

# The NMSSM and Gauge Mediation

Robert Ziegler (CERN)

based on

**arXiv:1304.1453, with L. Calibbi, P. Paradisi**

**arXiv:1502.05836, with B. Allanach, M. Badziak, C. Hugonie**

**arXiv:1502.05836, with B. Allanach, M. Badziak, G. Cottin, N. Desai, C. Hugonie**

**work in progress, with M. Badziak, N. Desai, C. Hugonie**

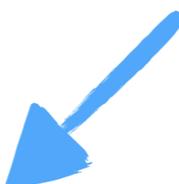
# Introduction

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The LHC legacy so far:

A 125 GeV Higgs, no new physics

Don't give up yet: best candidate to protect weak scale is still low-energy SUSY!



Still room for light sparticles:  
explore non-minimal models



Motivate new searches:  
explore non-standard signals



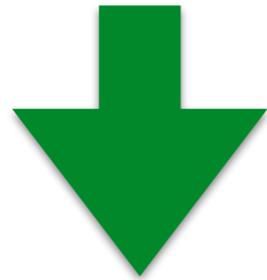
For heavy SUSY indirect tests important:  
explore non-minimal flavor structures

# Going top-down

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SUSY has too many parameters!

(~100 in the MSSM, mostly flavor)



Work with predictive  
scenario of SUSY breaking:

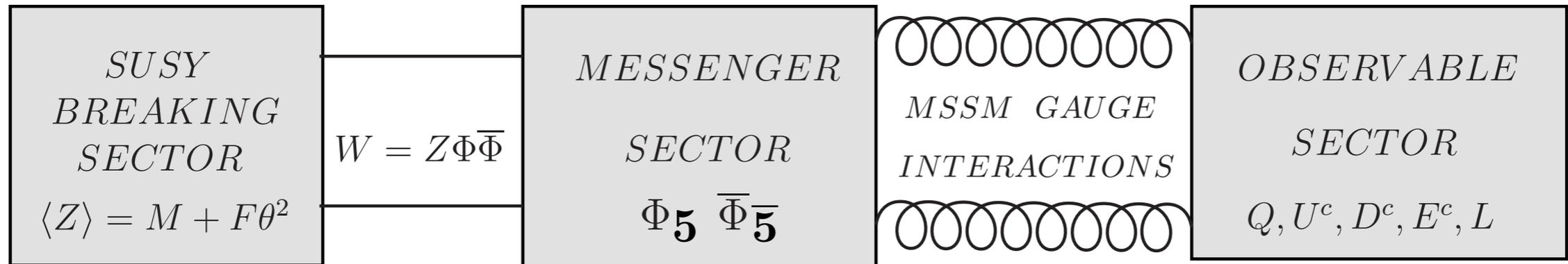
**Gauge Mediation**

(also takes care of SUSY flavor problem)



# Minimal Gauge Mediation

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- **Very predictive setup** (5 parameters) ✓
- **Solves SUSY Flavor Problem** (MFV) ✓
- **Heavy Higgs problematic** (small A-terms) ✗

# A Heavy Higgs in Gauge Mediation

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## How to get a heavy Higgs in GMSB ?

### PART I

Large boundary  $A_t$   
need matter-messenger couplings

Large radiative  $A_t$   
need heavy gluino & squarks  
{outside LHC reach}

### PART II

New tree-level  $\Delta m_h^2$   
need D-terms, F-terms, mixing

**NMSSM!**

# PART I

Large boundary  $A_t$   
need matter-messenger couplings

# Matter-Messenger Couplings

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Messengers have same quantum number as Higgs

$$W_{\text{yuk}} = y_{ij} Q_i U_j H_u \quad \longrightarrow \quad \Delta W_{\text{yuk}} = \lambda_{ij} Q_i U_j \Phi_{H_u}^5$$

Get new contributions to  $A$ -terms and soft masses

$$A_t \sim \frac{\Lambda}{16\pi^2} \lambda \lambda^\dagger y \qquad \Delta m_Q^2 \sim \frac{\Lambda^2}{256\pi^4} (\lambda \lambda^\dagger - g_3^2) \lambda \lambda^\dagger$$

**New sources of  
flavor violation**

**Need hierarchical flavor structure:**  $\lambda_{ij} \neq \mathcal{O}(1)_{ij}$

**...but very plausible since also**  $y_{ij} \neq \mathcal{O}(1)_{ij}$

# Natural Flavor Protection

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Relate new couplings to Yukawa hierarchies

$$\Delta W = \lambda_{ij} Q_i U_j \Phi_{H_u}^5$$

aligned to Yukawas  
through explicit mixing  
(Minimal Flavor Violation)

Evans, Ibe, Yanagida '11  $\lambda_{ij} \propto y_{ij}^u$

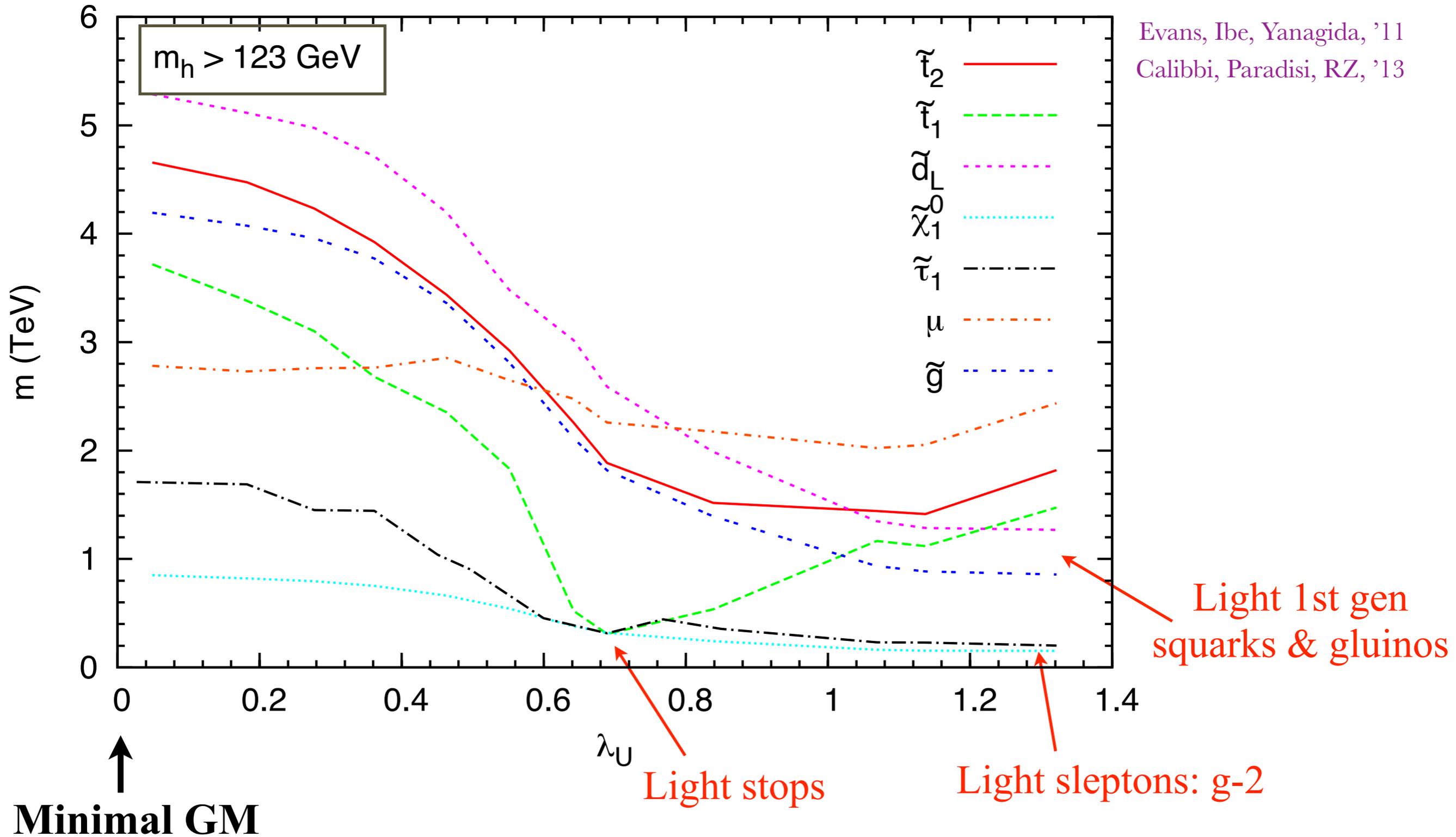
controlled by same  
underlying flavor model  
("Flavored Gauge Mediation")

Shadmi & Szabo '11  $\lambda_{ij} \sim y_{ij}^u$

**Same sparticle spectrum:** only  $\lambda_{33} \sim y_t$  is sizable

**Different sflavor structure:** close but beyond MFV!

# Low-energy Sparticle Spectrum



# Sflavor Structure in U(1) Flavor Model

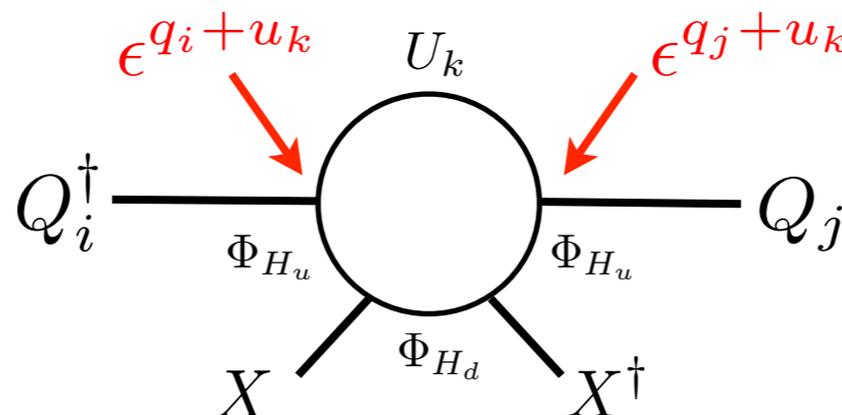
Calibbi, Paradisi, RZ, '13

If U(1) FN model controls couplings, can estimate sflavor structure in terms of masses and mixings

$$y_{ij}^u = \underbrace{x_{ij}}_{\mathcal{O}(1)} \epsilon^{q_i + u_j} \quad \longrightarrow \quad \lambda_{ij} = \underbrace{y_{ij}}_{\mathcal{O}(1)} \epsilon^{q_i + u_j}$$

Loop origin gives strong sflavor protection

$$(\tilde{m}_Q^2)_{ij} \sim \lambda_{ik} \lambda_{jk}^* \sim \epsilon^{q_i + q_j + 3u_3} \sim V_{i3} V_{j3}^* y_t^2$$



cf. Gravity  
Mediation+U(1):  
 $\sim V_{i3}/V_{j3}$

# Sflavor Structure between PC and MFV

	MFV	PC	$U(1)$	$FGM_{U,D} + U(1)$	$FGM_U + U(1)$
$(\delta_{LL}^u)_{ij}$	$V_{i3}V_{j3}^*y_b^2$	$(\epsilon_3^q)^2V_{i3}V_{j3}^*$	$\frac{V_{i3}}{V_{j3}}  i \leq j$	$V_{i3}V_{j3}^*y_t^2$	$V_{i3}V_{j3}^*y_t^2$
$(\delta_{LL}^d)_{ij}$	$V_{3i}^*V_{3j}y_t^2$	$(\epsilon_3^q)^2V_{i3}V_{j3}^*$	$\frac{V_{i3}}{V_{j3}}  i \leq j$	$V_{3i}^*V_{3j}y_t^2$	$V_{3i}^*V_{3j}y_t^2$
$(\delta_{RR}^u)_{ij}$	$y_i^U y_j^U V_{i3}V_{j3}^*y_b^2$	$\frac{y_i^U y_j^U}{V_{i3}V_{j3}^*} \frac{(\epsilon_3^u)^2}{y_t^2}$	$\frac{y_i^U V_{j3}}{y_j^U V_{i3}}  i \leq j$	$\frac{y_i^U y_j^U}{V_{i3}V_{j3}^*}$	$\frac{y_i^U y_j^U}{V_{i3}V_{j3}^*}$
$(\delta_{RR}^d)_{ij}$	$y_i^D y_j^D V_{3i}^*V_{3j}y_t^2$	$\frac{y_i^D y_j^D}{V_{i3}V_{j3}^*} \frac{(\epsilon_3^u)^2}{y_t^2}$	$\frac{y_i^D V_{j3}}{y_j^D V_{i3}}  i \leq j$	$\frac{y_i^D y_j^D}{V_{i3}V_{j3}^*}$	$y_i^D y_j^D V_{3i}^*V_{3j}y_t^2$
$(\delta_{LR}^u)_{ij}$	$y_j^U V_{i3}V_{j3}^*y_b^2$	$y_j^U \frac{V_{i3}}{V_{j3}^*}$	$y_j^U \frac{V_{i3}}{V_{j3}^*}$	$y_j^U V_{i3}V_{j3}^*y_t^2 + y_i^U \frac{y_i^U y_j^U}{V_{i3}V_{j3}^*}$ $y_j^U \frac{V_{i3}}{V_{j3}^*} y_t^6$	$y_j^U V_{i3}V_{j3}^*y_t^2 + y_i^U \frac{y_i^U y_j^U}{V_{i3}V_{j3}^*}$ $y_j^U \frac{V_{i3}}{V_{j3}^*} y_t^6$
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**Effects mainly in LR up-sector:  $\Delta C = 1$**

## PART II

New tree-level  $\Delta m_h^2$   
need D-terms, F-terms, mixing

**NMSSM!**

# Raise Higgs mass with mixing

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$$\begin{pmatrix} m_h^2 & m_{hs}^2 \\ m_{hs}^2 & m_s^2 \end{pmatrix} \xrightarrow{m_h^2 > m_s^2} m_{h_2}^2 \approx m_h^2 + \frac{m_{hs}^4}{m_h^2}$$

Can realize in NMSSM  $\Delta W = \lambda S H_u H_d + \frac{\kappa}{3} S^3$

Mixing angles constrained by LEP and LHC,  
maximal contribution to Higgs mass for:

$$m_{h_1} \approx 94 \text{ GeV} \quad \cos \theta \approx 0.88$$

# Embedding into GMSB

**Can one realize mixing scenario in GMSB?**

Nice to marry NMSSM to GMSB: solves  $\mu$ - $B_\mu$  problem

$$\text{MSSM} + \text{GM:} \quad \mu \sim a \frac{\Lambda}{16\pi^2} \quad B_\mu \sim a \frac{\Lambda^2}{16\pi^2} \sim 16\pi^2 \mu^2$$

$$\text{NMSSM} + \text{GM:} \quad \mu \sim \langle S \rangle \sim m_{\text{soft}} \quad B_\mu \sim \langle F_S \rangle \sim m_{\text{soft}}^2 \sim \mu^2$$

...not in minimal Gauge Mediation  
because singlet soft terms too small

Dine, Nelson '93

de Gouvea, Friedland, Murayama '13

# The DGS Model

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Simplest Model: MGM + two pairs of messengers and direct couplings to singlet

Delgado, Giudice, Slavich '07

$$W_{\text{DGS}} = S \left( \xi_D \bar{\Phi}_1^D \Phi_2^D + \xi_T \bar{\Phi}_1^T \Phi_2^T \right)$$

$$\xi_D(M_{\text{GUT}}) = \xi_T(M_{\text{GUT}}) \equiv \xi$$

1-loop  $A$ -terms

$$A_\lambda \sim A_\kappa \sim \tilde{m} \xi^2$$

2-loop singlet mass

$$\tilde{m}_S^2 \sim \tilde{m}^2 (\xi^4 - g_3^2 \xi^2)$$

$$\tilde{m} \equiv \Lambda / (16\pi^2) \sim m_{\tilde{g}}$$

Has only **4 parameters**  $\lambda, \tilde{m}, \xi, M$

[correct EWSB fixes  $\kappa$ ]

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$$\xi_D(M_{\text{GUT}}) = \xi_T(M_{\text{GUT}}) \equiv \xi$$

1-loop  $A$ -terms

$$A_\lambda \sim \lambda$$

2-loop singlet mass

$$(\xi^4 - g_3^2 \xi^2)$$

Singlet-Higgs coupling

$$\tilde{m} \equiv \Lambda / (16\pi^2) \sim m_{\tilde{g}}$$

Has only **4 parameters**  $\lambda, \tilde{m}, \xi, M$

[correct EWSB fixes  $\kappa$ ]

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1-loop A-terms

$$A_\lambda \sim A_\kappa \sim$$

2-loop singlet mass

Singlet-Messenger coupling

$$\left( \frac{2\xi^2}{3S^2} \right)$$

$$/(16\pi^2) \sim m_{\tilde{g}}$$

Has only **4 parameters**  $\lambda, \tilde{m}, \xi, M$

[correct EWSB fixes  $\kappa$ ]

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1-loop A-terms

$$A_\lambda \sim A_\kappa \sim \dots$$

2-loop singlet mass

$$\tilde{m}^2 (\xi^4 - g_3^2 \xi^2)$$

Soft SUSY scale

$$\tilde{m} \equiv \Lambda / (16\pi^2) \sim m_{\tilde{g}}$$

Has only **4 parameters**  $\lambda, \tilde{m}, \xi, M$

[correct EWSB fixes  $\kappa$ ]

# The DGS Model

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$$\xi_D(M_{\text{GUT}}) = \xi_T(M_{\text{GUT}}) \equiv \xi$$

1-loop  $A$ -terms

$$A_\lambda \sim A_\kappa \sim \dots$$

2-loop singlet mass

$$\tilde{m}^2 (\xi^4 - g_3^2 \xi^2)$$

Messenger scale

$$\tilde{m} \equiv \Lambda / (16\pi^2) \sim m_{\tilde{g}}$$

Has only **4 parameters**  $\lambda, \tilde{m}, \xi, M$

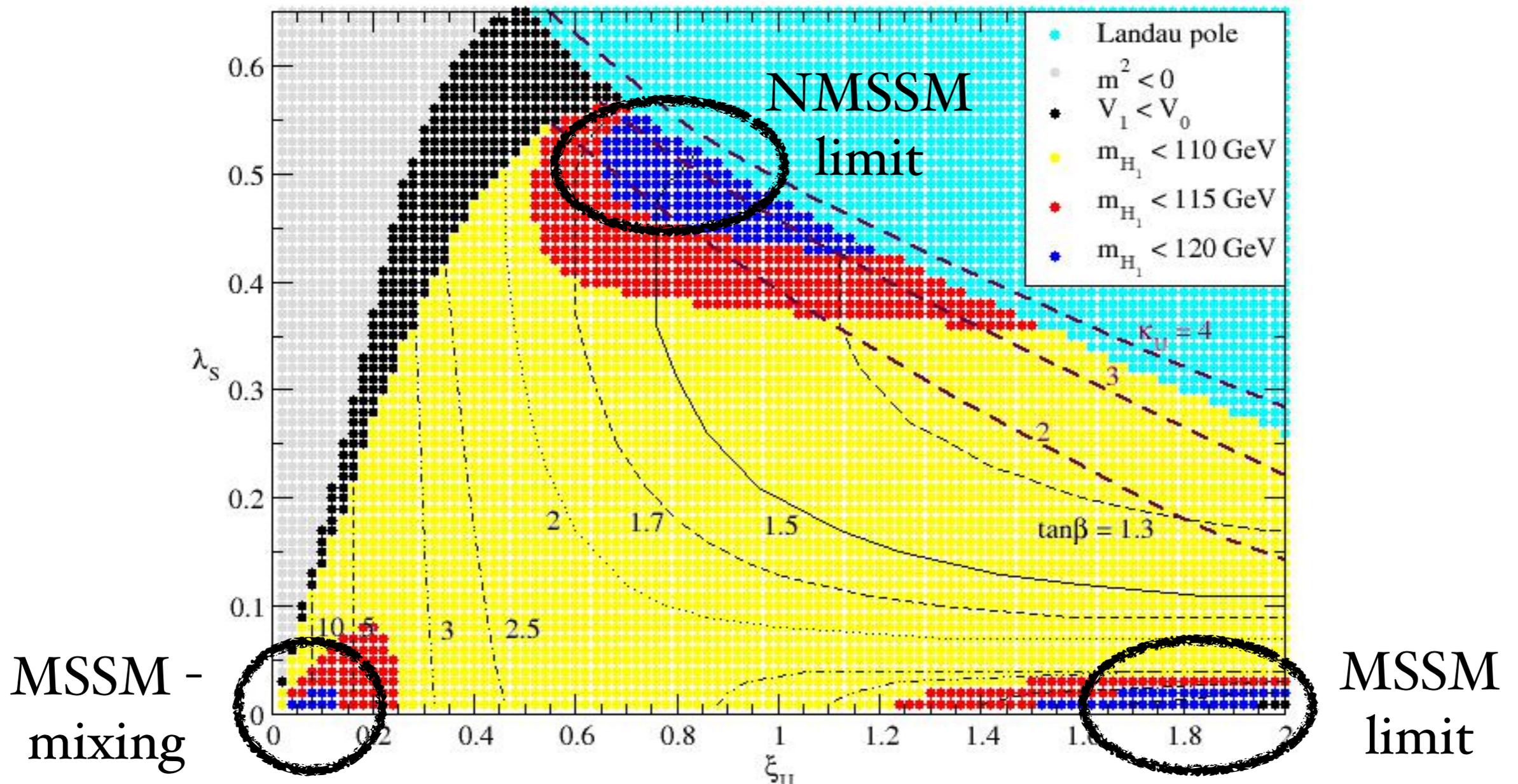
[correct EWSB fixes  $\kappa$ ]

# DGS Parameter Space

**Only 3 regions with sizable Higgs mass**

$$m_h^2 = M_Z^2 \cos^2 2\beta + \lambda^2 v^2 \sin^2 2\beta + m_{h,\text{mix}}^2 + m_{h,\text{loop}}^2$$

Sum bounded by  $M_Z^2$  (perturbativity)

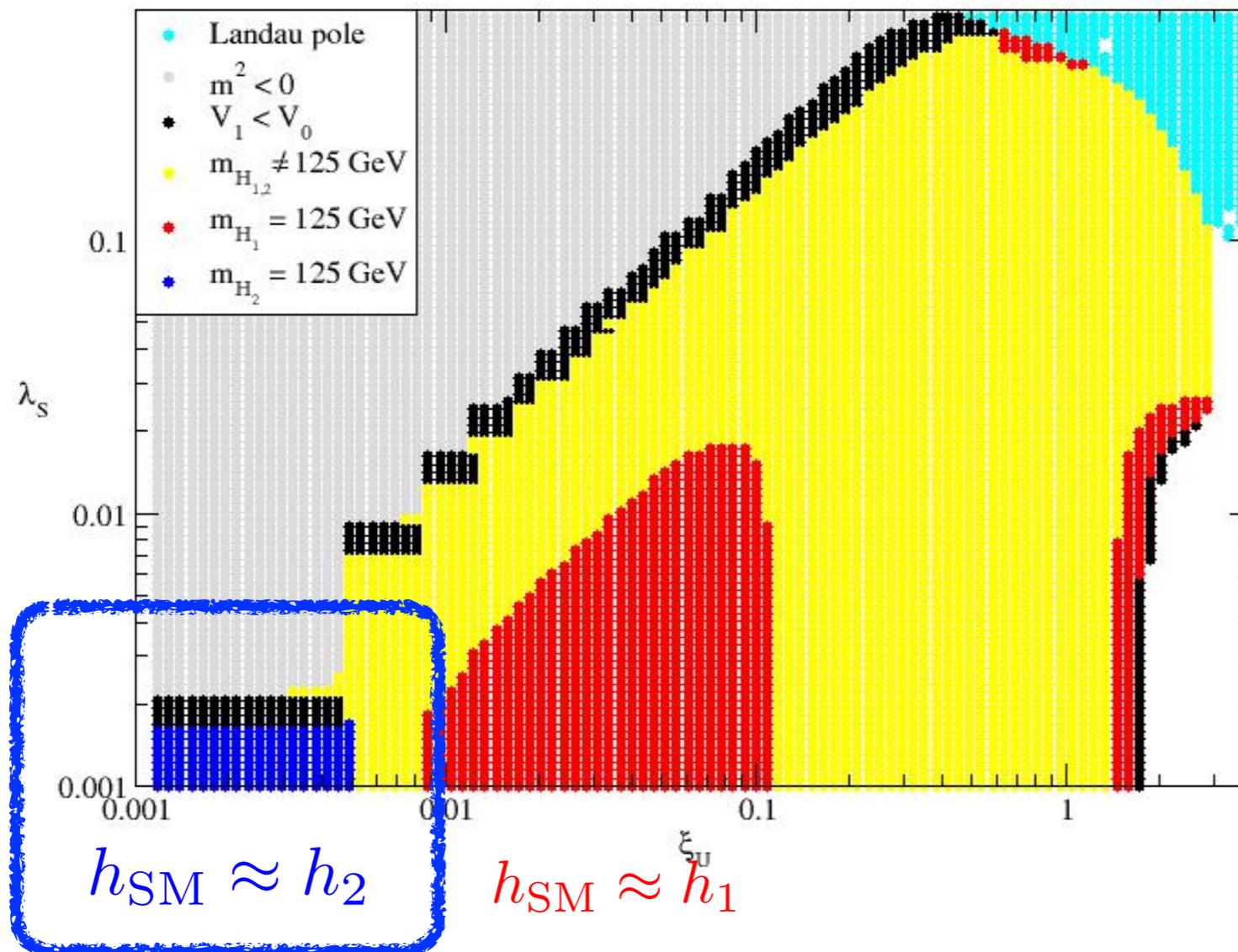


# The Push-up Region

“...even in these regions the lightest Higgs mass is not larger than the maximal value attainable in the usual GMSB.”

Delgado, Giudice, Slavich '07

**But MSSM + mixing region exists! (though hard to find)**



# The Maximal Push-up Region

Allanach, Badziak, Hugonie, RZ '15

Maximize tree-level Higgs contribution from mixing

$$m_{h_1} \approx 94 \text{ GeV} \quad \cos \theta \approx 0.88$$
$$m_{h_2} \approx 125 \text{ GeV}$$

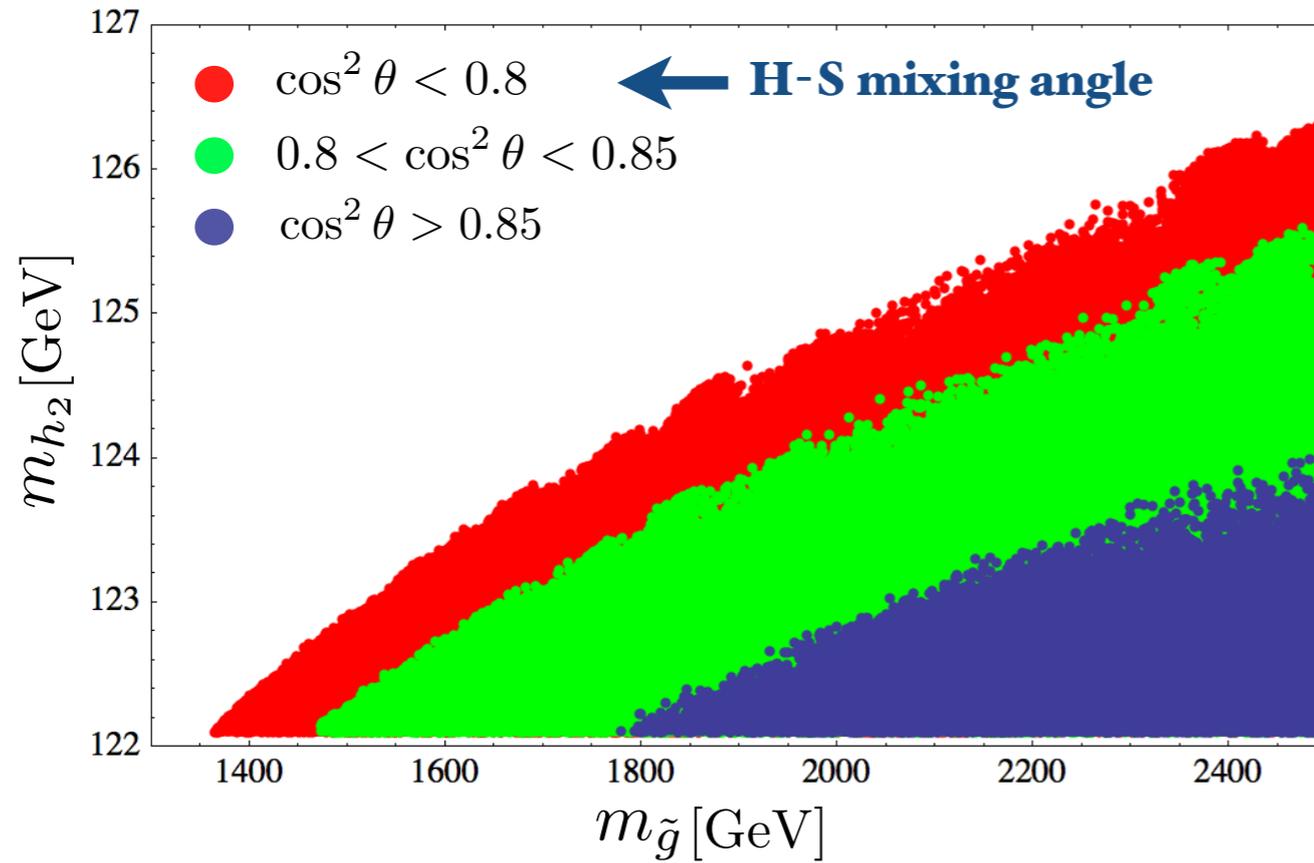
Fixes 3/4 parameters, only messenger scale free

$$\lambda, \xi \quad \tilde{m}$$
$$\sim 10^{-2} \quad \sim 1 \text{ TeV}$$

$M$   
determines gravitino phenomenology

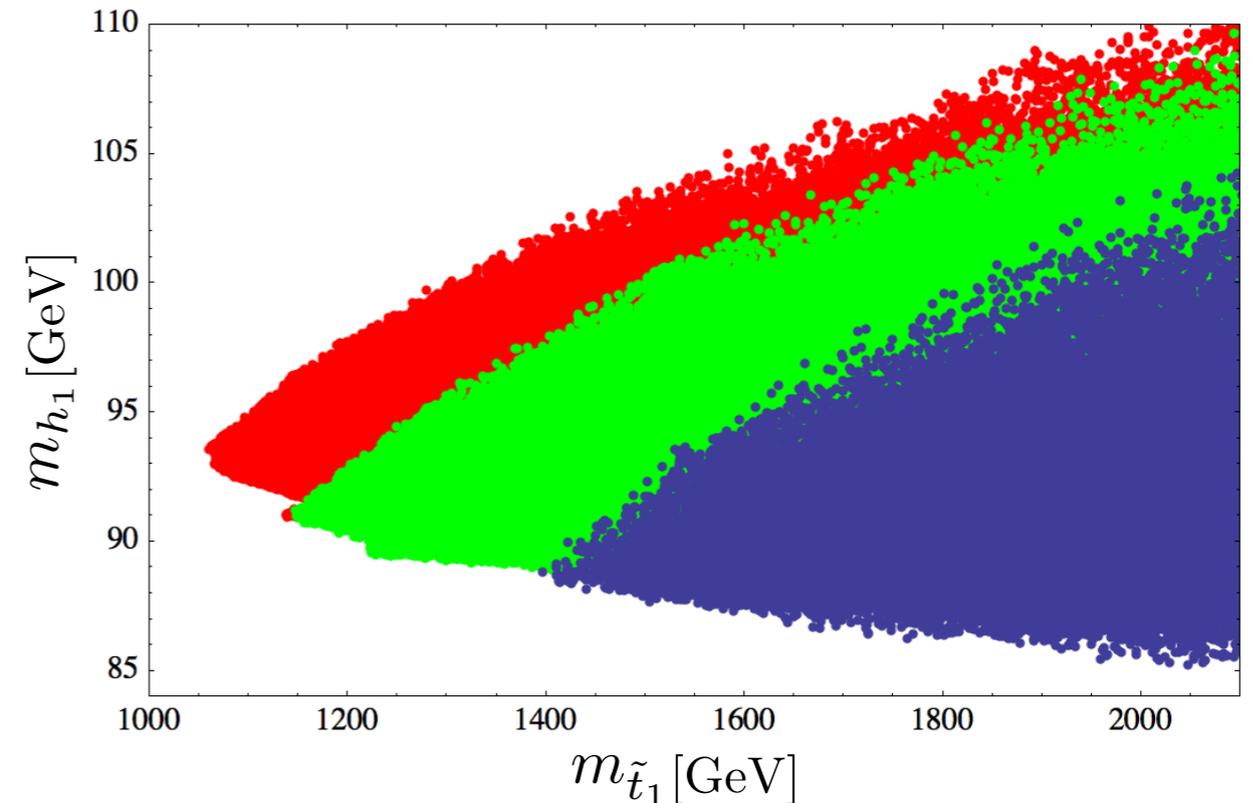
Higgs mass still drives lower bounds on sparticles, but can be close to direct exclusion bounds

# Light Gluinos and Stops



← Allows for 1.4 TeV Gluinos

...and 1.1 TeV stops →



# Phenomenology

	P1	P2	P3	P4	P5
$\tilde{m}$	$7.5 \cdot 10^2$	$8.7 \cdot 10^2$	$9.3 \cdot 10^2$	$5.9 \cdot 10^2$	$9.3 \cdot 10^2$
$M$	$1.4 \cdot 10^6$	$2.8 \cdot 10^6$	$3.3 \cdot 10^7$	$8.3 \cdot 10^{14}$	$3.4 \cdot 10^{14}$
$\lambda$	$1.0 \cdot 10^{-2}$	$9.3 \cdot 10^{-3}$	$6.7 \cdot 10^{-3}$	$9.2 \cdot 10^{-3}$	$6.9 \cdot 10^{-3}$
$\xi$	$1.2 \cdot 10^{-2}$	$1.1 \cdot 10^{-2}$	$1.2 \cdot 10^{-2}$	$2.2 \cdot 10^{-2}$	$2.0 \cdot 10^{-2}$
$\tan \beta$	25	28	24	26	21
$m_{h_1}$	92	93	98	94	94
$m_{h_2}$	122.1	123.4	122.9	122.1	125.0
$m_{a_1}$	26	26	28	40	32
$m_{\tilde{N}_1}$	101	102	106	104	104
$m_{\tilde{N}_2}$	322	377	400	251	379
$m_{\tilde{e}_1}$	303	358	406	449	676
$m_{\tilde{\tau}_1}$	284	333	376	432	637
$m_{\tilde{g}}$	1.73	1.98	2.09	1.37	2.06
$m_{\tilde{u}_R}$	1.79	2.06	2.15	1.36	2.07
$m_{\tilde{t}_1}$	1.64	1.87	1.90	1.06	1.63
$[m]$ $c\tau_{\tilde{N}_1}$	$6.4 \cdot 10^{-2}$	0.34	48	$1.9 \cdot 10^{16}$	$6.0 \cdot 10^{15}$
$\sigma_{\tilde{q}\tilde{q}}^{13\text{TeV}}$	9.35	2.99	1.98	59.7	2.63
$\sigma_{\tilde{q}\tilde{g}}^{13\text{TeV}}$	11.9	3.30	2.01	91.1	2.48
$\sigma_{\text{strong}}^{13\text{TeV}}$	25.2	7.28	4.58	190	5.95
$\sigma_{\text{strong}}^{8\text{TeV}}$	0.51	0.07	0.03	10.1	0.05
$\sigma_{\text{EW}}^{13\text{TeV}}$	27	12	7.5	6.7	5.6
$\sigma_{\text{EW}}^{8\text{TeV}}$	5.5	2.1	1.2	1.3	0.7

- Large tan beta

# Phenomenology

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- Large tan beta
- Light pseudoscalar ~25 GeV

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- Large tan beta
- Light pseudoscalar ~25 GeV
- **Singlino NLSP ~ 100 GeV**  
**displaced/stable**

# Phenomenology

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$\sigma_{\text{strong}}^{8\text{TeV}}$	0.51	0.07	0.03	10.1	0.05
$\sigma_{\text{EW}}^{13\text{TeV}}$	27	12	7.5	6.7	5.6
$\sigma_{\text{EW}}^{8\text{TeV}}$	5.5	2.1	1.2	1.3	0.7

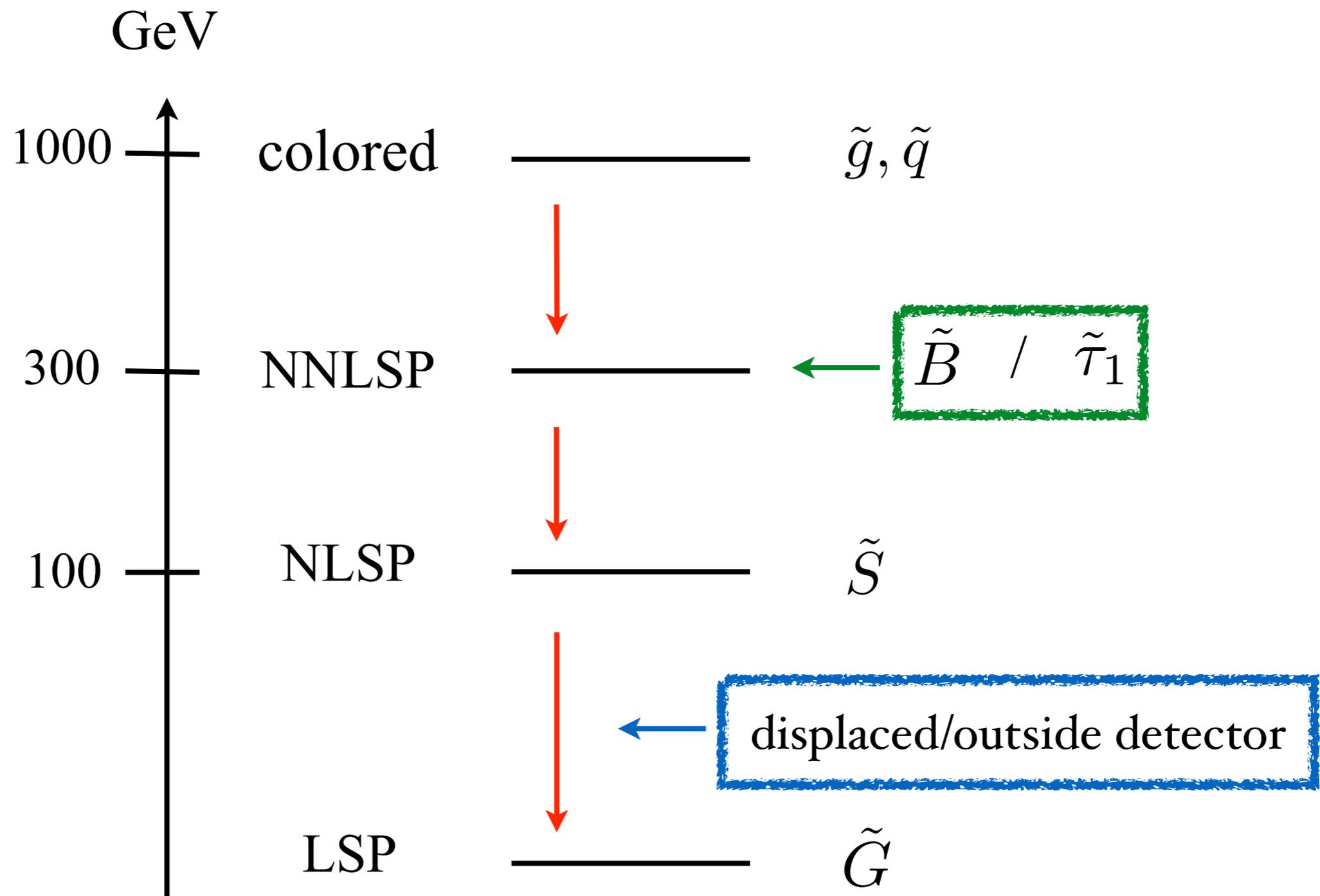
- Large tan beta
- Light pseudoscalar  $\sim 25$  GeV
- **Singlino NLSP  $\sim 100$  GeV**  
**displaced/stable**
- **stau/bino NNLSP  $\sim 300$  GeV**

# Phenomenology

	P1	P2	P3	P4	P5
$\tilde{m}$	$7.5 \cdot 10^2$	$8.7 \cdot 10^2$	$9.3 \cdot 10^2$	$5.9 \cdot 10^2$	$9.3 \cdot 10^2$
$M$	$1.4 \cdot 10^6$	$2.8 \cdot 10^6$	$3.3 \cdot 10^7$	$8.3 \cdot 10^{14}$	$3.4 \cdot 10^{14}$
$\lambda$	$1.0 \cdot 10^{-2}$	$9.3 \cdot 10^{-3}$	$6.7 \cdot 10^{-3}$	$9.2 \cdot 10^{-3}$	$6.9 \cdot 10^{-3}$
$\xi$	$1.2 \cdot 10^{-2}$	$1.1 \cdot 10^{-2}$	$1.3 \cdot 10^{-2}$	$3.2 \cdot 10^{-2}$	$2.0 \cdot 10^{-2}$
$\tan \beta$	25	28	24	26	21
$m_{h_1}$	92	93	98	94	94
$m_{h_2}$	122.1	123.4	122.9	122.1	125.0
$m_{a_1}$	26	26	28	40	32
$m_{\tilde{N}_1}$	101	102	106	104	104
$m_{\tilde{N}_2}$	322	377	400	251	379
$m_{\tilde{e}_1}$	303	358	406	449	676
$m_{\tilde{\tau}_1}$	284	333	376	432	637
$m_{\tilde{g}}$	1.73	1.98	2.09	1.37	2.06
$m_{\tilde{u}_R}$	1.79	2.06	2.15	1.36	2.07
$m_{\tilde{t}_1}$	1.64	1.87	1.90	1.06	1.63
$[m]$ $c\tau_{\tilde{N}_1}$	$6.4 \cdot 10^{-2}$	0.34	48	$1.9 \cdot 10^{16}$	$6.0 \cdot 10^{15}$
$\sigma_{\tilde{q}\tilde{q}}^{13\text{TeV}}$	9.35	2.99	1.98	59.7	2.63
$\sigma_{\tilde{q}\tilde{g}}^{13\text{TeV}}$	11.9	3.30	2.01	91.1	2.48
$\sigma_{\text{strong}}^{13\text{TeV}}$	25.2	7.28	4.58	190	5.95
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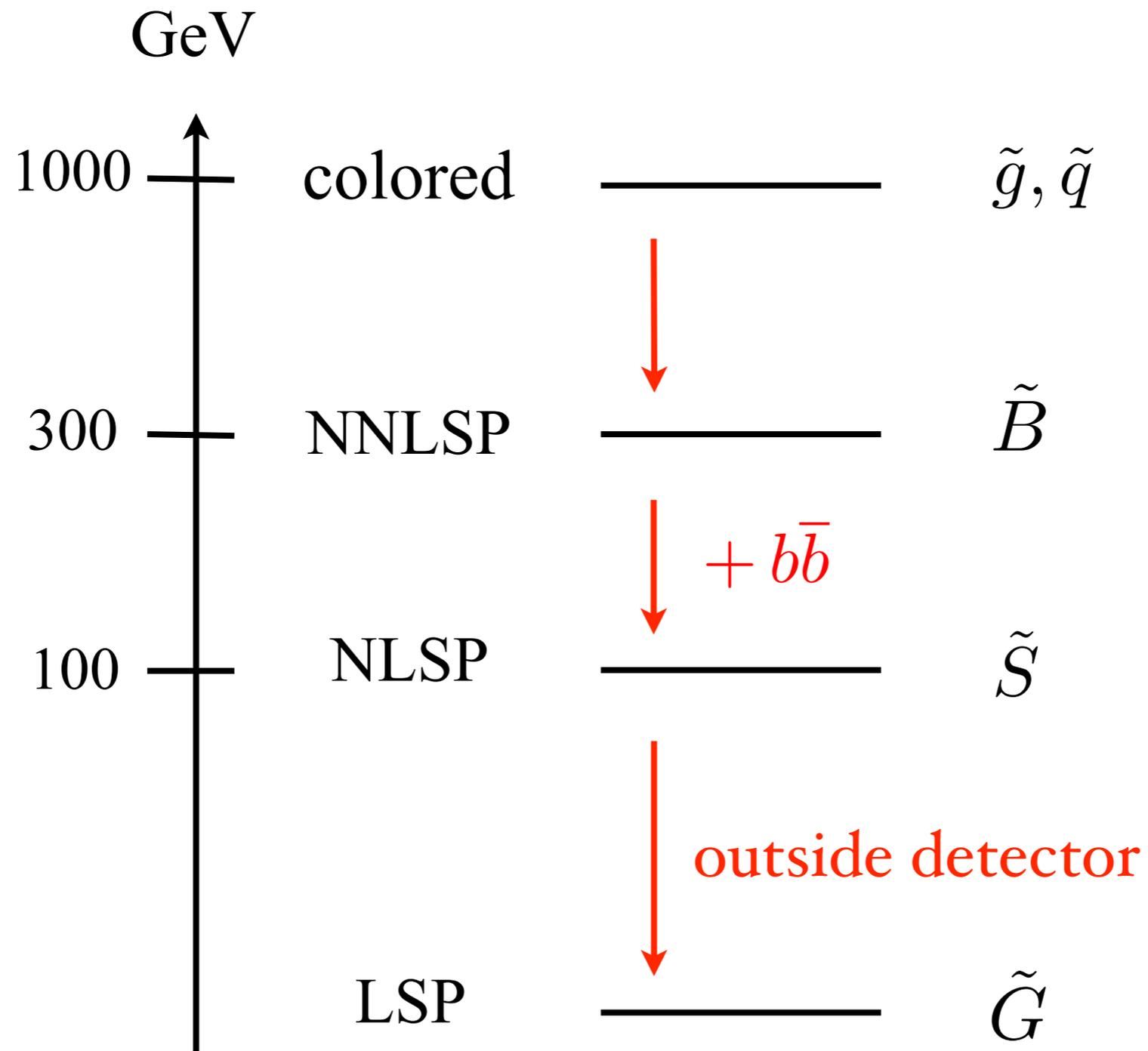
- Large tan beta
- Light pseudoscalar  $\sim 25$  GeV
- **Singlino NLSP  $\sim 100$  GeV**  
**displaced/stable**
- **stau/bino>NNLSP  $\sim 300$  GeV**
- Singlino essentially decoupled  
SUSY decays through NNLSP

# Signals: **NNLSP** & Singlino decay



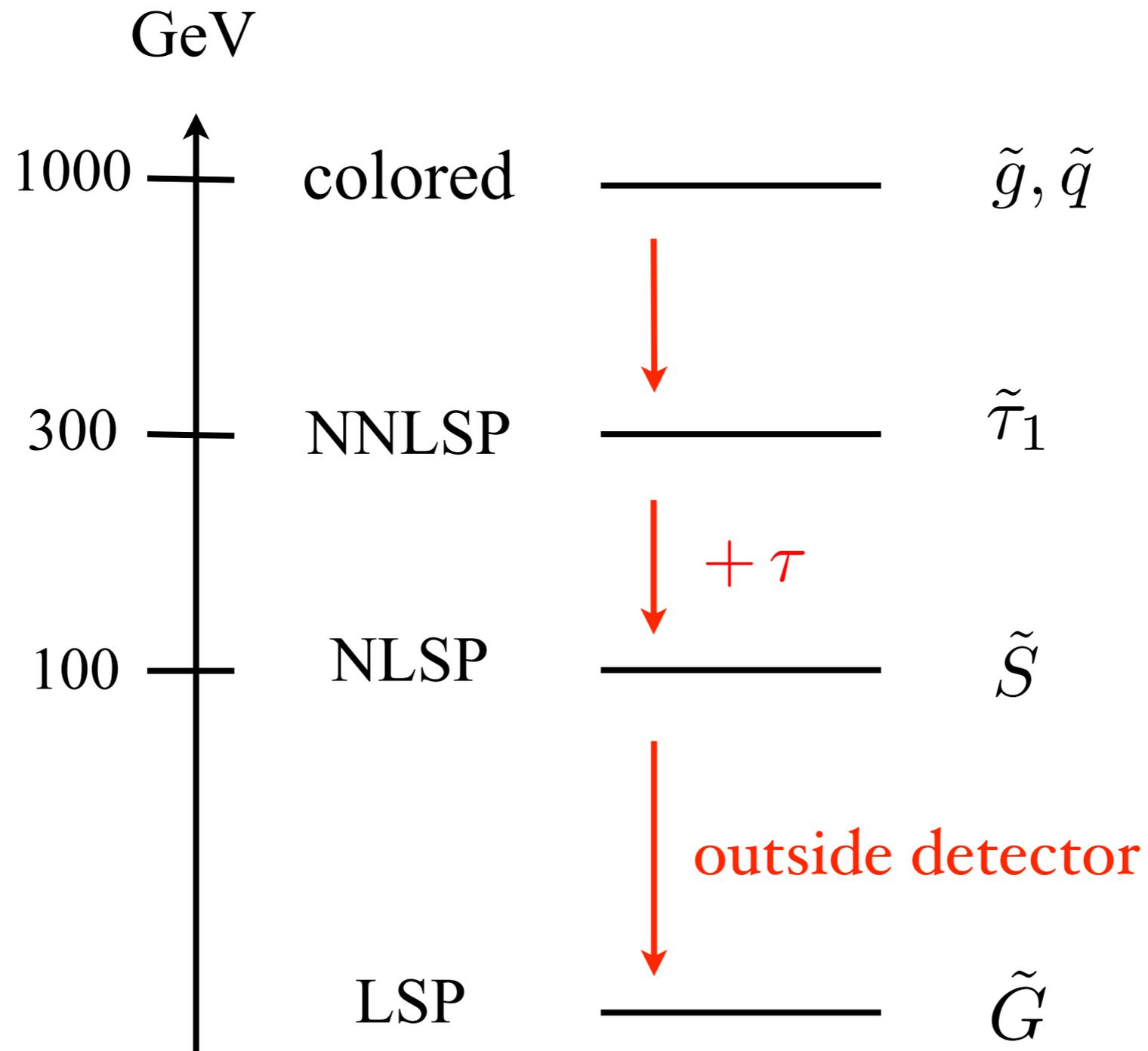
# Large-M Region: $M > 10^9 \text{ GeV}$

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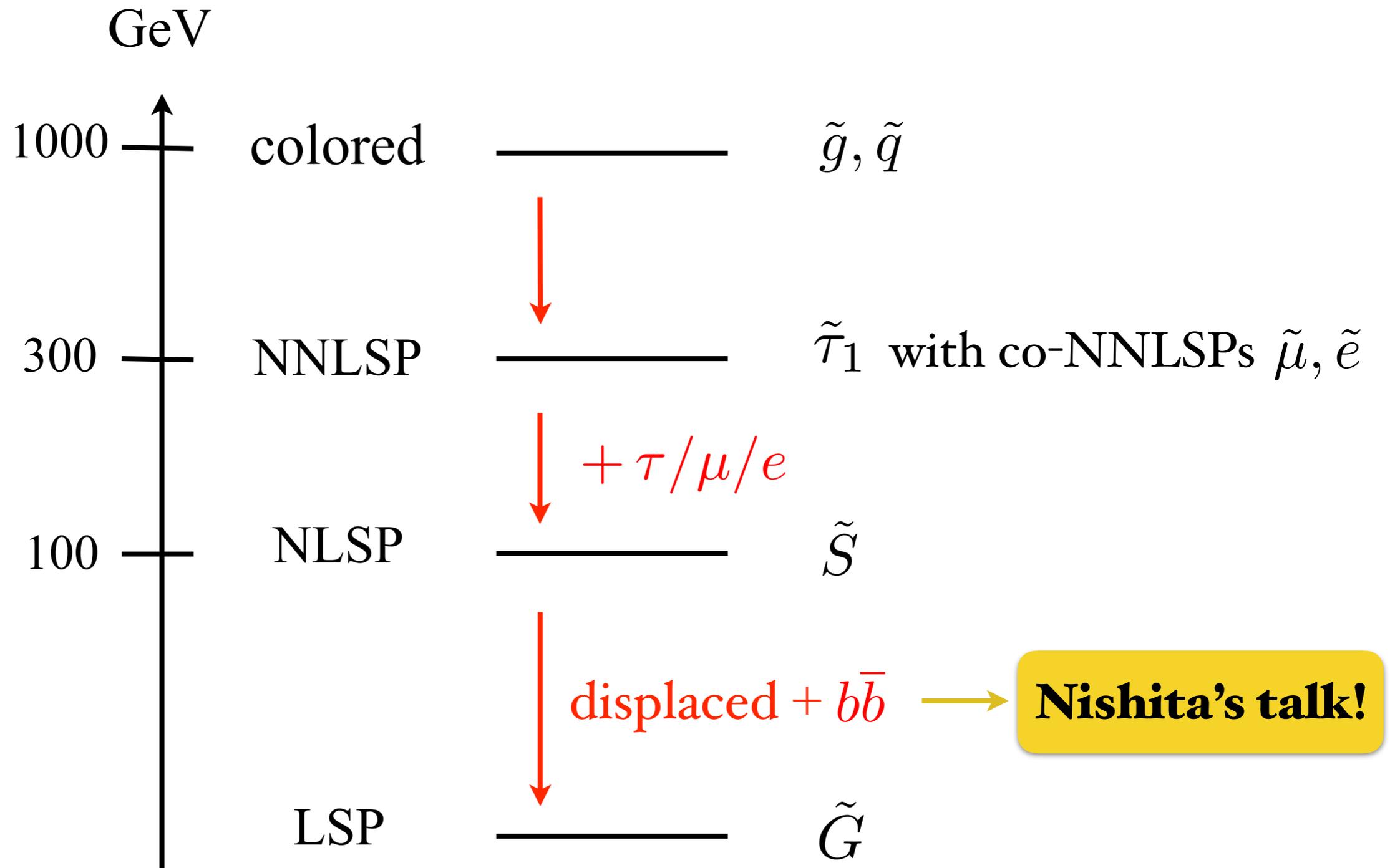
# Medium-M Region: $M \sim 10^7\text{-}9 \text{ GeV}$

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# Low-M Region: $M < 10^7 \text{ GeV}$

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# Outlook

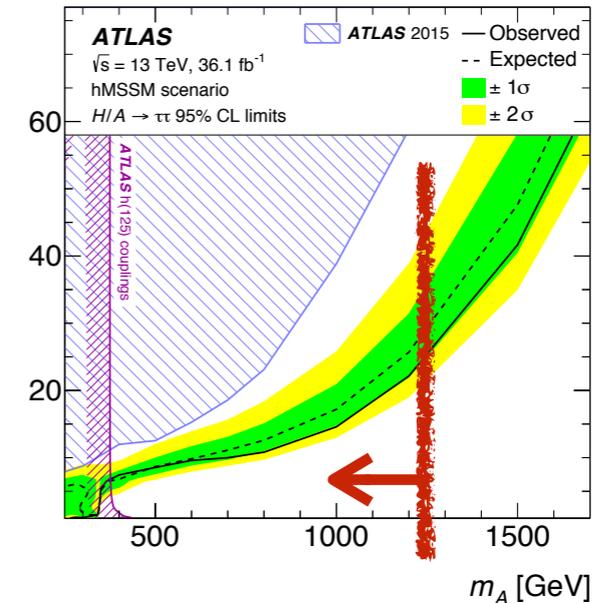
Problem for DGS: NLPS  $a_2$

Since  $\tan\beta$  large, need  $m_{a_2} \gtrsim 1.2 \text{ TeV}$

$$m_{a_2}^2 \approx m_{H_d}^2 - m_{H_u}^2$$



Heavy SUSY spectrum!



Need extra contribs to  $m_{a_2}$  !

1) Combine DGS with Messenger-Matter couplings

$$W_{\text{DGSU}} = S \left( \xi_D \Phi_u^{(1)} \Phi_d^{(2)} + \xi_T \Phi_T^{(1)} \Phi_{\bar{T}}^{(2)} \right) + \lambda_u Q_3 U_3 \Phi_u^{(2)} + \lambda_{S_d} S \Phi_u^{(2)} H_d$$

2) Just single Messenger-Matter couplings enough!?

$$W_U = X (\Phi_u \Phi_d + \Phi_T \Phi_{\bar{T}}) + \lambda_u Q_3 U_3 \Phi_u + \lambda_{S_d} S \Phi_u H_d$$

...

work in progress, with M. Badziak, N. Desai, C. Hugonie

# Summary

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- A 125 GeV Higgs in Minimal Gauge Mediation requires colored sparticles out of LHC reach: motivates extensions of minimal model
- Flavored messenger matter-couplings generate large  $A$ -terms: leads to rich (but viable) flavor pheno between MFV and SUSY PC
- Minimal model for NMSSM + Gauge Mediation allows for light sparticles thanks to Higgs-Singlet mixing: very predictive framework with new collider signatures

Backup

# Flavor constraints

---

Most constraints automatically satisfied for  $\tilde{m} \sim 1$  TeV

$(\delta_{XX}^D)_{12}$	$9.2 \times 10^{-2}$ [Re]	$1.2 \times 10^{-2}$ [Im]
$\langle \delta_{12}^D \rangle$	$1.9 \times 10^{-3}$ [Re]	$2.6 \times 10^{-4}$ [Im]
$(\delta_{LR}^D)_{12}$	$5.6 \times 10^{-3}$ [Re]	$4.0 \times 10^{-5}$ [Im]
$(\delta_{XX}^U)_{12}$	$1.0 \times 10^{-1}$ [Re]	$6.0 \times 10^{-2}$ [Im]
$\langle \delta_{12}^U \rangle$	$6.2 \times 10^{-3}$ [Re]	$4.0 \times 10^{-3}$ [Im]
$(\delta_{LR}^U)_{12}$	$1.6 \times 10^{-2}$ [Re]	$1.6 \times 10^{-2}$ [Im]
$(\delta_{XX}^D)_{13}$	$2.8 \times 10^{-1}$ [Re]	$6.0 \times 10^{-1}$ [Im]
$\langle \delta_{13}^D \rangle$	$4.2 \times 10^{-2}$ [Re]	$1.8 \times 10^{-2}$ [Im]
$(\delta_{LR}^D)_{13}$	$6.6 \times 10^{-2}$ [Re]	$1.5 \times 10^{-1}$ [Im]
$(\delta_{LR}^D)_{11}$	$2.0 \times 10^{-6}$	
$(\delta_{LR}^U)_{11}$	$4.0 \times 10^{-6}$	

$D - \bar{D}$  mixing

$$(\delta_{RR}^u)_{12} \sim (\lambda_U^*)_{31} (\lambda_U)_{32}$$

Neutron EDM

$$(\delta_{LR}^u)_{11} \sim (\lambda_U)_{13} (\lambda_U)_{31}$$