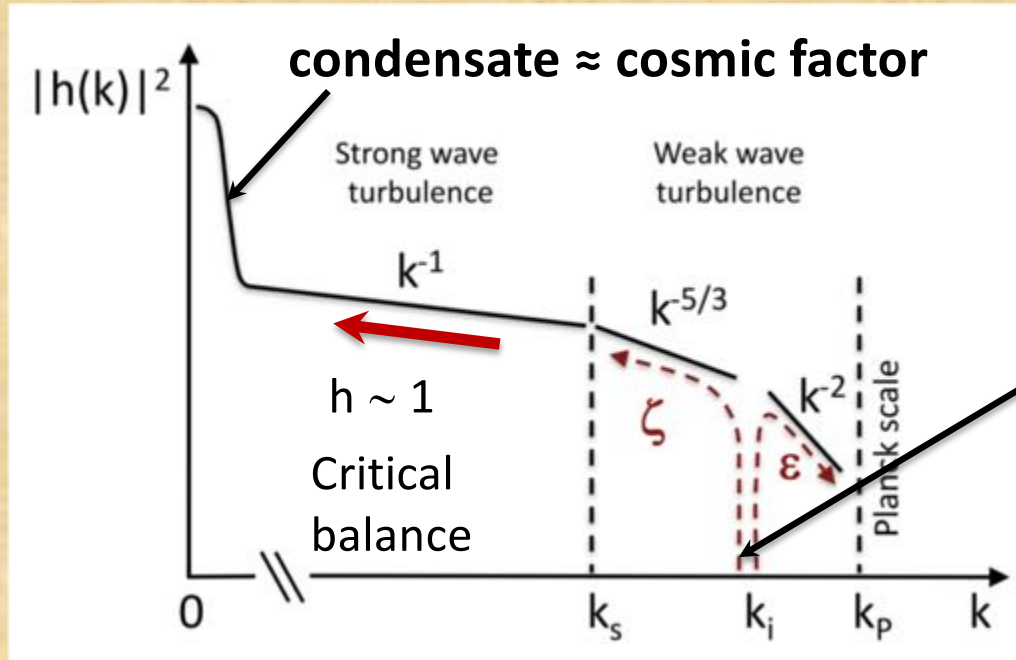
Numerical simulation in the local approximation

$$\frac{\partial N(k)}{\partial t} = \frac{\partial}{\partial k} \left[k^2 N^2(k) \frac{\partial (kN(k))}{\partial k} \right] - \nu k^4 N(k) - \eta \frac{N(k)}{k^4}$$

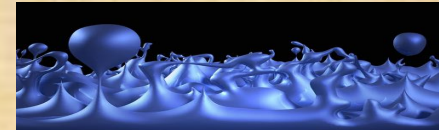


Big-Bang scenario driven by space-time turbulence

Conjecture



Quantum foam at Planck time



$$h \leq 0.1$$

$$h \sim |h_{\mu\nu}|$$

Formation of a condensate in a finite time ($10^{-39}\text{s} - 10^{-37}\text{s}$)
(compatible with the causal principle)

Explosive growth of the condensate as **a nonlinear mechanism of accelerated expansion** ('turbulent inflation')

Conclusion 1

- First analytical theory of **space-time turbulence**
- No tuning parameters
- The Riemann (4th order) curvature tensor is non-trivial
the Kretschmann scalar is non trivial
- Application to the **early universe** (peace of the story; conjecture)
- Strong wave turbulence as the main driver for the formation of a **condensate** in a finite time (compatible with the causal principle)
- A rapid growth of the condensate is interpreted as an **inflation**
- $|h(k)|^2 \sim k^{-1}$ is compatible with the Harrison-Zeldovich spectrum
- Can we perform **numerical simulations** of space-time turbulence?
- What is the origin of the inverse cascade in turbulent black holes?

Conclusion 2

Some comments:

- There is one Killing vector in this problem; this not a cylindrical GW
- Although an Alfvén wave is a nonlinear solution of incompressible MHD, weak Alfvén wave turbulence exists because the problem deals with a sea of waves