Colloque national dark Energy – 2ème édition

The dark side of gravity and the acceleration of the Universe

October 23-25 2018, Paris IAP

F Henry-Couannier CPPM & Aix-Marseille Univ

fppt.com

Dark Gravity theories are extensions of General Relativity aiming at a stable anti-gravitational sector

J.P. Petit, Twin Universe cosmology, Astrophys. Space Sci. Vol. 226, pp 273, 1995. and many other articles

- F. Henry-Couannier, Discrete symmetries and General Relativity, the Dark Side of Gravity, Int.J.Mod.Phys, vol. A20, no. NN, pp. 2341-2346, 2004.
- F. Henry-Couannier, Dark Gravity, GJSFR A. Vol 13, Issue 3, pp 1-53, 2013.

 S. Hossenfelder, Bimetric theory with exchange symmetry Phys. Rev. D 78,
- M. Milgrom, Matter and twin matter in bimetric MOND, MNRAS 405 (2), pp 1129-1139, 2010.

044015, 2008.

Laura Bernard, Luc Blanchet, Lavinia Heisenberg Bimetric gravity and dark matter 50th Rencontres de Moriond, "Gravitation: 100 years after GR", 2015

From background dependence to Dark Gravity (DG) How far can we go?

GR :
$$g_{\mu
u}$$

DG : $g_{\mu
u}$ and $\eta_{\mu
u}$

$$Riemm(\eta_{\mu\nu}) = 0$$

$$\Rightarrow g_{\mu
u}$$
 has a twin, « the inverse metric » $\tilde{g}_{\mu
u}$

$$\tilde{g}_{\mu\nu} = \eta_{\mu\rho}\eta_{\nu\sigma} \left[g^{-1}\right]^{\rho\sigma}$$

$$\Rightarrow (g_{\mu\nu}, \tilde{g}_{\mu\nu})$$
 is a Janus field



From the Action to DG field equations

The Action must respect the permutation symmetry between $g_{\mu\nu}$ and $\tilde{g}_{\mu\nu}$:

$$\int d^4x(\sqrt{g}R + \sqrt{\tilde{g}}\tilde{R}) + \int d^4x(\sqrt{g}L + \sqrt{\tilde{g}}\tilde{L})$$

$$\delta g_{\mu\nu} \Rightarrow \delta S = 0$$

$$\sqrt{g}\eta^{\mu\sigma}g_{\sigma\rho}G^{\rho\nu} - \sqrt{\tilde{g}}\eta^{\nu\sigma}\tilde{g}_{\sigma\rho}\tilde{G}^{\rho\mu} + \mu \leftrightarrow \nu = -8\pi G(\sqrt{g}\eta^{\mu\sigma}g_{\sigma\rho}T^{\rho\nu} - \sqrt{\tilde{g}}\eta^{\nu\sigma}\tilde{g}_{\sigma\rho}\tilde{T}^{\rho\mu} + \mu \leftrightarrow \nu)$$

Contracted form

$$\sqrt{g}R - \sqrt{\tilde{g}}\tilde{R} = 8\pi G(\sqrt{g}T - \sqrt{\tilde{g}}\tilde{T})$$

Implications of DG equations

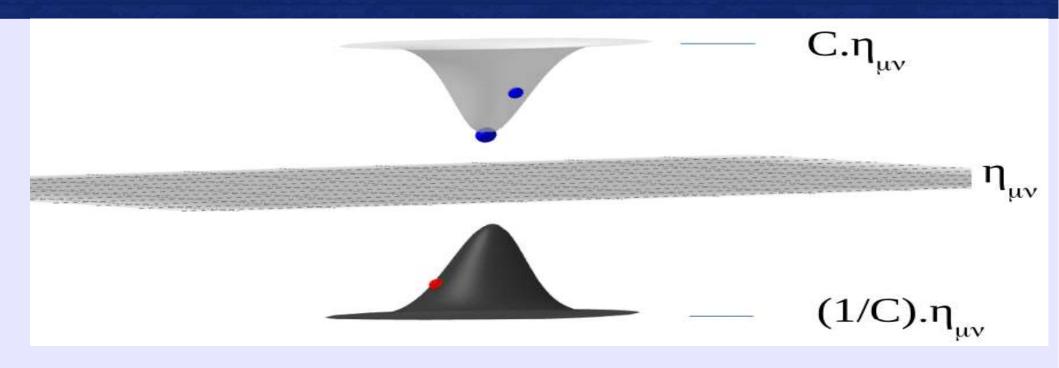
- DG is background dependent yet deviations from GR can remain arbitrarily small provided one side of the Janus Field dominates the other
- Ghost interaction between Janus and source fields but Janus field not understood to be a quantum field!
 - DG more natural than GR as a semiclassical* theory of gravity
 - Semiclassical DG stability : OK**
- New discrete (permutation) symmetry is very fundamental: will be interpreted as a global time reversal symmetry.

^{*} https://arxiv.org/abs/0802.1978 Mark Albers, Claus Kiefer, Marcel Reginatto, Measurement Analysis and Quantum Gravity : « Despite the many physical arguments which speak in favor of a quantum theory of gravity, it appears that the justification for such a theory must be based on empirical tests and does not follow from logical arguments alone »

^{**} https://arxiv.org/pdf/1401.4024.pdf V. A. Rubakov, page 8 : Gradient, tachyonic and ghost instabilities in scalar-tensor theories : « for ghosts, background is QM unstable but classically stable »

The static isotropic solution

Animggb



- Antigravity without run away!
- Asymptotic C matters : GR corresponds to C infinite

The static isotropic solution

$C=\infty$

DG:

$$g_{ii}(r) = A = e^{2MG/r} \approx 1 + 2\frac{MG}{r} + 2\frac{M^2G^2}{r^2}$$

$$-g_{00}(r) = \frac{1}{A} = e^{-2MG/r} \approx 1 - 2\frac{MG}{r} + 2\frac{M^2G^2}{r^2} - \frac{4}{3}\frac{M^3G^3}{r^3}$$

$$O(r) = \frac{\left(1 - \frac{MG}{2r}\right)^2}{\left(1 + \frac{MG}{2r}\right)^2} \approx 1 - 2\frac{MG}{r} + 2\frac{M^2G^2}{r^2} - \frac{3}{2}\frac{M^3G^3}{r^3}$$
No Horizon

- No Horizon
- Zero Gravitational Waves

$$\tilde{h}_{\mu\nu} = -h_{\mu\nu} + O(h^2)$$

$$2(R_{\mu\nu}^{(1)} - \frac{1}{2}\eta_{\mu\nu}R_{\lambda}^{(1)\lambda}) = -8\pi G(T_{\mu\nu} - \tilde{T}_{\mu\nu} + t_{\mu\nu} - \tilde{t}_{\mu\nu})$$
 Deviations from CP at DDN order only

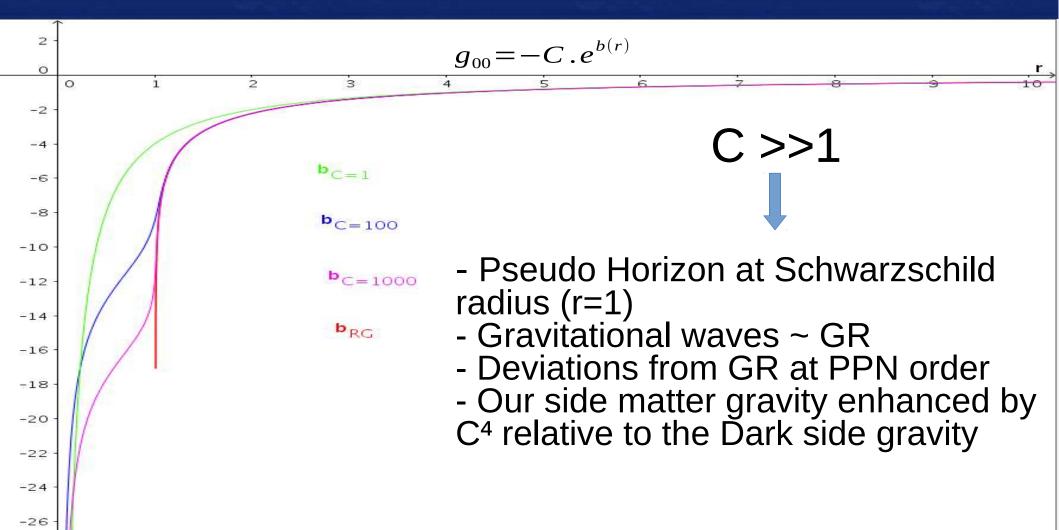
Deviations from GR at PPN order only

RG (Schwarzschild):

$$g_{ii}(r) = \left(1 + \frac{MG}{2r}\right)^4 \approx 1 + 2\frac{MG}{r} + \frac{3}{2}\frac{M^2G^2}{r^2}$$

$$u_{00}(r) = \frac{\left(1 - \frac{MG}{2r}\right)}{\left(1 + \frac{MG}{2r}\right)^2} \approx 1 - 2\frac{MG}{r} + 2\frac{M^2G^2}{r^2} - \frac{3}{2}\frac{M^3G^2}{r^3}$$

The static isotropic solution



Cosmological equation

- Homogeous & isotropic Janus solution is flat and static : C was indeed a constant!
 - \Rightarrow We need to introduce a separate scalar- η Janus field for cosmology :

$$g_{\mu\nu} = \Phi \eta_{\mu\nu} \text{ and } \tilde{g}_{\mu\nu} = \frac{1}{\Phi} \eta_{\mu\nu} \qquad \Phi(t) = a^2(t)$$

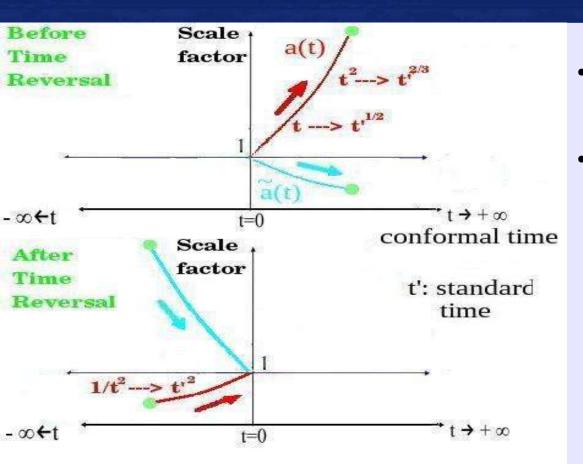
• Single scale factor equation :

$$a\ddot{a} - \tilde{a}\ddot{\tilde{a}} = \frac{4\pi G}{3}(a^4(\rho - 3p) - \tilde{a}^4(\tilde{\rho} - 3\tilde{p}))$$

$$\tilde{a}(t) = \frac{1}{a(t)}$$

Cosmological solutions

Anim ggb



- Janus scale factors are related by a global conformal time reversal symmetry T : $\tilde{a}(t) = \frac{1}{a(t)} = a(-t)$
- Both continuous evolution and discontinuous permutation T allowed when $\rho 3p = \tilde{\rho} 3\tilde{p}$

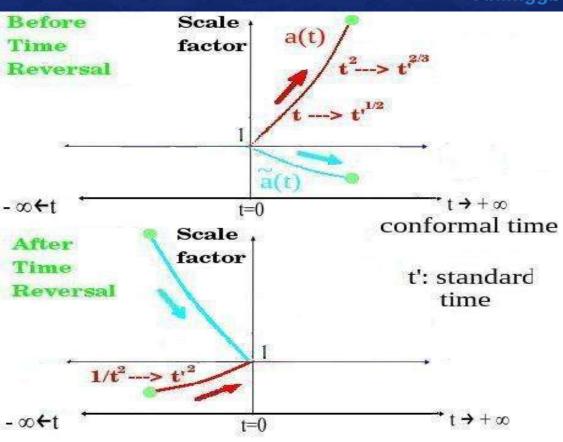


Global time reversal: not going backward in time, but jumping to the opposite time!

A cyclic Universe?

DG Cosmology

Animggb



Hyp : $\rho \simeq \rho - 3p = \tilde{\rho} - 3\tilde{p} \simeq \tilde{\rho}$ occured at transition redshift triggering T and a'(t')~ t'^2

With H(t) continuous at the transition and assuming same universe age as in LCDM:

$$a'(t') \sim t'^{\alpha} \Longrightarrow z_{tr} = \left(\frac{2/3 - \alpha}{1 - \alpha}\right)^{\alpha} - 1$$

- \Rightarrow z_{tr} = 0.78 vs observed z_{tr} =0.67+- 0.1
- ~ Same scale factor evolution as in LCDM
- Without DE
- Inflation not needed to get k=0
- Without Big Bang singularity
- Cosmological DM still needed
- Dark side effects only since t_{tr} or near t=0

Problem statement

- We have two separate theories :
 - Asymptotically static DG correctly describes all aspects of gravity except expansion
 - Scalar-η Janus field only correctly describes expansion
- How to get expansion effects on the largest scales and differential eoms non trivially mixing background and perturbations (GR like) as needed to reproduce CMB phenomenology?

Conclusion and outlooks

- DG avoids Big-Bang singularity and BH horizon very naturally
- Acceleration, k=0, large scale homogeneity, matter/antimatter asym
- Likely to cancel the gravity of vacuum energy
- Outlook:

Unification ⇒ New rich and effective phenomenology

(DM candidate, ...)

www.darksideofgravity.com/DG.pdf

How far could we go?

⇒ Fascinating phenomenological and theoretical implications!

^{*} EP violations (η effects) usually small, **harmless classical instabilities

Dynamical discrete symmetries

• Standard view:

Symmetries (cont & disc) ⇒ Action
Extreme action principle ⇒ Eoms & conservation equations
No dynamical processes associate with discrete symmetries

Extended view :

Symmetries (cont & disc) ⇒ Action

Extreme action principle ⇒ Eoms & conservation equations

Discrete symmetries ⇒ Discontinuous processes

Dynamical discrete symmetries

- 1) Discrete (permutation) symmetry and continuous symmetries already unified in DG framework
- 2) Just as discrete (T&P) and continuous spacetime symmetries already unified in the Lorentz group
- 1) and 2) turn out to be related : global T symmetry is permutation symmetry !

Dynamical discrete symmetries \Rightarrow discontinuous transitions in addition to usual continuous evolution processes deduced from differential eoms.

- ⇒ Fills the gap between the discrete and the continuous
- ⇒ Hopefully opens the way to a genuine unification (understanding) of QM discrete and non local laws to the rest of physics!

Vacuum energy terms in DG equations

DG vacuum source term:

$$(\sqrt{g}\Lambda - \sqrt{\tilde{g}}\tilde{\Lambda})g^{\mu\nu}$$

Cancels for $g_{\mu\nu} = \tilde{g}_{\mu\nu} = \eta_{\mu\nu}$ and $\Lambda = \tilde{\Lambda}$ (natural)

⇒ Might remain zero when Janus field starts to evolve, may be through the auto-adjustment of cut-offs to preserve compensation.

DG unification with adiabatic particles exchange? * adapted from original idea by Prigogin et al

 Matter and radiation fields conservation equations including adiabatic gravitationnally induced* transfers occuring between the two metrics:

$$\nabla_{\mathbf{v}} T_{\mathbf{\mu}}^{\mathbf{v}} \neq \mathbf{0} = \frac{\dot{\rho}}{H} = (\frac{\Gamma}{H} - 3)(\rho + p)$$

$$\widetilde{\nabla}_{\mathbf{v}} \widetilde{T}_{\mathbf{\mu}}^{\mathbf{v}} \neq \mathbf{0} = \frac{\dot{\tilde{\rho}}}{\tilde{H}} = (\frac{\tilde{\Gamma}}{\tilde{H}} - 3)(\tilde{\rho} + \tilde{p})$$

$$\widetilde{\nabla}_{\mathbf{v}} \widetilde{T}_{\mathbf{\mu}}^{\mathbf{v}} \neq \mathbf{0} = \frac{\dot{\tilde{\rho}}}{\tilde{H}} = (\frac{\tilde{\Gamma}}{\tilde{H}} - 3)(\tilde{\rho} + \tilde{p})$$

Replacing in DG_Friedmann equations

⇒ ~ usual solutions valid provided

$$a\ddot{a} = K(a^{4}(\rho - 3p) + \frac{1}{2}(C + \tilde{C})) \qquad C = a^{4}\frac{\Gamma}{H}(\rho + p)$$

$$\tilde{a}\ddot{\tilde{a}} = K(\tilde{a}^{4}(\tilde{\rho} - 3\tilde{p}) + \frac{1}{2}(C + \tilde{C})) \qquad \tilde{C} = \tilde{a}^{4}\frac{\Gamma}{H}(\tilde{\rho} + \tilde{p})$$

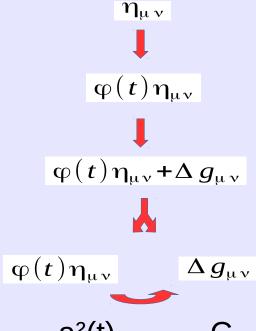
 $\Gamma \approx 2H \frac{a\ddot{a}}{a\ddot{a}}$ (for $a \ll \tilde{a}$)

$$\tilde{C}(\tilde{a}^4(\tilde{\rho}-3\tilde{p})+\frac{1}{2}(C+\tilde{C}))$$
 $\tilde{C}=\tilde{a}^4\frac{\Gamma}{H}(\tilde{\rho}+\tilde{p})$

DG unification with Emerging Dynamics (ED)

As the universe evolves new dynamical dofs are released:

- Non dynamical
- Homogenous scalar-eta
- Scalar-eta + non dynamical fluctuation
- Separate dynamics



ED: Early DG unification

• For a²(t) < Fundamental Threshold,

$$g_{\mu\nu} = \varphi(t) \eta_{\mu\nu} + \Delta g_{\mu\nu}$$

but only the scalar $\varphi(t)$ is dynamical \Rightarrow we again get a single equation

• Symmetries related to our privileged coordinate system (rather than isometries related to the sources) force the primordial metrics in the Newtonian Gauge form : $d\tau^2 = a^2(t)((1+2\Psi)dt^2 - (1-2\Psi)d\sigma^2)$

 ⇒ We get the same scale factor (order 0) and potential (order 1) eoms as in GR but rotational and radiative modes should be absent from the CMB.

ED: Late DG unification

- $a^2(t)$ > Fundamental Threshold breaks the primordial symmetries $\Rightarrow \varphi(t)\eta_{\mu\nu}$ and $\Delta g_{\mu\nu}$ start to play their dynamics independently
 - ⇒ Late DG unification required to account for expansion effects
- In the Linear domain, C (integration constant of $\Delta g_{\mu\nu}$) is driven step by step by the scale factor from $\varphi(t)\eta_{\mu\nu}$:
 - ⇒ expansion effects through discrete rules
 - ⇒ rich new and effective phenomenology related to field discontinuities
- In the Non Linear domain (solar system), we are asymptotically Minkowskian: C strictly constant!

Classical stability issues

- Background remains bounded thanks to global time reversal
- Linear inhomogeneous perturbations unstable in contracting phase but gravity from these is negligible: suppressed by C⁴ factor (~scale_factor⁸) before transition to acceleration.
- Linear inhomogeneous perturbations from the dark sector can start to grow under their own gravity after transition
- Strong gravity inhomogeneous pertubations presumably always stable on both sides thanks to C >1 at our side structures while C<1 at dark side structures

Problems with semiclassical Gravity

 Case I: Classical gravity triggers quantum collapses ⇒ no Energymomentum conservation violation, nor violation of uncertainty relations contrary to popular argument by Eppley & Hannah ...

https://arxiv.org/pdf/0802.1978.pdf

otherwise:

- Case 2A: No collapse interpretation of QM (MWI, decoherence ...) ruled out because classical gravity would see the uncollapsed superpositions
- Case 2B: Realistic collapse interpretation of QM leads to possible faster than light signaling. Either specific more local model of quantum collapse can solve this or ... DG: instantaneous signaling is not anymore a menace to causality as soon as there exists a unic privileged instantaneity frame for any collapse!