

# $B + L$ violation at colliders and new physics

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May 31, 2018



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SEIT 1386

based on JHEP 1804 (2018) 076 (arxiv:1801.03492) with  
D.G. Cerdeño, C. Tamarit, K. Sakurai

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# $B + L$ violation+Dark Matter $\rightarrow$ at Colliders?

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# Baryon Number Violating Processes

In the Standard Model:

- accidental baryon and lepton symmetries
- has been confirmed by all experiments so far
- broken by quantum anomalies  $\rightarrow$  Instanton Theory

Impact on physics of the early Universe:

- important in electroweak baryogenesis
- play a crucial role for baryogenesis through leptogenesis

**DM can take part in sphaleron processes! (seen later)**

# $B + L$ violation from $SU(2)$ Anomalies

- global anomaly for  $U(1)$  baryon and lepton charges:

$$\int \partial_\mu J_{B+L}^\mu = 6n_{\text{top}} = 6\Delta N_{\text{CS}}$$

- chiral anomaly on a quantum level

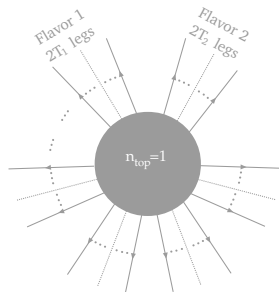
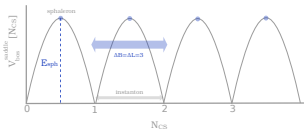
$$\Delta Q_{\text{chiral}} = \int \partial_\mu J_{\text{chiral}}^\mu = n_{\text{top}} \sum_k 2T_k$$

with the topological charge  $n_{\text{top}} = \frac{g^2}{16\pi^2} \int \text{Tr}(F_{\mu\nu} \tilde{F}^{\mu\nu}) = 1, 2, \dots$  and representation  $k$  with Dynkin index  $T_k$ .

General

## General Case:

- electric neutrality
- neutral under  $SU(3)_C$
- number of left-handed fermions determined by chiral anomaly
- “mixture” of particles determined by violation of additional global charges



## $B + L$ violation from $SU(2)$ Anomalies

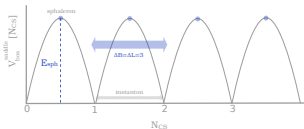
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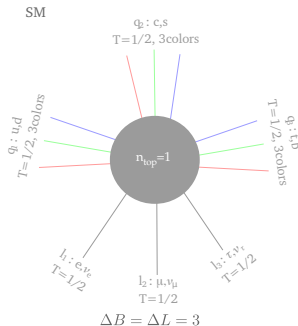
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### SM Case:

- $B + L$  violation by 6 (for  $n_{\text{top}} = 1$ )
- 12 left-handed Weyl-fermions in fundamental representation  $T_k = 1/2$
- SM vertex:

$$\mathcal{O}_{\text{SM}} = \prod_{i=1,2,3} (q_L q_L q_L l_L)_i$$



# $B + L$ violation from $SU(2)$ Anomalies

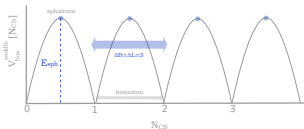
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How can BSM influence the vertices?

Possible BSM extensions:

- $\Delta(B + L) = 6$  plus additional charge violation?
- 12 left-handed Weyl-fermions in fundamental representation  
 $T(r) = 1/2$  plus more  $SU(2)$  fermions?
- additional fermion number violation strongly constrained
- But additional  $SU(2)$  fermions in certain representation  $k$  with  $T_k$  can extend the vertices!

## Anomalous vertices with massive fermions

Mass terms for Weyl fermions  $\psi_{k,l}$ :

$$\mathcal{L} \supset \underbrace{m_{kl}}_{-2} \underbrace{\psi_k}_{+1} \underbrace{\psi_l}_{+1} + \text{c.c.}$$

Classical symmetry restored by treating  $m$  as charged field.

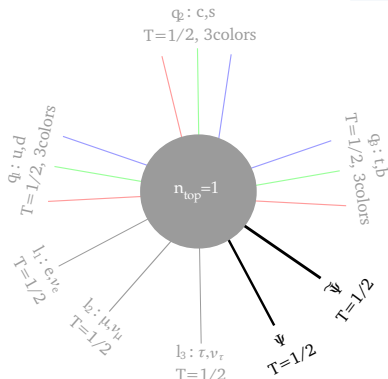
- symmetry violated as before by anomaly
- anomalous interaction vertices carry same charge as before
- spurious charge can be carried by mass insertions  $m_{kl}^*$



# Taking into Account more SU(2) particles

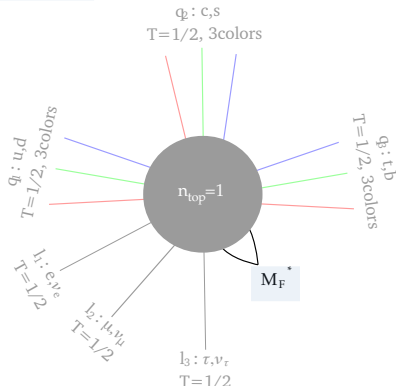
SM + Heavy Dirac in fundamental:

$$\mathcal{L} \supset -M\tilde{\Psi}\Psi + c.c.$$



BSM vertex

$$\Delta B = \Delta L = 3$$

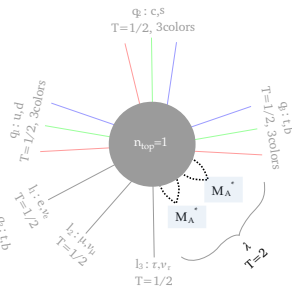
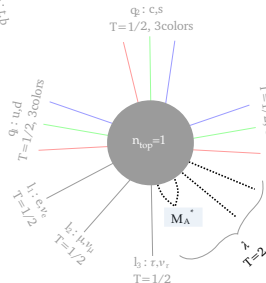
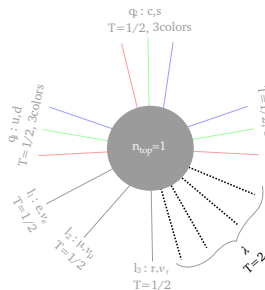


SM vertex!

# Taking into Account more SU(2) particles

SM + Heavy Weyl in adjoint:

$$\mathcal{L} \supset -M_A \lambda \lambda + c.c.$$

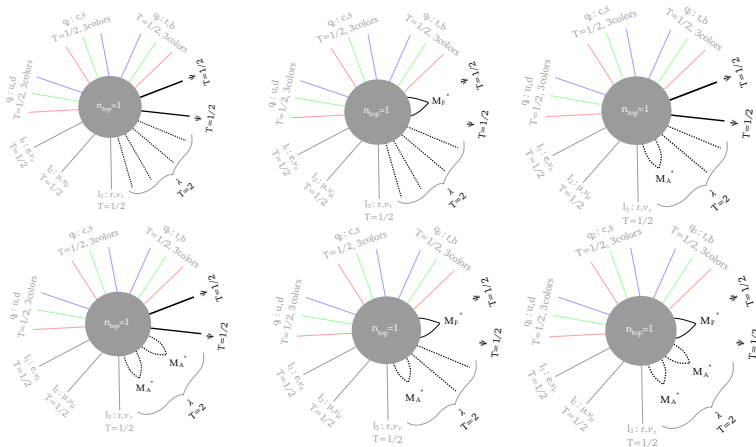


Up to 16 fermions!

$$\Delta B = \Delta L = 3$$

# Taking into Account more SU(2) particles

SM + Heavy Dirac in fundamental + Heavy Weyl in adjoint:



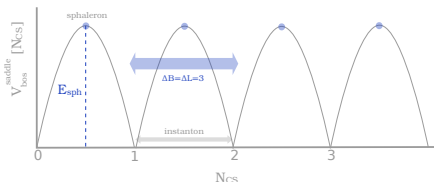
Up to 18 fermions!

$$\Delta B = \Delta L = 3$$

# Can we observe such processes at colliders?

## Dressing fermion amplitudes with gauge bosons:

First instanton estimates yield exponential corrections to rate when summing over gauge boson emission [Ringwald, O.Espinosa]



$$\sigma_{B+L, \text{leading}}^{2 \rightarrow \text{any}} = f(\hat{s}) e^{-\frac{4\pi}{\alpha_W} F[\sqrt{\hat{s}}/E_0]},$$

$$F\left[\frac{\sqrt{\hat{s}}}{E_0}\right] = 1 - \frac{9}{8} \left(\frac{\sqrt{\hat{s}}}{E_0}\right)^{4/3}, \quad E_0 = \frac{6\pi m_W}{\alpha_W} \gtrsim E_{\text{sph}}$$

$\Rightarrow$  unsuppressed above  $E_{\text{sph}}$ , but in conflict with unitarity arguments

## Instanton loop corrections:

$$F\left[\frac{\sqrt{\hat{s}}}{E_0}\right] = 1 - \frac{9}{8} \left(\frac{\sqrt{\hat{s}}}{E_0}\right)^{4/3} + \frac{9}{16} \left(\frac{\sqrt{\hat{s}}}{E_0}\right)^2 + \dots$$

breaks down above  $\sqrt{\hat{s}} > E_0 \dots$

# Can we observe such processes at colliders?

**Further progress** [Arnold & Mattis, Khlebnikov & Rubakov & Tinyakov, Mueller, Khoze & Ringwald]

- cross-section looks like

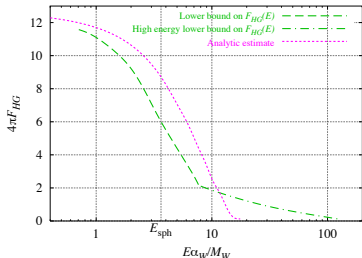
$$\sigma_{B+L, \text{leading}}^{2 \rightarrow \text{any}} = f(\hat{s}) e^{-\frac{4\pi}{\alpha_W} F[\sqrt{\hat{s}}/E_0]}$$

- with:

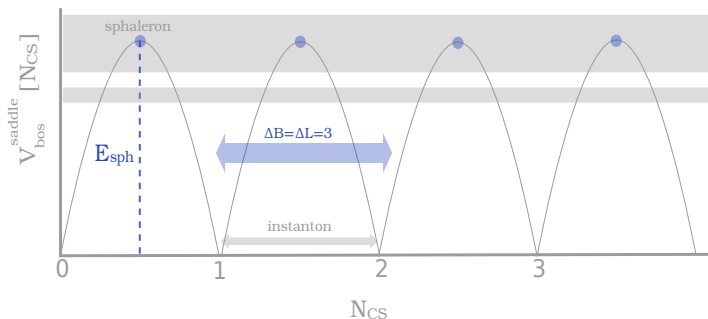
- 1 Unitarity bound  $F \gtrsim 0$   
[Zakharov, Maggiore & Shifman, Veneziano]
- 2 Dispersion relation  
[Zakharov, Porrati, Khoze & Ringwald]

$$f(\hat{s}) = \frac{1}{m_W^2} \left( \frac{2\pi}{\alpha_W} \right)^{7/2}$$

- 3 lower bound on  $F$   
[Rubakov & Tinyakov, Bezrukov et al]



## Can we observe such processes at colliders?



### Problem solved?

Take  $N_{\text{CS}}$  as kinetic variable and get 1D dynamics in the  $N_{\text{CS}}$  direction

$\Rightarrow$  one has a band structure and tunneling may be unsuppressed for  $\sqrt{\hat{s}} \geq E_{\text{sph}}$

[Tye & Wong, Funakubo et al]

$\Rightarrow$  arises attention: BaryoGEN [1805.02786], CMS study [1805.06013]

## How to Add More Particles

Calculate effective couplings with ordinary instanton techniques:

$$\langle \psi_{k_1} \dots \psi_{k_N} \rangle = \int \prod_m [d\psi]_m [d\psi^\dagger]_m \sum_{n_{\text{top}}=q} \int [dA_\mu^a]_q \exp(-S) \psi_{k_1} \dots \psi_{k_N}$$

- Saddle point expansion around extremal of  $S$
- 1-instanton dominates:  $\exp(-S_{\text{inst}}) = \exp\left(-\frac{4\pi}{\alpha_W} n_{\text{top}}\right)$
- Group Weyl fermions  $\psi_k, \psi_l$  into Dirac  $\Psi_{kl}$

$$\langle \bar{\Psi}_{kl} \Psi_{kl} \rangle \propto \det \left( \frac{\delta^2 S}{\delta \bar{\Psi} \delta \Psi} \right) \prod (\text{Fermion prop. in instanton bg.})$$

- Ensure compatibility with decoupling: Get SM-like vertices for heavy exotic fermions

We consider ratios of rates

$$E(\sqrt{\hat{s}}, \delta, M) = \frac{\Gamma_{\text{model}}^{10+\delta, n_h, n_W}}{\Gamma_{\text{model}'}^{10,0,0}} \sim \frac{\text{BSM}}{\text{SM-like}}$$

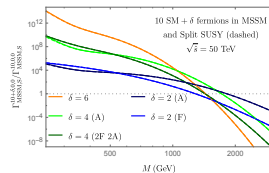
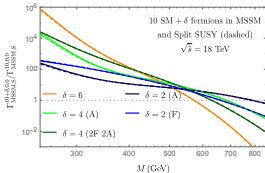
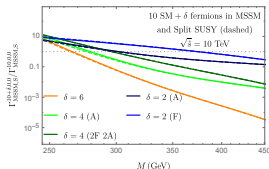
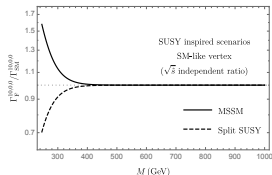
fermions have no impact on  $F[\sqrt{\hat{s}}] \rightarrow$  cancels in ratio

# Candidates

## BSM:

The “obvious” choice for models with additional  $SU(2)$  fermions is SUSY

- MSSM: **SM+ fundamental fermions (Higgsino) + adjoint fermions (Wino)**  
+ gluino+ bino+ SUSY scalars
- Split SUSY: **SM+ fundamental fermions + adjoint fermions**



**Sphaleron rate is significantly enhanced compared to the SM if the underlying theory has relatively light BSM fermions charged under  $SU(2)$ !**



# Candidates

## Dark Matter:

several DM candidates in various  $SU(2)$  representations

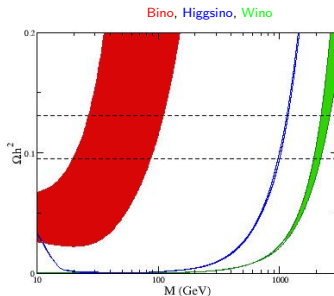
⇒ focus on doublet and triplet with Spin 1/2

Quantum numbers $SU(2)_L$ $U(1)_Y$ Spin	DM can decay into	DM mass in TeV	$m_{DM^\pm} - m_{DM}$ in MeV	Events at LHC $\int \mathcal{L} dt = 100/\text{fb}$	$\sigma_{SI}$ in $10^{-45} \text{ cm}^2$
2 1/2 0	$EL$	$0.54 \pm 0.01$	350	$320 \div 510$	0.2
2 1/2 1/2	$EH$	$1.1 \pm 0.03$	341	$160 \div 330$	0.2
3 0 0	$HH^*$	$2.0 \pm 0.05$	166	$0.2 \div 1.0$	1.3
3 0 1/2	$LH$	$2.4 \pm 0.06$	166	$0.8 \div 4.0$	1.3
3 1 0	$HH, LL$	$1.6 \pm 0.04$	540	$3.0 \div 10$	1.7
3 1 1/2	$LH$	$1.8 \pm 0.05$	525	$27 \div 90$	1.7
4 1/2 0	$HHH^*$	$2.4 \pm 0.06$	353	$0.10 \div 0.6$	1.6
4 1/2 1/2	$(LHH^*)$	$2.4 \pm 0.06$	347	$5.3 \div 25$	1.6
4 3/2 0	$HHH$	$2.9 \pm 0.07$	729	$0.01 \div 0.10$	7.5
4 3/2 1/2	$(LHH)$	$2.6 \pm 0.07$	712	$1.7 \div 9.5$	7.5
5 0 0	$(HHH^*H^*)$	$5.0 \pm 0.1$	166	$\ll 1$	12
5 0 1/2	—	$4.4 \pm 0.1$	166	$\ll 1$	12
7 0 0	—	$8.5 \pm 0.2$	166	$\ll 1$	46

[Cirelli, Fornengo, Strumia]

Sweet spot to obtain the DM relic density:

- 1 Higgsino ( doublet) with  $m_\chi \sim 1$  TeV
- 2 Wino (triplet)  $m_\chi \sim 2$  TeV  
+Sommerfeld effect  
⇒  $m_\chi \sim 2.7 - 3$  TeV [hep-ph/0610249]



[Arkani-Hamed, Delgado, Giudice]

# Cross Sections at Colliders

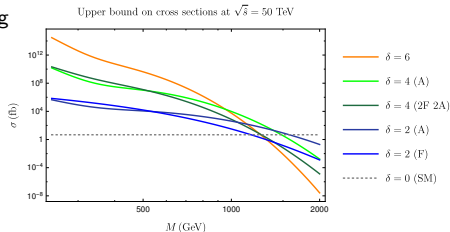
By comparing our results with former instanton calculations, we see that our results seem to be compatible with simple Ansatz

$$\sigma_{B+L, \text{leading}}^{2 \rightarrow \text{any}} = \frac{E(\sqrt{\hat{s}}, \delta, M)}{m_W^2} \left( \frac{2\pi}{\alpha_W} \right)^{7/2} e^{-\frac{4\pi}{\alpha_W} F[(\sqrt{\hat{s}} - \delta M)/E_0]}$$

for partonic cross sections.

Upper bounds for cross sections obtained by using

- ❶ lower bound on Holy Grail function from [Bezrukov et al]
- ❷ SM prefactor from [Khoze, Ringwald]
- ❸ our BSM enhancement factors



# Cross Sections at Colliders

Convoluting with pdfs extrapolated to 100 TeV collider (with  $\mathcal{L} = 30 \text{ ab}^{-1}$ ) yields

SM:

$$\sigma \lesssim 6 \times 10^{-5} \text{ fb} \rightarrow N \lesssim 1.8 \text{ events}$$

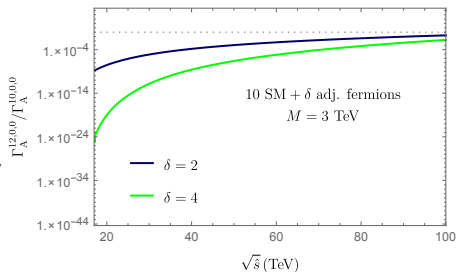
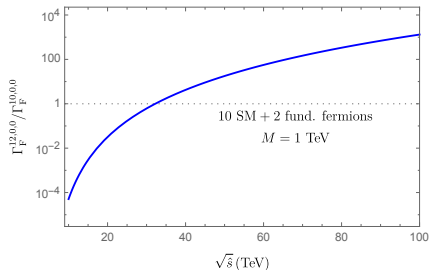
Higgsino DM with  $m_\chi = 1 \text{ TeV}$

$$\sigma \lesssim 4.9 \times 10^{-3} \text{ fb} \rightarrow N \lesssim 150 \text{ events}$$

Wino DM with  $m_\chi = 3 \text{ TeV}$

$$\sigma \lesssim 4 \times 10^{-10} \text{ fb} \rightarrow N \lesssim 1.2 \times 10^{-5} \text{ events}$$

If DM is 1 TeV Higgsino,  
and taking the lower bound of the Holy Grail  
function, sphalerons should be observable  
at 100 TeV  
 $\Rightarrow$  Holy Grail functions can be tested at 100 TeV



## Summary & Outlook

- $SU(2)$  anomalies predict new  $B + L$ -violating interaction vertices in the presence of BSM fermions charged under the gauge group
- Rates of new processes can be orders of magnitude above the SM rate!
- ratios of rates are independent of overall unknown cross section
- DM might make sphalerons observable and/or test Holy Grail functions
- If  $B + L$  violating interactions are ever seen at colliders, they could be tied to BSM physics