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1

SPectroscopy with Atom Neutrino: its principle and recent progress

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Okayama University

Introduction

known and unknown properties of neutrino

PDG

known properties

- PMNS mixing angle
- squared mass difference
- (Dirac CP phase)
- (Mass ordering)

$$\begin{aligned} \sin^2(\theta_{12}) &= 0.307 \pm 0.013 \\ \Delta m_{21}^2 &= (7.53 \pm 0.18) \times 10^{-5} \text{ eV}^2 \\ \sin^2(\theta_{23}) &= 0.421^{+0.033}_{-0.025} \quad (S = 1.3) \quad (\text{Inverted order, quad. I}) \\ \sin^2(\theta_{23}) &= 0.592^{+0.023}_{-0.030} \quad (S = 1.1) \quad (\text{Inverted order, quad. II}) \\ \sin^2(\theta_{23}) &= 0.417^{+0.025}_{-0.028} \quad (S = 1.2) \quad (\text{Normal order, quad. I}) \\ \sin^2(\theta_{23}) &= 0.597^{+0.024}_{-0.030} \quad (S = 1.2) \quad (\text{Normal order, quad. II}) \\ \Delta m_{32}^2 &= (-2.56 \pm 0.04) \times 10^{-3} \text{ eV}^2 \quad (\text{Inverted order}) \\ \Delta m_{32}^2 &= (2.51 \pm 0.05) \times 10^{-3} \text{ eV}^2 \quad (S = 1.1) \quad (\text{Normal order}) \\ \sin^2(\theta_{13}) &= (2.12 \pm 0.08) \times 10^{-2} \end{aligned}$$

unknown properties

- **absolute** mass (only upper limit)
- mass type (Dirac or Majorana)
- (Majorana CP phase)
- etc.

key parameters for particle physics and cosmology

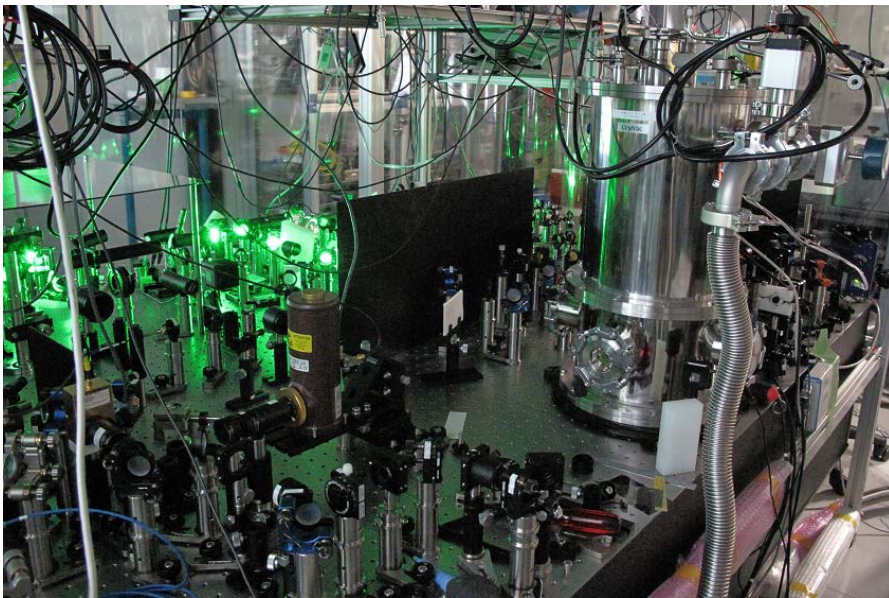
- BSM physics
- leptogenesis

Our approach to neutrino

- ✓ use **atomic or molecular** de-excitation process and techniques of **laser** spectroscopy

table-top experiment

laboratory in Okayama University



interdisciplinary science

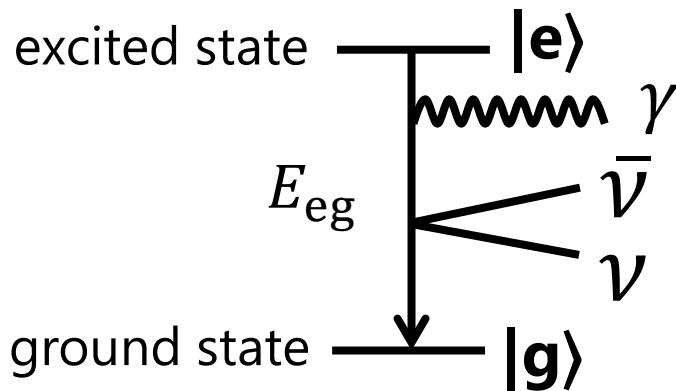
- areas of expertise
 - High energy physics
 - nuclear physics
 - AMO (atomic, molecular and optical) physics
 - chemistry
 - theorists

Experimental principle

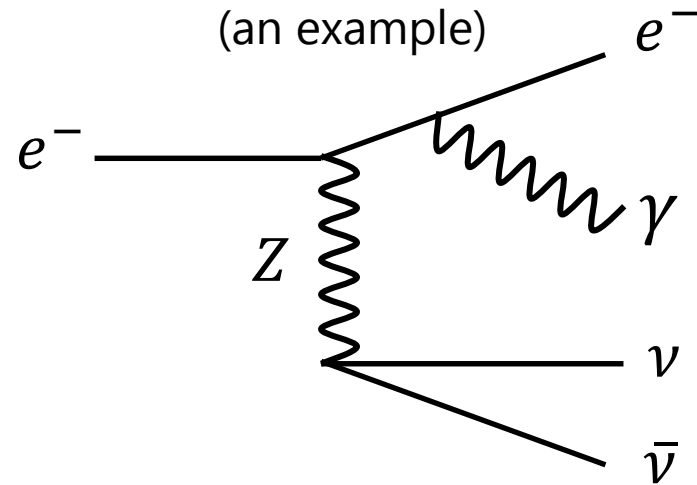
Radiative Emission of Neutrino Pair (RENPN)

$$|e\rangle \rightarrow |g\rangle + \gamma + \nu + \bar{\nu}$$

transition of atomic or
molecular state



Feynman diagram
(an example)

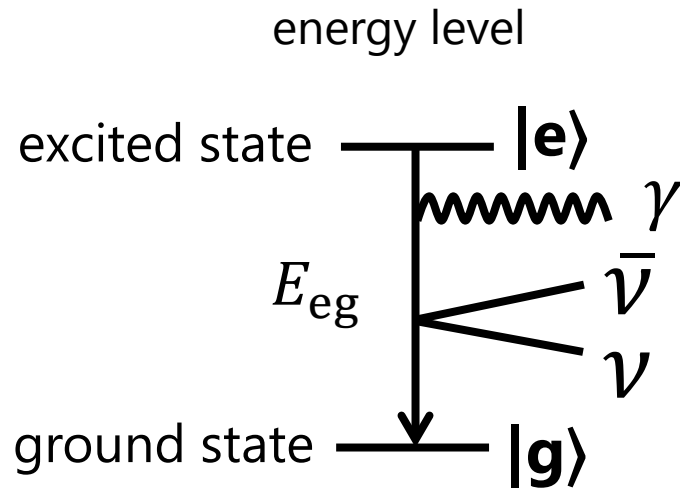


- ✓ detection of a photon (easy) instead of neutrino (difficult)

Experimental principle

Radiative Emission of Neutrino Pair (RENPN)

$$|e\rangle \rightarrow |g\rangle + \gamma + \nu + \bar{\nu}$$



- threshold energy of photon
neutrino absolute mass

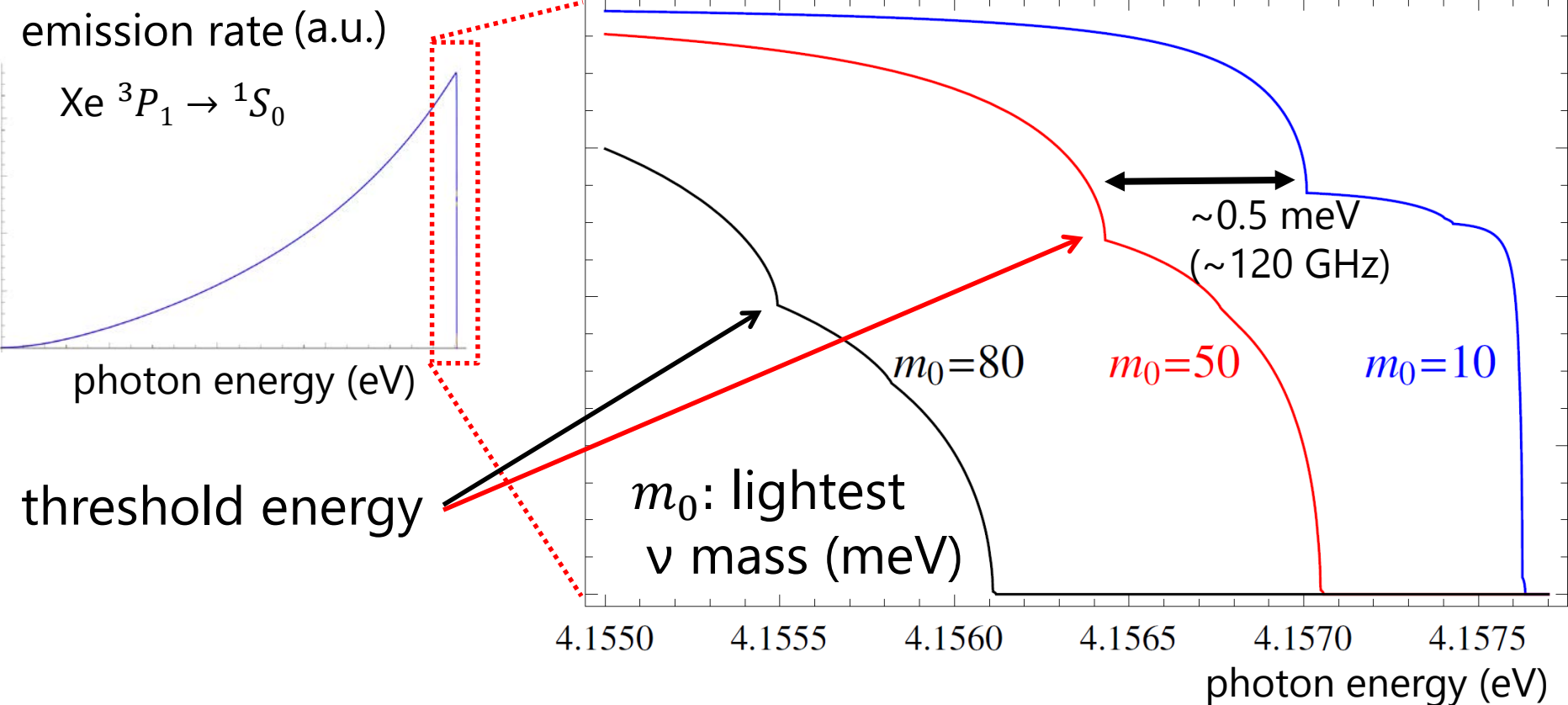
$$E_{th} = \frac{E_{eg}}{2} - \frac{(m_{\nu} + m_{\bar{\nu}})^2}{2E_{eg}}$$

- ✓ The emitted photon contains information of neutrinos

Energy spectra of emitted photon

- RENP process

Y. Miyamoto et al.
Prog. Theor. Exp. Phys. 2015 081C01



- ✓ Energy spectra are obtained by scanning the laser frequency
 - Frequency resolution is much better than 1 GHz
 - ➡ precise determination of neutrino absolute mass is possible

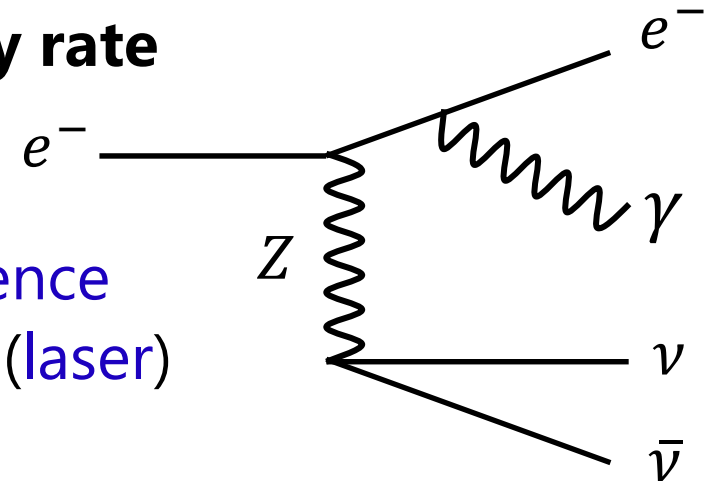
Amplification of emission rate

✓ Critical issue of this method: **very tiny rate**

- typical emission rate $\ll 10^{-30}$ Hz
- transition includes weak interaction

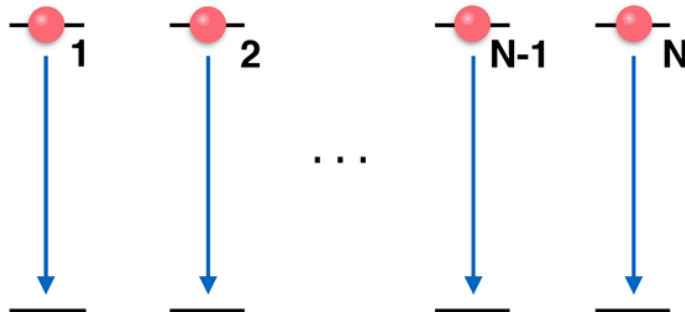
➔ **rate amplification** by **atomic coherence**

- obtained by using coherent photons (**laser**)



Simplified description

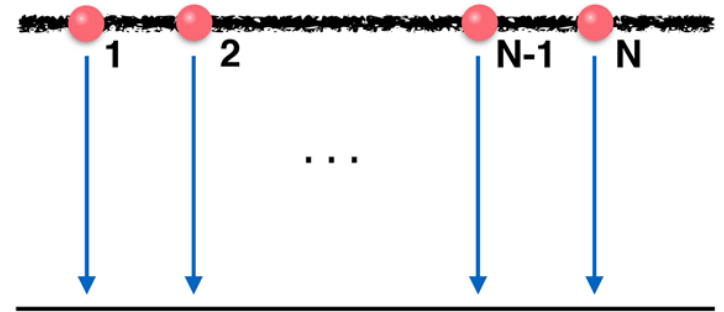
incoherent de-excitation



$$N|A|^2$$

$\times N$

coherent de-excitation



$$|NA|^2$$

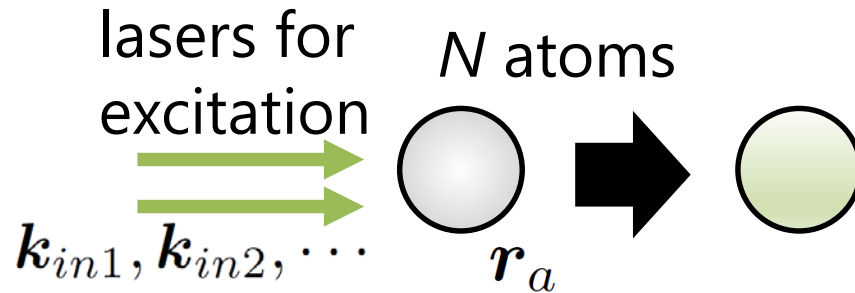
Huge enhancement can be achieved!

Rate Amplification

- Condition for rate amplification

proposed by
M. Yoshimura (2007)

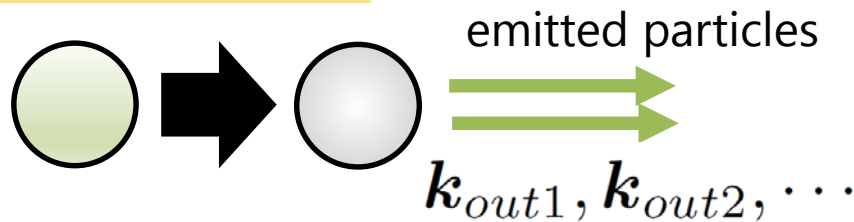
coherence generation



initial spatial phase

$$\sum_a \exp(i\mathbf{k}_{in} \cdot \mathbf{r}_a)$$

de-excitation process



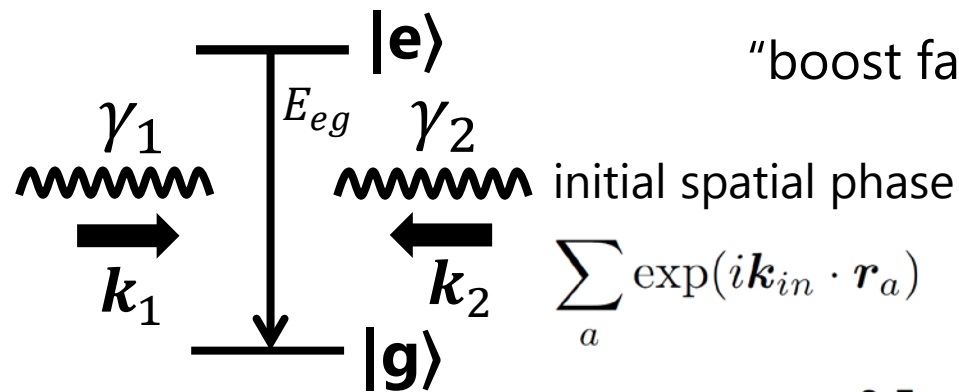
de-excitation rate \propto

$$\left| \sum_a \exp(i\Delta\mathbf{k} \cdot \mathbf{r}_a) \mathcal{M}_a \right|^2$$

$$\Delta\mathbf{k} = \mathbf{k}_{in} - \mathbf{k}_{out}$$

- If $\Delta\mathbf{k} = \mathbf{0}$ holds, the emission rate $\propto \mathbf{N}^2$ (**rate amplification**)
 - momentum conservation among initial and emitted particles
 - also coherence decay (decoherence) time should be long
- ✓ **Experimental demonstration of this principle is necessary**

Dirac/Majorana distinction



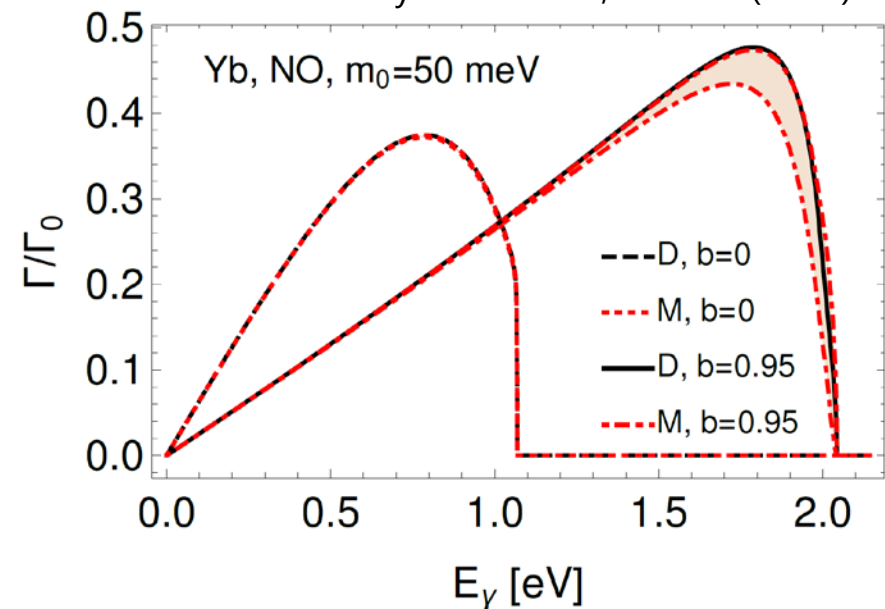
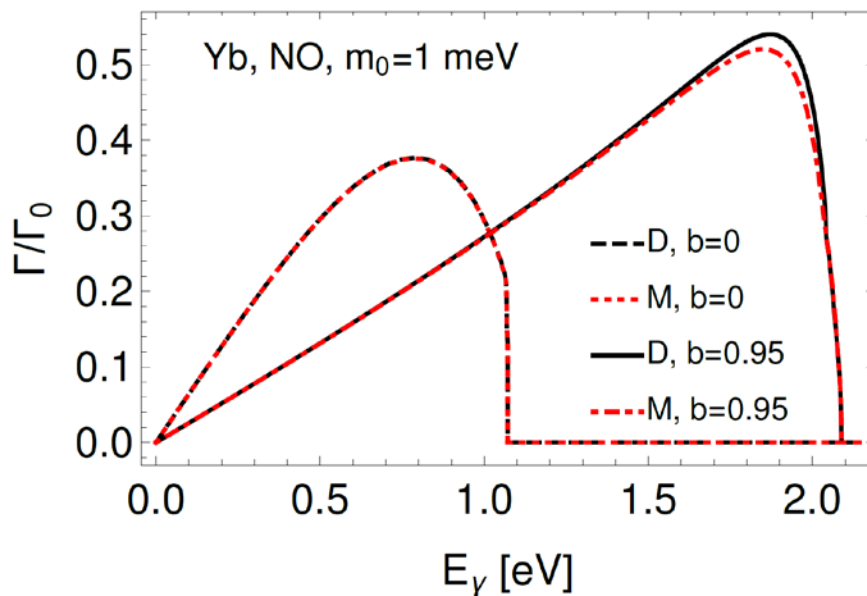
"boost factor" $b = \frac{|k_{in}|}{E_{eg}}$

$|k_1| = |k_2| \rightarrow b = 0$

$|k_1| \gg |k_2| \rightarrow b \sim 1$

M. Tanaka, *et al.*

Phys. Rev. D **96**, 113005 (2017)



- Dirac/Majorana distinction : principally possible
 - effect of identical particle emission (Majorana)

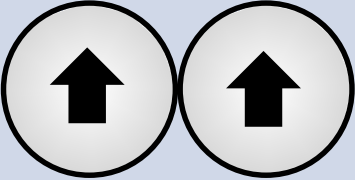
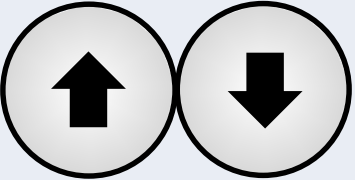
**Rate amplification experiment of
two-photon emission (TPE)
from para-H₂ molecule**

para-H₂ experiment

Motivation

- ✓ study rate amplification mechanism using **two-photon emission** processes
 - much easier to observe than neutrino emission process

Properties of para-H₂

	nuclear spin	rotational quantum number
ortho-H ₂ 	$I = 1$	$J = 1, 3, 5... \text{ (odd)}$
para-H ₂ 	$I = 0$	$J = 0, 2, 4... \text{ (even)}$

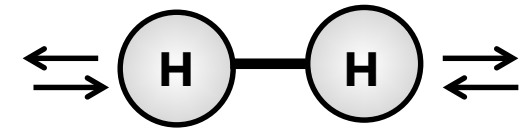
Pauli exclusion principle

dominant @ 78 K

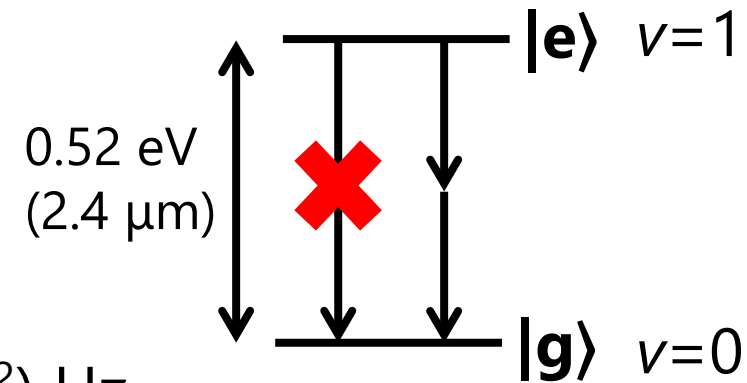
- ✓ $J=0$ (ground state) para-H₂: completely spherical wavefunction
 - ➔ weak intermolecular interaction
 - ➔ **long decoherence time** expected

coherent states of para-H₂

✓ coherence between vibrational states ($v=0, J=0 \leftrightarrow v=1, J=0$) of pH₂ molecules

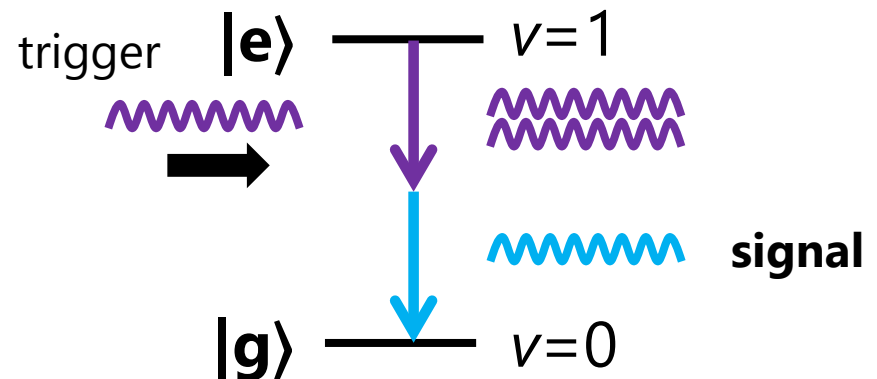


- 1-photon electric dipole transition (E1 transition): **forbidden**
- 2-photon E1 × E1 transition: **allowed**
- ✓ metastable excited state
 - spontaneous emission rate: $\mathcal{O}(10^{-12})$ Hz
- ✓ decoherence time: $\mathcal{O}(1)$ ns (gas)
 - mainly due to molecular collision



Induction of two-photon emission

- induce TPE by **trigger laser**
- $\omega_{\text{trigger}} + \omega_{\text{signal}} = E_{eg}$



Coherence generation scheme

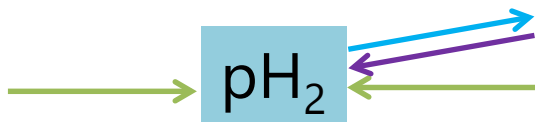
Requirement for rate amplification: 4-momentum conservation

one-side laser excitation



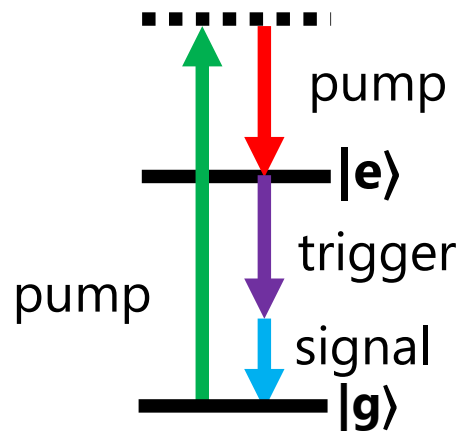
- simpler configuration
- invariant mass by pump photons: 0

counter-propagating laser excitation

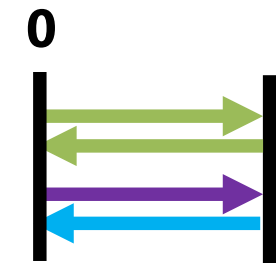
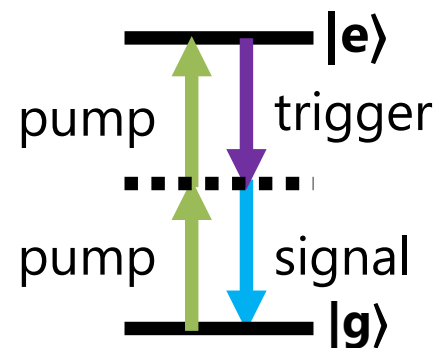
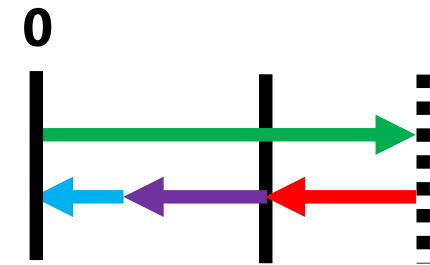


- **non-zero invariant mass**
- necessary for producing massive particle (**neutrino**)

Energy



momentum



Overview of the para-H₂ TPE experiment

one-side laser excitation

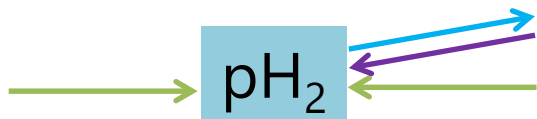


- ①: para-H₂ gas experiment
 - ②: para-H₂ solid experiment
- (difference is explained later)

gas para-H₂ cell
cell length: 15 cm
cooled to ~ 78 K

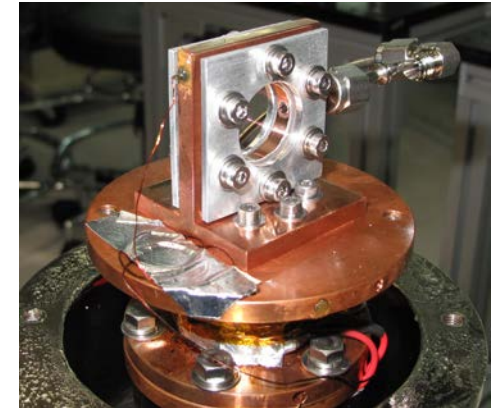


counter-propagating
laser excitation



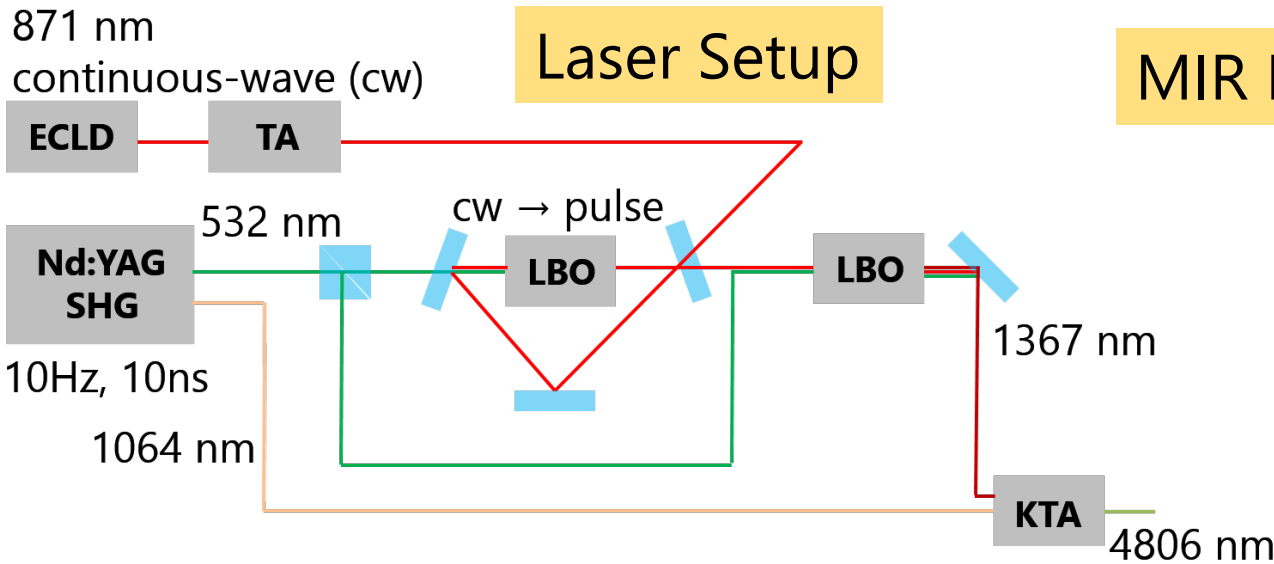
- ③: para-H₂ gas experiment

solid para-H₂ cell
cell length: 5 mm
cooled to ~ 4 K



Mid-infrared (MIR) laser generation

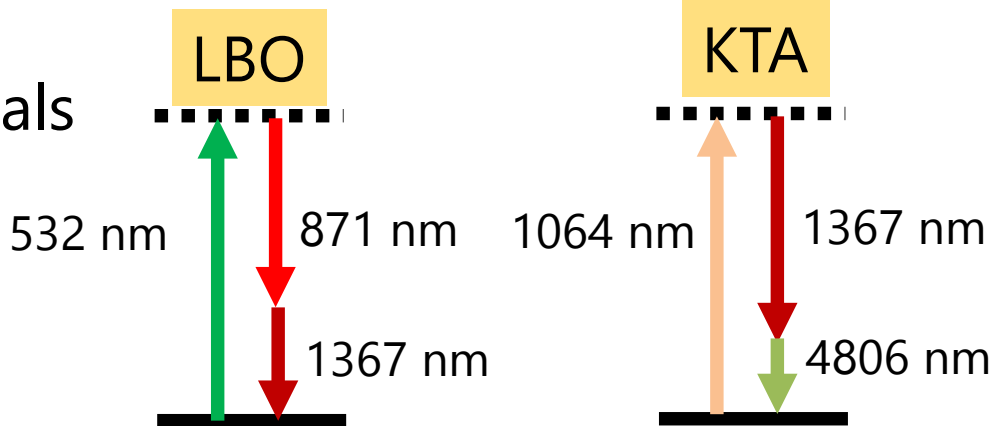
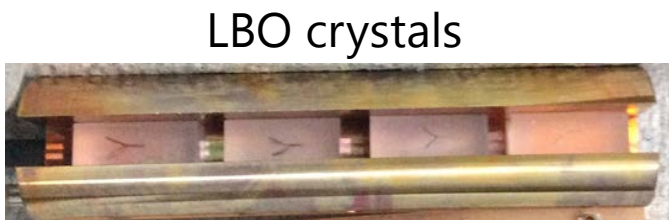
- ✓ developed **high-intensity and narrow-linewidth** MIR laser
- used for **counter-propagating excitation** experiment



MIR Laser specifications

energy:
~5 mJ/pulse
linewidth:
~150 MHz

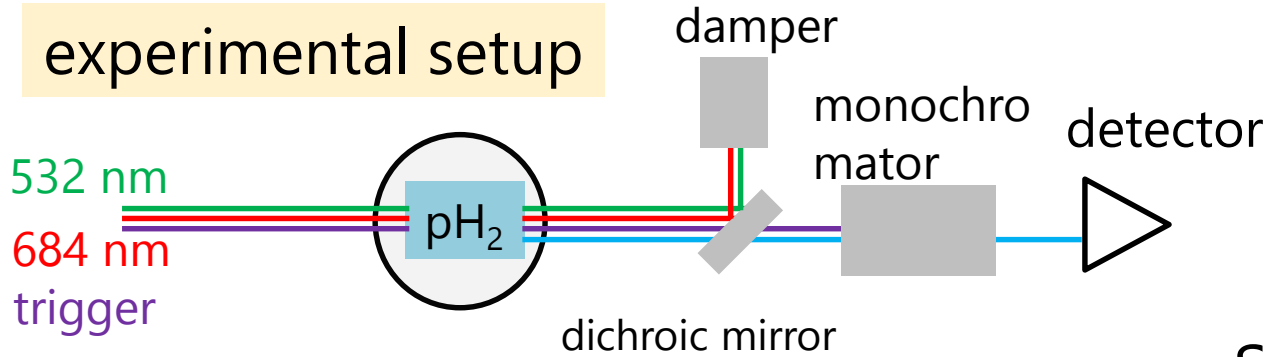
- Frequency conversion by nonlinear-optical crystals



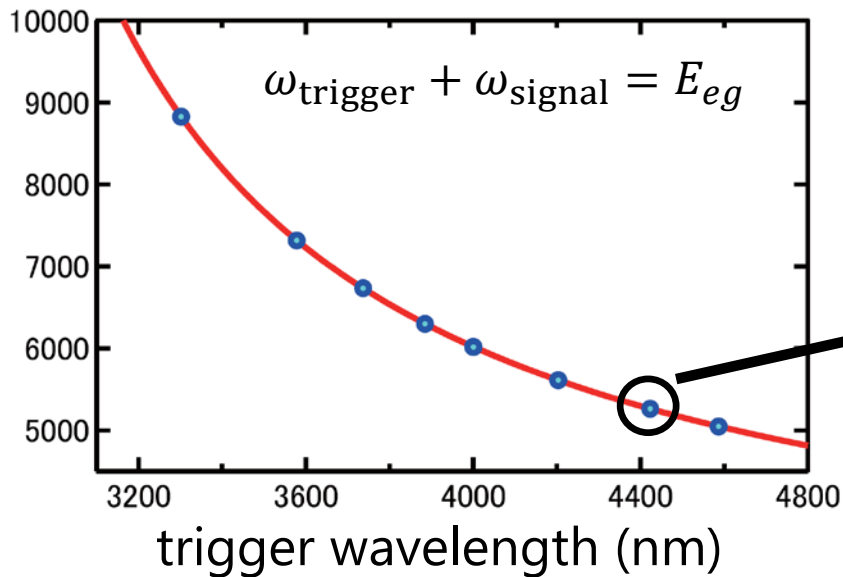
**para-H₂ one-side
laser excitation experiment**

① para-H₂ gas experiment

experimental setup



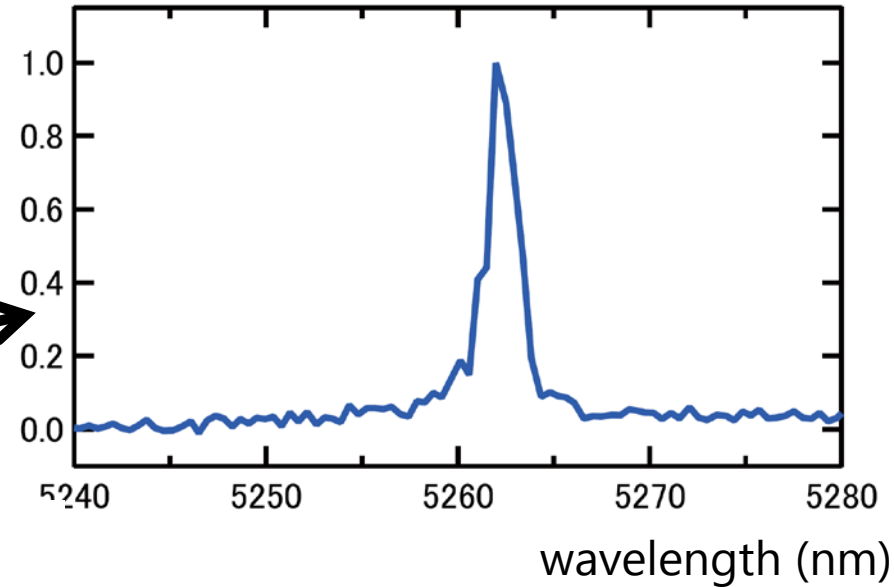
signal wavelength (nm)



signal (a.u.)

Signal spectrum

H. Hara et al. PRA **96**, 063827 (2017)

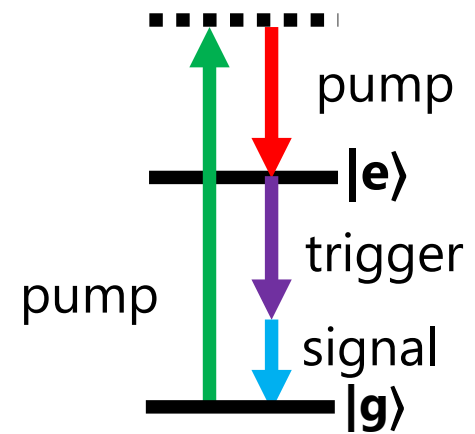


Rate amplification factor: (amplified rate)/(spontaneous emission rate) $> 10^{18}$

Y. Miyamoto et al. Prog. Theor. Exp. Phys. **2014**, 113C01

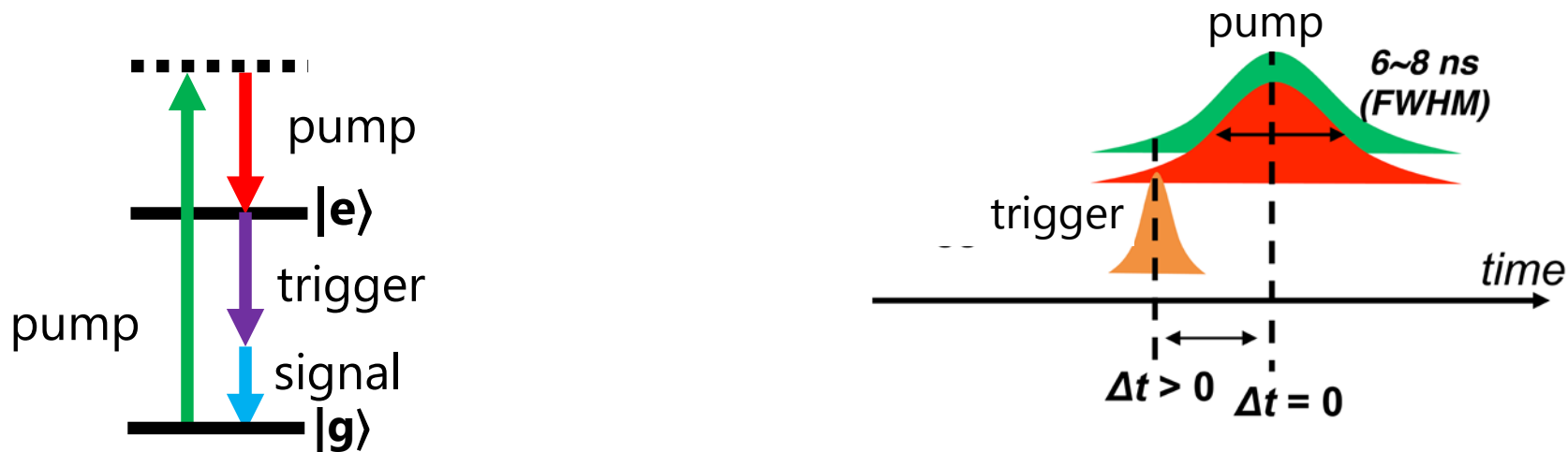
② para-H₂ solid experiment

- ✓ **solid** para-H₂: called "Quantum solid"
rotational and vibrational excited states exist even in solid
- weak intermolecular interaction:
- no collisional broadening (cause of decoherence)



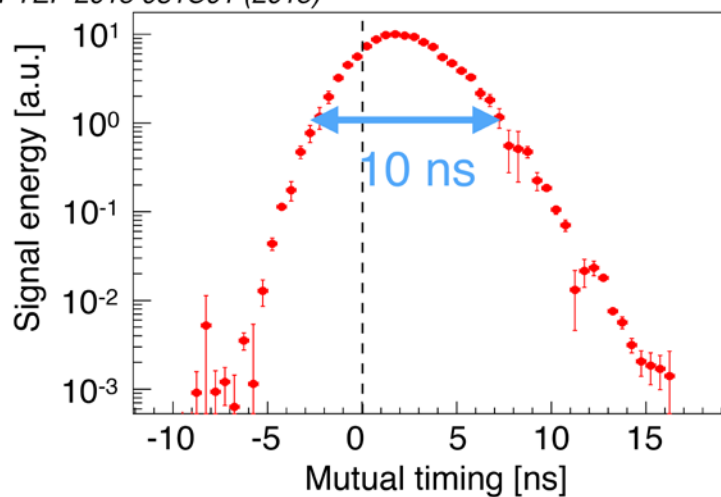
	gas para-H ₂	solid para-H ₂
density	$\sim 10^{20}$ /cm ³	2.6×10^{22} /cm ³
damage threshold	High	low
coherence time	$\mathcal{O}(1)$ ns	$\mathcal{O}(10)$ ns

gas V.S. solid timing dependence



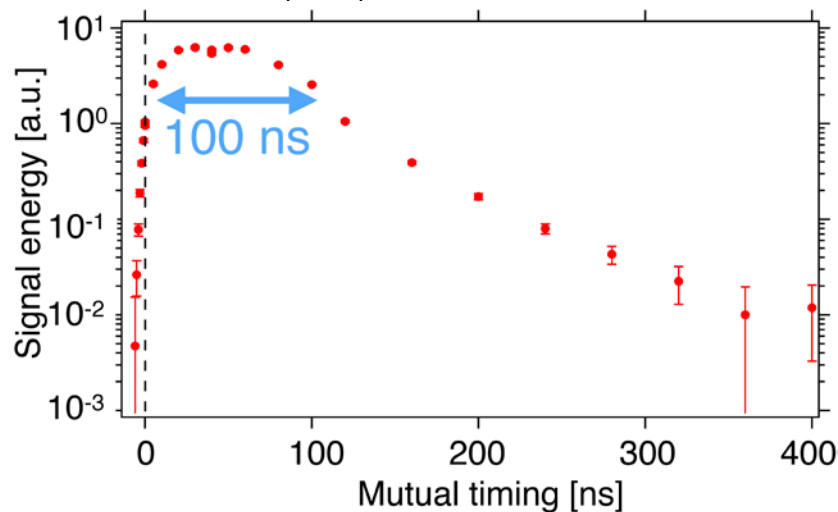
Y. Miyamoto et al.,
PTEP 2015 081C01 (2015)

gas



Y. Miyamoto et al.
J. Phys. Chem. A **121**, 3943 (2017)

solid

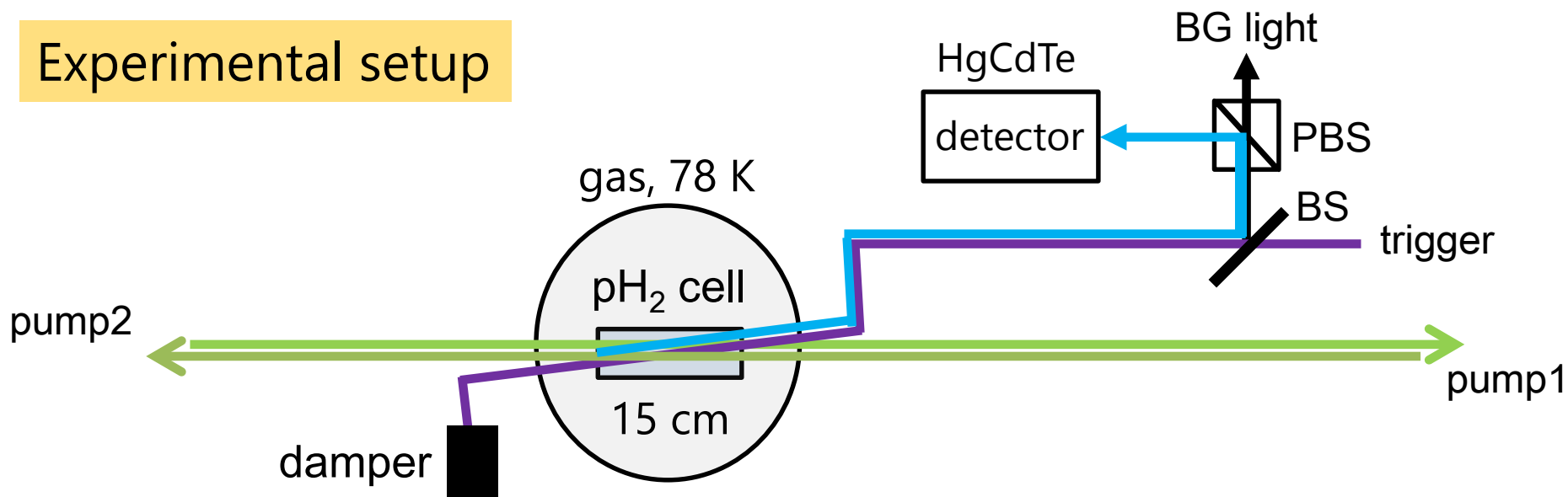


✓ Coherence develops after the pump lasers exist

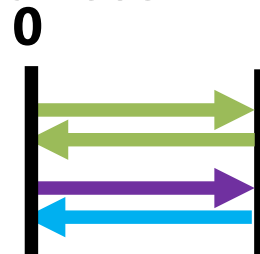
para-H₂ counter-propagating laser excitation experiment

③ Counter-propagating experiment

Experimental setup



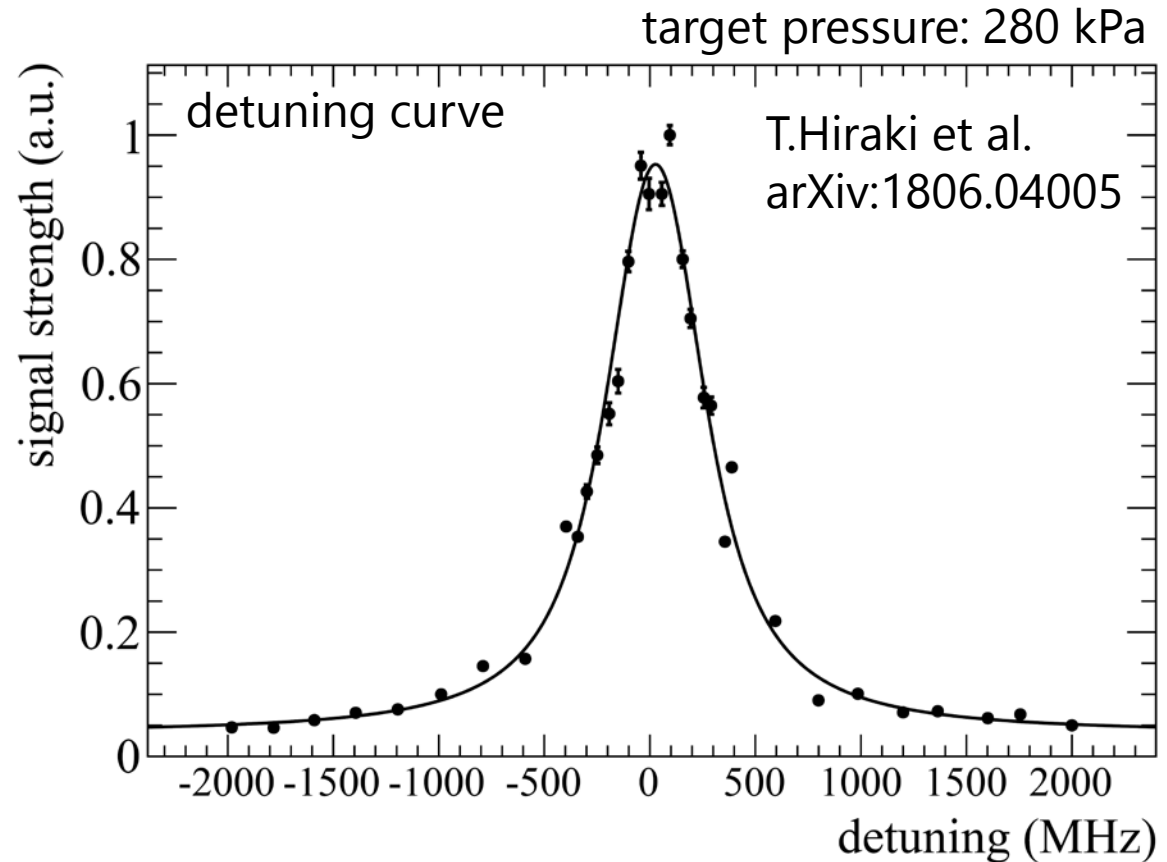
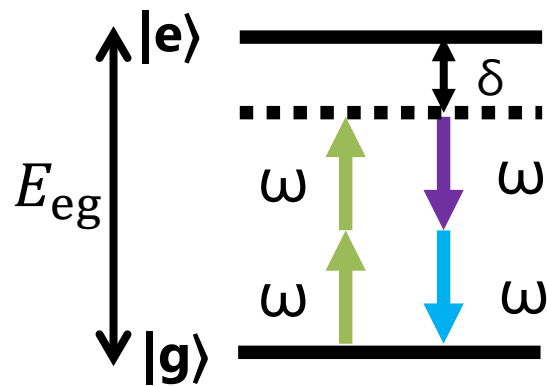
- Signal light is generated by the trigger laser and advances in the backward direction
- amplification condition
- Wrong-polarization component of the background scattering light is reduced by using a polarized beam splitter.



Results: detuning dependence

- use the new mid-infrared laser as both pumps and trigger
 - pump energy: ~ 1 mJ/pulse, trigger energy: ~ 0.6 mJ/pulse

✓ vary the detuning δ

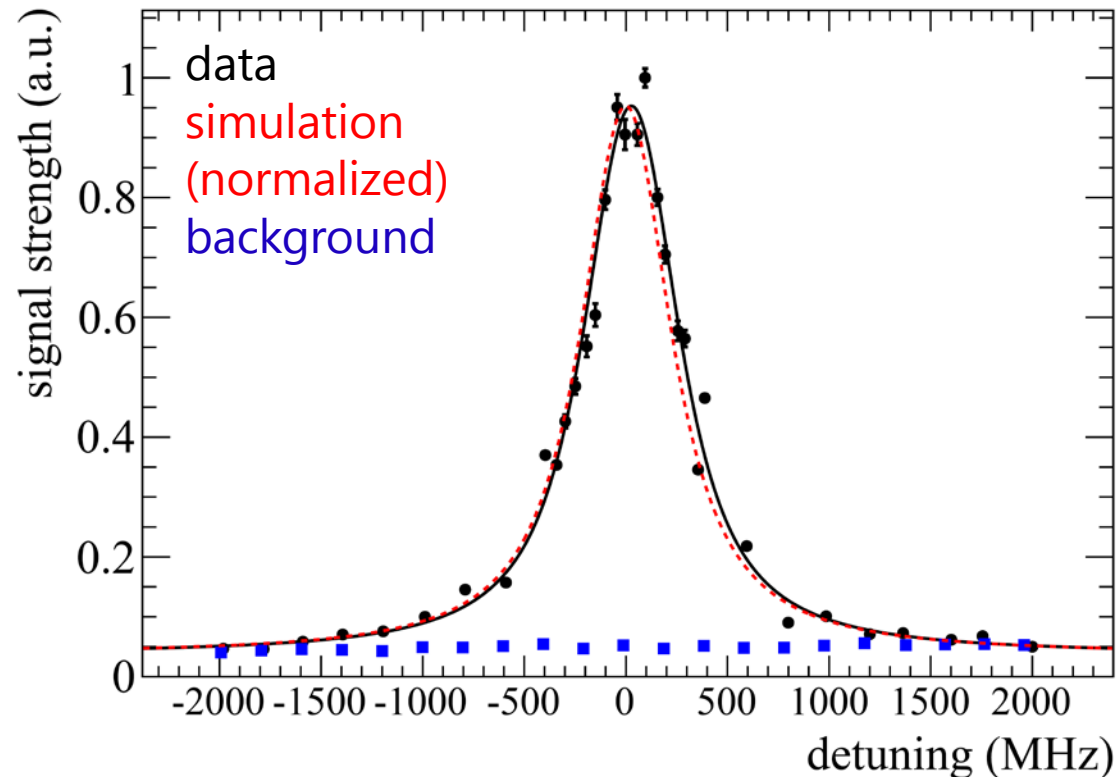
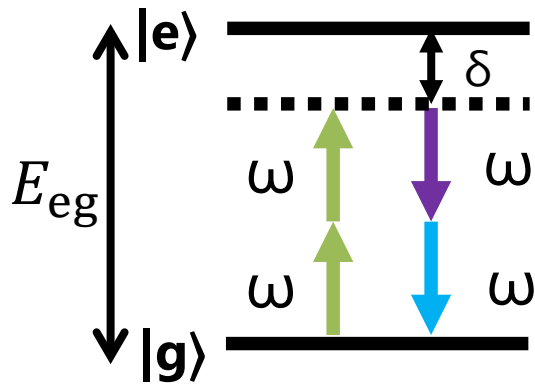


✓ Successfully observed a clear signal peak!

Results: detuning dependence

target pressure: 280 kPa

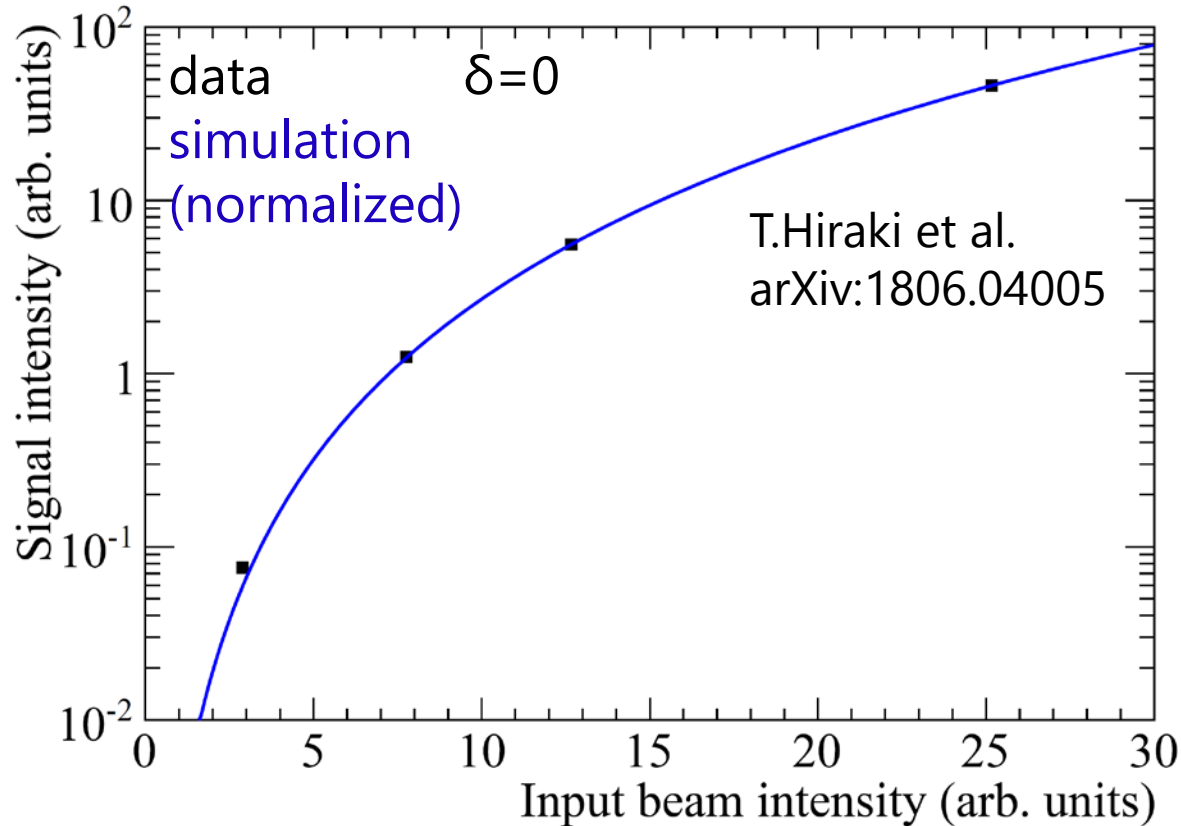
- vary the detuning δ



- comparison with simulation based on Maxwell-Bloch equations
 - describe development of laser fields and coherence
- Though it is difficult to reproduce absolute signal intensity, curve shape is consistent between data and simulation.

Results: input energy dependence

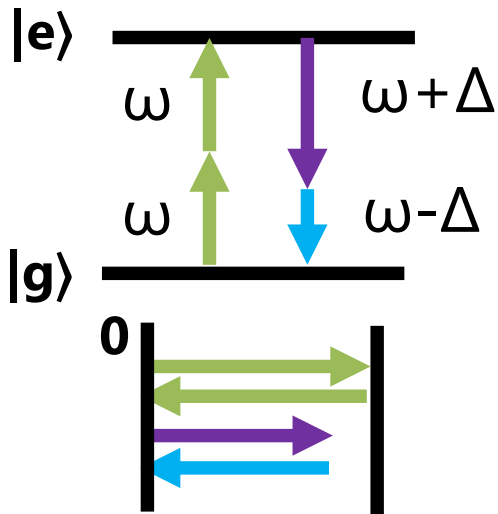
- ✓ vary the pump and trigger beam energies at the same time



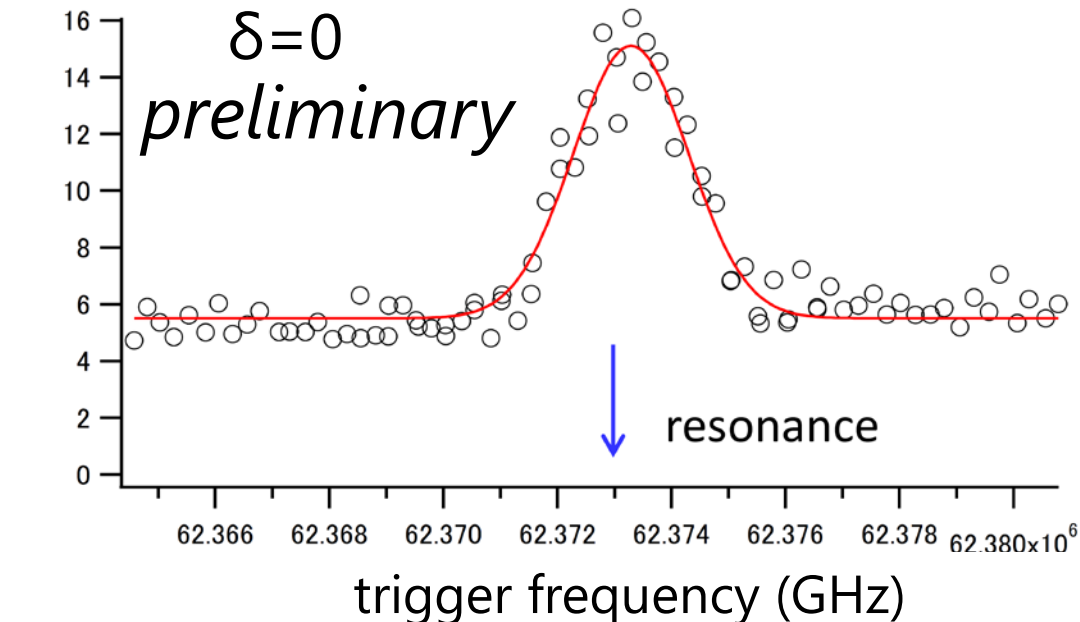
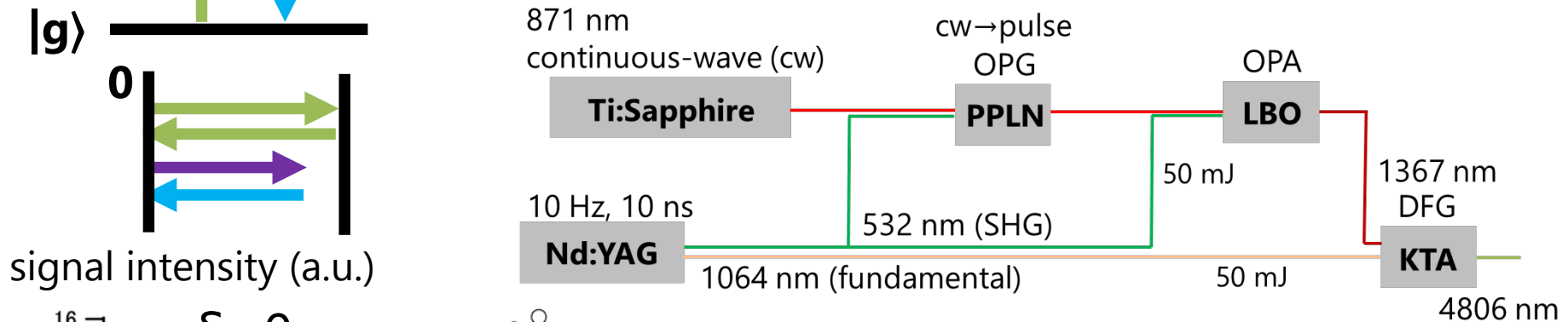
- Signal intensity is proportional to both pump beams and trigger beam

$$I_{\text{signal}} \propto I_{\text{pump1}} I_{\text{pump2}} I_{\text{trigger}} \propto I^3$$

trigger frequency dependence



- vary **only** the frequency of the **trigger** laser
- amplification condition requires $\Delta=0$
- another MIR laser is necessary



- ✓ Signal peak is observed!
- amplification condition is confirmed

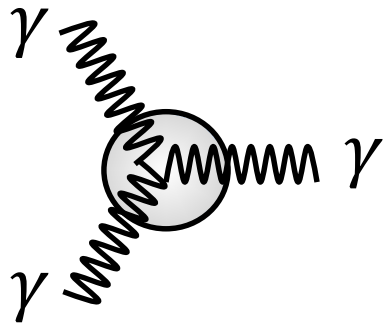
paper in preparation

Next step

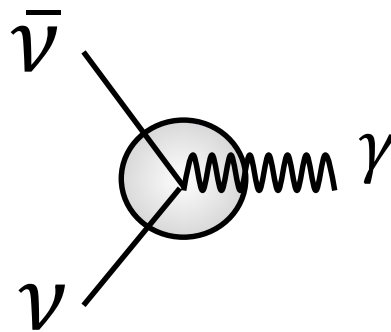
Higher QED process

- study of coherent amplification of **higher QED** process
 - 2-photon $E1 \times M1$ (magnetic dipole), 3-photon $E1 \times E1 \times E1$

3-photon emission

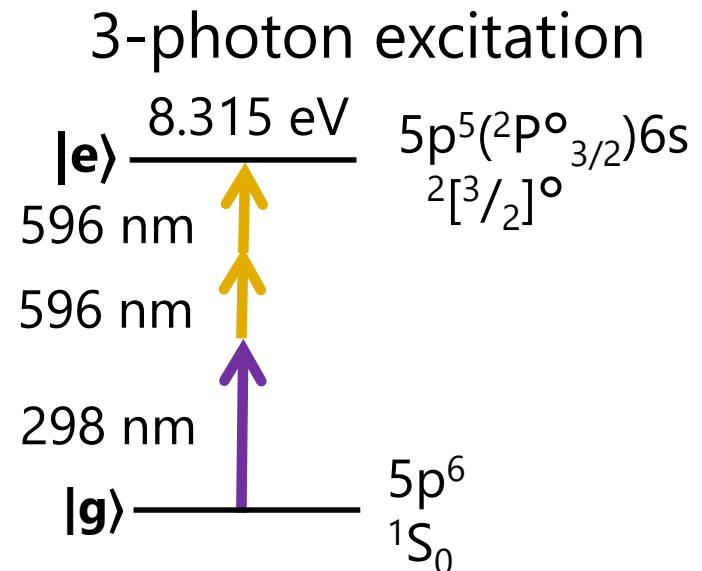


RENP



similar kinematics

- ✓ Xe target:
 - one of the candidates of the RENP experiment
- use metastable excited state
 - $E1$, $E1 \times E1$: forbidden
 - $E1 \times M1$, $E1 \times E1 \times E1$: allowed



876 nm
continuous-wave (cw)

Laser setup (Xe)

ECDL

TA

cw → pulse

Nd:YAG

LBO

Amp

Ti:S

DFG

LBO

SHG

LBO

298 nm

532 nm (SHG)

OPG

596 nm

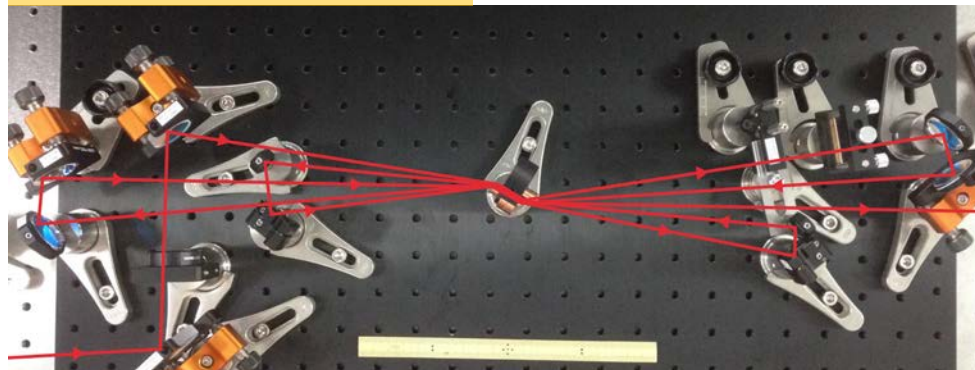
Nd:YAG

355 nm (THG)

ECDL, LBO (OPG)



Ti:Sapphire (OPA)



✓ Experiment will start soon!

Summary

para-H₂ experiment

- Rate amplification of two-photon emission process
- observed TPE signal and verified rate amplification mechanism experimentally
- further study ongoing (counter-propagating solid experiment)

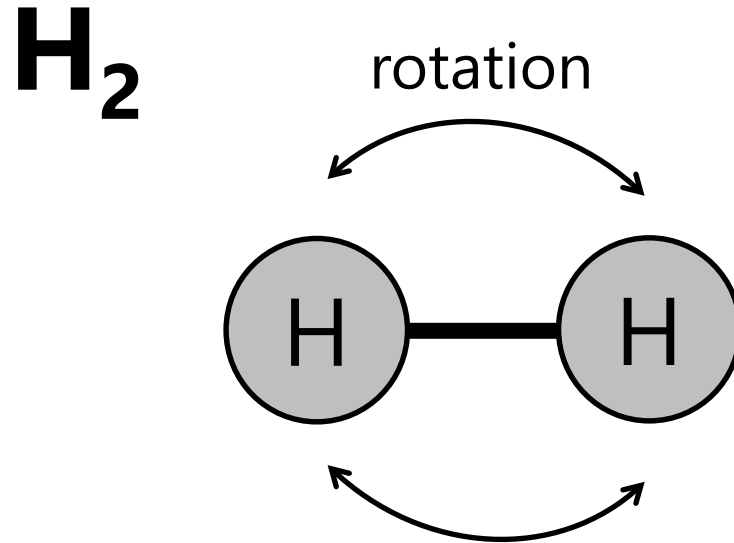
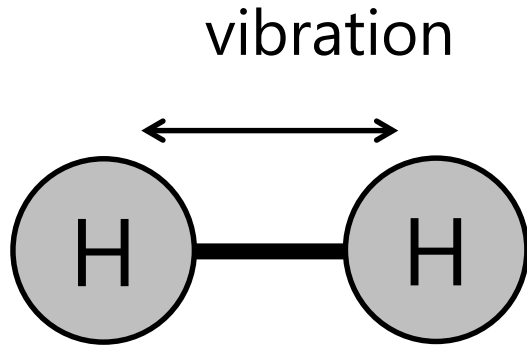
Xe experiment

- coherent amplification of higher-order QED processes
- Experiment will start soon

Future prospects

- Background study and reduction (2 or 3 photon emission)
- obtain higher emission rate
- RENP experiment

Back up



- Vibration or rotation of H_2 molecule are quantized

✓ wavefunction of H_2

$$\psi = \psi_{electron} \psi_{vibration} \psi_{rotation} \psi_{nuclearspin}$$

↑
antisymmetric
under nucleus
exchange

↑ ↑
symmetric @
ground state

	$\psi_{nuclearspin}$	$\psi_{rotation}$
ortho- H_2	symmetric	antisymmetric
para- H_2	antisymmetric	symmetric

Coherent amplification condition

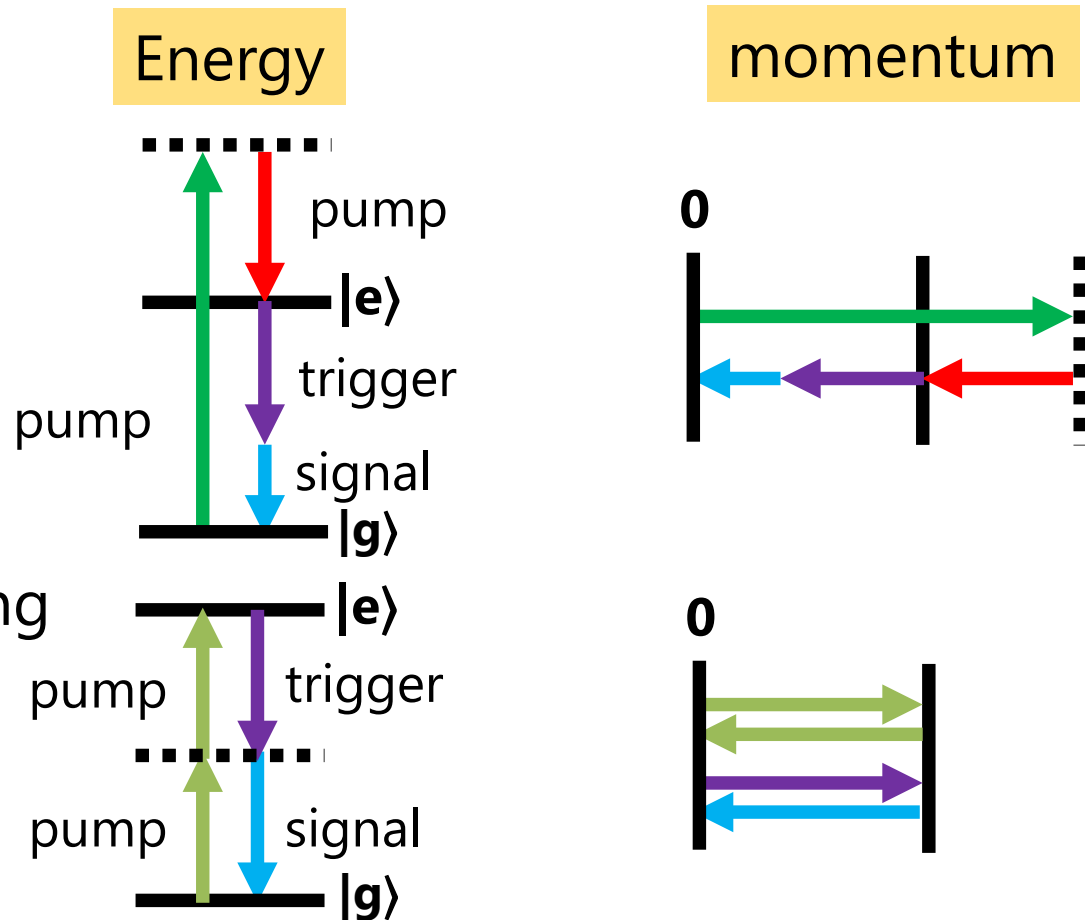
- ✓ Energy-momentum conservation among photons
- ✓ Process: **Two-photon emission (TPE)**

no dispersion case

- one-side excitation



- counter-propagating excitation



The condition is satisfied in both cases

Coherent amplification condition

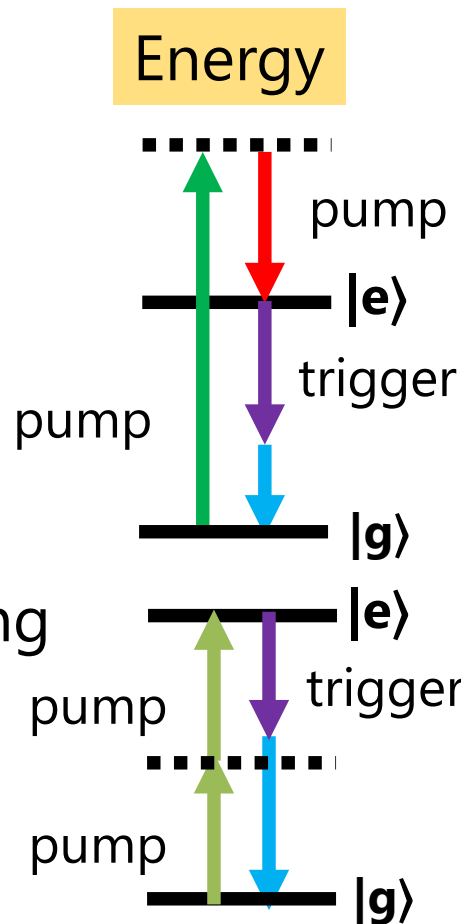
- ✓ Energy-momentum conservation among photons+ ν
- ✓ Process: Radiative emission of neutrino pair (RENAP)

no dispersion case

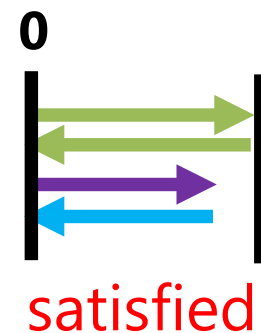
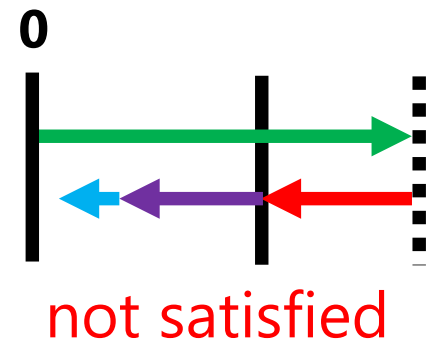
- one-side excitation



- counter-propagating excitation



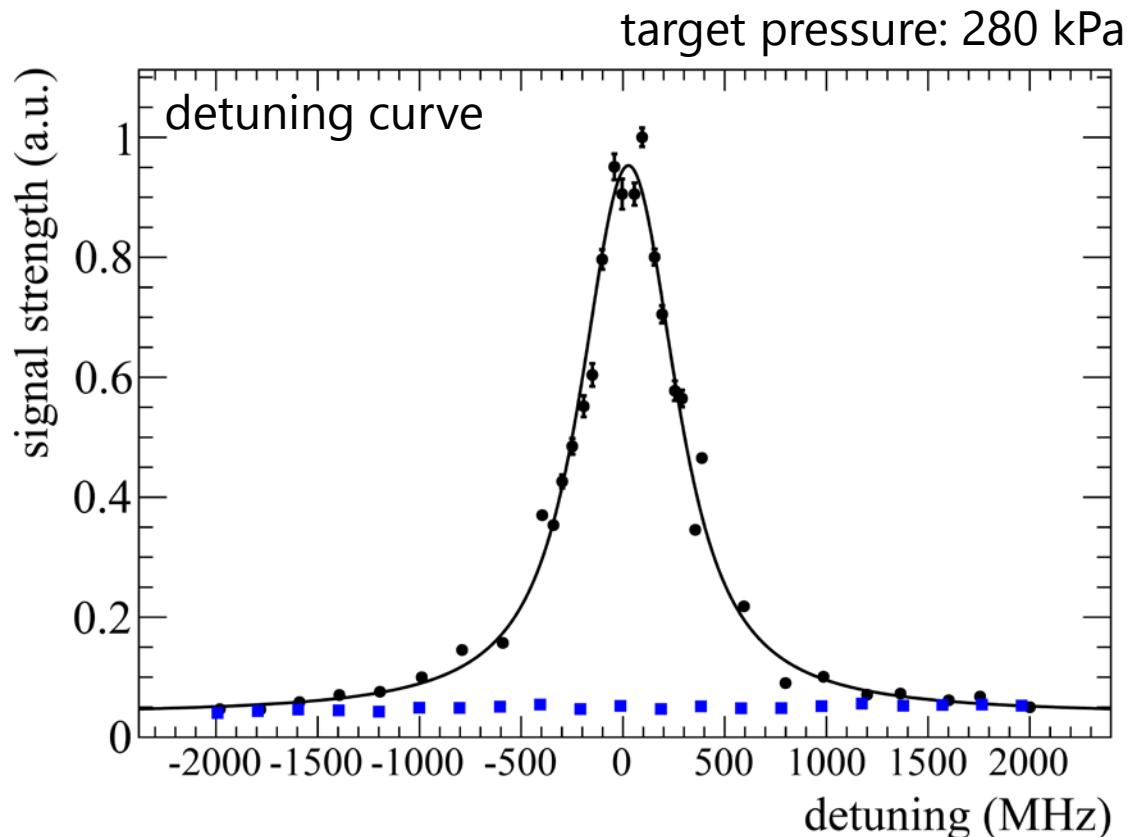
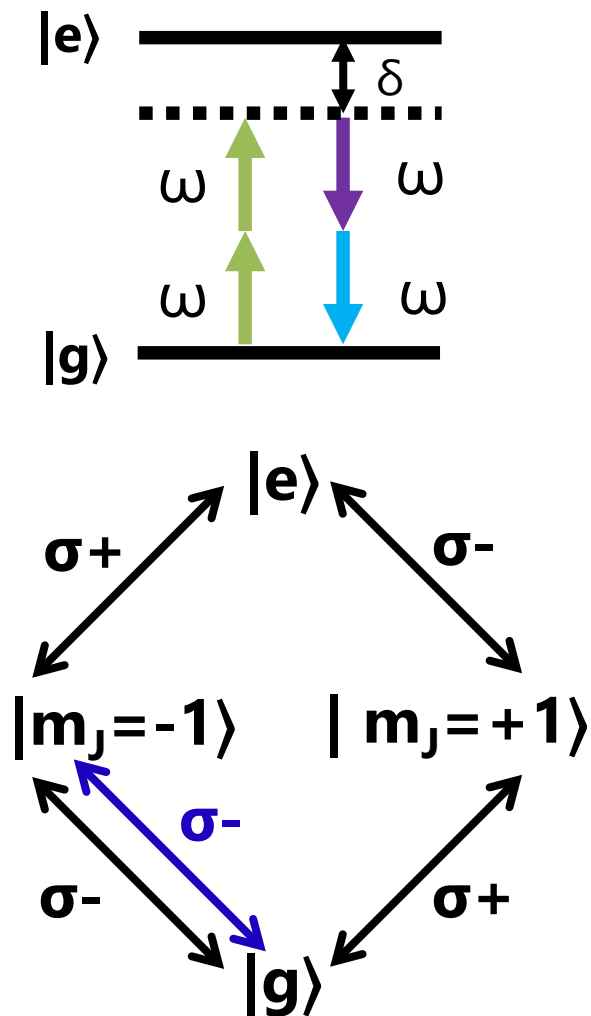
momentum



- ✓ High-quality mid-infrared (4806 nm) laser is required.

Results: detuning dependence

- vary the detuning δ

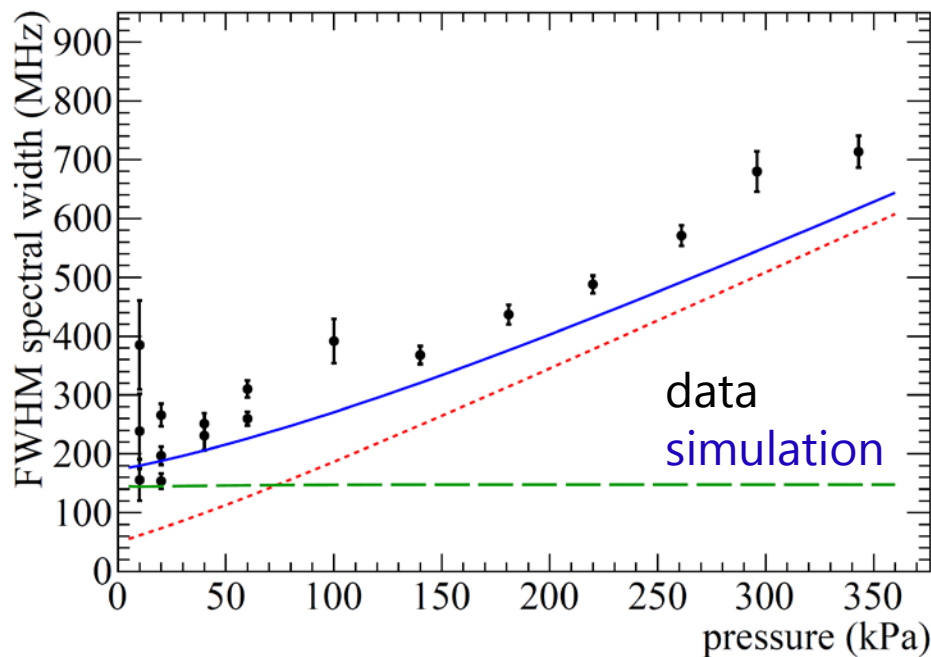


- Circular polarization of a pump is flipped.
- ➔ no signal is observed.
- confirmation of the excitation by counter-propagating lasers

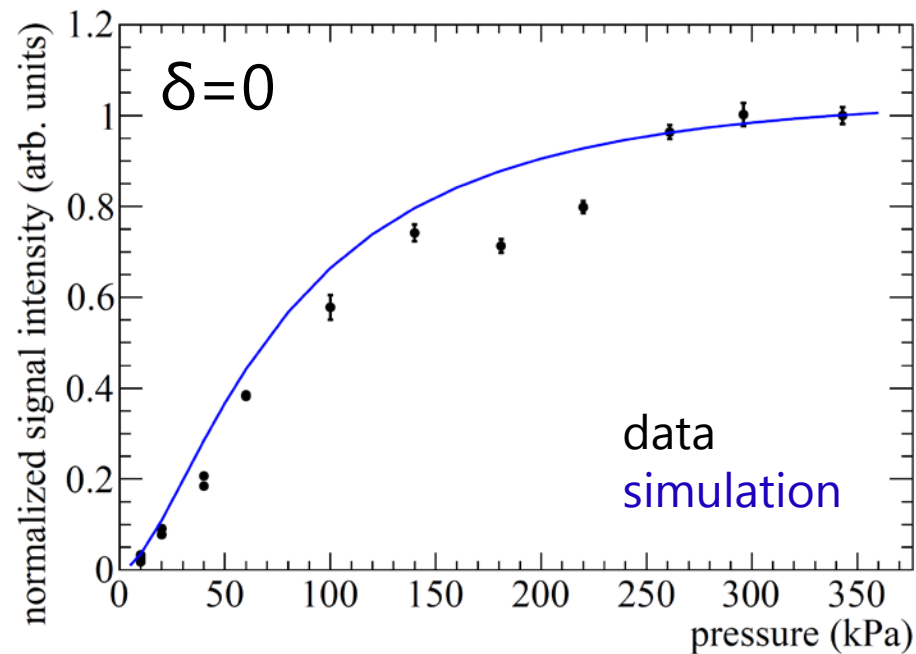
Results: Pressure dependence

- vary the pH_2 target pressure

detuning curve
width (FWHM)



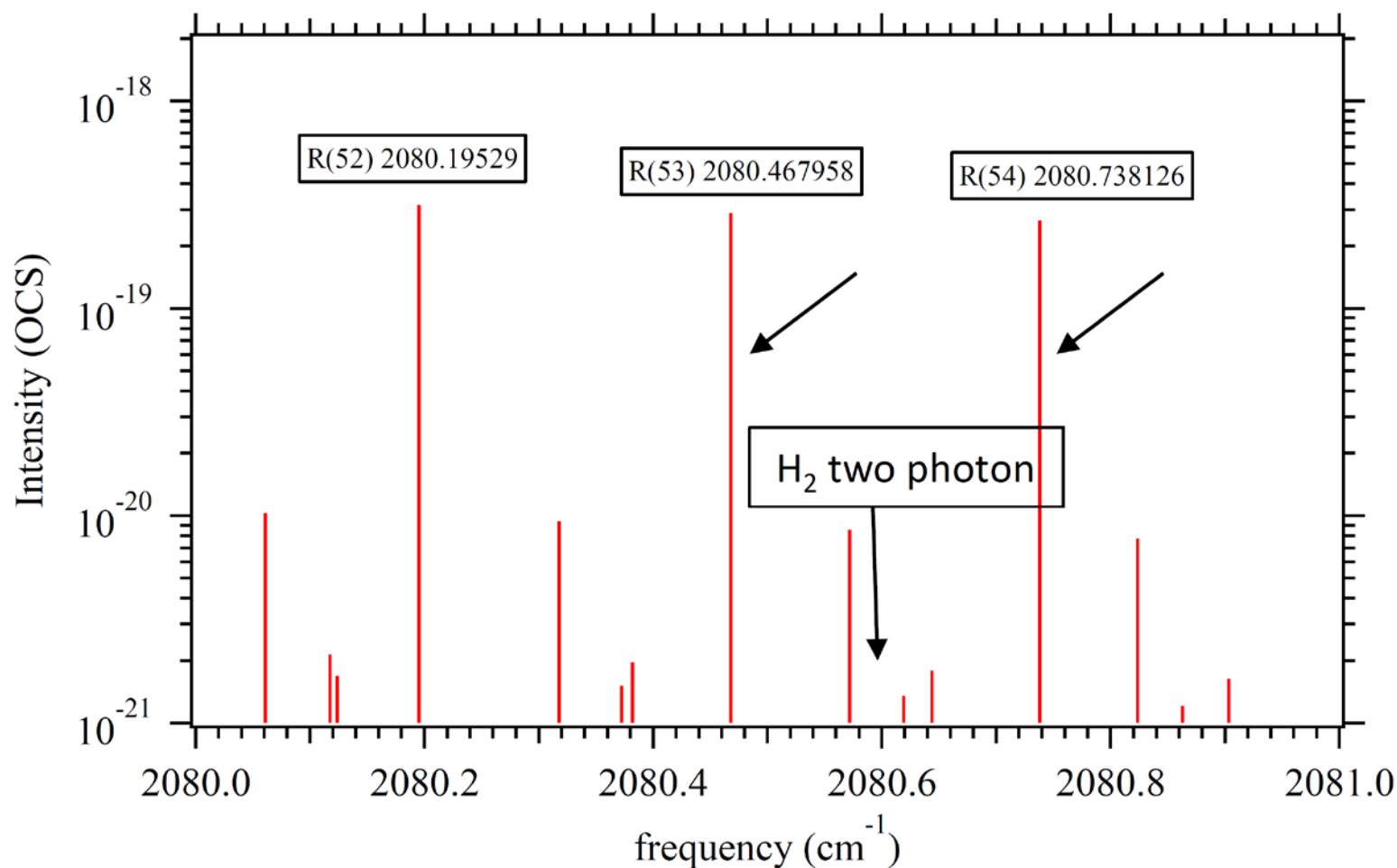
normalized intensity



- Laser linewidth and pressure broadening determine the width
- Signal intensity increases as the target density larger.
- ✓ Consistent tendency is obtained between data and simulation.

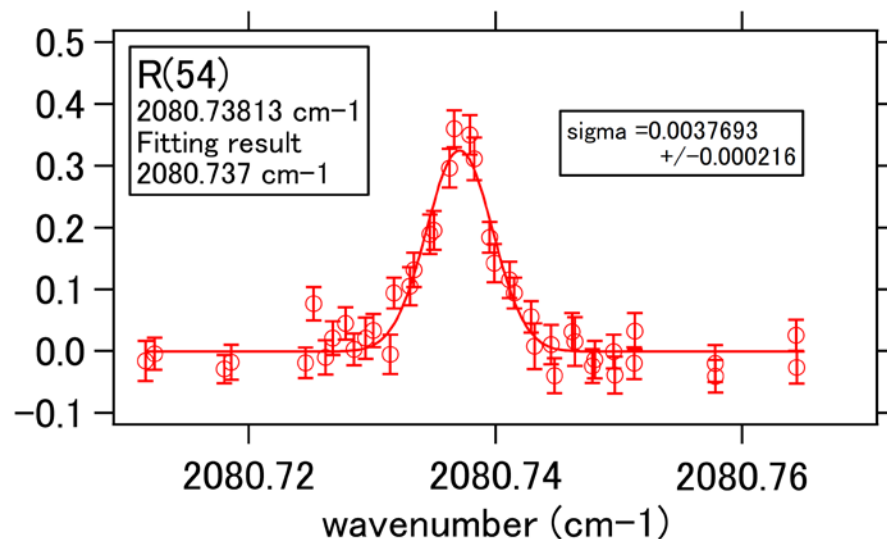
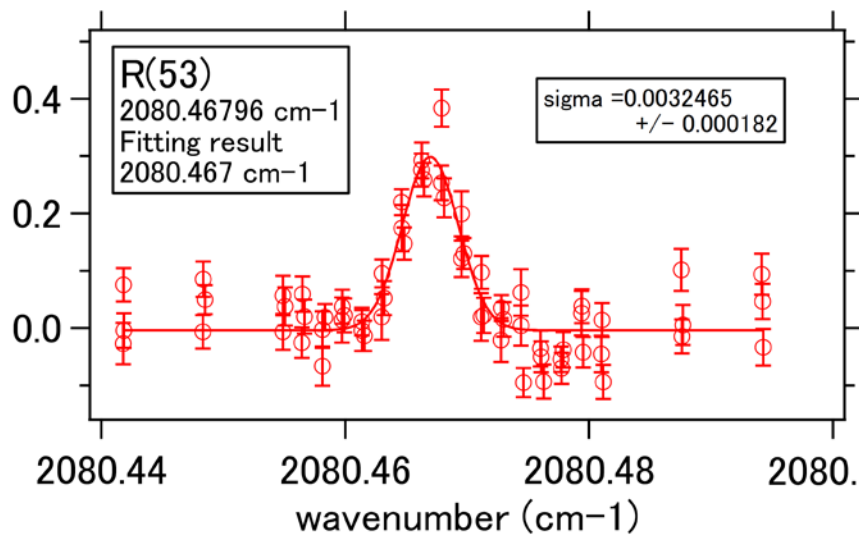
Laser linewidth measurement

- measurement of the narrow-linewidth MIR laser
- method: absorption spectroscopy of carbonyl sulfide (OCS)



Laser linewidth

- observed absorption spectra



	width (FWHM)
Observed linewidth	175 (13)
Doppler width	99
MIR Laser linewidth	145 (16)

✓ narrow laser linewidth ($\sim 1.6 \times \text{FT-limit}$) is achieved.

Maxwell-Bloch equations

Development of the density matrix

$$\frac{\partial \rho_{gg}}{\partial t} = i(\Omega_{ge}\rho_{eg} - \Omega_{eg}\rho_{ge}) + \gamma_1\rho_{ee},$$

$$\frac{\partial \rho_{ee}}{\partial t} = i(\Omega_{eg}\rho_{ge} - \Omega_{ge}\rho_{eg}) - \gamma_1\rho_{ee},$$

$$\frac{\partial \rho_{ge}}{\partial t} = i(\Omega_{gg} - \Omega_{ee} + \delta)\rho_{ge} + i\Omega_{ge}(\rho_{ee} - \rho_{gg}) - \gamma_2\rho_{ge}.$$

ρ : density matrix

$\Omega_{gg(ee)}$: two-photon
Rabi frequency

$\Omega_{eg(ge)}$: AC Stark shift

γ_1, γ_2 : relaxation rates

δ : detuning

Development of the electric fields

$$\left(\frac{\partial}{\partial t} - c\frac{\partial}{\partial z}\right) E_{p1} = \frac{i\omega_l N_t}{2} \left((\alpha_{gg}\rho_{gg} + \alpha_{ee}\rho_{ee}) E_{p1} + 2\alpha_{eg}\rho_{eg} E_{p2}^* \right),$$

$$\left(\frac{\partial}{\partial t} + c\frac{\partial}{\partial z}\right) E_{p2} = \frac{i\omega_l N_t}{2} \left((\alpha_{gg}\rho_{gg} + \alpha_{ee}\rho_{ee}) E_{p2} + 2\alpha_{eg}\rho_{eg} E_{p1}^* \right),$$

$$\left(\frac{\partial}{\partial t} - c\frac{\partial}{\partial z}\right) E_{\text{trig}} = \frac{i\omega_l N_t}{2} \left((\alpha_{gg}\rho_{gg} + \alpha_{ee}\rho_{ee}) E_{\text{trig}} + 2\alpha_{eg}\rho_{eg} E_{\text{sig}}^* \right),$$

$$\left(\frac{\partial}{\partial t} + c\frac{\partial}{\partial z}\right) E_{\text{sig}} = \frac{i\omega_l N_t}{2} \left((\alpha_{gg}\rho_{gg} + \alpha_{ee}\rho_{ee}) E_{\text{sig}} + 2\alpha_{eg}\rho_{eg} E_{\text{trig}}^* \right).$$

ω_l : laser
frequency

N_t : target density

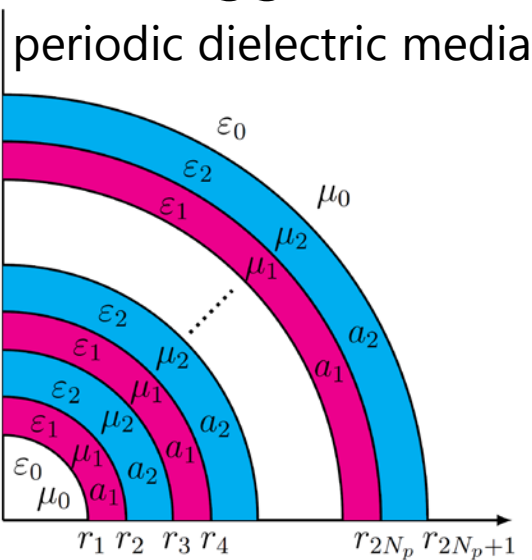
α : polarizability

Theoretical studies

- Towards Background-free RENP using a Photonic Crystal Waveguide

M.Tanaka et al. Prog. Theor. Exp. Phys. **2017** 043B03

Bragg fiber



lasers

3-photon emission suppression factor

