



# Neutrino Physics Overview

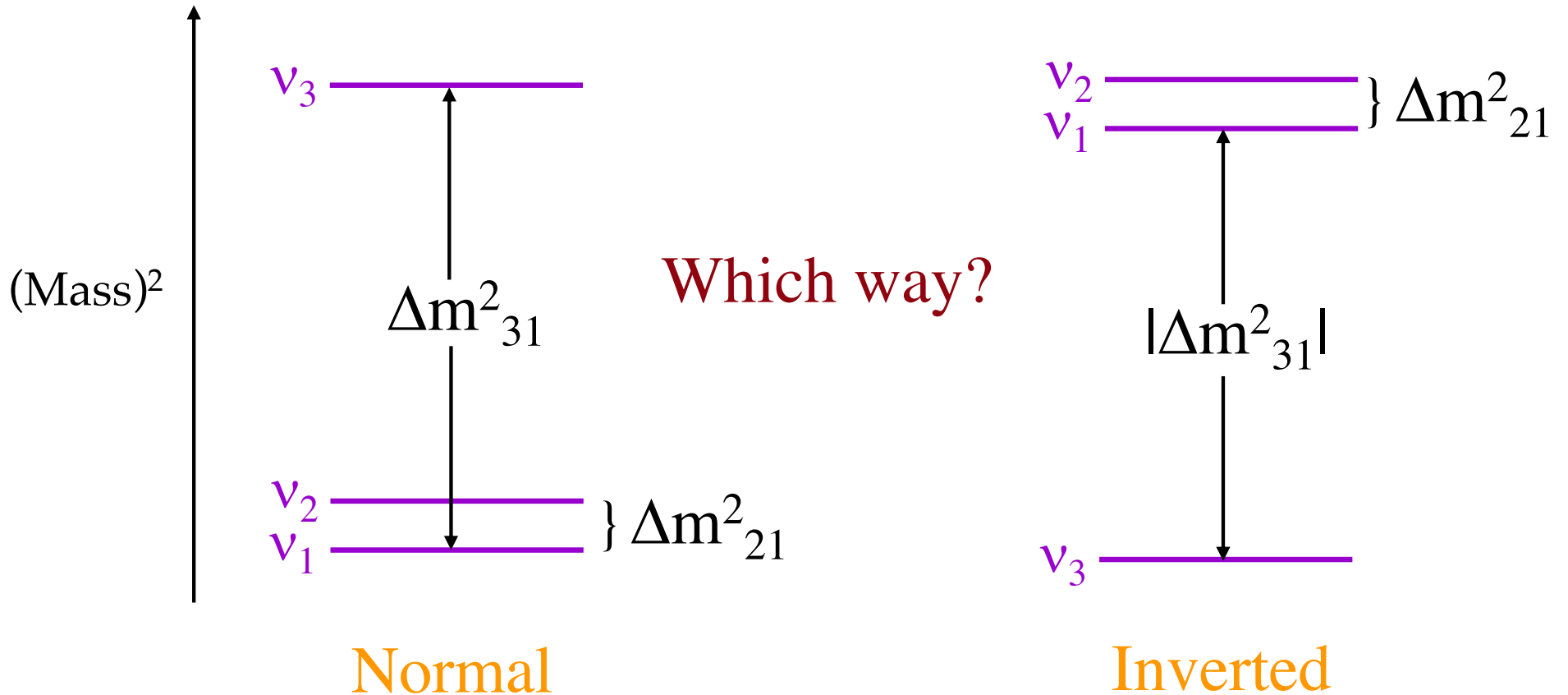
Boris Kayser  
ICISE  
July 16, 2018



# What We Have Learned

(c) Kamioka Observatory, ICRR(Institute for Cosmic Ray Research), The University of Tokyo

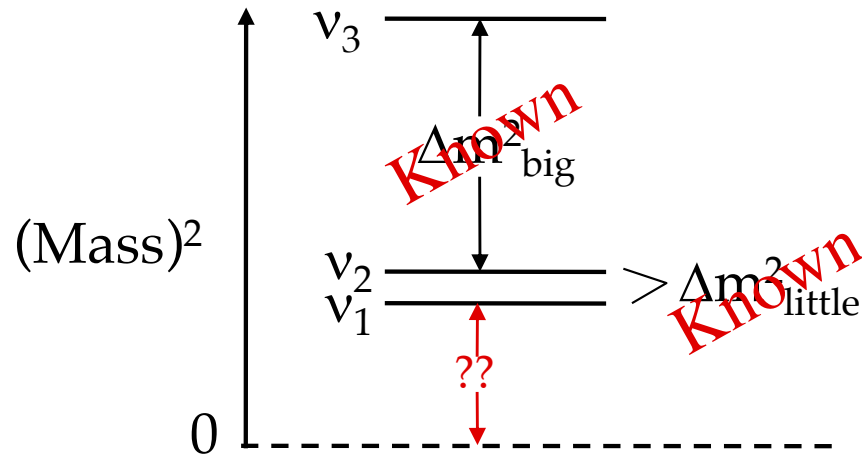
# The (Mass)<sup>2</sup> Spectrum



$$\Delta m_{21}^2 \cong 7.5 \times 10^{-5} \text{ eV}^2, \quad |\Delta m_{31}^2| \cong 2.5 \times 10^{-3} \text{ eV}^2$$

Are there *more* mass eigenstates?

# Constraints On the Absolute Scale of Neutrino Mass



How far above zero is the whole pattern?

Cosmology, *under certain assumptions*  $\longrightarrow \sum m(\nu_i) < 0.2 \text{ eV}$   
*All i*

Tritium beta decay  $\longrightarrow \sqrt{0.69m^2(\nu_1) + 0.29m^2(\nu_2) + 0.02m^2(\nu_3)} < 2 \text{ eV}$

Oscillation  $\longrightarrow \text{Mass}[\text{Heaviest } \nu_i] > \sqrt{\Delta m^2_{\text{big}}} > 0.05 \text{ eV}$

# Leptonic Mixing

Mixing means that —

**PMNS mixing matrix**

$$|\nu_\alpha\rangle = \sum_i U_{\alpha i}^* |\nu_i\rangle .$$

Neutrino of flavor

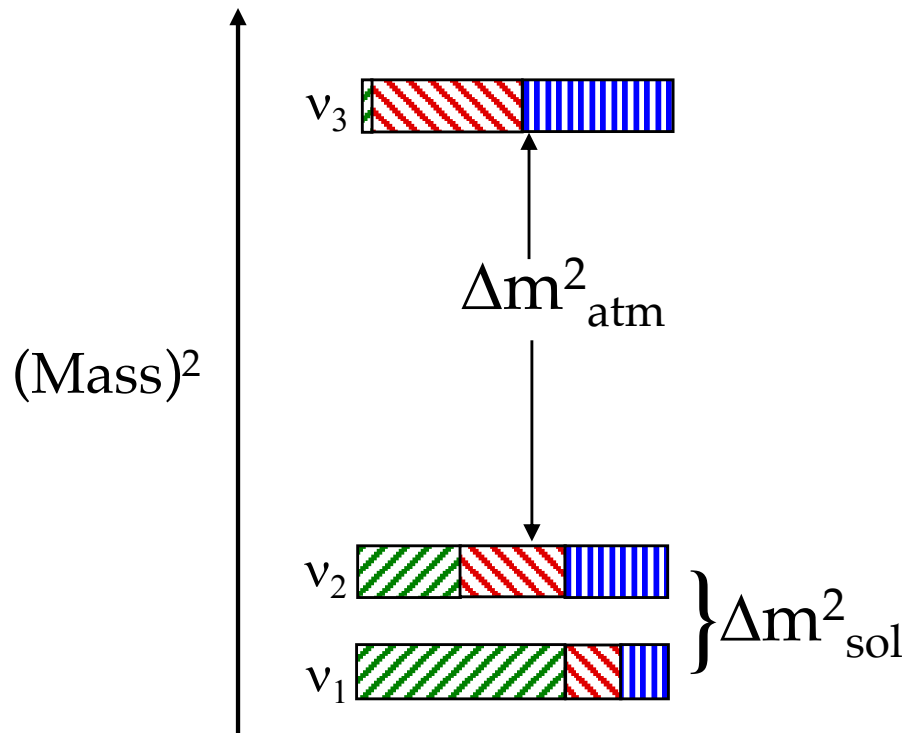
$\alpha = e, \mu, \text{ or } \tau$

Neutrino of definite mass  $m_i$

$$\text{Inversely, } |\nu_i\rangle = \sum_\alpha U_{\alpha i} |\nu_\alpha\rangle . \quad (\text{if } U \text{ is unitary})$$

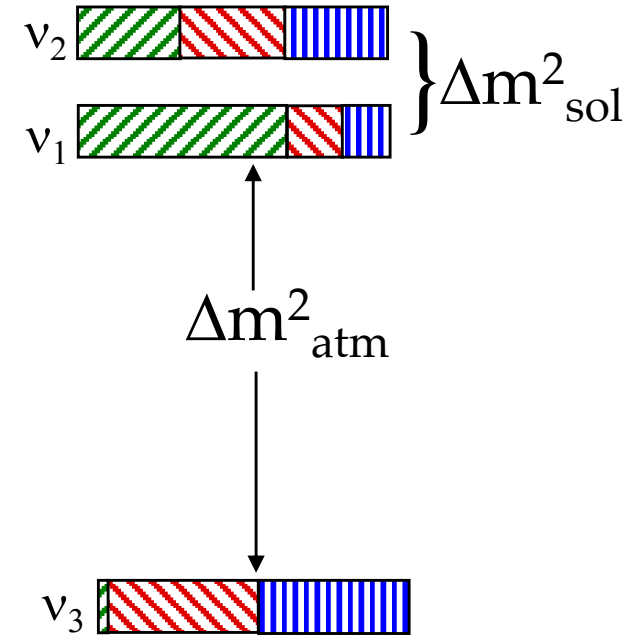
Flavor- $\alpha$  fraction of  $\nu_i = |U_{\alpha i}|^2$  .

Experimentally, the flavor fractions are —



Normal

or



Inverted

  $\nu_e [ |U_{ei}|^2 ]$

  $\nu_\mu [ |U_{\mu i}|^2 ]$

  $\nu_\tau [ |U_{\tau i}|^2 ]$

# The Leptonic Mixing Matrix $U$

$$U = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \times \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \times \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$\times \underbrace{\begin{bmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{bmatrix}}_{\text{Majorana phases}}$$

$c_{ij} \equiv \cos \theta_{ij}$   
 $s_{ij} \equiv \sin \theta_{ij}$

$$\theta_{12} \approx 34^\circ, \theta_{23} \approx 42-51^\circ, \theta_{13} \approx 8.5^\circ \quad \delta = ??$$

(Tortola at  $\nu$  2018)

# Precision measurements

<https://globalfit.astroparticles.es/>

parameter	best fit $\pm 1\sigma$	$3\sigma$ range	
$\Delta m_{21}^2$ [ $10^{-5}\text{eV}^2$ ]	$7.55^{+0.20}_{-0.16}$	7.05–8.14	2.4%
$ \Delta m_{31}^2 $ [ $10^{-3}\text{eV}^2$ ] (NO)	$2.50 \pm 0.03$	2.41–2.60	1.3%
$ \Delta m_{31}^2 $ [ $10^{-3}\text{eV}^2$ ] (IO)	$2.42^{+0.03}_{-0.04}$	2.31–2.51	1.3%
$\sin^2 \theta_{12}/10^{-1}$	$3.20^{+0.20}_{-0.16}$	2.73–3.79	5.5%
$\sin^2 \theta_{23}/10^{-1}$ (NO)	$5.47^{+0.20}_{-0.30}$	4.45–5.99	4.7%
$\sin^2 \theta_{23}/10^{-1}$ (IO)	$5.51^{+0.18}_{-0.30}$	4.53–5.98	4.4%
$\sin^2 \theta_{13}/10^{-2}$ (NO)	$2.160^{+0.083}_{-0.069}$	1.96–2.41	3.5%
$\sin^2 \theta_{13}/10^{-2}$ (IO)	$2.220^{+0.074}_{-0.076}$	1.99–2.44	3.5%
$\delta/\pi$ (NO)	$1.32^{+0.21}_{-0.15}$	0.87–1.94	10%
$\delta/\pi$ (IO)	$1.56^{+0.13}_{-0.15}$	1.12–1.94	9%

relative  $1\sigma$  uncertainty

deSalas et al, 1708.01186 (May 2018)

**From M. Tortola, at  $\nu$  2018.**





# What We Would Like To Find Out

NASA Hubble Photo

- Is the physics behind the masses of neutrinos different from that behind the masses of all other known particles?
- Are neutrinos their own antiparticles?

• What is the absolute scale of neutrino mass?

• Is the spectrum like  $\begin{matrix} \text{=} \\ \text{=} \end{matrix}$  or  $\begin{matrix} \text{=} \\ \text{=} \end{matrix}$  ?

• Is  $\theta_{23}$  maximal?

• Do neutrino interactions  
violate CP?

Is  $P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta) \neq P(\nu_\alpha \rightarrow \nu_\beta)$  ?

• Is CP violation involving neutrinos  
the key to understanding the baryon –  
antibaryon asymmetry of the universe?

- What can neutrinos and the universe tell us about one another?

- What are the electromagnetic properties of neutrinos?

- Are there *more* than 3 mass eigenstates?
  - Are there “sterile” neutrinos that don’t couple to the W or Z?

- Do neutrinos have Non-Standard-Model interactions?

- Is coherent neutrino – nucleus scattering anomalous?

- Do neutrinos break the rules?
  - Violation of Lorentz invariance?
  - Violation of CPT invariance?
  - Departures from quantum mechanics?

# Questions With Current Hints

Does the leptonic mixing matrix  $U$  violate CP?

That is, is the CP phase  $\delta \neq 0$  or  $\pi$ , so that  $U$  is not real?

Is the order of the mass eigenstates  $\overline{\text{NO}}$  or  $\overline{\text{IO}}$  ?

# New results on the CP phase

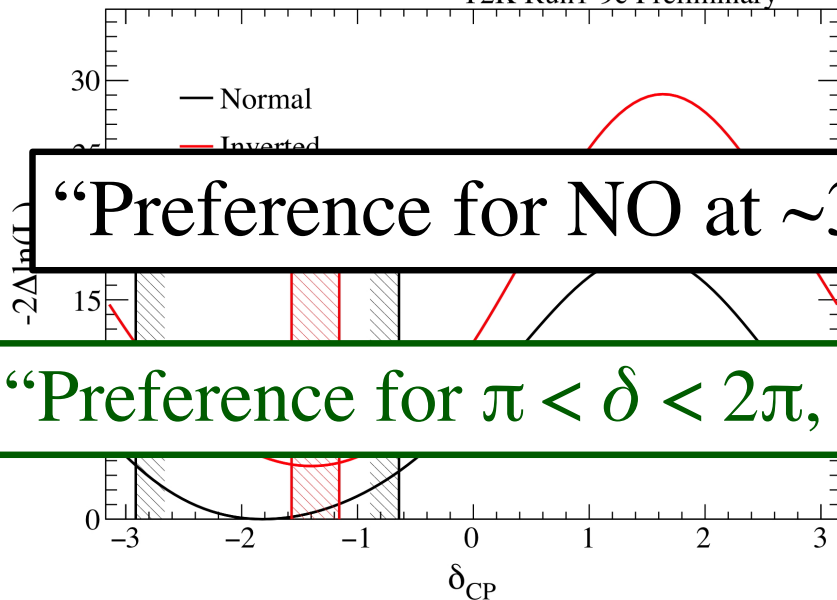
T2K

Talk by M. Wascko

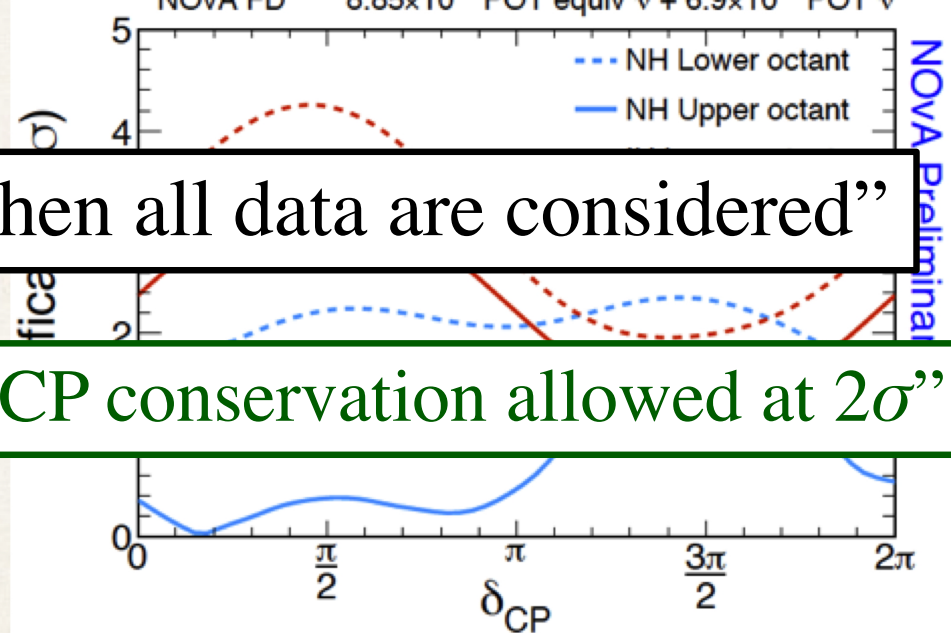
NOvA

Talk by M. Sánchez

T2K Run1-9c Preliminary



NOvA FD 8.85x10<sup>20</sup> POT equiv  $\nu$  + 6.9x10<sup>20</sup> POT  $\bar{\nu}$



“Preference for NO at  $\sim 3\sigma$  when all data are considered”

“Preference for  $\pi < \delta < 2\pi$ , with CP conservation allowed at  $2\sigma$ ”

(fit with reactor constraint)

CP conservation more disfavoured  
 $\Rightarrow \delta = 0, \pi$  outside  $2\sigma$  region for NO & IO

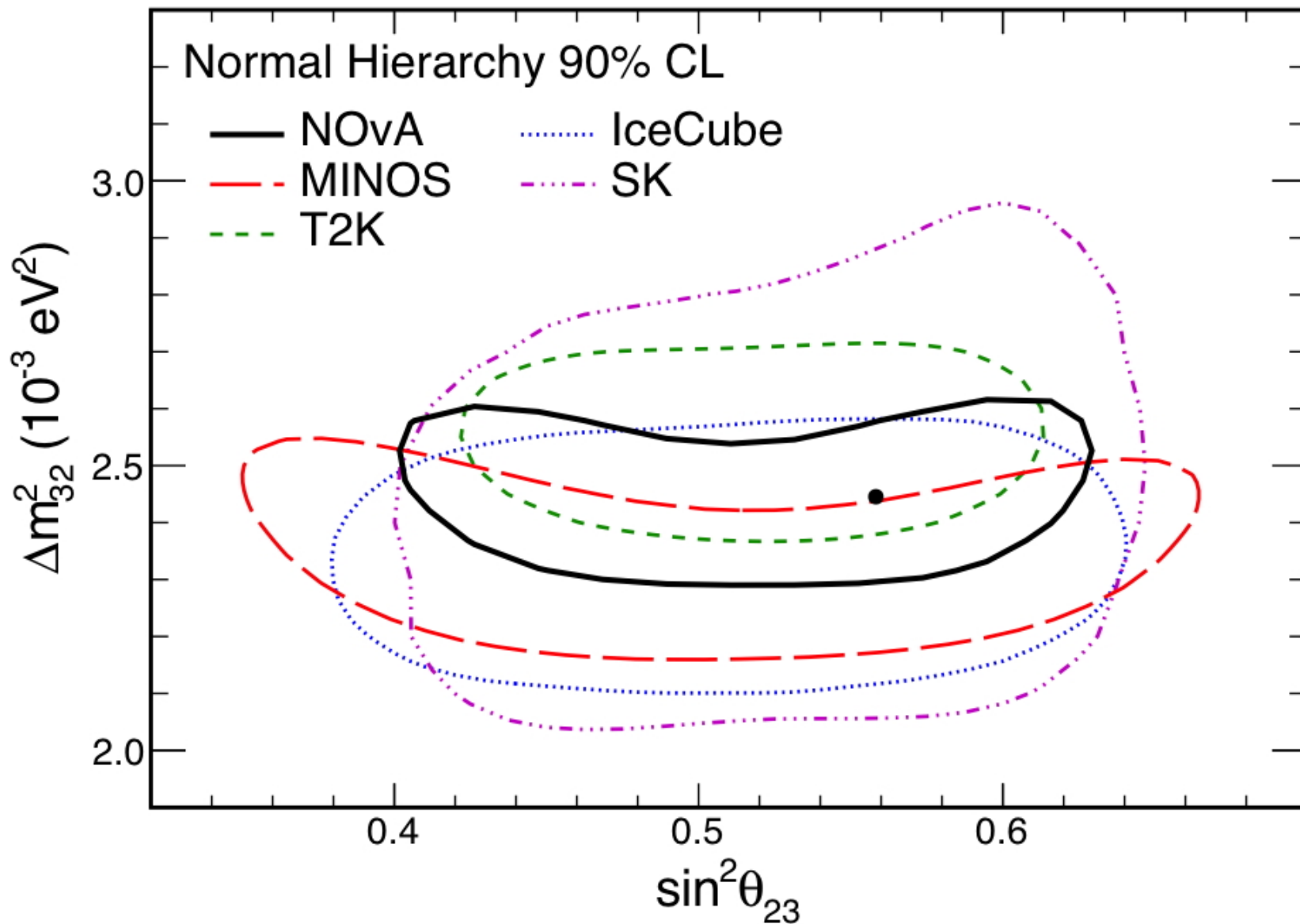


Best fit:  $\delta_{NO} = 0.17\pi$   $\delta_{IO} \sim 1.5\pi$

$\Rightarrow$  improved status for  $\delta = 0$

M. Tortola at  $\nu$  2018

# Is $\theta_{23}$ maximal ( $45^\circ$ )?



(NOvA; 1806.00096)



Are there eV-mass sterile neutrinos?

$$\text{Probability (Oscillation)} \propto \sin^2 \left[ 1.27 \Delta m^2 (\text{eV}^2) \frac{L(\text{m})}{E(\text{MeV})} \right]$$

There are several hints of oscillation with  $L(\text{m})/E(\text{MeV}) \sim 1$ .

These suggest that there is a  $\Delta m^2$  larger than the 2 established ones. If so, there must be a 4th mass eigenstate, hence a 4th flavor.

Since only 3 neutrino flavors couple to the Z or W, this new flavor must be sterile ( $\equiv$  does not couple to Z or W).

# The Hints of eV<sup>2</sup>-Scale $\Delta m^2$

<u>Experiment</u>	<u>Possible Oscillation</u>	<u>Comment</u>
LSND	$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$	Interesting
MiniBooNE	$\nu_\mu \rightarrow \nu_e$	Low energy excess?
MiniBooNE	$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$	Low energy excess?
Reactor Exps.	$\bar{\nu}_e \rightarrow$ Not $\bar{\nu}_e$	Flux uncertainty $\sim$ 6% size of effect
<sup>51</sup> Cr and <sup>37</sup> Ar Source Exps.	$\nu_e \rightarrow$ Not $\nu_e$	Detection efficiency?

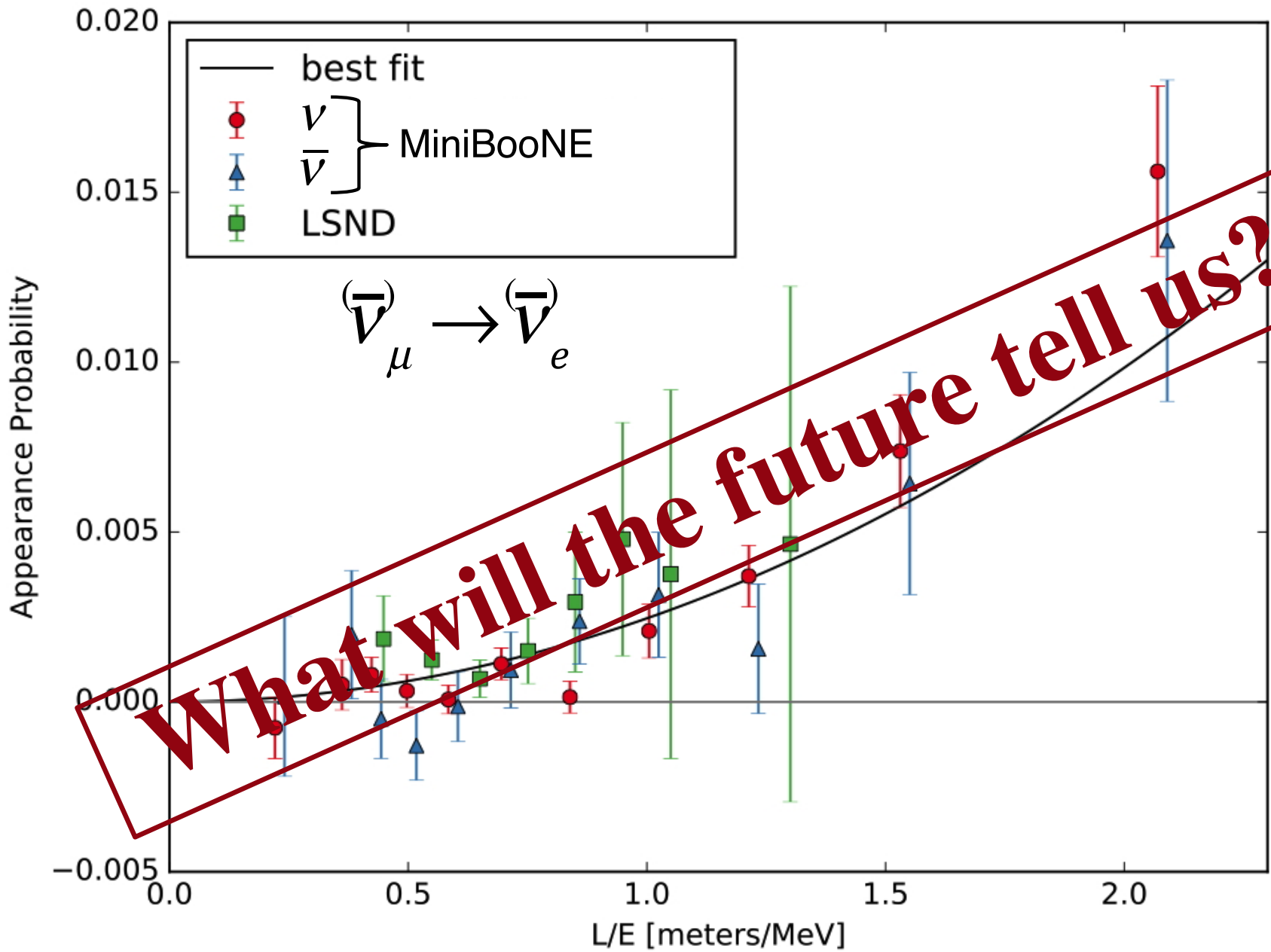
There is strong tension between the evidence for rapid  $\nu_{\mu} \rightarrow \nu_e$  and  $\bar{\nu}_{\mu} \rightarrow \bar{\nu}_e$ , and the limits on rapid

$$\nu_{\mu} \rightarrow \nu_{\mu} \text{ and } \bar{\nu}_e \rightarrow \bar{\nu}_e .$$

(M. Dentler et al., 1803.10661)

A recent paper shows that the MiniBooNE and LSND positive indications of *something* going on (*a sterile neutrino???*) are not inconsistent.

(MiniBooNE, 1805.12028)



1 eV scale sterile neutrinos, if real, could greatly affect the interpretation of the CP-violation studies of the long-baseline experiments.

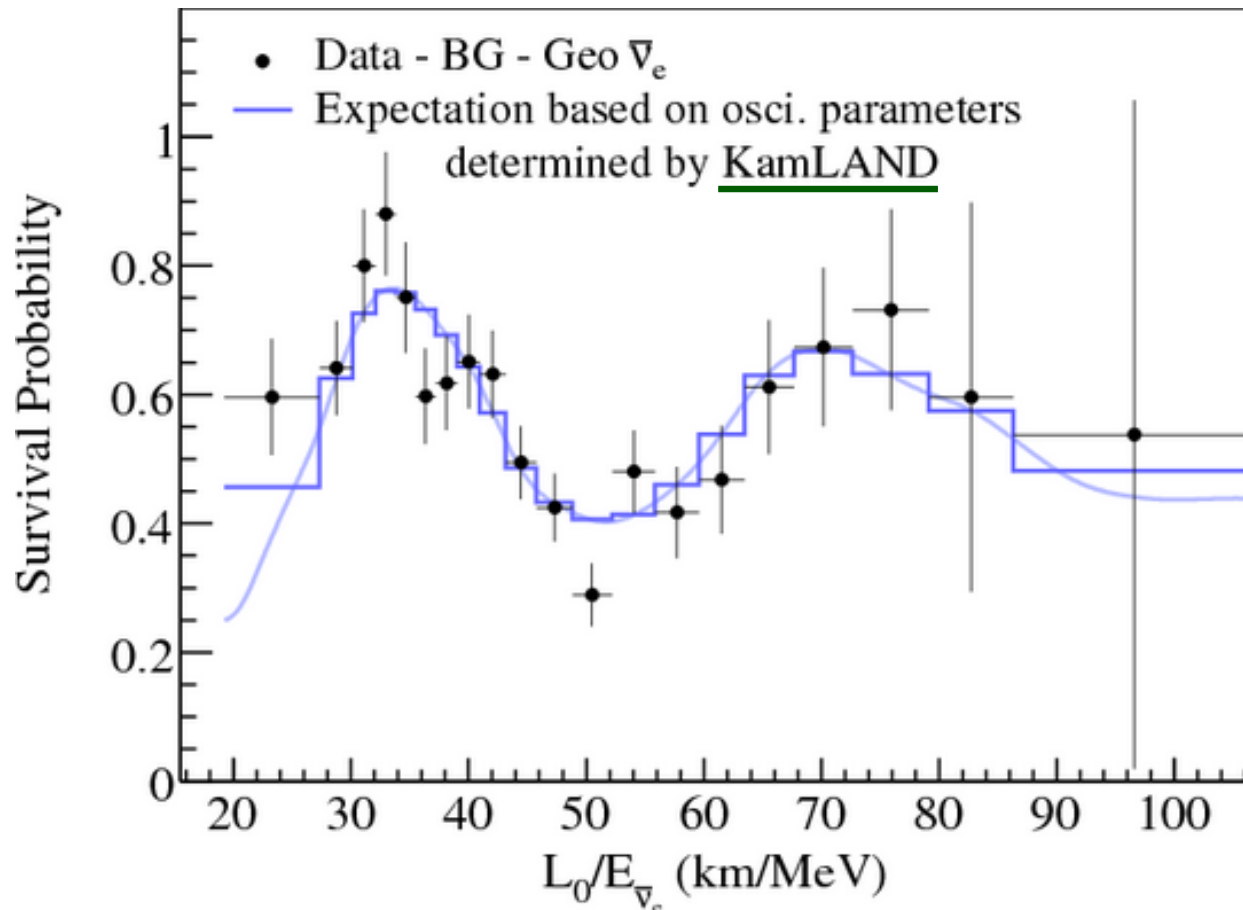
*For example, long-baseline data that, when analyzed disregarding the possibility of sterile neutrinos, indicate that CP violation is very small or absent, could in fact hide quite large CP violation.*

**(R. Gandhi, B.K., M. Masud, S. Prakash)**

***The presence or absence of 1 eV scale sterile neutrinos needs to be settled experimentally, and many experiments to do that, using neutrinos from sources, reactors, accelerators, or elsewhere, are in progress or planned.***

To confirm their existence, it would be nice to see actual

# Oscillatory Behavior a la



DANSS and NEOS do see hints of wiggles, but make no claims.

A theoretical issue that  
should not cause any worries

Much of what we know about the neutrinos was  
learned from studies of neutrino flavor change,  
using such formulas as —

# Probability of Flavor Change In Vacuum

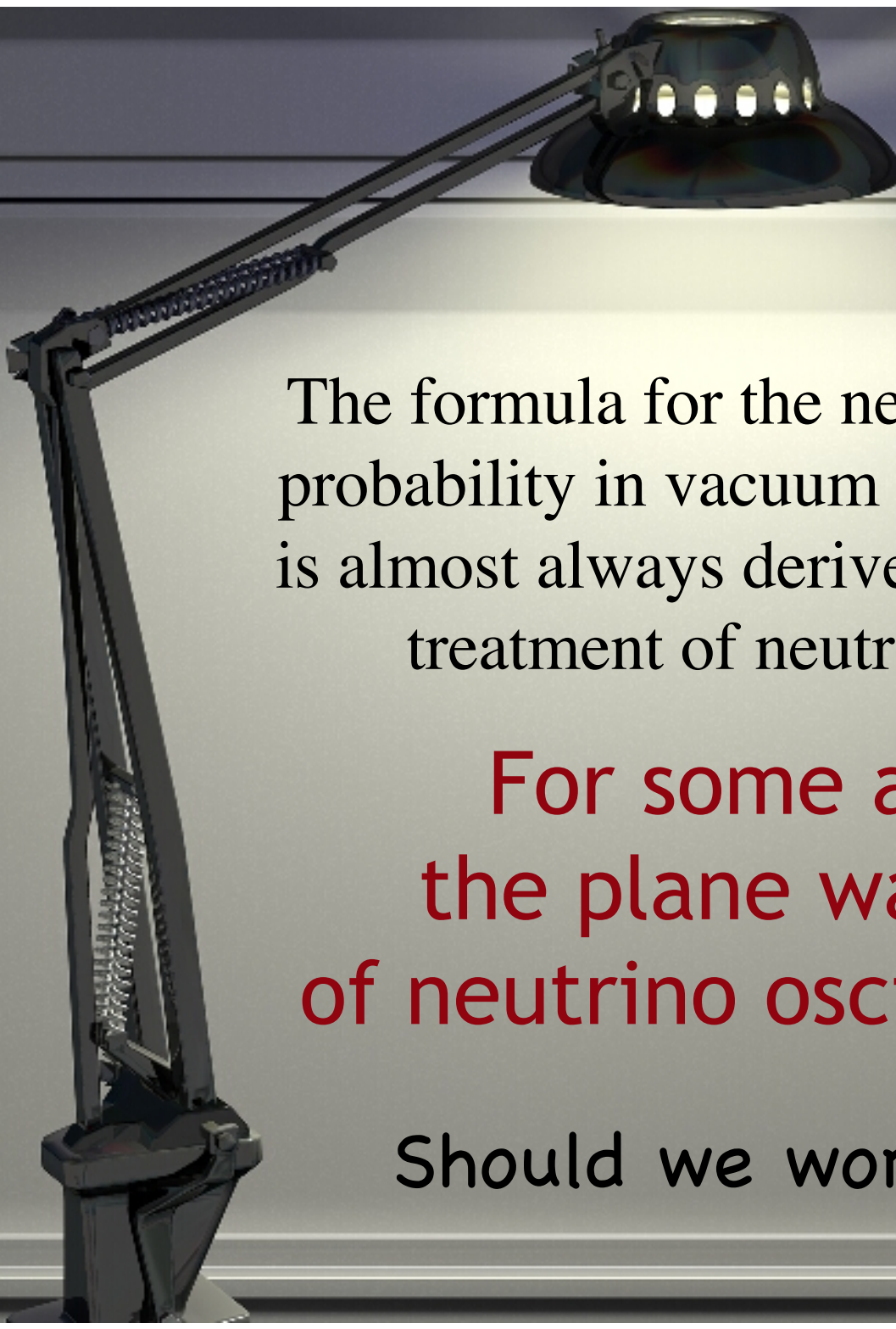
$$\begin{aligned}
 P\left(\bar{\nu}_{\alpha} \rightarrow \bar{\nu}_{\beta}\right) &= \\
 &= \delta_{\alpha\beta} - 4 \sum_{i>j} \text{Re}\left(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*\right) \sin^2\left(\Delta m_{ij}^2 \frac{L}{4E}\right) \\
 &\quad \pm 2 \sum_{i>j} \text{Im}\left(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*\right) \sin\left(\Delta m_{ij}^2 \frac{L}{2E}\right)
 \end{aligned}$$

Diagrammatic annotations:
 

- Arrows point from  $\alpha$  and  $\beta$  to the text  $e, \mu, \text{ or } \tau$ .
- An arrow points from "Mixing matrix" to the  $U$  terms in the first sum.
- An arrow points from "Distance" to  $L$  in the first sine argument.
- An arrow points from "Energy" to  $E$  in the first sine argument.

where  $\Delta m_{ij}^2 \equiv m_i^2 - m_j^2$ .





The formula for the neutrino flavor-change probability in vacuum (neutrino oscillation) is almost always derived using a plane-wave treatment of neutrino propagation.

For some applications, the plane wave treatment of neutrino oscillation is wrong.

Should we worry about that?

The probability of neutrino oscillation depends on the distance  $L$  between the neutrino source and the point of detection.

To determine  $L$ , we must know where the neutrino started, and where it was detected.

A plane wave has a definite, precise momentum  $p$ .

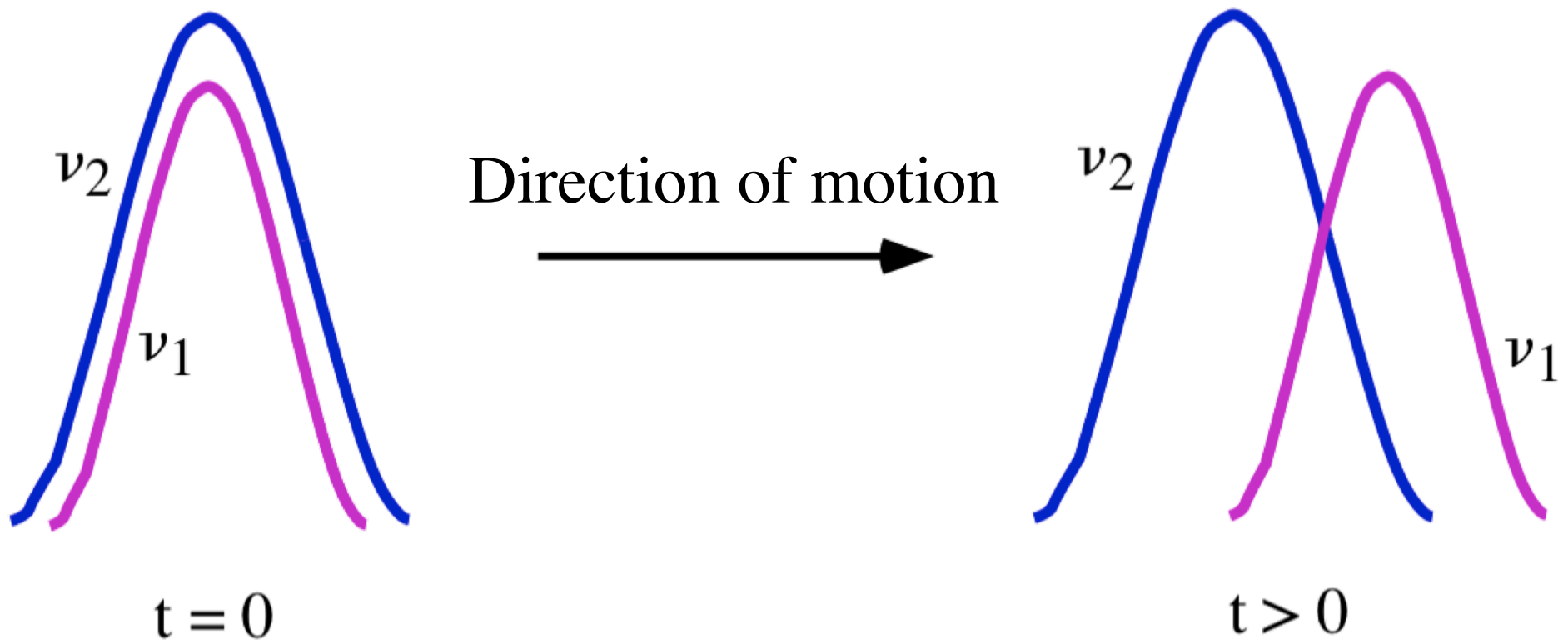
$$\text{Heisenberg: } \Delta x \Delta p \geq \hbar/2.$$

*If we know precisely the momentum with which a neutrino was born, we know nothing about where it was born.*

# The Wave Packet Picture

Each mass eigenstate is described by a wave packet.

Suppose  $\nu_2$  is heavier than  $\nu_1$ .



*Eventually the wave packets will separate.*

*No more oscillation.*

## *How soon do the wave packets separate??*

For accelerator neutrinos with energy  $E = 1$  GeV, and a wave packet width equal to the length of the pion decay region where the neutrinos are born, the bigger  $\Delta m^2 = 2.5 \times 10^{-3}$  eV<sup>2</sup> leads to wave packet separation in

**$10^{20}$  km.**

**This separation may be safely ignored!**

However, for supernova neutrinos from SN 1987A, with energy  $E \sim 10$  MeV, and a wave packet width equal to an *estimated* inter-nucleon distance within the star, separation occurs in

**$10^3$  km.**

*Supernova neutrinos are no longer oscillating  
when they reach us.*

*Different mass eigenstates produced at the  
same instant arrive at separate times,  
depending on their individual speeds.*

*The arrival time difference for the SN 1987 A  
neutrinos could have been  $\sim 10^{-4}$  sec.*

# Probability of Flavor Change In Vacuum

*For experiments with terrestrial neutrinos, use this formula with confidence.*

$$\begin{aligned}
 P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta) &= \\
 &= \delta_{\alpha\beta} - 4 \sum_{i>j} \text{Re}(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin^2 \left( \Delta m_{ij}^2 \frac{L}{4E} \right) \\
 &\quad \pm 2 \sum_{i>j} \text{Im}(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin \left( \Delta m_{ij}^2 \frac{L}{2E} \right)
 \end{aligned}$$

# Theoretical Ideas

Most theorists suspect that neutrinos, and only neutrinos, have **Majorana masses**.

If this suspicion is right, then the origin of neutrino masses is different from the origin of the masses of all other known particles.

A majorana mass term, such as —

$$\mathcal{L}_{\text{Mass}} = -\frac{m_R}{2} \left[ \overline{(\nu_R)^c} \nu_R + \overline{\nu_R} (\nu_R)^c \right],$$

Mass charge conjugation

causes the transitions  $\nu \rightarrow \bar{\nu}$  and  $\bar{\nu} \rightarrow \nu$  between the underlying neutrino  $\nu$  and its antiparticle.

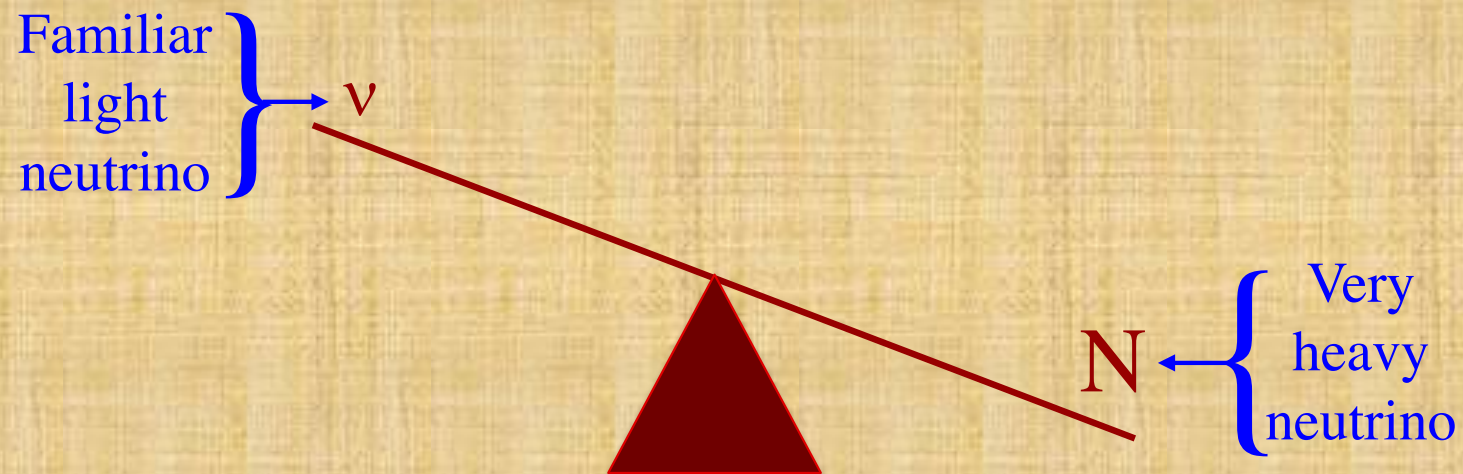
The mass eigenstate is then  $\nu + \bar{\nu}$ , since  $\nu + \bar{\nu} \rightarrow \bar{\nu} + \nu$ .

We see that the mass eigenstate is a  
Majorana (self-conjugate) neutrino.

Majorana mass terms make possible the simple version of the **See-Saw Mechanism** for neutrino masses.



# The See-Saw Mechanism



(Gell-Mann, Ramond, and Slansky; Yanagida;  
Mohapatra and Senjanovic; Minkowski)

The straightforward (type-I) See-Saw model adds to the SM 3 heavy neutrinos  $N_i$ , with —

$$\mathcal{L}_{\text{new}} = -\frac{1}{2} \sum_i m_{N_i} \overline{N_{iR}^c} N_{iR} + \sum_{\substack{\alpha=e,\mu,\tau \\ i=1,2,3}} y_{\alpha i} \left[ \overline{\nu}_{\alpha L} \overline{H^0} - \overline{\ell}_{\alpha L} H^- \right] N_{iR} + h.c.$$

Large Majorana masses

Charge conjugate

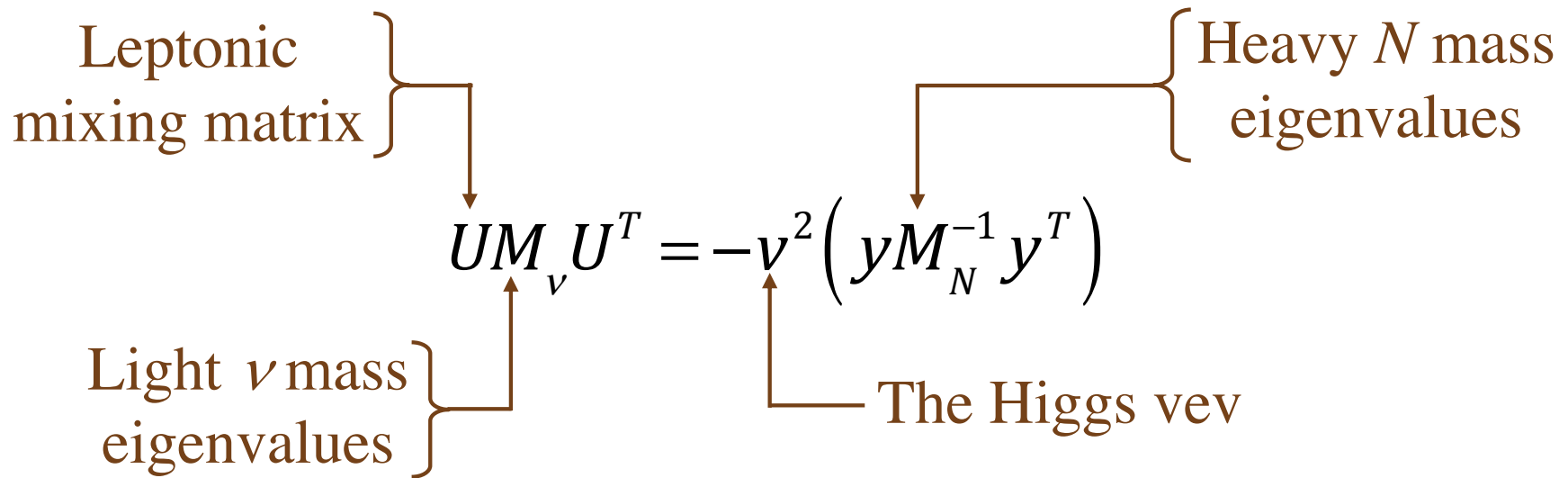
SM lepton doublet

SM Higgs doublet

Yukawa coupling matrix

Exploring the physics of this model leads to  
The See-Saw Relation.

# The See-Saw Relation



$$\left( \underbrace{UM_\nu U^T}_{\text{Outputs}} = -v^2 \underbrace{(yM_N^{-1}y^T)}_{\text{Inputs, in } \mathcal{L}} \right)$$

The See-Saw partner neutrinos  $N_i$ , although heavy, would have been made during the *hot* Big Bang.

This makes possible —

## *Leptogenesis*

(an explanation of the baryon-antibaryon asymmetry of the universe)

First, CP violation in  $N_i$  decays converts  $L = 0$  into  $L \neq 0$ .

Then, the Standard Model Sphaleron process converts part of this non-zero lepton number into a non-zero baryon number  $B$ .

The key ingredients of Leptogenesis are —

*CP violation among the leptons*

*Confirm via observation of  
CP violation in neutrino oscillation.*

*Non-conservation of Lepton Number  $L$*

*Confirm via observation of  
neutrinoless double beta decay.*

*The heavy neutrinos  $N_i$  may well be far too heavy to observe experimentally.*

*However, generically, leptogenesis and light-neutrino ~~CP~~ imply each other.*

*They both come from phases in the same Yukawa coupling matrix  $y$ .*

# The Oscillation — Leptogenesis Connection

(B.K.)

## The See-Saw Relation

Leptonic mixing matrix

Light  $\nu$  mass eigenvalues

Heavy  $N$  mass eigenvalues

The Higgs vev, a real number

$$UM_{\nu}U^T = -v^2 \left( yM_N^{-1}y^T \right)$$

Through  $\mathbf{U}$ , the phases in  $\mathbf{y}$  lead to  $\cancel{\mathcal{CP}}$  in light neutrino oscillation.

$$\begin{aligned}
 P\left(\bar{\nu}_{\alpha} \rightarrow \bar{\nu}_{\beta}\right) &= \\
 &= \delta_{\alpha\beta} - 4 \sum_{i>j} \operatorname{Re}\left(U_{\alpha i}^{*} U_{\beta i} U_{\alpha j} U_{\beta j}^{*}\right) \sin^2\left(\Delta m_{ij}^2 \frac{L}{4E}\right) \\
 &\quad \left(\begin{array}{c} + \\ - \end{array}\right) 2 \sum_{i>j} \operatorname{Im}\left(U_{\alpha i}^{*} U_{\beta i} U_{\alpha j} U_{\beta j}^{*}\right) \sin\left(\Delta m_{ij}^2 \frac{L}{2E}\right)
 \end{aligned}$$



*If the oscillation  $\mathcal{CP}$  phase  $\delta$  proves to be large, it could explain almost the entire Baryon – Antibaryon asymmetry by itself.  
(Pascoli, Petcov, Riotto)*

# Future and Ongoing Experiments

## Long Baseline Accelerator Neutrinos

CP Violation, Neutrino Mass Ordering,  $\theta_{23}$ ,  
Non-Standard Interactions, Sterile Neutrinos,  
Extra Dimensions, Lorentz or CPT Violation,  
Atmospheric, Solar, and Supernova Neutrinos

## Long Baseline Reactor Neutrinos

Neutrino Mass Ordering, Atmospheric,  
Solar, and Supernova Neutrinos

## Short Baseline Accelerator, Reactor, and Radioactive-Source Neutrinos

1 eV Sterile Neutrinos: Yes or No?

## Neutrinoless Double Beta Decay

Neutrinos: Dirac or Majorana?

## Beta Spectrum in Beta Decay

Neutrino Mass Determination

## Coherent Neutrino-Nucleus Scattering

Non-Standard Interactions?

## Searches for Heavy Neutrinos

keV, MeV, GeV, TeV Neutrinos?

## Neutrino Telescopes

Astronomy with Neutrinos,  
General Neutrino Physics

## Cosmological Observations

Sum of Neutrino Masses,  
Number of Neutrino Flavors

*The future program  
is rich.*

*We look forward  
to the results.*