

# Heavy neutral lepton searches at the LHC

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*on behalf of the  
ATLAS and CMS collaborations*



**COEPP**  
ARC Centre of Excellence for  
Particle Physics at the Terascale



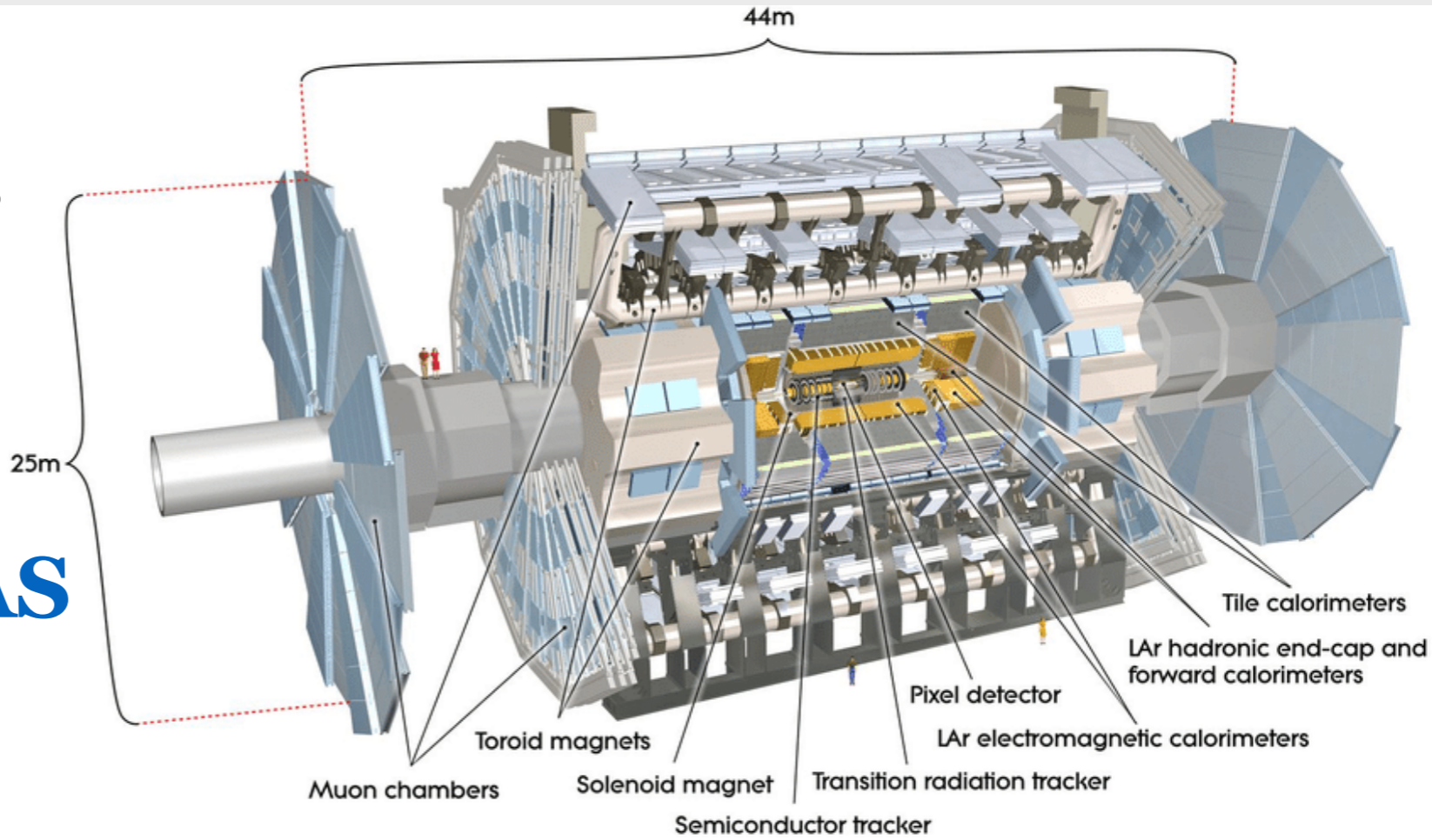
THE UNIVERSITY OF  
MELBOURNE

# ATLAS and CMS

- Multi-purpose experiments surrounding the interaction regions where energetically symmetric proton beams collide.
- pp collisions from LHC at  $\sqrt{s} = 13$  TeV in Run-II.

**CMS**

**ATLAS**



**CMS DETECTOR**

Total weight : 14,000 tonnes  
Overall diameter : 15.0 m  
Overall length : 28.7 m  
Magnetic field : 3.8 T

STEEL RETURN YOKE  
12,500 tonnes

SILICON TRACKERS  
Pixel (100x150  $\mu\text{m}$ )  $\sim 16\text{m}^2 \sim 66\text{M}$  channels  
Microstrips (80x180  $\mu\text{m}$ )  $\sim 200\text{m}^2 \sim 9.6\text{M}$  channels

SUPERCONDUCTING SOLENOID  
Niobium titanium coil carrying  $\sim 18,000\text{A}$

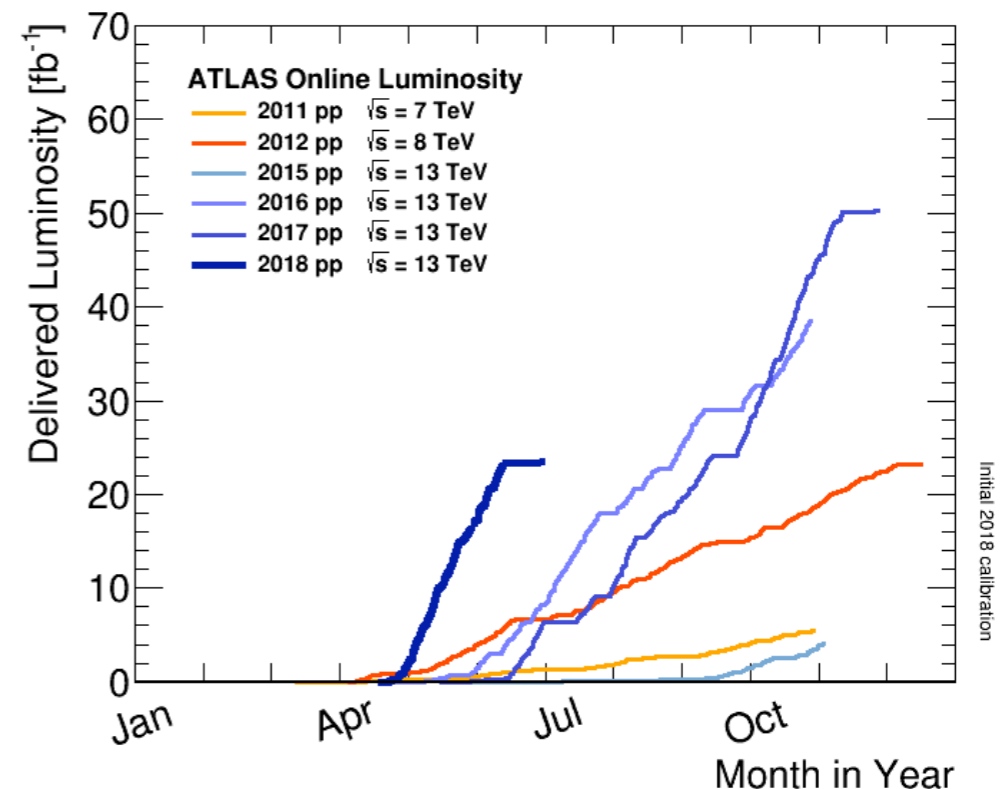
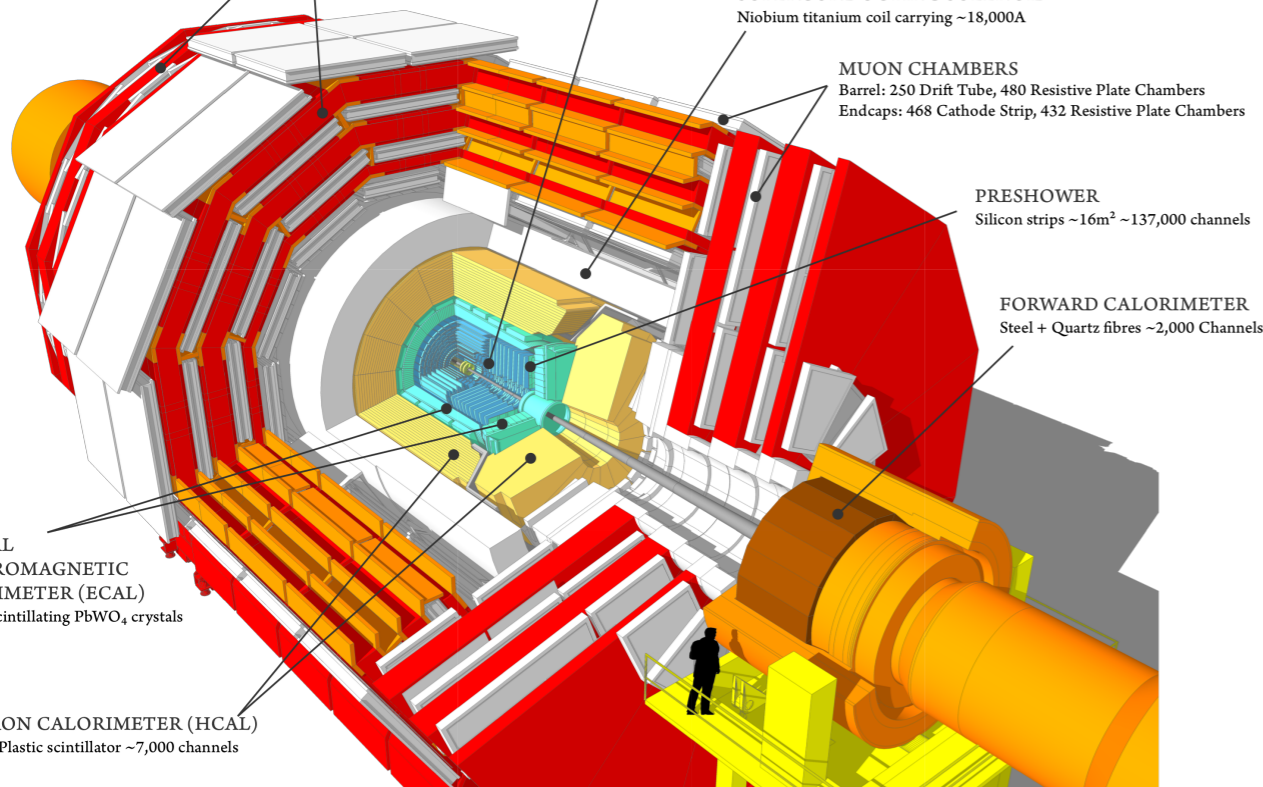
MUON CHAMBERS  
Barrel: 250 Drift Tube, 480 Resistive Plate Chambers  
Endcaps: 468 Cathode Strip, 432 Resistive Plate Chambers

PRESHOWER  
Silicon strips  $\sim 16\text{m}^2 \sim 137,000$  channels

FORWARD CALORIMETER  
Steel + Quartz fibres  $\sim 2,000$  Channels

CRYSTAL ELECTROMAGNETIC CALORIMETER (ECAL)  
 $\sim 76,000$  scintillating PbWO<sub>4</sub> crystals

HADRON CALORIMETER (HCAL)  
Brass + Plastic scintillator  $\sim 7,000$  channels



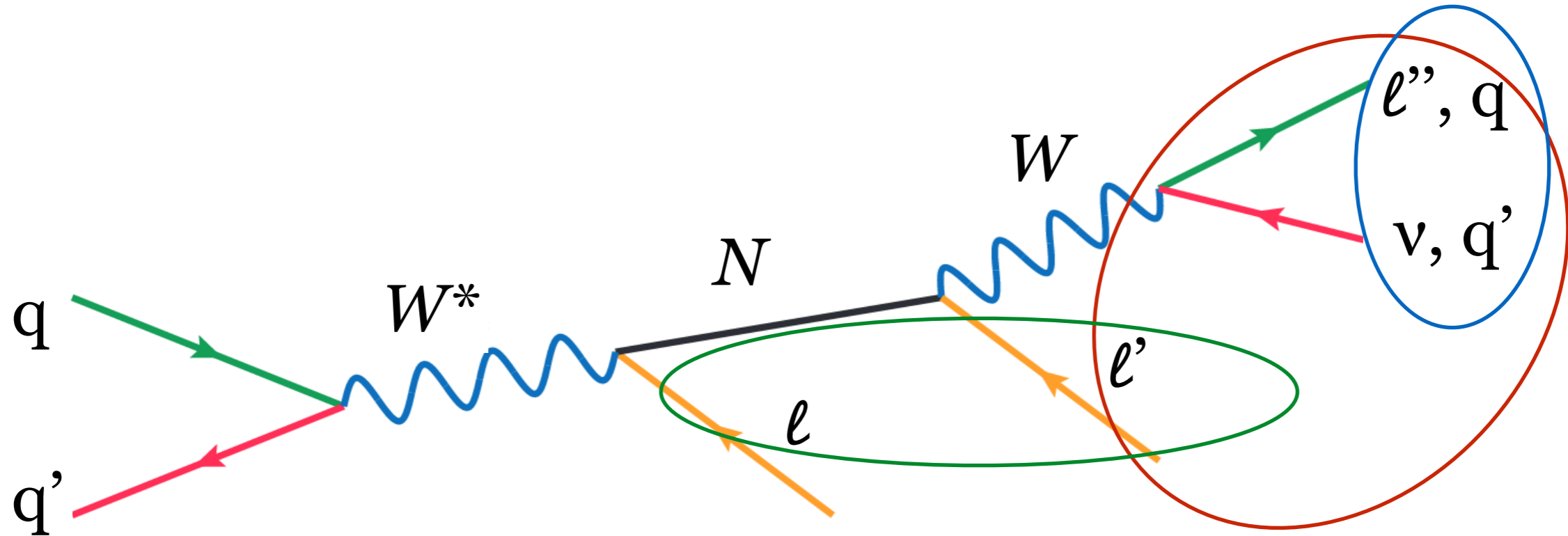
# Theory overview: See-Saw models

- Neutrino masses can be accommodated by EWSB, but a  $\nu_R$  singlet is trivial under the SM gauge group.
- If EWSB holds then Yukawa couplings must be much smaller than for other SM particles,  $\lambda_\nu \approx 10^{-12}$ .
- Since  $\nu_R$  is neutral, Majorana mass terms are possible and decoupled from SM EWSB.
- **See-saw mechanisms** based on introducing Majorana masses  $m_M$  in addition to Dirac ones  $m_D$ .

$$\mathbf{M} = \begin{pmatrix} 0 & m_D \\ m_D & m_M \end{pmatrix} \longrightarrow m_D \ll m_M \longrightarrow m_{\nu, N} = \frac{m_M}{2} \mp \sqrt{\left(\frac{m_M}{2}\right)^2 + m_D^2}$$

- **Type-I**: SU(2) singlet fermion  $N$ .
- **Type-II**: SU(2) triplet scalar  $\Phi^{0, \pm, \pm\pm}$ .
- **Type-III**: SU(2) triplet fermion  $\Sigma^{0, \pm}$ .

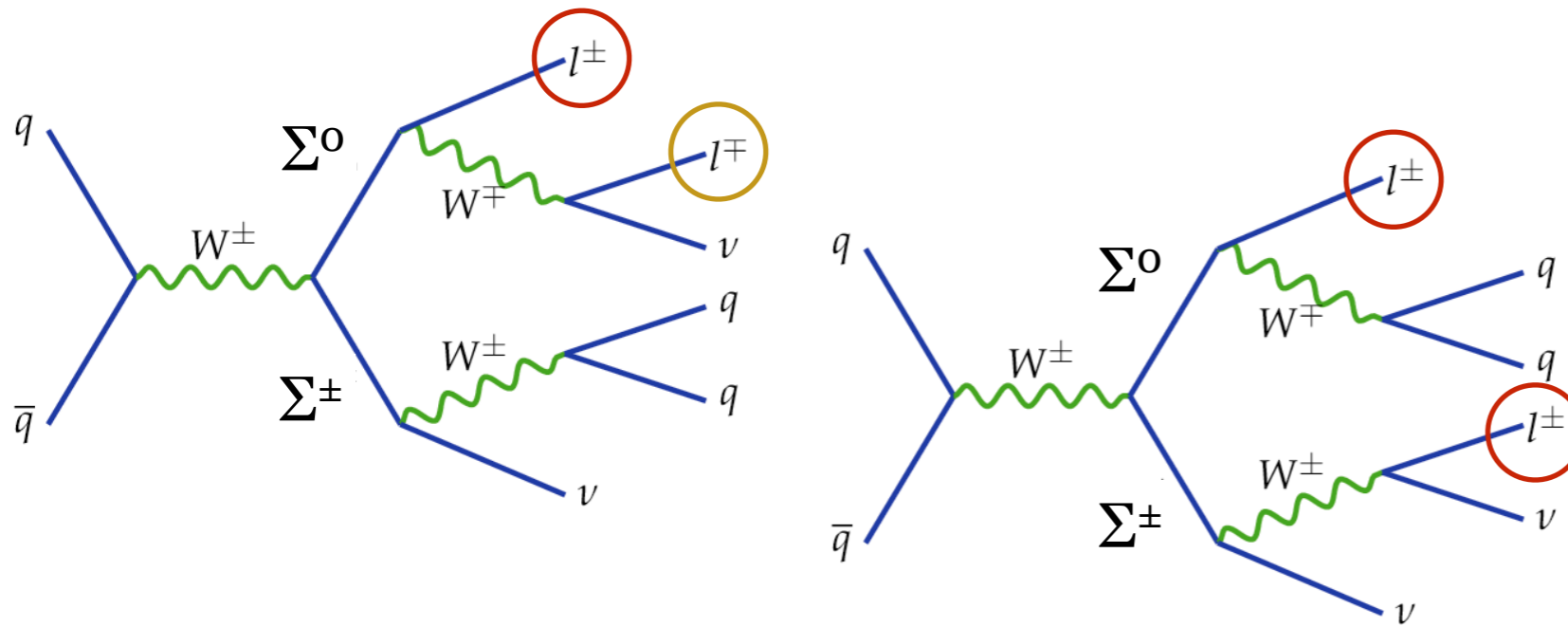
# Type-I signatures



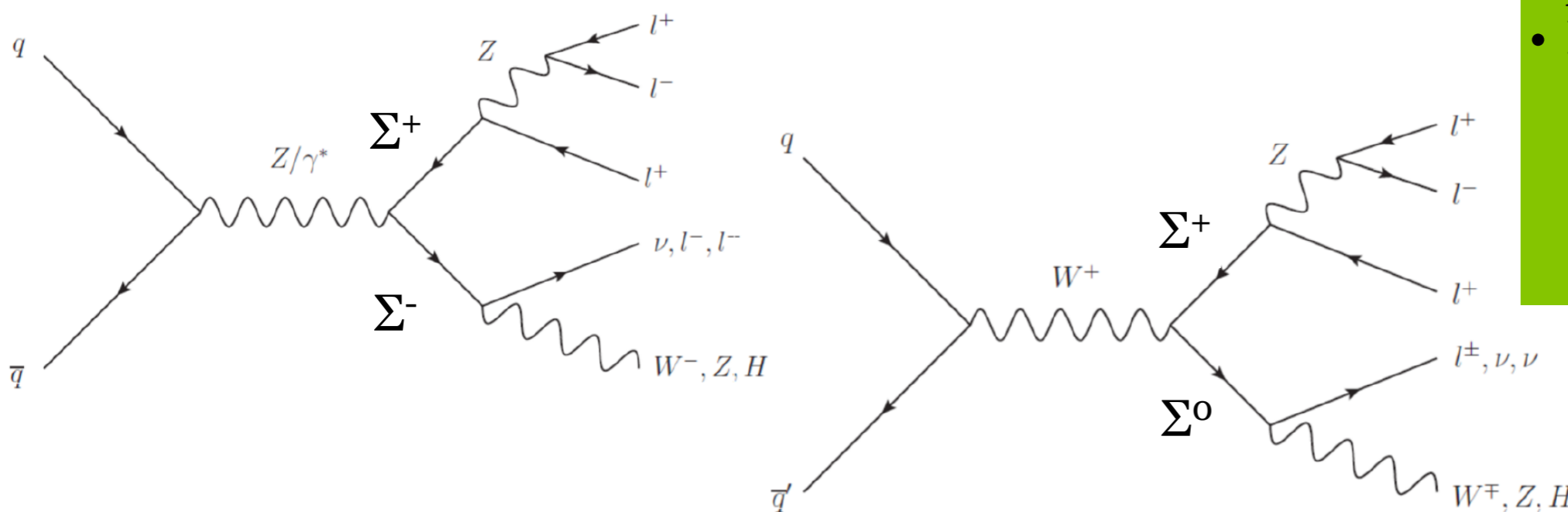
- In minimal Type-I, the heavy neutrino parameters,  $\mathbf{m}_N$  and mixing matrix elements  $|\mathbf{V}_{\nu N}|^2$  are free.
- Heavy neutrino produced via mixings with SM neutrinos.
- Search strategies based on  $m_N$  vs  $m_W$  hypothesis.
- $W^*$  is **on-shell** at very low  $m_N$ . **Off-shell** otherwise. Also:
  - $\mathbf{m}_N \ll \mathbf{m}_W$ : soft and displaced  $N$  decay products.
  - $\mathbf{m}_N < / \approx / > \mathbf{m}_W$ : hierarchies of  $p_T(\ell)$  vs  $p_T(\ell')$ . E.g.  $\ell'$  dominates at high  $\mathbf{m}_N$ .
  - $\mathbf{m}_N \gg \mathbf{m}_W$ : boosted decay products (jets).

# Type-III signatures

- Mass degenerate  $\Sigma^0$  and  $\Sigma^\pm$  due to gauge invariance. Only one free parameter.
- Production of  $\Sigma$  via gauge interaction.



- ATLAS**
- $pp \rightarrow \Sigma^0 \Sigma^\pm \rightarrow llqq$
  - Two leptons in final state with same or opposite charge.
  - Other decay modes found negligible in final selection.

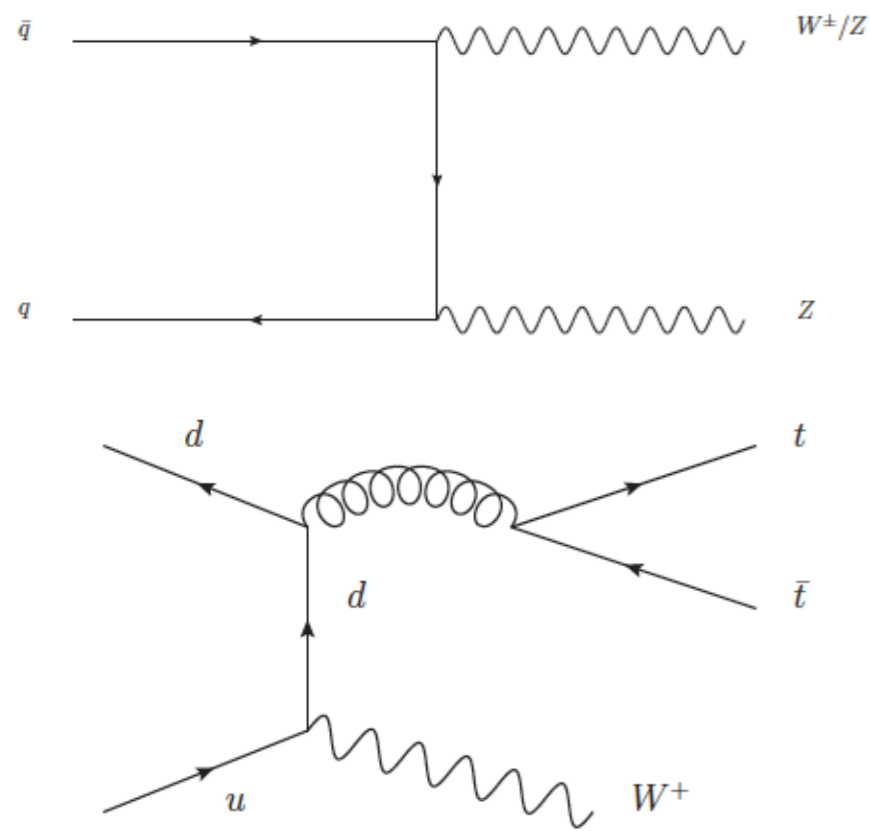


- CMS**
- $pp \rightarrow \Sigma^0 \Sigma^\pm / \Sigma^\mp \Sigma^\pm$
  - Purely multi-leptonic
    - $\Sigma^0 \rightarrow W^\mp l^\pm$
    - $\Sigma^0 \rightarrow Z / H \nu_l$
    - $\Sigma^\pm \rightarrow Z / H l^\pm$
    - $\Sigma^\pm \rightarrow W^\pm \nu_l$

# Backgrounds

## Prompt

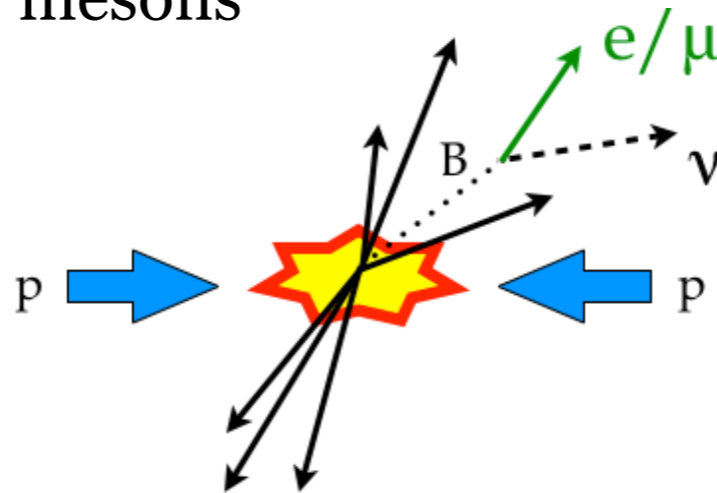
Real prompt leptons:  
 $ZW, ZZ, ttW, ttZ, ttH, W^\pm W^\pm$



Estimated with simulation

## Mis-ID leptons

- Real electrons or muons from non-prompt decays, e.g. from heavy flavoured mesons

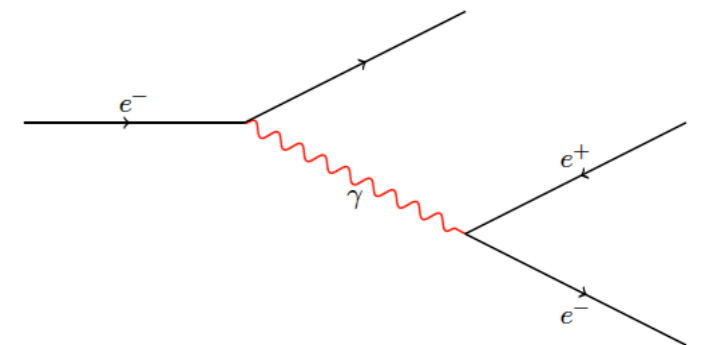


- Jets mis-reconstructed as electrons

Data-driven estimation

## Mis-ID charge

Oppositely charged leptons with charge mis-identification:  
 $Z/\gamma^*, ZZ, W^\pm W^\mp, tt$



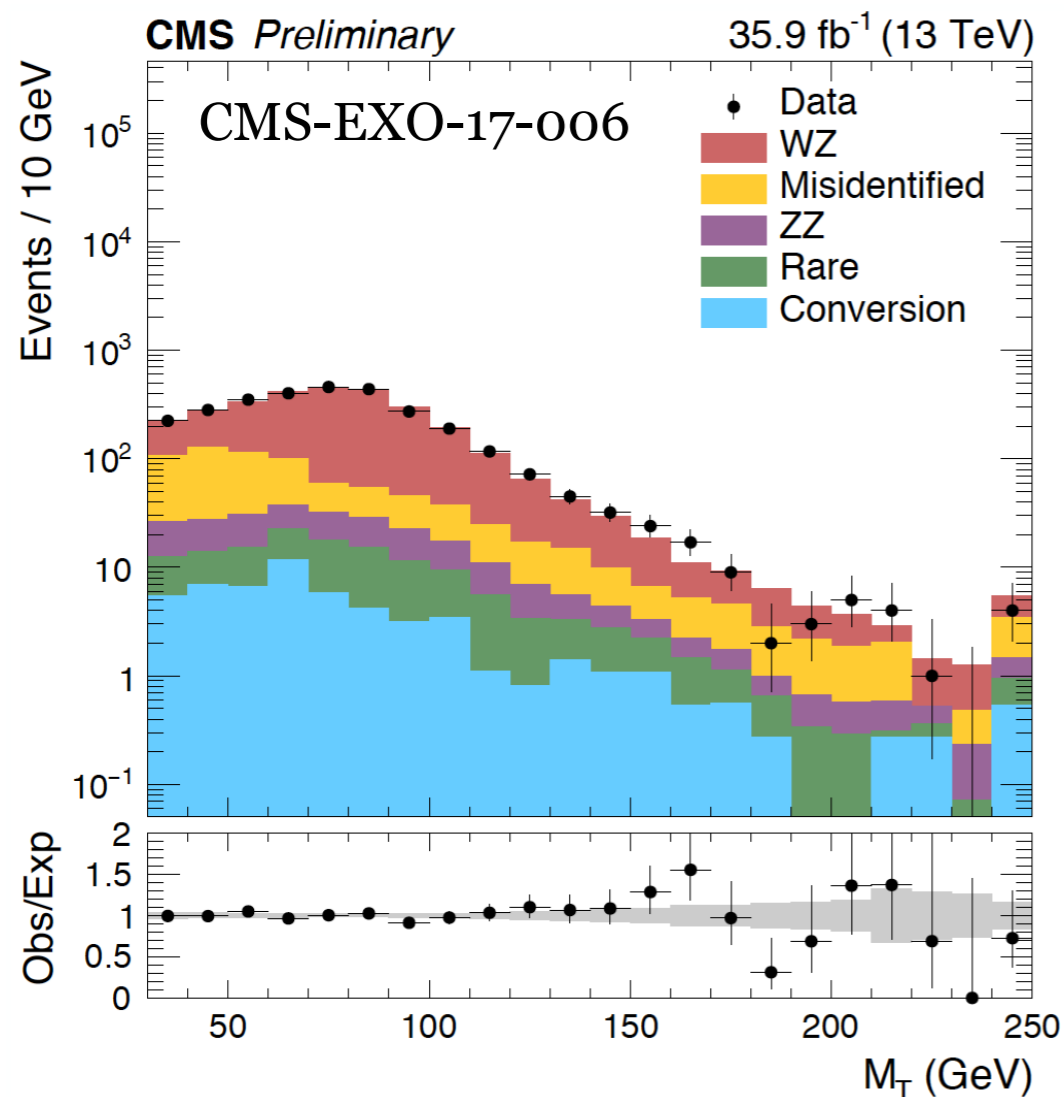
Mis-ID probability estimated from data

# Background estimation

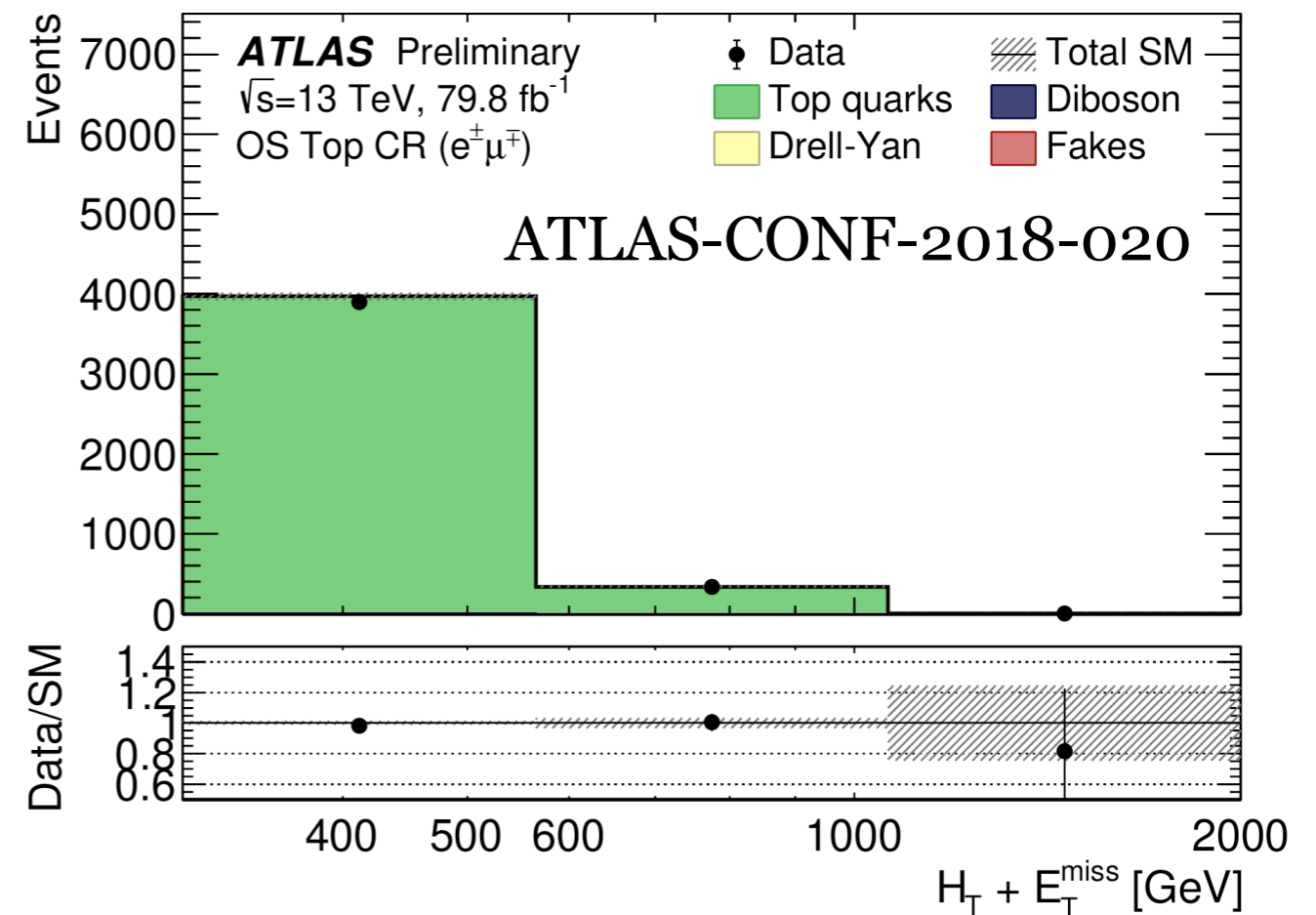
# Prompt leptons

- Prompt contributions are estimated using simulated samples.
- Usually simulated @NLO accuracy
- Major contributions are normalised in dedicated control regions.
- Derive correction factors or include the control regions in the fit.

WZ  $\rightarrow 3\ell\nu$  control region



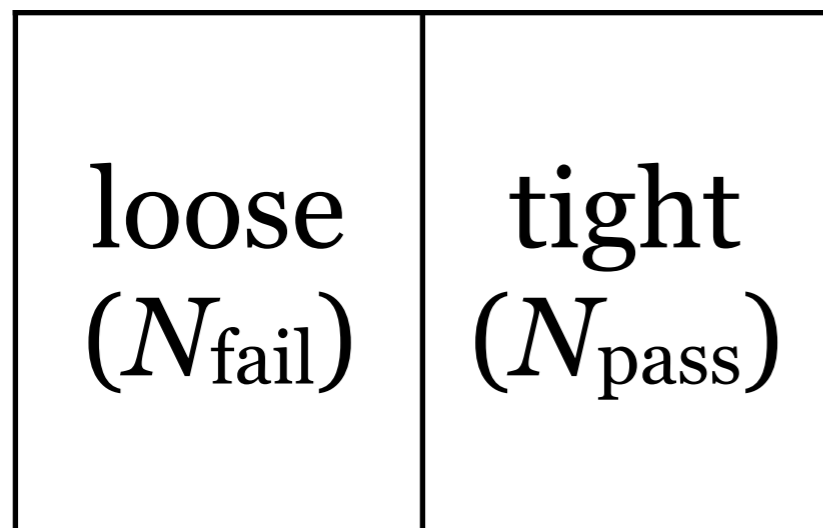
Top control region





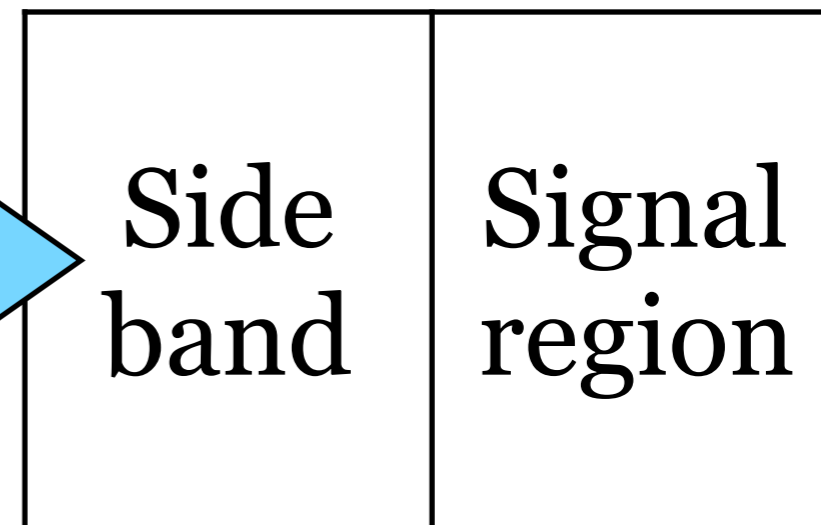
# Mis-identified leptons

Fakes enriched region



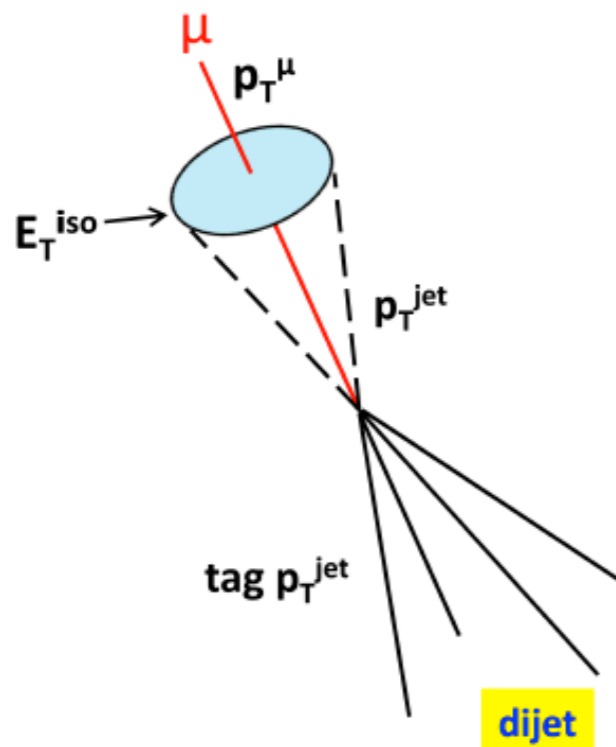
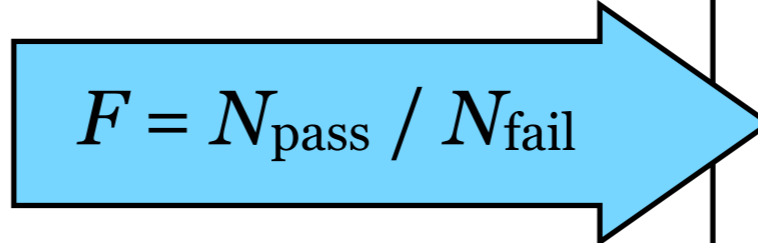
isolation threshold

Main selection



isolation threshold

“fake-factor”



Tag-and-probe  
selections on di-jet  
events

- Signal region extrapolation: e.g. two lepton case:

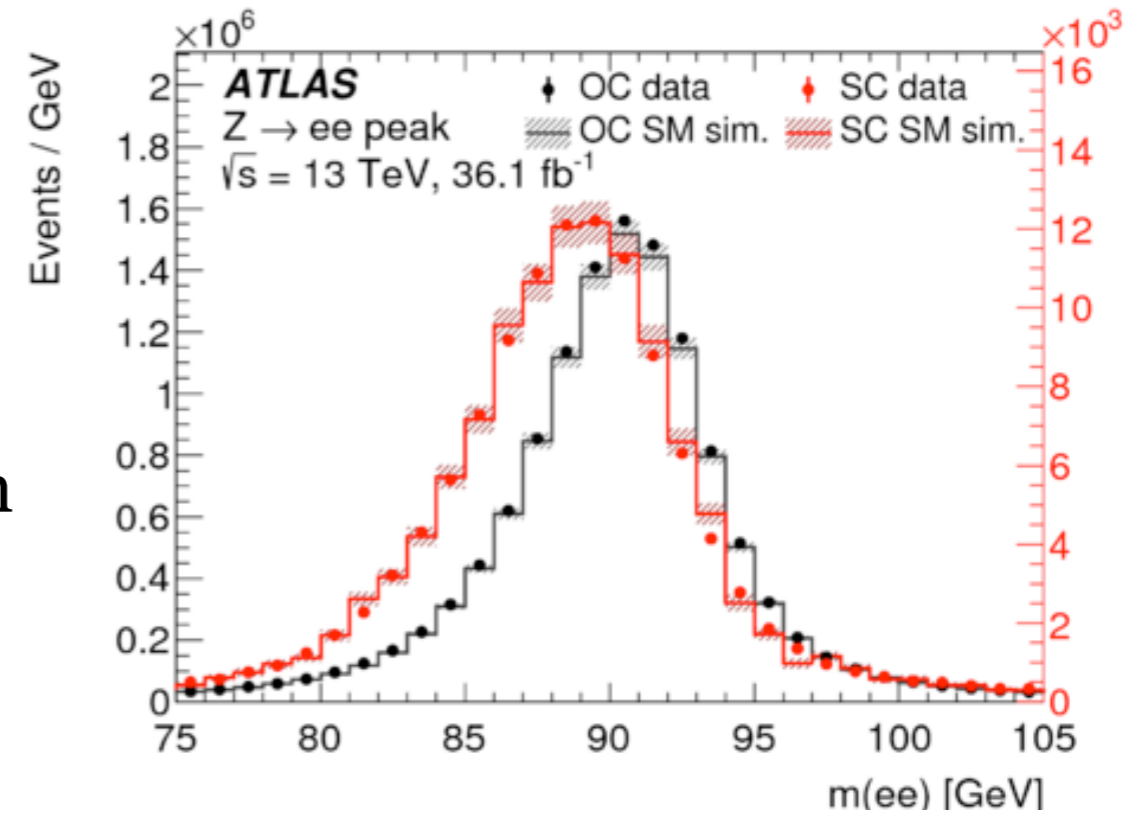
$$N_{TT}^{\text{fakes}} = \left[ \sum_{TL} F_2 + \sum_{LT} F_1 - \sum_{LL} F_1 F_2 \right]_{\text{data}} - \left[ \sum_{TL} F_2 + \sum_{LT} F_1 - \sum_{LL} F_1 F_2 \right]_{\text{prompt simulation}}$$

- Can be extended for more than two leptons
- Parameterisation of  $F$  based on lepton kinematic

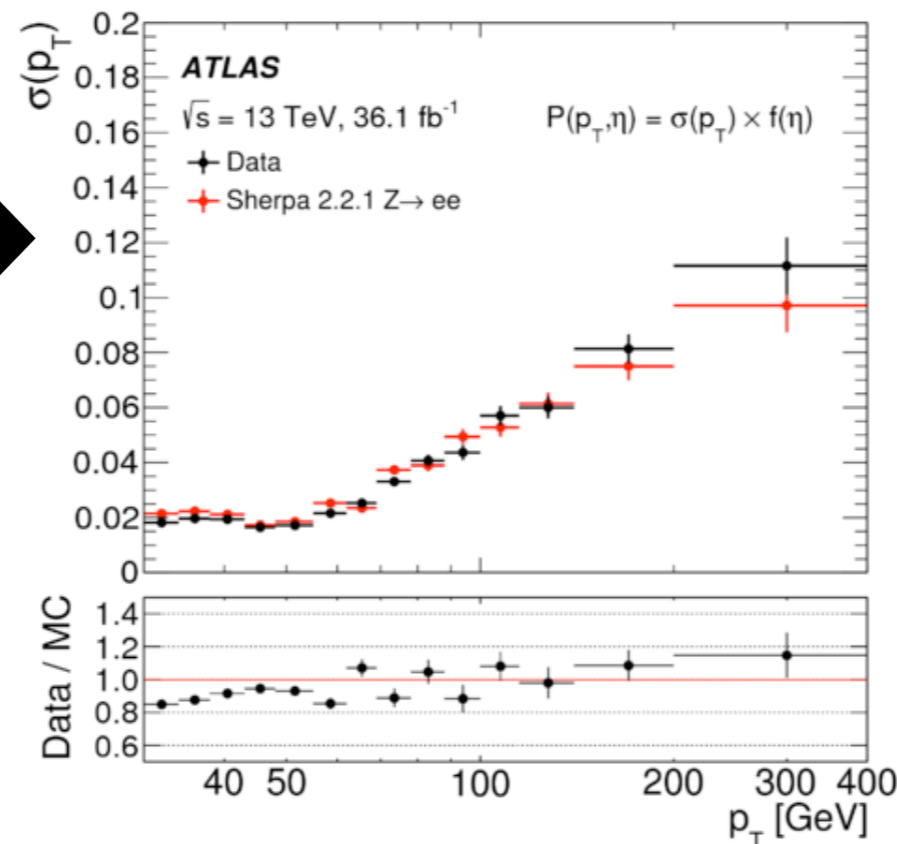
# Mis-identified charge

**ATLAS** example: ATLAS-CONF-2018-020

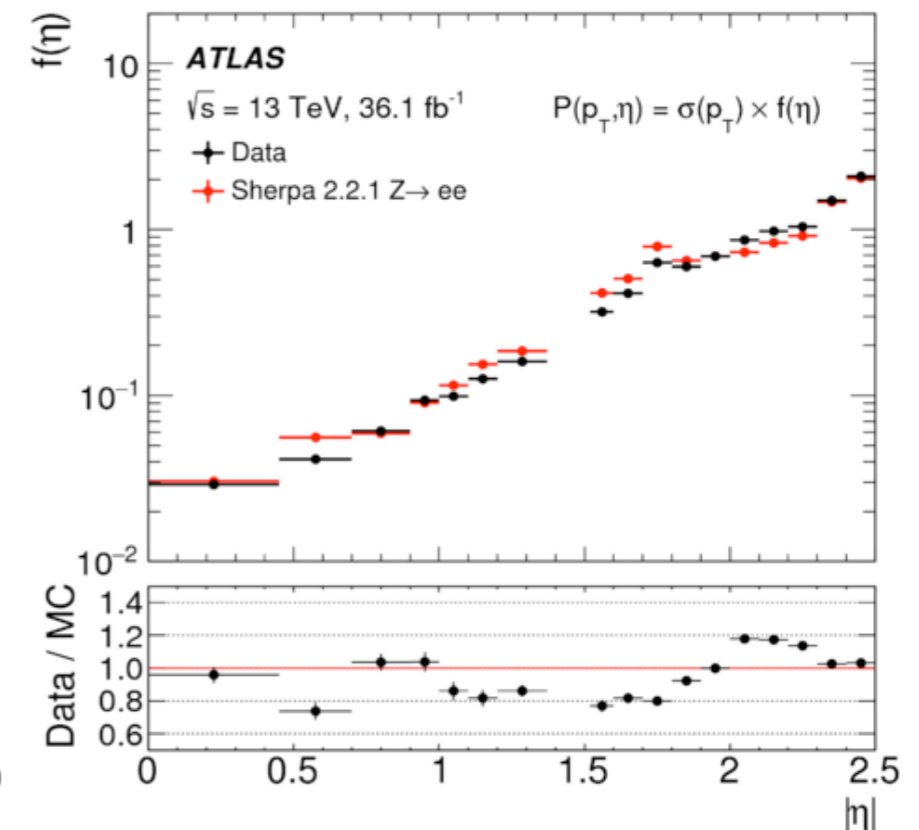
- $Z \rightarrow ee$  events used from data and simulation.
- Fit simultaneously opposite and same -charge events and separately for data and simulation.
- Derive parameterised probabilities and measure a correction based on data/simulation trends.
- Correction is applied to any simulated electron with mis-identified charge.



Mis-id probability:  
 $P(p_T, \eta) = \sigma(p_T) \times f(\eta)$



corrections

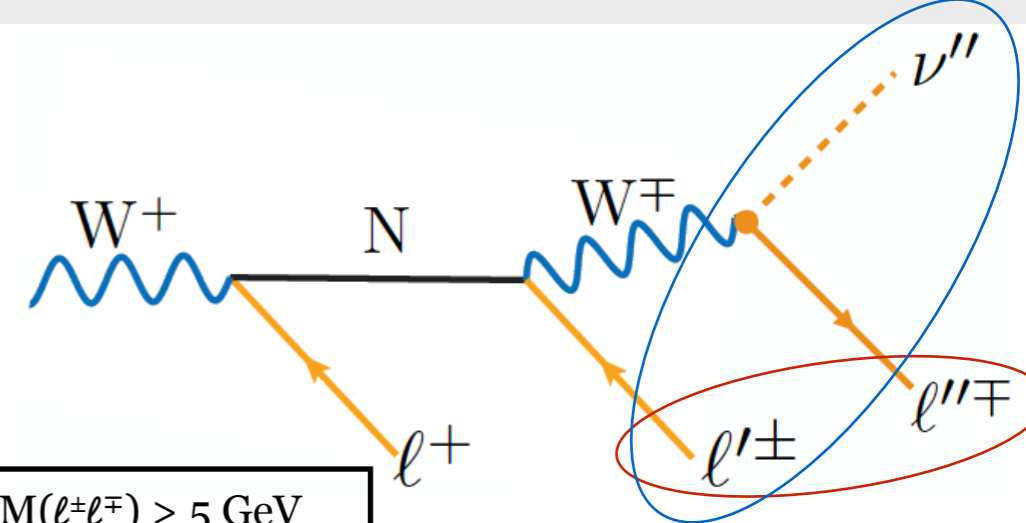


# Type-I searches

# Type-I - trilepton channel

CMS:  
1802.02965

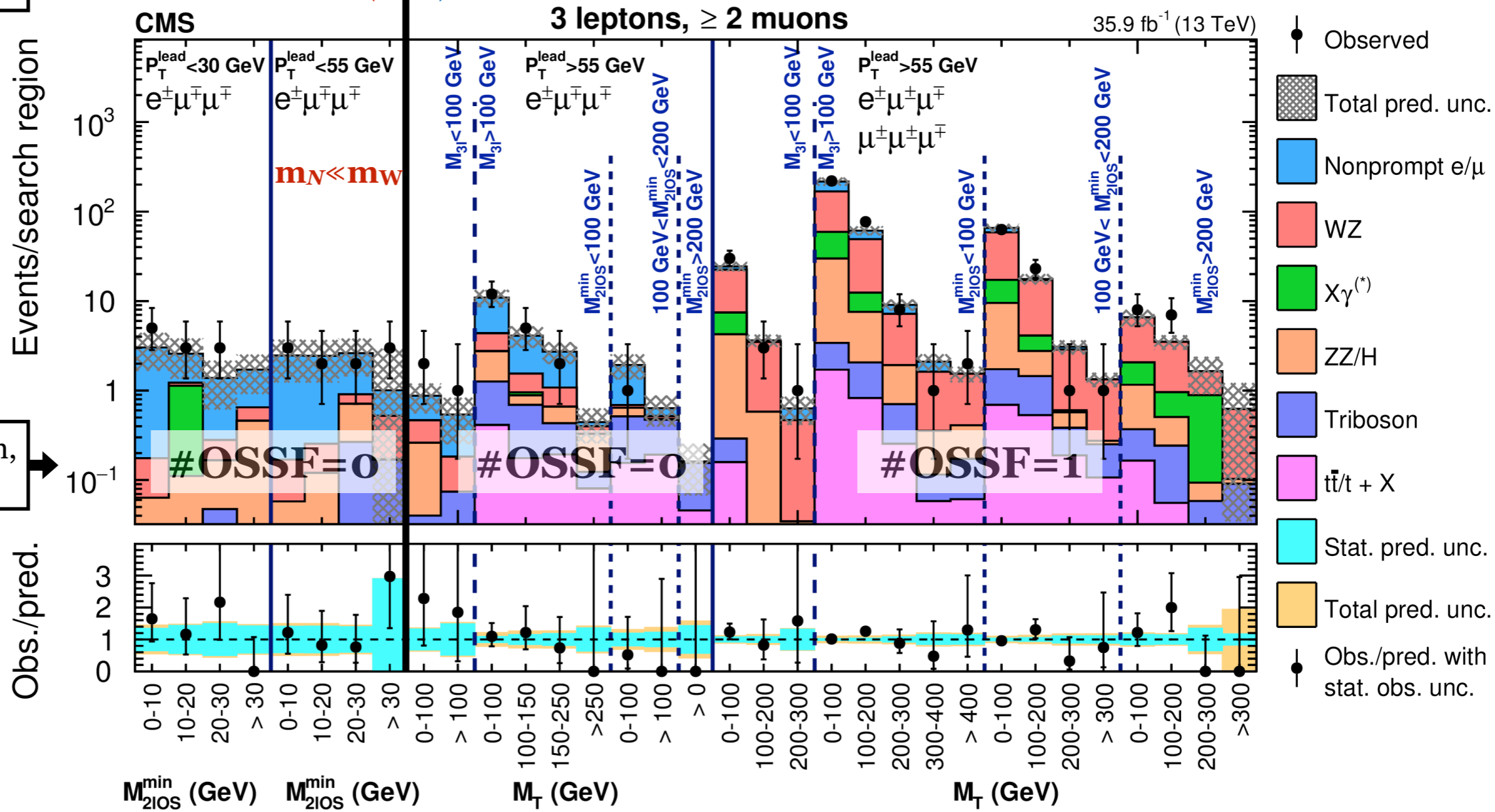
- Any mix of charge and flavour for  $e$  and  $\mu$ .
- No mass peak due to neutrinos in  $W$  decays.
- Veto events with same-sign three leptons,  $b$ -jets or with lepton invariant masses consistent with  $m_Z$ .
- Binned in  $p_T(\ell_{\text{lead}})$ ,  $M(\ell^\pm \ell^\mp \ell^\pm)$ ,  $M_{\text{min}}(\ell^\pm \ell^\mp)$ ,  $M_T(\ell', E_T^{\text{miss}})$ .



$$M(\ell^\pm \ell^\mp \ell^\pm) < 80 \text{ GeV}, \\ E_T^{\text{miss}} < 75 \text{ GeV}$$

$$m_N < m_W \quad \longleftrightarrow \quad m_N > m_W$$

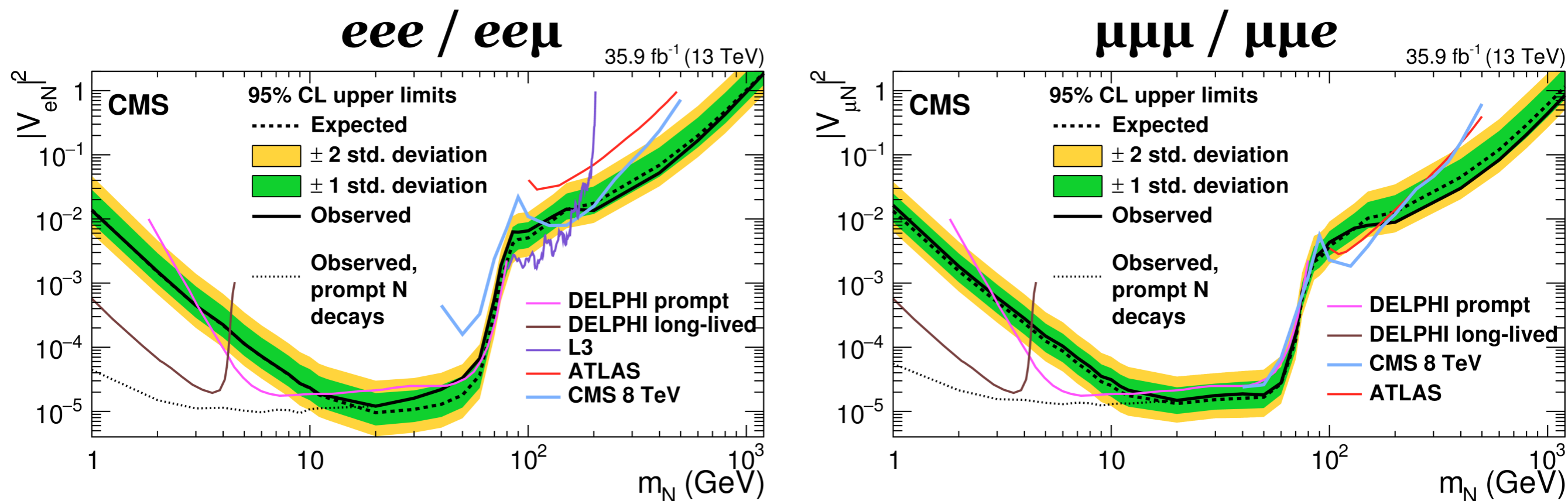
$$M(\ell^\pm \ell^\mp) > 5 \text{ GeV}$$



OSSF: opposite-sign, same-flavour

# Type-I - trilepton channel

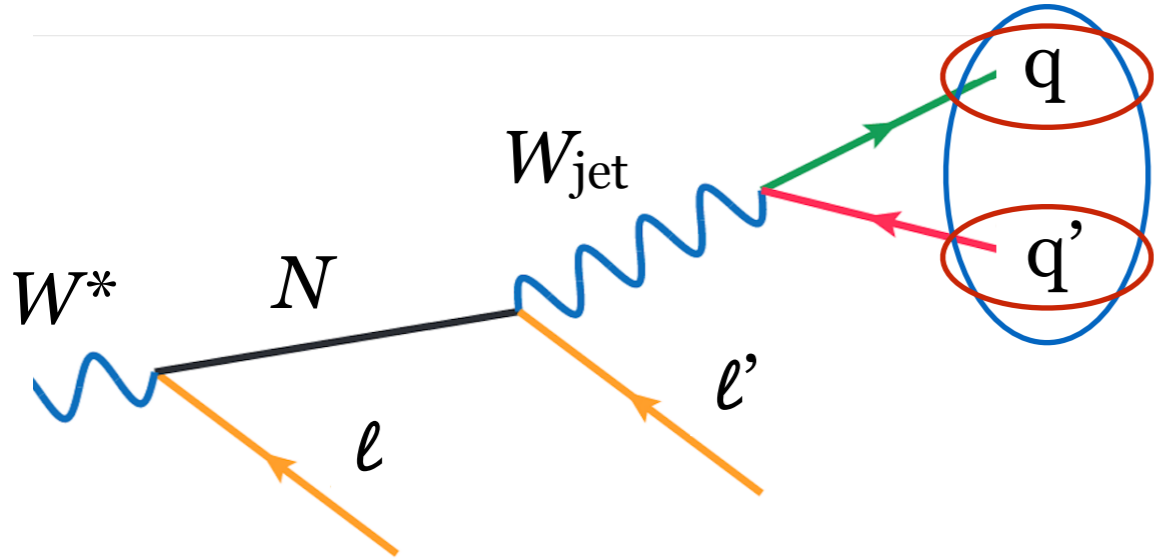
CMS:  
1802.02965



- Elements of the mixing matrix  $|V_{\nu N}|$  are varied independently and are non-zero for a single flavour at a time.
- Displaced  $N$  decays occur for  $m_N \approx 20$  GeV.
  - $\tau_N \sim |V_{\nu N}|^{-2} m_N^{-5}$ .
  - Efficiency correction taken into account at lower masses.

# Type-I - dilepton channel

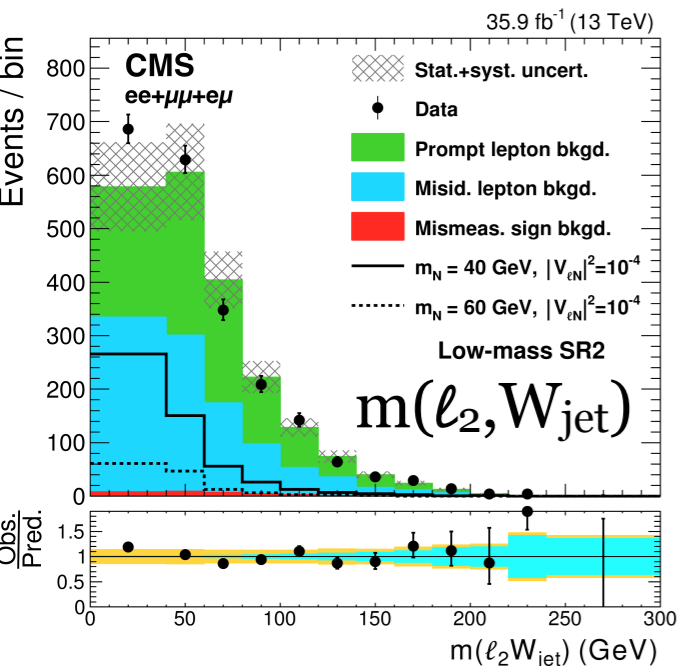
CMS:  
1806.10905



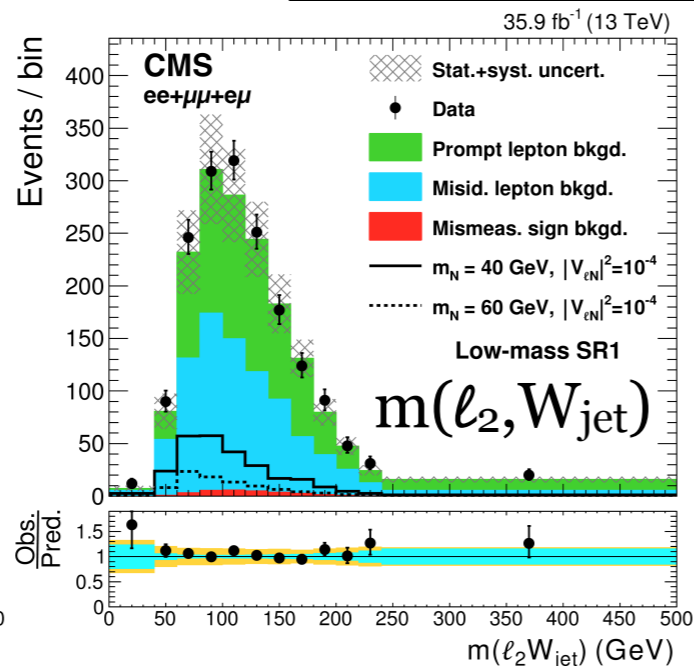
	Hierarchy	N res. jets	N merg. jets
low-mass SR2	$m_N < m_W$	1	0
low-mass SR1	$m_N \approx m_W$	$\geq 2$	0
high-mass SR1	$m_N > m_W$	$\geq 2$	0
high-mass SR2	$m_N \gg m_W$	$\geq 0$	$\geq 1$

- Includes  $gg/qq' \rightarrow N\ell^\pm$  and  $q\gamma \rightarrow Nq'\ell^\pm$
- Same-sign dileptons only.
- $p_T(\ell) \gtrsim 20$  GeV.
- All  $e/\mu$  combinations.
- Optimised using  $W_{jet}$  and  $\ell$  kinematics.

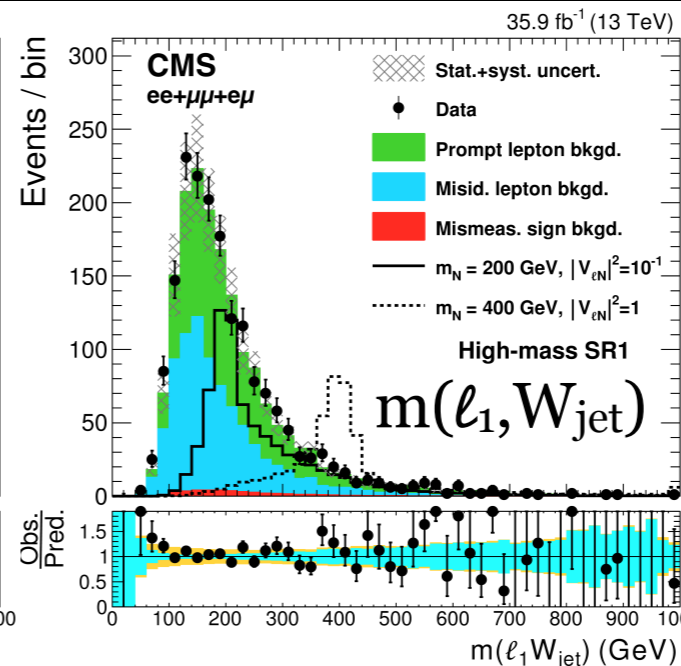
Region	$p_T^{\text{miss}}$ (GeV)	$(p_T^{\text{miss}})^2 / S_T$ (GeV)	$m(\ell^\pm \ell^\pm W_{jet})$ (GeV)	$m(W_{jet})$ (GeV)	$p_T^j$ (GeV)
Low-mass SR1+SR2	$< 80$	—	$< 300$	—	$> 20$
High-mass SR1	—	$< 15$	—	30–150	$> 25$
High-mass SR2	—	$< 15$	—	40–130	$> 200$



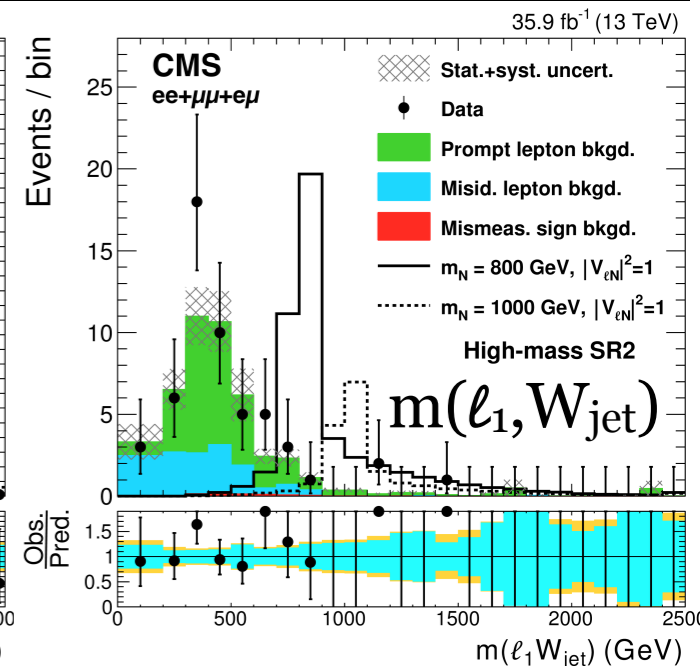
Low-mass SR2



Low-mass SR1

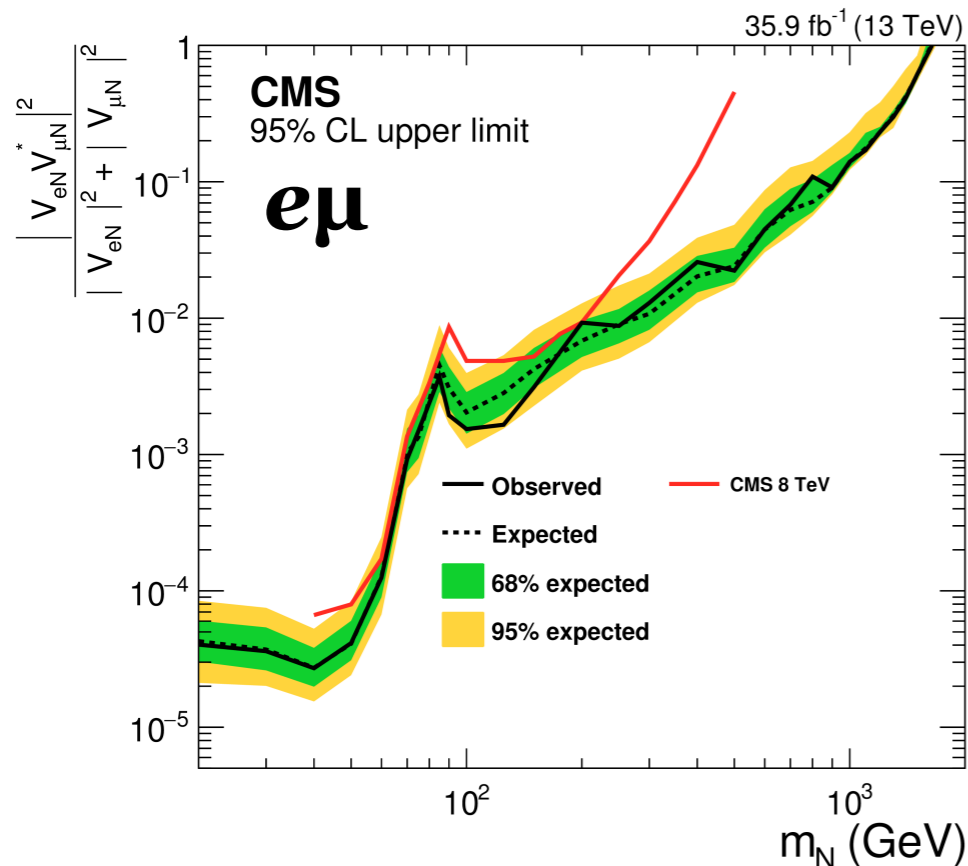
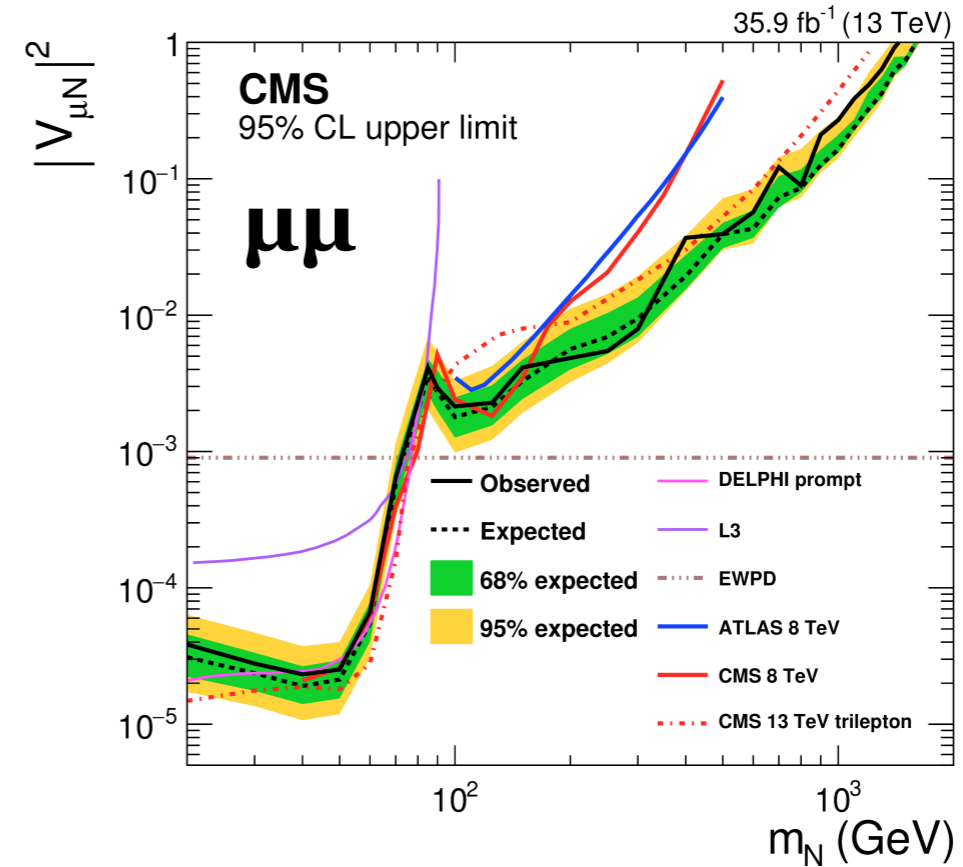
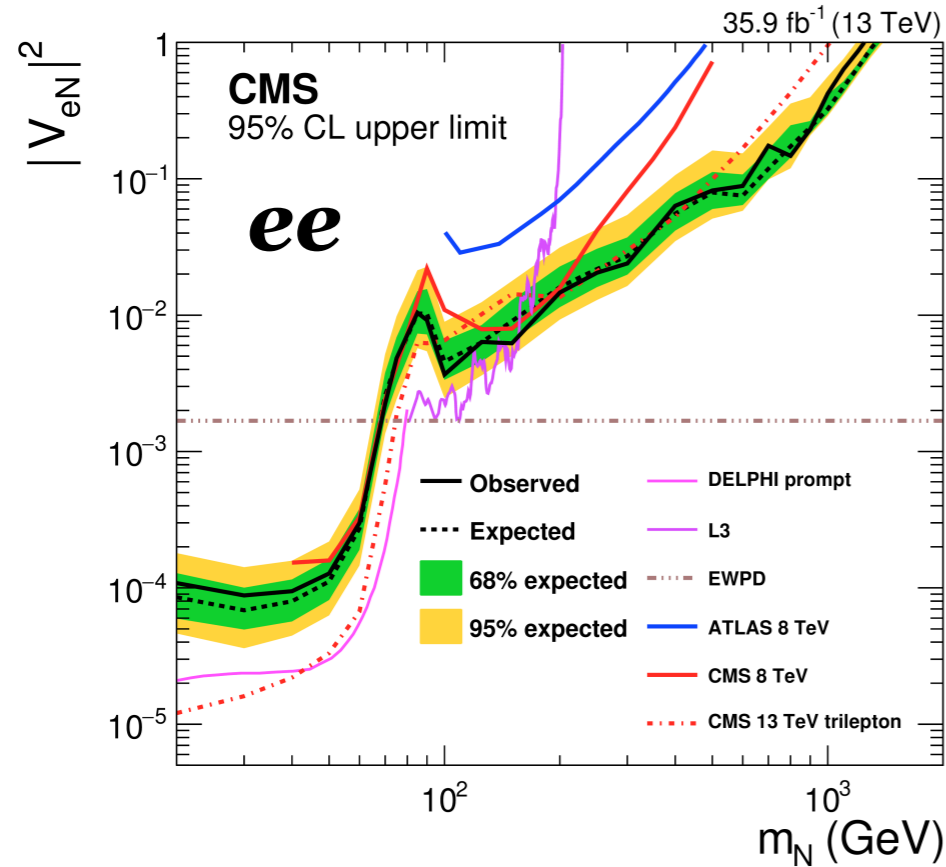


High-mass SR1



High-mass SR2

# Type-I - dilepton channel



- High-mass limit dominated by dilepton channel over trilepton:  $BR(W \rightarrow qq') > BR(W \rightarrow \ell\nu)$ .
- Low-mass sensitivity dominated by trilepton: minimum  $p_T(\ell) < \text{minimum } p_T(\text{jet})$ .

# Type-III searches

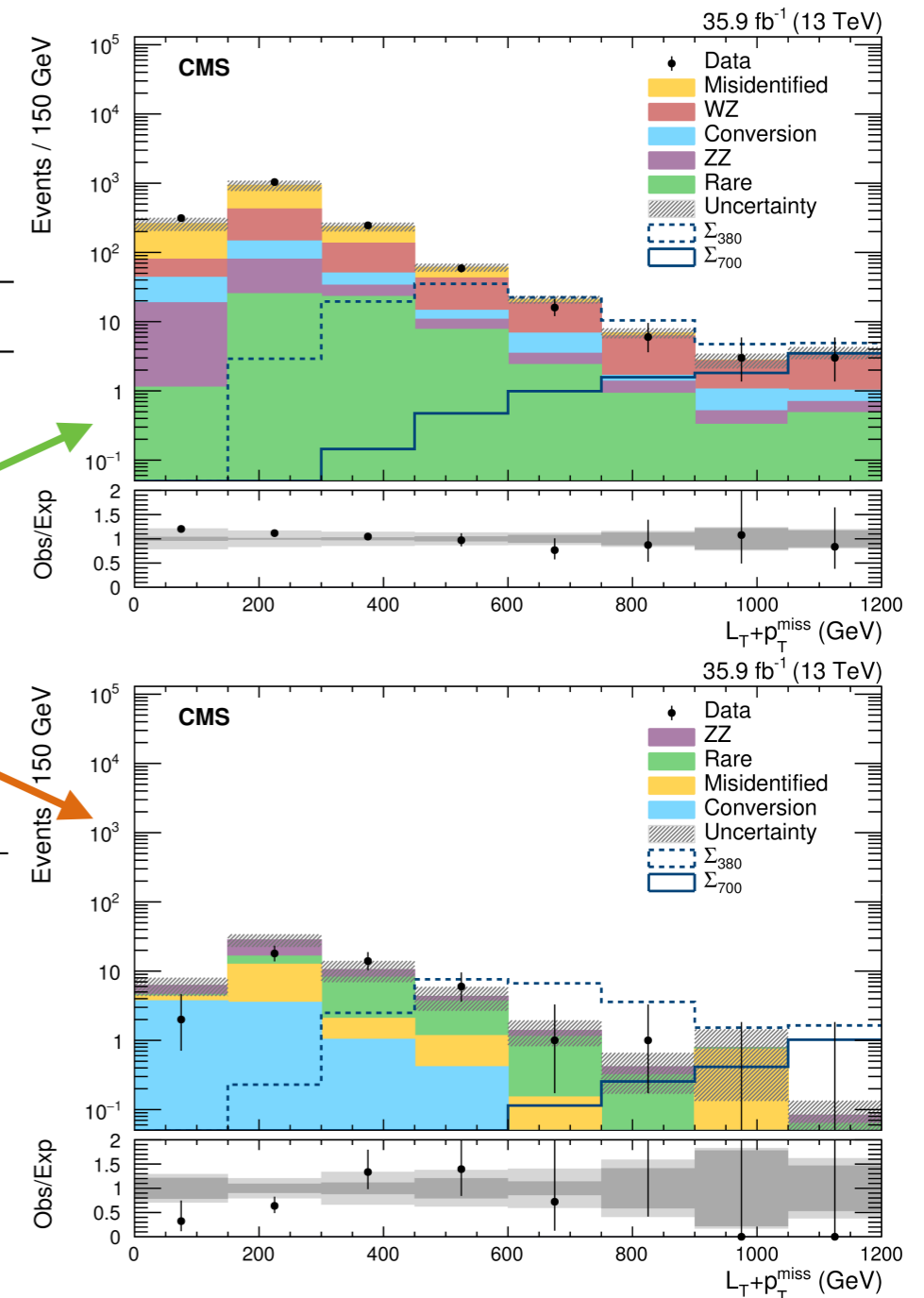


# Type-III - multilepton channel

- $pp \rightarrow \Sigma^0 \Sigma^\pm / \Sigma^\mp \Sigma^\pm \rightarrow$  Complicated decay chain leads to 27 different channels.
- Scalar sum of  $p_T(\ell)$ , called  $L_T$ , and  $E_T^{\text{miss}}$  are combined as primary discriminant. Contribution from neutrinos and charged leptons.
- Count lepton multiplicities based on charge product and multiplicity of same-flavour pairs.
- Categorisation for  $N_{\text{leptons}} = 3$  and  $\geq 4$ .
- $p_T(\ell) > 25, 15, 10$  GeV.

$N_{\text{leptons}}$	OSSF & mass	Variable	$p_T^{\text{miss}}$ requirement
3	OSSF1, on-Z	$M_T$	$p_T^{\text{miss}} > 100$ GeV
	OSSF1, above-Z	$L_T + p_T^{\text{miss}}$	—
3	OSSF1, below-Z	$L_T + p_T^{\text{miss}}$	$p_T^{\text{miss}} > 50$ GeV
	OSSF0	$L_T + p_T^{\text{miss}}$	—
$\geq 4$	OSSF1	$L_T + p_T^{\text{miss}}$	—
	OSSF2	$L_T + p_T^{\text{miss}}$	$p_T^{\text{miss}} > 50$ GeV if on-Z

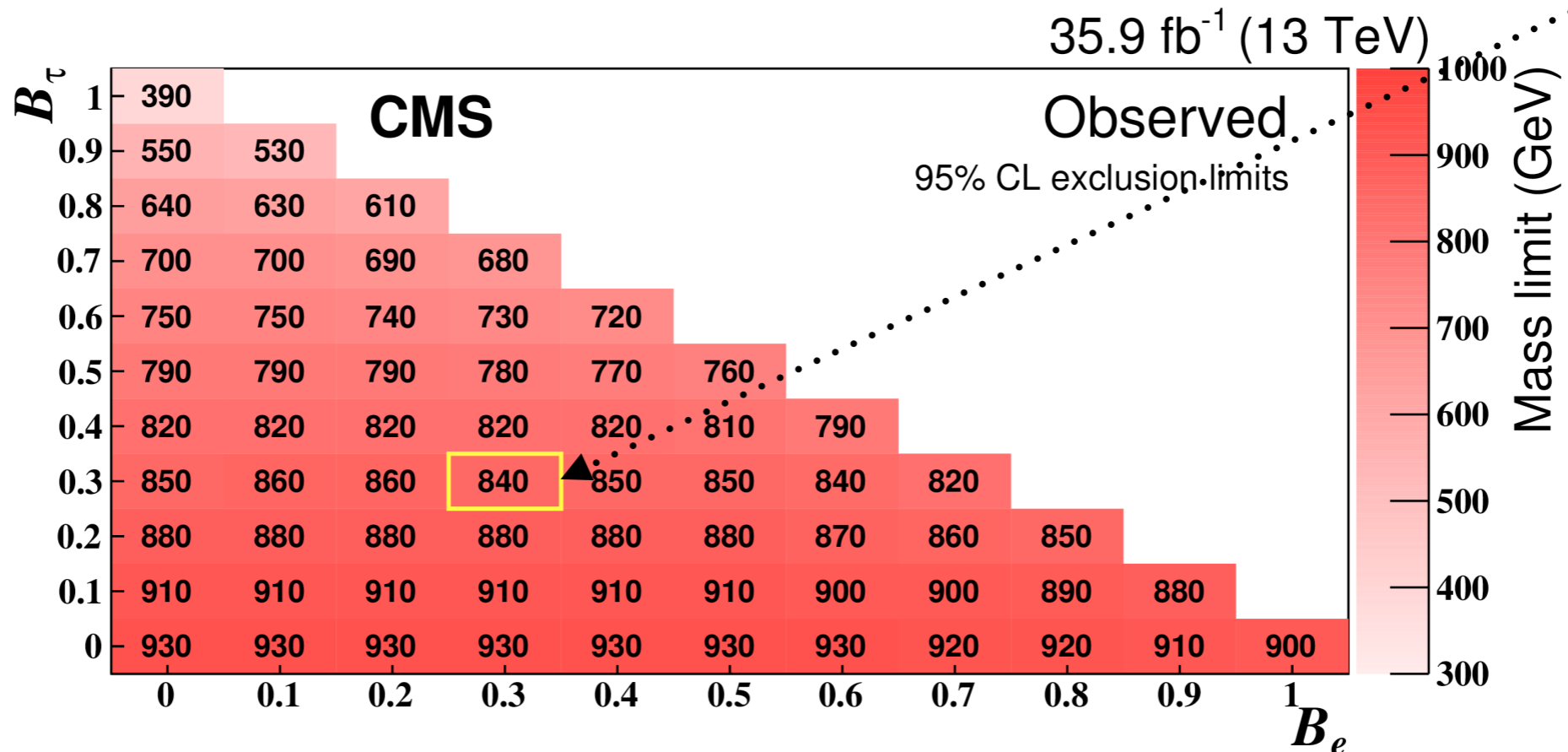
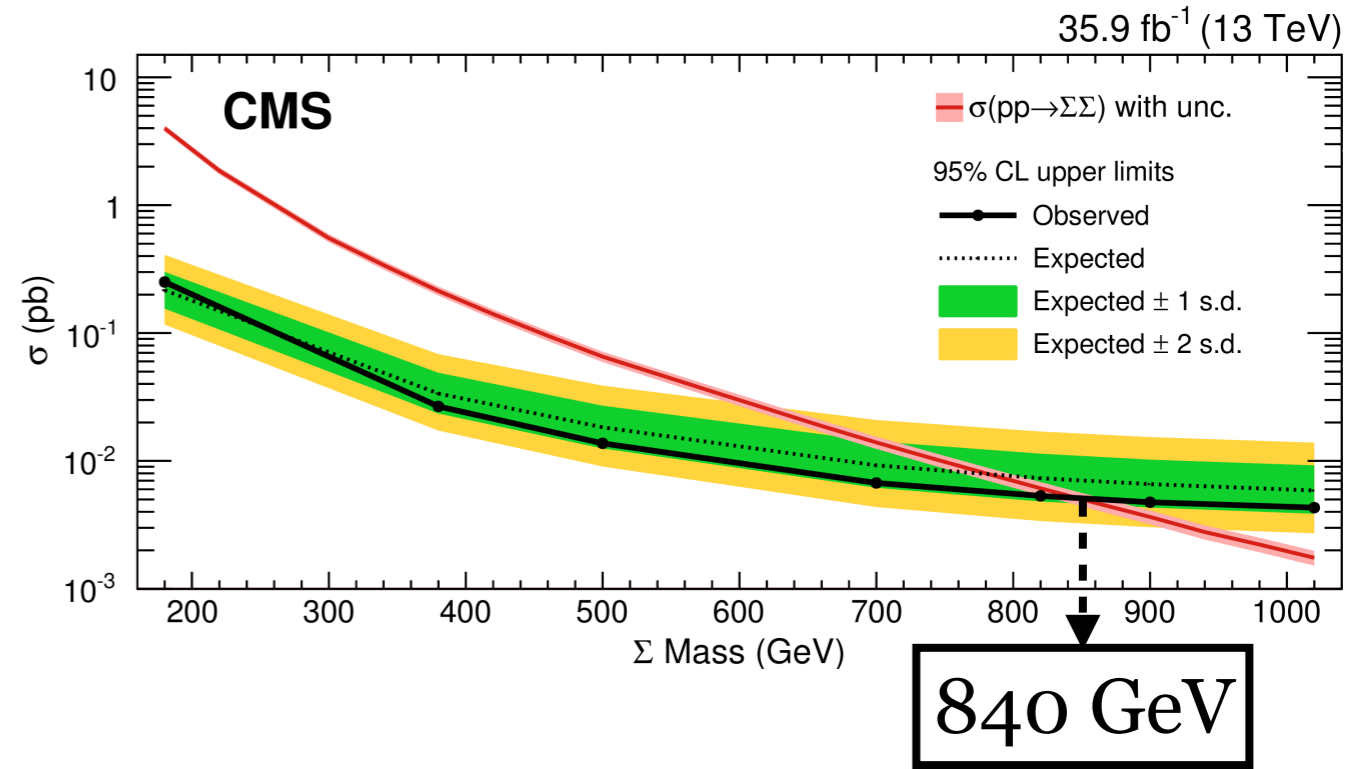
OSSF $n$  = number of opposite-sign, same-flavour pairs



# Type-III - multilepton channel

CMS:  
1708.07962

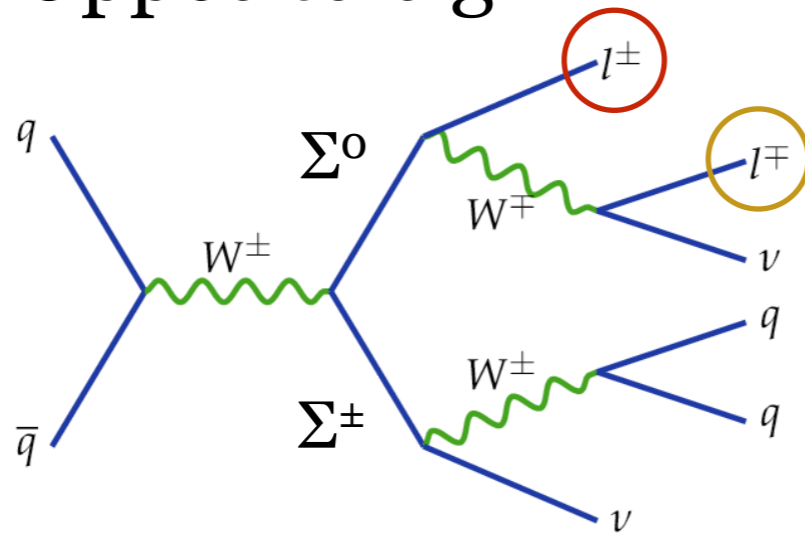
- Upper limits on the cross-section for  $pp \rightarrow \Sigma^0 \Sigma^\pm / \Sigma^\mp \Sigma^\pm$
- “Flavour democratic” scenario: branching ratios to all leptons equal.
- 2d limits on  $BR_\ell \sim |V_\ell|^2 / (|V_\mu|^2 + |V_e|^2 + |V_\tau|^2)$  where  $V_\ell$  is a mixing angle b/w the heavy and light fermions.
- Taus included via their leptonic decay.



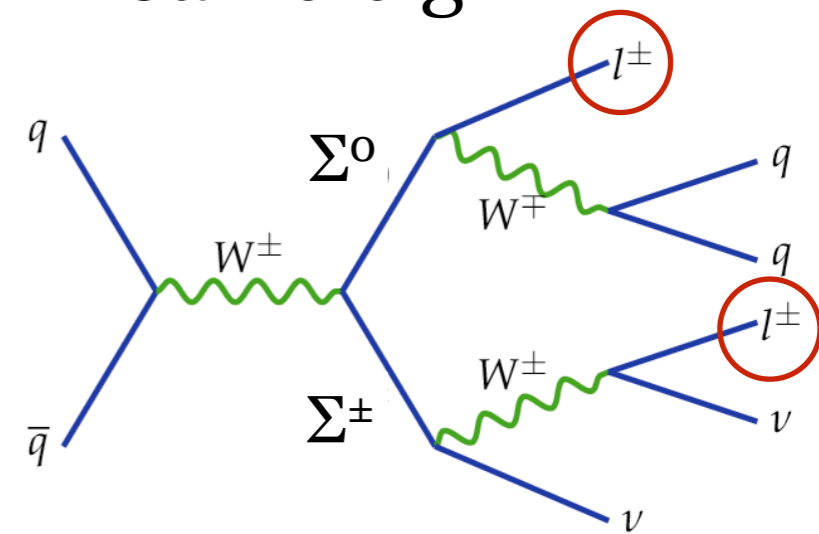
# Type-III - dilepton channel

- $pp \rightarrow \Sigma^0 \Sigma^\pm$
- Opposite and same -sign optimised independently.
- Two resolved jets in final state.  $M(j,j)$  consistent with  $W$  mass and  $E_T^{\text{miss}}$ .
- Scalar sum of  $p_T(\ell)$ , called  $H_T$ , and  $E_T^{\text{miss}}$  are combined as primary discriminant.

## Opposite-sign



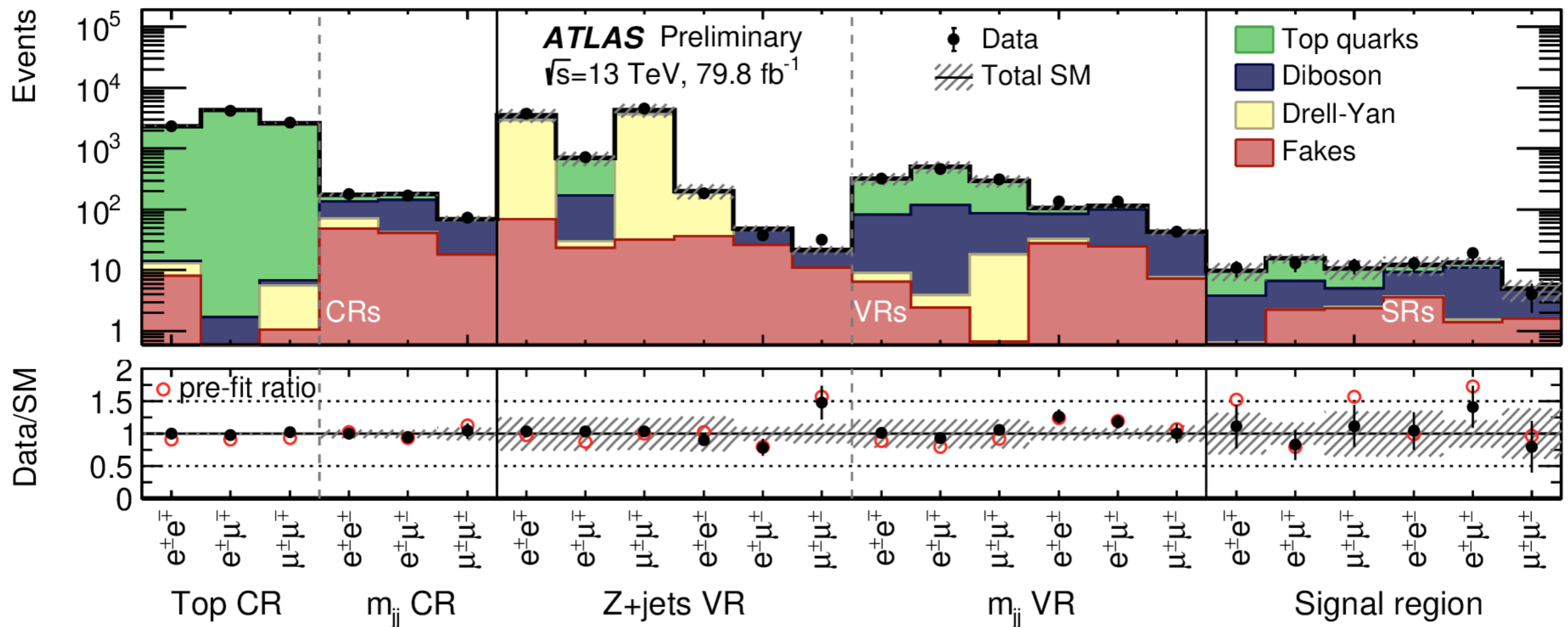
## Same-sign



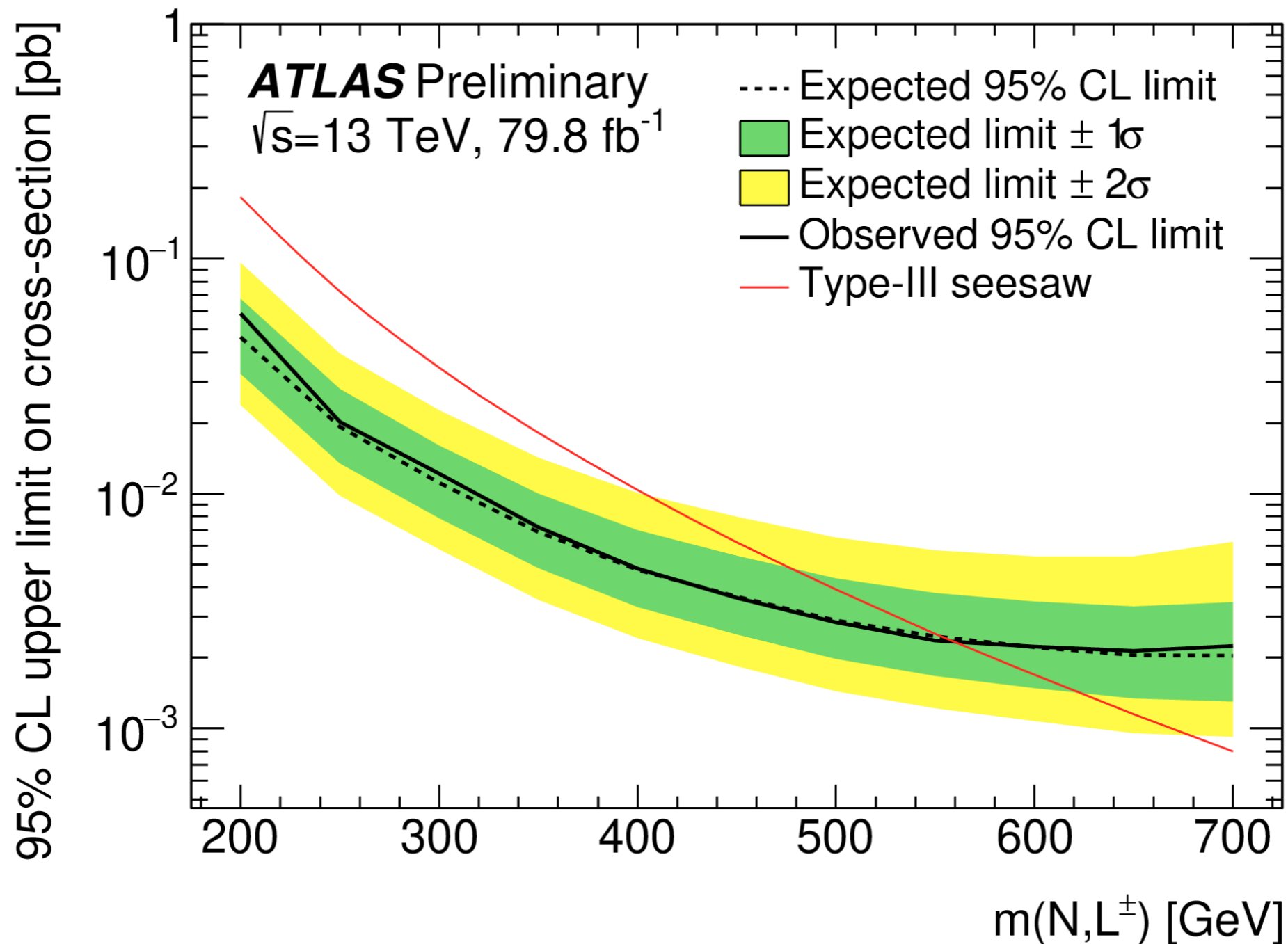
	OS ( $l^+l^- = e^+e^-, e^\pm\mu^\mp, \mu^+\mu^-$ )				SS ( $l^\pm l^\pm = e^\pm e^\pm, e^\pm\mu^\pm, \mu^\pm\mu^\pm$ )			
	Top CR	Z + jets VR	$m_{jj}$ VR	SR	Z + jets VR	$m_{jj}$ VR	$m_{jj}$ CR	SR
$N(\text{jet})$	$\geq 2$	$\geq 2$	$\geq 2$	$\geq 2$	$\geq 2$	$\geq 2$	$\geq 2$	$\geq 2$
$N(b\text{-jet})$	$\geq 2$	0	0	0	0	0	0	0
$m_{jj} [\text{GeV}]$	[60, 100)	[60, 100)	<b>[35, 60) <math>\cup</math> [100, 125)</b>	[60, 100)	[60, 100)	<b>[0, 60) <math>\cup</math> [100, 300)</b>	<b>[0, 60) <math>\cup</math> [100, 300)</b>	[60, 100)
$m_{\ell\ell} [\text{GeV}]$	[110, $\infty$ )	<b>[70, 110)</b>	[110, $\infty$ )	[110, $\infty$ )	<b>[70, 100)</b>	[100, $\infty$ )	[100, $\infty$ )	[100, $\infty$ )
$\text{Sig}(E_T^{\text{miss}})$	$\geq 5$	$\geq 5$	$\geq 10$	$\geq 10$	$\geq 5$	$\geq 5$	$\geq 5$	$\geq 7.5$
$\Delta\phi(E_T^{\text{miss}}, l)_{\text{min}}$				$\geq 1$				
$p_T(jj) [\text{GeV}]$				[100, $\infty$ )				[60, $\infty$ )
$p_T(\ell\ell) [\text{GeV}]$				[100, $\infty$ )				[100, $\infty$ )
$H_T + E_T^{\text{miss}} [\text{GeV}]$	[300, $\infty$ )	[300, $\infty$ )	[300, $\infty$ )	[300, $\infty$ )		[500, $\infty$ )	[300, 500)	[300, $\infty$ )

# Type-III - dilepton channel

- Full shape information used in the final fit.
- CRs used to normalise simulated backgrounds directly in the fit procedure.
- Independent validation regions defined to validate the modelling.
- Z+jets validation regions to validate charge misidentification modelling.



# Type-III - dilepton channel



- All channels are combined to derive the final limit.
- Assume “flavour democratic” scenario with equal branching ratios to all leptonic decays.

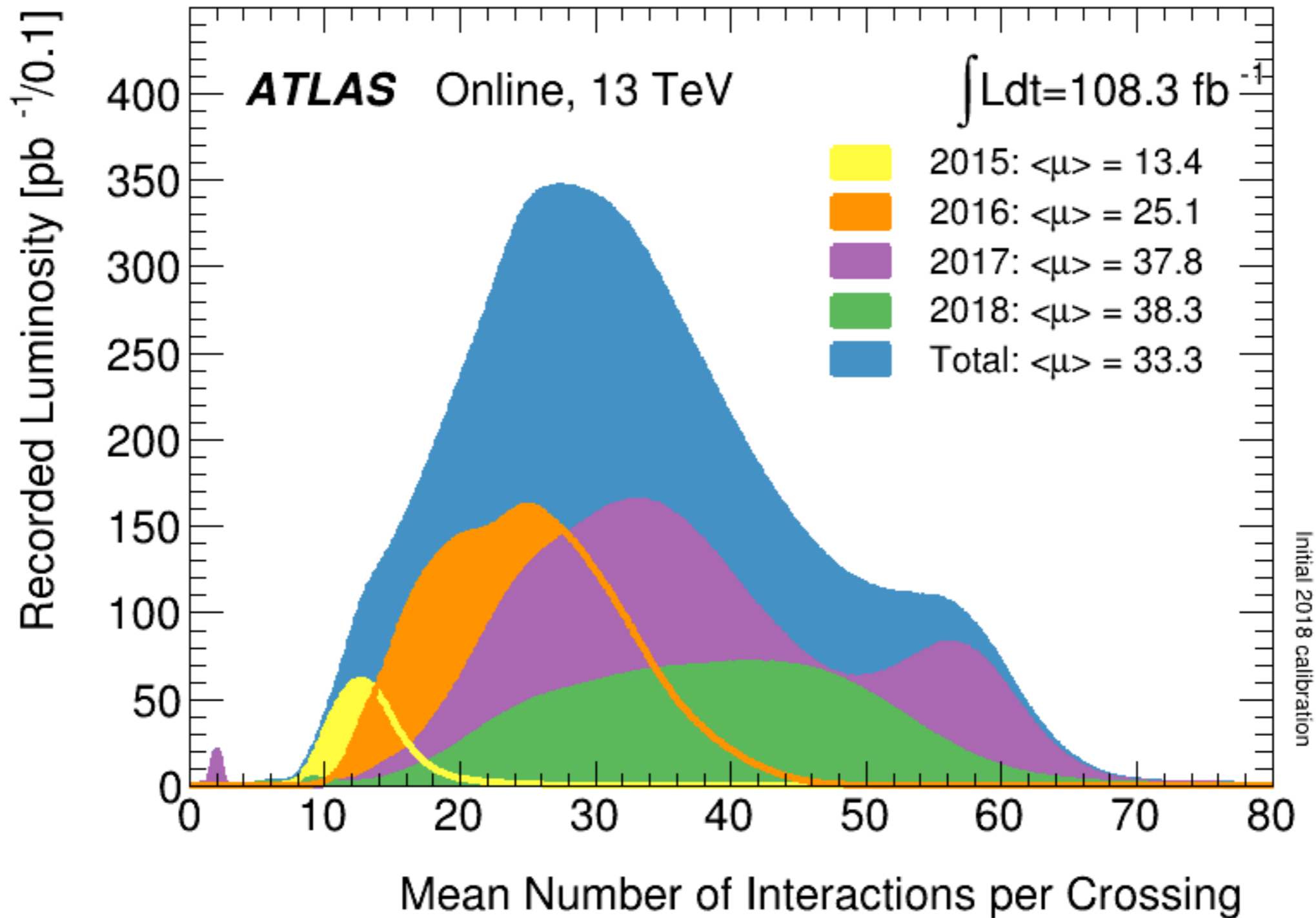
# Conclusions

- The see-saw mechanism is a promising paradigm for neutrino mass generations embedded in many theories of new physics.
- Collider experiments do have access to different incarnations of it: Type-I/II/III.
- ATLAS and CMS have a rich plan for the end of Run-II. Stay tuned!

Thank you!

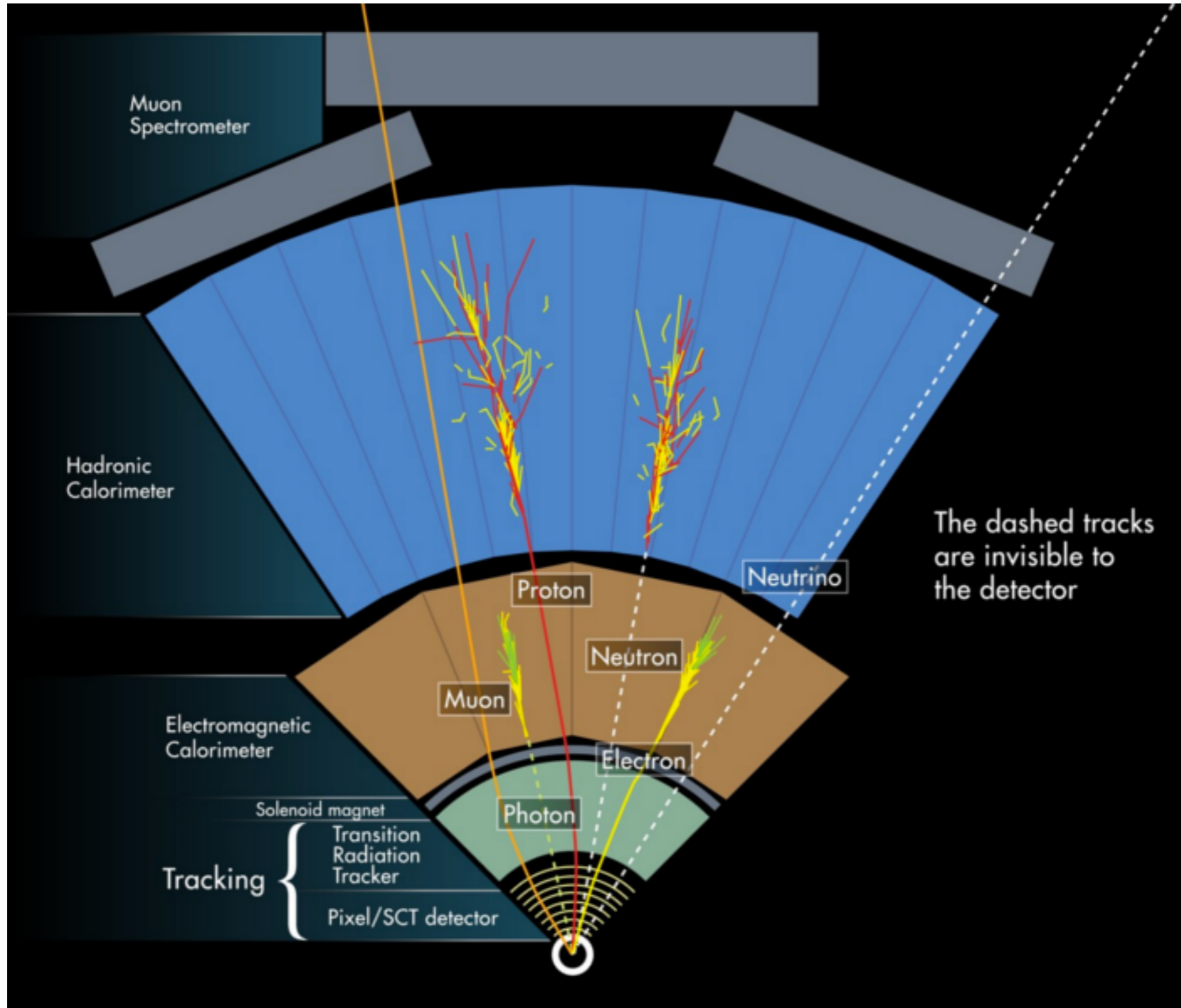
# Backup

# ATLAS and CMS

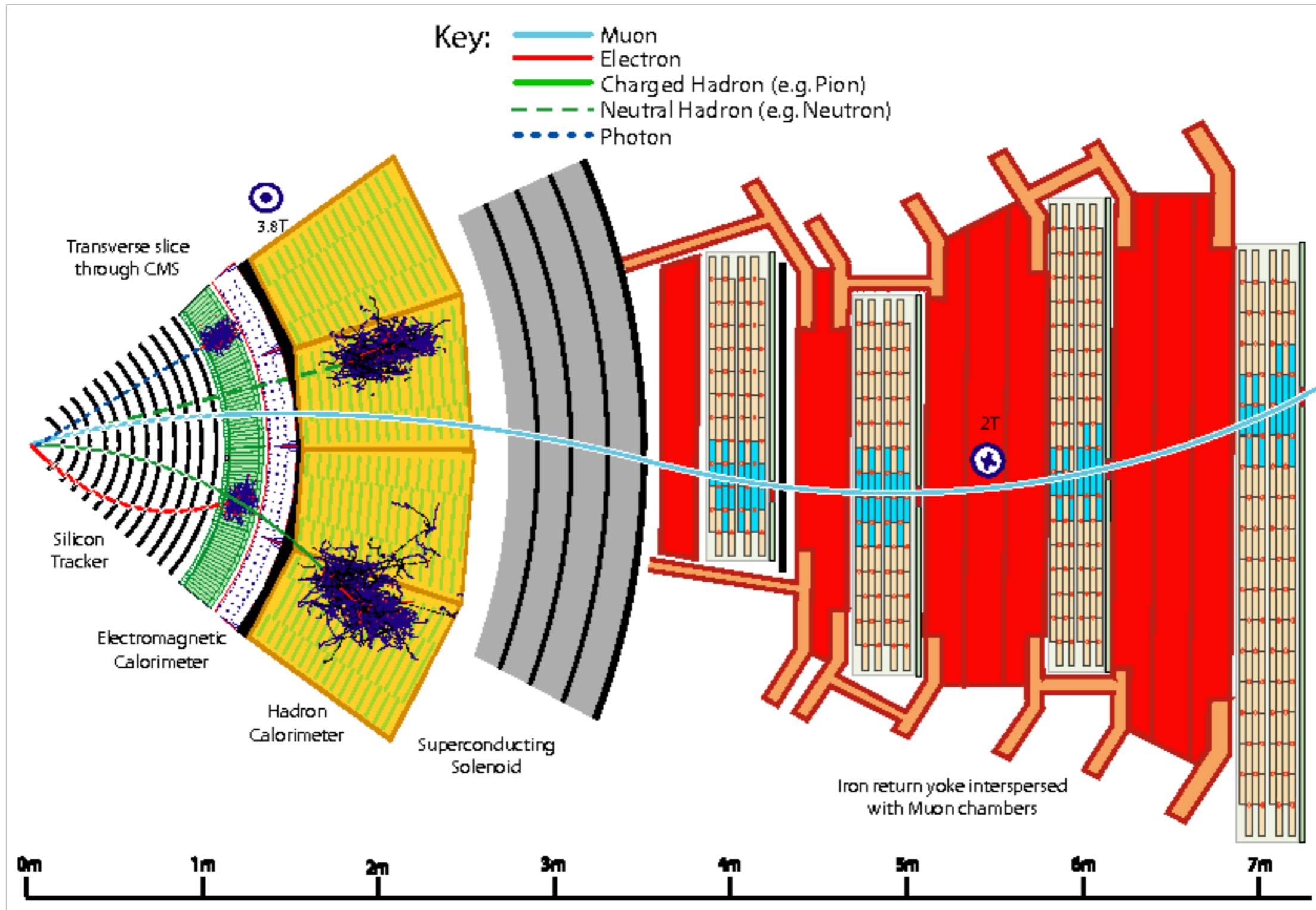




# ATLAS and CMS



# ATLAS and CMS



# ATLAS and CMS

## CMS Integrated Luminosity, pp

Data included from 2010-03-30 11:22 to 2018-06-30 02:58 UTC

