

Massive black hole triplets as LISA sources

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With:

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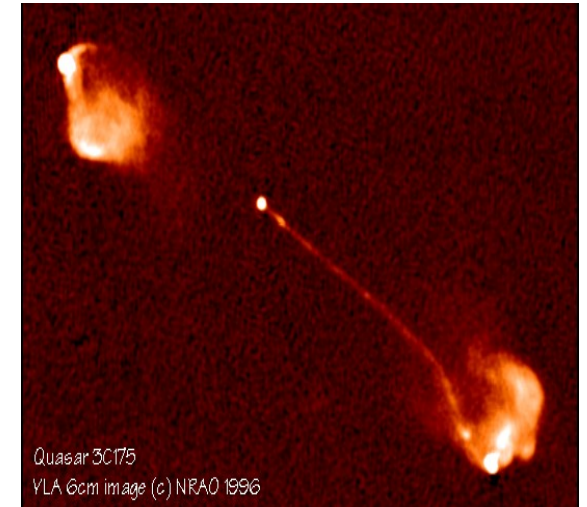
Massive black hole binaries

Galaxies merge

Galaxies host MBHs

Formation of Massive
Black Hole Binaries
(MBHBs)

Source of GWs



Standard scenario

Two MBHs in order to get close enough to emit detectable GWs have to cover a **huge spatial range**.
They must rely on several astrophysical processes

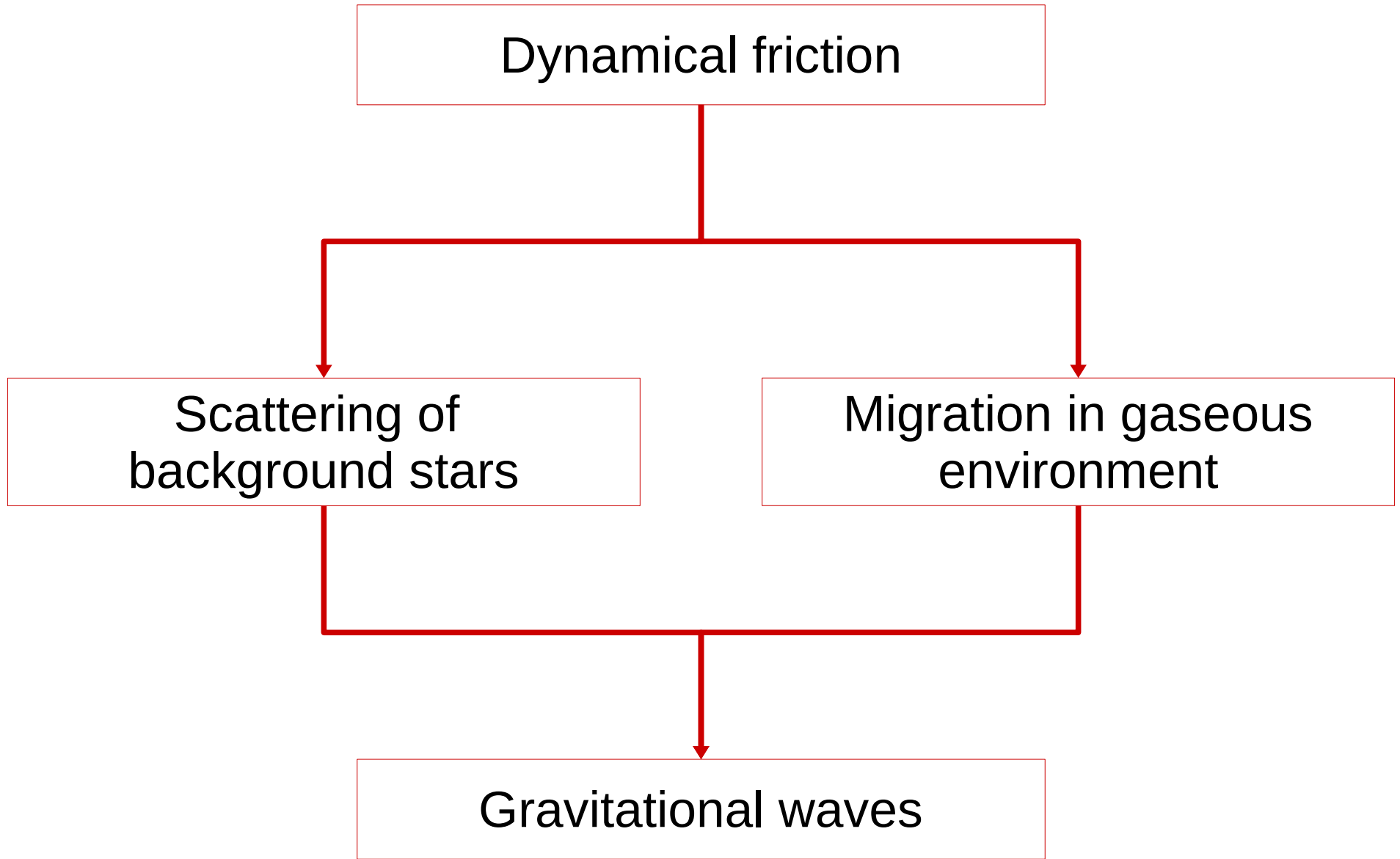
Standard scenario

Dynamical friction

Scattering of
background stars

Migration in gaseous
environment

Gravitational waves



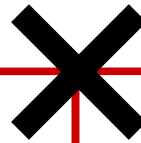
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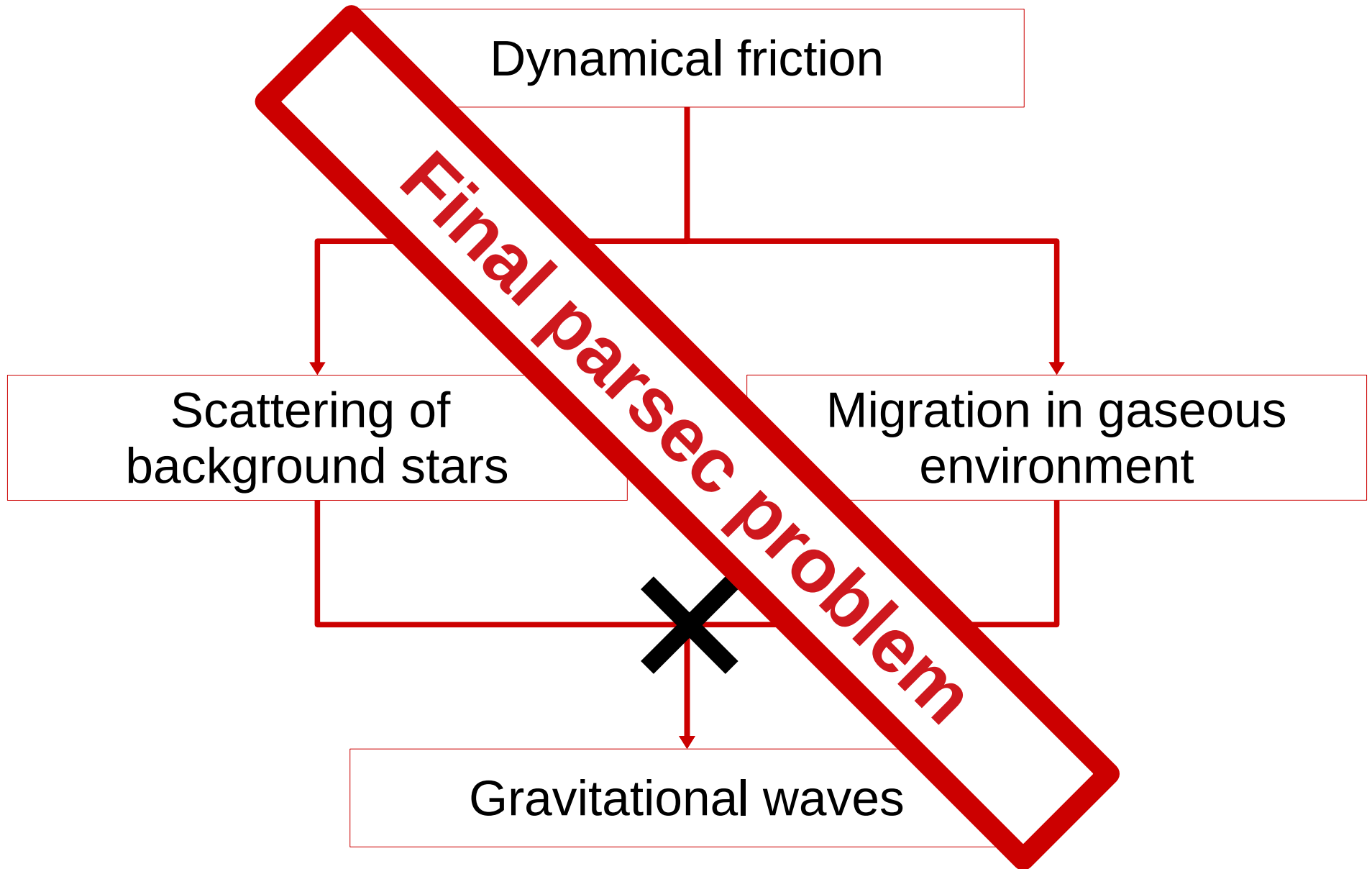
Scattering of
background stars

Migration in gaseous
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Standard scenario



What if...

Hierarchical growth of
cosmic structure

Binary stalling

- Successive galaxy mergers
- Galaxies host at least one MBH

Formation of triplets of MBHs

Our Goal

Assess the implications of a sizeable population of MBH triplets

↓ How?

Simulating MBH triplets in galactic nuclei with astrophysically and cosmologically motivated initial conditions

Part 1:

Integrate equations of motion of a wide set of triplets with different initial conditions.

Always a stalled binary + a third body

Part 2:

Embed the results in a cosmological framework

Part 1: setup

Simulate MBH triplets in galactic nuclei with astrophysically and cosmologically motivated initial conditions

3-body Newtonian dynamics + GR corrections up to 2.5PN

Part 1: Hamiltonians

Newtonian

$$H_0 = \frac{1}{2} \sum_{\alpha} \frac{|\vec{p}_{\alpha}|^2}{m_{\alpha}} - \frac{G}{2} \sum_{\alpha} \sum_{\beta \neq \alpha} \frac{m_{\alpha} m_{\beta}}{r_{\alpha\beta}}$$

1PN

$$H_1 = -\frac{1}{8} \sum_{\alpha} m_{\alpha} \left(\frac{|\vec{p}_{\alpha}|^2}{m_{\alpha}^2} \right)^2 - \frac{G}{4} \sum_{\alpha} \sum_{\beta \neq \alpha} \frac{1}{r_{\alpha\beta}} \left[6 \frac{m_{\beta}}{m_{\alpha}} |\vec{p}_{\alpha}|^2 - 7 \vec{p}_{\alpha} \cdot \vec{p}_{\beta} - (\vec{n}_{\alpha\beta} \cdot \vec{p}_{\alpha})(\vec{n}_{\alpha\beta} \cdot \vec{p}_{\beta}) \right] + \frac{G^2}{2} \sum_{\alpha} \sum_{\beta \neq \alpha} \sum_{\gamma \neq \alpha, \beta} \frac{m_{\alpha} m_{\beta} m_{\gamma}}{r_{\alpha\beta} r_{\alpha\gamma}}$$

2.5PN

$$H_{2.5} = \frac{G}{45} \dot{\chi}_{(4)ij}(\vec{x}_{\alpha'}, \vec{p}_{\alpha'}; t) \chi_{(4)ij}(\vec{x}_{\alpha}, \vec{p}_{\alpha})$$

$$H_2 = \frac{1}{16} \sum_{\alpha} m_{\alpha} \left(\frac{|\vec{p}_{\alpha}|^2}{m_{\alpha}^2} \right)^3 + \frac{G}{16} \sum_{\alpha} \sum_{\beta \neq \alpha} \frac{(m_{\alpha} m_{\beta})^{-1}}{r_{\alpha\beta}} \left[10 \left(\frac{m_{\beta}}{m_{\alpha}} |\vec{p}_{\alpha}|^2 \right)^2 - 11 |\vec{p}_{\alpha}|^2 |\vec{p}_{\beta}|^2 - 2 (\vec{p}_{\alpha} \cdot \vec{p}_{\beta})^2 + 10 |\vec{p}_{\alpha}|^2 (\vec{n}_{\alpha\beta} \cdot \vec{p}_{\beta})^2 - 12 (\vec{p}_{\alpha} \cdot \vec{p}_{\beta})(\vec{n}_{\alpha\beta} \cdot \vec{p}_{\alpha})(\vec{n}_{\alpha\beta} \cdot \vec{p}_{\beta}) - 3 (\vec{n}_{\alpha\beta} \cdot \vec{p}_{\alpha})^2 (\vec{n}_{\alpha\beta} \cdot \vec{p}_{\beta})^2 \right] + \frac{G^2}{8} \sum_{\alpha} \sum_{\beta \neq \alpha} \sum_{\gamma \neq \alpha, \beta} \frac{1}{r_{\alpha\beta} r_{\alpha\gamma}} \left[18 \frac{m_{\beta} m_{\gamma}}{m_{\alpha}} |\vec{p}_{\alpha}|^2 + 14 \frac{m_{\alpha} m_{\gamma}}{m_{\beta}} |\vec{p}_{\beta}|^2 - 2 \frac{m_{\alpha} m_{\gamma}}{m_{\beta}} (\vec{n}_{\alpha\beta} \cdot \vec{p}_{\beta})^2 - 50 m_{\gamma} (\vec{p}_{\alpha} \cdot \vec{p}_{\beta}) + 17 m_{\alpha} (\vec{p}_{\beta} \cdot \vec{p}_{\gamma}) - 14 m_{\gamma} (\vec{n}_{\alpha\beta} \cdot \vec{p}_{\alpha})(\vec{n}_{\alpha\beta} \cdot \vec{p}_{\beta}) + 14 m_{\alpha} (\vec{n}_{\alpha\beta} \cdot \vec{p}_{\beta})(\vec{n}_{\alpha\beta} \cdot \vec{p}_{\gamma}) + m_{\alpha} (\vec{n}_{\alpha\beta} \cdot \vec{n}_{\alpha\gamma})(\vec{n}_{\alpha\beta} \cdot \vec{p}_{\beta})(\vec{n}_{\alpha\gamma} \cdot \vec{p}_{\gamma}) \right] + \frac{G^2}{8} \sum_{\alpha} \sum_{\beta \neq \alpha} \sum_{\gamma \neq \alpha, \beta} \frac{1}{r_{\alpha\beta}^2} \left[2 m_{\beta} (\vec{n}_{\alpha\beta} \cdot \vec{p}_{\alpha})(\vec{n}_{\alpha\gamma} \cdot \vec{p}_{\gamma}) + 2 m_{\beta} (\vec{n}_{\alpha\beta} \cdot \vec{p}_{\beta})(\vec{n}_{\alpha\gamma} \cdot \vec{p}_{\gamma}) + \frac{m_{\alpha} m_{\beta}}{m_{\gamma}} (5 (\vec{n}_{\alpha\beta} \cdot \vec{n}_{\alpha\gamma}) |\vec{p}_{\gamma}|^2 - (\vec{n}_{\alpha\beta} \cdot \vec{n}_{\alpha\gamma})(\vec{n}_{\alpha\gamma} \cdot \vec{p}_{\gamma})^2 - 14 (\vec{n}_{\alpha\beta} \cdot \vec{p}_{\gamma})(\vec{n}_{\alpha\gamma} \cdot \vec{p}_{\gamma})) \right] + \frac{G^2}{4} \sum_{\alpha} \sum_{\beta \neq \alpha} \frac{m_{\alpha}}{r_{\alpha\beta}^2} \left[\frac{m_{\beta}}{m_{\alpha}} |\vec{p}_{\alpha}|^2 + \frac{m_{\alpha}}{m_{\beta}} |\vec{p}_{\beta}|^2 - 2 (\vec{p}_{\alpha} \cdot \vec{p}_{\beta}) \right] + \frac{G^2}{2} \sum_{\alpha} \sum_{\beta \neq \alpha} \sum_{\gamma \neq \alpha, \beta} \frac{(n_{\alpha\beta}^i + n_{\alpha\gamma}^i)(n_{\alpha\beta}^j + n_{\alpha\gamma}^j)}{(r_{\alpha\beta} + r_{\beta\gamma} + r_{\gamma\alpha})^2} \left[8 m_{\beta} (p_{\alpha i} p_{\gamma j}) - 16 m_{\beta} (p_{\alpha j} p_{\gamma i}) + 3 m_{\gamma} (p_{\alpha i} p_{\beta j}) + 4 \frac{m_{\alpha} m_{\beta}}{m_{\gamma}} (p_{\gamma i} p_{\gamma j}) + \frac{m_{\beta} m_{\gamma}}{m_{\alpha}} (p_{\alpha i} p_{\alpha j}) \right] + \frac{G^2}{2} \sum_{\alpha} \sum_{\beta \neq \alpha} \sum_{\gamma \neq \alpha, \beta} \frac{m_{\alpha} m_{\beta} m_{\gamma}}{(r_{\alpha\beta} + r_{\beta\gamma} + r_{\gamma\alpha}) r_{\alpha\beta}} \left[8 \frac{\vec{p}_{\alpha} \cdot \vec{p}_{\gamma} - (\vec{n}_{\alpha\beta} \cdot \vec{p}_{\alpha})(\vec{n}_{\alpha\beta} \cdot \vec{p}_{\gamma})}{m_{\alpha} m_{\gamma}} - 3 \frac{\vec{p}_{\alpha} \cdot \vec{p}_{\beta} - (\vec{n}_{\alpha\beta} \cdot \vec{p}_{\alpha})(\vec{n}_{\alpha\beta} \cdot \vec{p}_{\beta})}{m_{\alpha} m_{\beta}} - 4 \frac{|\vec{p}_{\gamma}|^2 - (\vec{n}_{\alpha\beta} \cdot \vec{p}_{\gamma})^2}{m_{\gamma}^2} - \frac{|\vec{p}_{\alpha}|^2 - (\vec{n}_{\alpha\beta} \cdot \vec{p}_{\alpha})^2}{m_{\alpha}^2} \right] - \frac{G^3}{2} \sum_{\alpha} \sum_{\beta \neq \alpha} \left(\sum_{\gamma \neq \alpha, \beta} \frac{m_{\alpha}^2 m_{\beta} m_{\gamma}}{r_{\alpha\beta}^2 r_{\beta\gamma}} + \frac{1}{2} \sum_{\gamma \neq \beta} \frac{m_{\alpha}^2 m_{\beta} m_{\gamma}}{r_{\alpha\beta}^2 r_{\beta\gamma}} \right) - \frac{3G^3}{8} \sum_{\alpha} \sum_{\beta \neq \alpha} \left(\sum_{\gamma \neq \alpha, \beta} \frac{m_{\alpha}^2 m_{\beta} m_{\gamma}}{r_{\alpha\beta}^2 r_{\alpha\gamma}} + \sum_{\gamma \neq \alpha, \beta} \frac{m_{\alpha}^2 m_{\beta} m_{\gamma}}{r_{\alpha\beta}^2 r_{\beta\gamma}} \right) - \frac{3G^3}{8} \sum_{\alpha} \sum_{\beta \neq \alpha} \sum_{\gamma \neq \alpha, \beta} \frac{m_{\alpha}^2 m_{\beta} m_{\gamma}}{r_{\alpha\beta} r_{\alpha\gamma} r_{\beta\gamma}} - \frac{G^3}{64} \sum_{\alpha} \sum_{\beta \neq \alpha} \sum_{\gamma \neq \alpha, \beta} \frac{m_{\alpha}^2 m_{\beta} m_{\gamma}}{r_{\alpha\beta} r_{\alpha\gamma}^3 r_{\beta\gamma}} \left[18 r_{\alpha\gamma}^2 - 60 r_{\beta\gamma}^2 - 24 r_{\alpha\gamma} (r_{\alpha\beta} + r_{\beta\gamma}) + 60 \frac{r_{\alpha\gamma} r_{\beta\gamma}^2}{r_{\alpha\beta}} + 56 r_{\alpha\beta} r_{\beta\gamma} - 72 \frac{r_{\beta\gamma}^3}{r_{\alpha\beta}} + 35 \frac{r_{\beta\gamma}^4}{r_{\alpha\beta}^2} + 6 r_{\alpha\beta}^2 \right] - \frac{G^3}{4} \sum_{\alpha} \sum_{\beta \neq \alpha} \frac{m_{\alpha}^2 m_{\beta}^2}{r_{\alpha\beta}^3}$$

2PN

Part 1: setup

Simulate MBH triplets in galactic nuclei with astrophysically and cosmologically motivated initial conditions

- 3-body Newtonian dynamics + GR corrections up to 2.5PN
+ interaction with stellar environment, i.e.,
- bulge potential (spherically symmetric)
 - Dynamical Friction + stellar hardening (dissipative forces)

Part 1: exploration through simulations

Systematic survey of the parameter space

$$m_1 = 10^5 - 10^{10} M_\odot$$

$$\log q_{\text{in}} = -1.5, -1.0, -0.5, 0.0$$

$$\log q_{\text{out}} = -1.5, -1.0, -0.5, 0.0$$

$$e_{\text{in}} = 0.2, 0.4, 0.6, 0.8$$

$$e_{\text{out}} = 0.3, 0.6, 0.9$$

$$\cos \iota = 13 \text{ values equally spaced in } (-1, 1)$$

14976
simulations

$\log m_1$ [M_\odot]	% Mergers			Total
	m_1-m_2	m_1-m_3	m_2-m_3	
5	16.8	0.9	0.8	18.5(1.6)
6	16.2	1.4	1.0	18.5(1.9)
7	15.4	2.5	1.4	19.4(4.4)
8	14.7	4.0	2.5	21.2(6.3)
9	15.2	4.1	3.2	22.5(11.2)
10	21.1	7.6	3.3	31.9(12.7)

We found that ~20-30% of otherwise stalled binaries actually undergo a merger. Main driver is the high eccentricity.

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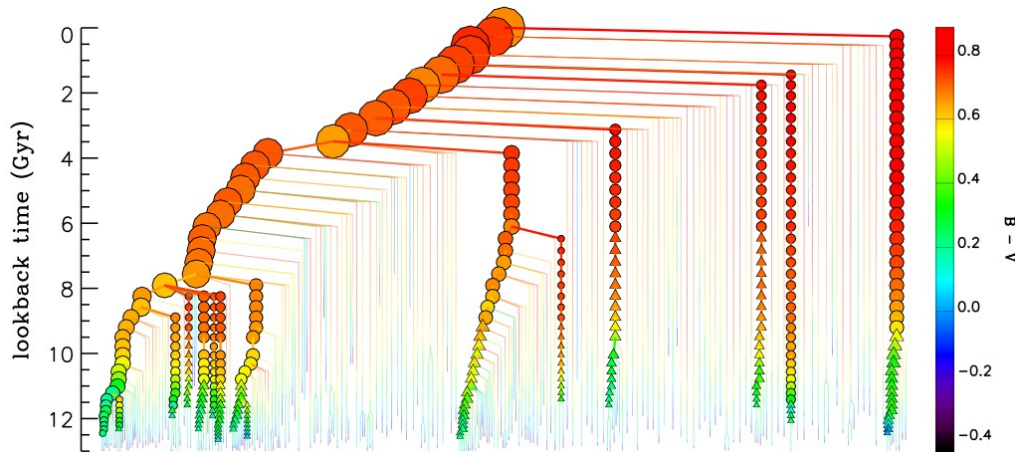
- Secular 3-body dynamics (Kozai-Lidov mechanism)
- Chaotic 3-body dynamics (ejections, exchanges, strong encounters)

Part 2: SAM

Systematic survey of the parameter space



Coupling the results to a cosmological merger tree + SAM



Infer the
“cosmological weight”
of each combination of
surveyed parameters

Part 2: SAM

Systematic survey of the parameter space



Coupling the results to a cosmological merger tree + SAM

SAM of Barausse (2012)
and later expansions



Ad-hoc recipe to include triple
interactions in a cosmological
framework

Part 2: SAM

Systematic survey of the parameter space

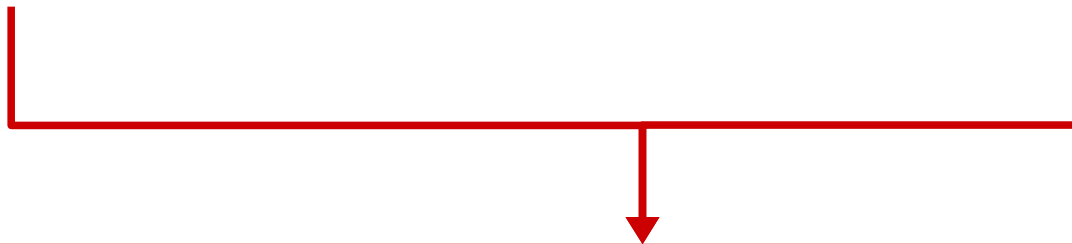


Coupling the results to a cosmological merger tree + SAM

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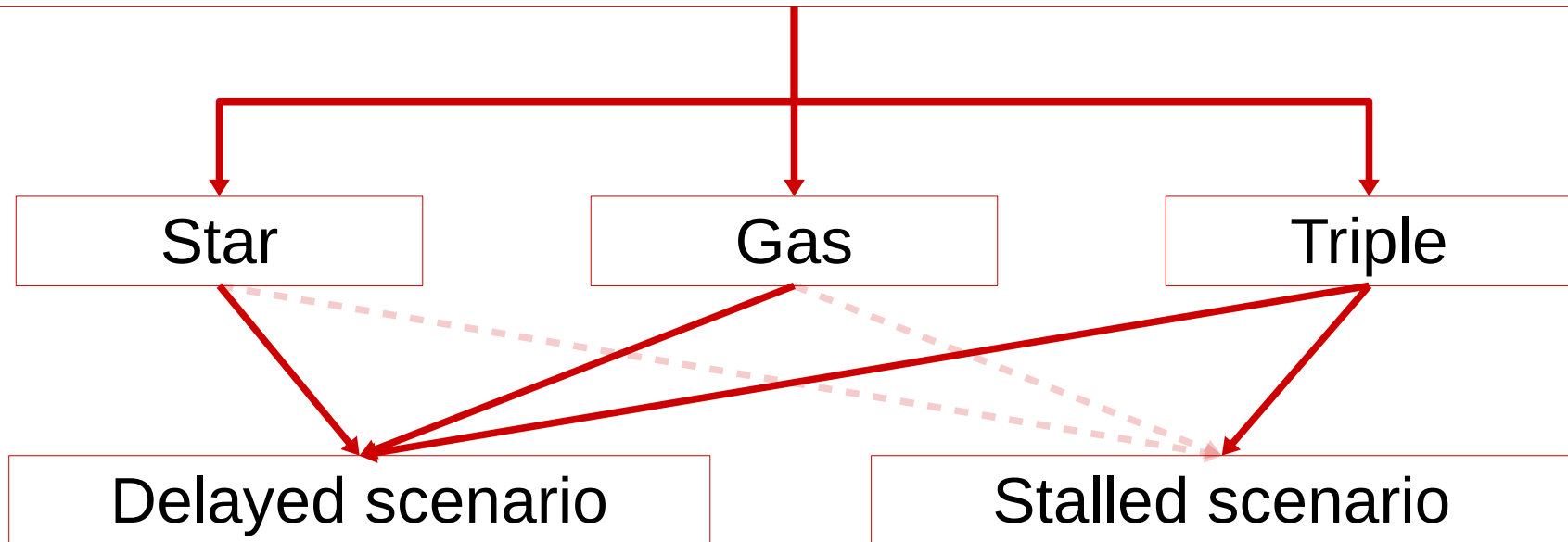
Ad-hoc recipe to include triple
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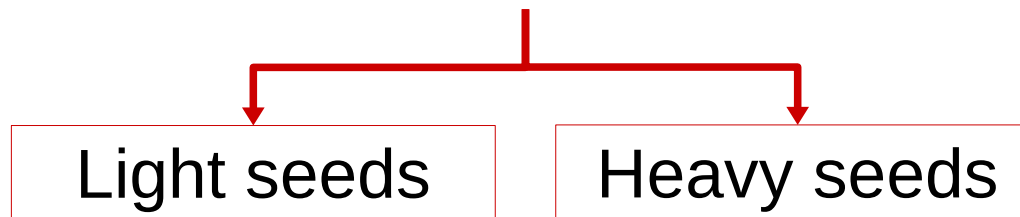
Catalogues of coalescing MBHBs,
keeping track of the mechanism that forced the coalescence

Part 2: SAM

Catalogues of coalescing MBHBs,
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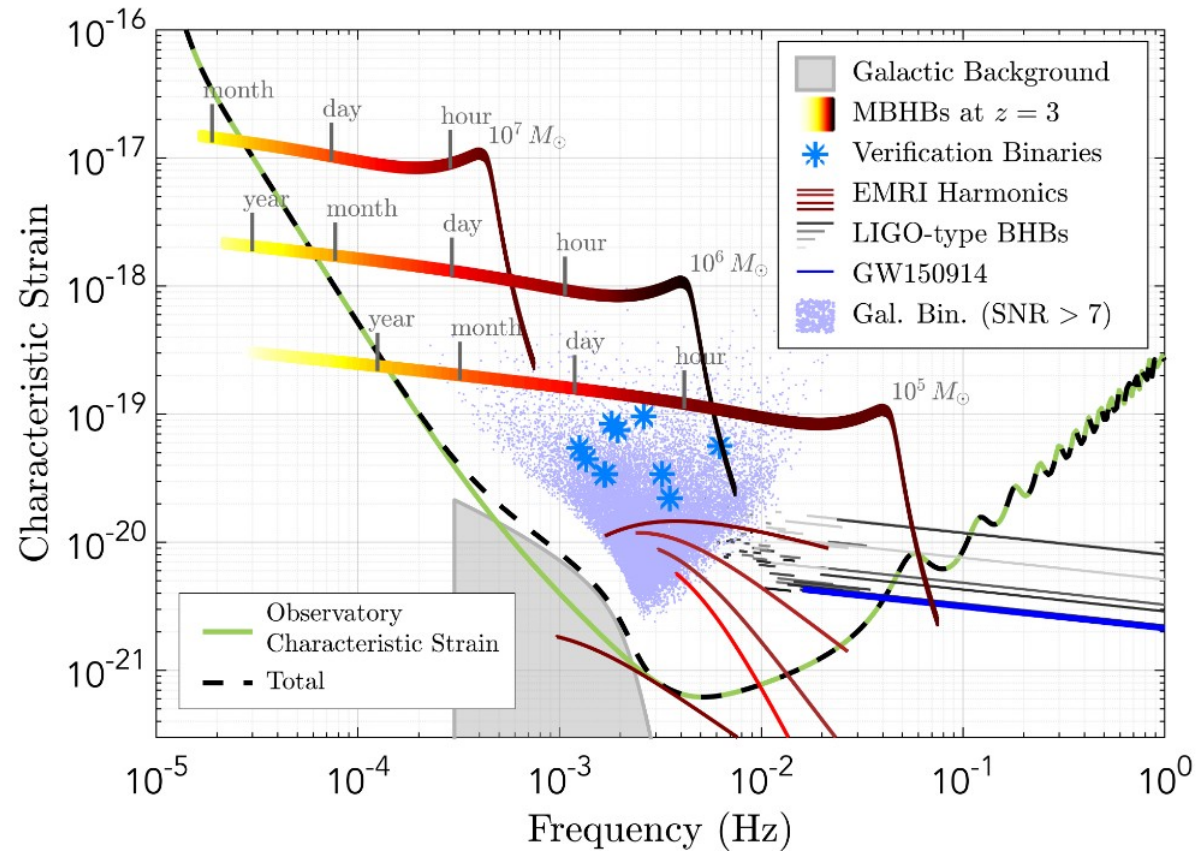
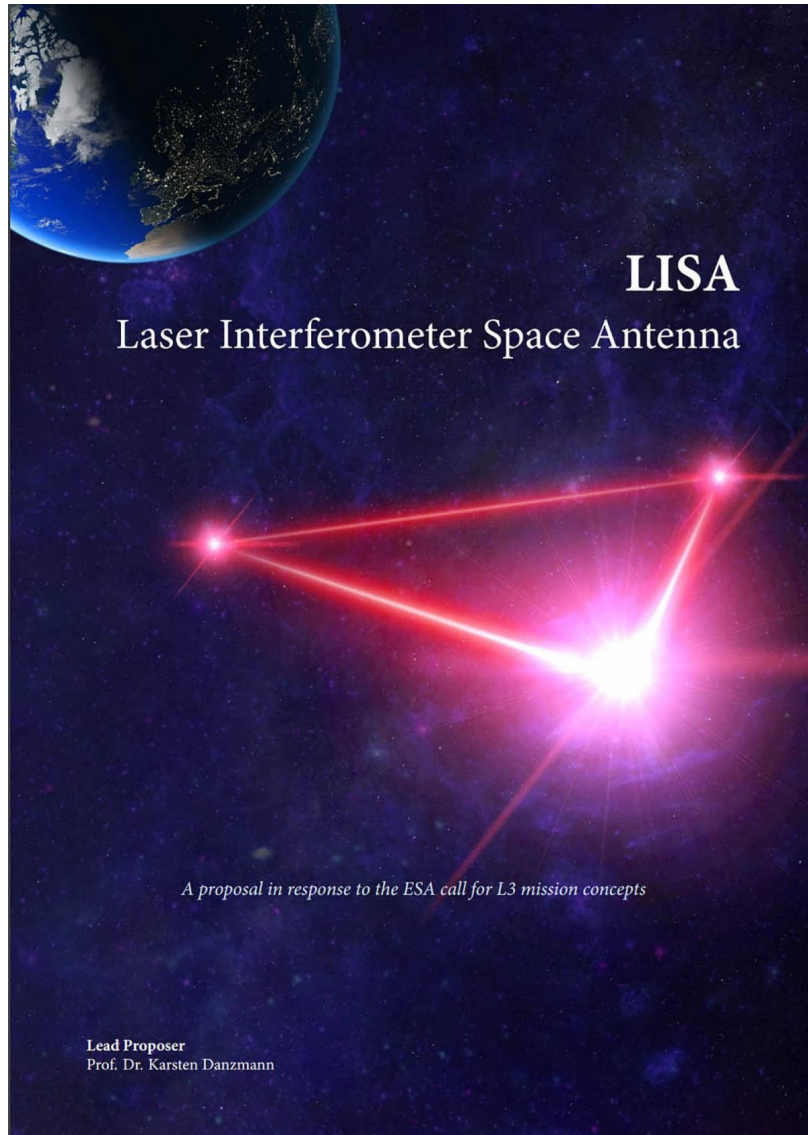


For both scenarios we consider two different MBH seeding recipes



**4 different Universe
realisations**

Implications for LISA



Implications for LISA: mergers

We are mainly interested in single source detection

From the SAM we can infer the merger rate

$$\frac{dN}{dt} \times 4 \text{ yr}$$

Average number of mergers in 4 yr				
	HS-stalled	LS-stalled	HS-delayed	LS-delayed
Triple	~42	~86	~36	~15
Total	~42	~86	~90	~890

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How many events are effectively observable with LISA?

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Implications for LISA: detections

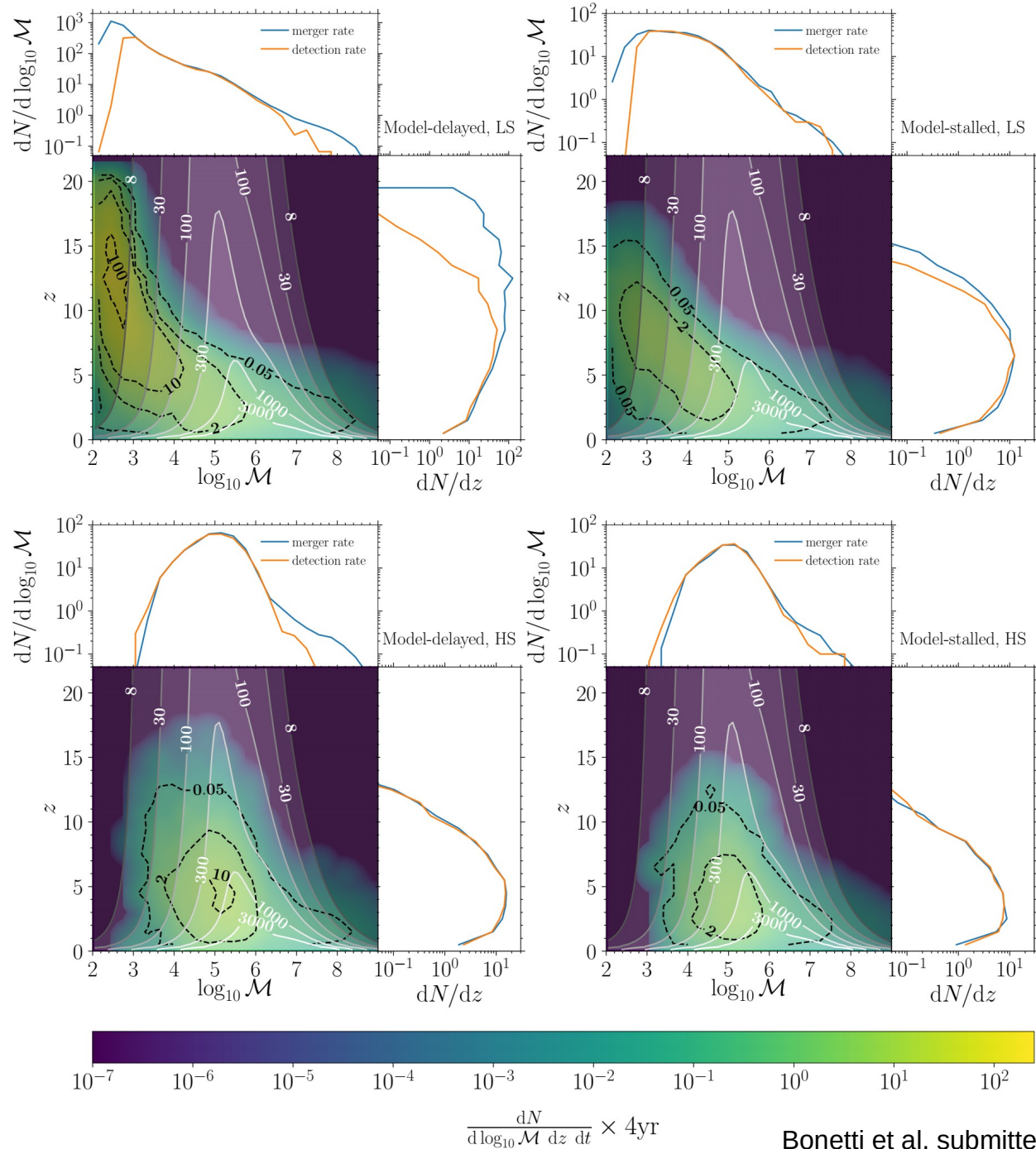
Qualitative answer

Note the difference according to seeding model

HS: all mergers are detected

LS: some mergers will be missed

$$\frac{d^3 N}{dz d \log_{10} \mathcal{M} dt} \times 4 \text{ yr}$$

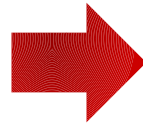


Implications for LISA: detections

Quantitative answer

To infer the detection rate we have to switch from a “time-domain” to a “frequency-domain” description

$$\frac{d^3 N}{dz d\mathcal{M} dt}$$



$$\frac{d^4 N}{dz d\mathcal{M} df de}$$

Implications for LISA: detections

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$$\frac{d^3 N}{dz d\mathcal{M} dt} \quad \rightarrow \quad \frac{d^4 N}{dz d\mathcal{M} df de}$$

$$\frac{d^4 N}{dz d\mathcal{M} df de} = \frac{d^3 N}{dz d\mathcal{M} dt} \times \tau(f, e)$$

Time spent by binaries in each frequency bin

Implications for LISA: detections

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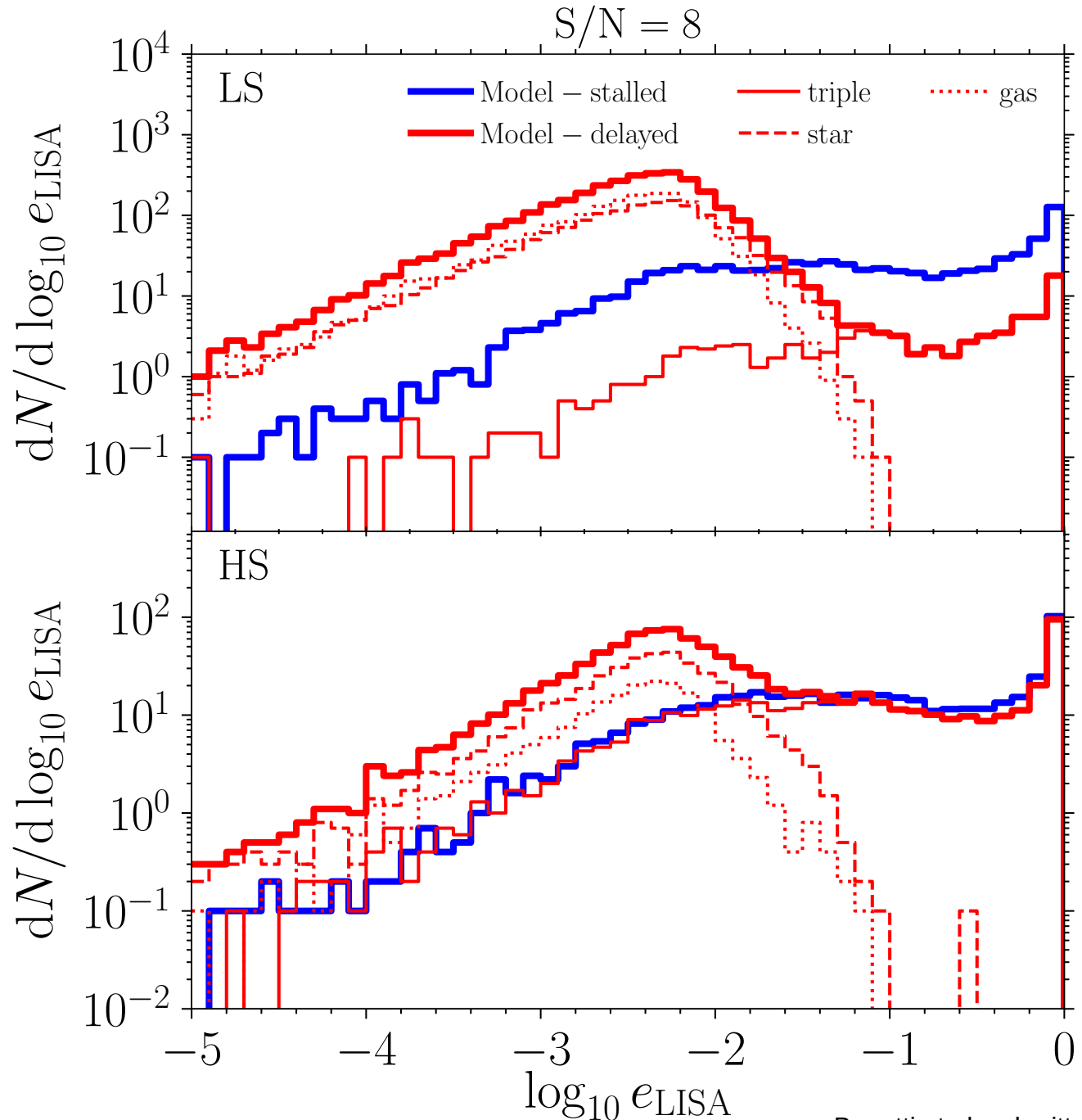
Through a Montecarlo sampling we obtain the total population of binaries emitting GW in the Universe

Time spent by binaries in each frequency bin

Implications for LISA: eccentricity

The eccentricity in the LISA band (i.e. when $S/N = 8$) can be quite high

The dynamical channel is the only evolutionary path (apart from very extreme case) that can leave such imprint



Conclusions & main results

- **Triple interactions can be a viable evolutionary channel**
- **Even in the most pessimistic scenario the merger rate is not heavily suppressed**
- **Triple induced mergers can enter the LISA band with high eccentricity, requiring specific waveform templates**

How to fit in the main LISA science target:

Astrophysically motivated catalogues of GW sources for

- Setting requirements for accurate waveforms
- Informing data analysis people about additional possible sources