Can laboratory experiments help to characterize Dark Energy?

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GC observation @Keck

Colloque nat. Dark Energy 24 October 2018



Systèmes de Référence Temps-Espace

Global picture & motivations

- Some of the "greatest challenges" in theoretical physics:
 - what are Dark Matter and Dark Energy ?
 - how can we develop a quantum theory of gravity and/or unify it with the Standard Model of particles ?



Astronomy & cosmology

(Grav. waves, SNIa, CMB, structure formation, galactic dynamics, ...)

Local physics

(Solar System, **lab tests**, GNSS, ...)



Quantum Gravity

Unification

DM and **DE**

High energy

(particle physics: CERN-LHC, Fermilab, DESY, ...)



Picture inspired by Altschul et al, Adv, in Space Res. 55, 501, 2015

Why looking at lab scales?





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Models of DM or DE typically introduce new fields in this part of the action

Why looking at lab scales?





What are the impacts of these additional fields on local experiments?

A priori no reason why there would be no coupling with SM



Laboratory experiments have several advantages

- Very high accuracy/stability
- Extremely good control over the systematics
- Reproducibility (even by different labs)
- Can easily be optimized to search for specific signatures



"Affordable"

One example: UltraLight DM



Dark Matter can be made out of a bosonic scalar particles

A light scalar Dark Matter model

• A massive scalar field $V(arphi) \propto m^2 arphi^2$

$$S = \frac{1}{c} \int d^4x \frac{\sqrt{-g}}{2\kappa} \left[R - 2g^{\mu\nu} \partial_{\mu}\varphi \partial_{\nu}\varphi - V(\varphi) \right] + S_{\text{mat}} \left[g_{\mu\nu}, \varphi \right]$$

For low masses (< 0.1 eV), it behaves as a classical field and at cosmo scales:

$$\varphi = \varphi_0 \cos\left[\frac{mc^2}{\hbar}t\right]$$

• similar to pressureless fluid with $ho \propto m^2 arphi_0^2$

see e.g. Arvanitaki et al PRD, 2015 or Stadnik and Flambaum, PRL 2015

• can produce structure formation if $m > 10^{-23} eV$

see e.g. Marsh, Phys. Reports, 2016

This DM is expected to break the equivalence principle

• An effective Lagrangian for the scalar-matter coupling

$$\mathcal{L}_{\mathrm{mat}}[g_{\mu\nu},\Psi] = \mathcal{L}_{\mathrm{SM}}[g_{\mu\nu},\Psi] + \left(\frac{\varphi^{i}}{i}\left[\frac{d_{e}}{4e^{2}}F_{\mu\nu}F^{\mu\nu} - \frac{d_{g}\beta_{3}}{2g_{3}}F^{A}_{\mu\nu}F^{\mu\nu}_{A} - \sum_{i=e,u,d}d_{m_{i}} + \gamma_{m_{i}}d_{g}m_{i}\bar{\psi}_{i}\bar{\psi}_{i}\right]$$

see Damour and Donoghue, PRD, 2010

- Most usual couplings: linear (cfr Damour-Donoghue) or quadratic (cfr Stadnik et al)
- This leads to a space-time dependance of some constants of Nature to the scalar field

$$\alpha(\varphi) = \alpha \left(1 + d_e^{(i)} \frac{\varphi^i}{i} \right) ,$$

$$m_j(\varphi) = m_j \left(1 + d_{m_j}^{(i)} \frac{\varphi^i}{i} \right) \quad \text{for } j = e, u, d$$

$$\Lambda_3(\varphi) = \Lambda_3 \left(1 + d_g^{(i)} \frac{\varphi^i}{i} \right) ,$$

A signature of a violation of the Einstein Equivalence Principle that can be searched with atomic clocks!

Violation of the EEP can be searched with atomic sensors

 use of more than 8 years of Cs/Rb FO2 atomic fountain data from SYRTE: high accuracy and high stability

see J. Guéna et al, Metrologia, 2012 and J. Guéna et al., IEEE UFFC, 2012

• Search for a periodic signal in the data using Scargle's method, see Scargle ApJ, 1982







 $\log_{10} f_{\varphi}$ [Hz] 2 -10 -8 -6 -4 12 0 4 10 6 8 -2 95% C.L. excluded area $\log_{10} | d_e^{(1)} | [-]$ SYRTE Rb/Cs Stanford Dy Eöt-Wash Be/Ti MICROSCOPE Ti/Pt MICROSCOPE - osc. -9 $\log_{10} |a_{m_e}^{(1)} - a_g^{(1)}| [-] \log_{10} |a_m^{(1)} - a_g^{(1)}| [-]$ - ' 95% C.L. excluded area -5 -7 -9 3 95% C.L. excluded area -3 -5 -20 -18 -16 -14 -12 -10 -8 -24 -22 -6 -2 -4 $\log_{10} m_{\varphi} \,[{\rm eV}/c^2]$

see A. Hees et al, PRL, 2016 A. Hees et al, PRD, 2018



Systèmes de Référence Temps-Espace

The phenomenology for the quadratic coupling is richer $\frac{d_e}{4e^2}F_{\mu\nu}F^{\mu\nu} - \frac{d_g\beta_3}{2g_3}F^A_{\mu\nu}F^{\mu\nu}_A - \sum_{i=e,u,d}d_{m_i} + \gamma_{m_i}d_gm_i\bar{\psi}_i\bar{\psi}_i$

$$\mathcal{L}_{\mathrm{mat}}[g_{\mu\nu},\Psi] = \mathcal{L}_{\mathrm{SM}}[g_{\mu\nu},\Psi] + \varphi^2$$

$$\varphi^{(2)}(t, \mathbf{x}) = \varphi_0 \cos\left(\frac{m_{\varphi}c^2}{\hbar}t + \delta\right) \left[1 - s_A^{(2)} \frac{GM_A}{c^2 r}\right]$$

Rich phenomenology: scalarization / screening





Systèmes de Référence Temps-Espace

Example II: DM as a scalar topological default



 Topological defect: carry energy and can be a DM candidate

 Frequencies of the clocks will change when topological defaults cross the Earth



From Derevianko and Pospelov, Nature Phys., 2014

Example II: DM as a scalar topological default

- Fiber network: high accuracy longdistance clocks comparison
- Different clocks: Hg/Sr/Yb





- Simulations show that this dataset is very promising
- Data analysis on-going



Work by Benjamin Roberts @SYRTE

Temporal variations of fundamental constants seem unaffected by "screening mechanism"

- Screening mechanism: deviations from GR "hidden" in the Solar System $~\gamma_{\rm PPN}\approx 1$

PRL 107, 251102 (2011)

PHYSICAL REVIEW LETTERS

week ending 16 DECEMBER 2011

Constraints on Shift-Symmetric Scalar-Tensor Theories with a Vainshtein Mechanism from Bounds on the Time Variation of *G*

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• "This mechanism, if efficient to hide the effects of the scalar field at short distance and in the static approximation can in general not alter the cosmological time evolution of the scalar field"

Are there other signatures to be searched in lab data that can help constraining DE models?



Another example: the chameleon

$$S = \int d^4x \sqrt{-g} \left[\frac{R}{2\kappa} - \frac{1}{2} \left(\partial \phi \right)^2 - V(\phi) \right] + S_{\text{mat}} \left[\Psi, e^{2\phi/M} g_{\mu\nu} \right]$$

with $V(\phi) = \Lambda^4 + \frac{\Lambda^{4+n}}{M_{pl}^n \phi^n}$

 Screening mechanism: explain cosmic expansion while hiding the scalar field in region of "high density" Khoury and Weltmann, PRL and PRD, 2004

Atom-interferometry constraints on dark energy

P. Hamilton,^{1*} M. Jaffe,¹ P. Haslinger,¹ Q. Simmons,¹ H. Müller,^{1,2}[†] J. Khoury³

Science, 2013

Jaffe, M., et al, Nature Physics, 2017

Some of the best constraints on chameleons are from laboratory experiments



Another example: a scalar model of DE

from Martins et al, JCAP 08, 47, 2015

$$\mathcal{L} = \frac{R}{16\pi G} - \frac{1}{2} (\partial \phi)^2 - V(\phi) - \frac{1 + d_e \phi}{4} F_{\mu\nu} F^{\mu\nu} + \dots$$

Constraints on DE in a spècific model

from Martins et al, JCAP 08, 47, 2015

- Considered to model Dark Energy. The DE equation of state is characterized by $w(z) = \frac{\phi^2}{p} = \frac{\phi^2}{2} + \frac{1}{2} + \frac{1$
- Combining cosmological data with clocks (1-2-3 σ conf. level)

