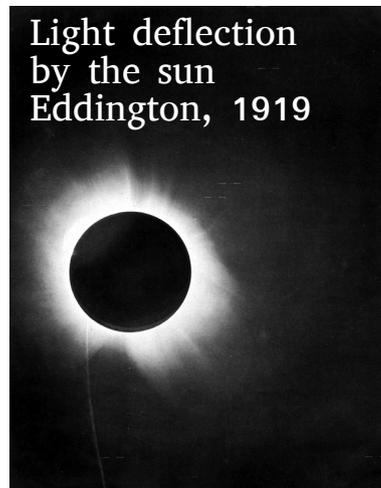
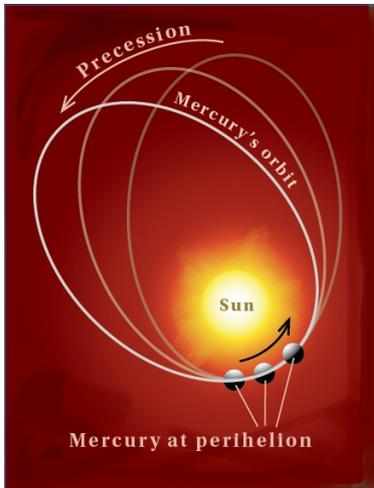
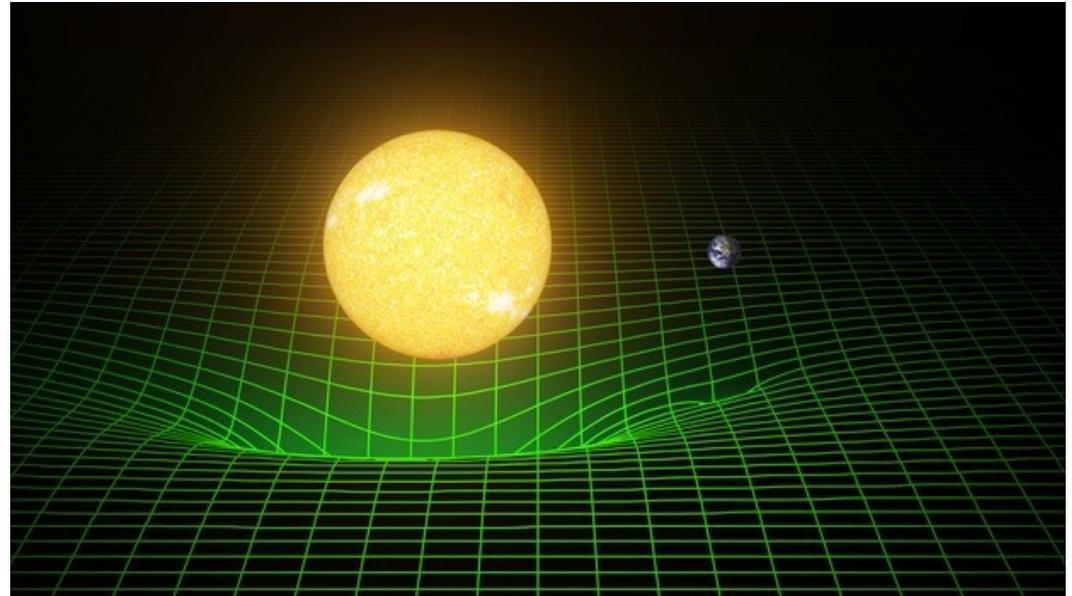


# Data analysis for gravitational-wave searches

- Gravitational-wave detectors
- Gravitational-wave sources
- Signal extraction methods
- Parameter estimation
- Open data

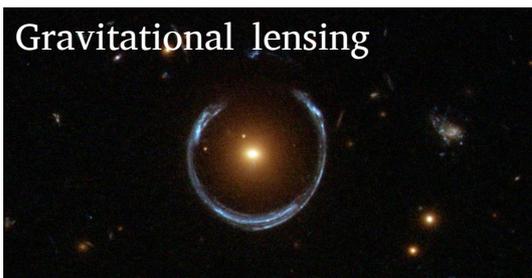
# General Relativity

- Gravity is no longer a force (Newton)
- Gravity = geometric property of space and time (Einstein)
- Space-time curvature dictated by the mass distribution and/or energy radiation
- Einstein's equation: 
$$G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$



## Specific predictions:

- Fine description of orbiting bodies (Mercury anomalous perihelion shift)
- Light propagation affected by gravity (gravitational lensing, gravitational redshift, Shapiro delay...)
- Black holes
- Gravitational waves

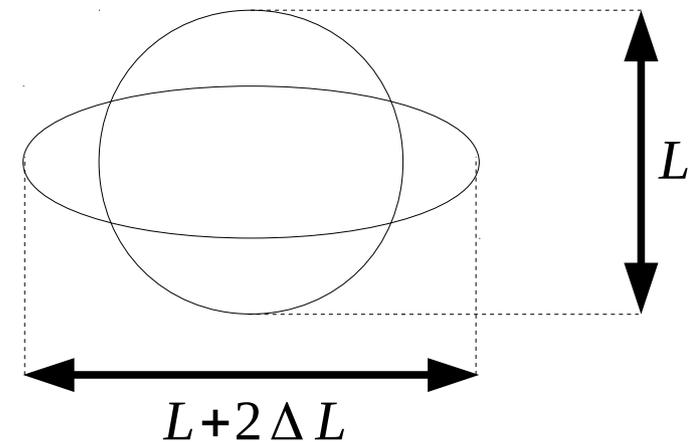
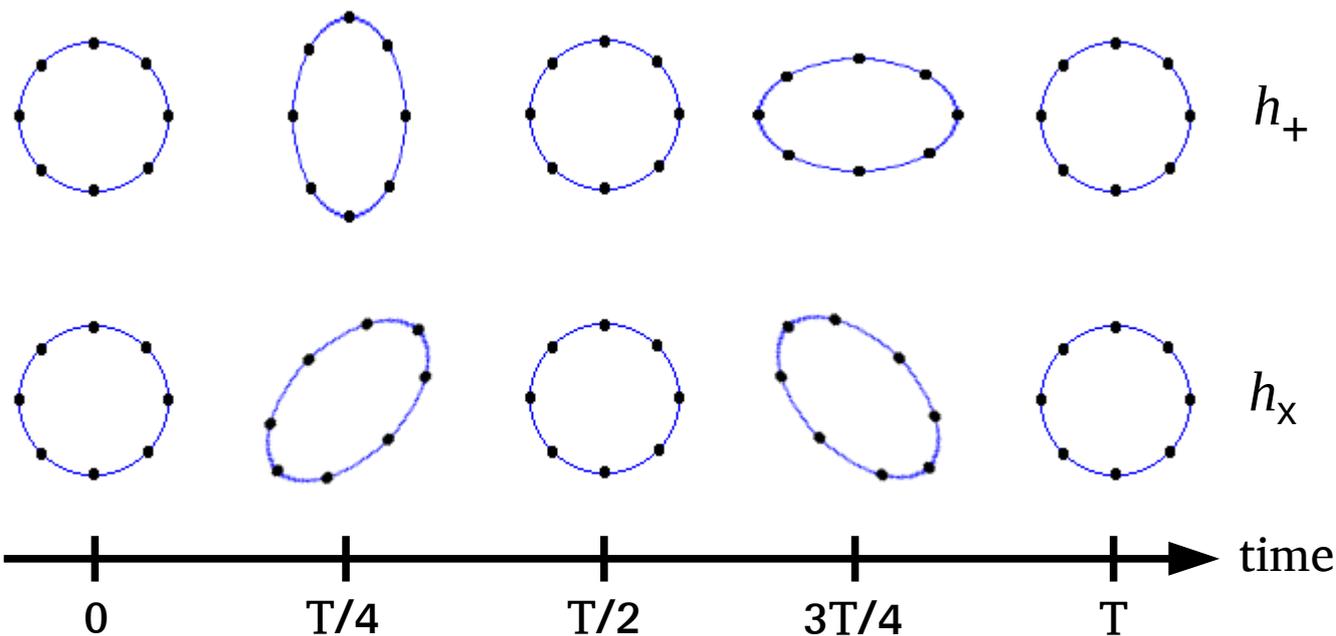


# Gravitational waves

Add a small perturbation to the Minkowski metric:  $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$       $|h_{\mu\nu}| \ll 1$

- $h$  obeys a plane-wave equation
- the wave propagates at the speed of light
- 2 degrees of freedom:  $h_+$  and  $h_x$

## → Gravitational waves



$$h = 2 \frac{\Delta L}{L}$$

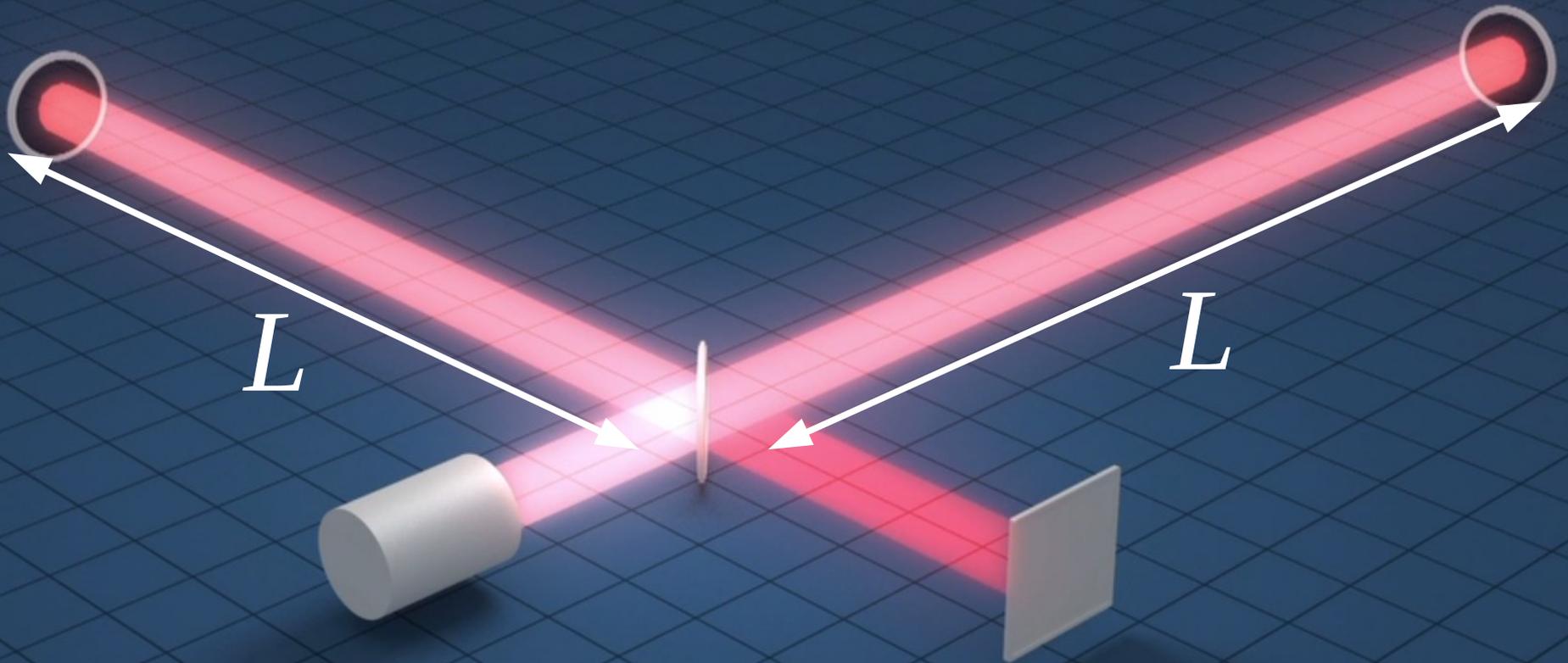
When they propagate, gravitational waves

- do not interact with matter
- are attenuated by  $1/r$

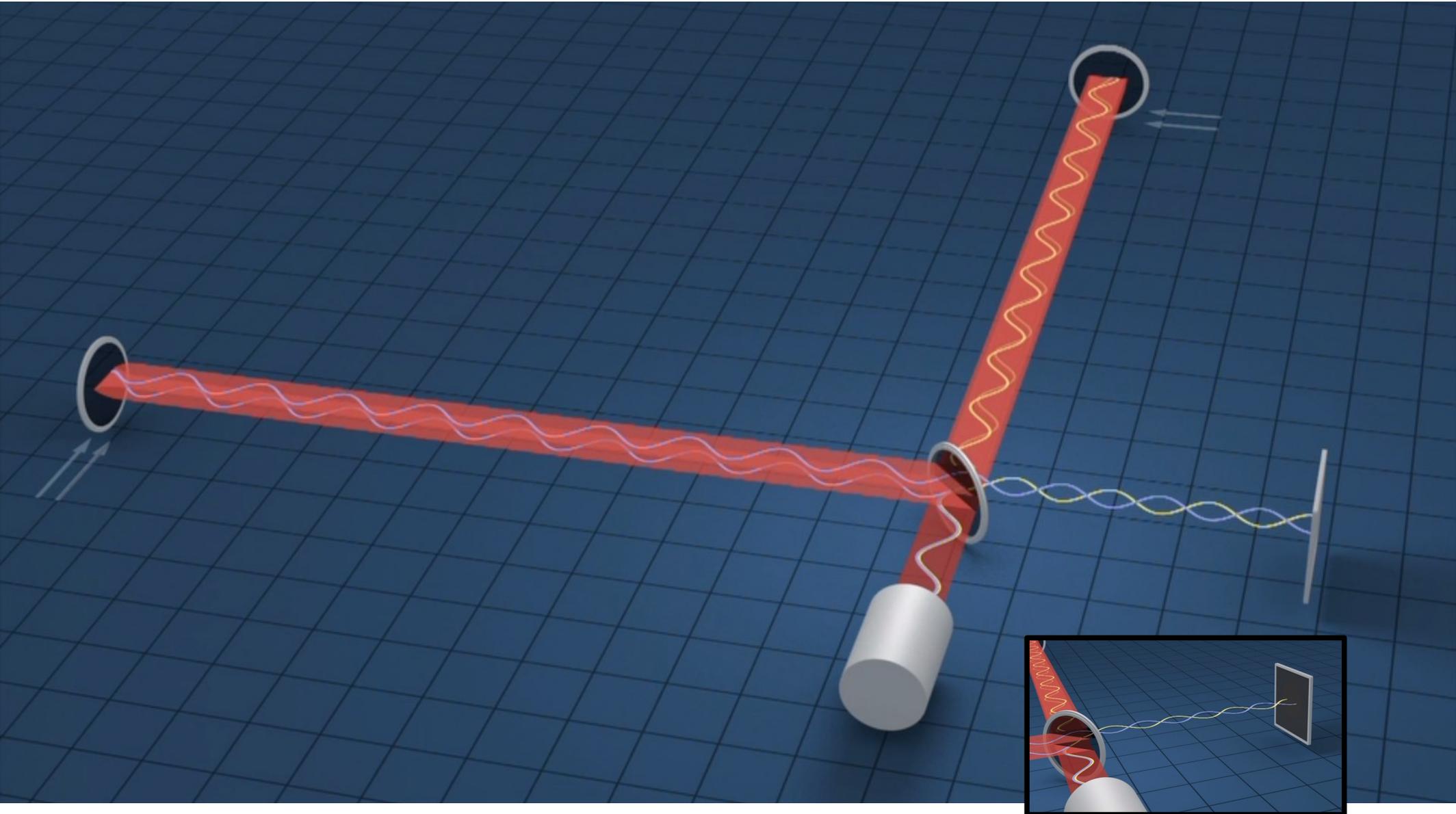
→ **Gravitational waves are the perfect probe!**

→ **BUT...  $h \sim 10^{-21}$**

## Michelson interferometer

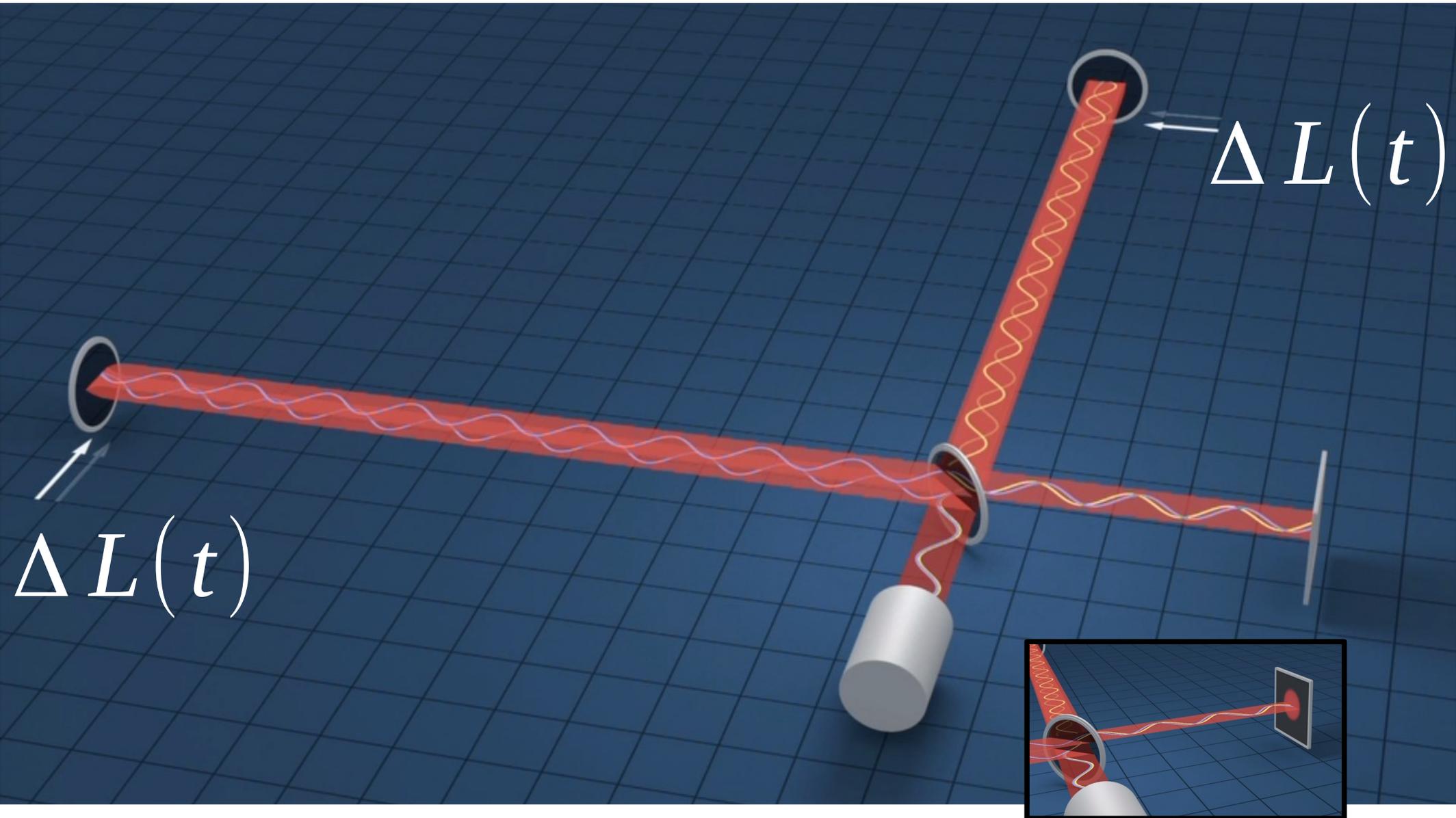


# Gravitational-wave detection: interferometry



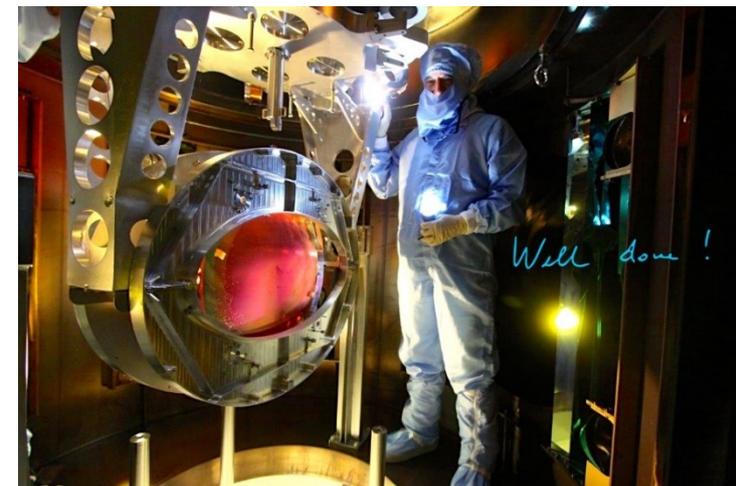
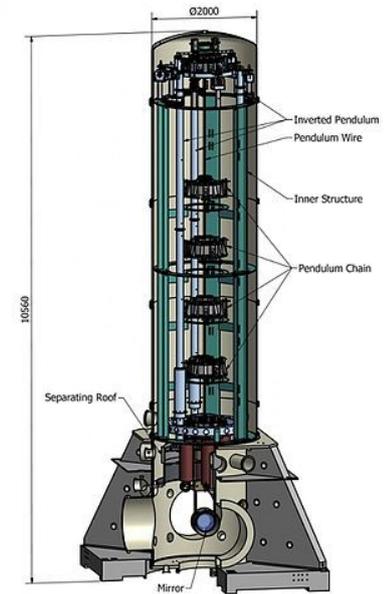
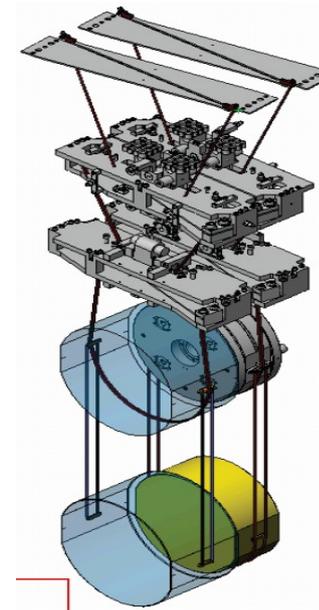
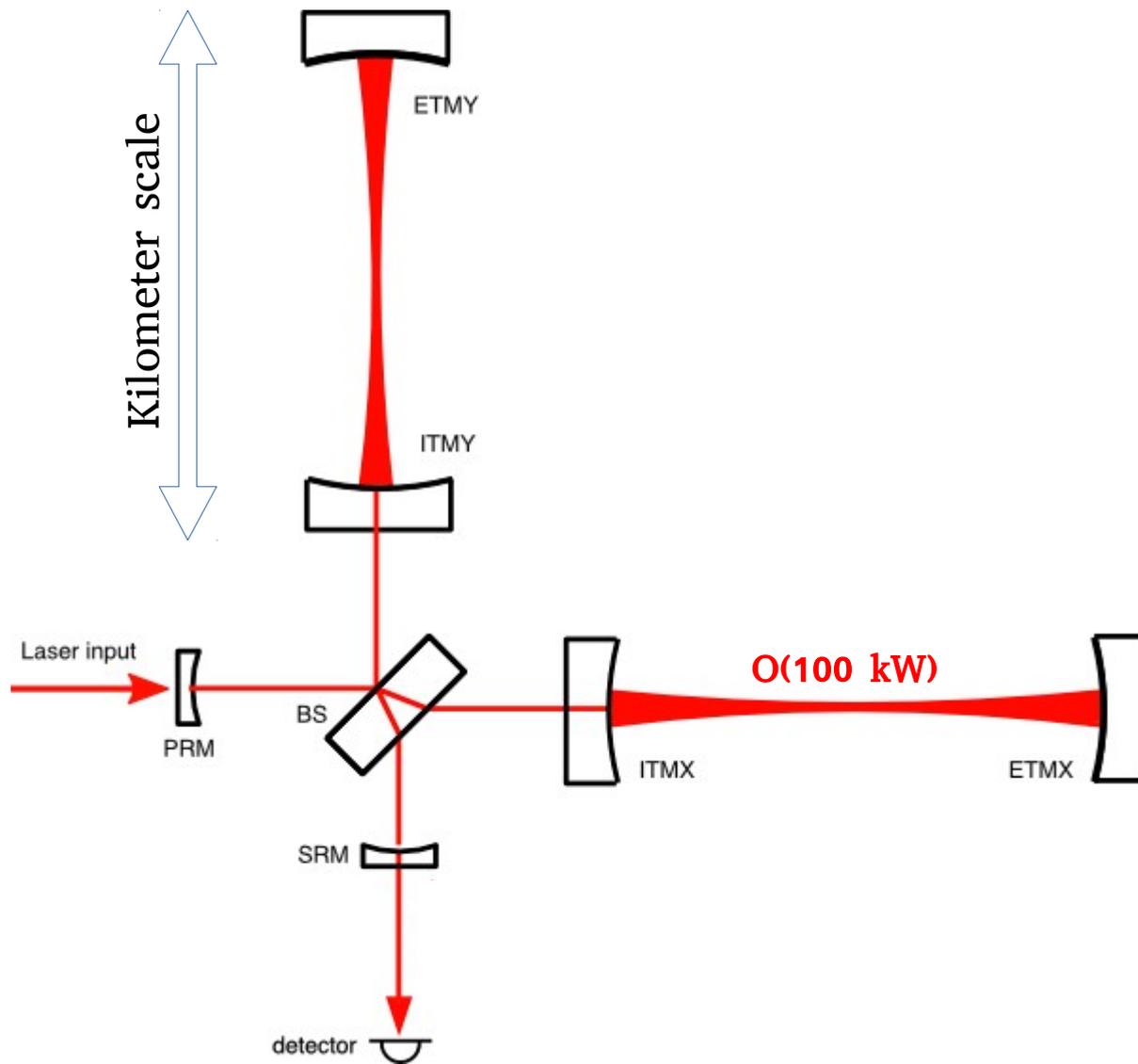
Destructive interference<sup>5</sup>  
→ no light

# Gravitational-wave detection: interferometry



Constructive interference<sup>6</sup>  
→ bright light

# Gravitational-wave detection: ground-based interferometers



# Gravitational-wave detection: ground-based interferometers

LIGO Livingston (USA)



LIGO Hanford (USA)



Virgo (Italy)



Kagra (Japan)

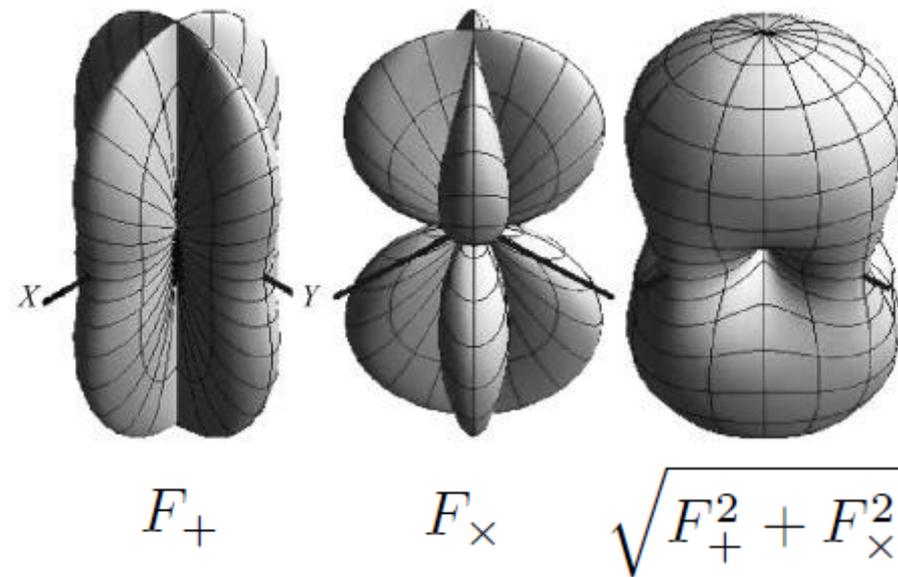


LIGO India



# Gravitational-wave detection: ground-based interferometers

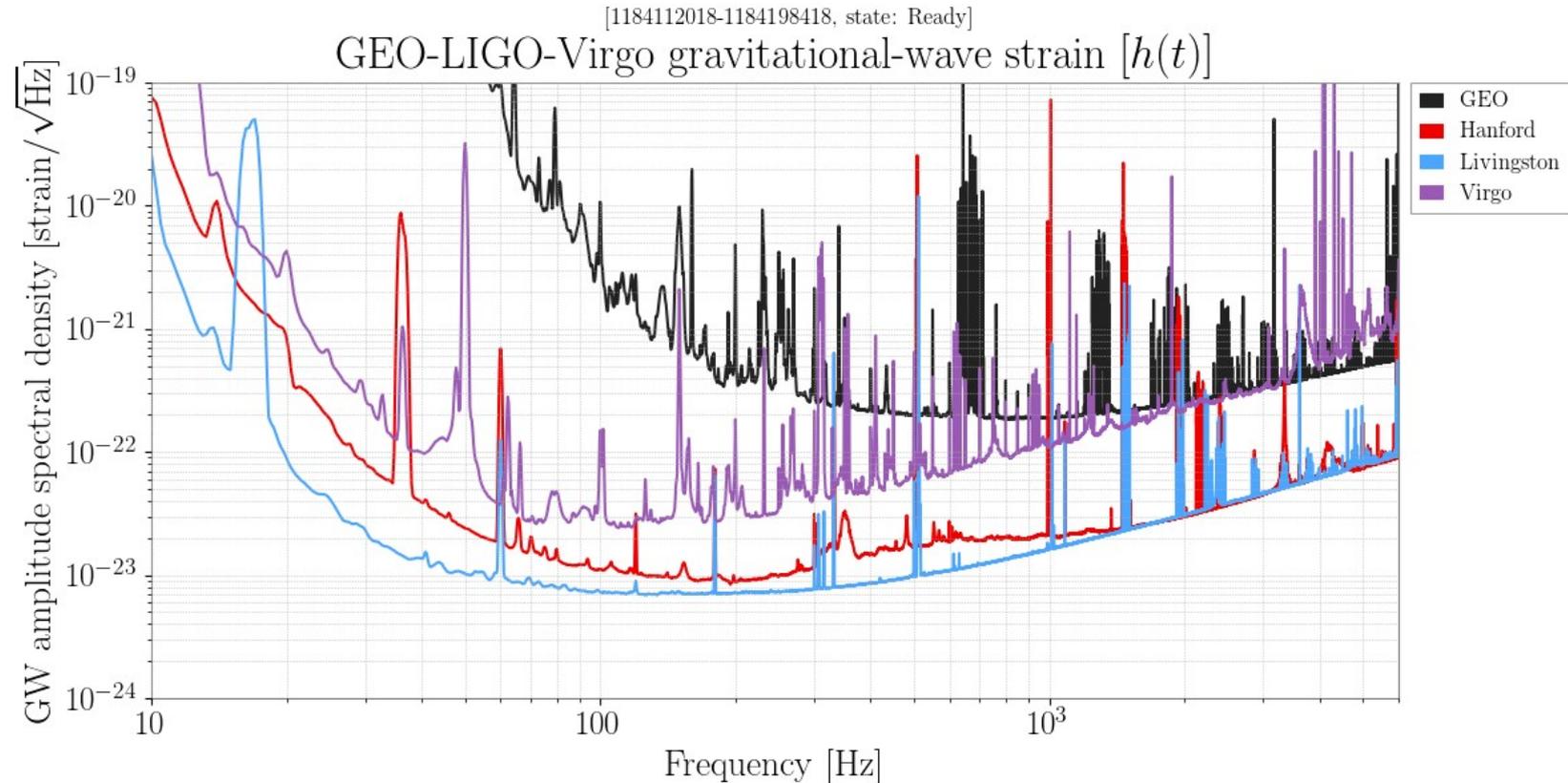
The sensitivity of the detector is not uniform over the sky



$$GW(t) = F_+(t, \underline{ra}, \underline{dec}, \underline{\Psi}) \times h_+(t) + F_x(t, \underline{ra}, \underline{dec}, \underline{\Psi}) \times h_x(t)$$

Source position      Source polarization angle

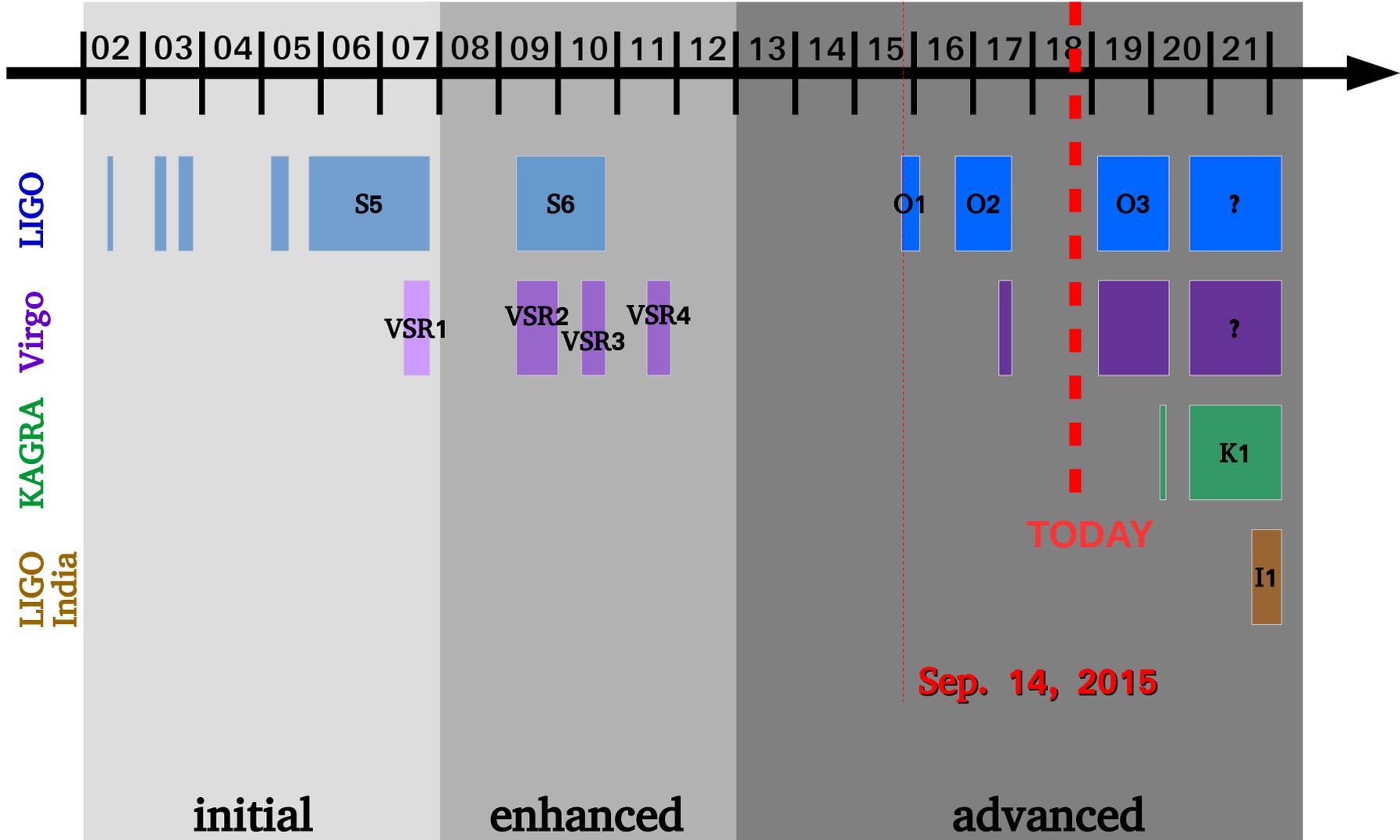
# Gravitational-wave detection: ground-based interferometers



The sensitivity of ground-based interferometers is limited by

- 3 fundamental noise sources
  - seismic noise ( $< 10$  Hz)
  - thermal noise ( $10 \rightarrow 100$  Hz)
  - quantum noise ( $> 100$  Hz)
- Technical noise
  - control noise
  - scattered light
  - electronic noise
  - frequency noise
  - ...

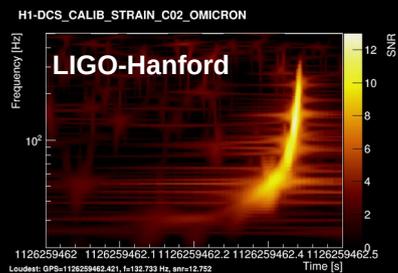
# Gravitational-wave detection: ground-based interferometers



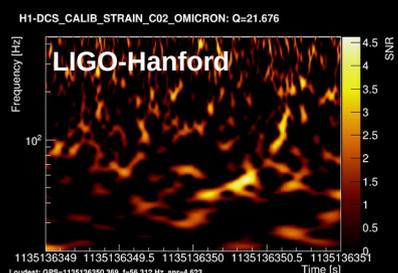
Data analysis :

- O1 : ~50 days of data, 2 detectors
- O2 : ~100 days of data, 2 detectors
- O3 : ~200 days of data, 3 detectors

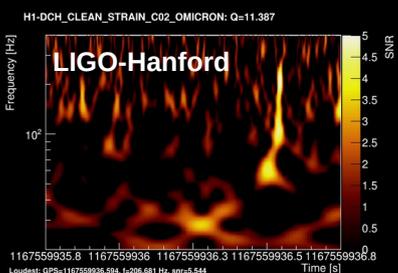
# O1 & O2 detections



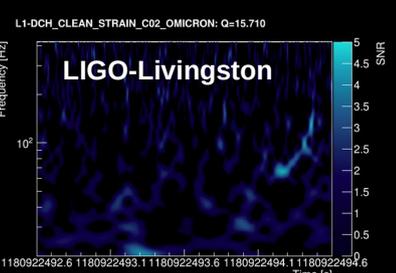
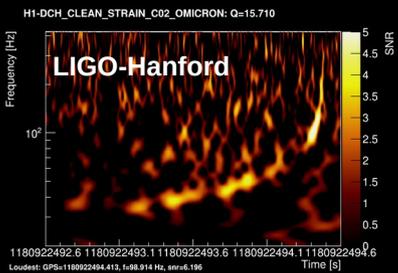
**GW150914**



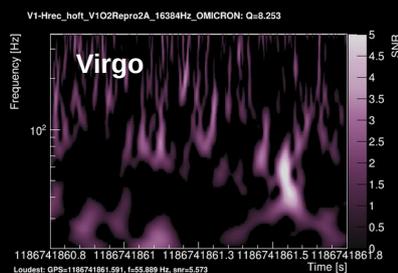
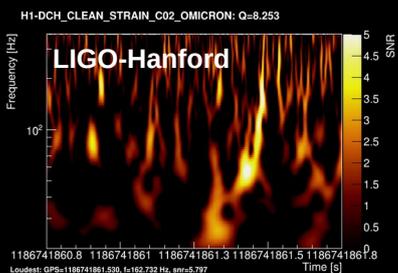
**GW151226**



**GW170104**



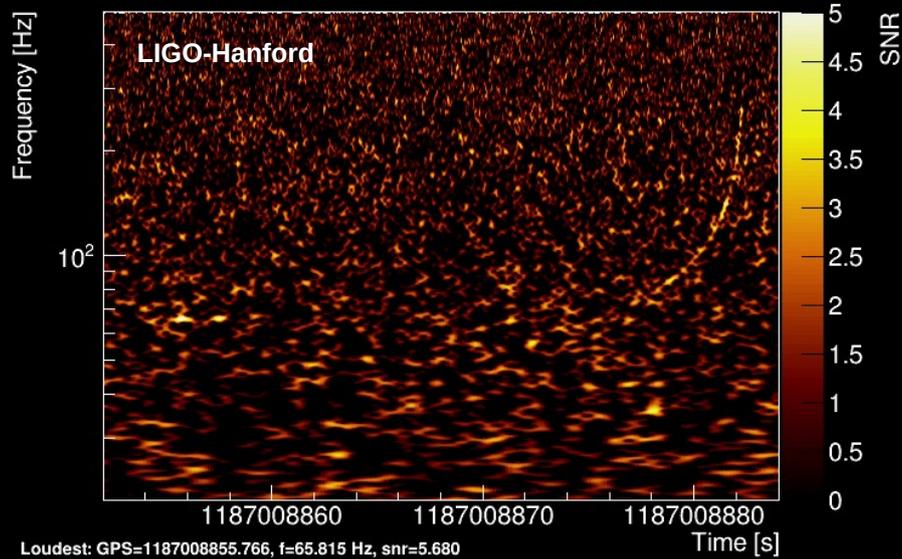
**GW170608**



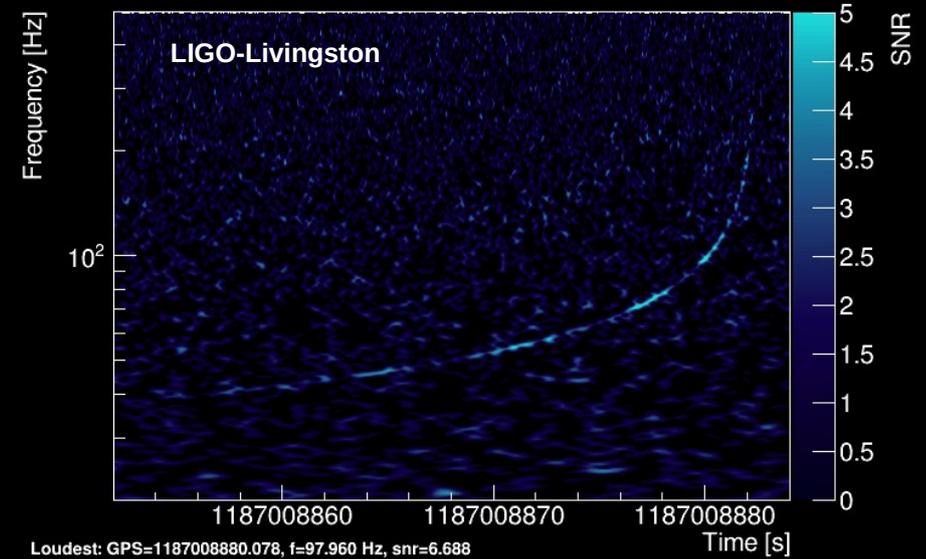
**GW170814**

# GW170817

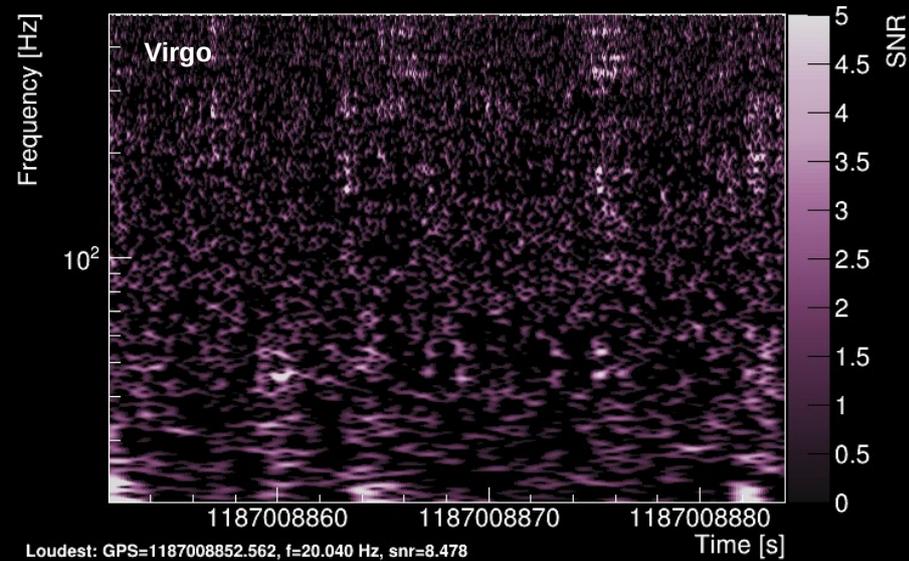
H1-DCH\_CLEAN\_STRAIN\_C02\_OMICRON: Q=99.459



L1-DCH\_CLEAN\_STRAIN\_C02\_T1700406\_v3\_OMICRON: Q=99.459



V1-Hrec\_hoft\_V1O2Repro2A\_16384Hz\_OMICRON: Q=99.459



What's next ?

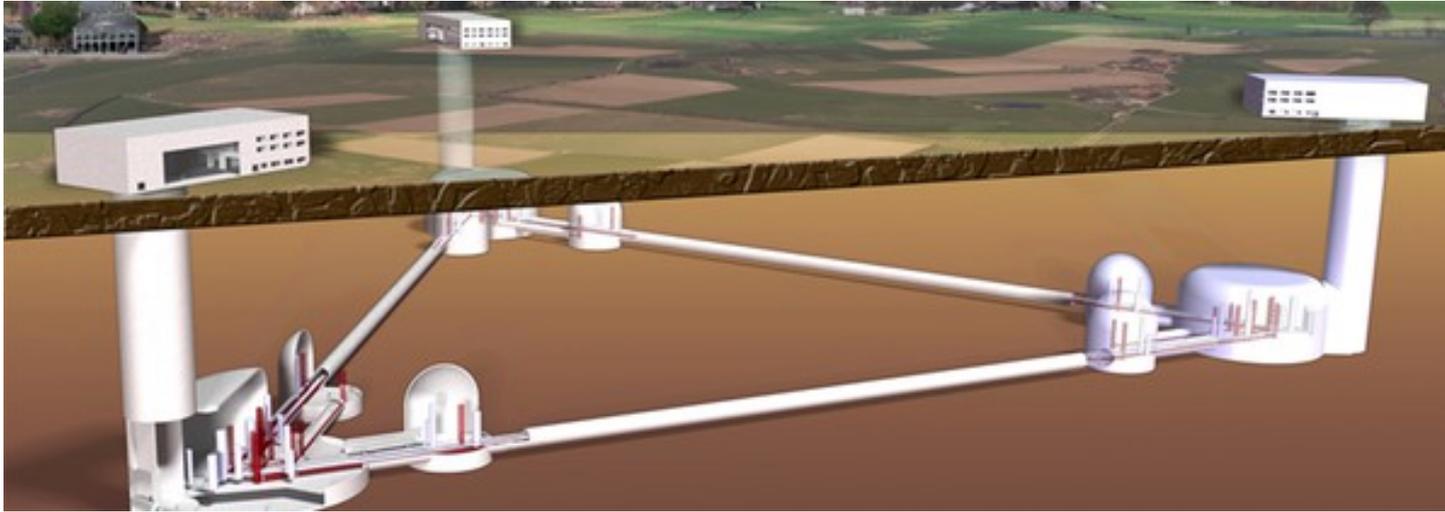
# Gravitational-wave detection: Advanced Virgo +

## Advanced Virgo +

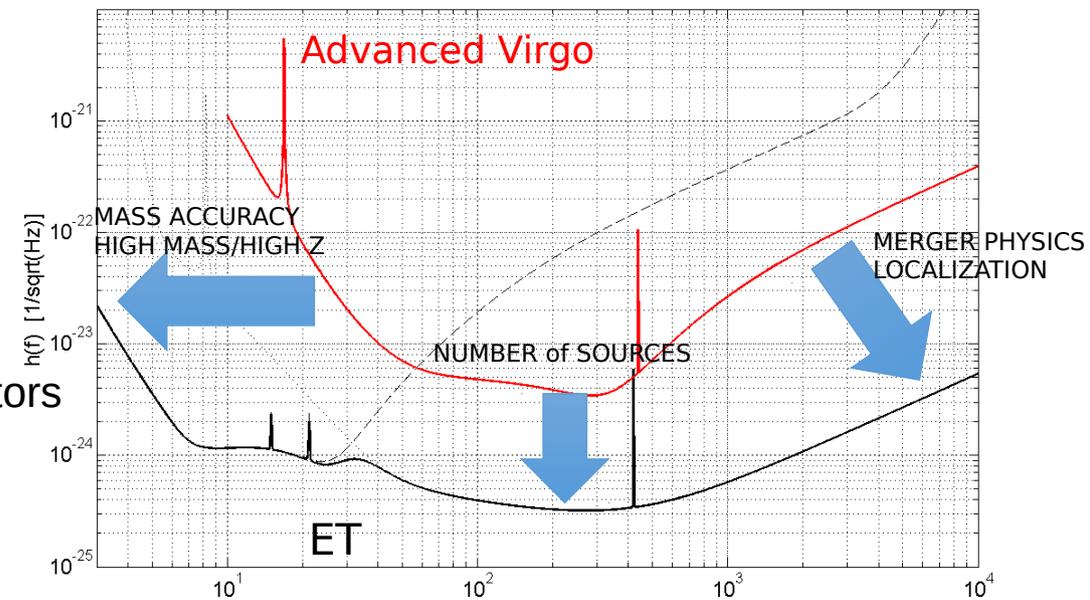


- Use Virgo infrastructure → push further technologies (laser power, mirror size, coatings...)
- Improve sensitivity by a factor of a few
- Widening the frequency band

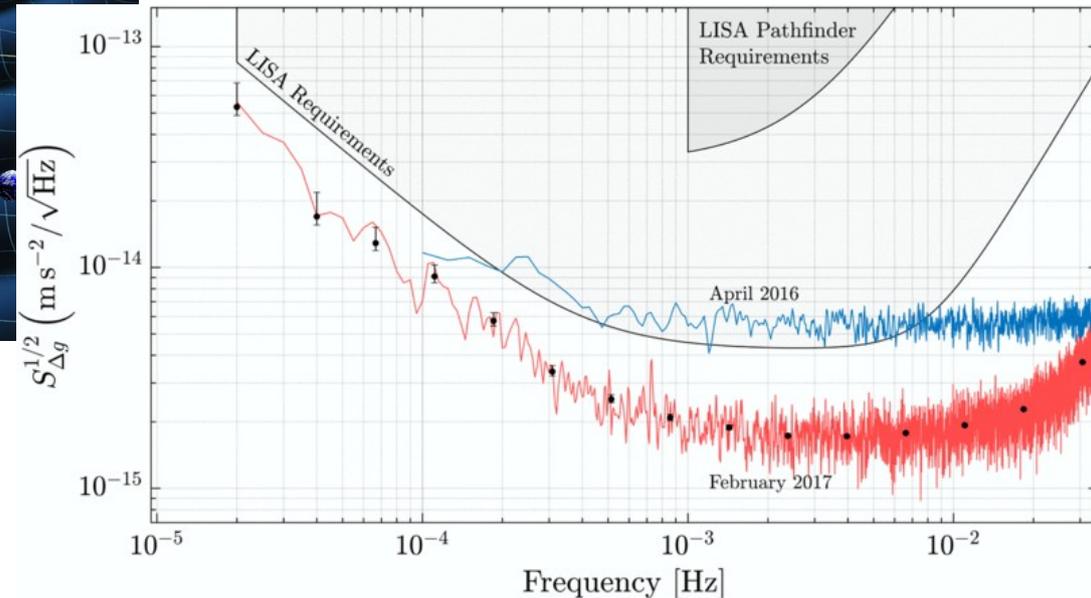
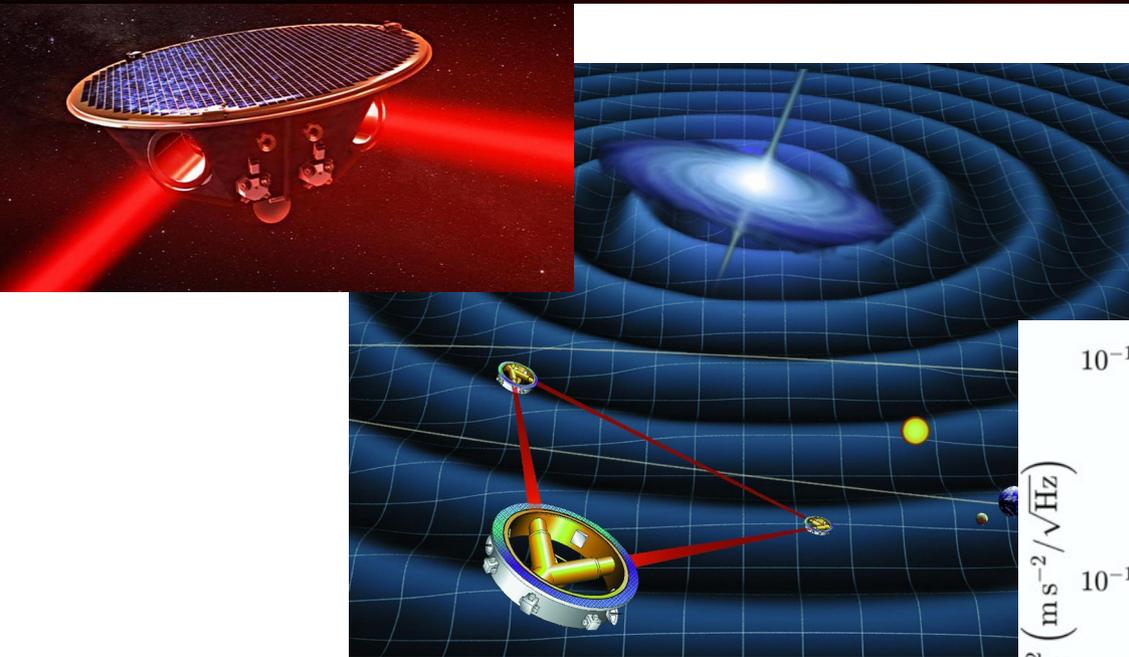
# Gravitational-wave detection: Einstein telescope



- Underground to reduce seismic noise
- 10 km arms
- Cryogenic mirrors
- Lower frequency limit – 1 Hz
- 10 x better sensitivity than 2nd generation detectors
- Farther back in the universe

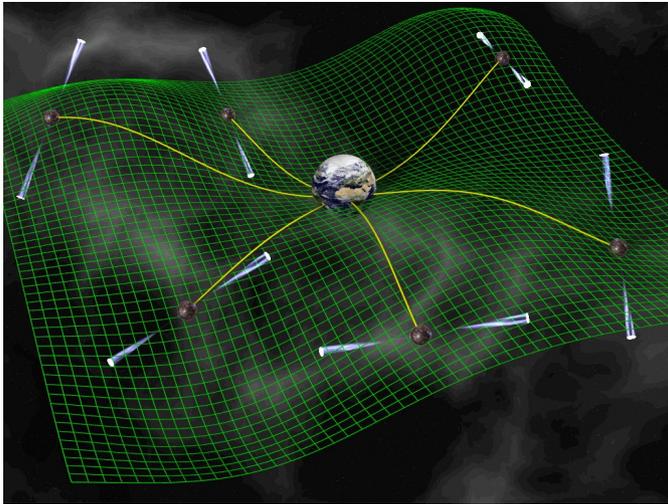


# Gravitational-wave detection: LISA



- ESA – All Systems GO!
- NASA coming back
- LIGO GW events and Lisa Pathfinder success have helped significantly
- Tremendous activity at present
- Planned launch 2034
- Earlier launch? 2030?
- 4 year mission → 10 years?

# Gravitational-wave detection: pulsar timing



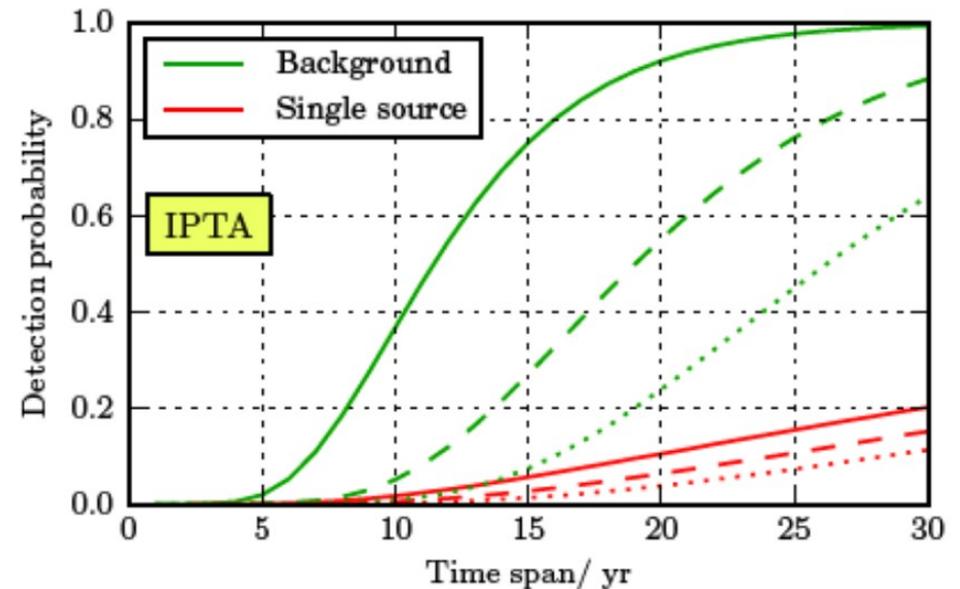
Distant pulsars send regular radio pulses

→ highly accurate clocks.

A passing gravitational wave would change the arrival time of the pulse.

Numerous collaborations around the world. Interesting upper limits and likely detections in the near future.

## European pulsar timing array



# Gravitational-wave sources

~~Monopole~~

$$\cancel{m = \int \rho d^3 \vec{r}}$$

~~Dipole~~

$$\cancel{P_i = \int \rho x_i d^3 \vec{r}}$$

Quadrupole (traceless)

$$Q_{ij} = \int \rho \left( x_i x_j - \frac{1}{3} r^2 \delta_{ij} \right) d^3 \vec{r}$$

Einstein quadrupole formula (radiated power)

$$\frac{dE}{dt} = -\frac{G}{5c^5} \left\langle \frac{d^3 Q^{ij}}{dt^3} \frac{d^3 Q_{ij}}{dt^3} \right\rangle$$

Estimate using the source parameters

$$Q \sim \varepsilon M R^2$$

$$\frac{d^3 Q}{dt^3} \sim \varepsilon M R^2 \omega^3$$

$$\frac{dE}{dt} \sim -\frac{G}{c^5} \varepsilon^2 M^2 R^4 \omega^6 \sim -\frac{c^5}{G} \varepsilon^2 \left( \frac{R_s}{R} \right)^2 \left( \frac{v}{c} \right)^6$$

$\simeq 10^{52} \text{ W}$

→ Important source characteristics:

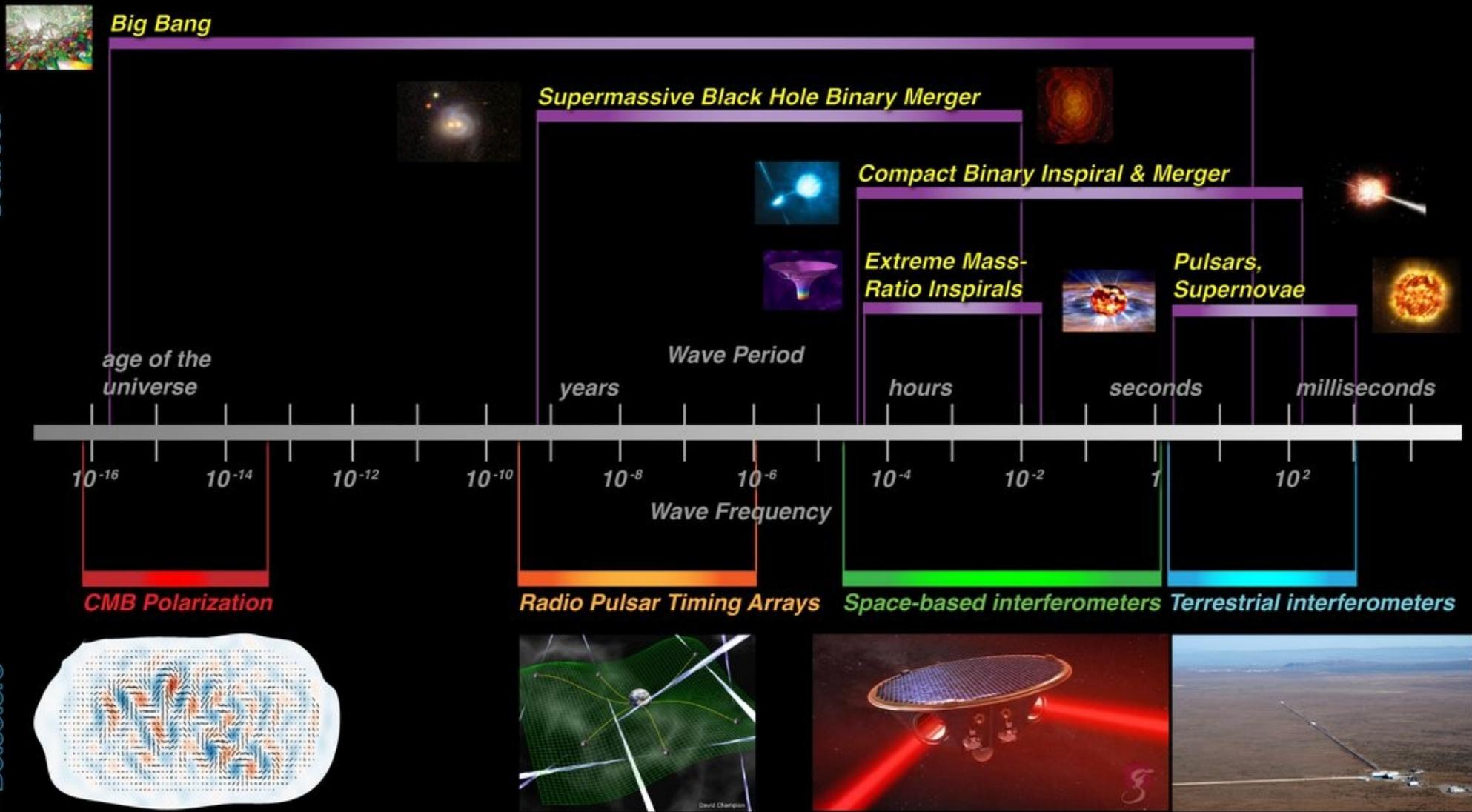
- asymmetric
- compact
- relativistic

# Gravitational-wave sources

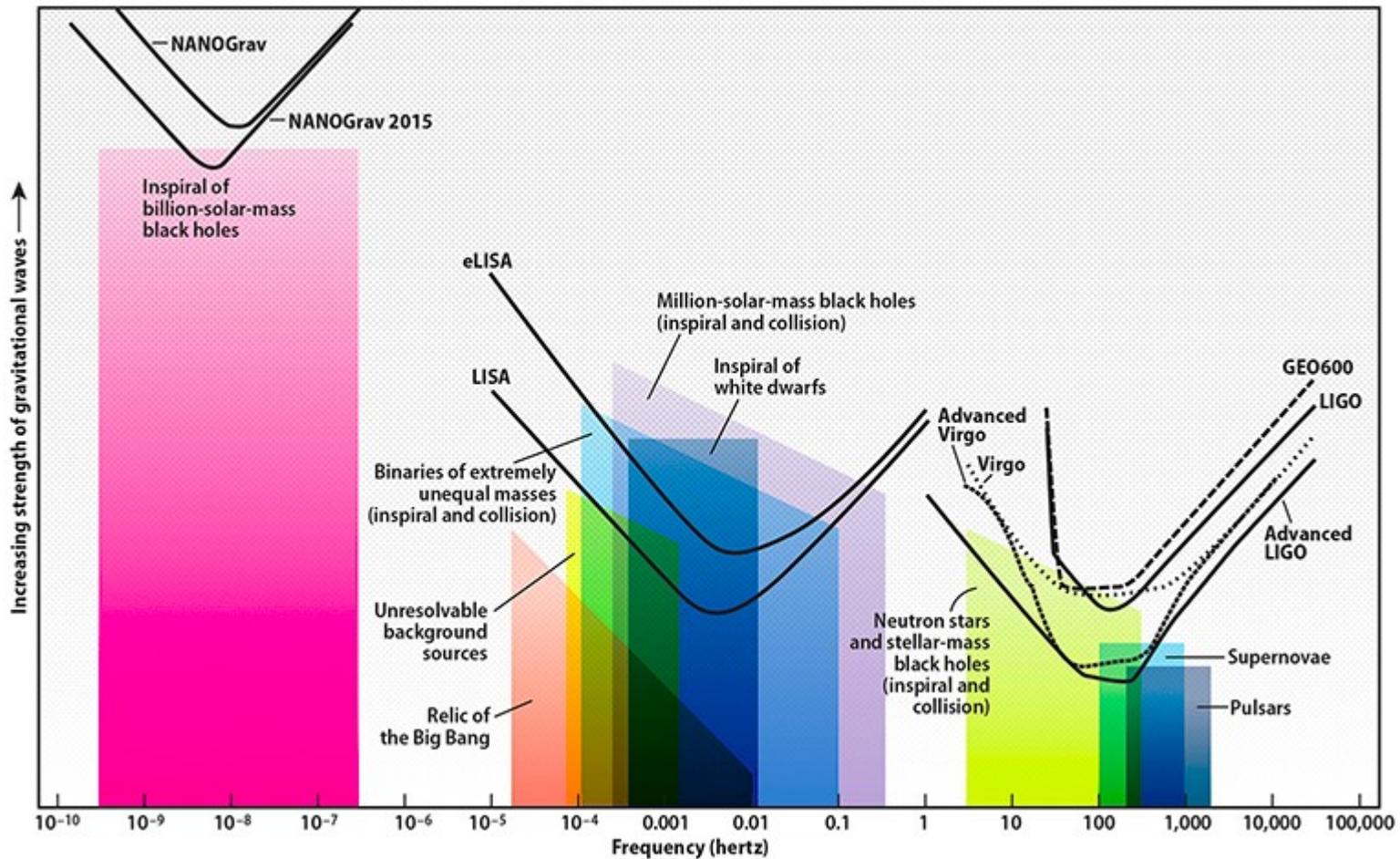
The Gravitational Wave Spectrum

Sources

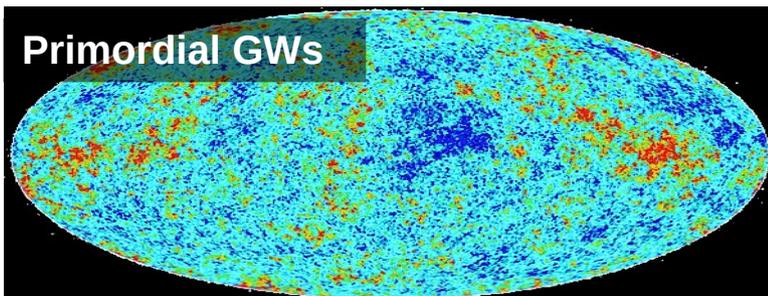
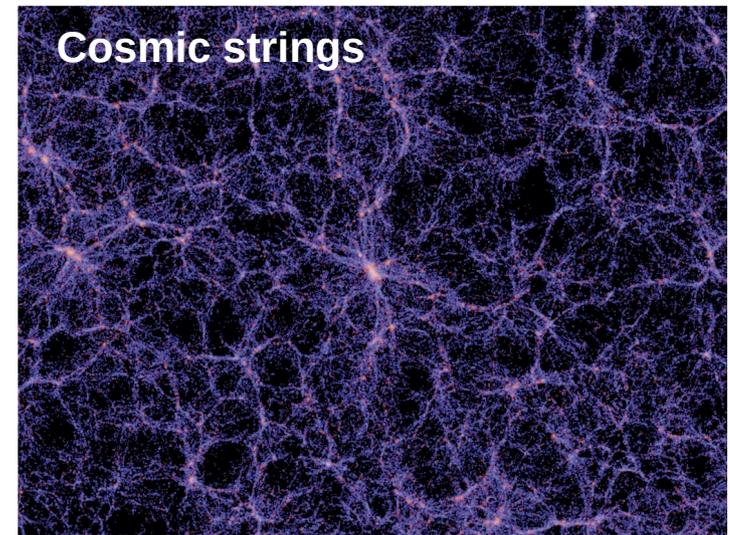
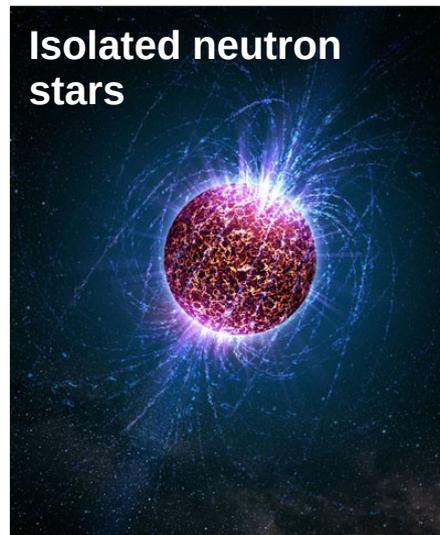
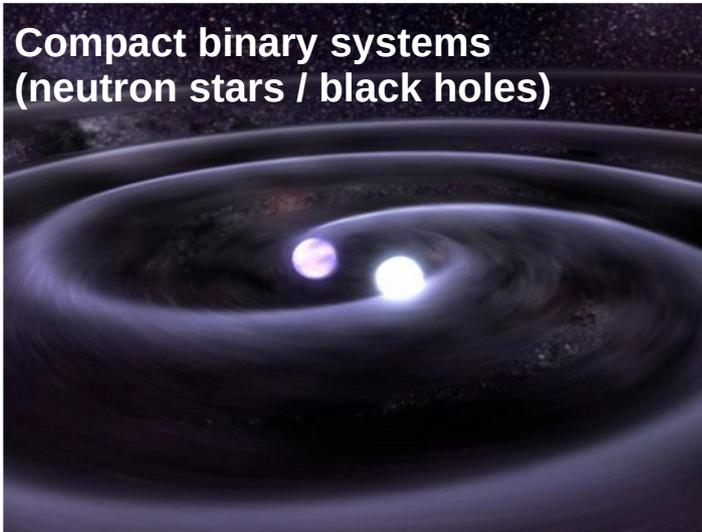
Detectors



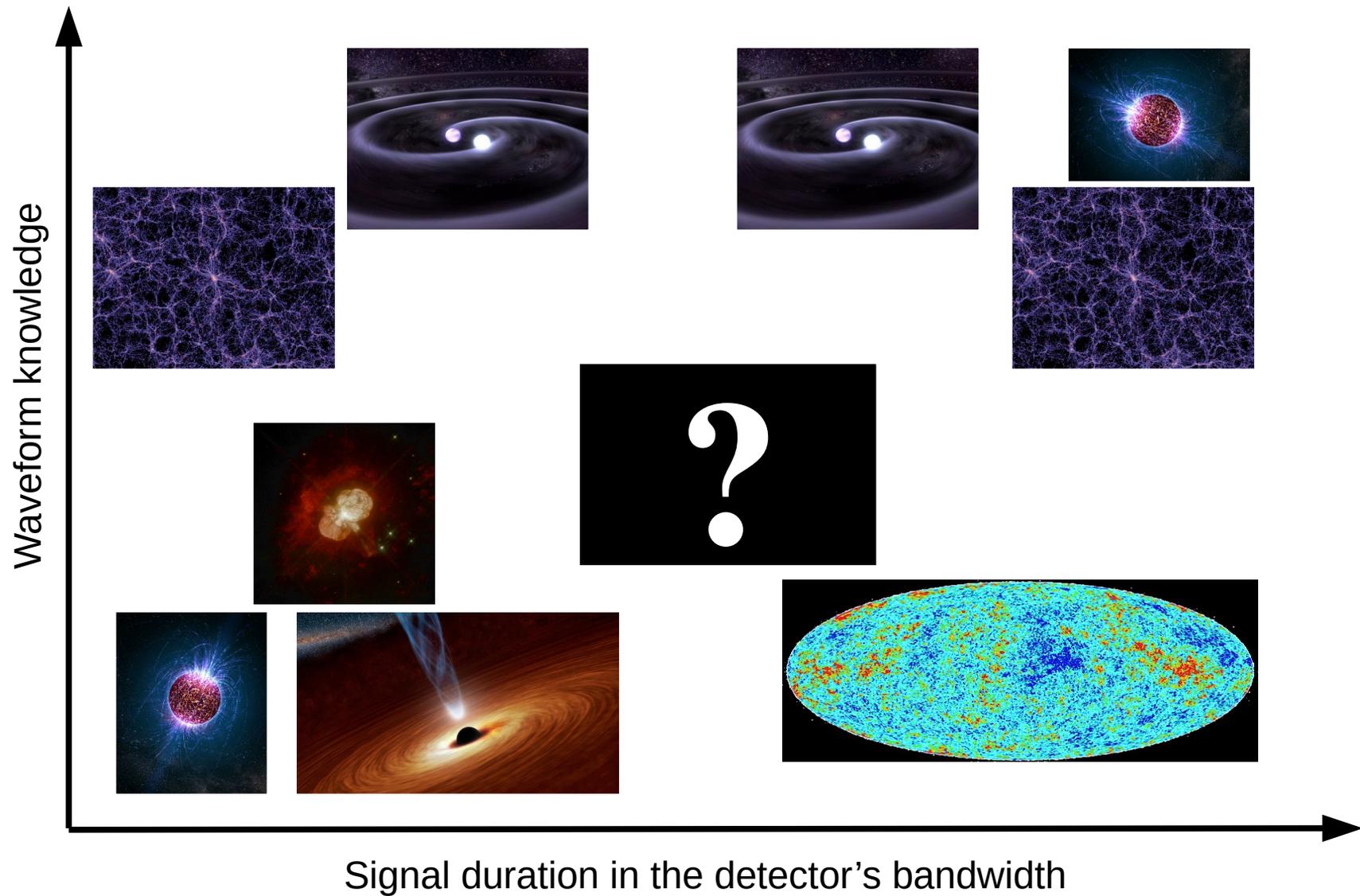
# Gravitational-wave sources



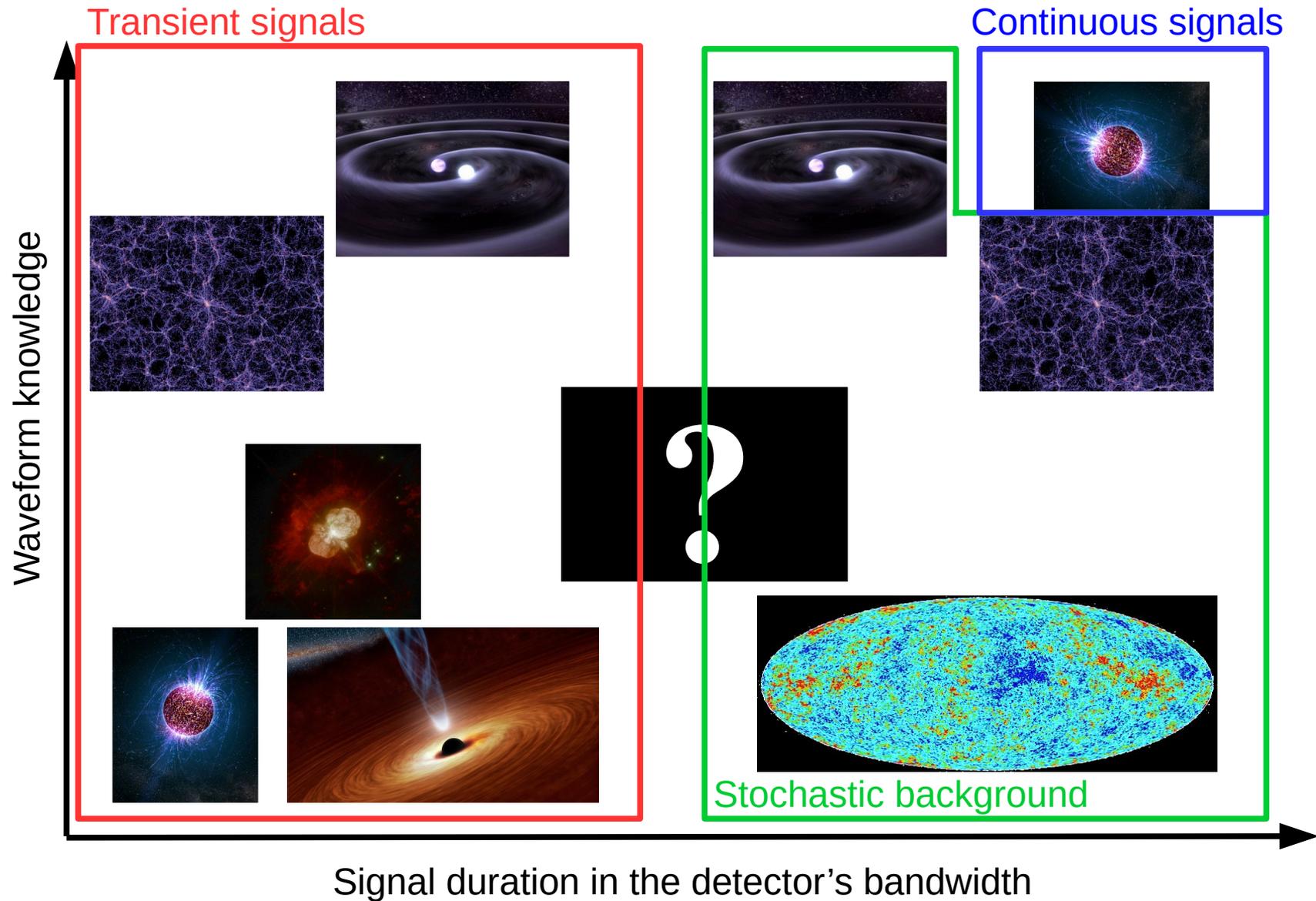
# Relevant gravitational-wave sources ( $> 10$ Hz)



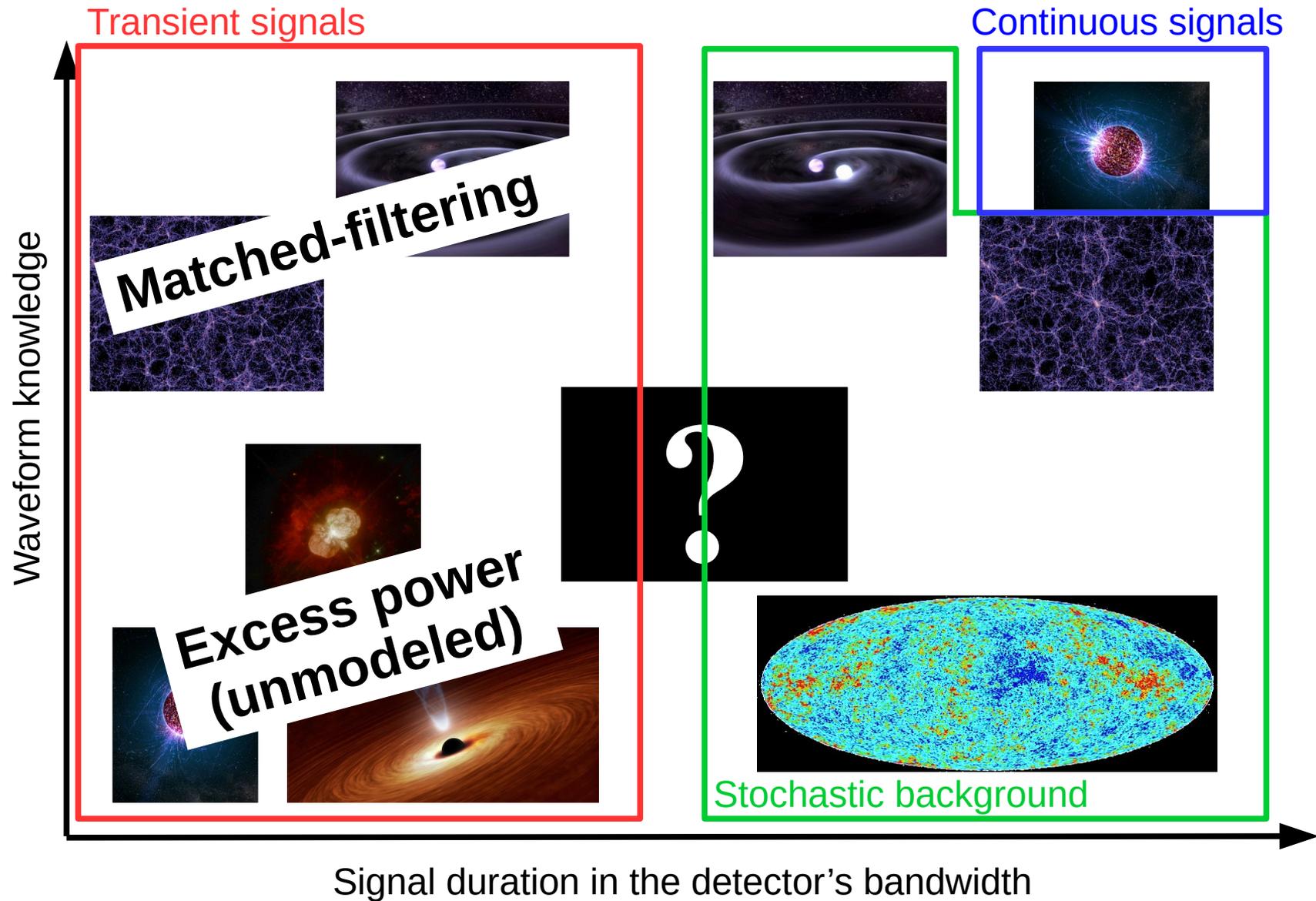
# Source classification : analysis methods



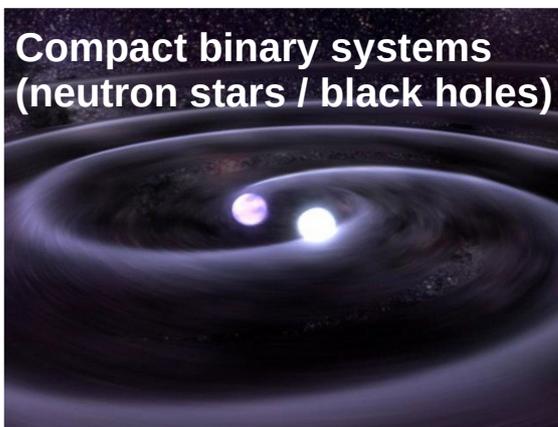
# Source classification : analysis methods



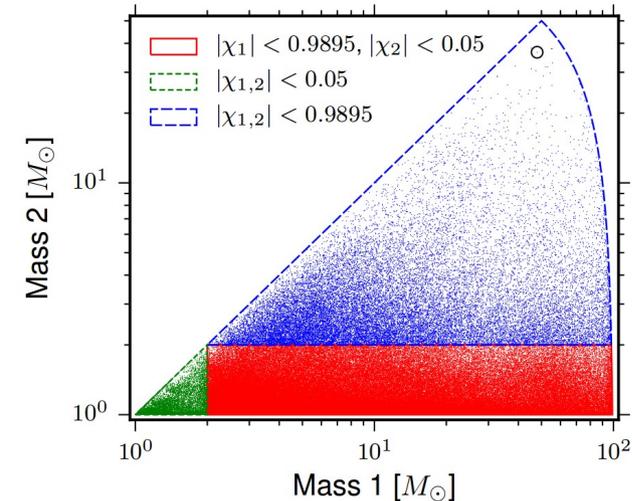
# Source classification : analysis methods



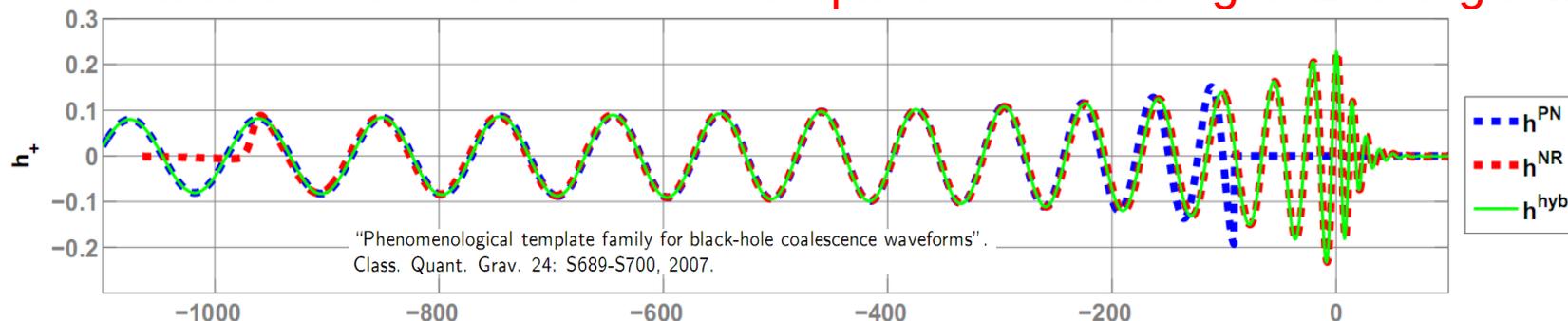
# CBC searches



- Compact binary objects: Two neutron stars and/or black holes.
- Inspiral toward each other. Emit gravitational waves as they inspiral.
- Amplitude and frequency of the waves increases over time, until the merger.
- Waveform relatively well understood,  $\rightarrow$  matched-filtering template searches ( $\sim 250,000$  templates).
- Unique way to study strong field gravity and the structure of the nuclear matter in the most extreme conditions

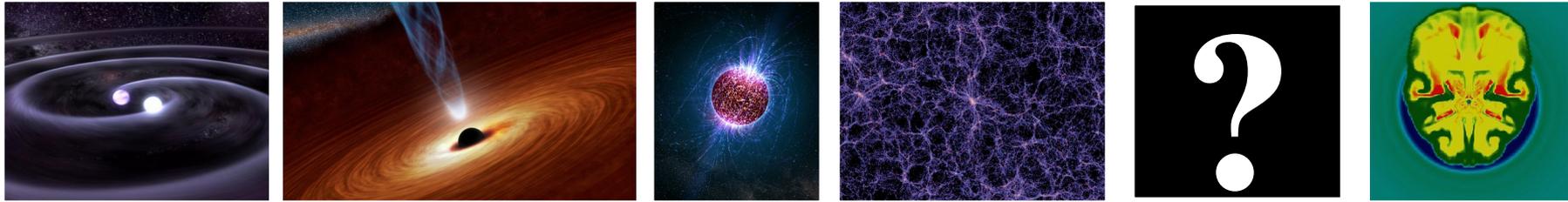


Gravitational waveform: inspiral merger BH-ringdown



Waveform carries lots of information about binary masses, orbit, merger, spins, ...

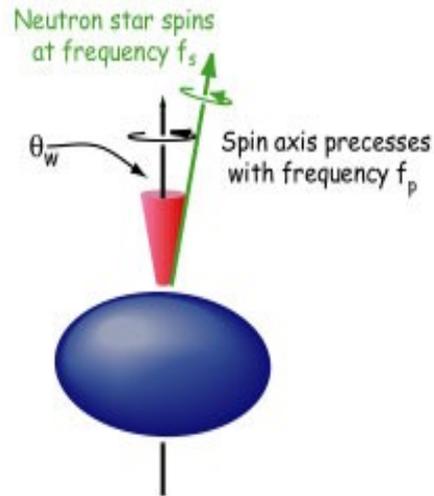
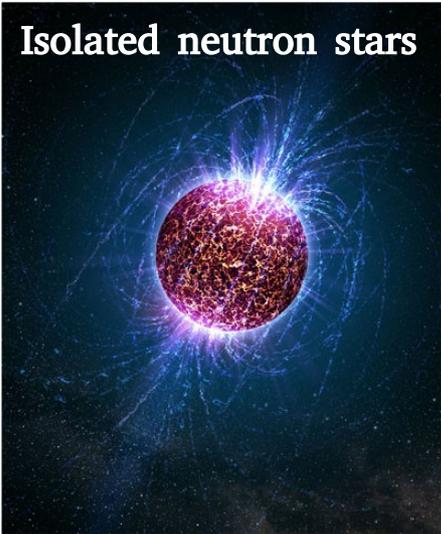
# Burst searches



- Many transient sources:
  - CBC
  - Supernovae: probe the explosion mechanisms.
  - Gamma Ray Bursts: collapse of rapidly rotating massive stars or neutron star mergers.
  - Pulsar glitches.
  - Cosmic strings cusps and kinks.
- Models are ok, but not essential:
  - Search for power excess in the data.
  - Search for any short signal with measurable strain signal.

# Continuous-wave searches

Isolated neutron stars



Persistent signals associated to sources with mass quadrupole moment varying in time in a nearly periodic way

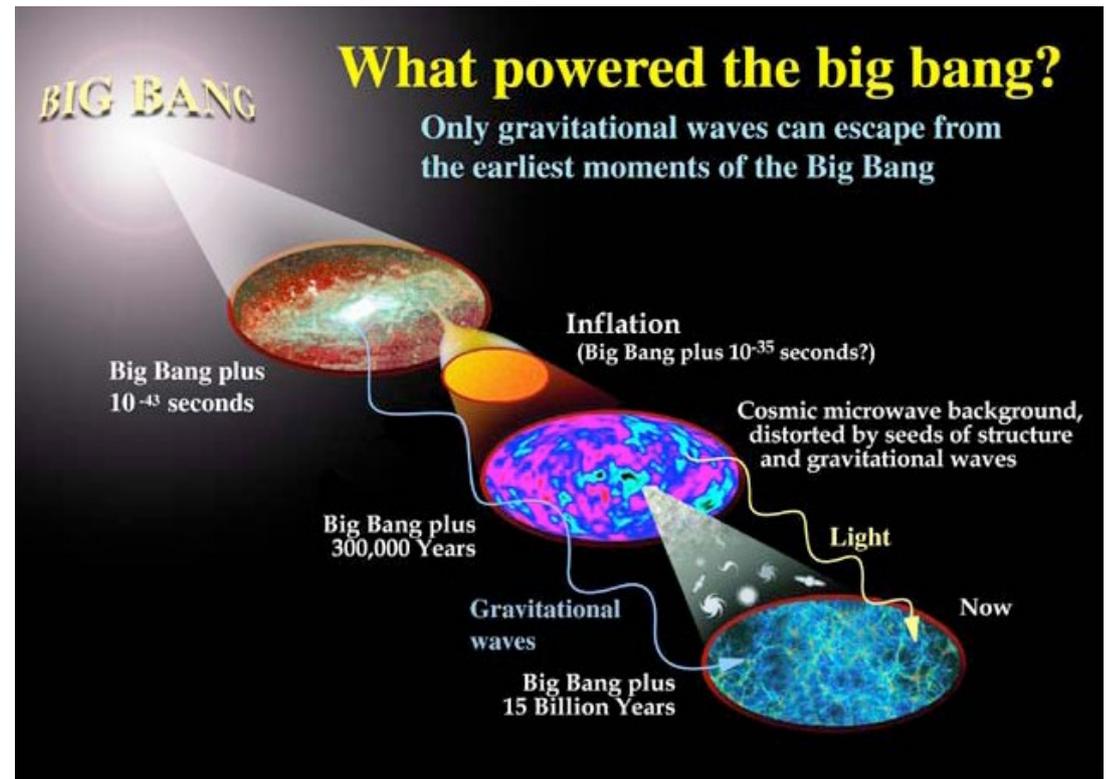
- Pulsars with mass non-uniformity:
  - distortion due to elastic stresses or magnetic field
  - distortion due to matter accretion
  - free precession around rotation axis
  - excitation of long-lasting oscillations (e.g. r-modes)
- Produce gravitational-waves, often at twice the rotational frequency.
- Waveform well-understood:
  - Sinusoidal but Doppler modulated
- Continuous source

Signal amplitude: 
$$h_0 \cong 10^{-27} \left( \frac{I_{zz}}{10^{38} \text{kg} \cdot \text{m}^2} \right) \left( \frac{10 \text{kpc}}{r} \right) \left( \frac{f}{100 \text{Hz}} \right)^2 \left( \frac{\epsilon}{10^{-6}} \right)$$

$\epsilon$  : ellipticity (adimensional number measuring the star's degree of asymmetry)  
 $f$  : signal frequency, proportional to star rotation frequency

# Stochastic background searches

- Incoherent superposition of many unresolved sources.
- Cosmological:
  - Inflationary epoch, preheating, reheating
  - Phase transitions
  - Cosmic strings
  - Alternative cosmologies
- Astrophysical:
  - Supernovae
  - Magnetars
  - Binary black holes
- Potentially could probe physics of the very-early Universe.

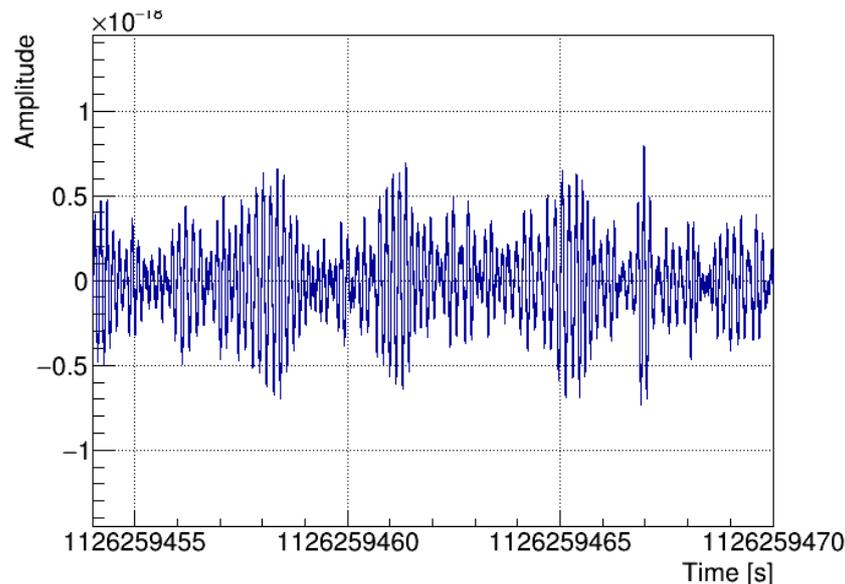


$$\Omega_{GW}(f) = \frac{f}{\rho_c} \frac{d\rho_{GW}}{df}$$

# Gravitational-wave data

GW detectors' readout system provides at any instant an estimate of strain: a quantity that is sensitive to arms' length difference:

- Digitized discrete time series:  $raw(t)$  (sampled at 16384 Hz or 20000 Hz) and synchronized with GPS clocks.
- Calibration of  $raw(t)$ : apply a frequency dependent factor [in reality this is a bit more complicated ...]



→  $h_{det}(t)$  time series that is detector noise plus all hypothetical GW signals

$$h_{det}(t) = n(t) + \mathbf{GW}(t)$$

# The detector's noise

Power spectral density:  
(PSD)  $\lim_{T \rightarrow \infty} \frac{1}{T} |\tilde{x}_T(f)|^2$

Power spectral density estimator for finite data set: Periodogram =  $\frac{1}{T} |\tilde{x}_T(f)|^2$

Improved estimator:

- average multiple periodograms ( $M$ ) to reduce the variance
- noise is non-stationary:  $T$  should not be too long (a few minutes)
- use windowed data to limit spectral leakage
- Welch approach: average of periodograms computed over overlapping windowed data segments

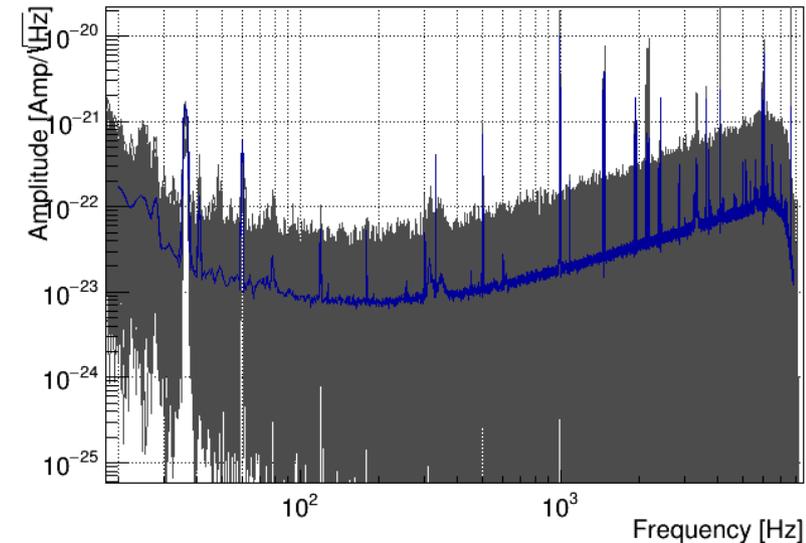
Sensitivity measured using the noise power spectral density :

$$S_n(k) = \text{Median}_{0 \leq m < M} \left\{ \frac{1}{Nf_s} \left| \sum_{j=0}^{N-1} x_m[j] w[j] e^{-2i\pi jk/N} \right|^2 \right\}$$

+ median-to-mean correction

One-sided / Two sided PSDs

Amplitude power spectral density:  $\sqrt{S_n(k)}$

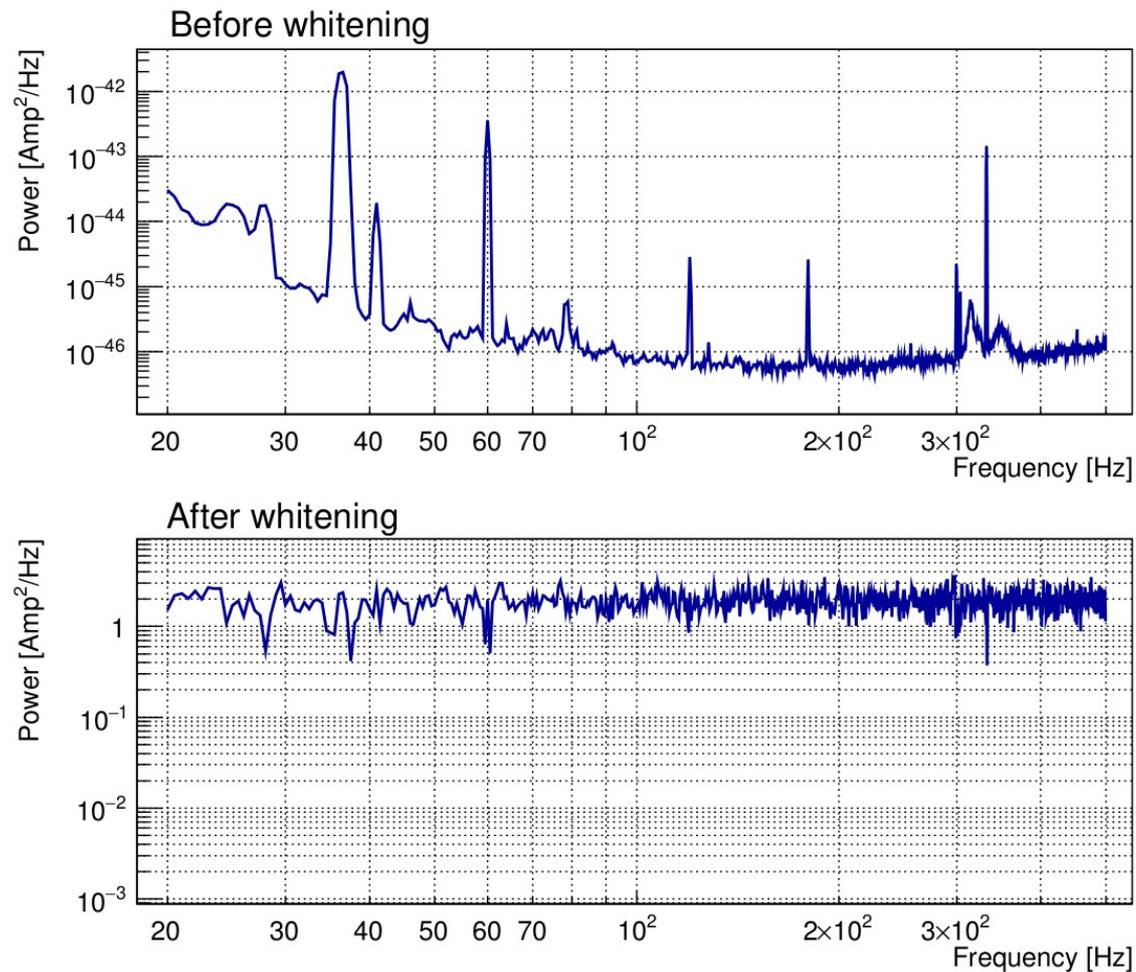


# Data whitening

GW data must be whitened. Several methods are used :

- reweighting of frequency bins
- linear prediction

→ white noise is mandatory for statistical interpretation of the data



# Stochastic background of gravitational waves

Assumption : stationary, unpolarized, and Gaussian stochastic background

→ Cross correlate the output of detector pairs to eliminate the noise

$$h_i = n_i + GW_i$$

$$\langle h_1, h_2 \rangle = \langle GW_1, GW_2 \rangle + \underbrace{\langle n_1, GW_2 \rangle}_0 + \underbrace{\langle GW_1, n_2 \rangle}_0 + \underbrace{\langle n_1, n_2 \rangle}_0$$

With  $\langle x_1, x_2 \rangle = \int_{-\infty}^{+\infty} \tilde{x}_1^*(f) \tilde{Q}(f) \tilde{x}_2(f) df$

Optimal filter:

$$\tilde{Q}(f) \propto \frac{\gamma(f) \Omega_{GW}(f)}{f^3 S_{n,1}(f) S_{n,2}(f)}$$

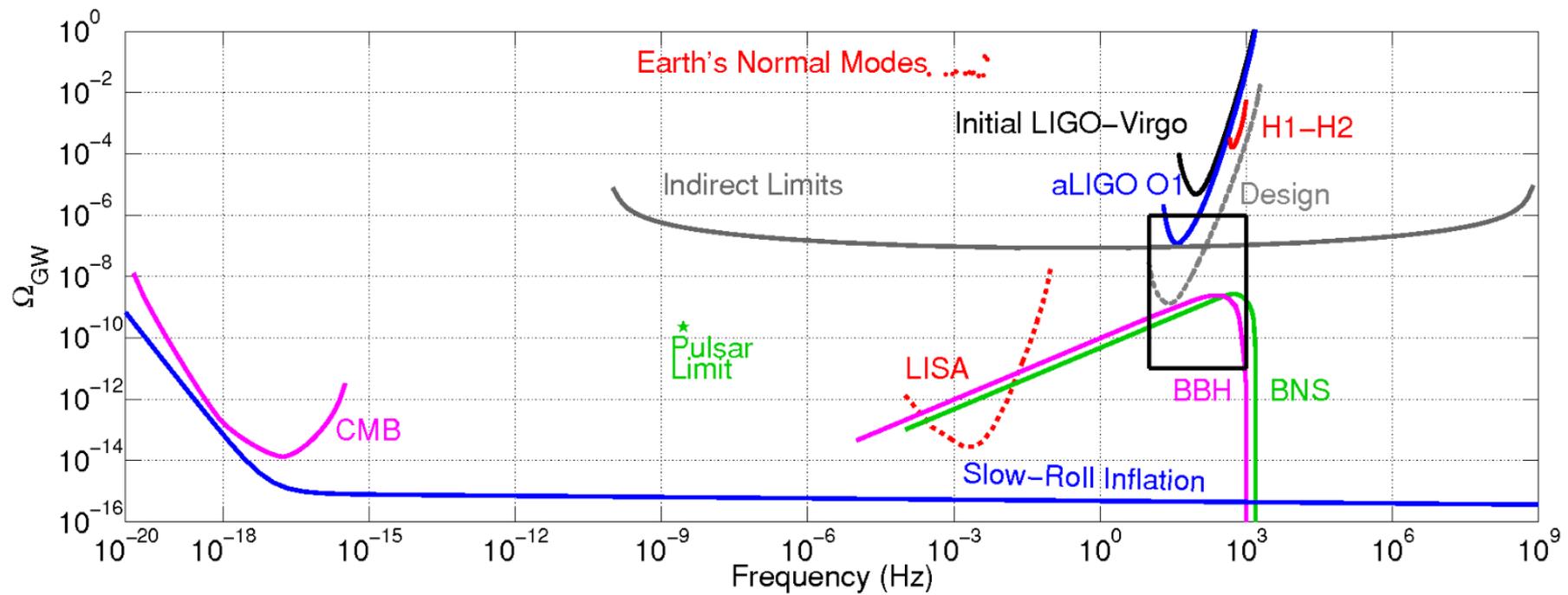
← overlap of antenna pattern
← GW spectrum

$\Omega_{GW}(f) = \Omega_\alpha f^\alpha$

↙ ↘
 Detector PSDs

O1 isotropic search, for  $\alpha = 0$  :  $\Omega_{GW}(25 \text{ Hz}) < 1.7 \times 10^{-7}$

# Stochastic background of gravitational waves



# Transient searches : excess power

Example : Q-transform

$$X(\tau, \phi, Q) = \int_{-\infty}^{+\infty} h_{det}(t) w(t - \tau, \phi, Q) e^{-2i\pi\phi\tau} dt$$

→ window width  $\sim 1/\phi$

~ short Fourier transform with a Gaussian window

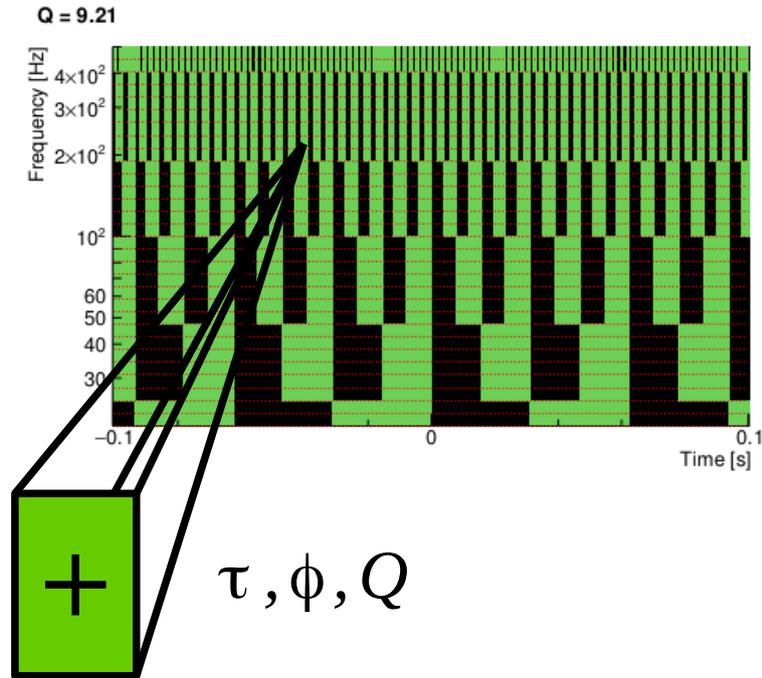
→ Goal : cover a parameter space as large as possible

$$\text{Noise only: } \langle |N(\tau, \phi, Q)|^2 \rangle = \int_{-\infty}^{+\infty} |\tilde{w}(\phi - f, \phi, Q)|^2 S_n(f) df$$

$$\text{Whitened noise + window normalization: } \langle |N^w(\tau, \phi, Q)|^2 \rangle = 1$$

→ Signal-to-noise ratio estimator

$$\hat{\rho}^2(\tau, \phi, Q) = |X^w(\tau, \phi, Q)|^2 - \langle |N^w(\tau, \phi, Q)|^2 \rangle = |X^w(\tau, \phi, Q)|^2 - 1$$



Statistical interpretation: noise is Gaussian-distributed with unit variance

# GW150914

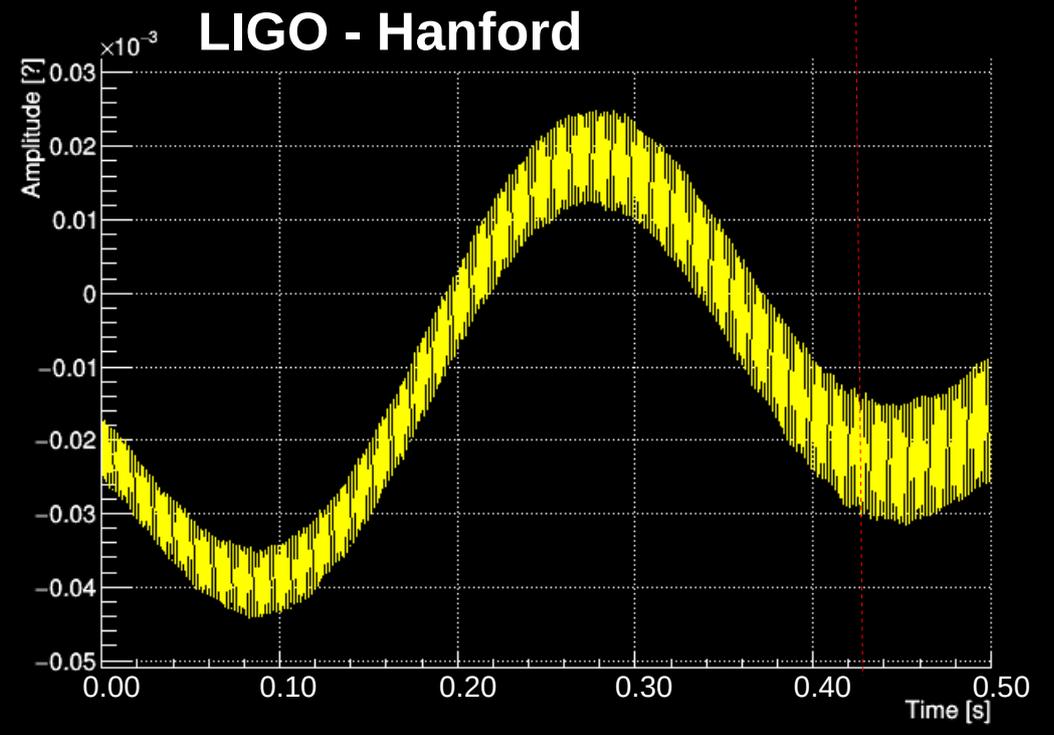
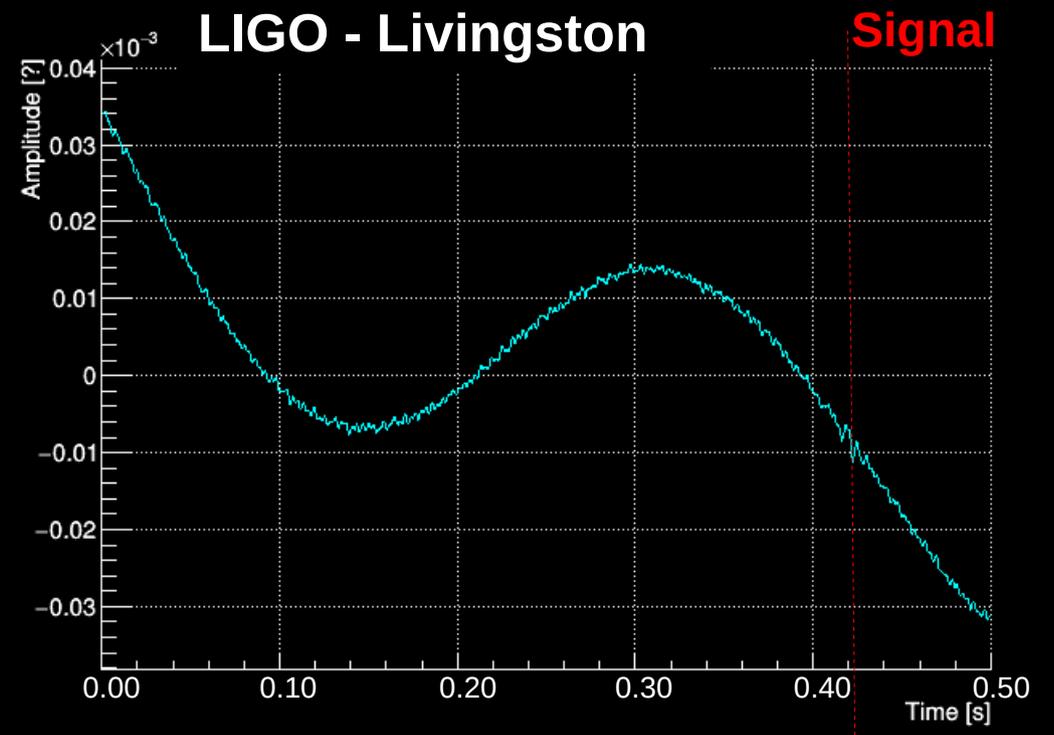
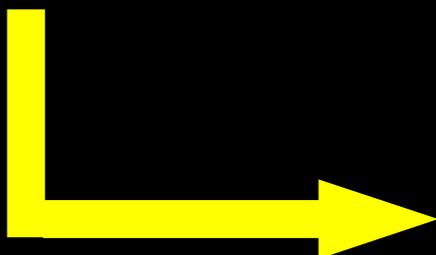
Output power



Livingston

Hanford

Output power

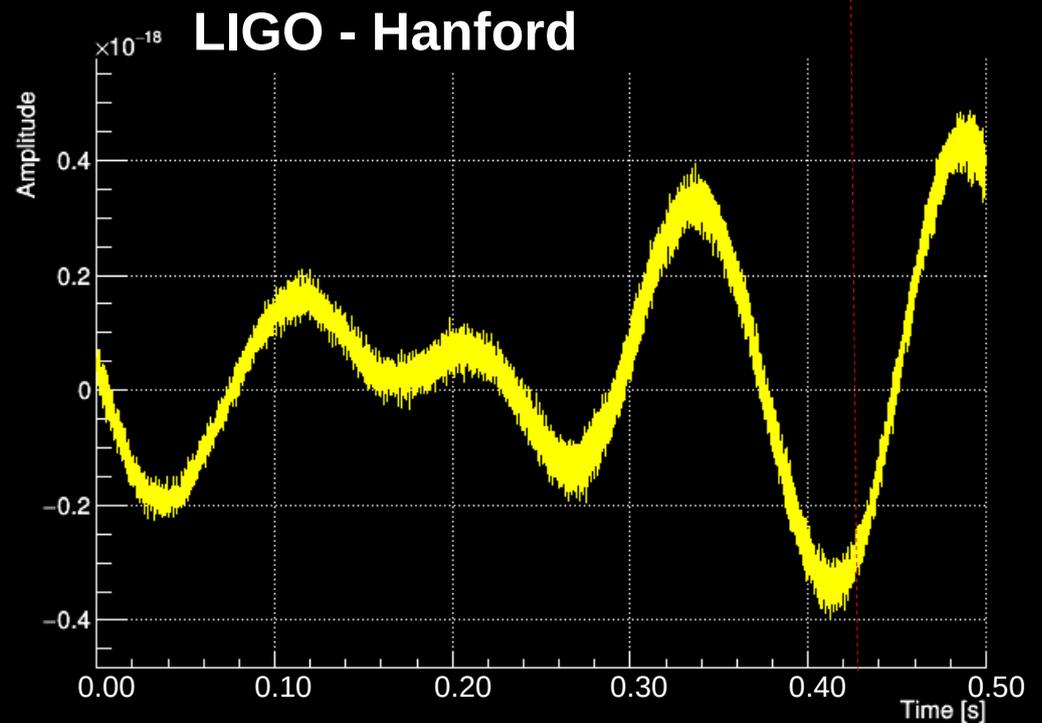
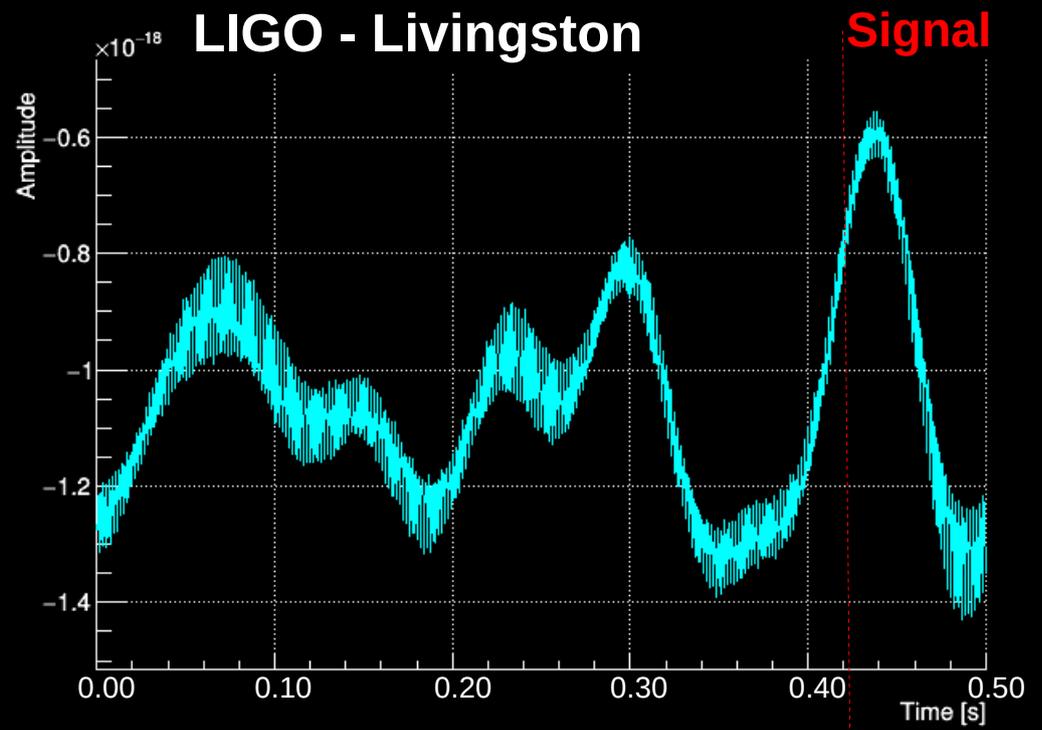


# GW150914

$$h(t)$$

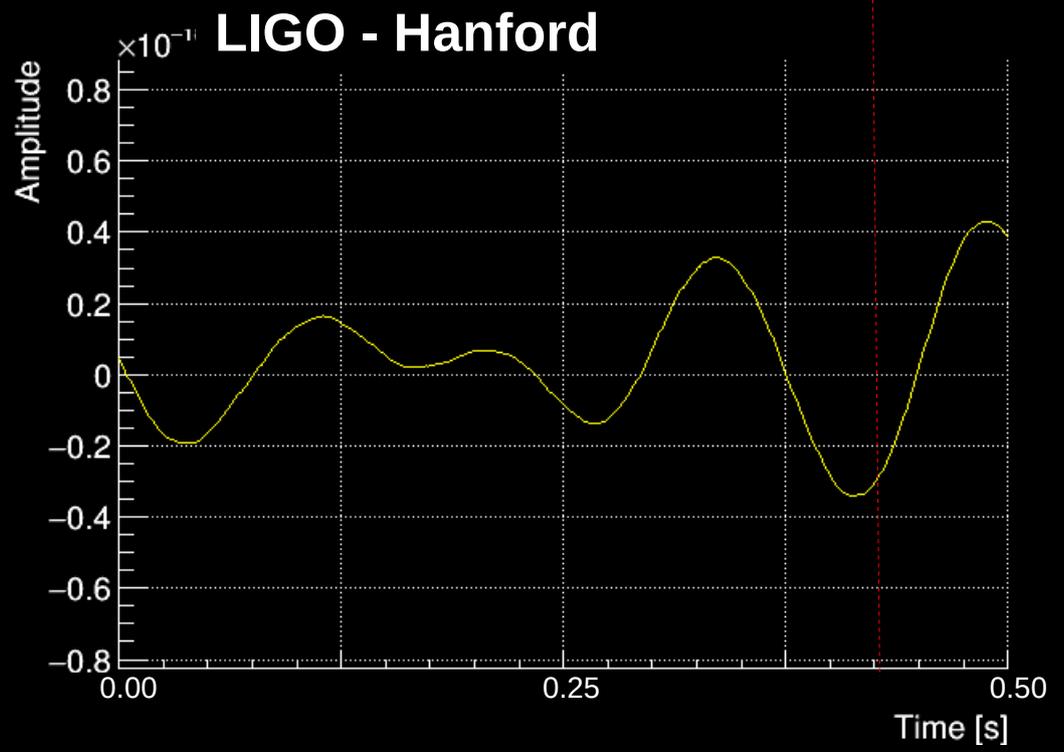
Data is calibrated

→ GW strain amplitude  $h(t)$   
(including high-pass filter  $f > 10$  Hz)



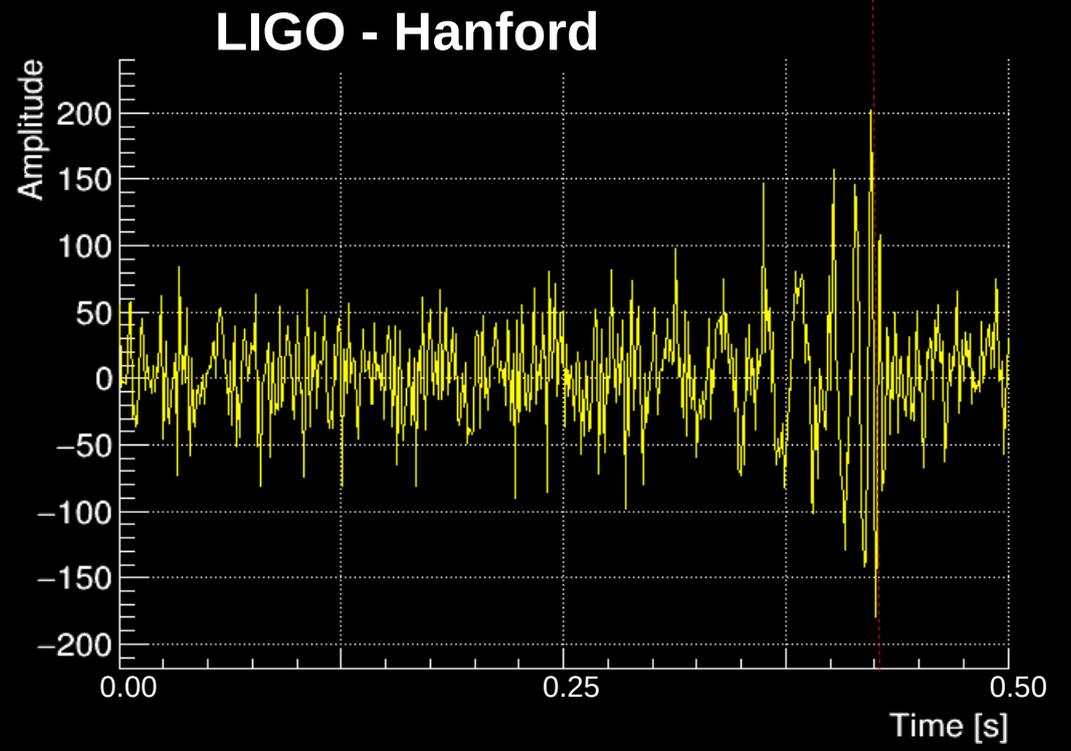
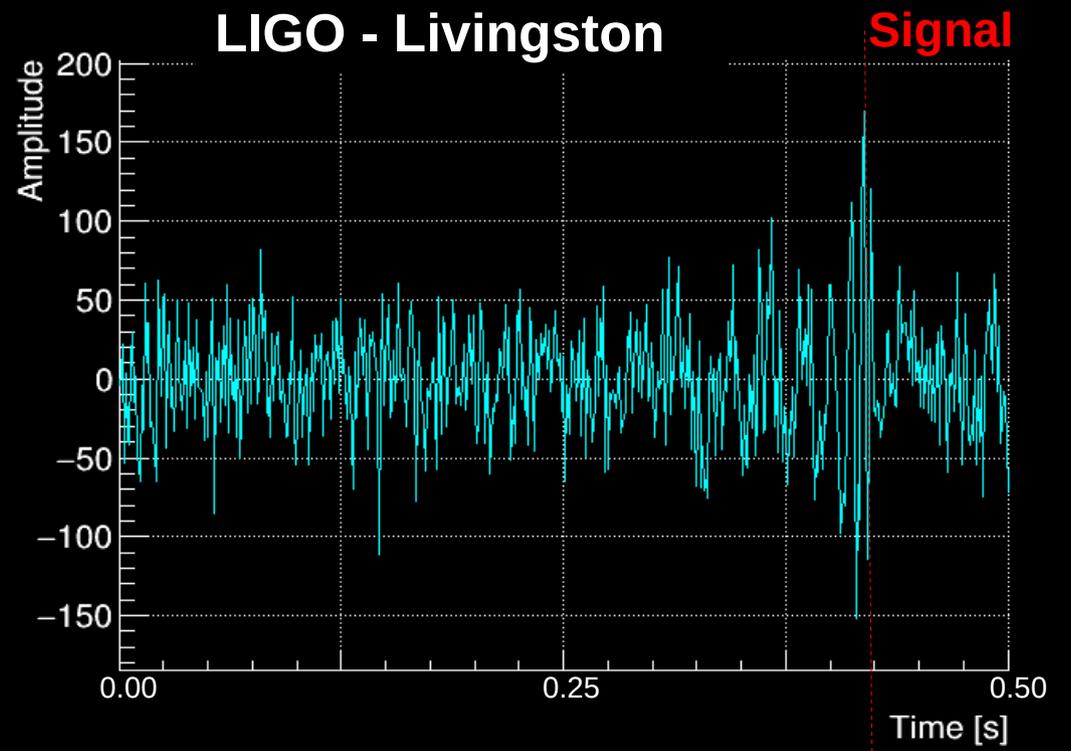
# GW150914

Data are low-pass filtered  
(here, < 500 Hz)



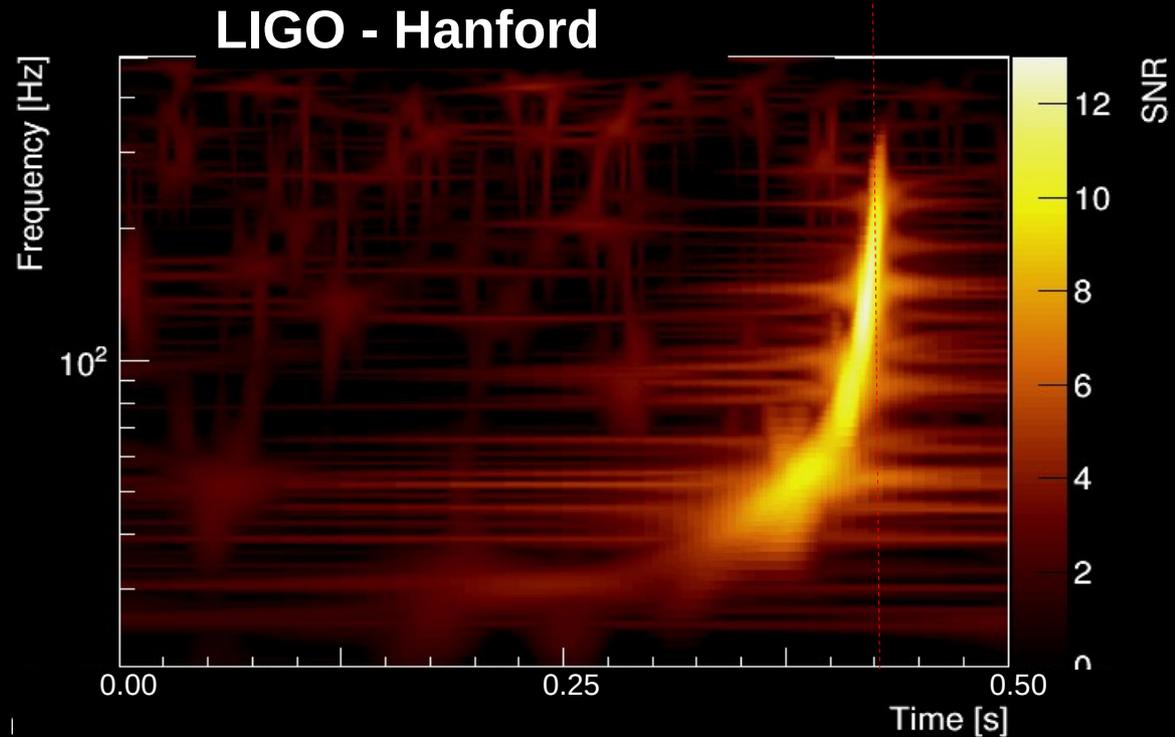
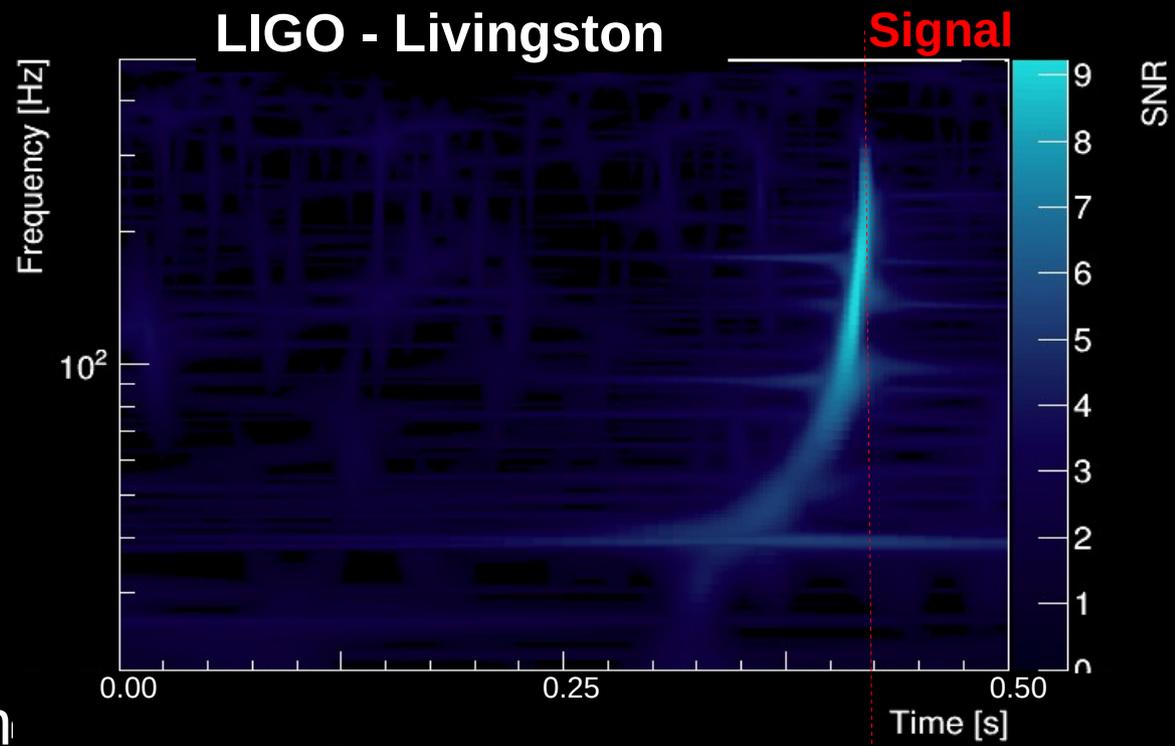
# GW150914

Data are whitened



# GW150914

Time-frequency decomposition  
(Short Fourier transforms)



## Known waveform → matched-filtering technique

Simplest linear filter: correlation  $C(t) = \int_{-\infty}^{+\infty} h_{det}(t')k(t-t')dt = \int_{-\infty}^{+\infty} \tilde{h}_{det}(f)\tilde{k}^*(f)e^{2i\pi ft}df$

$k(t)$  is the impulse response function of the filter :  $h_{det}(t) = \delta(t) \Rightarrow C(t) = k(t)$

**Matched-filter:** optimal filter maximizing the SNR in presence of additive noise

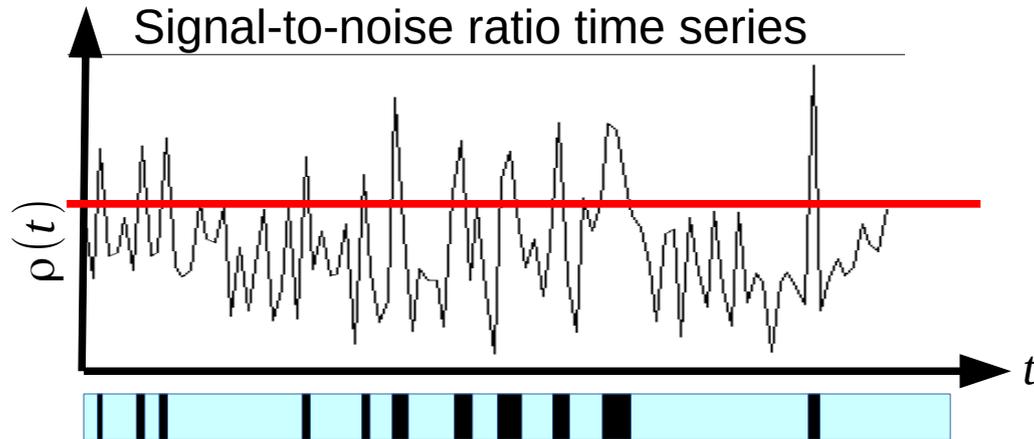
$$h_{det}(t) = n(t) + \mathbf{GW}(t)$$

$$\rho(t) = \frac{C(t)}{\sqrt{\langle N^2(t) \rangle}} \quad \text{with} \quad \langle N^2(t) \rangle = \int_{-\infty}^{+\infty} |\tilde{k}(f)|^2 S_n(f) df$$

The SNR is maximized if  $\tilde{k}(f) \propto \frac{\mathbf{G}\tilde{W}^*(f)}{S_n(f)}$

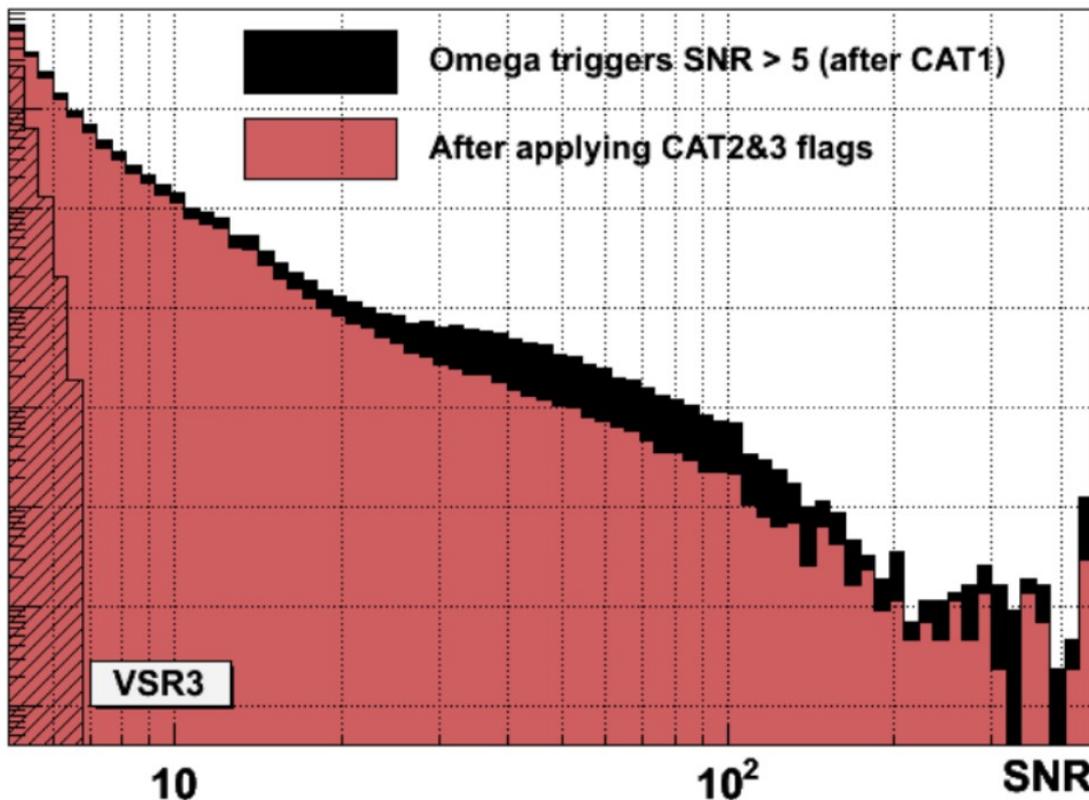
$$\rho(t) = \int_{-\infty}^{+\infty} \frac{\mathbf{G}\tilde{W}^*(f)\tilde{h}_{det}(f)}{S_n(f)} e^{2i\pi ft} df$$

# Transient searches : triggers



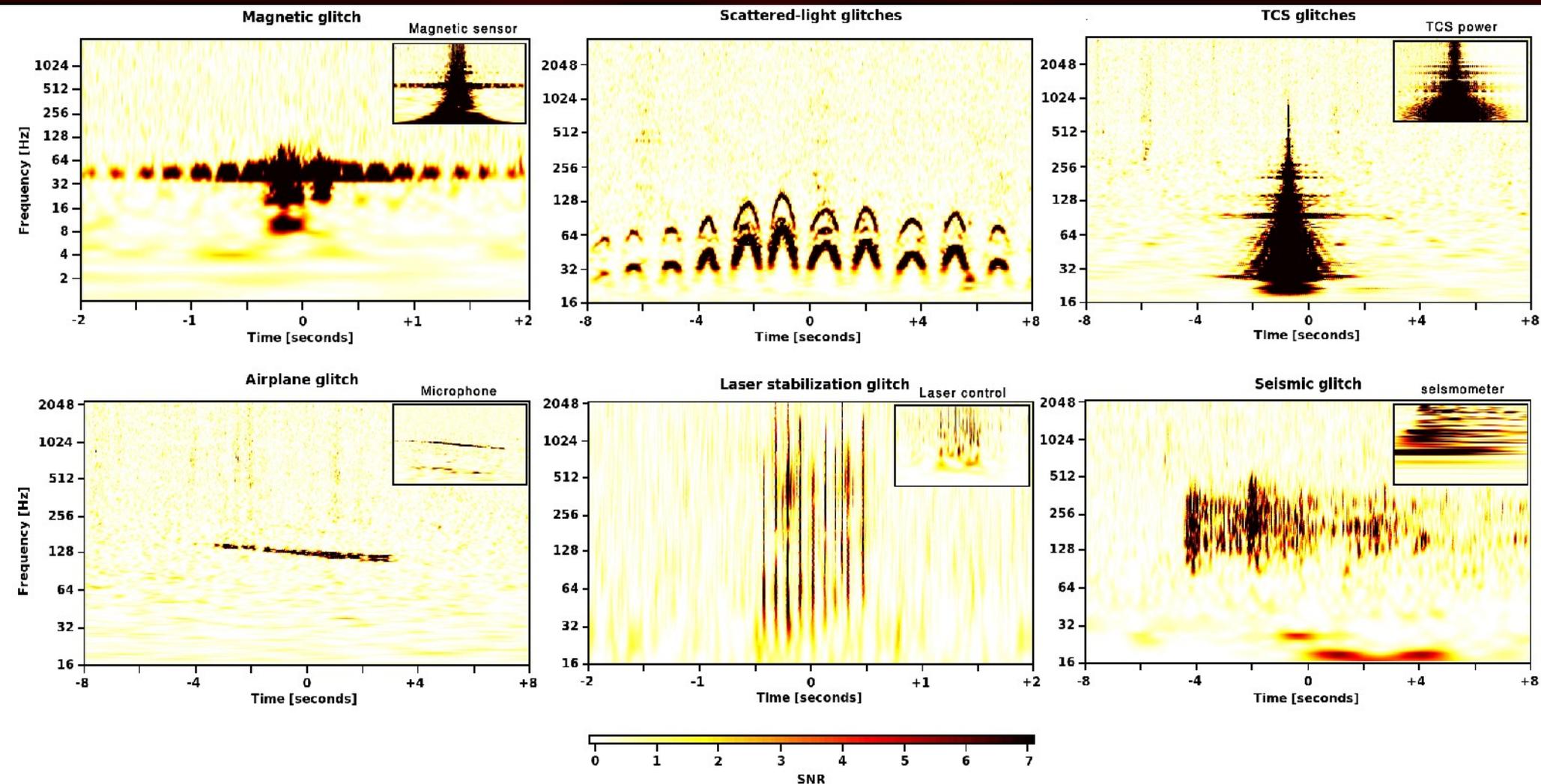
- A list of events is produced:
- start/end/peak times
  - SNR
  - template parameters (masses, spins)

Now, the challenge is to reject noise events to better isolate true signals

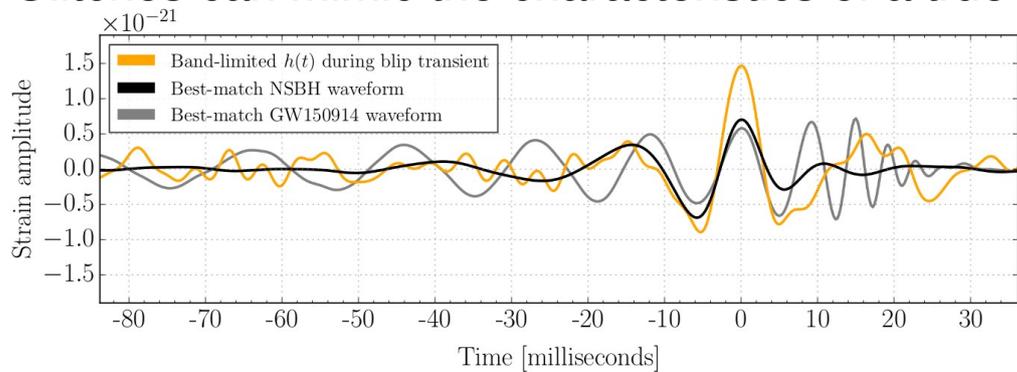


The noise distribution is highly non-Gaussian !

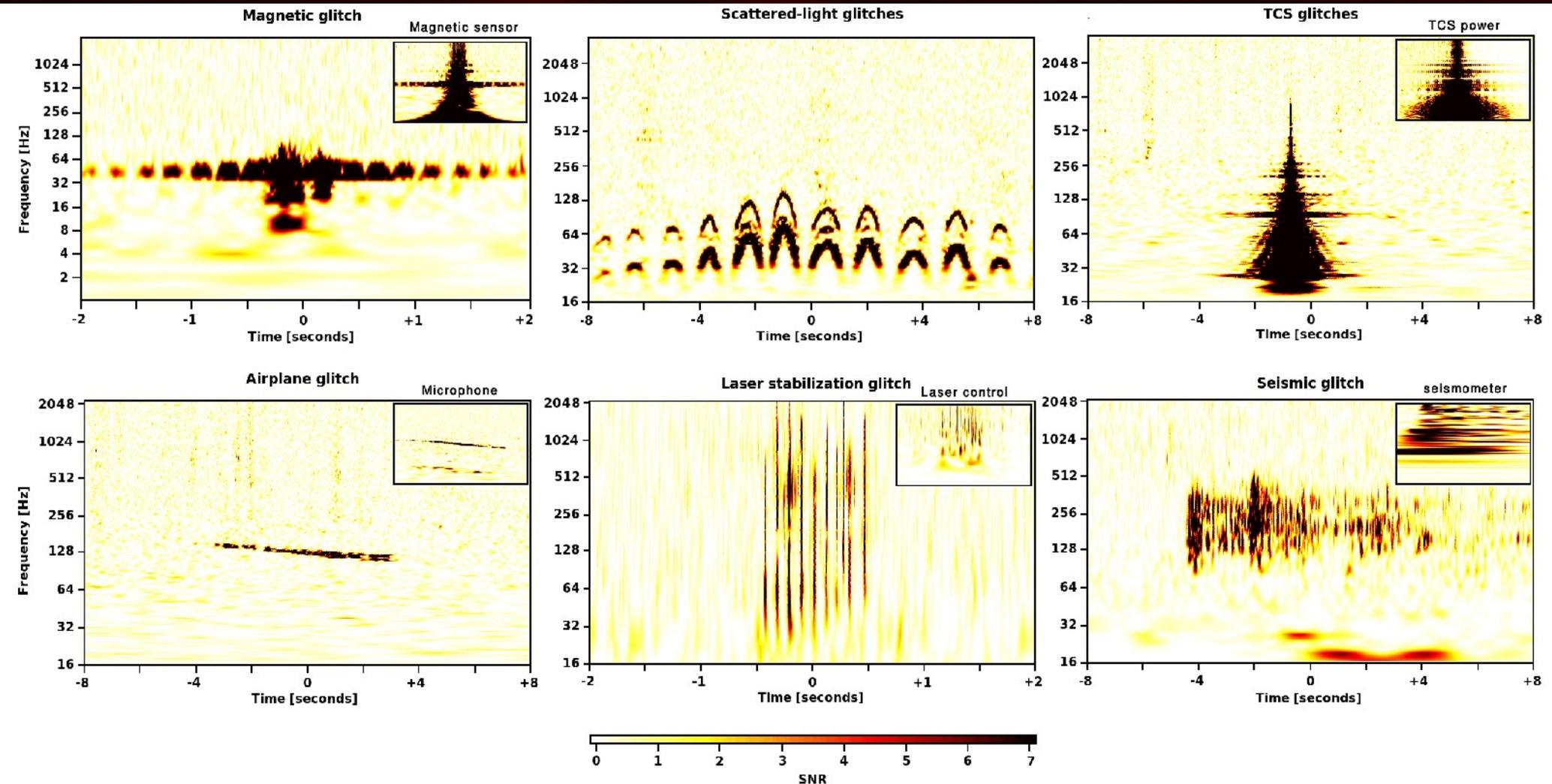
# Transient searches : “glitches”



Glitches can mimic the characteristics of a true signal



# Transient searches : “glitches”

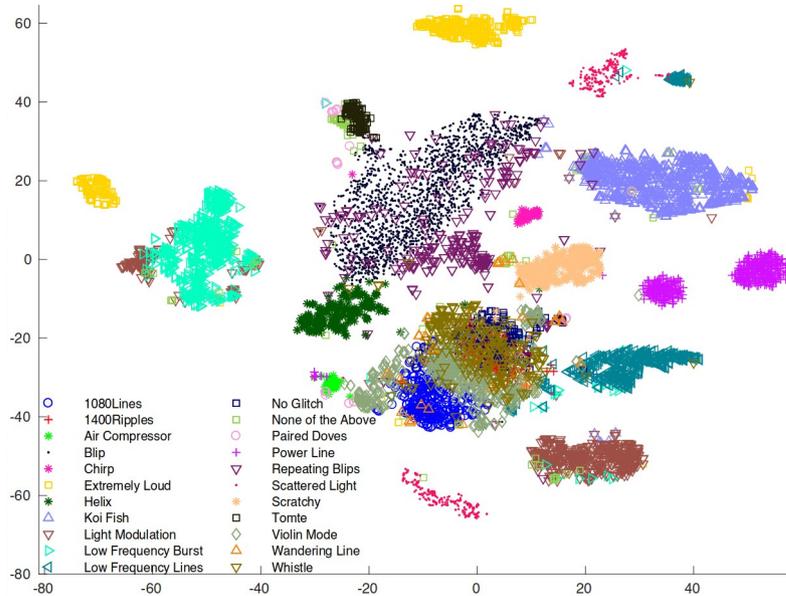


Thousands of auxiliary channels are used to monitor the instruments and witness glitches

- environmental sensors
- detector sub-systems
- detector control

→ **Data quality flags to veto glitches**

# Transient searches : glitch classification



- t-distributed stochastic neighbor embedding (t-SNE) for reducing high-dimensional data into a space of two or three dimensions
- Identifies clustering of glitches based on morphology, and similarity between differing glitch types

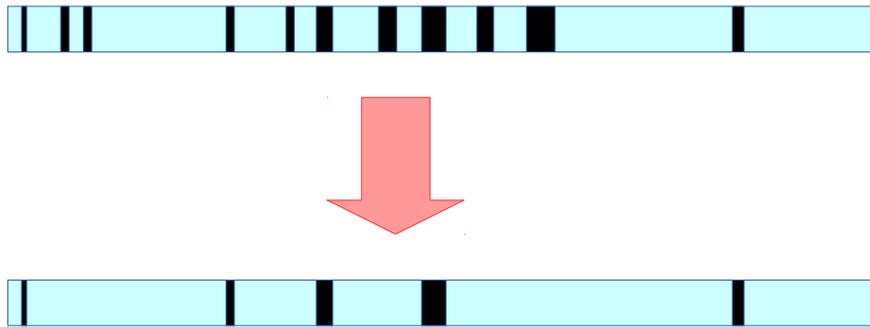
Citizen project "GravitySpy"

<https://www.zooniverse.org/projects/zooniverse/gravity-spy>

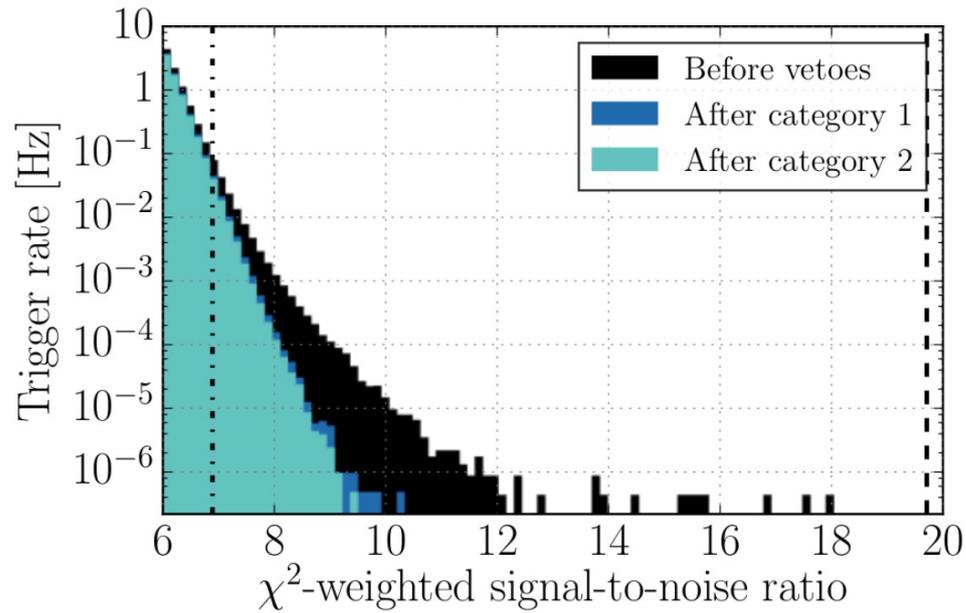
Ongoing efforts of MLA to classify and reject glitches

→ See Agata Trovato's talk

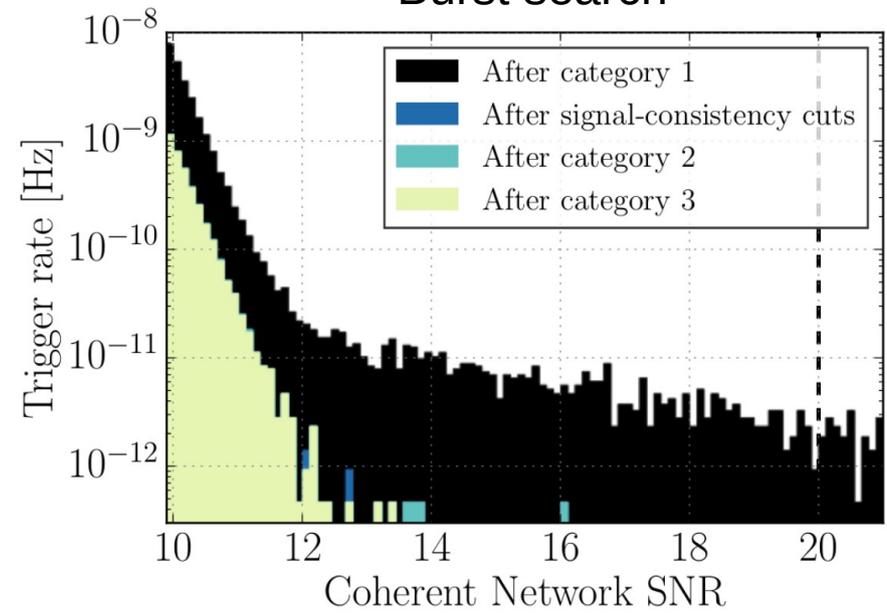
# Transient searches : vetoes



### CBC search



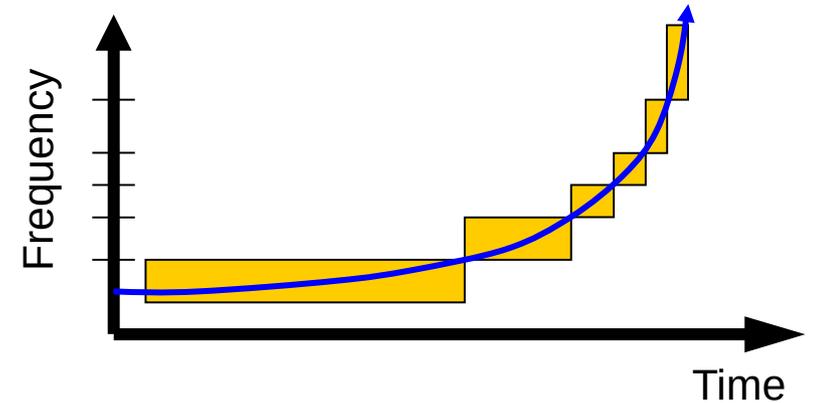
### Burst search



# CBC searches : waveform consistency test

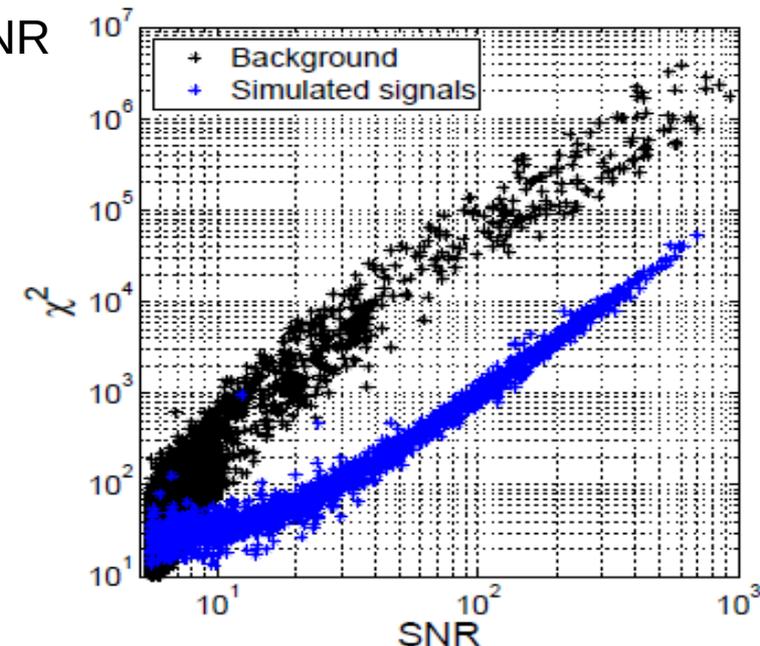
- Divide the “selected” template into  $p$  parts
- The frequency intervals are chosen so that for a true signal, the SNR is uniformly shared among the frequency bands.

$$\chi^2(t) = p \sum_{j=1}^p \left| \rho_j - \frac{\rho}{p} \right|^2$$



- For a stationary and Gaussian noise  $\chi^2$  has an expectation value:  $\langle \chi^2 \rangle = p - 1$
- In practice  $\chi^2$  values are larger than expected for large SNR (discrete template banks effect)  $\rightarrow$  cut in  $(\text{SNR}, \chi^2)$  plane
- Weighted SNR

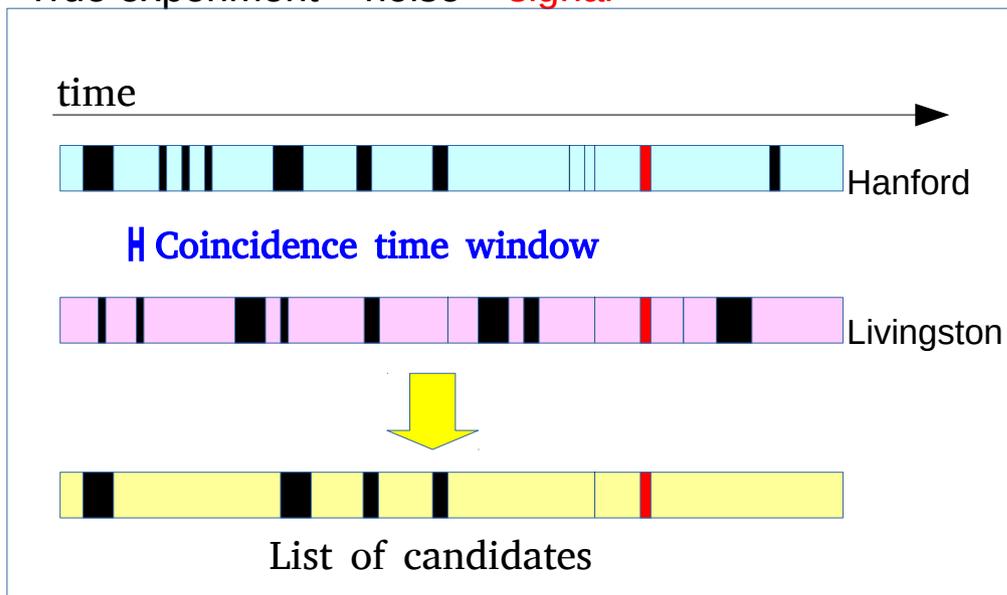
$$\rho_{\text{new}} = \begin{cases} \rho, & \chi^2 \leq n_{\text{dof}} \\ \frac{\rho}{\left[ \left( 1 + \frac{\chi^2}{n_{\text{dof}}} \right)^{4/3} / 2 \right]^{1/4}}, & \chi^2 > n_{\text{dof}} \end{cases}$$



# Coincidence between detectors

A gravitational-wave signal is detected by multiple detectors almost simultaneously

True experiment = noise + **signal**



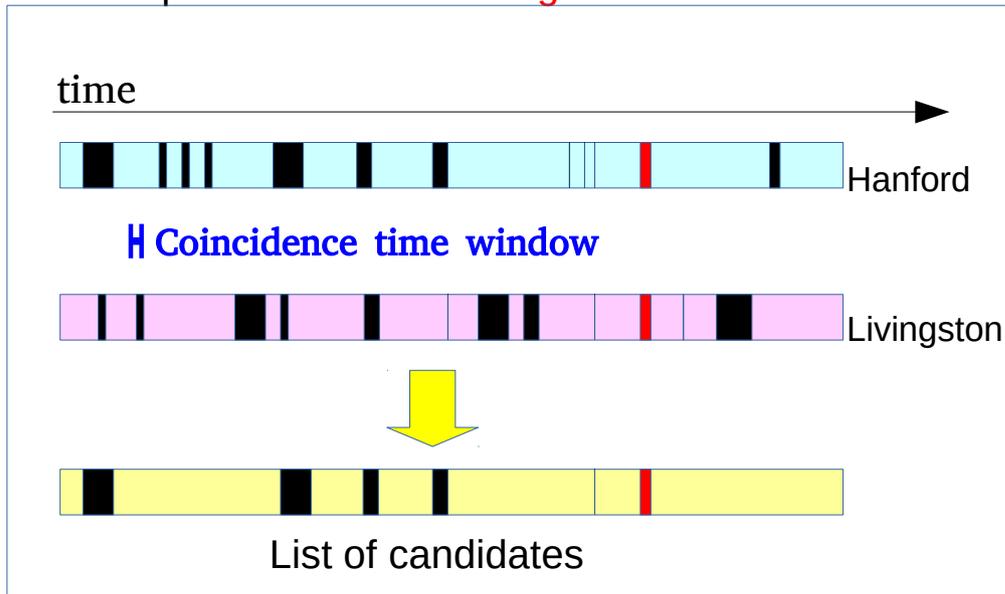
Coincidence rate:

$$R_{\text{coinc}} \sim R_H R_L \Delta t_{\text{win}}$$
$$\sim (1 \text{ Hz}) \times (1 \text{ Hz}) \times (10^{-2} \text{ s}) = 10^{-2} \text{ Hz}$$

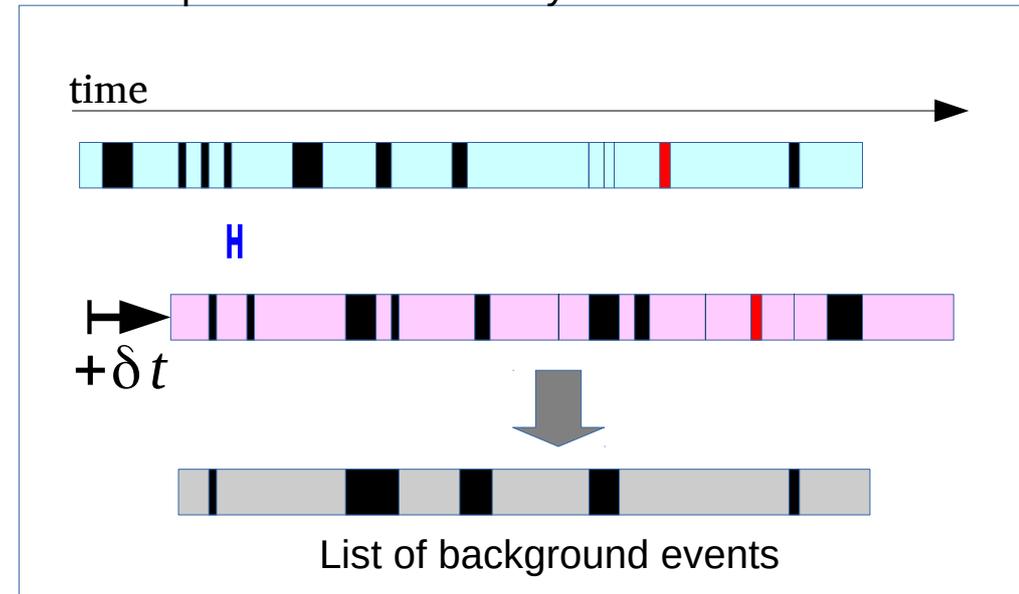
# Background estimation

The background of a gravitational-wave search is estimated using the time-slide technique  
Assumption = uncorrelated noise between detectors

True experiment = noise + **signal**



Fake experiment = noise only



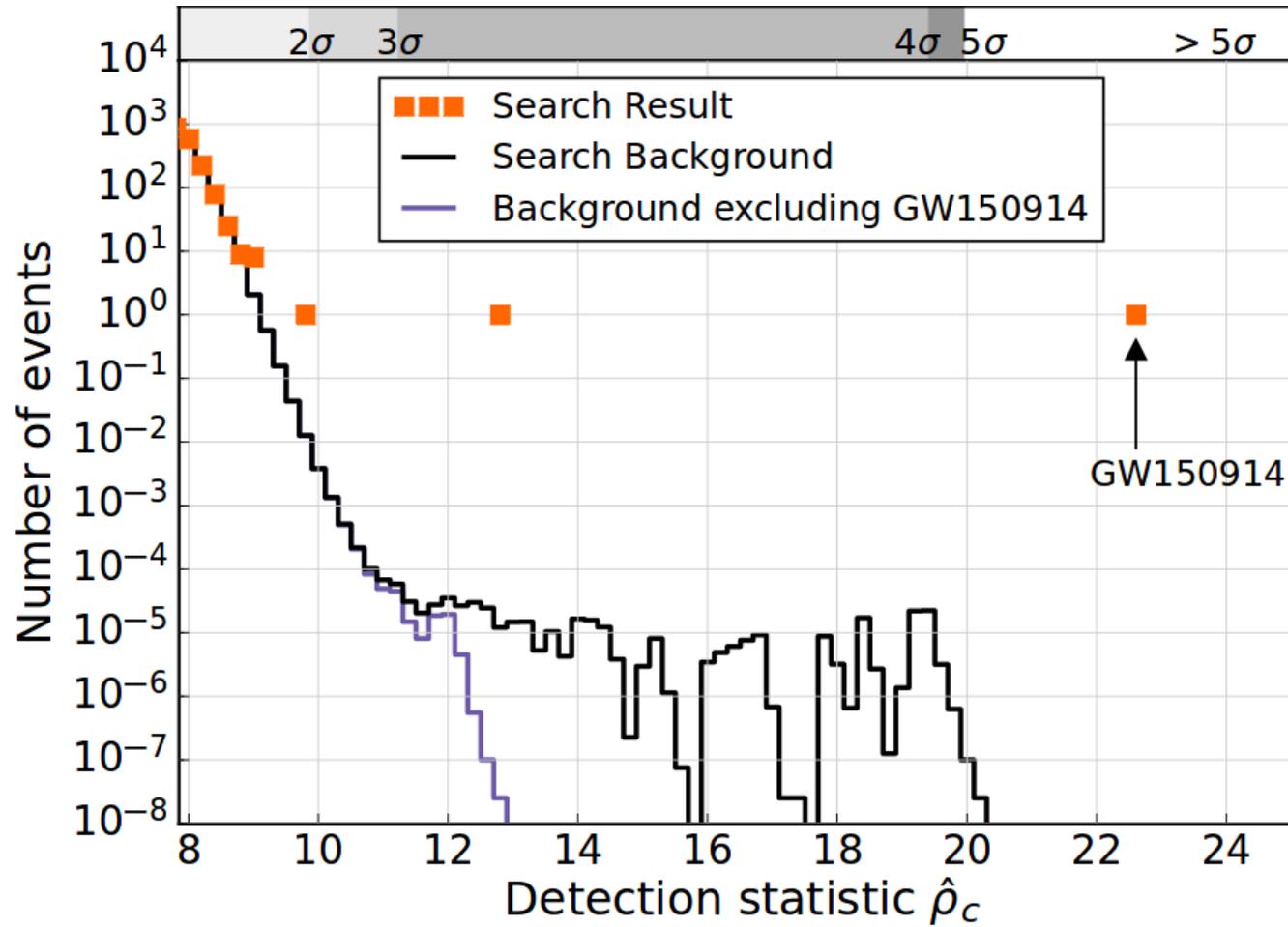
A very large number of fake experiments can be simulated using multiple offsets

**LIGO O1 analysis:**

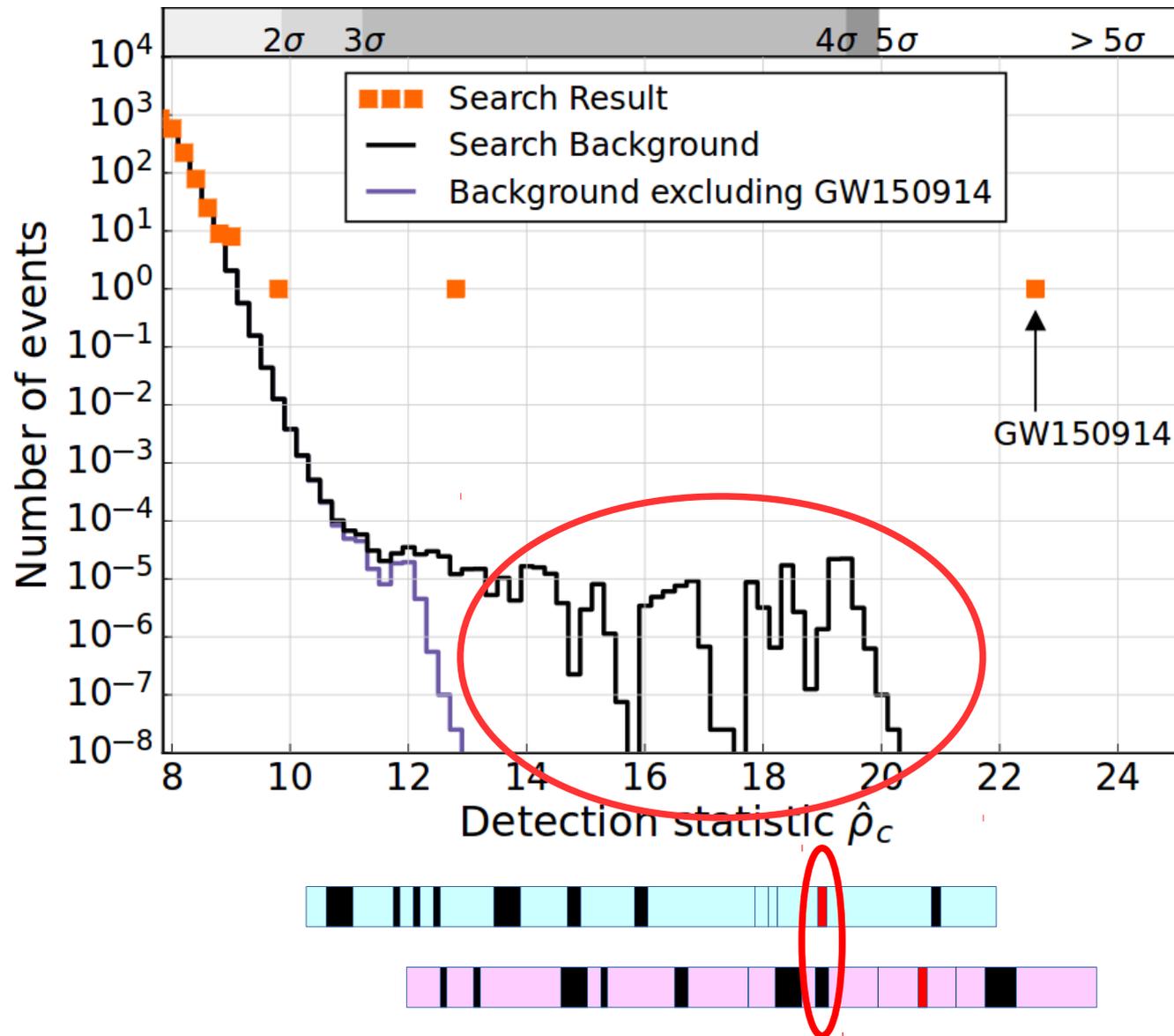
–  $O(10^6)$  time offsets

→ **background estimated using a fake experiment of  $O(100,000)$  years**

# Event significance



# Event significance



## GraceDB — Gravitational Wave Candidate Event Database

HOME	SEARCH	CREATE	REPORTS	RSS	LATEST	OPTIONS	DOCUMENTATION	AUTHENTICATED AS: FLORENT ROBINET		
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### Basic Info

UID	Labels	Group	Pipeline	Search	Instruments	UTC Event Time	FAR (Hz)	Links	UTC Submitted
G211117	H1OK L1OK ADVOK EM_READY	CBC	gstlal	HighMass	H1,L1	2015-12-26 03:38:53 UTC	3.333e-11	<a href="#">Data</a>	2015-12-26 03:40:00 UTC

### Coinc Tables

End Time (GPS)	1135136350.6478 s
Total Mass	26.3501 $M_{\odot}$
Chirp Mass	9.5548 $M_{\odot}$
SNR	11.7103
False Alarm Probability	1.120e-04
Log Likelihood Ratio	22.5996

### Single Inspiral Tables

IFO	L1	H1
Channel	GDS-CALIB_STRAIN	GDS-CALIB_STRAIN
End Time (GPS)	1135136350.646883043 s	1135136350.647757924 s
Template Duration	2.25322770554 s	2.25322770554 s
Effective Distance	472.93436 Mpc	461.88879 Mpc
COA Phase	2.7356486 rad	0.13969257 rad
Mass 1	19.924686 $M_{\odot}$	19.924686 $M_{\odot}$
Mass 2	6.4254546 $M_{\odot}$	6.4254546 $M_{\odot}$
$\eta$	0.18438664	0.18438664
F Final	1024.0 Hz	1024.0 Hz
SNR	7.3947201	9.0802174
$\chi^2$	1.0857431	1.0069774
$\chi^2$ DOF	1	1
spin1z	0.33962944	0.33962944
spin2z	-0.1238557	-0.1238557

### Neighbors [-5,+5]

UID	Labels	Group	Pipeline	Search	Instruments	UTC Event Time	$\Delta$ gptime	FAR (Hz)	Links	UTC Submitted
<a href="#">G211182</a>		Burst	CWB2G	AllSky	H1,L1	2015-12-26 03:38:53 UTC	-0.018658		<a href="#">Data</a>	2015-12-26 09:44:37 UTC
<a href="#">G211115</a>		CBC	gstlal	HighMass	H1,L1	2015-12-26 03:38:53 UTC	-0.007229	1.032e-09	<a href="#">Data</a>	2015-12-26 03:39:59 UTC
<a href="#">G211118</a>		CBC	gstlal	HighMass	H1,L1	2015-12-26 03:38:53 UTC	-0.000043	3.279e-08	<a href="#">Data</a>	2015-12-26 03:40:00 UTC
<a href="#">G216856</a>		CBC	gstlal	HighMass	H1,L1	2015-12-26 03:38:53 UTC	0.000278	1.187e-12	<a href="#">Data</a>	2016-01-15 14:31:22 UTC
<a href="#">G211116</a>		CBC	gstlal	HighMass	H1,L1	2015-12-26 03:38:53 UTC	0.000780	4.507e-09	<a href="#">Data</a>	2015-12-26 03:40:00 UTC

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### Coinc Tables

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Chirp Mass	9.5548 $M_{\odot}$
SNR	11.7103
False Alarm Probability	1.120e-04
Log Likelihood Ratio	22.5996

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$\eta$	0.18438664	0.18438664
F Final	1024.0 Hz	1024.0 Hz
SNR	7.3947201	9.0802174
$\chi^2$	1.0857431	1.0069774
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spin1z	0.33962944	0.33962944
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Low-latency detection

### Neighbors [-5,+5]

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<a href="#">G211115</a>		CBC	gstlal	HighMass	H1,L1	2015-12-26 03:38:53 UTC	-0.007229	1.032e-09	<a href="#">Data</a>	2015-12-26 03:39:59 UTC
<a href="#">G211118</a>		CBC	gstlal	HighMass	H1,L1	2015-12-26 03:38:53 UTC	-0.000043	3.279e-08	<a href="#">Data</a>	2015-12-26 03:40:00 UTC
<a href="#">G216856</a>		CBC	gstlal	HighMass	H1,L1	2015-12-26 03:38:53 UTC	0.000278	1.187e-12	<a href="#">Data</a>	2016-01-15 14:31:22 UTC
<a href="#">G211116</a>		CBC	gstlal	HighMass	H1,L1	2015-12-26 03:38:53 UTC	0.000780	4.507e-09	<a href="#">Data</a>	2015-12-26 03:40:00 UTC

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<a href="#">G211118</a>		CBC	gstlal	HighMass	H1,L1	2015-12-26 03:38:53 UTC	-0.000043	3.279e-08	<a href="#">Data</a>	2015-12-26 03:40:00 UTC
<a href="#">G216856</a>		CBC	gstlal	HighMass	H1,L1	2015-12-26 03:38:53 UTC	0.000278	1.187e-12	<a href="#">Data</a>	2016-01-15 14:31:22 UTC
<a href="#">G211116</a>		CBC	gstlal	HighMass	H1,L1	2015-12-26 03:38:53 UTC	0.000780	4.507e-09	<a href="#">Data</a>	2015-12-26 03:40:00 UTC

Multiple triggers

# Parameter estimation

Full analysis of the data surrounding the event

- only input from searches: time of the event
- fully explore the parameter space
- fully coherent search
- include calibration uncertainty

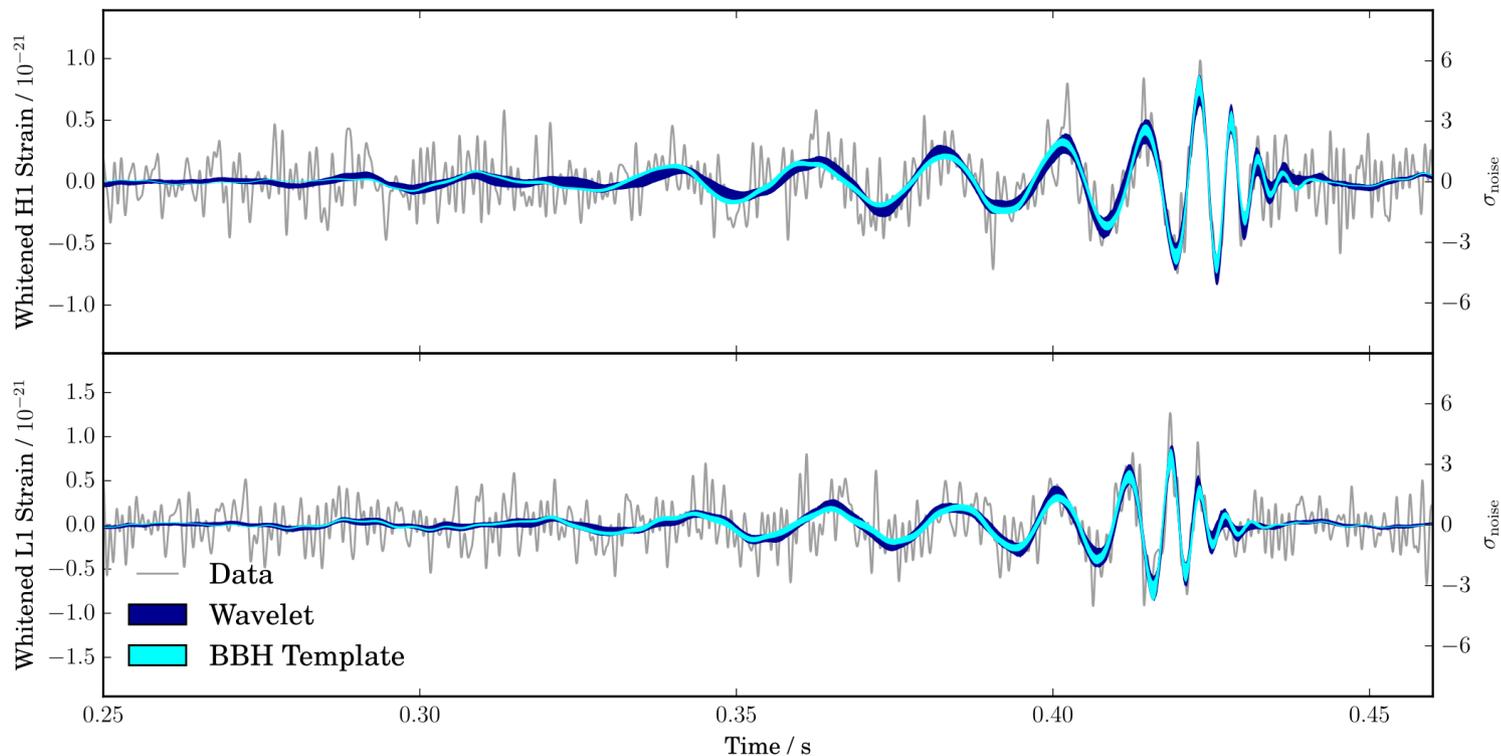
8 intrinsic parameters (masses and spins)

9 extrinsic parameters (distance, position, orientation, coalescence time and phase)

Orbital ellipticity is neglected

Dimensionless spin:  $a = \frac{c |\vec{S}|}{Gm^2} \leq 1$

Frequency is redshifted → masses must be rescaled by a factor  $(1+z)$



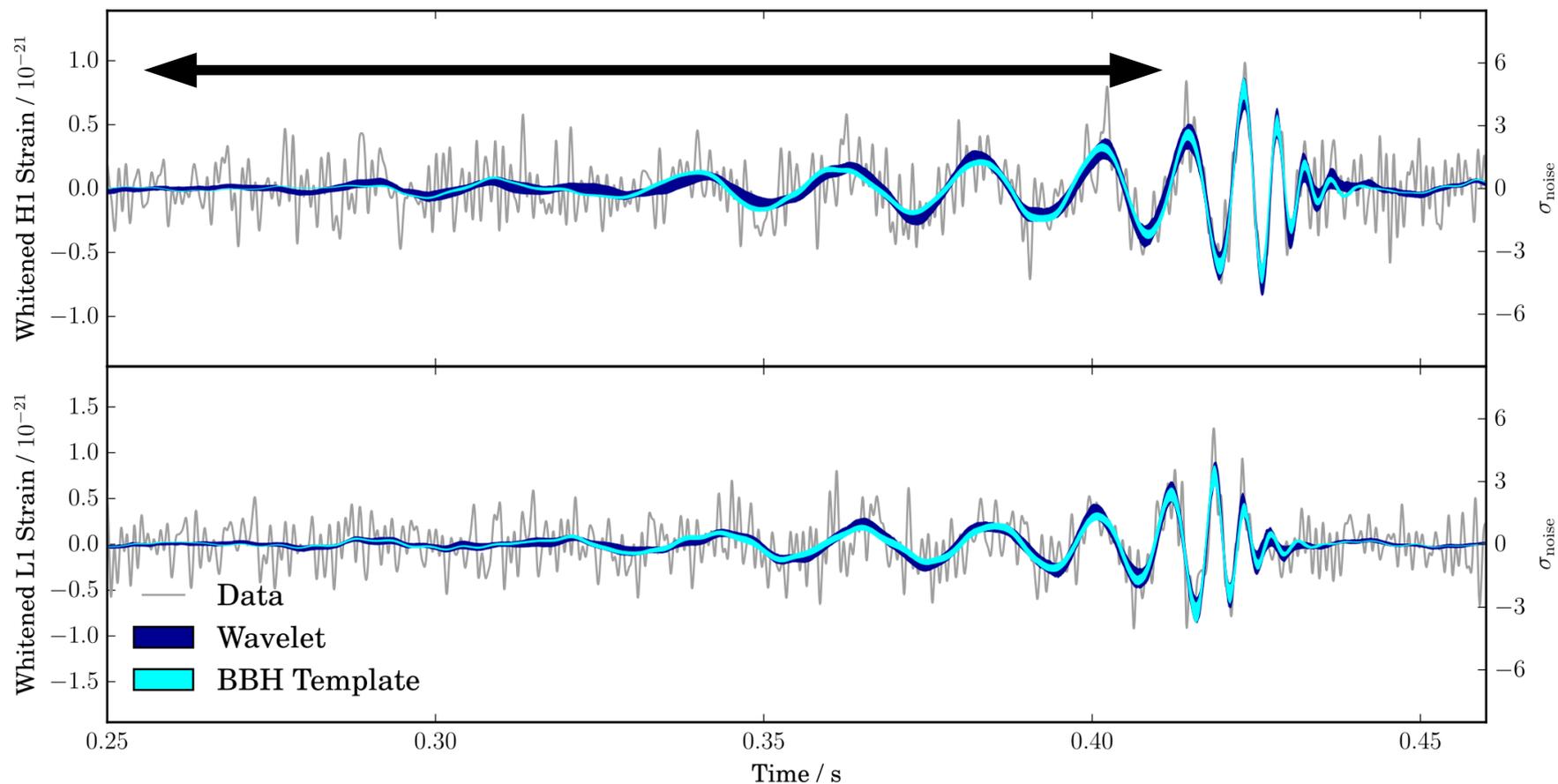
# Parameter estimation

Inspiral phase: PN perturbative expansion ( $v/c$ )

Leading order  $\rightarrow$  phase evolution driven by the chirp mass  
(tight constraints)

Next order  $\rightarrow$   $m_2/m_1$  and spins  $\parallel \mathbf{L}$

Next orders  $\rightarrow$  full spins

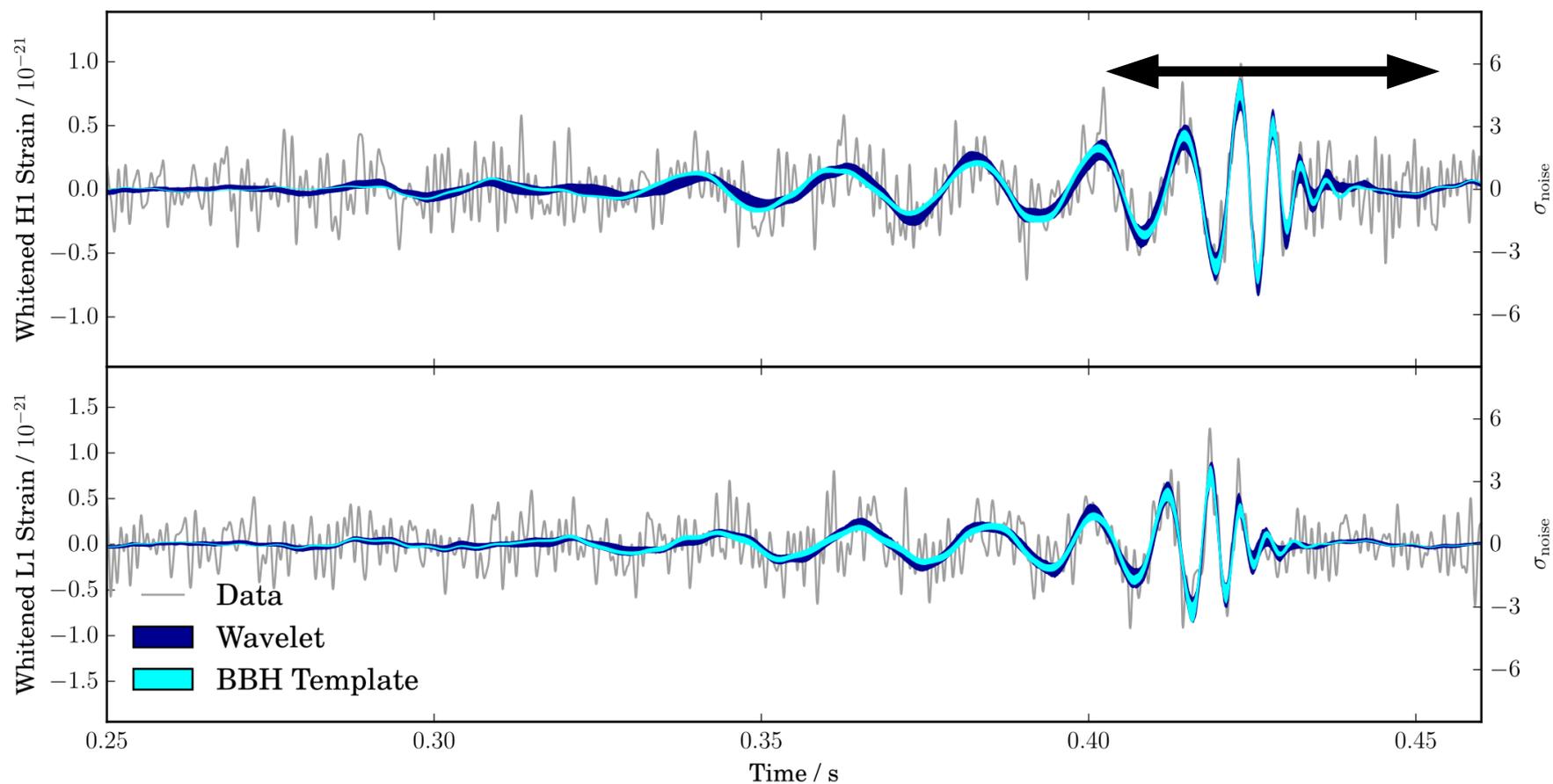


# Parameter estimation

Late inspiral – merger – ringdown: numerical relativity waveforms

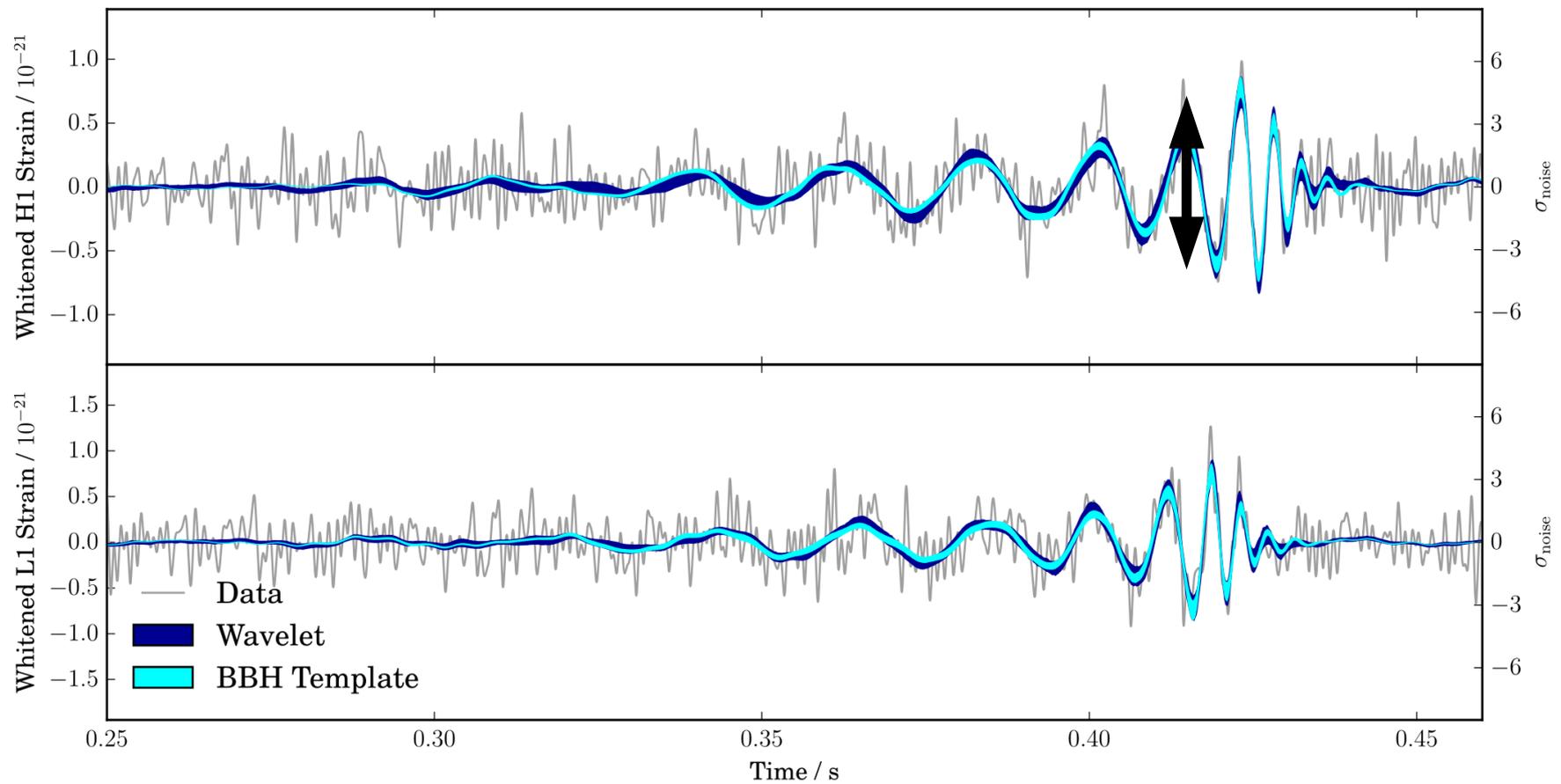
Late inspiral → total mass (+chirp mass +  $m_1/m_2$ ) → individual masses

Ringdown → final BH mass and spin



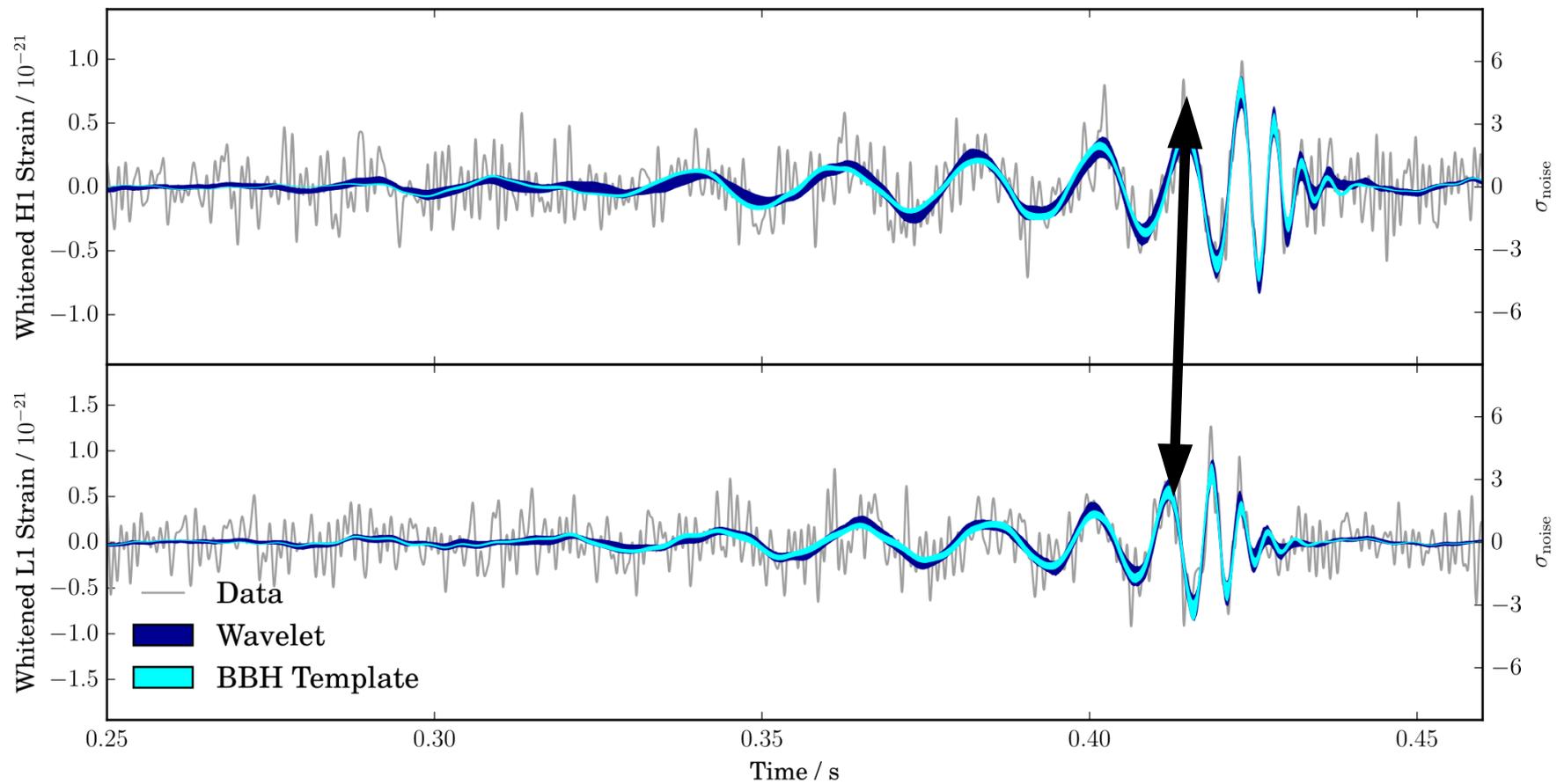
# Parameter estimation

Amplitude: inversely proportional to the distance

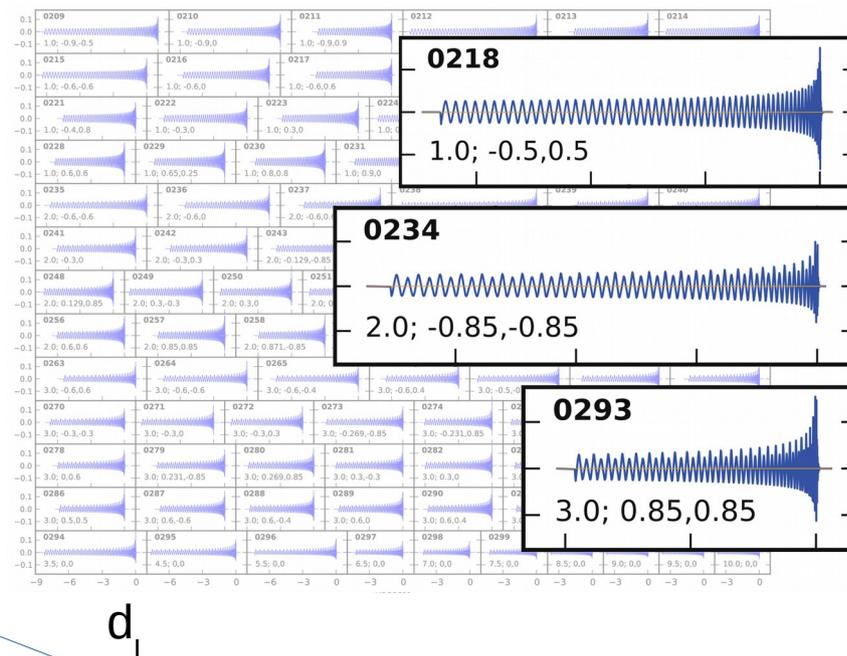
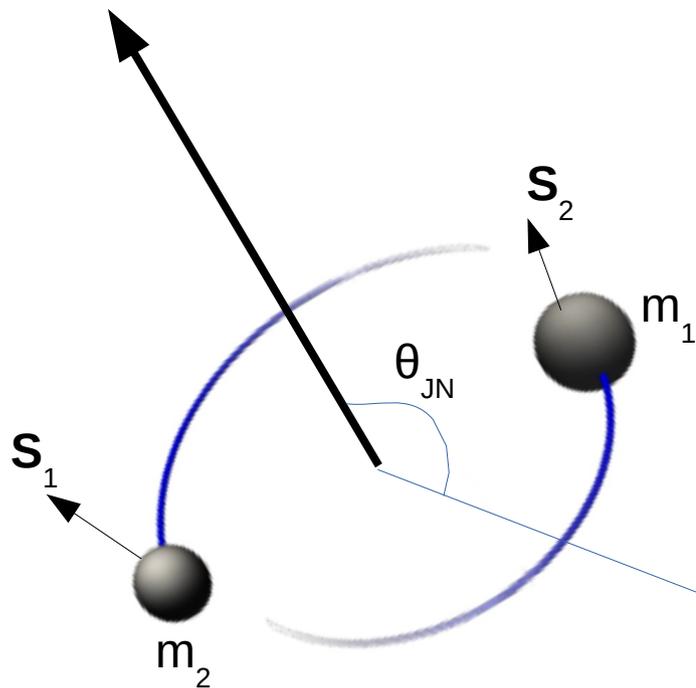


# Parameter estimation

Amplitude and phase difference between sites  $\rightarrow$  sky location  
+ Amplitude and phase consistency



# Parameter estimation



GWs (?)

noise

data from multiple instruments

$$\vec{d}^{(D)}(f) = \vec{h}^{(D)}(f) + \vec{n}^{(D)}(f)$$

$$\vec{d} = \{d^{(H)}, d^{(L)}, \dots\}$$

Bayes theorem:

$$p(\vec{\theta}|\vec{d}, \mathcal{H}, I) = \frac{p(\vec{\theta}|\mathcal{H}, I)p(\vec{d}|\vec{\theta}, \mathcal{H}, I)}{p(\vec{d}|\mathcal{H}, I)}$$

posterior

prior

likelihood

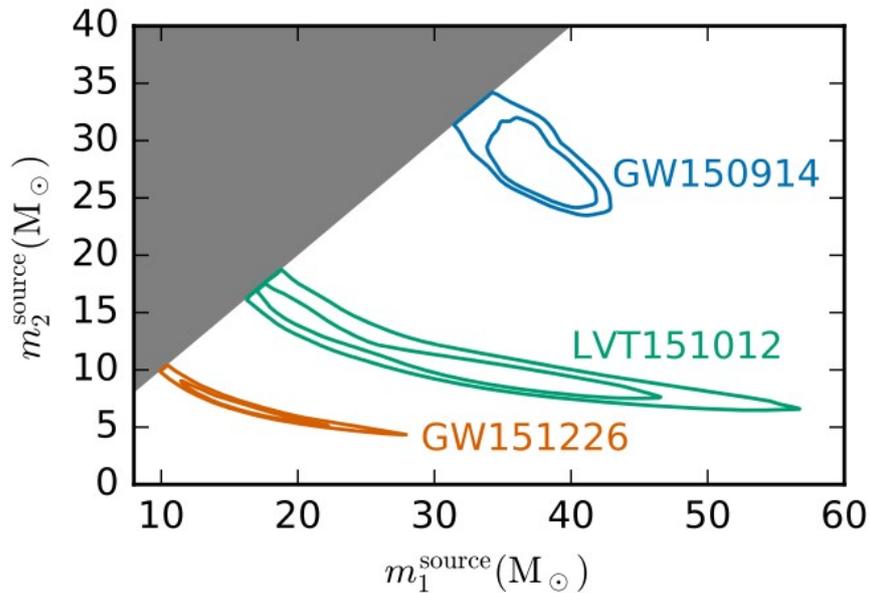
Marginalized PDF:

$$p(\vec{\theta}_A|\vec{d}, \mathcal{H}, I) = \int_{\Theta_B} p(\vec{\theta}|\vec{d}, \mathcal{H}, I)d\vec{\theta}_B$$



→ Marc Arène's talk

# Parameter estimation : ex. masses



Mostly sensitive to the chirp mass  
→  $m_1, m_2$  degeneracy

$$M_c = \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}}$$

## GW150914

$$m_1 = 36.2^{+5.2}_{-3.2} M_{sun}$$

$$m_2 = 29.1^{+3.7}_{-4.4} M_{sun}$$

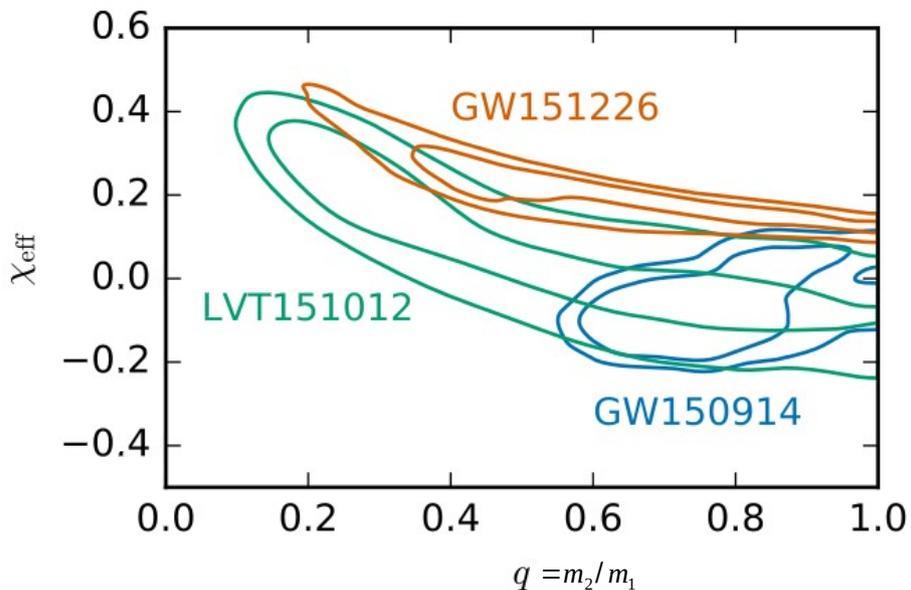
## GW151226

$$m_1 = 14.2^{+8.3}_{-3.7} M_{sun}$$

$$m_2 = 7.5^{+2.3}_{-2.3} M_{sun}$$

- All the components are black holes
- Very high masses for GW150914

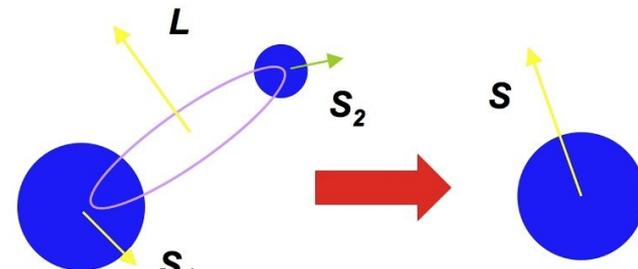
# Parameter estimation : ex. spins



$$\chi_{\text{eff}} = \frac{m_1 a_{1z} + m_2 a_{2z}}{m_1 + m_2}$$

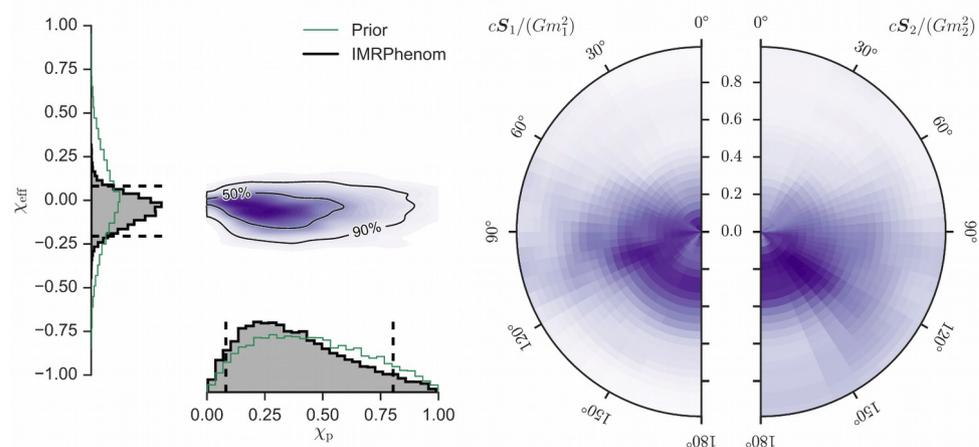
→ not well constrained

**GW151226:** at least one black hole is a Kerr black hole  
spin  $> 0.2$

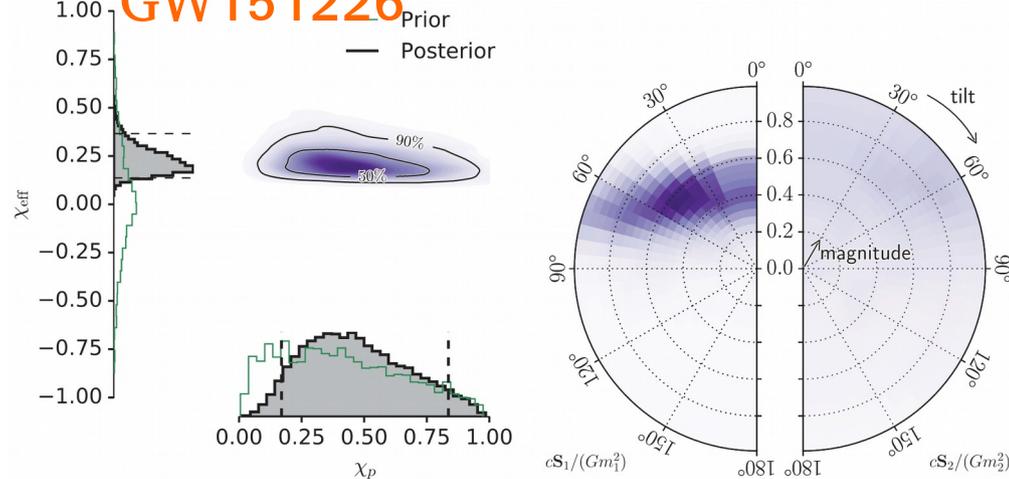


Uninformative about precession

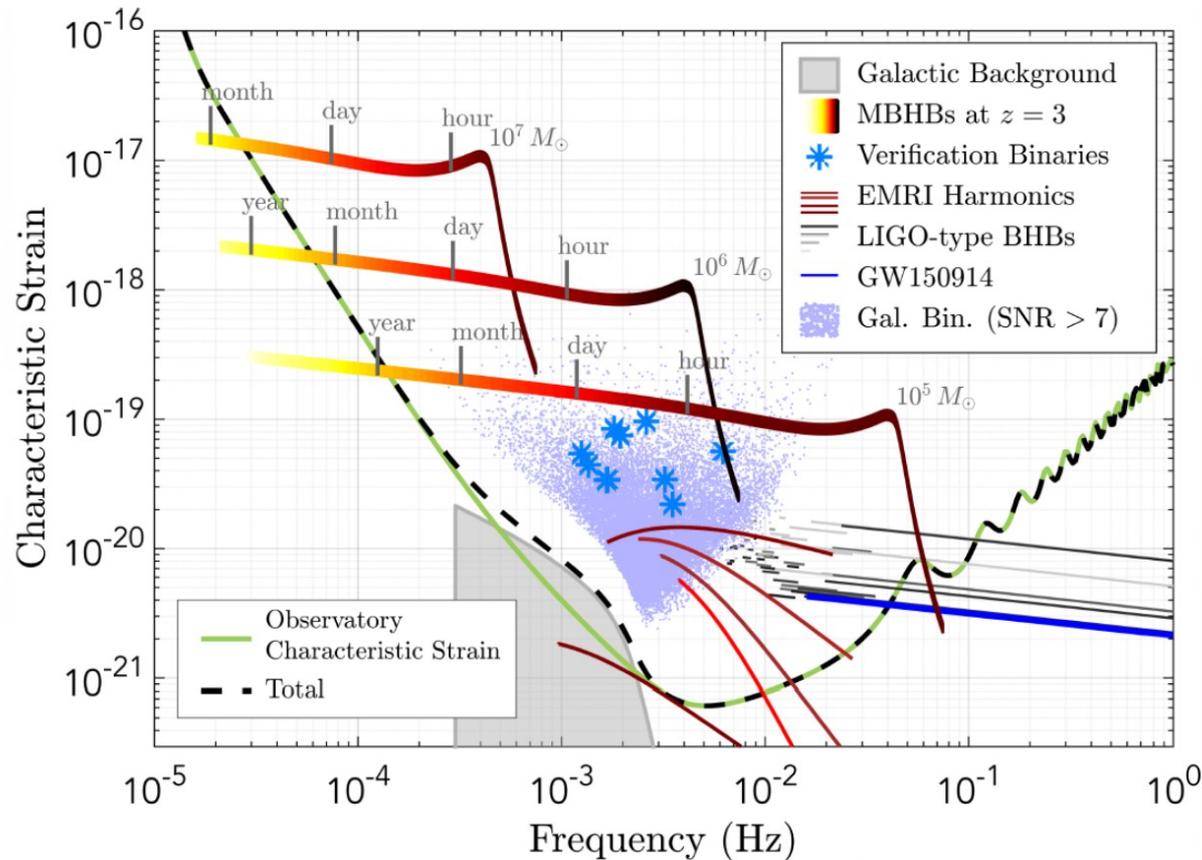
## GW150914



## GW151226



# Signal-dominated analysis → LISA



- $6 \times 10^7$  galactic binaries
- 10-100/year super-massive black hole binaries
- 10-1000/year extreme mass ratio inspirals
- large number of stellar origin black hole binaries (LIGO/Virgo)
- cosmological backgrounds

Challenge: physical background

→ See Nikolaos Karnesis's talk

https://www.gw-openscience.org

Gravitational Wave Open Science Center

The Gravitational Wave Open Science Center provides data from gravitational-wave observatories, along with access to tutorials and software tools.

[Get started!](#)

[Download data](#)

[See parameter estimation samples](#)

[Join the email list](#)

[Explore the open data web course](#)

The LIGO observatories are built and operated by the LIGO Laboratory (California Institute of Technology and Massachusetts Institute of Technology) with participation by the LIGO Scientific Collaboration, and are supported by the U.S. National Science Foundation. The Virgo detector is designed, built and operated by a collaboration that includes the Centre National de la Recherche Scientifique (France), the Istituto Nazionale di Fisica Nucleare (Italy) and Nikhef (Netherlands), with Polish, Hungarian and Spanish institutes and the European Gravitational Observatory (EGO) consortium. For general information, visit [ligo.org](http://ligo.org) and [virgo-gw.eu](http://virgo-gw.eu). The LIGO Laboratory's Data Management Plan describes the scope and timing of LIGO data releases.

LIGO Open Science Center

Data Releases for Observed Transients

Data Releases: Compact Object Mergers

Click icons below for data and documentation:

GW150914 LVT151012 GW151226 GW170104 GW170608

GW170814 GW170817

Audio files

Listen to audio files from LIGO detections.

Rapid Triggers from LIGO Data

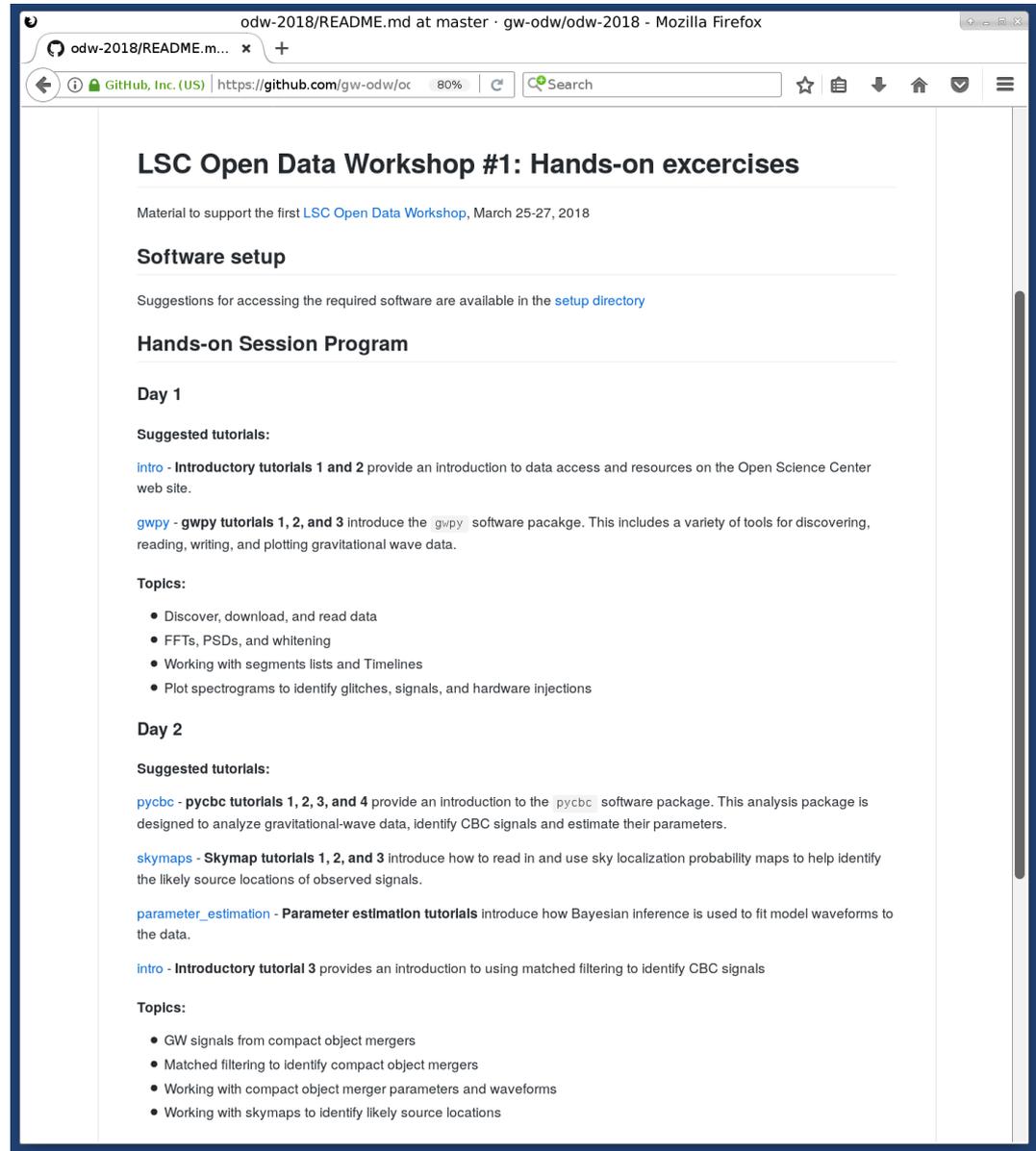
During O1 and O2, information about gravitational-wave detections is made available to astronomers as GCN notices. This exchange is archived:

- GW150914
- LVT151012
- GW151226
- GW170104
- GW170608
- GW170814
- GW170817

~1 heure autour de chaque événement publié

<https://www.gw-openscience.org>

Online tutorials and hands-on exercises



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odw-2018/README.m... x +

GitHub, Inc. (US) | <https://github.com/gw-odw/oc> 80% | Search

## LSC Open Data Workshop #1: Hands-on exercises

Material to support the first [LSC Open Data Workshop](#), March 25-27, 2018

### Software setup

Suggestions for accessing the required software are available in the [setup directory](#)

### Hands-on Session Program

#### Day 1

**Suggested tutorials:**

[intro](#) - **Introductory tutorials 1 and 2** provide an introduction to data access and resources on the Open Science Center web site.

[gwpv](#) - **gwpv tutorials 1, 2, and 3** introduce the `gwpv` software package. This includes a variety of tools for discovering, reading, writing, and plotting gravitational wave data.

**Topics:**

- Discover, download, and read data
- FFTs, PSDs, and whitening
- Working with segments lists and Timelines
- Plot spectrograms to identify glitches, signals, and hardware injections

#### Day 2

**Suggested tutorials:**

[pycbc](#) - **pycbc tutorials 1, 2, 3, and 4** provide an introduction to the `pycbc` software package. This analysis package is designed to analyze gravitational-wave data, identify CBC signals and estimate their parameters.

[skymaps](#) - **Skymap tutorials 1, 2, and 3** introduce how to read in and use sky localization probability maps to help identify the likely source locations of observed signals.

[parameter\\_estimation](#) - **Parameter estimation tutorials** introduce how Bayesian inference is used to fit model waveforms to the data.

[intro](#) - **Introductory tutorial 3** provides an introduction to using matched filtering to identify CBC signals

**Topics:**

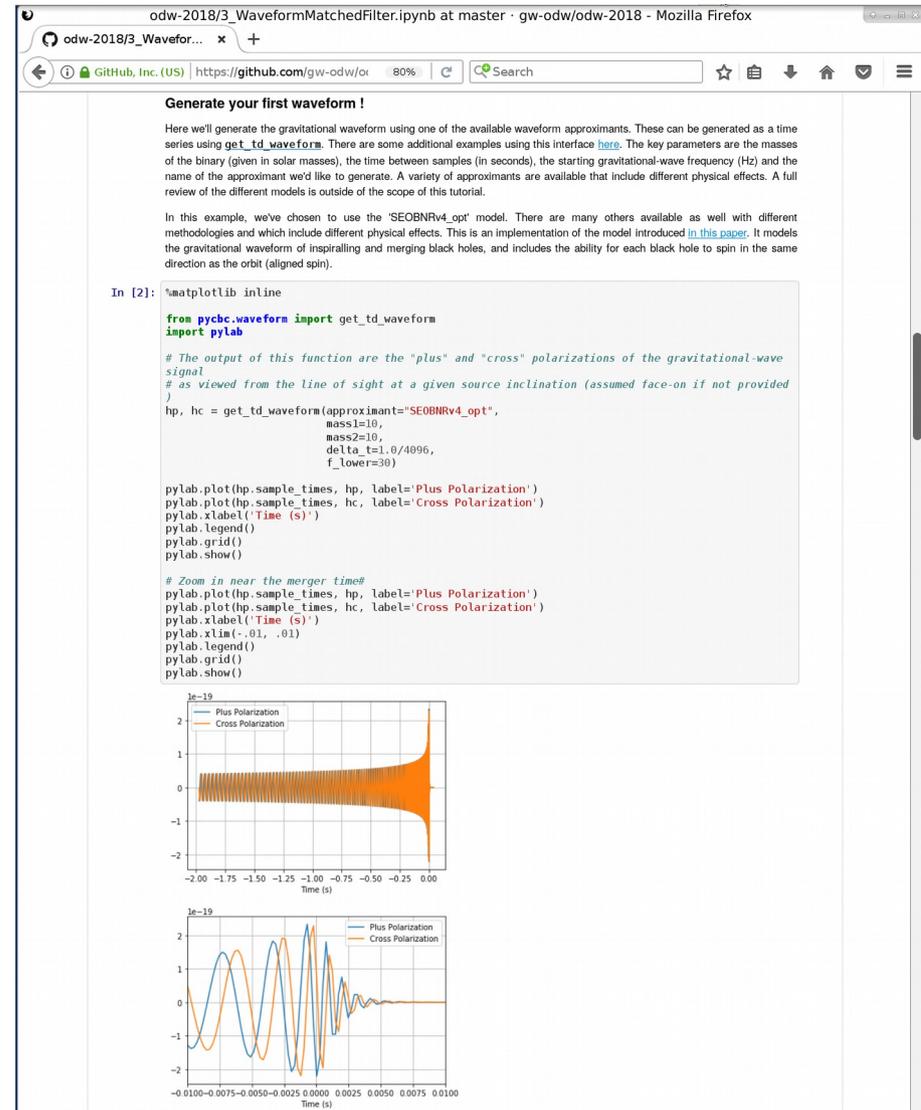
- GW signals from compact object mergers
- Matched filtering to identify compact object mergers
- Working with compact object merger parameters and waveforms
- Working with skymaps to identify likely source locations

<https://www.gw-openscience.org>

## Jupyter notebook Software libraries

```
pip install gwpy  
pip install lalsuite  
pip install pycbc
```

Data access,  
Waveform generation  
Filtering, pre-processing  
And more...



## Data analysis challenge and open questions:

- rare signals buried in detectors' noise
- highly non-Gaussian noise distribution (glitches)
- no model to describe the noise components
- Noise subtraction
- Detect gravitational-waves with a single detector
- a full coherent CBC analysis is computing demanding
- simplifications must be used for spins and orbit eccentricity of binaries
- low-latency analysis (parameter estimation) for multi-messenger follow-up
- source rate not well constrained
- physical background (ex. galactic binaries for LISA)
- control systematic effects (pulsar timing)