



# Axion-like Particle Dark Matter & Small-scale Structure

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# Two models of CDM

## WIMPs

Predicted ~1970's

Naturally come from SUSY.

Heavy, nucleon mass.

Thermal.

“Traditional” search strategies.

...

Cold due to small thermal  
velocities

Collisions necessary for  
production

## Axions

Predicted ~1970's

Come from strong-CP.

v. light, sub-neutrino mass.

Non-thermal.

“Non-traditional” search strategies.

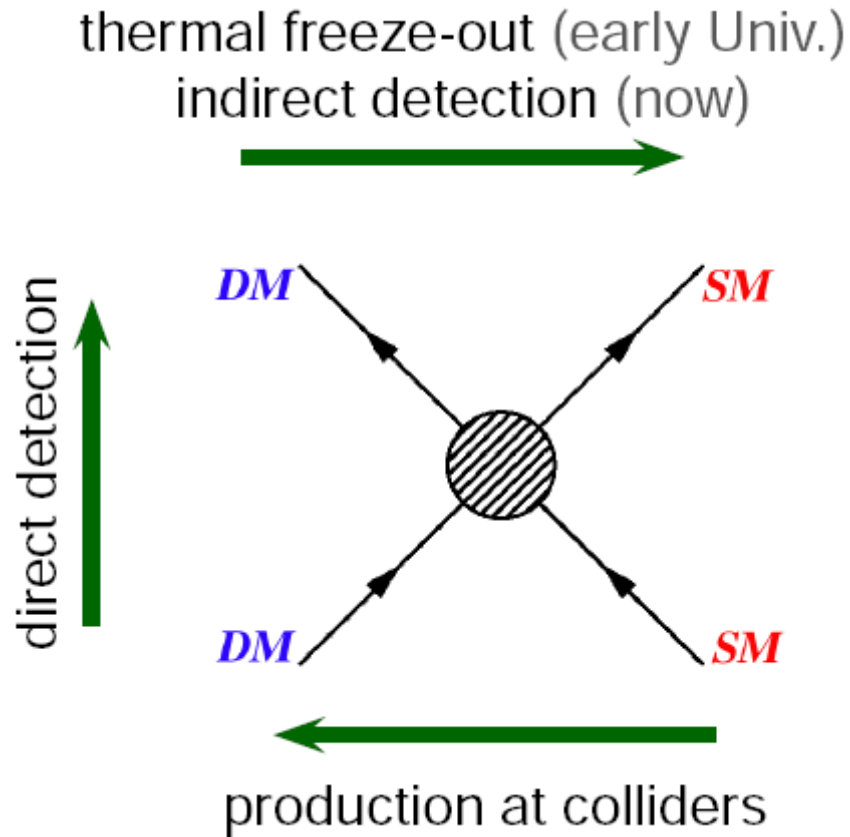
...

“Cold” due to coherence  
and small wavelength

Interactions via wave  
equations

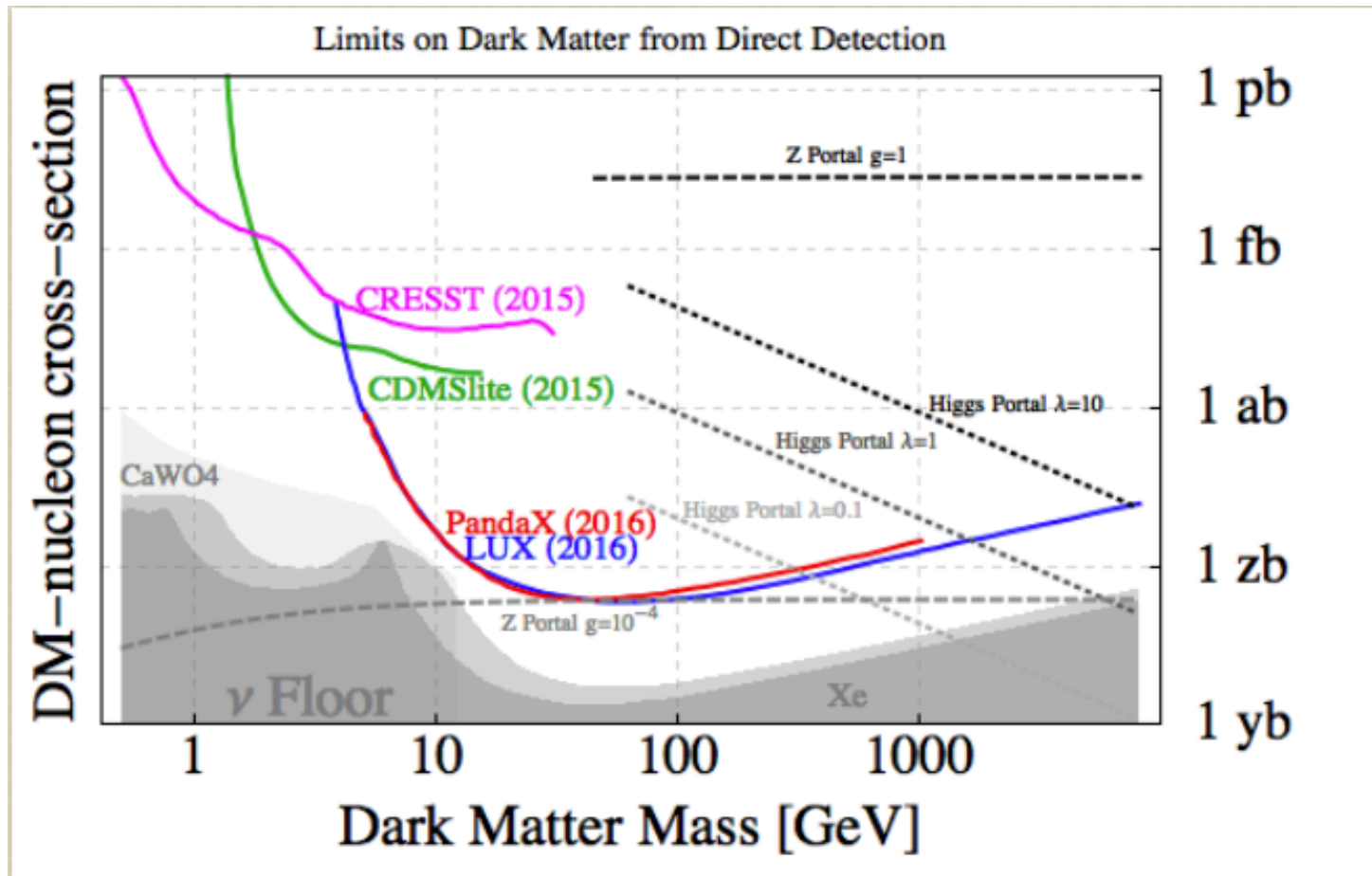
*Axions are just a type of CDM, with different  
extreme limits than thermal DM*

# On Axions and WIMPs



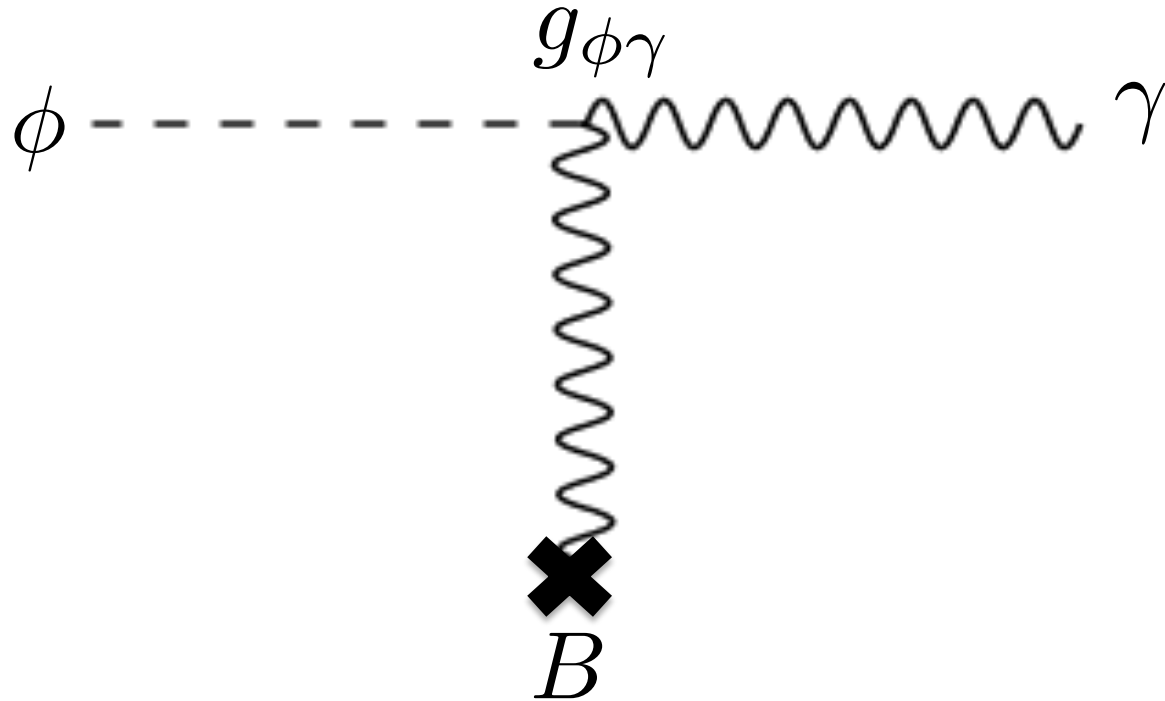
$$Z_{\mu} \bar{\psi} \gamma^{\mu} \psi$$

# On Axions and WIMPs



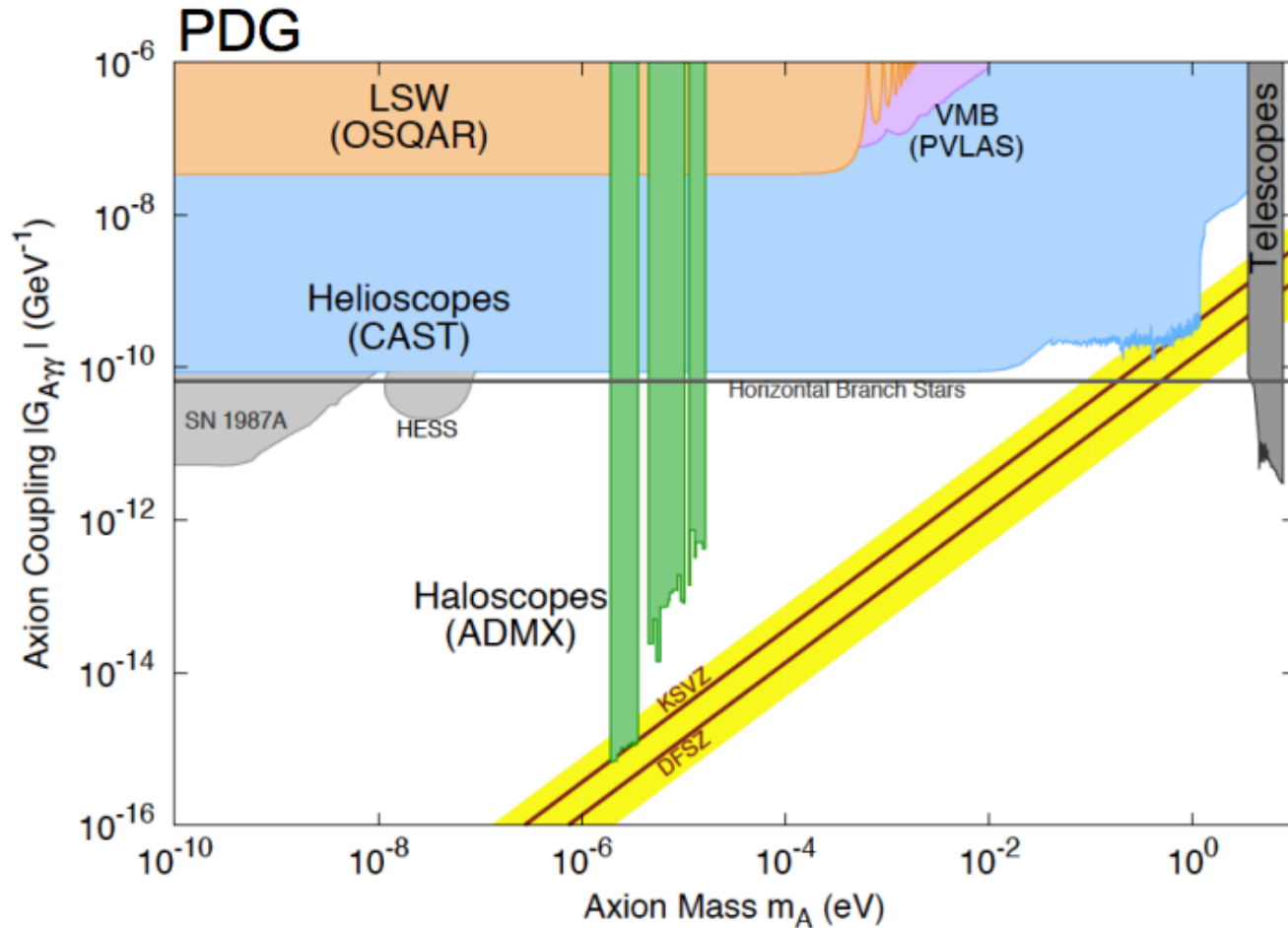
Predicted interaction strength for WIMPs has already been excluded. Soon experiments will hit the “neutrino floor”.

# On Axions and WIMPs



$$\phi \epsilon_{\mu\nu\alpha\beta} F^{\mu\nu} F^{\alpha\beta}$$

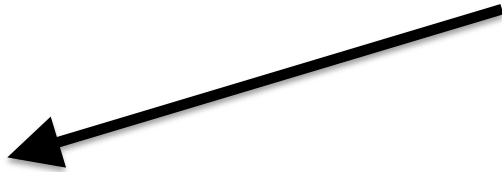
# On Axions and WIMPs



The QCD axion is essentially a one-parameter model.  
Only one experiment has got into the range to probe DM.

# What makes an ALP?

An ALP is a **non-thermally produced classical scalar field**.  
ALPs differ from CDM in **initial conditions & dynamics**.



**Symmetry breaking** leading to axion relic density

- Dense DM relics
- “**Miniclusters**”
- **Microlensing** constraints a la PBHs
- Important for the standard **QCD axion**
- Smallest DM structures

**Uncertainty principle**

- Small-scale coherence and dynamic halos
- Suppression of structure
- Formation of solitonic “**axion stars**”
- Pronounced for ultralight “**Fuzzy DM**”,  $m \sim 10^{-22}$  eV
- Lightest DM particle

“Miniclusters”: dense clumps from initial conditions (c.f. MACHOs). Sub-lunar mass. Classic QCD axion window.

“Fuzzy DM”: diffuse due to macroscopic wavelength (c.f. warm DM). Dwarf galaxy scales. String theory axions?



# The Life of Axions



# Spontaneous Symmetry Breaking

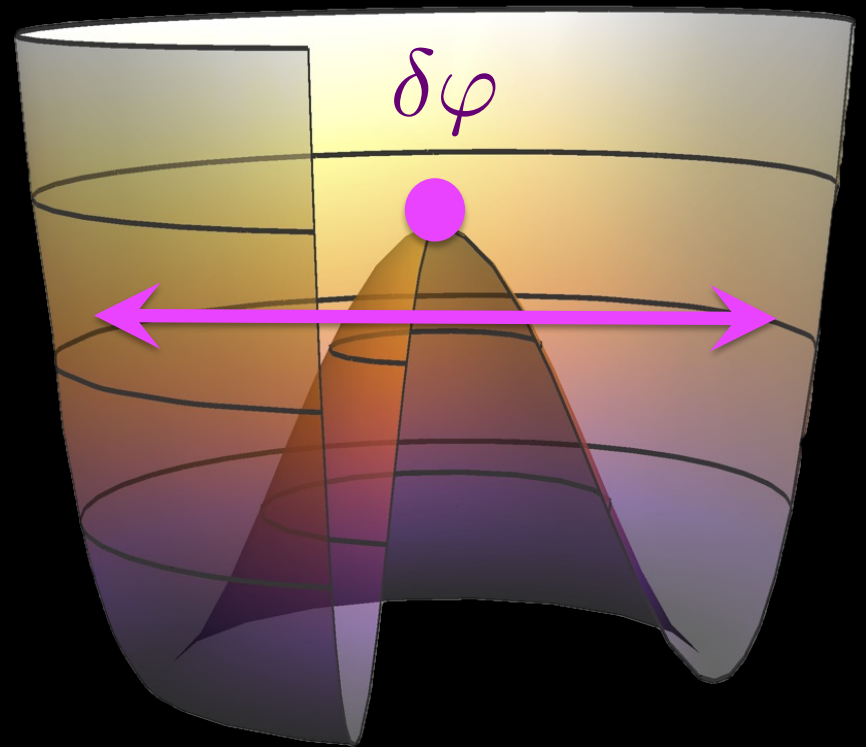
Spontaneous Symmetry Breaking  $\rightarrow$  (p)NGB

The “decay constant” determines temperature of phase transition:

$$\delta\varphi \sim T \sim \frac{H_I}{2\pi}$$

$$T \gg f_a$$

$$\langle\varphi\rangle = 0$$



# Spontaneous Symmetry Breaking

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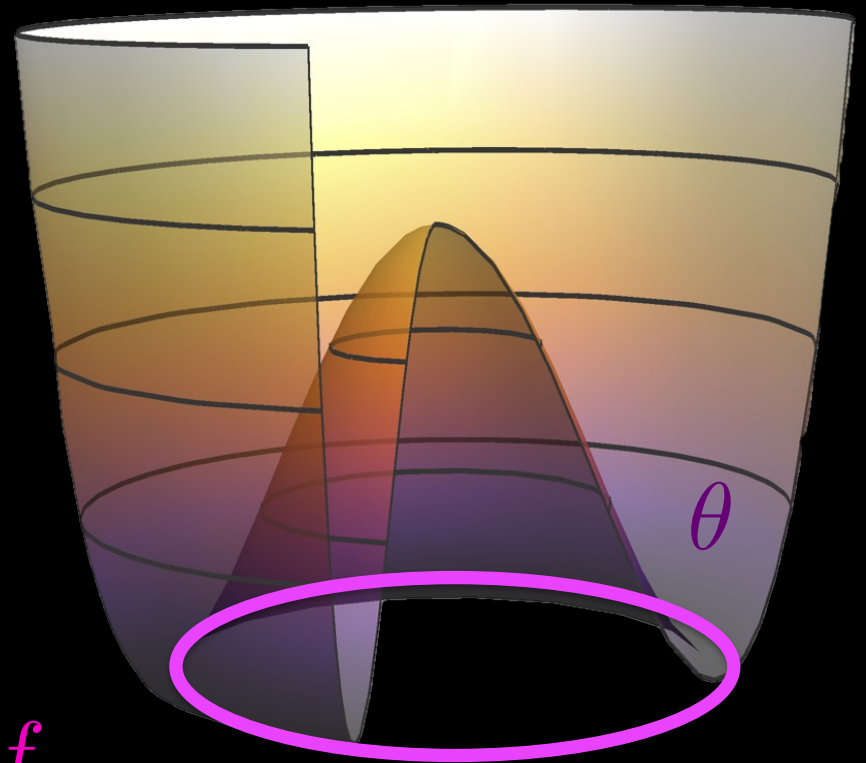
$$T \ll f_a$$

$$\langle\varphi\rangle = f_a/\sqrt{2}$$

The axion is born:  $\theta = \phi/f_a$

Symmetry breaking  $\rightarrow$  relics

$$\theta \in \mathcal{U}[-\pi, \pi]$$



# Vacuum Realignment

Axion acquires mass, evolves according to Klein-Gordon:

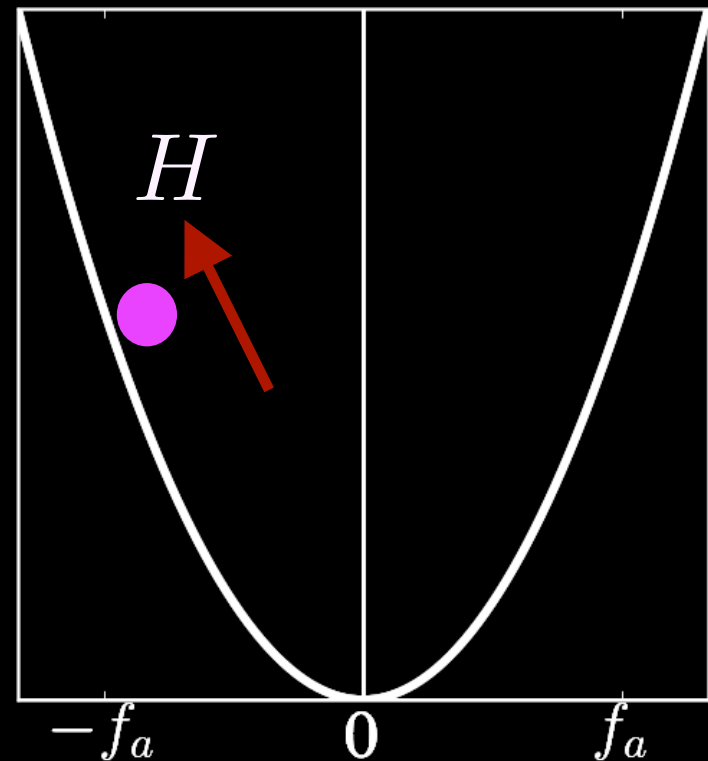
$$\ddot{\phi} + 3H\dot{\phi} + m_a^2\phi = 0$$

$$H \gg m_a$$

Axion is “frozen” by Hubble friction term.

$$\Rightarrow \rho_a \approx \text{const.}$$

$$\Rightarrow w_a \approx -1$$



# Vacuum Realignment

Axion acquires mass\*, evolves according to Klein-Gordon:

$$\ddot{\phi} + 3H\dot{\phi} + m_a^2\phi = 0$$

$$H \ll m_a$$

Field oscillates & damps.

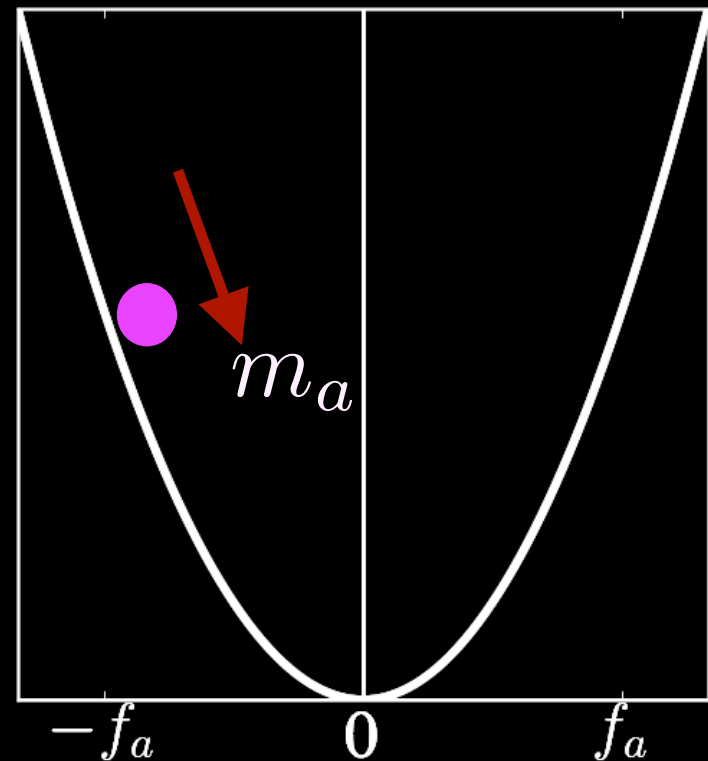
WKB (or exact)  $\rightarrow$

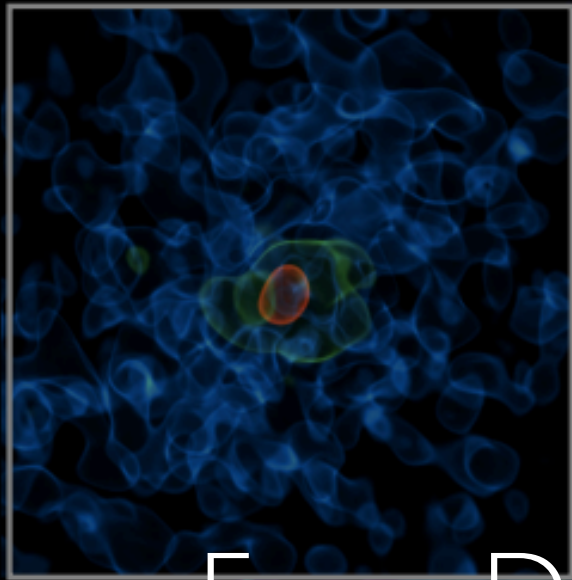
$$\rho_a \approx \rho_a(a_{\text{osc}})a^{-3}$$

Homogeneous scalar  $\sim$   
matter

Inhomogeneities  $\rightarrow$

gradients  $\rightarrow$  pressure





# Fuzzy Dark Matter



9 kpc/h



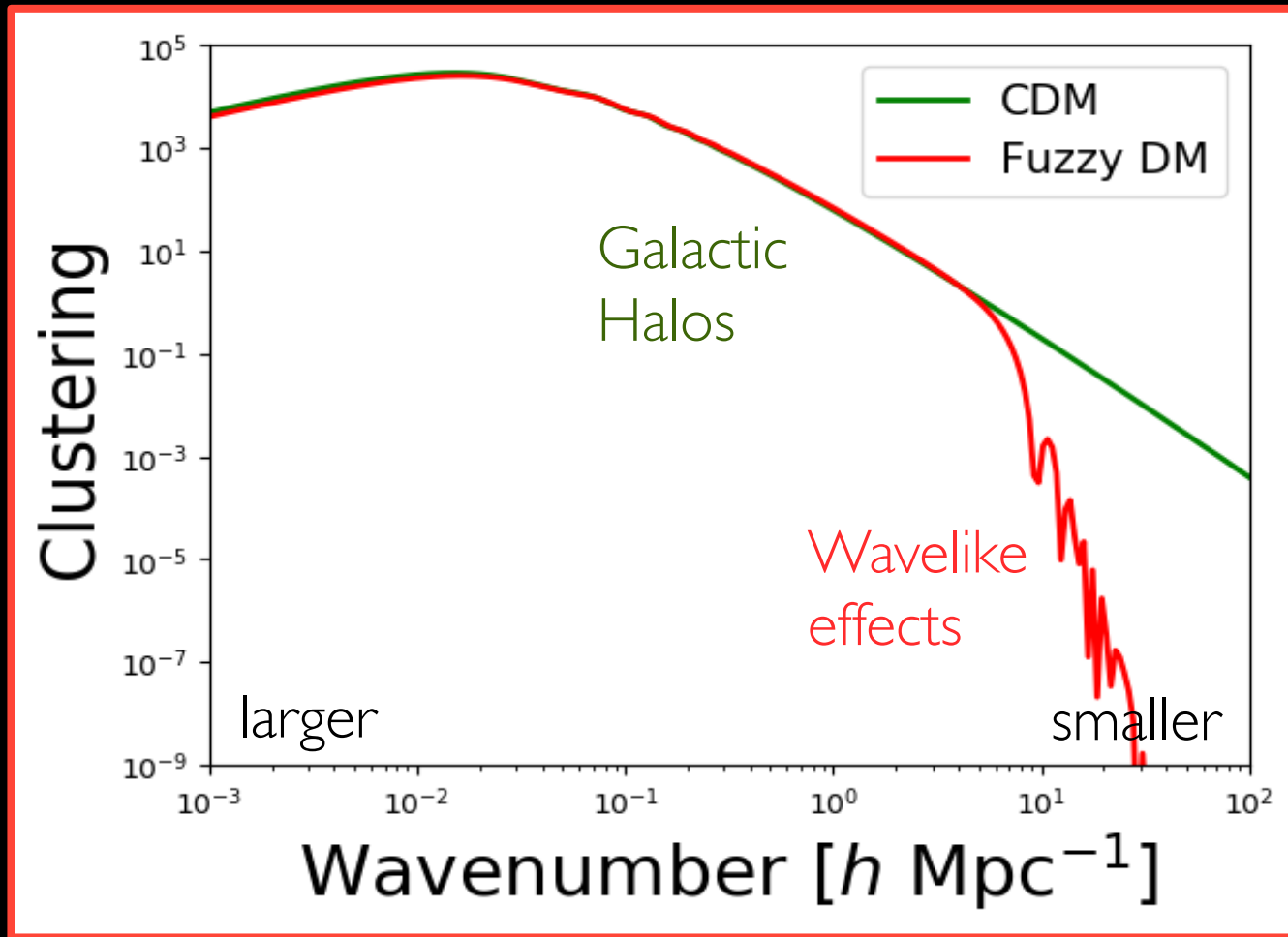
DJEM, Physice Reports (2016)

Hui et al PRD (2017)

Image: Veltmaat et al (2018); FDM Simulation

$z = 1.07$   
2.5 Mpc/h

Light axions depart from standard CDM in their dynamics. De Broglie wavelength suppresses formation of structure.

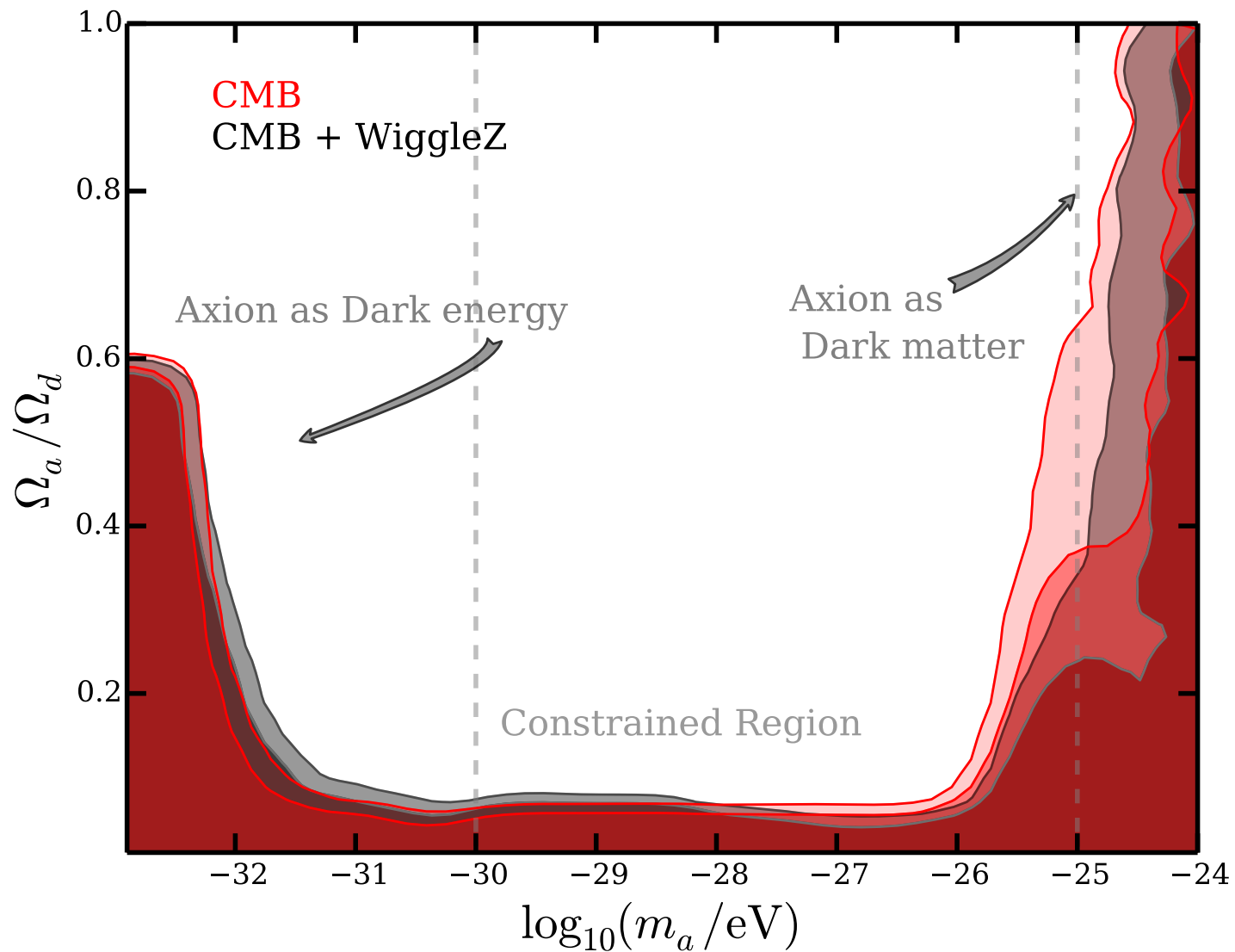


e.g. DJEM (2016)

$$k_{J,\text{eq}} = 9(m_a/10^{-22} \text{ eV})^{1/2} \text{ Mpc}^{-1}$$

# CMB Baseline

Amendola & Barbieri (2006)  
Hlozek, DJEM et al (2015) *axionCAMB*





# Non-linear Scales

e.g. Widrow & Kaiser (1993); Chavanis (2011+);  
DJEM (2015,2016); Hui et al (2016)

Fundamentally different from CDM/WDM/SIDM

Non-rel limit of Klein-Gordon Einstein  $\rightarrow$  Schrodinger-Poisson

$$i\dot{\psi} + \frac{1}{2m_a^2} \nabla^2 \psi - m_a \Phi \psi = 0; \quad \nabla^2 \Phi = 4\pi G_N |\psi|^2$$

Related to the smoothed Vlasov equation. Field equation not a particle distribution function  $\rightarrow$  “non-linear optics” regime.

Madelung transformation (polar co-ords)  $\rightarrow$  fluid system:

$$\dot{\delta} + \vec{v} \cdot \nabla \delta = (1 + \delta) \nabla \cdot \vec{v}$$

continuity

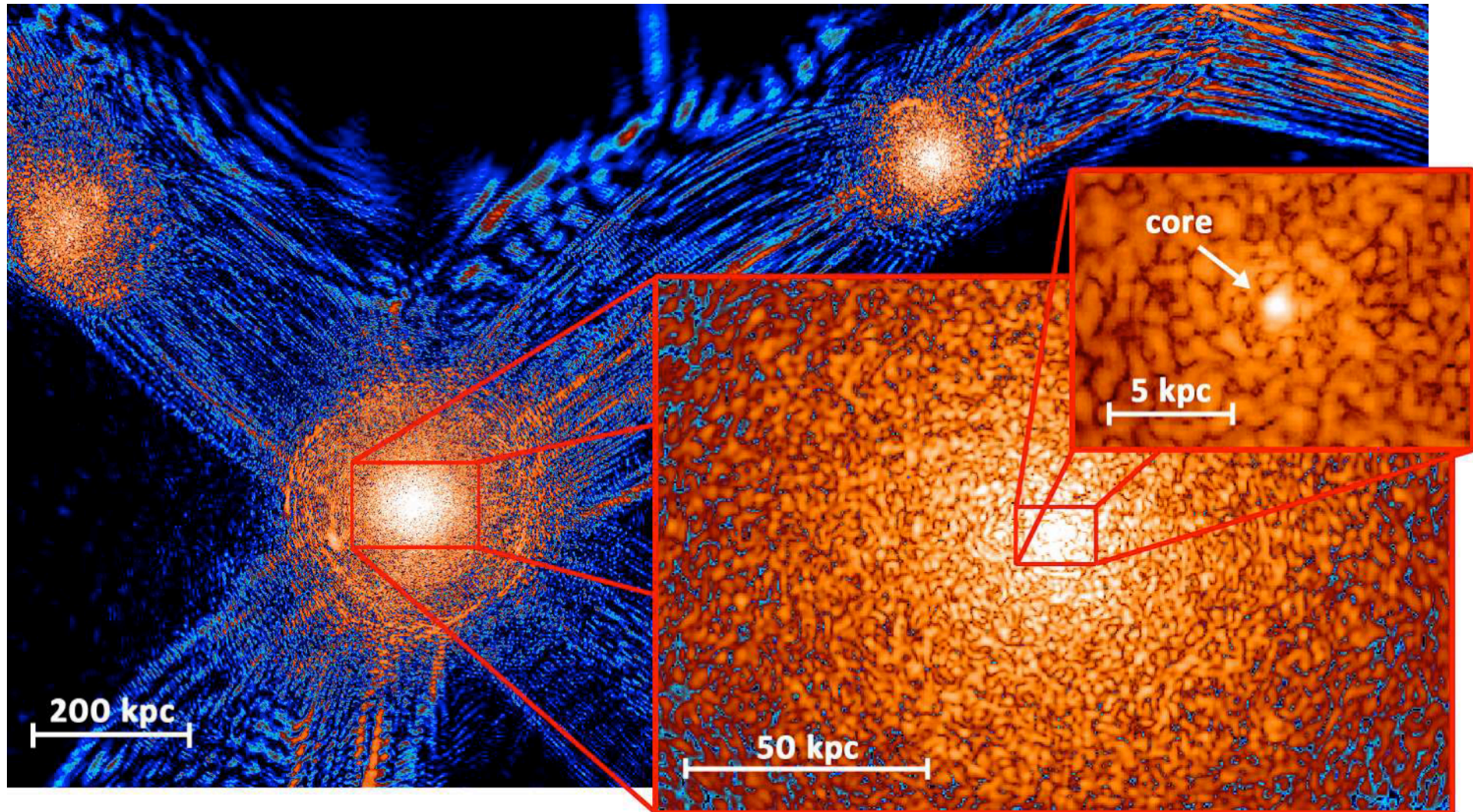
$$\dot{\vec{v}} + (\vec{v} \cdot \nabla) \vec{v} = -\nabla(\Phi + Q)$$

Euler

$$Q = -\frac{1}{2m^2} \frac{\nabla^2 \sqrt{1 + \delta}}{\sqrt{1 + \delta}}$$

Quantum Pressure : source  
of interference effects

# FDM Simulations



Schive et al (2014)

# FDM Simulations



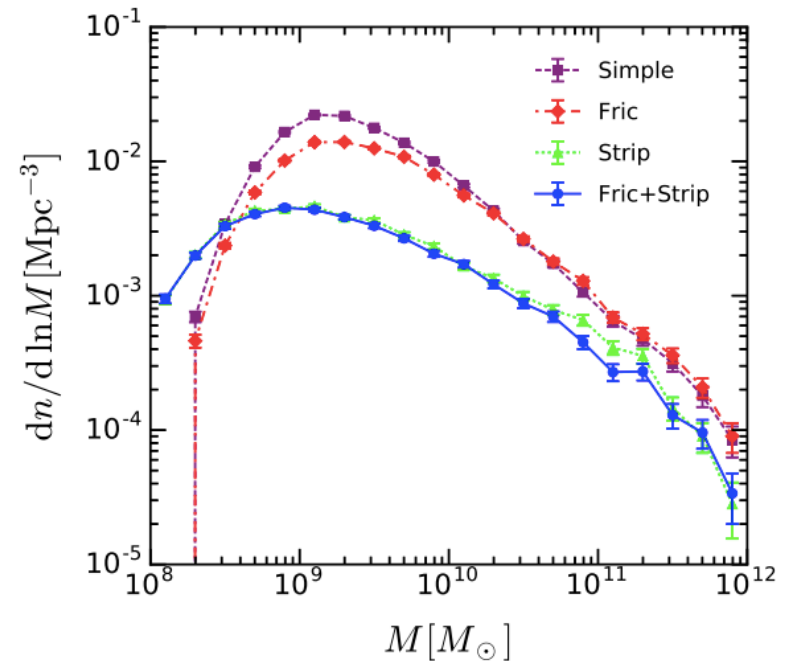
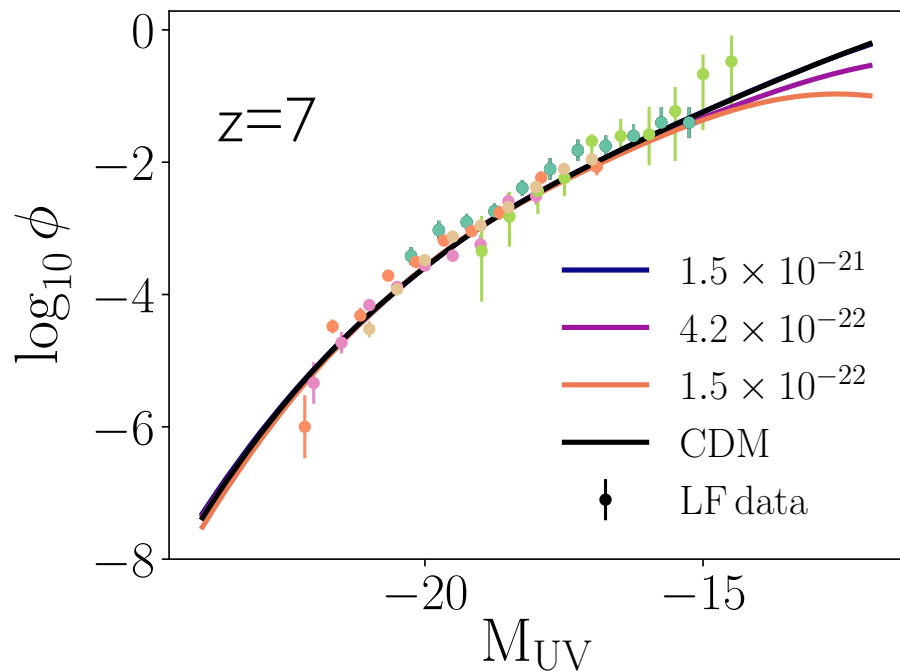
Schwabe et al (2016). Blob size  $\sim$  kpc, lifetime  $\sim 10^6$  years ( $10^{-22}$  eV/m)

# Halo Mass Function

DJEM & Silk (2014); Du et al (2016)  
Corasaniti, DJEM et al (2017)

Suppressed power  $\rightarrow$  no low mass halos/subhalos.

$$M_c \approx 10^8 (m/10^{-22} \text{eV})^{-1.35} M_\odot$$



In the field: no high- $z$  low mass halos  $\rightarrow$  reduced early star formation  $\rightarrow$  later reionization.

In the MW: no low mass satellites / subhalos. Constrained by tidal streams?

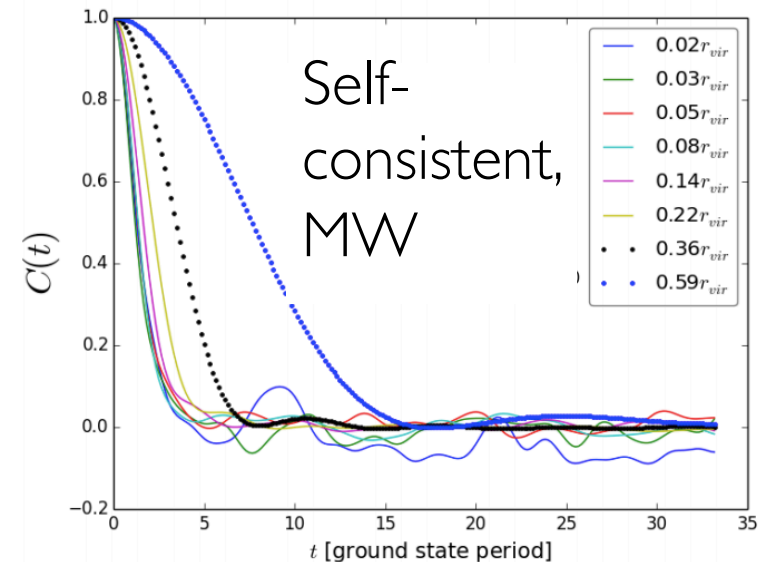
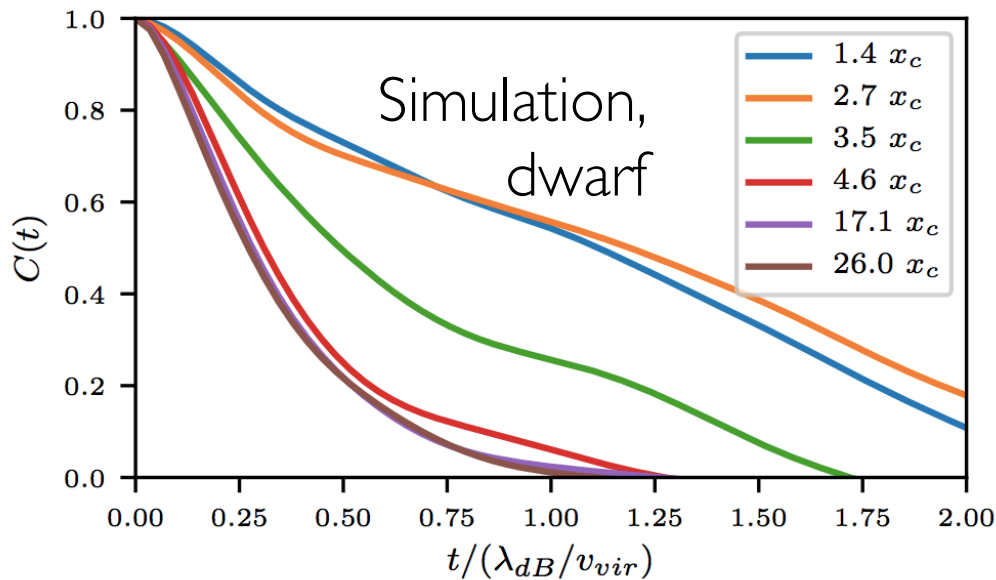
# Dynamical Effects

Hui et al (2017); Veltmaat et al (2018)  
 Lin et al (2018); Khlemeninsky & Rubikov (2014)

The FDM halo is **not static** on times  $< 1/mv^2 \sim 10^6$  years.

**Pressure oscillations** on Compton times  $\rightarrow$  pulsar timing.

“Wavelets”  $\rightarrow$  quasiparticles and **dynamical relaxation**.

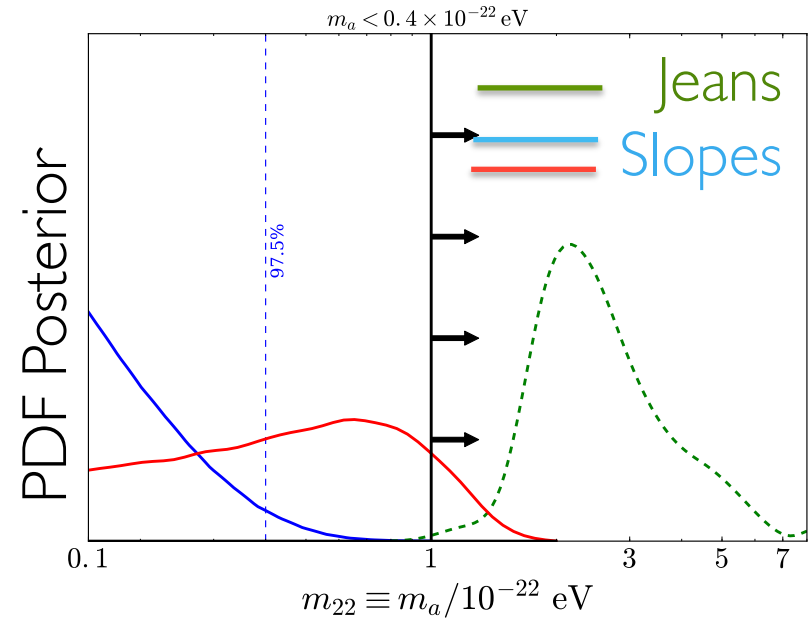
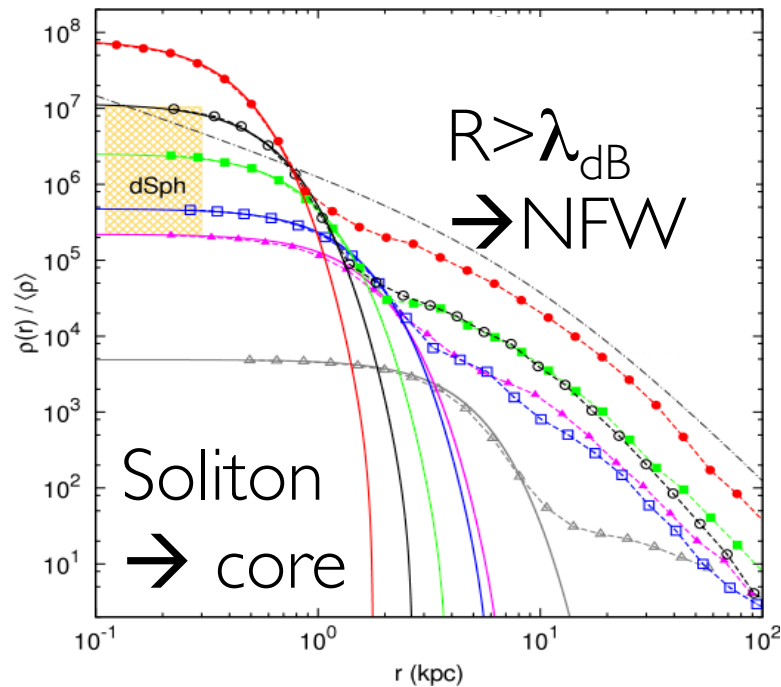


$$t_{\text{relax}}(r) \sim \frac{0.4}{f_{\text{relax}}} \frac{m^3 v^2 r^4}{\pi^3 \hbar^3} \sim \frac{1 \times 10^{10} \text{ yr}}{f_{\text{relax}}} \left( \frac{v}{100 \text{ km s}^{-1}} \right)^2 \left( \frac{r}{5 \text{ kpc}} \right)^4 \left( \frac{m}{10^{-22} \text{ eV}} \right)^3$$

# Axion Stars and Cores

Schive et al (2014+); Veltmaat et al (2018)  
Gonzalez-Morales, DJEM et al (2017)

Density profiles show **prominent cores**. Pressure supported solitons.

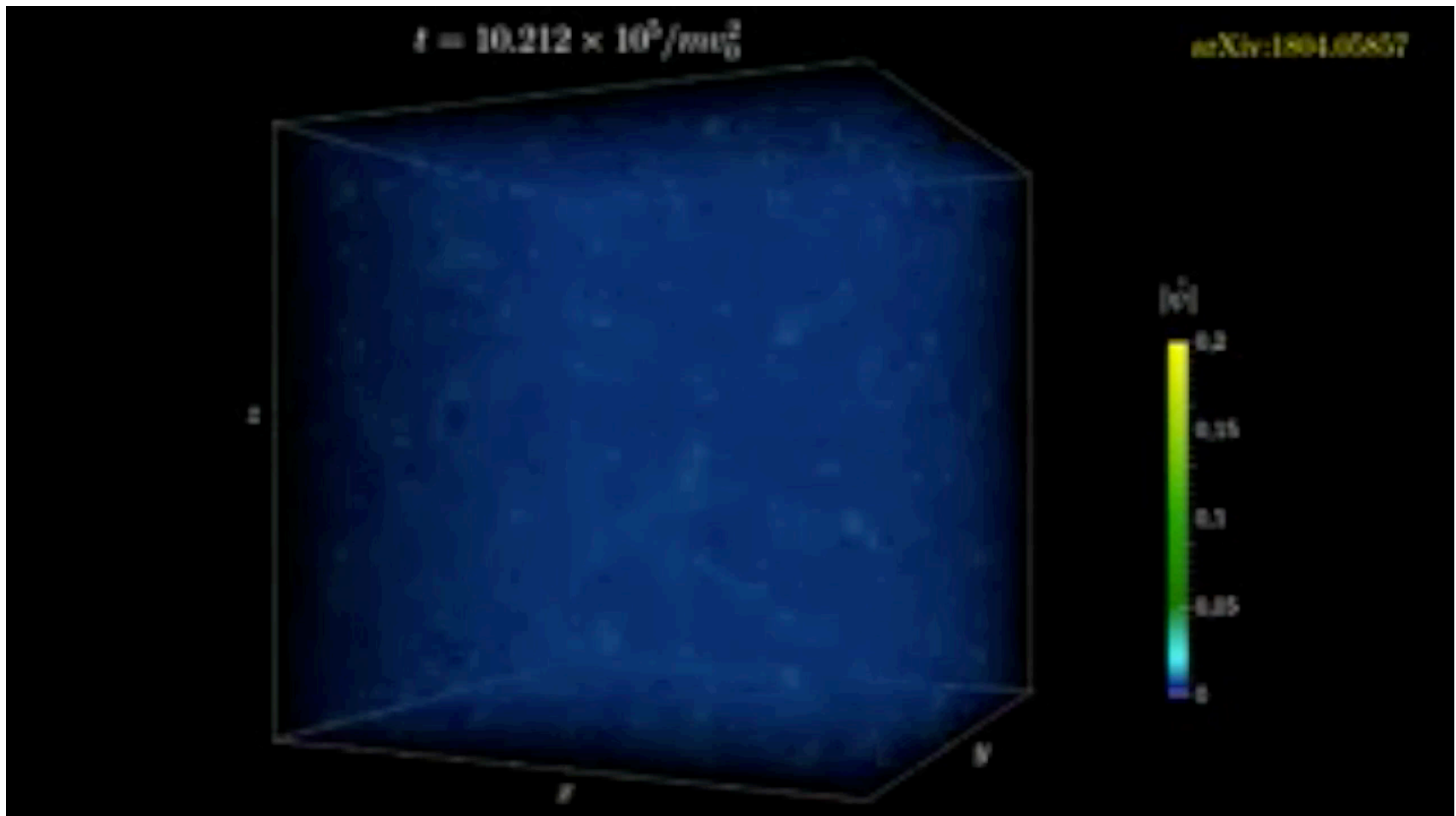


Solitons alone cannot explain dSph cores: "Catch 22" as WDM.  
Other observational effects of enhanced density?  
Recent simulations seem to show **strong core oscillations**.

# Axion Star Birth & Death

Levkov et al (2016)  
Helfer, DJEM et al (2016)

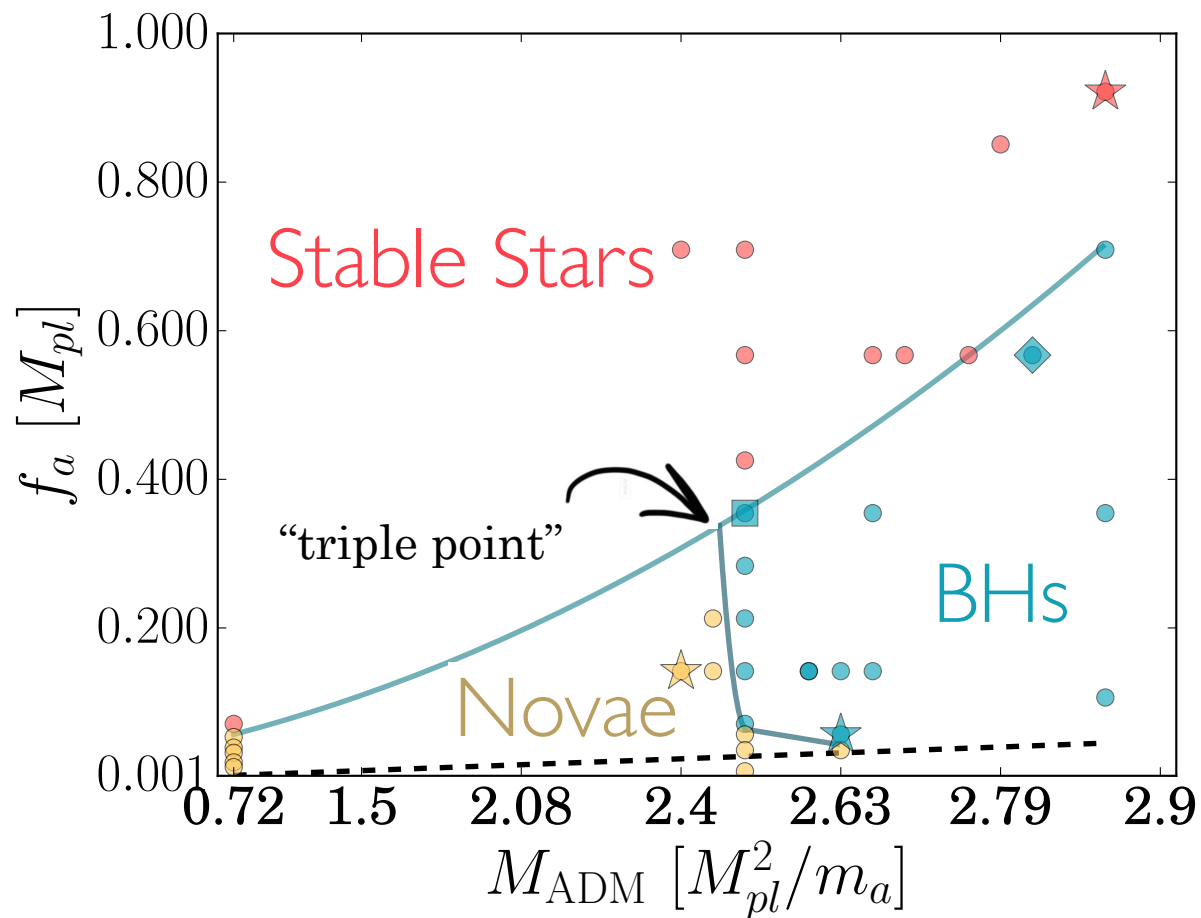
Axion stars are **relevant beyond FDM**. Should form in all models.



# Axion Star Birth & Death

Levkov et al (2016)  
Helfer, DJEM et al (2016)

Core-halo mass relation  $\rightarrow$  growth quenches for FDM. Are there relativistic axion stars? Axion novae? ECOs and GWs?

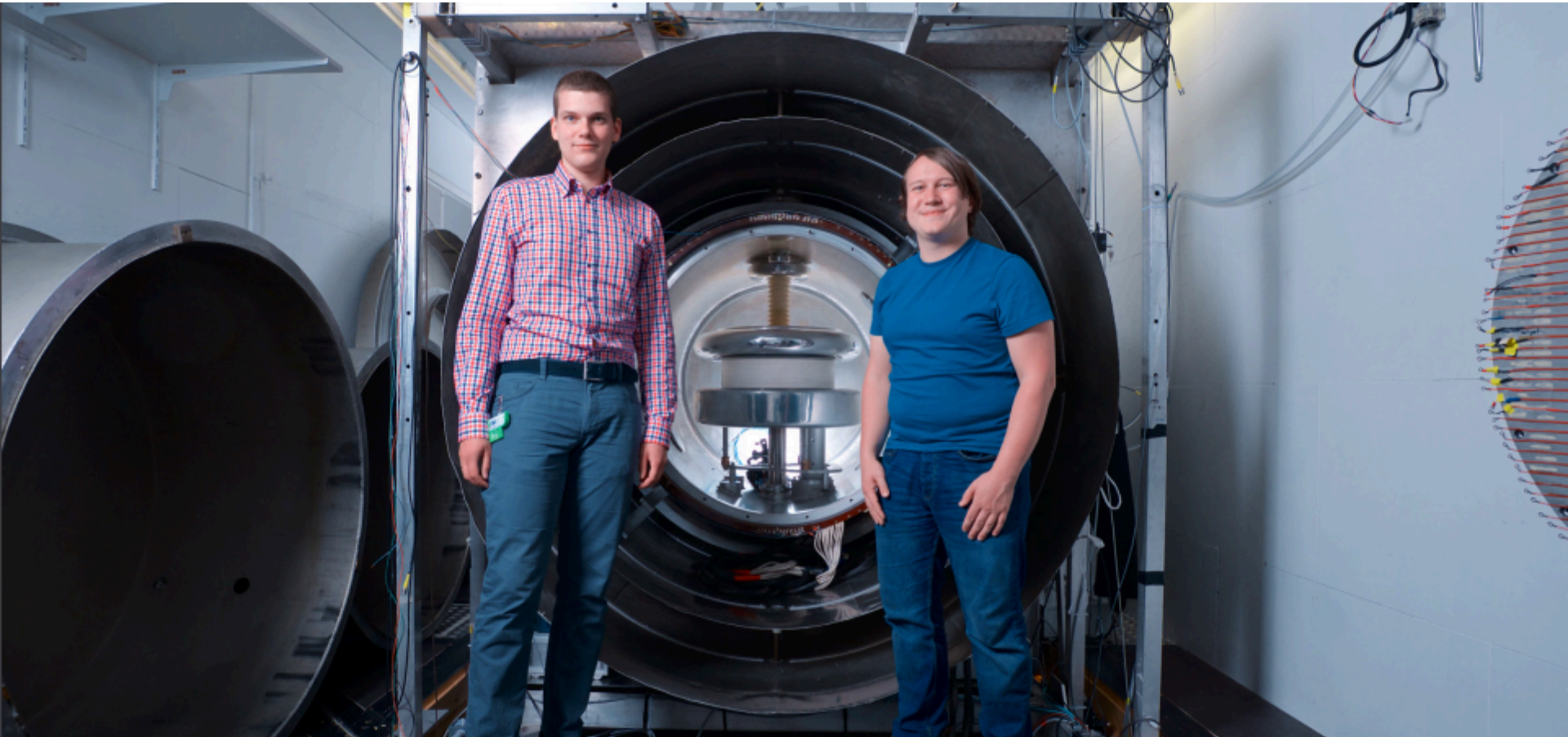




# FDM Direct Detection

Abel, DJEM et al (2017)  
Students: Michal Rawlik, Nick Ayres

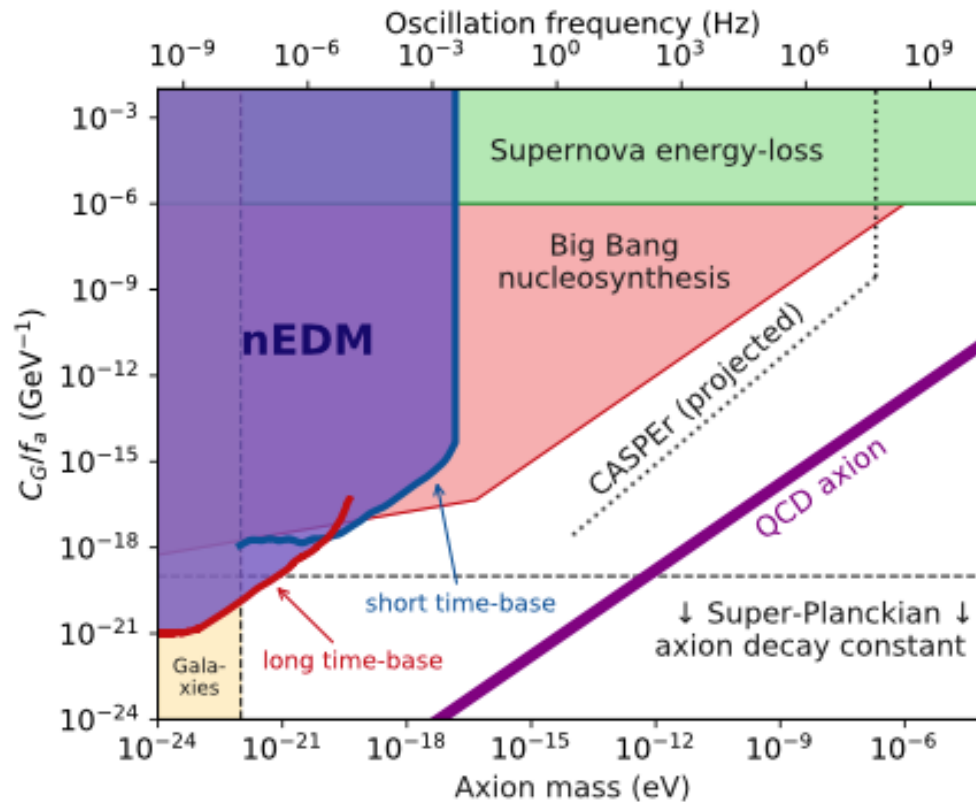
Detection relies on mass → Compton frequency. FDM  $\sim 10^{-7}$  Hz.  
Neutron EDM @ PSI and ILL measured for '98-'02 and '15-'16.  
→ First lab constraints on axions at this frequency (scalar DM easier)



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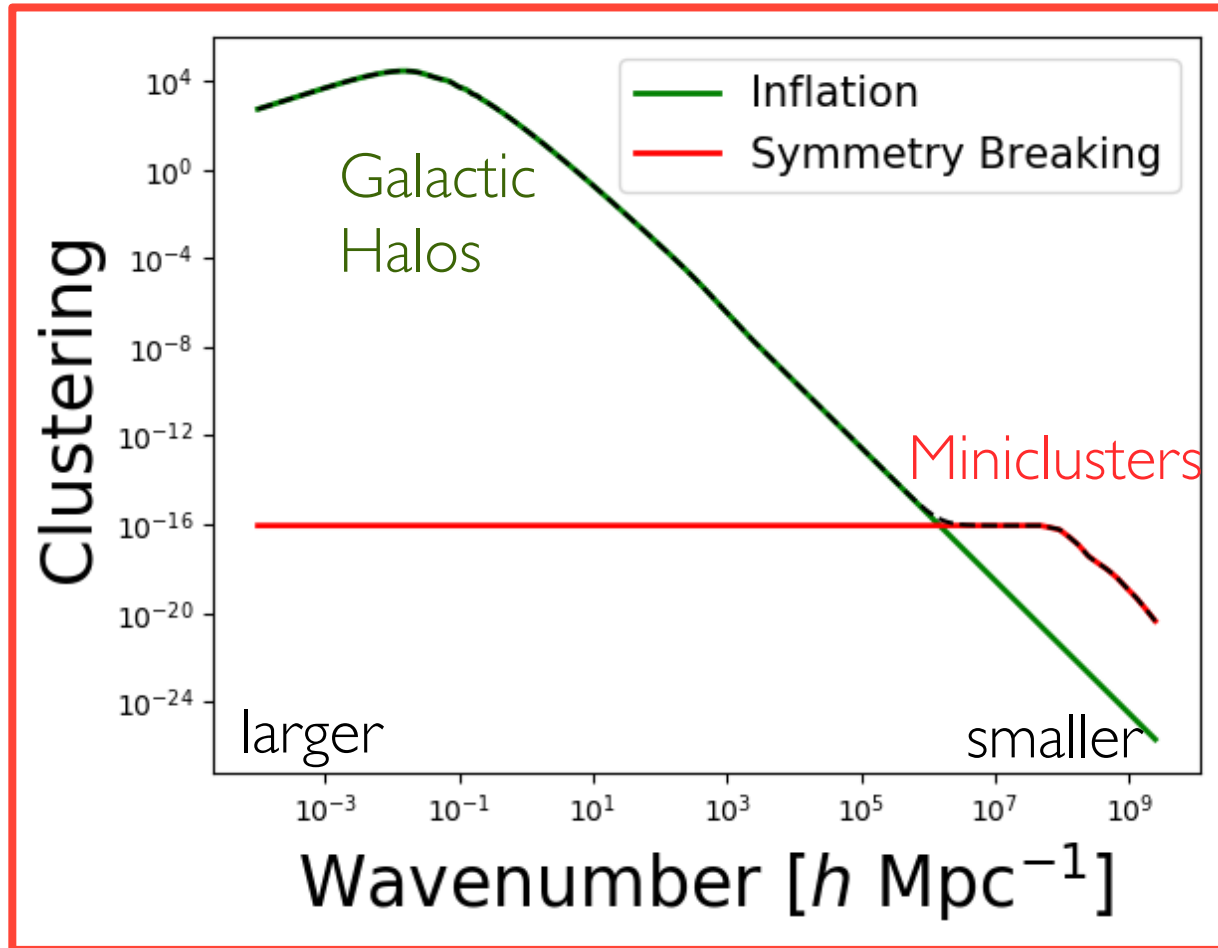


# Miniclusters & Microlensing

Fairbairn, DJEM, Quevillon, PRL (2017);  
Fairbairn, DJEM, Quevillon, Rozier, PRD (2017)

Image: the Subaru Hyper Suprime Cam

Miniclusters depart from standard CDM in **initial conditions**.  
Post-inflation symmetry breaking  $\rightarrow$  **large field fluctuations** + relics.  
This extra source of fluctuations produces axion relics + structure.

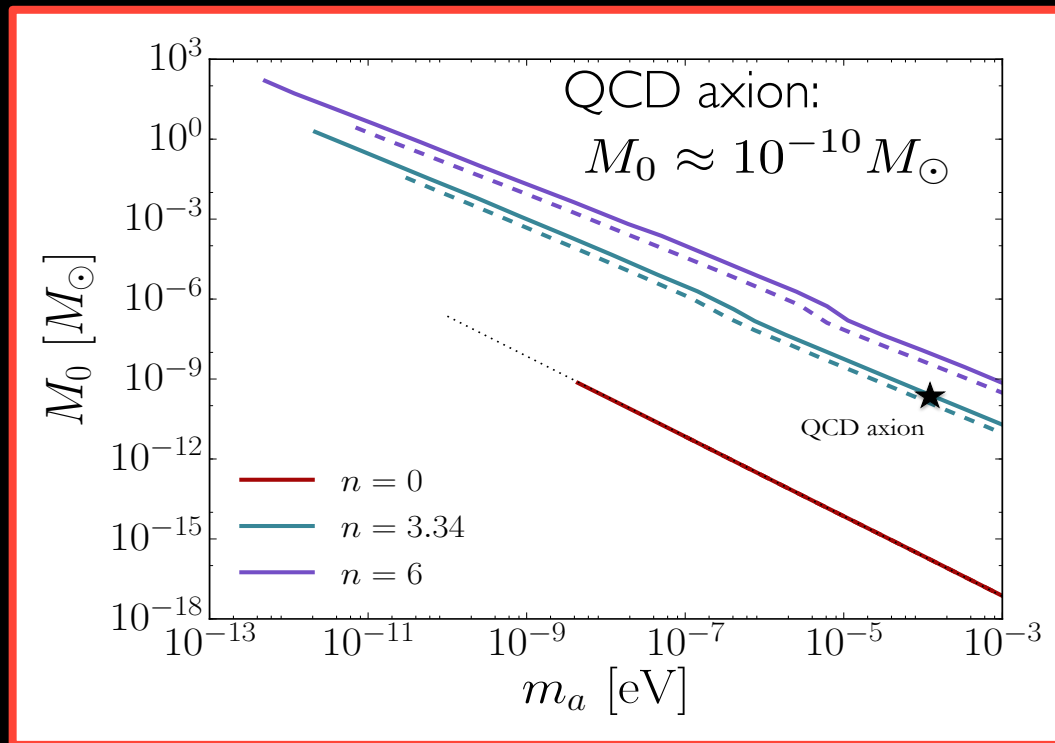


# Minicluster Mass Scale

Hogan & Rees (1988)

Axion randomises beyond horizon  $\rightarrow$  Large isocurvature fluctuations  
 $\rightarrow$  Mass inside horizon when  $m \sim H$  collapses early.

$$M_0 = (4/3)\pi(\pi/k)^3$$



Smaller than smallest WIMP structures ( $10^{-6}$ ). Axions are very cold.

The amount of DM in compact objects is strongly constrained:

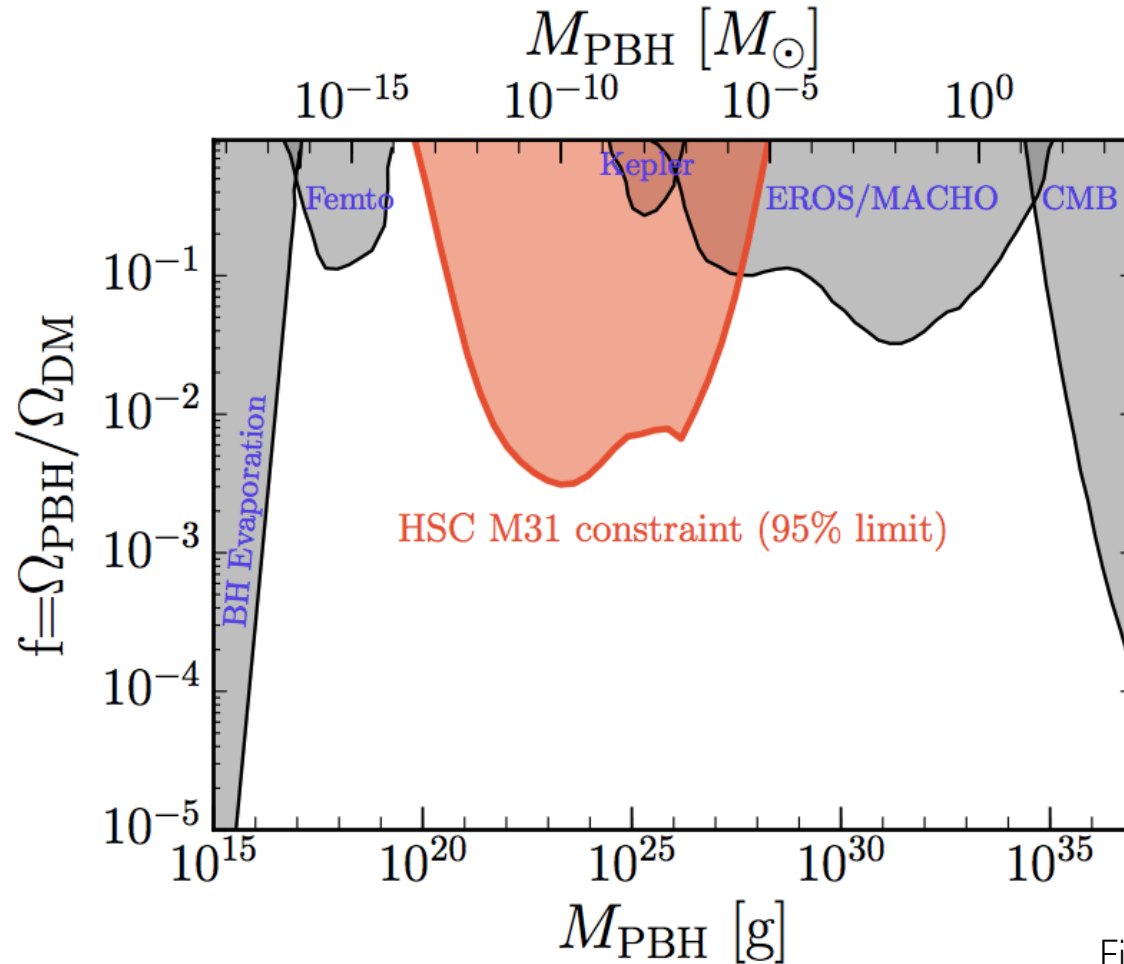
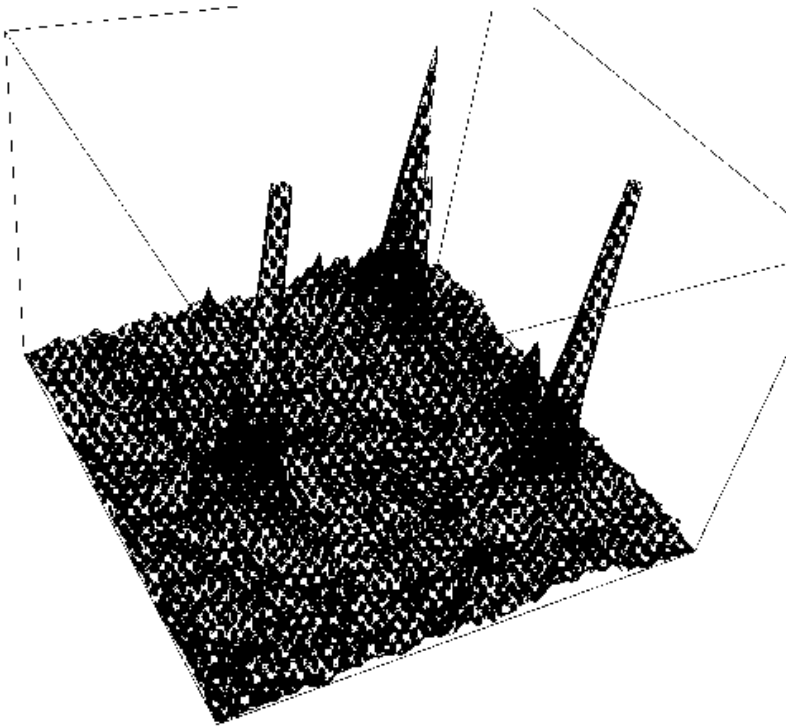


Fig: Niikura et al (2017)

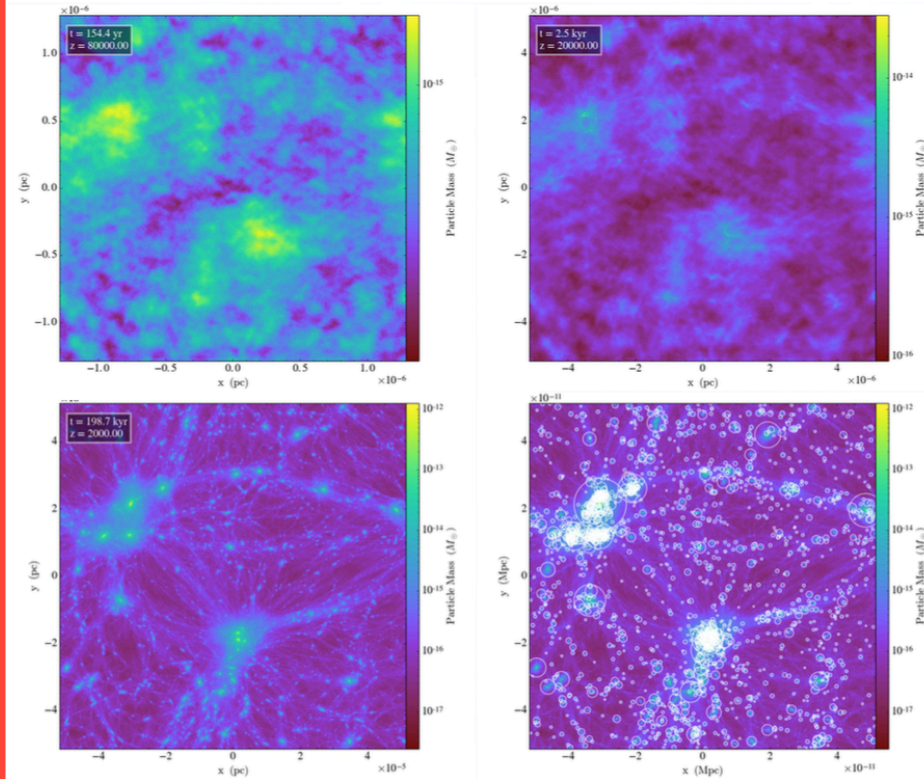
How do these constraints translate to miniclusters (or e.g. UCMHs)?  
 What is the relic density, mass, and size distribution of miniclusters?

# Numerical Simulations

See also Hardy (2017),  
Zurek et al (2007)



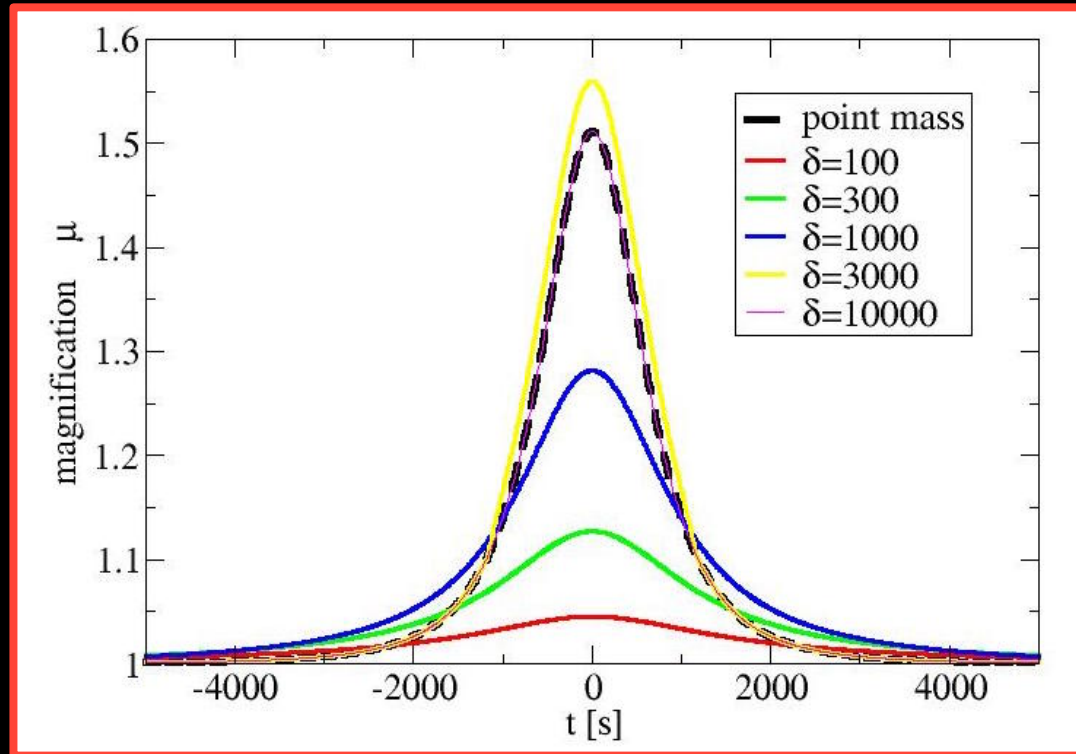
Kolb & Tkachev (1990's): field  
simulations w/ no gravity.  
→ Mass and radius scale.



Wiebe, Redondo, Niemeyer  
(2017): SSB i.c.'s w/ strings + N-  
body during rad. era.

# Microlensing: non-pointlike objects

Parameterise the density profile based on initial overdensity,  $\delta$ .



Effects described by a rescaling of the “microlensing tube”:

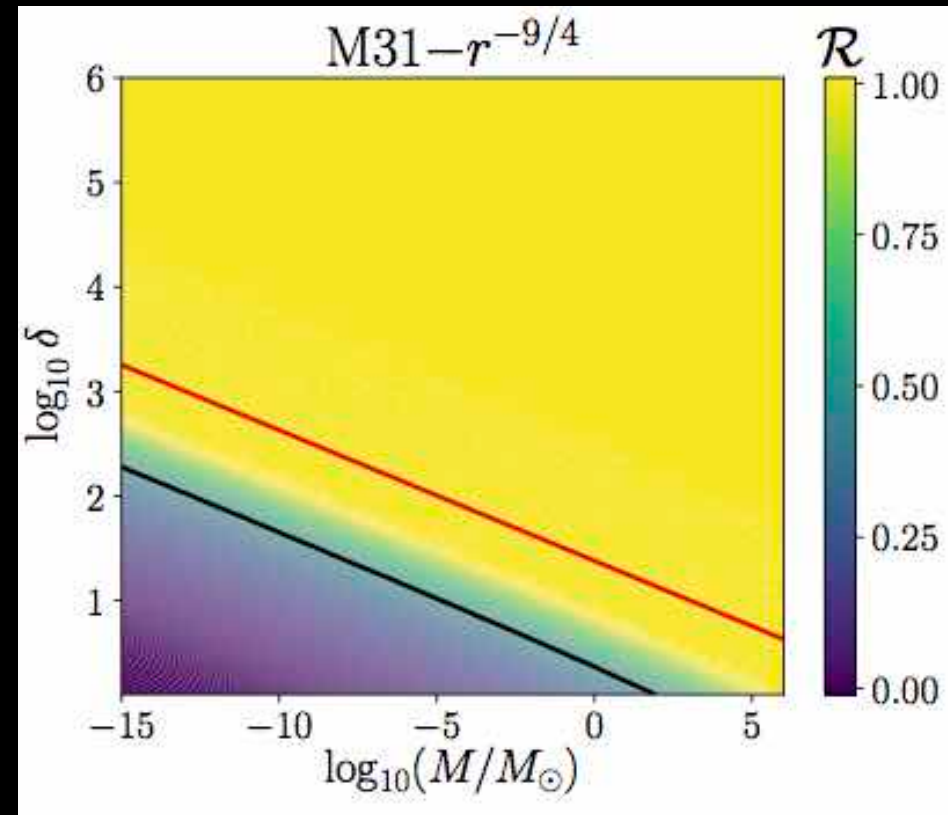
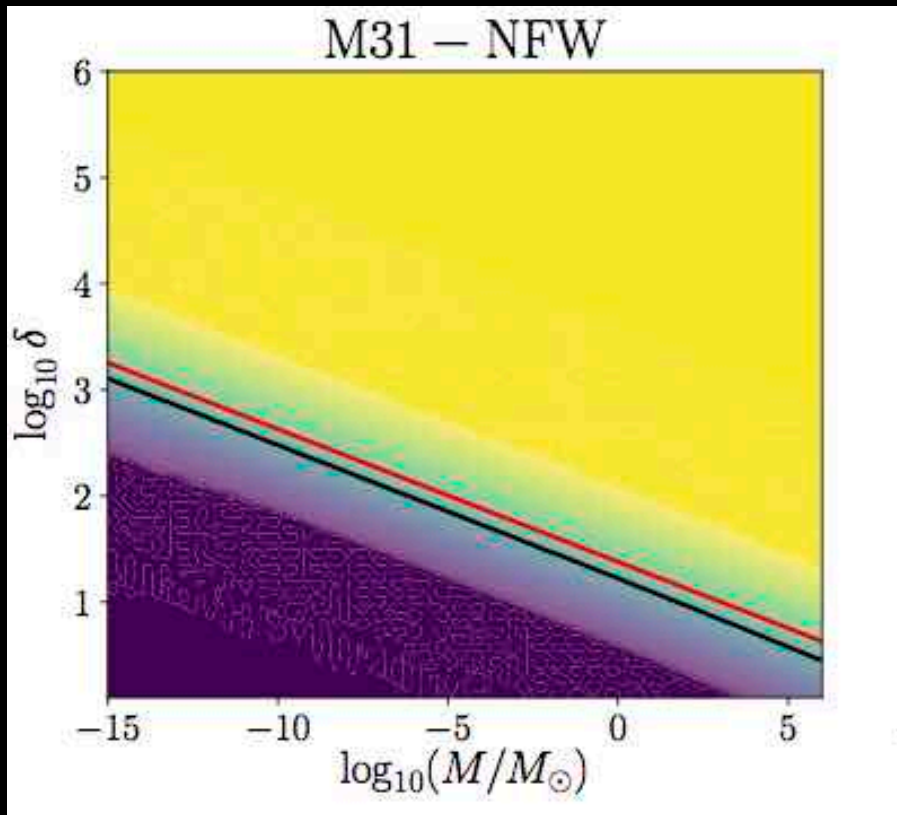
$$R_{\text{MC}}(\delta, M, x) = \mathcal{R}(\delta, M) R_{\text{E}}(M, x)$$

# lensing events depends on density profile and distribution of sizes.



# Density Profiles

Above some value of  $\delta$  miniclusters are **effectively point-like lenses**.  
The transition depends on the assumed density profile.



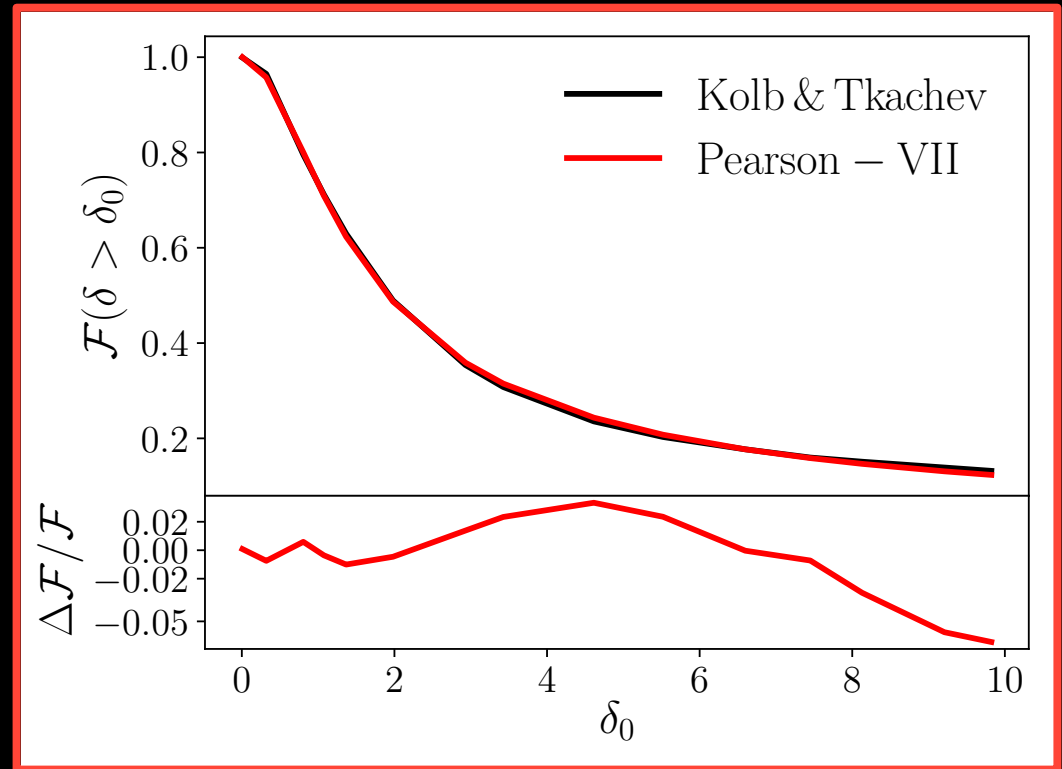
Self-similar infall  $\rightarrow$  power law density profile for isolated objects.  
Mergers + environment  $\rightarrow$  NFW density profile.

# Size Distribution

Size distribution: see also Ensander et al (2017)  
Extrapolation: Javier Redondo private comm.

Minicluster density field is **non-Gaussian** due to axion interactions.

Kolb & Tkachev sims → wide distribution for characteristic density:



$$\rho_{\text{MC}} = 140\delta^3(1 + \delta)\bar{\rho}_a(1 + z_{\text{eq}})^3$$

We use this to set the radius of minicluster density profiles

→ **Non-Gaussian** distribution of sizes. Key to results.

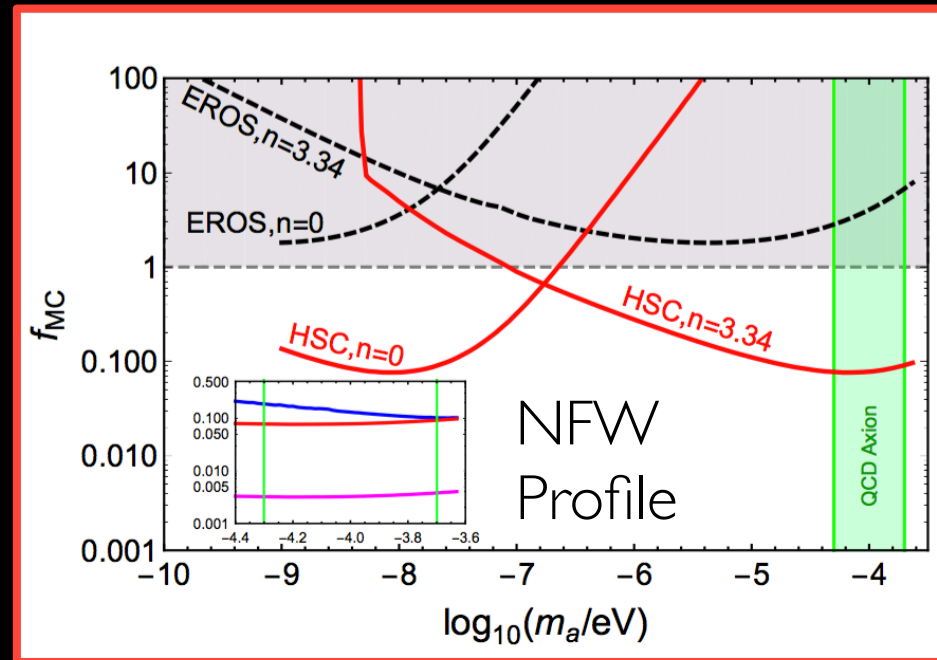
# Constraints: EROS and HSC

We constrain the allowed fraction of DM bound in MCs,  $f_{\text{MC}}$ .

Mass lower limit: telescope cadence. Upper limit: observing time.

Subaru HSC observed M31. Cadence  $\sim 2$  minutes  
 $\rightarrow$  access to very low masses.

Constraints from single night observing!

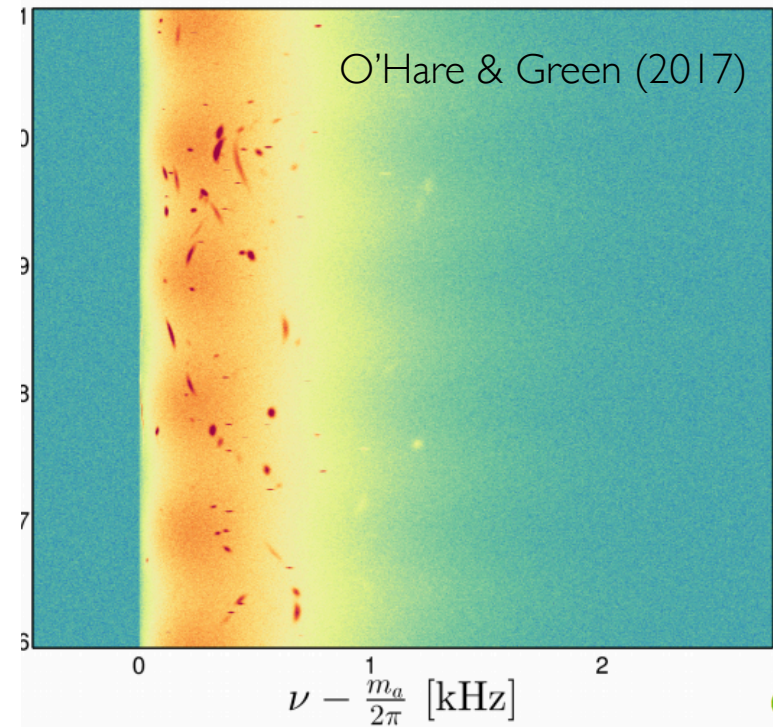
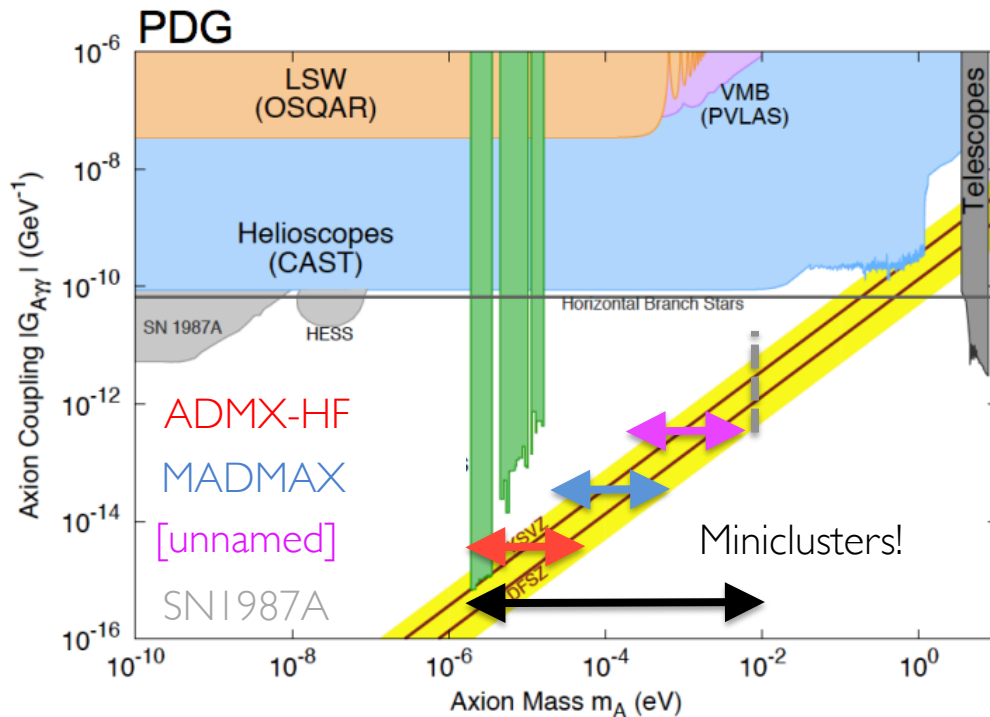


$$f_{\text{MC}} < 0.083(m_a/100\mu\text{eV})^{0.12}$$

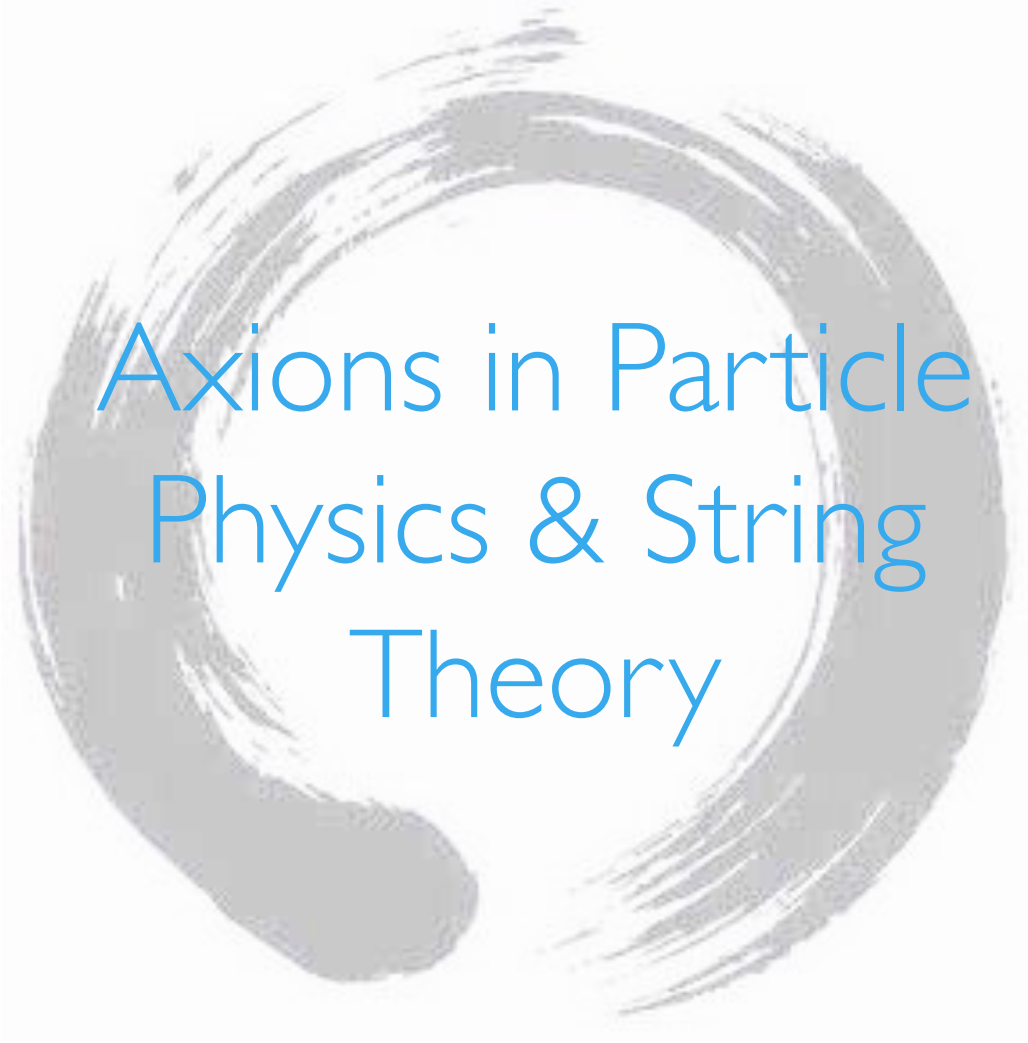
Lots of assumptions in this number: improved modeling needed.

# Miniclusters & Axion Astronomy

Many new **high frequency** (>GHz) axion experiments are being built.



If miniclusters fraction is large  $\rightarrow$  **drastic effect on direct detection.**  
 But, MC streams could be detected in spectra  $\rightarrow$  **measure halo.**  
 Axion experiments can measure the whole DM phase space dist.!



Axions in Particle  
Physics & String  
Theory

# Fine tuning in theoretical Physics

$$\mathcal{L} = (\text{number}) \times (\text{operator})$$

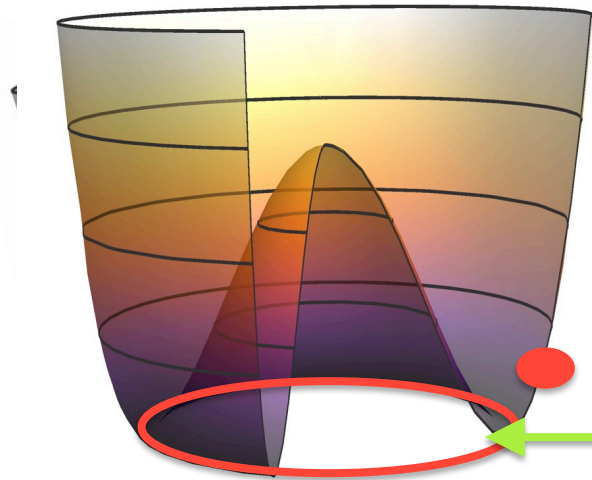
↓  
=the value you measure  
for the effect.

↓  
=physical effect. The  
neutron EDM.

$$(\text{number}) = \sum_{(\text{physics})} c_i \longrightarrow \begin{array}{l} \text{Contributions from} \\ \text{different sources.} \\ \text{Strong and weak} \\ \text{nuclear forces.} \end{array}$$

The nEDM is measured consistent with zero. This implies a delicate cancellation at  $\sim 10^{-10}$ . We don't like coincidences.

# Cleaning up the mess



Spontaneous symmetry breaking: a Goldstone Boson,  $\theta$ , has a continuous “shift symmetry”, i.e. the underlying rotation.

“instantons”

Trick: couple a Goldstone to the problematic operator.

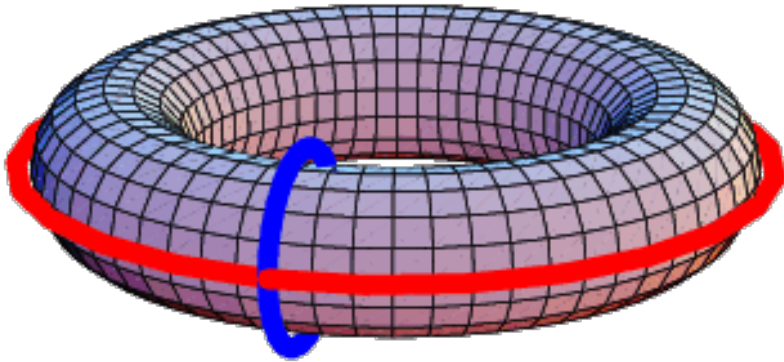
$$(\text{number}) \rightarrow (\text{number}) + \theta \rightarrow \theta$$

Now “tilt the wine bottle”:  $\theta$  will now dynamically move to a fixed value = 0 by symmetry  $\rightarrow$  no problematic nEDM.

BONUS: the oscillations of  $\theta$  carry energy density  $\rightarrow$  DM!  
BUT: axions are some  $10^{15}$  times lighter than “WIMPs”!

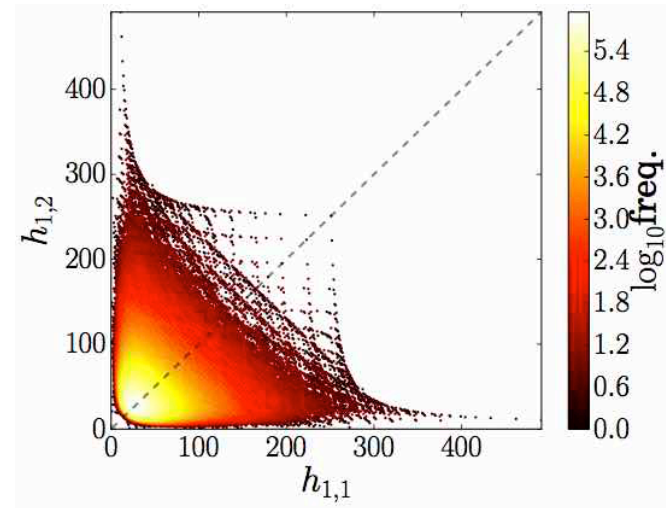
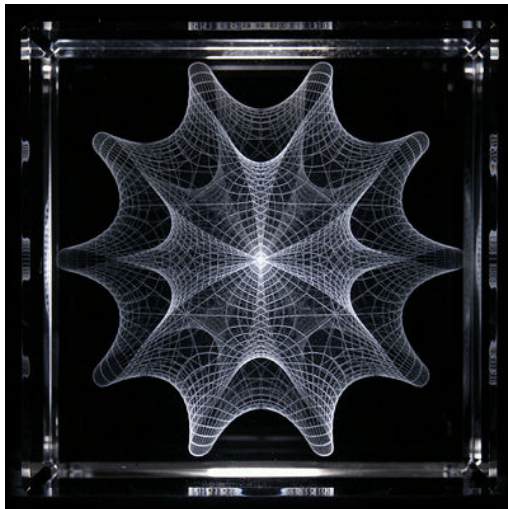
# Kaluza-Klein and String Theory

GR: fields describe the geometry of space. Cosmology: the scale factor.



A torus has two “moduli” that describe its shape.  
In general, # of fields given by topology.

String theory  $\rightarrow$  SUSY  $\rightarrow$  moduli are “paired” with AXIONS.  
10 dimensions can be a lot more complicated than a donut.





Extra dimensions and SUSY can both “hide” at high energies near the Planck scale. Many “flavours” of axions a [generic, low-energy, prediction.](#)

**AXION**  
PENGHAPUS MINYAK

**100% Effective**  
**Cosmology**



**Dark Matter**



**Dark Energy**



**Inflation**



**Baryogenesis**

[Learn More](#)

\*Based on 2011 Consumer Research