Search for new resonances in yy channel with the ATLAS detector

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1. Theoretical context

Extension of Standard Model

Two Higgs-doublet models (2HDM)

TWO HIGGS-DOUBLET MODELS ARE SIMPLE EXTENSION OF THE SM WITH

AN ENRICHED SCALAR SECTOR

 \rightarrow Introduction of an **additional scalar field** => 2 doublet scalar fields Φ_1 and Φ_2 in the SM lagrangian (8 degrees of freedom).

→ After symmetry breaking => Prediction of physical 5 states/Higgs Boson :

Two CP-even bosons h and H



 \rightarrow Taken from Camille's <u>talk</u>

Same motivation, different experiment!

2 neutral CP-even bosons: H, h with $m_H > m_h$

H or h might be the 125 GeV boson discovered 2012... \rightarrow what about the other one?

Theory and phenomenology of two-Higgs-doublet models

2. Experiment

ATLAS Detector





→ EM calorimeter and inner tracking detector are most relevant for reconstruction of photon candidates.

Electromagnetic calorimeter



3. A photon candidate in ATLAS

Photon reconstruction

- Photon energy calibration
- Photon identification
- Photon isolation

Photons in ATLAS detector are reconstructed through:

- ➤ interactions with the EM calorimeters → energy deposited in a cluster of calorimeter cells
- ➤ interactions upstream of the calorimeter (possible) → conversions to electron pairs: leaving tracks matched to a EM cluster



- **Electrons:** clusters matched to ID track from a vertex in interaction region
- **Converted photons:** clusters matched to a track from a conversion vertex
- Unconverted photons: clusters without matching tracks

- Photon Reconstruction
- Photon energy calibration
- Photon identification
- Photon isolation

(See Hicham' <u>talk</u> yesterday)

- ➤ interactions with the EM calorimeters → energy deposited in a cluster of calorimeter cells
- ➤ interactions upstream of the calorimeter (possible) → conversions to electron pairs: leaving tracks matched to a EM cluster

Photon energy corrected by a dedicated energy calibration:

MC based calibration
Calibrate the cluster energy to the original electron or photon energy



➢ absolute energy scale (data-driven)
Correct for data/MC difference using Z → ee samples



- validation with $J/\psi \rightarrow ee$ and $Z \rightarrow ll\gamma$ samples
 - Photon specific uncertainties



Electron and photon energy calibration with the ATLAS detector using 2015-2016 LHC protonproton collision data

- Photon Reconstruction
- Photon energy calibration
- Photon identification
- Photon isolation

Reject background candidates after reconstruction using:

- Leakage in the hadronic calorimeter
- Cluster shape in 2nd layer of EM calorimeter
- Cluster shape in 1st layer of EM calorimeter (tighter)



Photon identification with ATLAS detector

- Photon Reconstruction
- Photon energy calibration
- Photon identification
- Photon isolation

Photons from hard process (e.g. from resonance decays) expected to be well isolated from hadronic activity.

- Track based isolation variable: scalar sum of transverse momenta of tracks in a cone around γ candidate.
- Calorimeter-based isolation variable: sum of the transverse energy of the clusters in a cone around γ candidate (subtracting candidate contribution)



4. Search for new resonance in γγ channel

Analysis status

Search for a new resonances in $\gamma\gamma$ channel:



Run 1 result

Run 1 (2011-2012):

 \rightarrow search for narrow resonances in [65;600] GeV

Run 2 (2015-2018): <u>A low mass [65;110] GeV CONF note (ICHEP 2018) with 2015 – 2017 data</u>

Analysis overview



Gluon fusion production mode:

- LHC: gluon-gluon collider
- Largest cross-section for SM Higgs



Run 2 low-mass search: 65-110 GeV





Signal Modeling

Narrow-width resonance: shape dominated by the detector resolution

Shape of signal: modeled using a double-sided Crystal Ball (DSCB) function.

6 parameters describing:

- A Gaussian core
- power-law low-end and high-end tails

Parameters extracted from MC.



Background modeling

Continuum:

- real γγ events: irreducible (from MC samples)
- Jet faking photons (γ+jet, multi-jet): reducible (from data-driven control regions)

Resonant:

- Drell Yan Z/γ*→ee events misidentified as γγ (from di-electron data sample)
- SM Higgs (negligible)

→ Build background templates of each components

Background modeling

Continuum:

- real yy events: irreducible (from MC samples)
- Jet faking photons (γ+jet, multi-jet): reducible (from data-driven control regions)

Resonant:

- ➢ Drell Yan Z/γ*→ee events misidentified as yy (from di-electron data sample)
- SM Higgs (negligible)

Continuum (γγ, γ+jet):

- two components added together according to their respective fraction measured in data
- described by an analytic function

Resonant (Drell-Yan):

- normalized to the amount of di-electron events faking diphoton events
- modeled using a DSCB function

Background modeling

Check the quality of background modelling: we hope there's no "spurious signal"

→ Signal+background fit to a backgroundonly template

> Continuum ($\gamma\gamma$, γ +jet):

- two components added together according to their respective fraction measured in data
- described by an analytic function
- Resonant (Drell-Yan):
 - In normalized to the amount of di-electron events faking diphoton events
 - modeled using a DSCB function

Ideal: A good oyster template! ③



No room for spurious signal

Reality: a bad oyster template...? 😕



Spurious signal everywhere!

5. Results

Post-fit distributions

Background-only fit:

DY peak is clearly visible. Most prominent in the CC category, as expected.

No structure seen in the residuals.



Summary of systematics

Source	Uncertainty [%]	Remarks	
Signal yield			
Luminosity	±2		
Trigger eff.	$\pm 1.4 - 1.7$	mx-dependent	
Photon identification eff.	$\pm 1.5 - 2.3$	m_X -dependent	
Isolation eff.	± 4	•	
Photon energy scale	$\pm 0.13 - 0.49$	m_X -dependent	
Photon energy resolution	$\pm 0.053 - 0.28$	m_X -dependent	
Pile-up	$\pm 1.8 - 4.1$	m_X -dependent	
Production mode	$\pm 2.4 - 25$	m_X -dependent	
Signal modeling			
Photon energy scale	$\pm 0.3 - 0.5$	m_X - and category-dependent	
Photon energy resolution	$\pm 2 - 8$	m_X - and category-dependent	
Migration between categories			
Material	-2.0/+1.0/+4.1	category-dependent (UU/CU/CC)	
Non-resonant Background			
Spurious Signal	128/104/79	ratio to the expected spurious signal uncertainty	
	(604/496/181 events)	(category-dependent)	
DY Background modeling			
Peak position	$\pm 0.1 - 0.2$	category-dependent	
Peak width	$\pm 2 - 3$	category-dependent	
Normalization	±9-21	category-dependent	

Results



No significant excess with respect to the background-only hypothesis is observed.

An upper limit at the 95% CL is set on $\sigma_{fid} \cdot \mathcal{B}$ from 30 to 101 fb in the range 65 < m < 110 GeV.

CMS results: comparison



 \sim 2.9 σ local excess at 96 GeV, not seen by ATLAS.

Conclusion and further plan

In search for a new resonances <u>below the Higgs mass</u> in $\gamma\gamma$ channel:

- ATLAS sees no significant excesses above 1σ
- Not confirming the CMS excess (but can't exclude it yet)

Started analysis with Full 2015 – 2018 dataset and Full mass range:

- ➤ Very low mass range: below 65 GeV?
- Low-mass range: [65-110] GeV
- ➢ Intermediate mass range: [110-200] GeV
- High-mass range: above 200 GeV for spin0 and spin2

Optimizations (systematics, templates, etc) ongoing...



Back up

Standard Model

A theory of fundamental particles and how they interact.

- Elementary fermions (half-integer spin): 3 generations of quarks and leptons
- **Gauge bosons** (integer spin): 4 force carriers of fundamental interactions
 - Gluon (strong interaction)
 - Photon (electromagnetic interaction)
 - W and Z boson (weak interaction)
- **Higgs boson**: One last missing piece of SM, discovered in 2012 at the LHC

 \succ h(125), scalar boson, spin = 0

→ Currently our best description of elementary particles and their interaction. However, the standard model is incomplete.

Standard Model of Elementary Particles



Standard Model

Extension of Standard Model

• Standard Model: only one SU(2) doublet

$$\Phi = \begin{pmatrix} \Phi^+ \\ \Phi^0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi_1 + i\phi_2 \\ \phi_3 + i\phi_4 \end{pmatrix}, \phi \text{ are normalized real scalar fields}$$

• Two-Higgs-doublet model(2HDM): simple possible extension of SM

$$\Phi_j = \begin{pmatrix} \Phi_j^+ \\ (v_j + \rho_j + i\eta_j)/\sqrt{2} \end{pmatrix}, j=1,2$$

• 2 complex scalar SU(2) doublets \rightarrow 8 fields:

> 3 get eaten to give mass to W and Z gauge bosons

5 are physical scalar (Higgs) fields

H or h might be the H(125) discovered 2012 → What about the other? ▶1 neutral CP-odd: A
▶2 neutral CP-even: H, h with $m_H > m_h$ ▶2 charged: H[±]

Lagrangian giving mass terms in 2HDM

• With minimization of the potential, the mass terms are given by:

•
$$\mathcal{L}_{\Phi^{\pm} mass} = [m_{12}^2 - (\lambda_4 + \lambda_5)v_1v_2](\phi_1^- \phi_2^-) \begin{pmatrix} \frac{v_2}{v_1} & -1\\ -1 & \frac{v_1}{v_2} \end{pmatrix} \begin{pmatrix} \phi_1^+\\ \phi_2^+ \end{pmatrix}$$

• $\mathcal{L}_{\eta mass} = \frac{m_A^2}{v_1^2 + v_2^2} (\eta_1 - \eta_2) \begin{pmatrix} v_2^2 & -v_1v_2\\ -v_1v_2 & v_1^2 \end{pmatrix} \begin{pmatrix} \eta_1\\ \eta_2 \end{pmatrix}$
• $\mathcal{L}_{\rho mass} = -(\rho_1 - \rho_2) \begin{pmatrix} m_{12}^2 \frac{v_2}{v_1} + \lambda_1v_1^2 & -m_{12}^2 + \lambda_{345}v_1v_2\\ -m_{12}^2 + \lambda_{345}v_1v_2 & m_{12}^2 \frac{v_1}{v_2} + \lambda_2v_2^2 \end{pmatrix} \begin{pmatrix} \rho_1\\ \rho_2 \end{pmatrix}$ with $\lambda_{345} = \lambda_3 + \lambda_4 + \lambda_5$

Yukawa couplings

Yukawa Lagrangian: (f couple to Higgs in SM: $\frac{m_f}{n}$)

 $\mathcal{L}_{Yukawa}^{2HDM}$

$$= -\sum_{f=u,d,l} \frac{m_f}{v} \left(\xi_h^f \bar{f} fh + \xi_H^f \bar{f} fH - i\xi_A^f \bar{f} \gamma_5 fA\right) - \left\{\frac{\sqrt{2}V_{ud}}{v} \bar{u} \left(m_u \xi_A^u P_L + m_d \xi_A^d P_R\right) dH^+ + \frac{\sqrt{2}m_l \xi_A^l}{v} \bar{v}_l l_R H^+ + H.c.\right\}$$

 $P_{L/R}$: projection operators for left-/right-handed fermions

The coupling of the neutral Higgs bosons to the W and Z are the same in all models.

	Type I	Type II	Lepton-specific	Flipped
ξ_h^u	$\cos \alpha / \sin \beta$	$\cos \alpha / \sin \beta$	$\cos \alpha / \sin \beta$	$\cos \alpha / \sin \beta$
ξ_h^d	$\cos \alpha / \sin \beta$	$-\sin \alpha / \cos \beta$	$\cos \alpha / \sin \beta$	$-\sin \alpha / \cos \beta$
ξ_h^ℓ	$\cos \alpha / \sin \beta$	$-\sin \alpha / \cos \beta$	$-\sin \alpha / \cos \beta$	$\cos \alpha / \sin \beta$
ξ^u_H	$\sin \alpha / \sin \beta$	$\sin \alpha / \sin \beta$	$\sin \alpha / \sin \beta$	$\sin \alpha / \sin \beta$
ξ_H^d	$\sin \alpha / \sin \beta$	$\cos \alpha / \cos \beta$	$\sin \alpha / \sin \beta$	$\cos \alpha / \cos \beta$
ξ_H^ℓ	$\sin \alpha / \sin \beta$	$\cos \alpha / \cos \beta$	$\cos \alpha / \cos \beta$	$\sin \alpha / \sin \beta$
ξ^u_A	\coteta	\coteta	\coteta	$\cot eta$
ξ^d_A	$-\cot\beta$	aneta	$-\cot\beta$	an eta
ξ^{ℓ}_A	$-\cot\beta$	aneta	aneta	$-\cot\beta$

Two-Higgs-doublet models

Most general **potential** for two doublets Φ_1 and Φ_2 :

•
$$V = m_{11}^2 \Phi_1^{\dagger} \Phi_1 + m_{22}^2 \Phi_2^{\dagger} \Phi_2 - m_{12}^2 (\Phi_1^{\dagger} \Phi_2 + \Phi_2^{\dagger} \Phi_1) + \frac{\lambda_1}{2} (\Phi_1^{\dagger} \Phi_1)^2 + \frac{\lambda_2}{2} (\Phi_2^{\dagger} \Phi_2)^2 + \lambda_3 \Phi_1^{\dagger} \Phi_1 \Phi_2^{\dagger} \Phi_2 + \lambda_4 \Phi_1^{\dagger} \Phi_2 \Phi_2^{\dagger} \Phi_1 + \frac{\lambda_5}{2} [(\Phi_1^{\dagger} \Phi_2)^2 + (\Phi_2^{\dagger} \Phi_1)^2]$$

All the parameters are real (5 independent coupling λ and 3 mass parameters m).

• Scalar doublets:
$$\Phi_j = \begin{pmatrix} \Phi_j^+ \\ (v_j + \rho_j + i\eta_j)/\sqrt{2} \end{pmatrix}$$
 j=1,2

2 complex scalar SU(2) doublets \rightarrow 8 fields:

> 3 get eaten to give mass to W and Z gauge bosons

➤ 5 are physical scalar (Higgs) fields.

Orthogonal combinations of $\rho_j \rightarrow$ physical scalars:

- $h = \rho_1 sin\alpha \rho_2 cos\alpha$
- $H = -\rho_1 cos \alpha \rho_2 sin \alpha$

See backup: Lagrangian for mass terms

▶1 neutral CP-odd: A
▶2 neutral CP-even: H, h with m_H > m_h
▶2 charged: H[±]

SM Higgs boson:

$$H^{SM} = \rho_1 cos\beta + \rho_2 sin\beta$$

= $hsin(\alpha - \beta) - Hcos(\alpha - \beta)$

Experiments at LHC



Seven experiments at the Large Hadron Collider (LHC) use detectors to analyse the myriad of particles produced by collisions in the accelerator.

ATLAS, CMS: general-purpose detectors, investigate the largest range of physics possible.
 ALICE, LHCb: detectors specialized for focusing on specific phenomena.
 TOTEM, LHCf: focus on "forward particles".
 MOEDAL: search for a hypothetical particle: the magnetic monopole.

LHC experiments

Designed parameters

Detector component	Required resolution	η coverage	
		Measurement	Trigger
Tracking	$\sigma_{p_T}/p_T = 0.05\% p_T \oplus 1\%$	±2.5	
EM calorimetry	$\sigma_E/E = 10\%/\sqrt{E} \oplus 0.7\%$	±3.2	±2.5
Hadronic calorimetry (jets)			
barrel and end-cap	$\sigma_E/E = 50\%/\sqrt{E} \oplus 3\%$	± 3.2	±3.2
forward	$\sigma_E/E = 100\%/\sqrt{E} \oplus 10\%$	$3.1 < \eta < 4.9$	$3.1 < \eta < 4.9$
Muon spectrometer	$\sigma_{p_T}/p_T = 10\%$ at $p_T = 1$ TeV	±2.7	±2.4

Higgs production



search in $\gamma\gamma$ channel



Branching ratio predicted by 2HDM: similar as SM with a fraction depending on model type

The $\gamma\gamma$ decay channel has the advantage of a clean experimental signature:

- excellent mass resolution
- smoothly falling background (diphoton production by QCD processes)



Event selection

Event selection:

- 2g20_tight trigger for <u>2015+2016 periods A-D3</u>
- 2g22_tight trigger for <u>2016 periods D4-after</u>
- 2g20_tight_icalovloose trigger for 2017
- ET{leading,subleading} > 22 GeV
- Tight photon ID
- Photon isolation: FixedCutLoose
- Invariant mass range :
- 60-120 GeV and Search in 65-110 GeV

Photon identification

Variable	Definition	Description	
Leakage in the hadronic calorimeter	$R_{had} = \frac{E_T^{had}}{E_T}$	Leakage in the hadronic calorimeter	
Middle η energy ratio	$\mathbf{R}_{\eta} = \frac{E_{3\times7}^{s2}}{E_{7\times7}^{s2}}$	$\phi E_{3 \times 7}$	
Middle ϕ energy ratio	$\mathbf{R}_{\phi} = \frac{E_{3\times3}^{s2}}{E_{3\times7}^{s2}}$	$E_{3\times 3}$	
Middle lateral width	$\omega_{\eta^2} = \sqrt{\frac{\sum E_i \eta_i^2}{\sum E_i} - (\frac{\sum E_i \eta_i}{\sum E_i})^2}$	Shower width in middle layer	
Front side energy ratio	$F_{side} = \frac{E(\pm 3) - E(\pm 1)}{E(\pm 1)}$	$egin{array}{c c c c c c c c c c c c c c c c c c c $	
Front lateral width (3 strips)	$\omega_{s3} = \sqrt{\frac{\sum E_i (i - i_{\max})^2}{\sum E_i}}$	Shower width in 3 strips around the hottest strip	
Front lateral width (total)	$\omega_{s,tot} = \sqrt{\frac{\sum E_i (i - i_{\max})^2}{\sum E_i}}$	Shower width in all strips Γ^{S1}	
Front second maximum difference	$\Delta E = \left[E_{2^{nd}max}^{s1} - E_{min}^{s1} \right]$	E ₁ stmax	
Front maxima relative ratio	$E_{ratio} = \frac{E_{1^{st}max}^{s1} - E_{2^{nd}max}^{s1}}{E_{1^{st}max}^{s1} + E_{2^{nd}max}^{s1}}$		

Signal Modeling

Narrow width resonance: shape dominated by the detector re

While there is no assumption on the production mode of the resonance, the shape of signal is modelled using a **double-sided crystal ball (DSCB)** function.



6 parameters of DSCB:

- $\succ \Delta m_X, \alpha_{Low}, \alpha_{High}, \sigma_{CB} \rightarrow$ mass dependent
- \succ *n*_{Low}, *n*_{High} → mass independent (constant)

Each parameter is determined in a fit to the fixed-mass simulated samples, and is parametrized as a function of mass separately for each conversion category.



Pre-fit distributions

Search range: 65~110 GeV (width of signal~1.5GeV)



- > Data in good agreement with the background template within uncertainties.
- Small excess around the DY region, but covered by the systematics.

Continuum backgrounds

irreducible (γγ):
 taken from high-statistics MC samples
 reducible (γ+jet, multi-jet):
 taken from data-driven control regions

Step 1: build a **template** (irreducible and reducible) representative of the non-resonant background.

Step 2: add the two parts together according to their **respective fraction** measured in data.



Resonant Drell-Yan backgrounds

Crucial ingredient for background estimation: $Z \rightarrow ee$ misidentified as $\gamma\gamma$

- Using a di-electron data sample to build a Drell-Yan template
- Normalize the Drell-Yan template to the amount of di-electrons events faking diphoton events



Spurious signal test









Statistical model

• The data are described using an extended PDF expressed as:

$$\mathcal{L} = \prod_{c=1}^{n_c} e^{-N_c^{total}} \prod_{i=1}^{n_c^{data}} \mathcal{L}_c(m_{\gamma\gamma}(i,c))$$

 $n_c = 3$: number of categories n_c^{data} : number of data events n_c^{total} : number of fitted events

• The per-event term is expressed as:

$$\mathcal{L}_{c}(m_{\gamma\gamma};\sigma_{fid},m_{X},N_{uu,c},N_{uc,c},N_{cu,c},N_{cc,c},N_{bkg,c},c_{c},\theta) =$$

 σ_{fid} : fiducial production cross-section of the new resonance $N_{xx,c}$: number of DY background events identified as (and contribute to) UU, UC or CC

N_{bkg,c}: fitted number of background events

 c_c : collectively refers to the background parameters used to describe its shape

 θ : collectively designates the nuisance parameters used to describe the systematic uncertainties

= $N_{X,c}(\sigma_{fid}, m_X, \theta_{N_X}, \theta_{SS}) f_X(m_{\gamma\gamma}, m_X, x_X(m_X), \theta_{\sigma})$

- + $N_{uu,c}(\theta_{N_{uu,c}})f_{uu,c}(m_{\gamma\gamma}, x_{uu,c}, \theta_{uu,c})$
- + $N_{uc,c}(\theta_{N_{uc,c}})f_{uc,c}(m_{\gamma\gamma}, x_{uc,c}, \theta_{uc,c})$
- + $N_{cu,c}(\theta_{N_{cu,c}})f_{cu,c}(m_{\gamma\gamma}, x_{cu,c}, \theta_{cu,c})$
- + $N_{cc,c}(\theta_{N_{cc,c}})f_{cc,c}(m_{\gamma\gamma}, x_{cc,c}, \theta_{cc,c})$
- + $N_{bkg,c}f_{bkg,c}(m_{\gamma\gamma},c_c)$

The continuum background PDF $f_{bkg,c}$ is described by the function chosen for each category mentioned before.