Colloque national dark Energy – 2ème édition

The dark side of gravity and the acceleration of the Universe

October 23-25 2018, Paris IAP

F Henry-Couannier CPPM & Aix-Marseille Univ Dark Gravity theories are extensions of General Relativity aiming at a stable anti-gravitational sector

J.P. Petit, Twin Universe cosmology, Astrophys. Space Sci. Vol. 226, pp 273, 1995. and many other articles

F. Henry-Couannier, Discrete symmetries and General Relativity, the Dark Side of Gravity, Int.J.Mod.Phys, vol. A20, no. NN, pp. 2341-2346, 2004.

F. Henry-Couannier, Dark Gravity, GJSFR A. Vol 13, Issue 3, pp 1-53, 2013.

S. Hossenfelder, Bimetric theory with exchange symmetry Phys. Rev. D 78, 044015, 2008.

M. Milgrom, Matter and twin matter in bimetric MOND, MNRAS 405 (2), pp 1129-1139, 2010.

Laura Bernard, Luc Blanchet, Lavinia Heisenberg Bimetric gravity and dark matter 50th Rencontres de Moriond, "Gravitation: 100 years after GR", 2015

From background dependence to Dark Gravity (DG) How far can we go ?

$$GR: g_{\mu\nu}$$

$$DG: g_{\mu\nu} \text{ and } \eta_{\mu\nu}$$

$$Riemm(\eta_{\mu\nu}) = 0$$

$$\Rightarrow g_{\mu\nu} \text{ has a twin, $``} the inverse metric $``} \tilde{g}_{\mu\nu}$$

$$\tilde{g}_{\mu\nu} = \eta_{\mu\rho}\eta_{\nu\sigma} \left[g^{-1}\right]^{\rho\sigma}$$

 $\Rightarrow (g_{\mu\nu}, \tilde{g}_{\mu\nu})$ is a Janus field



From the Action to DG field equations

The Action must respect the permutation symmetry between $g_{\mu\nu}$ and $\tilde{g}_{\mu\nu}$:

$$\int d^4x (\sqrt{g}R + \sqrt{\tilde{g}}\tilde{R}) + \int d^4x (\sqrt{g}L + \sqrt{\tilde{g}}\tilde{L})$$
$$\delta g_{\mu\nu} \Rightarrow \delta S = 0$$

 $\sqrt{g}\eta^{\mu\sigma}g_{\sigma\rho}G^{\rho\nu} - \sqrt{\tilde{g}}\eta^{\nu\sigma}\tilde{g}_{\sigma\rho}\tilde{G}^{\rho\mu} + \mu \leftrightarrow \nu = -8\pi G(\sqrt{g}\eta^{\mu\sigma}g_{\sigma\rho}T^{\rho\nu} - \sqrt{\tilde{g}}\eta^{\nu\sigma}\tilde{g}_{\sigma\rho}\tilde{T}^{\rho\mu} + \mu \leftrightarrow \nu)$ Contracted form

$$\sqrt{g}R - \sqrt{\tilde{g}}\tilde{R} = 8\pi G(\sqrt{g}T - \sqrt{\tilde{g}}\tilde{T})$$

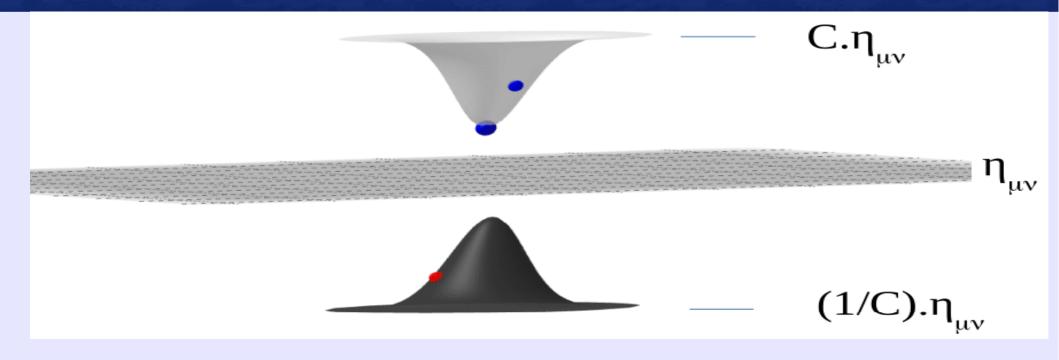
Implications of DG equations

- DG is background dependent yet deviations from GR can remain arbitrarily small provided one side of the Janus Field dominates the other
- Ghost interaction between Janus and source fields but Janus field not understood to be a quantum field !
 - DG more natural than GR as a semiclassical* theory of gravity
 - Semiclassical DG stability : OK**
- New discrete (permutation) symmetry is very fundamental : will be interpreted as a global time reversal symmetry.

* https://arxiv.org/abs/0802.1978 Mark Albers, Claus Kiefer, Marcel Reginatto, Measurement Analysis and Quantum Gravity : « Despite the many physical arguments which speak in favor of a quantum theory of gravity, it appears that the justification for such a theory must be based on empirical tests and does not follow from logical arguments alone »

** https://arxiv.org/pdf/1401.4024.pdf V. A. Rubakov, page 8 : Gradient, tachyonic and ghost instabilities in scalar-tensor theories : « for ghosts, background is QM unstable but classically stable »

The static isotropic solution



- Antigravity without run away !
- Asymptotic C matters : GR corresponds to C infinite

The static isotropic solution

C=∞ RG (Schwarzschild) :

DG:

$$g_{ii}(r) = A = e^{2MG/r} \approx 1 + 2\frac{MG}{r} + 2\frac{M^2G^2}{r^2}$$

$$-g_{00}(r) = \frac{1}{A} = e^{-2MG/r} \approx 1 - 2\frac{MG}{r} + 2\frac{M^2G^2}{r^2} - \frac{4}{3}\frac{M^3G^3}{r^3}$$

- No Horizon
- Zero Gravitational Waves

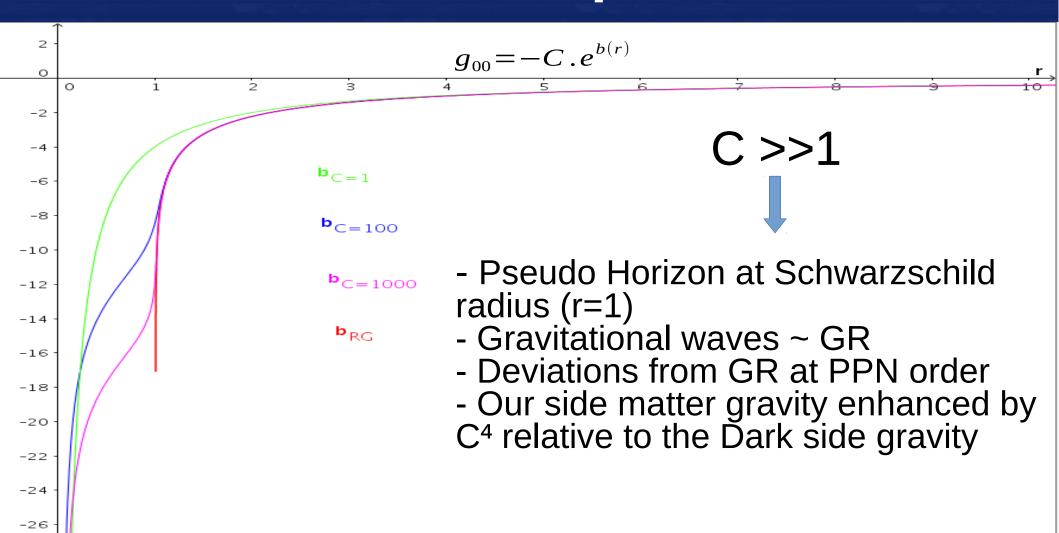
$$h_{\mu\nu} = -h_{\mu\nu} + O(h^2)$$

$$2(R^{(1)}_{\mu\nu} - \frac{1}{2}\eta_{\mu\nu}R^{(1)\lambda}_{\lambda}) = -8\pi G(T_{\mu\nu} - \tilde{T}_{\mu\nu} + t_{\mu\nu} - \tilde{t}_{\mu\nu})$$

Deviations from GR at PPN order only

$$g_{ii}(r) = \left(1 + \frac{MG}{2r}\right)^4 \approx 1 + 2\frac{MG}{r} + \frac{3}{2}\frac{M^2G^2}{r^2}$$
$$g_{00}(r) = \frac{\left(1 - \frac{MG}{2r}\right)^2}{\left(1 + \frac{MG}{2r}\right)^2} \approx 1 - 2\frac{MG}{r} + 2\frac{M^2G^2}{r^2} - \frac{3}{2}\frac{M^3G^3}{r^3}$$

The static isotropic solution



Cosmological equation

- Homogeous & isotropic Janus solution is flat and static : C was indeed a constant !
 - \Rightarrow We need to introduce a separate scalar- η Janus field for cosmology :

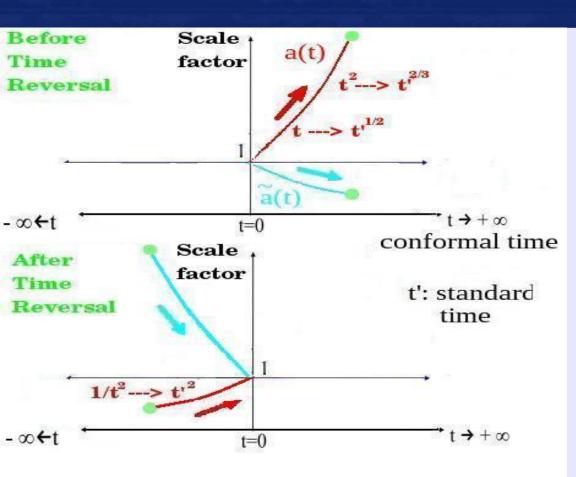
$$g_{\mu\nu} = \Phi \eta_{\mu\nu}$$
 and $\tilde{g}_{\mu\nu} = \frac{1}{\Phi} \eta_{\mu\nu}$ $\Phi(t) = a^2(t)$

• Single scale factor equation :

$$a\ddot{a} - \tilde{a}\ddot{\tilde{a}} = \frac{4\pi G}{3}(a^4(\rho - 3p) - \tilde{a}^4(\tilde{\rho} - 3\tilde{p}))$$

 $\tilde{a}(t) = \frac{1}{a(t)}$

Cosmological solutions



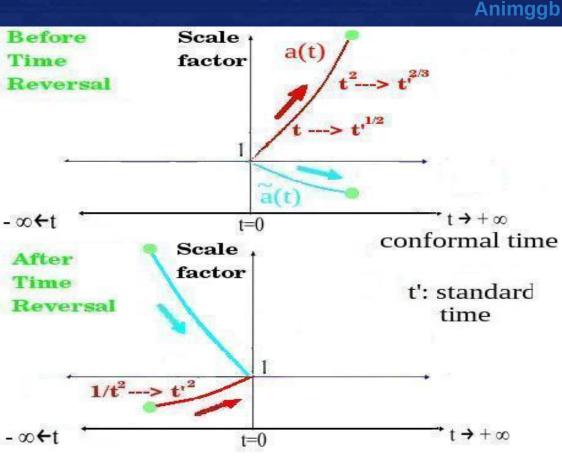
- Janus scale factors are related by a global conformal time reversal symmetry T : $\tilde{a}(t) = \frac{1}{a(t)} = a(-t)$
- Both continuous evolution and discontinuous permutation T allowed when $\rho 3p = \tilde{\rho} 3\tilde{p}$



Global time reversal : not going backward in time, but jumping to the opposite time !

A cyclic Universe ?

DG Cosmology



Hyp : $\rho \simeq \rho - 3p = \tilde{\rho} - 3\tilde{p} \simeq \tilde{\rho}$ occured at transition redshift triggering T and a'(t')~ t'^2

With H(t) continuous at the transition and assuming same universe age as in LCDM:

a'(t')~
$$t'^{\alpha} \Longrightarrow z_{\rm tr} = \left(\frac{2/3 - \alpha}{1 - \alpha}\right)^{\alpha} - 1$$

conformal time \Rightarrow $z_{tr} = 0.78$ vs observed $z_{tr} = 0.67 + 0.1$

- ~ Same scale factor evolution as in LCDM
- Without DE
- Inflation not needed to get k=0
- Without Big Bang singularity
- Cosmological DM still needed
- Dark side effects only since t_{tr} or near t=0

Problem statement

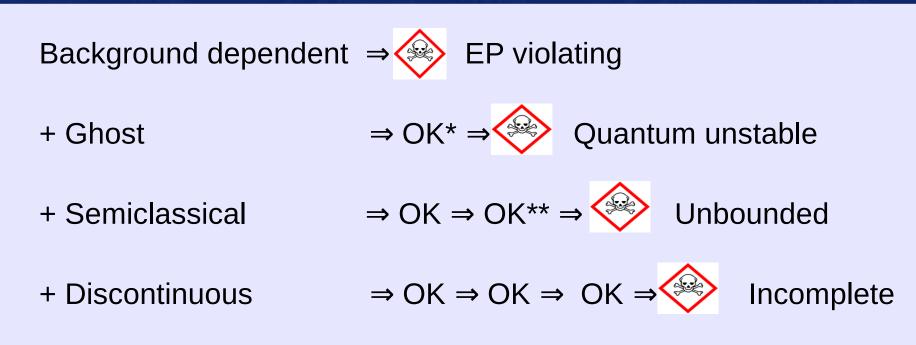
- We have two separate theories :
 - Asymptotically static DG correctly describes all aspects of gravity except expansion
 - Scalar-η Janus field only correctly describes expansion
- How to get expansion effects on the largest scales and differential eoms non trivially mixing background and perturbations (GR like) as needed to reproduce CMB phenomenology ?

Conclusion and outlooks

- DG avoids Big-Bang singularity and BH horizon very naturally
- Acceleration, k=0, large scale homogeneity, matter/antimatter asym
- Likely to cancel the gravity of vacuum energy
- Outlook :

Unification ⇒ New rich and effective phenomenology (DM candidate, ...) www.darksideofgravity.com/DG.pdf

How far could we go?



+ Emerging dynamics \Rightarrow OK \Rightarrow OK \Rightarrow OK \Rightarrow OK

* EP violations (η effects) usually small, **harmless classical instabilities

 \Rightarrow Fascinating phenomenological and theoretical implications !

Dynamical discrete symmetries

• Standard view :

Symmetries (cont & disc) \Rightarrow Action Extreme action principle \Rightarrow Eoms & conservation equations No dynamical processes associate with discrete symmetries

• Extended view :

Symmetries (cont & disc) → Action Extreme action principle → Eoms & conservation equations Discrete symmetries → Discontinuous processes

Dynamical discrete symmetries

- 1) Discrete (permutation) symmetry and continuous symmetries already unified in DG framework
- 2)Just as discrete (T&P) and continuous spacetime symmetries already unified in the Lorentz group
- 1) and 2) turn out to be related : global T symmetry is permutation symmetry !

Dynamical discrete symmetries \Rightarrow discontinuous transitions in addition to usual continuous evolution processes deduced from differential eoms.

- \Rightarrow Fills the gap between the discrete and the continuous
- \Rightarrow Hopefully opens the way to a genuine unification (understanding) of QM discrete and non local laws to the rest of physics !

Vacuum energy terms in DG equations

DG vacuum source term :

$$(\sqrt{g}\Lambda \!=\! \sqrt{ ilde{g}} ilde{\Lambda} \;) \, g^{\mu
u}$$

Cancels for $g_{\mu\nu} = \tilde{g}_{\mu\nu} = \eta_{\mu\nu}$ and $\Lambda = \tilde{\Lambda}$ (natural)

 \Rightarrow Might remain zero when Janus field starts to evolve, may be through the auto-adjustment of cut-offs to preserve compensation.

DG unification with adiabatic particles exchange? * adapted from original idea by Prigogin et al

• Matter and radiation fields conservation equations including adiabatic gravitationnally induced* transfers occuring between the two metrics :

$$\nabla_{\mathbf{v}} T^{\mathbf{v}}_{\mu} \neq \mathbf{0} \implies \frac{\dot{\rho}}{H} = (\frac{\Gamma}{H} - 3)(\rho + p)$$

$$\widetilde{\nabla}_{\mathbf{v}} \widetilde{T}^{\mathbf{v}}_{\mu} \neq \mathbf{0} \implies \frac{\dot{\tilde{\rho}}}{\tilde{H}} = (\frac{\tilde{\Gamma}}{\tilde{H}} - 3)(\tilde{\rho} + \tilde{p})$$

$$\widetilde{\Gamma} = -\Gamma, \tilde{H} = -H$$

• Replacing in DG_Friedmann equations

DG unification with Emerging Dynamics (ED)

As the universe evolves new dynamical dofs are released :

- Non dynamical
- Homogenous scalar-eta
- Scalar-eta + non dynamical fluctuation
- Separate dynamics

 $\eta_{\mu \nu}$ $\varphi(t)\eta_{\mu
u}$ $\varphi(t)\eta_{\mu\nu}$ + $\Delta g_{\mu\nu}$ $\varphi(t)\eta_{\mu
u}$ $\Delta g_{\mu
u}$ $a^{2}(t)$

ED : Early DG unification

• For a²(t) < Fundamental Threshold,

 $g_{\mu\nu}=\varphi(t)\eta_{\mu\nu}+\Delta g_{\mu\nu}$

but only the scalar $\varphi(t)$ is dynamical \Rightarrow we again get a single equation

 Symmetries related to our privileged coordinate system (rather than isometries related to the sources) force the primordial metrics in the Newtonian Gauge form :

$$d\tau^{2} = a^{2}(t)((1+2\Psi)dt^{2} - (1-2\Psi)d\sigma^{2})$$

 ⇒ We get the same scale factor (order 0) and potential (order 1) eoms as in GR but rotational and radiative modes should be absent from the CMB.

ED : Late DG unification

- a²(t) > Fundamental Threshold breaks the primordial symmetries
 - $\Rightarrow \varphi(t)\eta_{\mu\nu}$ and $\Delta g_{\mu\nu}$ start to play their dynamics independently
 - \Rightarrow Late DG unification required to account for expansion effects
- In the Linear domain, C (integration constant of $\Delta g_{\mu\nu}$) is driven step by step by the scale factor from $\varphi(t)\eta_{\mu\nu}$:

 \Rightarrow expansion effects through discrete rules

 \Rightarrow rich new and effective phenomenology related to field discontinuities

• In the Non Linear domain (solar system), we are asymptotically Minkowskian: C strictly constant !

Classical stability issues

- Background remains bounded thanks to global time reversal
- Linear inhomogeneous perturbations unstable in contracting phase but gravity from these is negligible : suppressed by C⁴ factor (~scale_factor⁸) before transition to acceleration.
- Linear inhomogeneous perturbations from the dark sector can start to grow under their own gravity after transition
- Strong gravity inhomogeneous pertubations presumably always stable on both sides thanks to C >1 at our side structures while C<1 at dark side structures

Problems with semiclassical Gravity

 Case I : Classical gravity triggers quantum collapses ⇒ no Energymomentum conservation violation, nor violation of uncertainty relations contrary to popular argument by Eppley & Hannah ...

https://arxiv.org/pdf/0802.1978.pdf

otherwise :

- Case 2A : No collapse interpretation of QM (MWI, decoherence ...) ruled out because classical gravity would see the uncollapsed superpositions
- Case 2B : Realistic collapse interpretation of QM leads to possible faster than light signaling. Either specific more local model of quantum collapse can solve this or ... DG : instantaneous signaling is not anymore a menace to causality as soon as there exists a unic privileged instantaneity frame for any collapse !