



Workshop on future instruments for  
gravity-based earthquake early warning

# Gradiometers based on atom interferometers

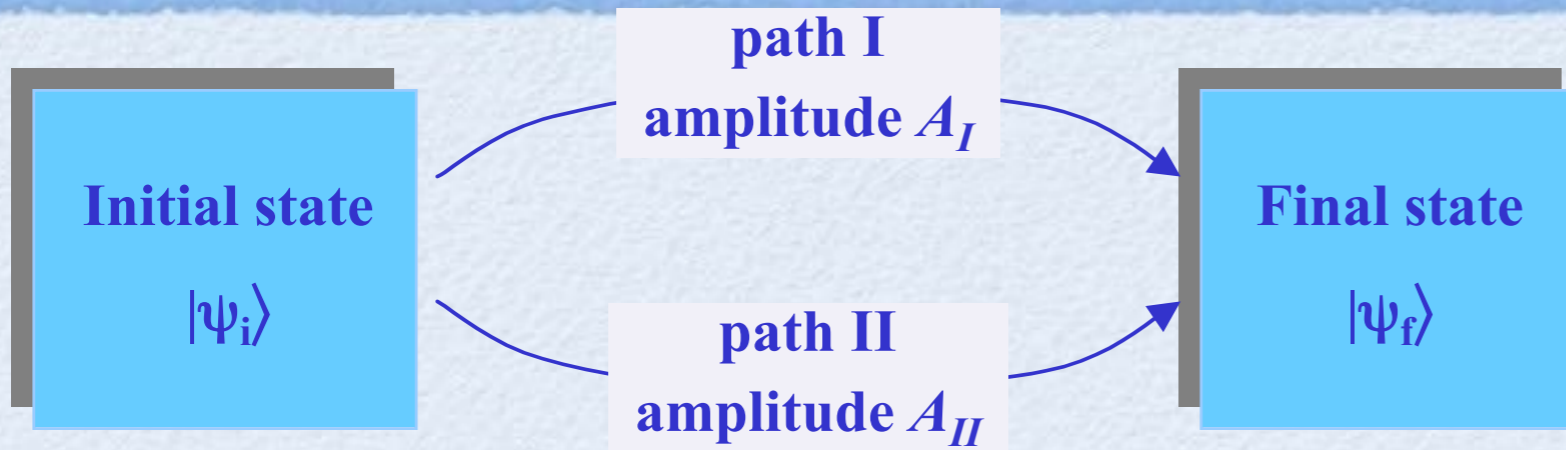
Fiodor Sorrentino  
INFN - Genova

# Outline



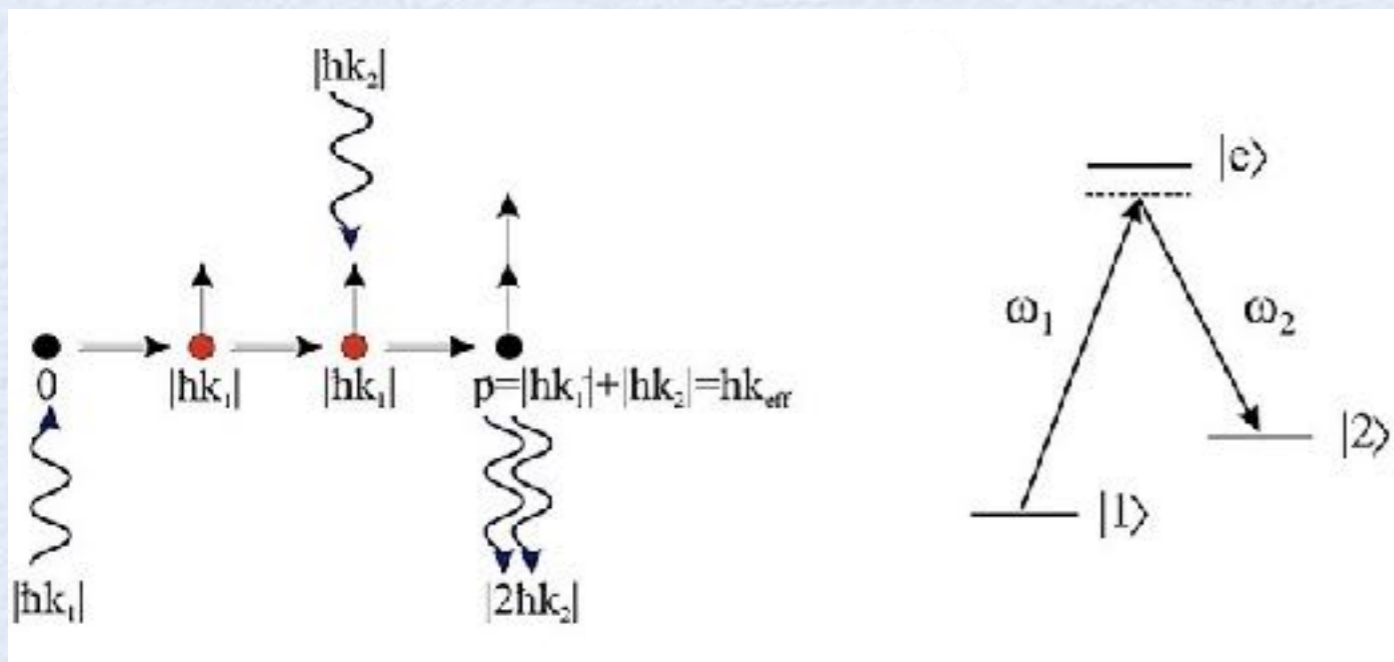
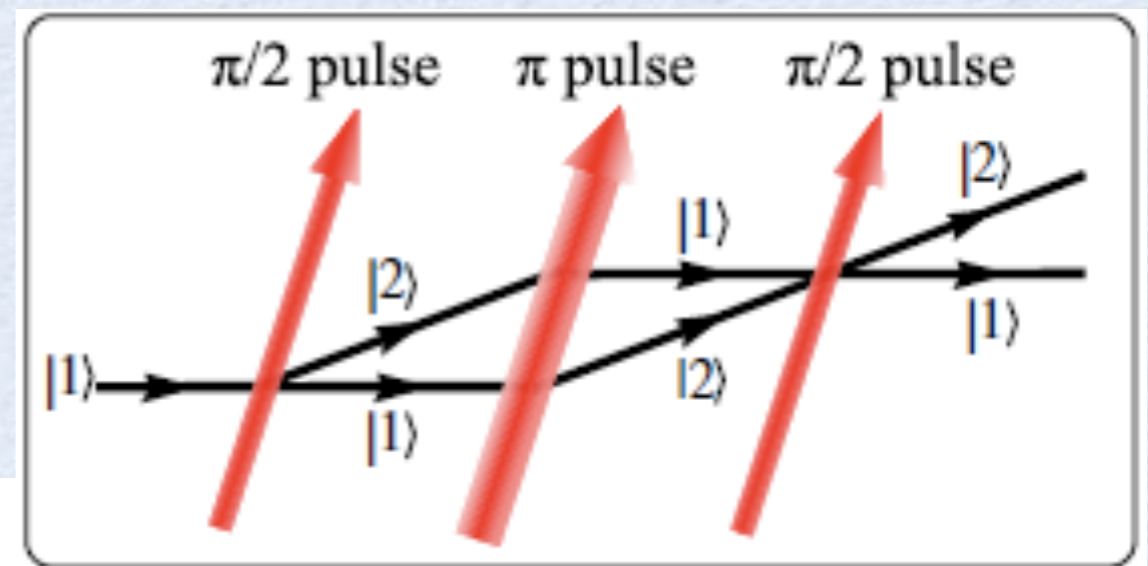
- Brief recap on AI gravity sensors:
  - atom optics, ultracold atoms & atom interferometers
  - AI gravity gradiometers: seism CMRR & sensitivity limits
  - state of the art of atomic sensors
    - laboratory experiments
    - compact AI sensors: TRL
- Perspectives and potential improvements
  - medium-large scale experiments
  - atom optics progress
  - alternative schemes
    - trapped atoms
    - fiber links

# Atom interferometry



Interference of de Broglie amplitudes

Light-pulse beam splitters  
+ fluorescence detection



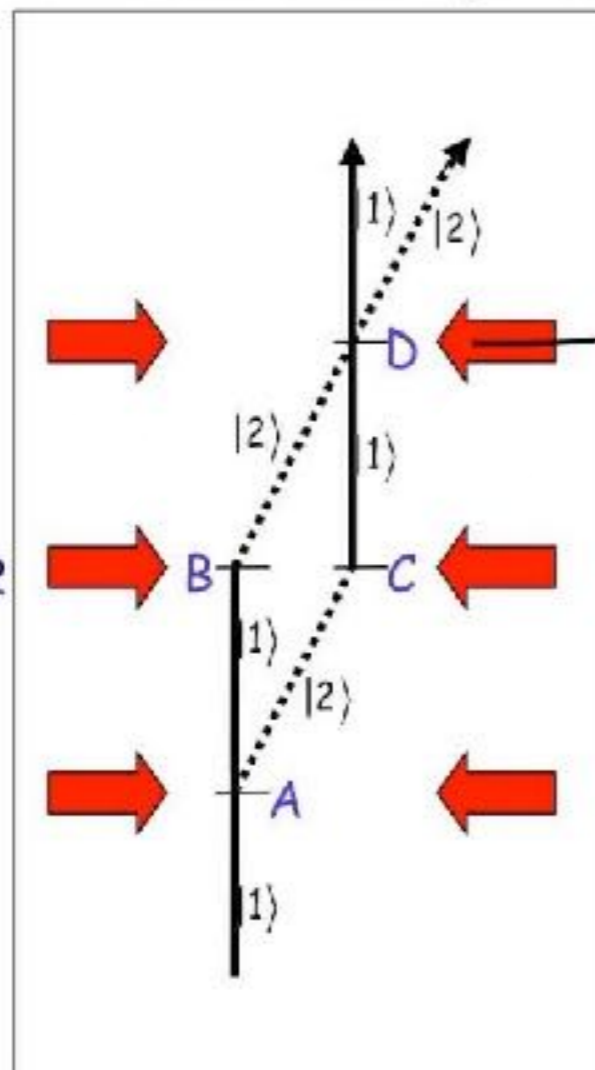
Output phase selectively sensitive to different effects (inertial, gravitational, external fields, laser phase/frequency, etc) via choice of quantum states

# Light-pulse AI inertial sensors



## TRANSVERSAL PULSES

- the interferometer encloses an area
- used to measure rotations (GYROSCOPES)



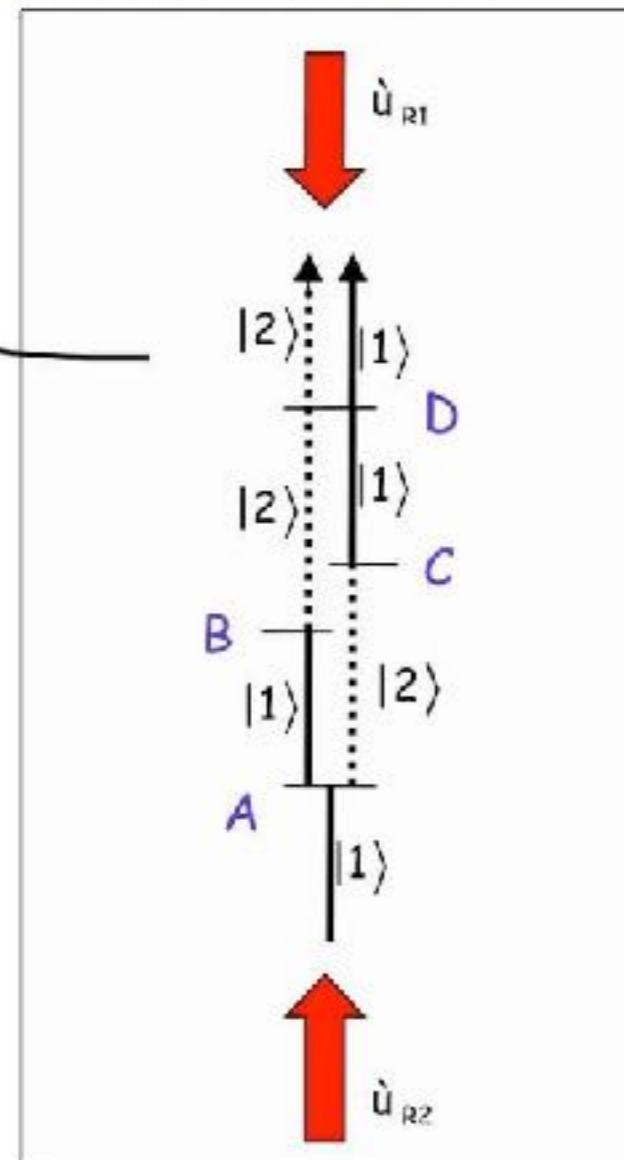
With an acceleration  $g$ ,  
the phase difference

$$\Delta\phi = 2k_{\text{eff}} \cdot (a - 2(\Omega \times v)) T^2$$

where  $k$  is the laser  
wavenumber and  $T$   
the time interval  
between laser pulses

## LONGITUDINAL PULSES

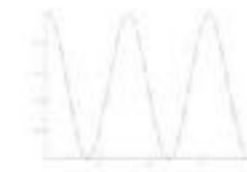
- no area enclosed
- used to measure accelerations (GRAVIMETERS)



With an acceleration  $g$ ,  
the phase difference

$$\Delta\phi = k_{\text{eff}} g T^2$$

where  $k$  is the laser  
wavenumber and  $T$   
the time interval  
between laser pulses



Fringes  
detected

# Current performance



- Absolute gravimeters

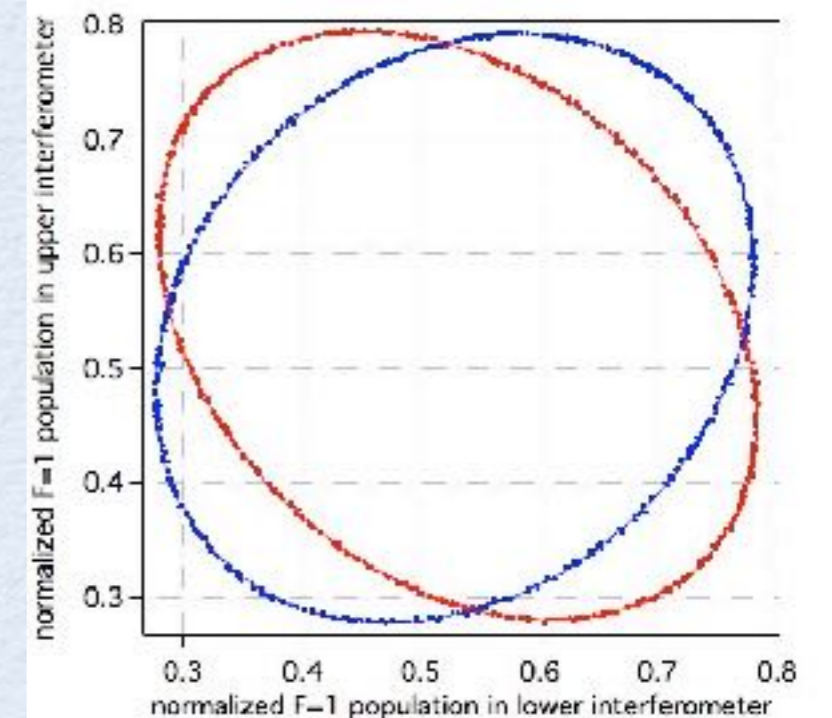
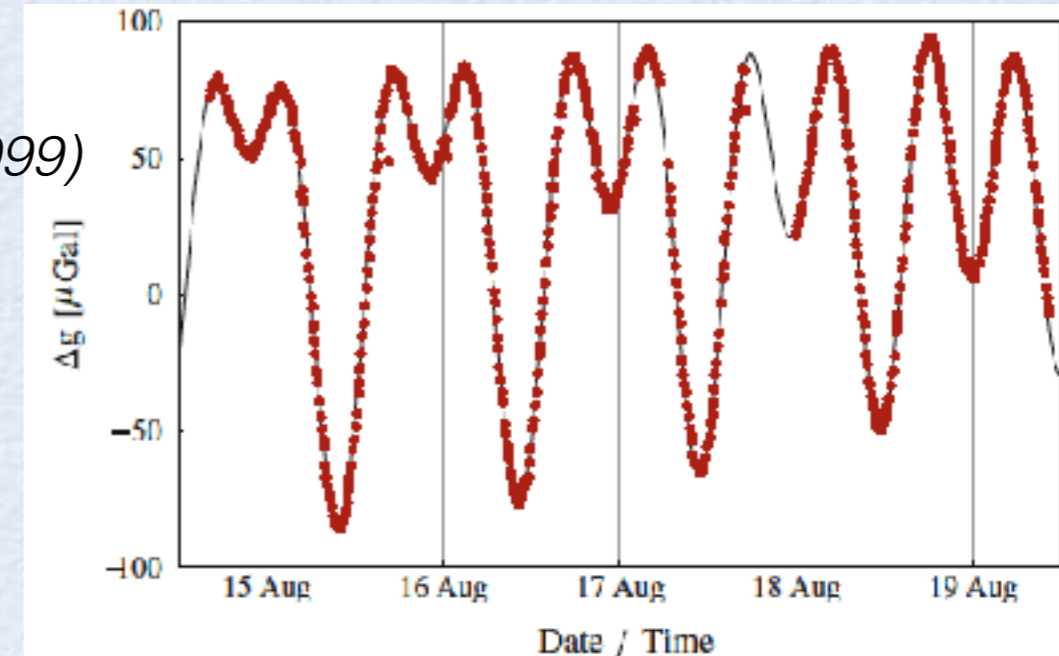
- resolution:  $3 \times 10^{-9}$  g in 1 second (SYRTE)
- averaging down to  $2 \times 10^{-10}$  g after 30 min (SYRTE)
- accuracy:  $10^{-9} \div 10^{-10}$  g, limited by tidal models
  - *A. Peters, K.Y. Chung and S. Chu, Nature 400, 849 (1999)*
  - *H. Müller et al., Phys. Rev. Lett 100, 031101 (2008)*
  - *M. Hauth et al., Appl. Phys. B 113, 49 (2013)*
  - *P. Gillot et al., Metrologia 51, L15 (2014)*

- Gravity gradiometers

- 1 m fountain: differential acceleration sensitivity:  $5 \times 10^{-9}$  g in 1 s ( $5 \times 10^{-11}$  g after  $10^4$  s)
  - *F. Sorrentino et al., Phys. Rev. A 89, 023607 (2014)*
- 10 m fountain: differential acceleration sensitivity:  $5 \times 10^{-10}$  g in 1 s
  - *P. Asenbaum et al., Phys. Rev. Lett. 118, 183602 (2017)*

- Rotation sensors

- sensitivity:  $6 \times 10^{-10}$  rad/s/ $\sqrt{\text{Hz}}$
- scale factor stability  $< 5$  ppm
- bias stability  $< 70$   $\mu\text{deg/h}$ 
  - *T. L. Gustavson et al., Class. Quantum Grav. 17, 2385 (2000)*
  - *D. S. Durfee et al., Phys. Rev. Lett. 97, 240801 (2006)*

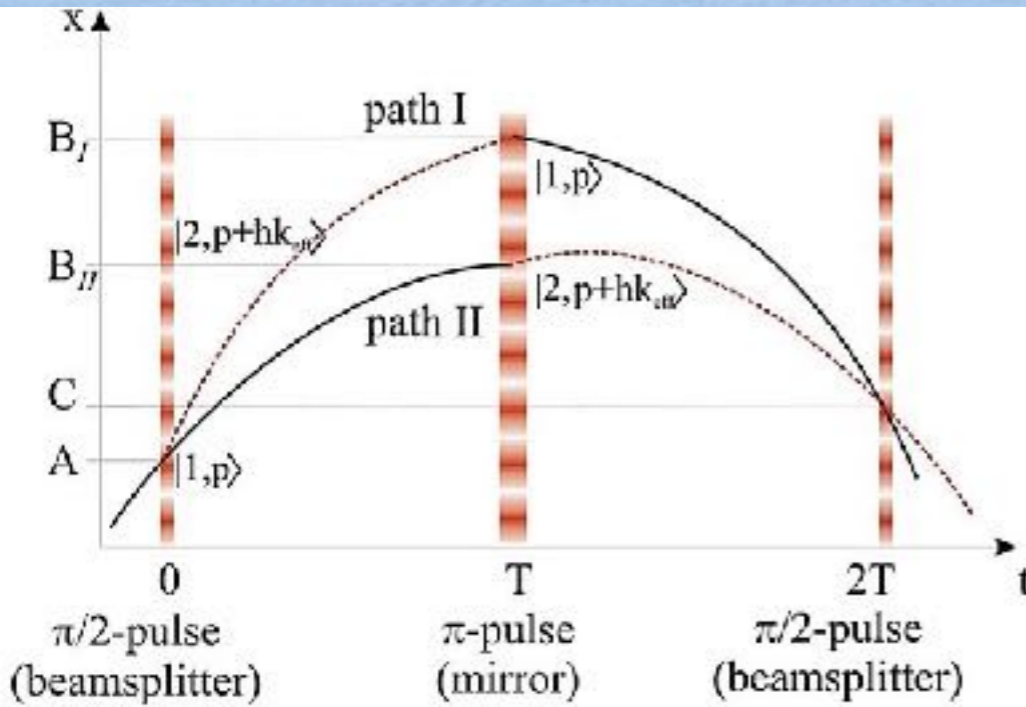


# Sensitivity limits



- Short-term sensitivity in state-of-the art atom interferometers is usually limited by **quantum noise**
- For long integration times, and / or for major sensitivity improvements one should consider several effects
  - Acceleration noise
  - Laser wavefront
  - Atomic motion
  - Laser phase and frequency noise
  - Aliasing
  - External fields
  - Other technical noise sources

# Quantum projection noise



Phase difference between the paths:  
 $\Delta\Phi = k_e[z(0) - 2z(T) + z(2T)] + \Phi_e$

$$k_e = k_1 - k_2$$

with  $z(t) = -gt^2/2 + v_0t + z_0$  &  $\Phi_e = 0$

$$\rightarrow \Delta\Phi = k_e g T^2$$

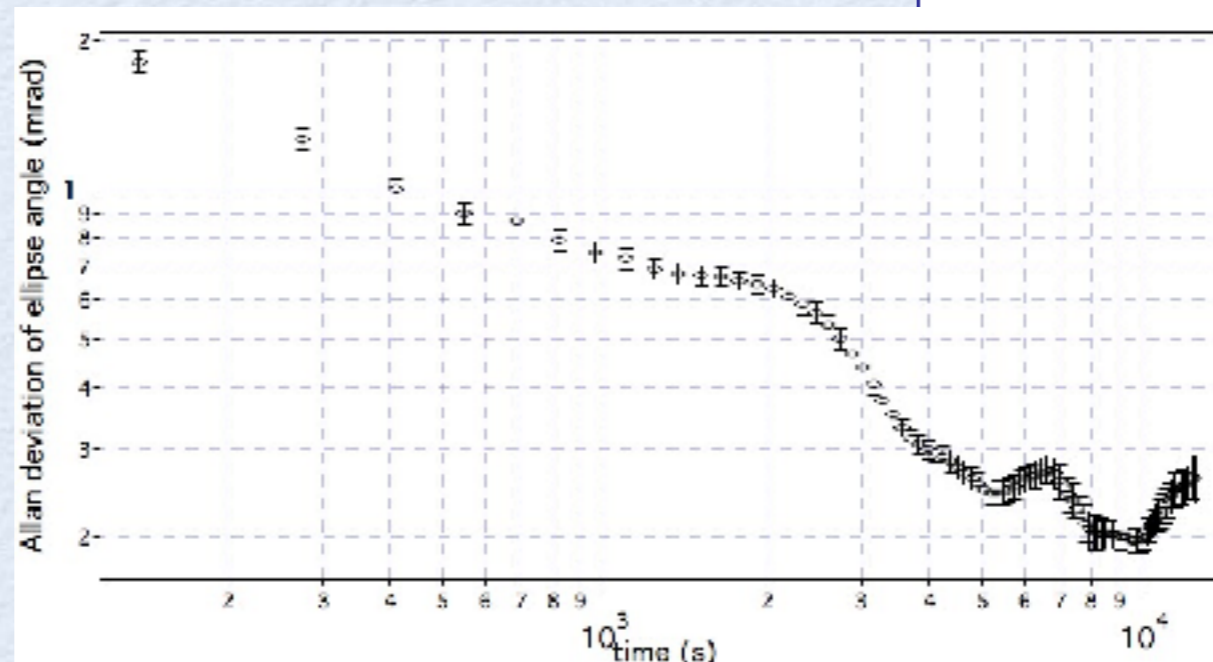
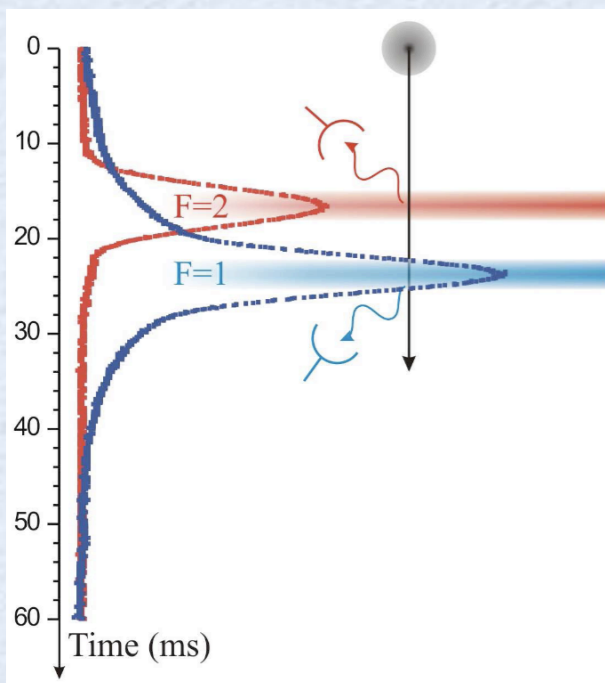
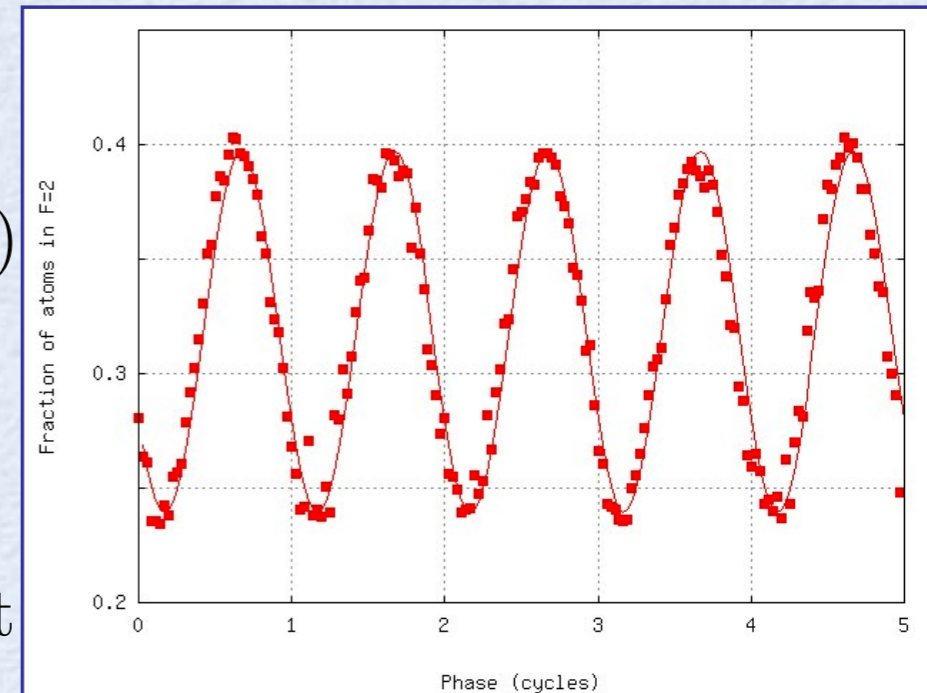
Final population:

$$N_a = N/2(1 + \cos[\Delta\Phi])$$

$$\delta\Phi \sim \frac{1}{\sqrt{(N_{at})}}$$

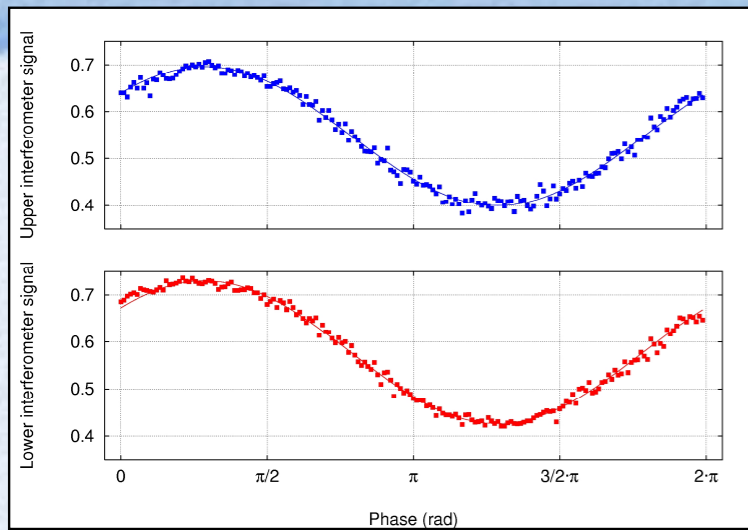
$$T = 150 \text{ ms} \rightarrow 2\pi = 10^{-6} \text{ g}$$

$$S/N=1000 \rightarrow \text{Sensitivity } 10^{-9} \text{ g/shot}$$



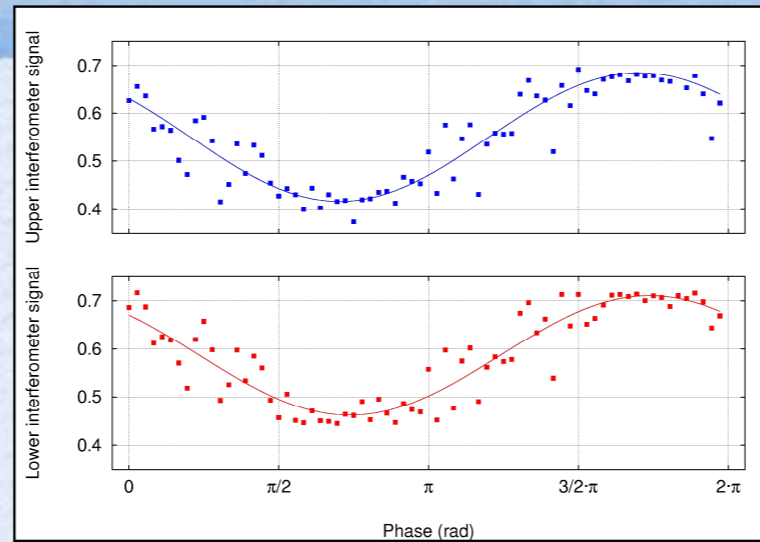
metry gradiometers

# Acceleration noise



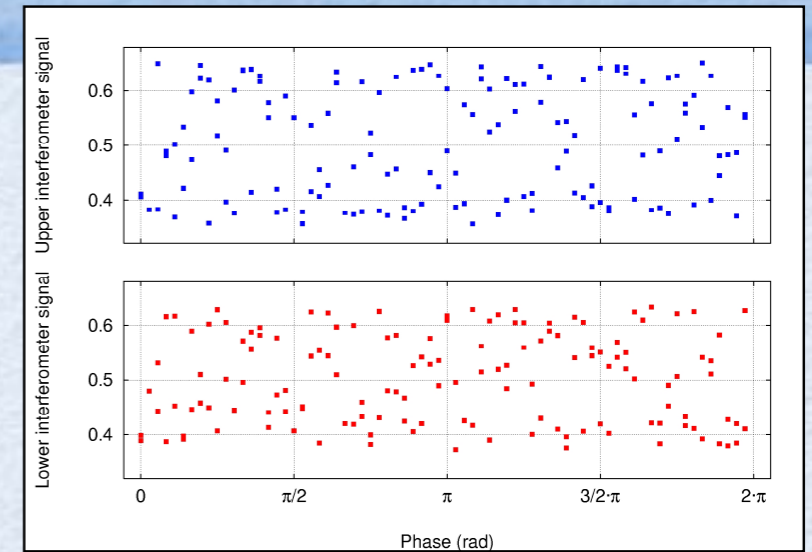
$T=5$  ms

resol. =  $2.3 \times 10^{-5}$  g/shot



$T=50$  ms

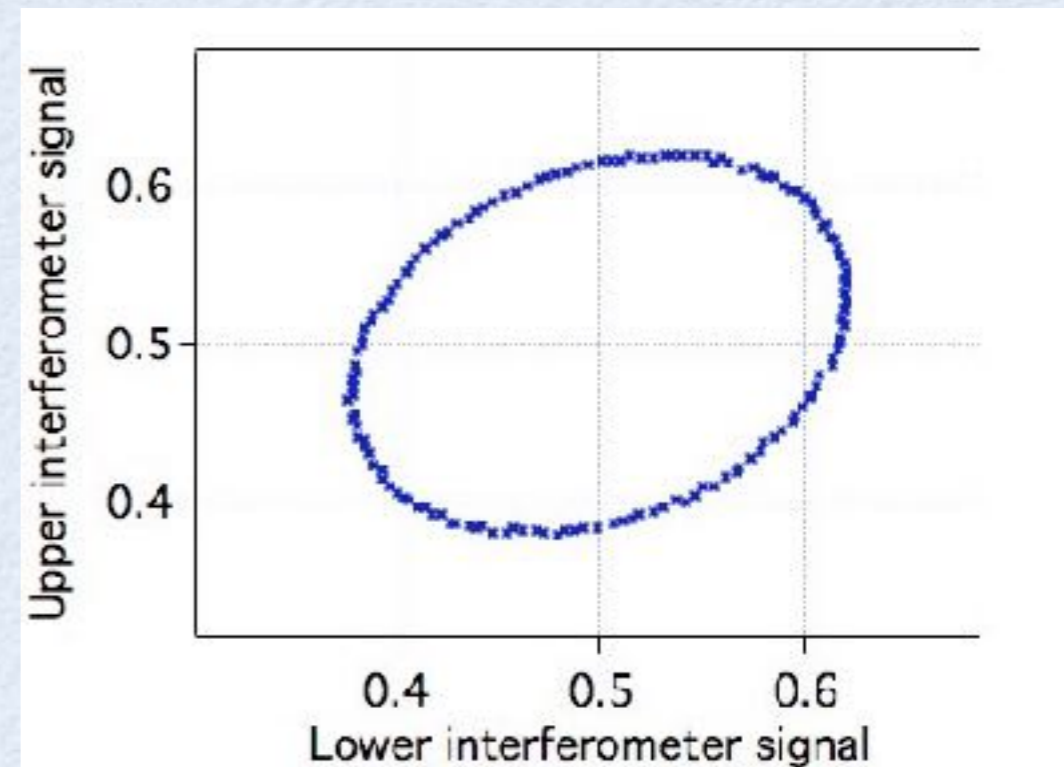
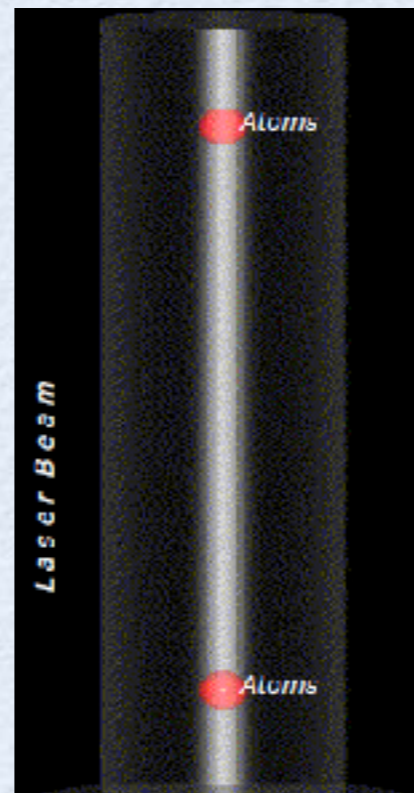
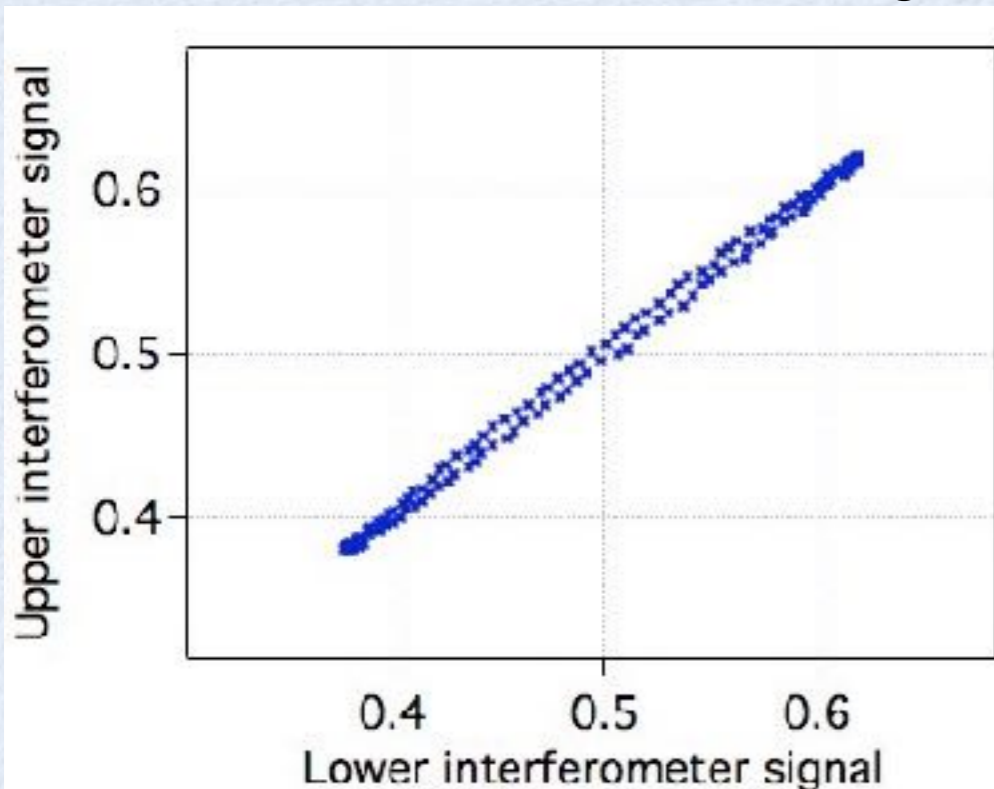
resol. =  $1.0 \times 10^{-6}$  g/shot



$T=150$  ms

resol. =  $3.2 \times 10^{-8}$  g/shot

$$\Delta\Phi = k_e g T^2$$

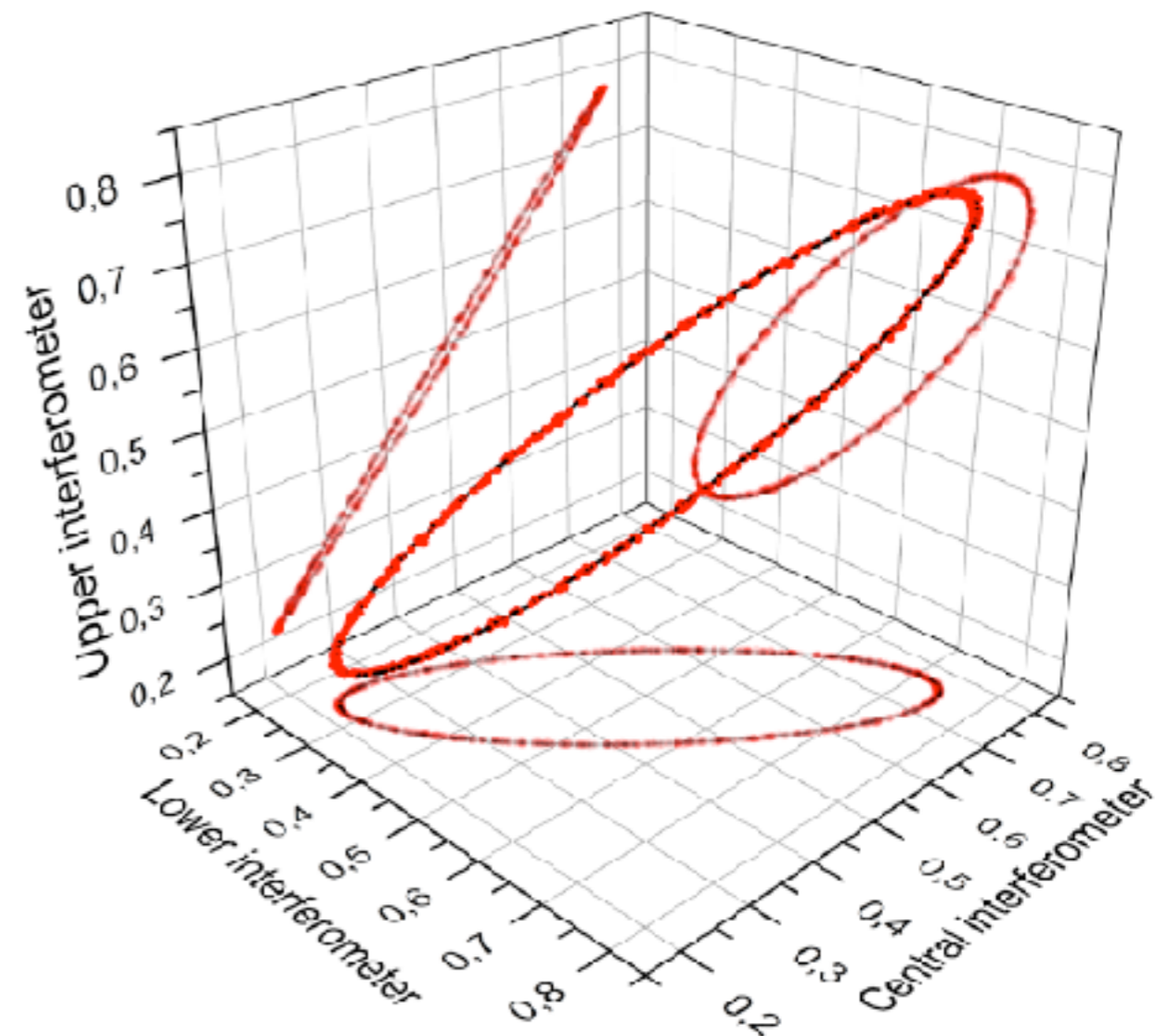
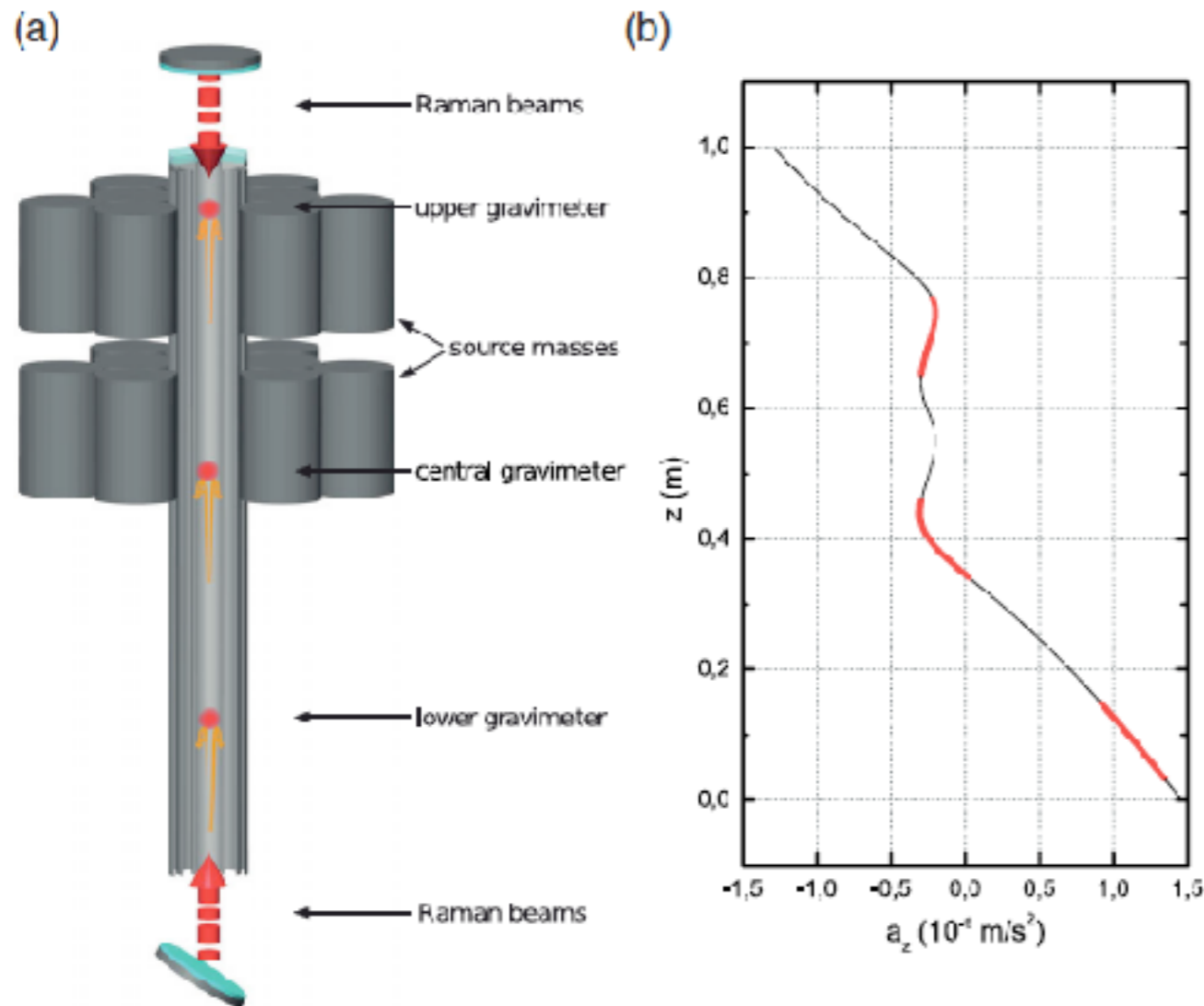


up to 140 dB CMRR with 1 m separation

G. T. Foster *et al.*, *Opt. Lett* **27**, 951 (2002)



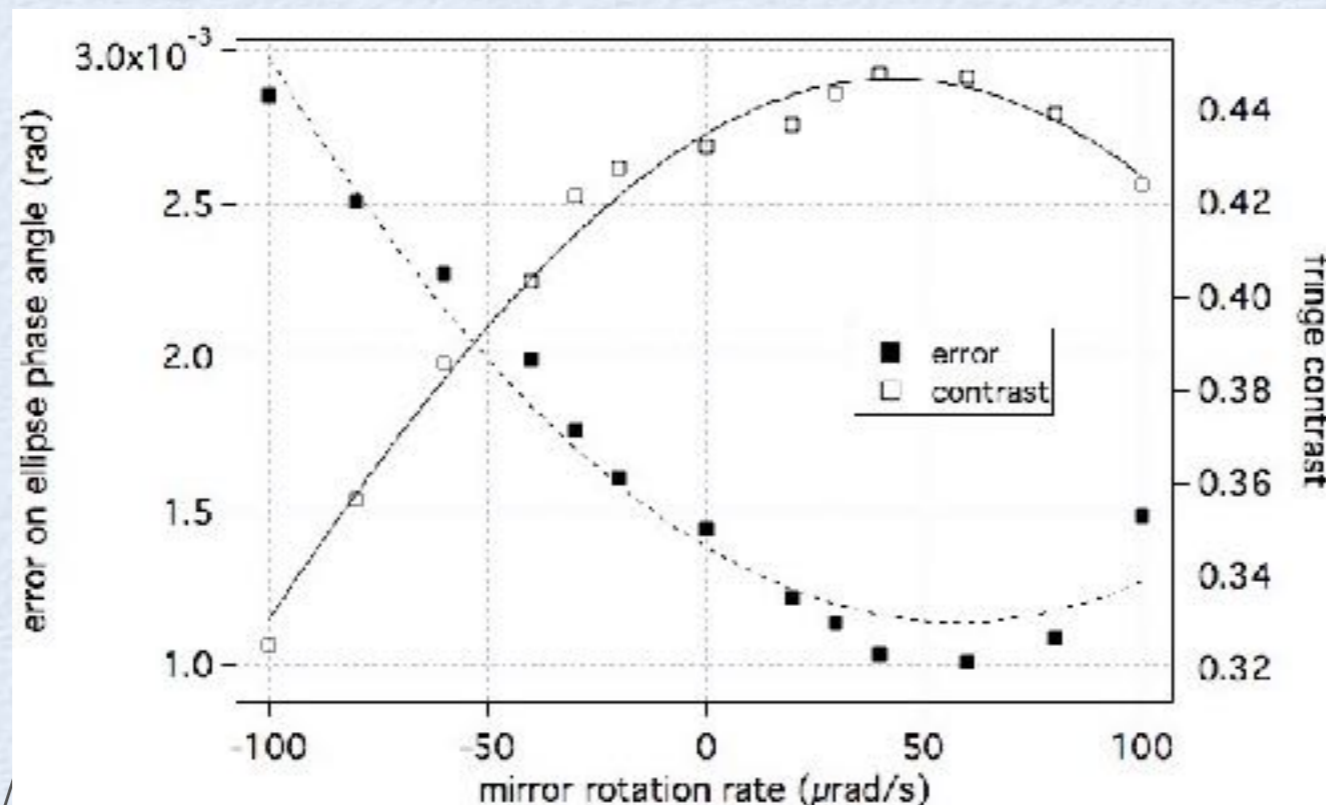
# Multiple samples



- Use three atomic clouds to measure the vertical derivative of vertical gravity gradient
  - *G. Rosi et al., Phys. Rev. Lett. 114, 013001 (2015)*
- Scalable to arbitrary large number of samples
- Simultaneous, correlated AI can improve g measurements as well
  - *F. Sorrentino et al., Appl. Phys. Lett. 101, 114104 (2012)*
- Multiple, correlated AIs to mitigate NN
  - *W. Chaibi et al, Phys. Rev. D 93, 021101 (2016)*
- See talk by R. Geiger

# Atomic motion

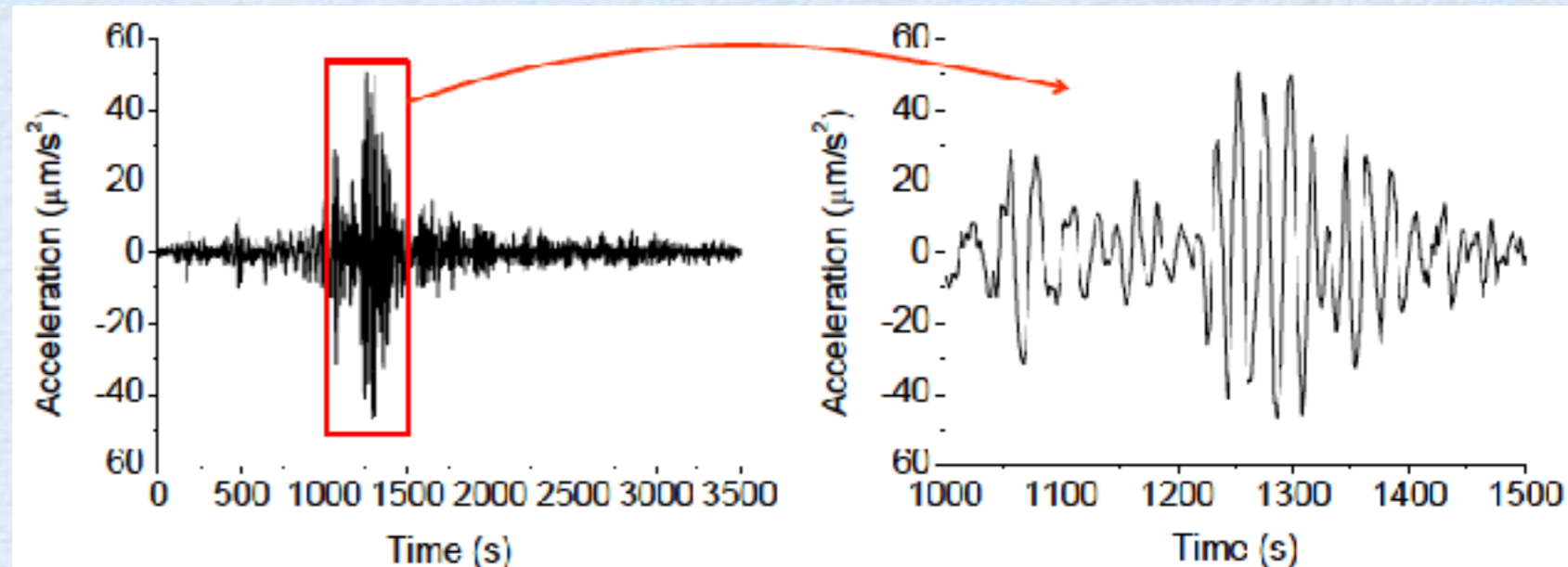
- AI phase is affected by residual atomic motion
- Velocity spread due to finite temperature
- Shot to shot jitter in initial position and mean velocity
- Atomic motion couples with
  - laser wavefront distribution
  - Coriolis acceleration & higher order inertial effects
  - gravitational and magnetic gradients
- Coupled constraints on atomic temperature, launching velocity, and wavefront jitter
- Limits are particularly stringent for very long baseline (space GW detectors)
- For ground applications, nK temperatures and  $\lambda/100$  wavefront errors would be ok



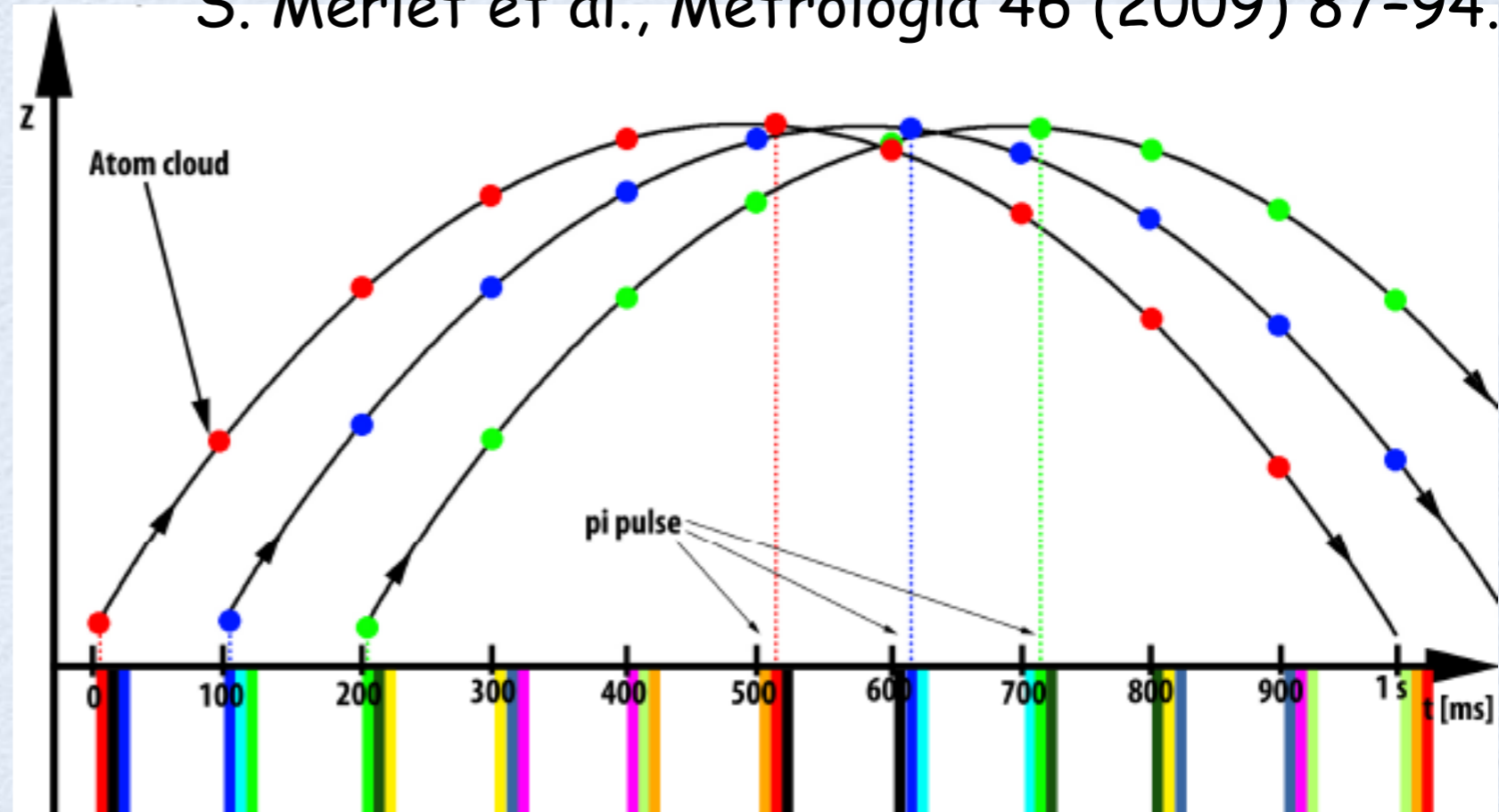
# Bandwidth & aliasing



- Sensitivity  $\sim T^2$
- Bandwidth  $\sim T^{-1}$
- Highly sensitive gravimeters detect earthquakes
  - $T \sim 50$  ms
  - rep rate  $\sim 5$  Hz
- Improving  $T$  without loosing bandwidth
  - interleaved interferometers
  - main limit: Doppler separation of different channels
  - see e.g. *I. Dutta et al., PRL 116, 183003 (2016)*



S. Merlet et al., Metrologia 46 (2009) 87-94.

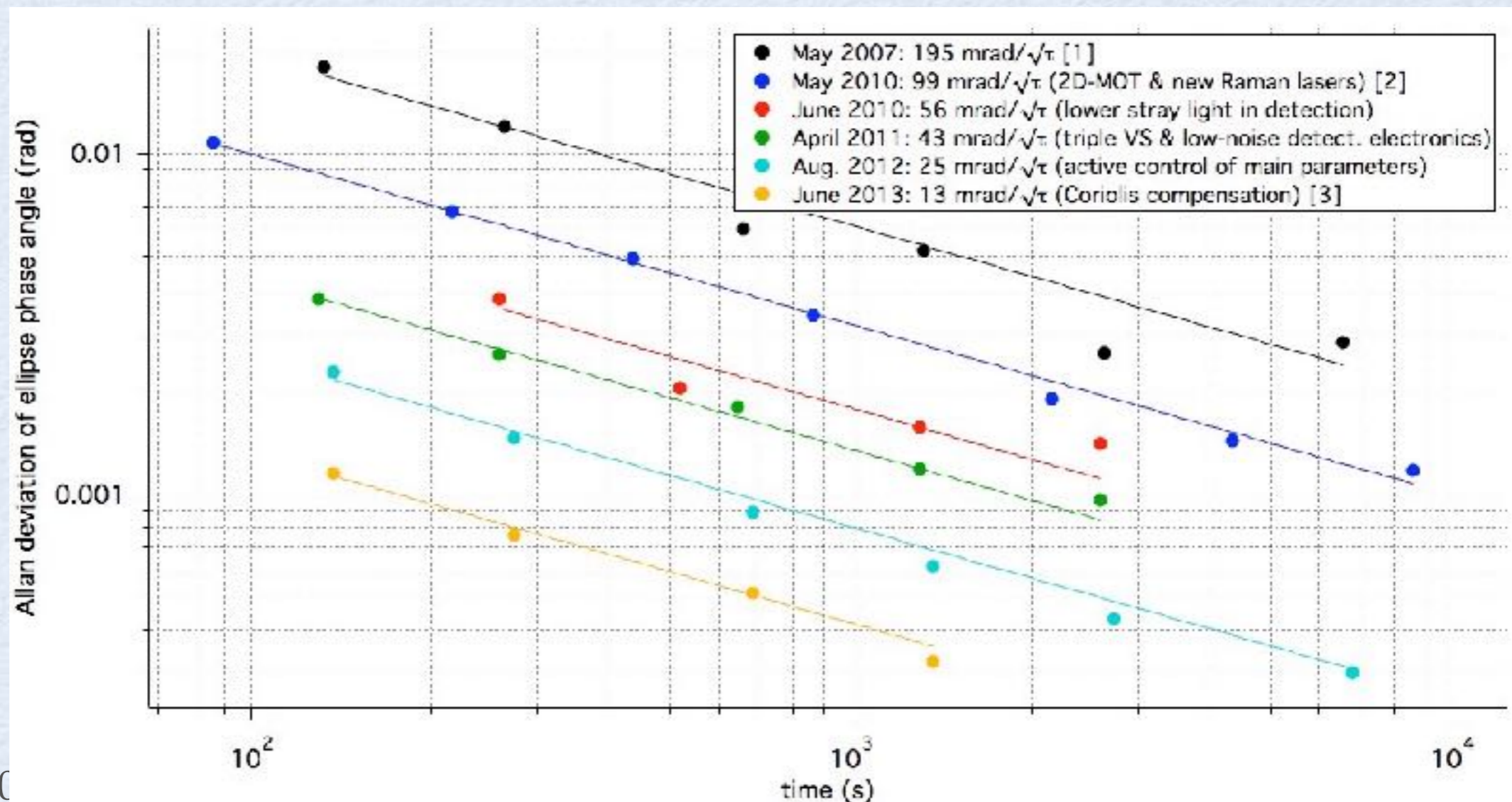


# Technical noise sources



- magnetic fields
- stray light
- photon shot noise
- tilt noise

- [1] G. Lamporesi et al., Phys. Rev. Lett 100, 050801 (2008)  
[2] F. Sorrentino et al., New J. Phys. 12, 095009 (2010)  
[3] F. Sorrentino et al., Phys. Rev. A 89, 023607 (2014)



# Potential improvements

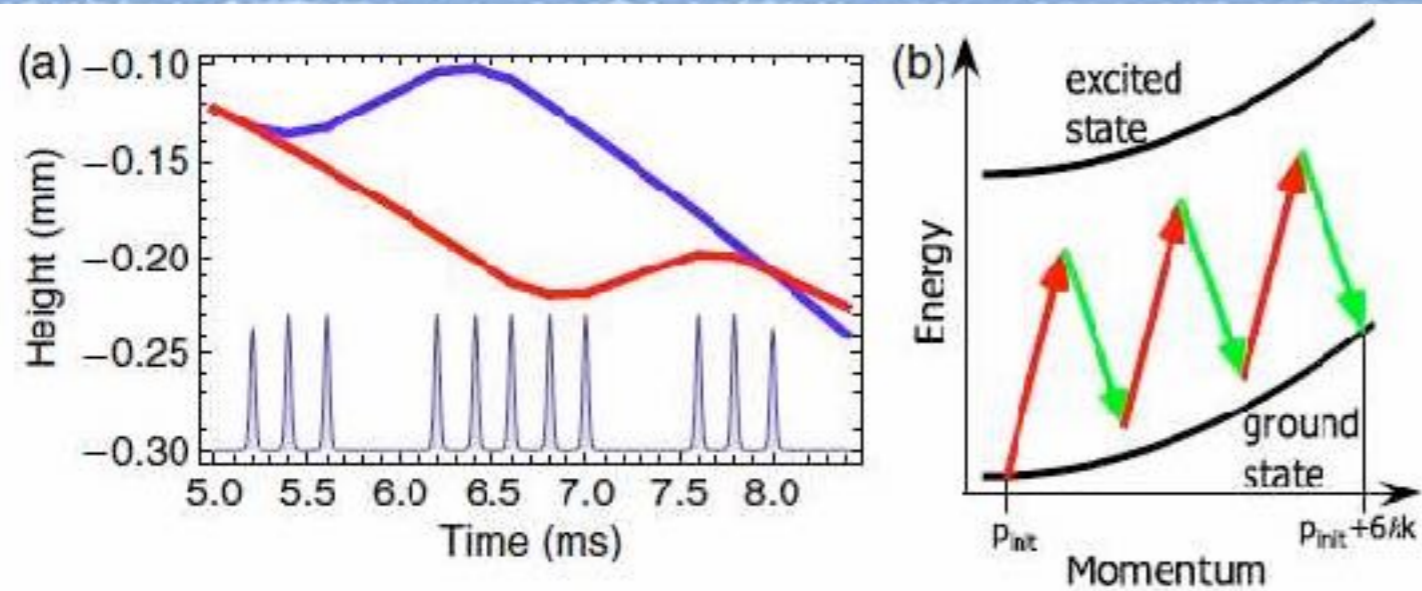


Single-shot, quantum noise limited gradiometer sensitivity

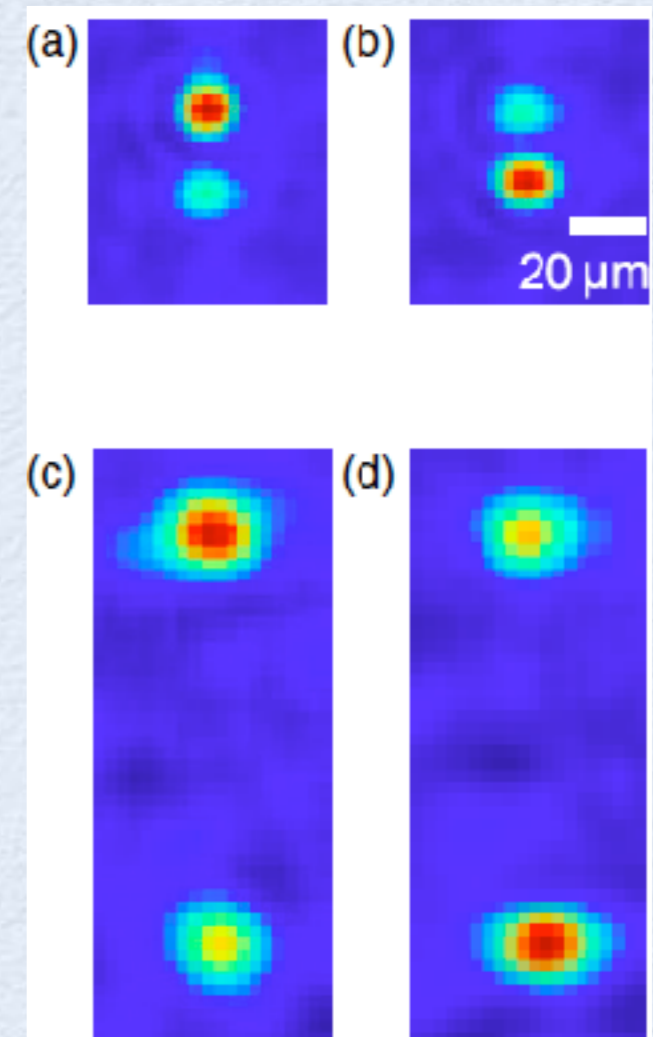
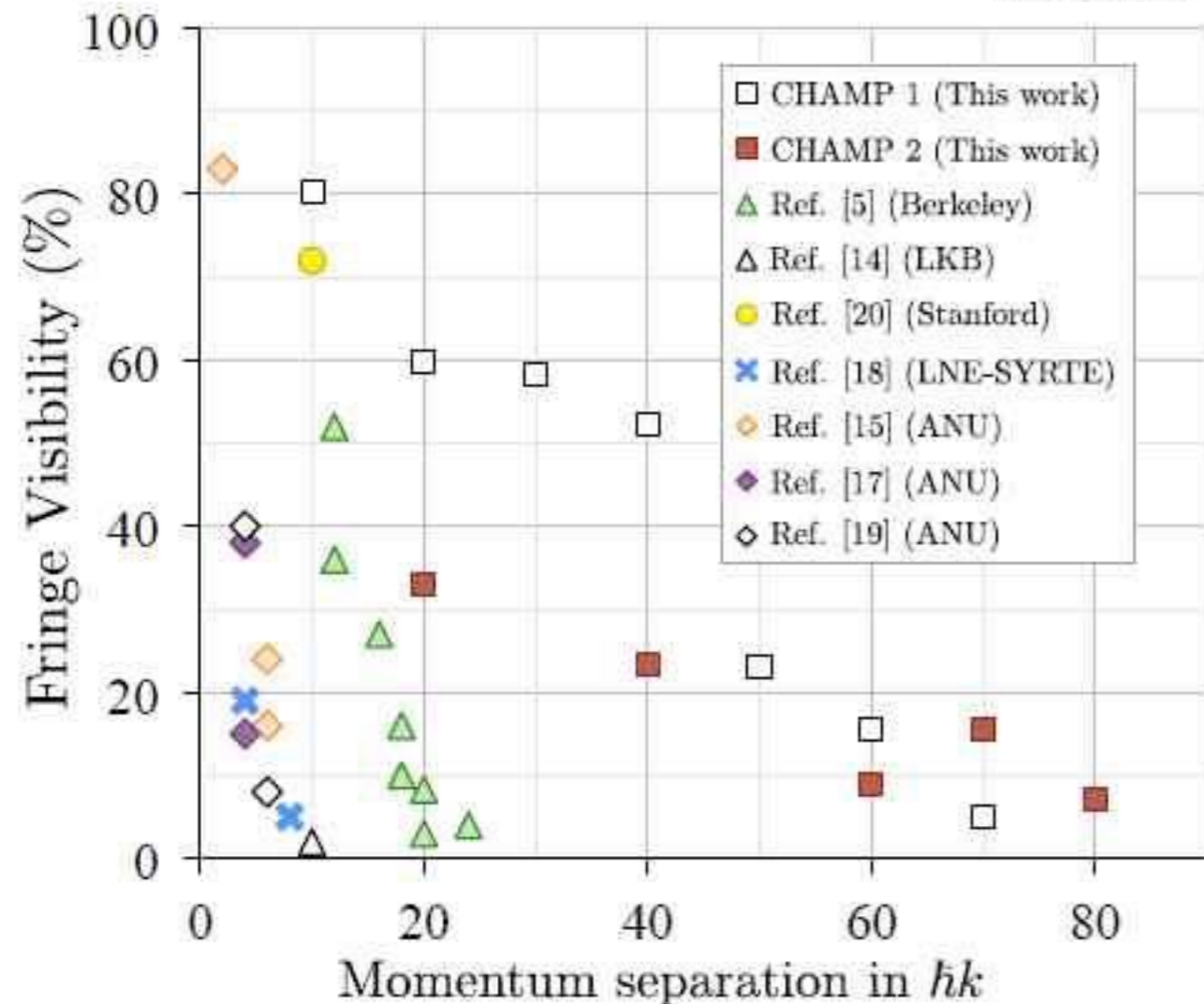
$$\delta \nabla g = \frac{\delta \phi}{nkLT^2}$$

- Increasing the scale factor
  - Large momentum transfer splitters (increase  $nk$ )
  - Large scale AI (increase  $T$ )
  - Long baseline (increase  $L$ )
- Reducing phase noise
  - High flux atomic sources (reduce quantum noise)
  - Squeezing (phase noise below standard quantum limit)
- Choice of atomic species
- Interferometer topology

# Large momentum transfer beam splitters



- H. Müller et al., PRL **102**, 240403 (2009)
- S.-W. Chiow et al., PRL **107**, 130403 (2011)
- G. D. McDonald et al., PRA **88**, 053620 (2013)

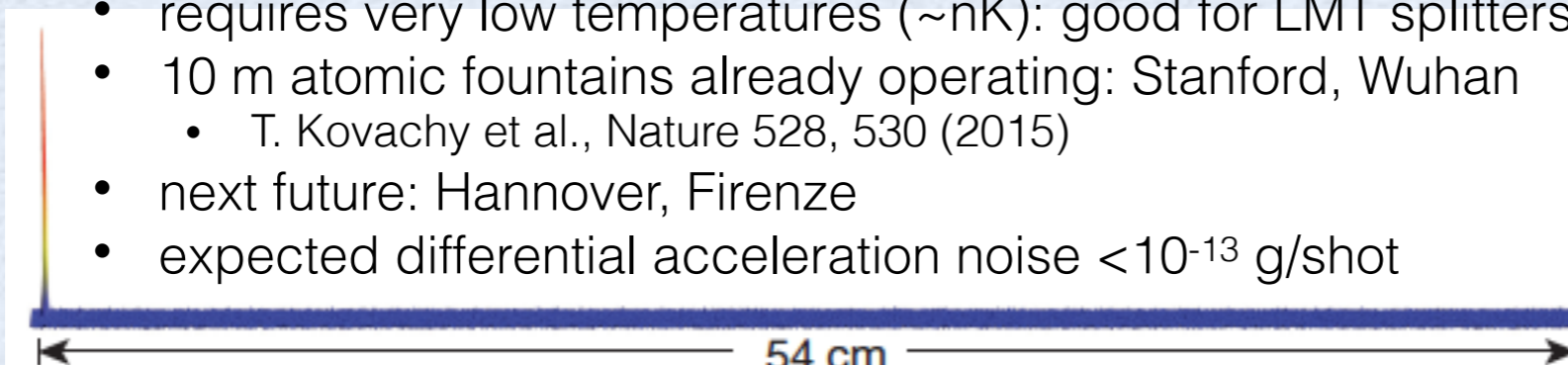
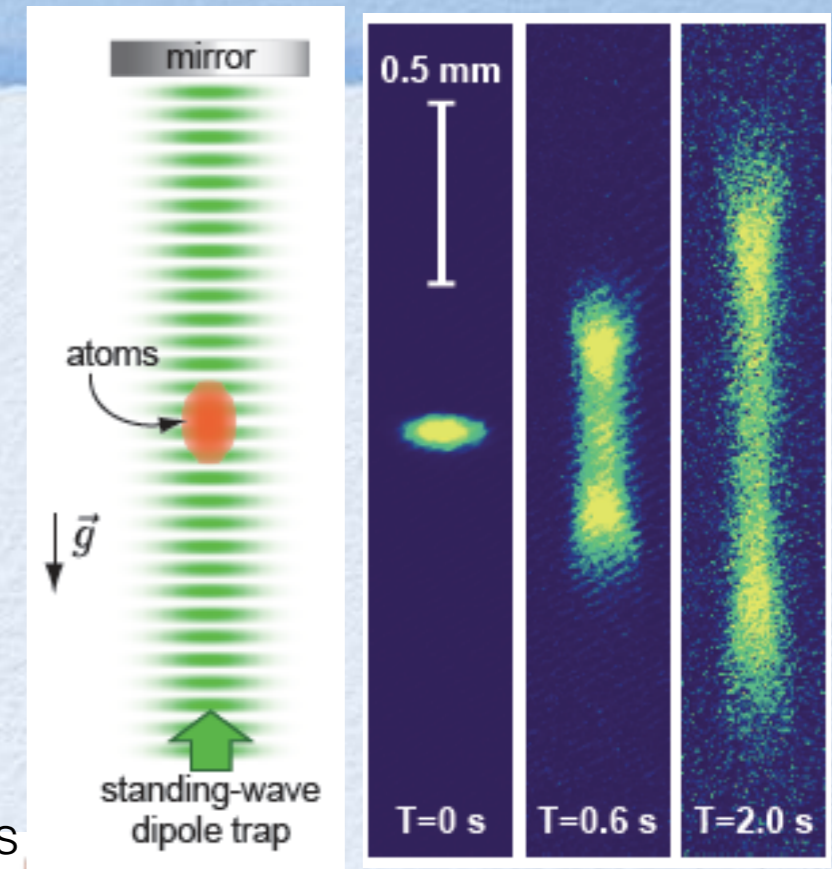


Atom interferometry gradiometers

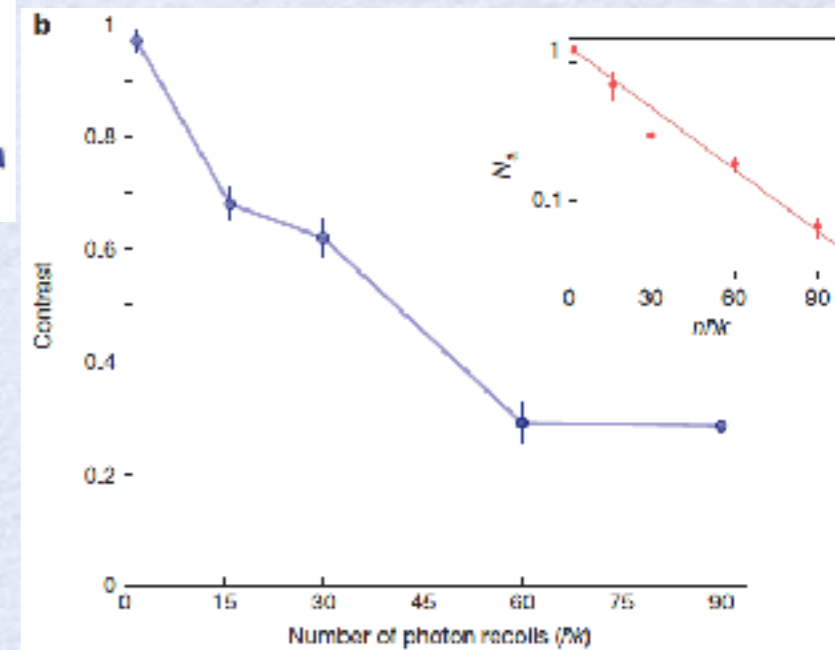
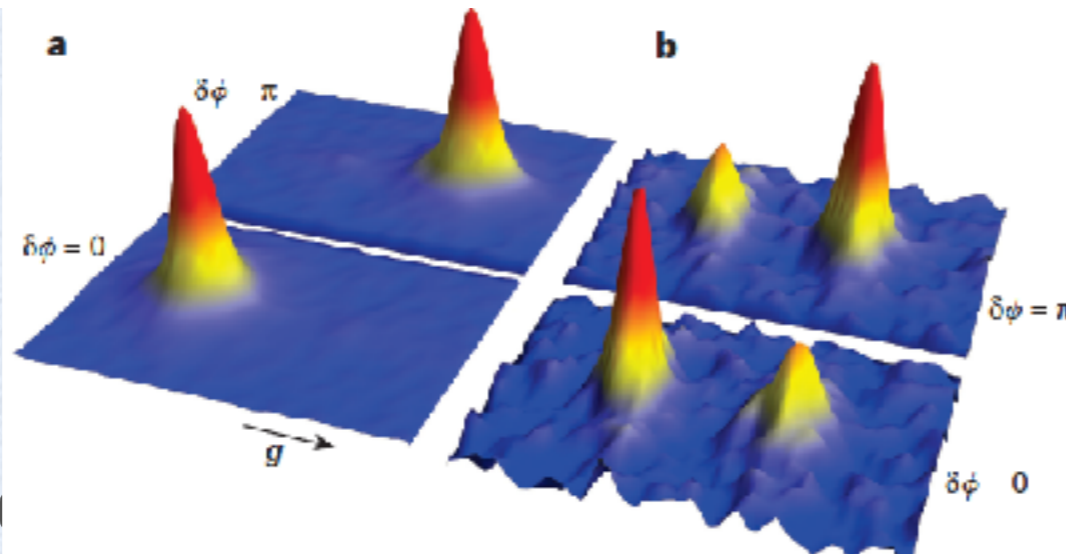
# Long interaction time



- Microgravity: in principle  $T > 10$  s
- Increasing  $T$  on Earth
  - trapped atoms (Sr)
    - up to 15000 coherent photon recoils
    - decoherence time  $> 500$  s
      - *G. Ferrari et al., PRL 97, 060402 (2006)*
      - *V. V. Ivanov et al., PRL 100, 043602 (2008)*
      - *F. Sorrentino et al., Phys. Rev. A 79, 013409 (2009)*
      - *M. Tarallo et al., PRA 86, 033615 (2012)*
  - free fall
    - requires large vertical size
    - requires very low temperatures ( $\sim$ nK): good for LMT splitters
    - 10 m atomic fountains already operating: Stanford, Wuhan
      - *T. Kovachy et al., Nature 528, 530 (2015)*
    - next future: Hannover, Firenze
    - expected differential acceleration noise  $< 10^{-13}$  g/shot



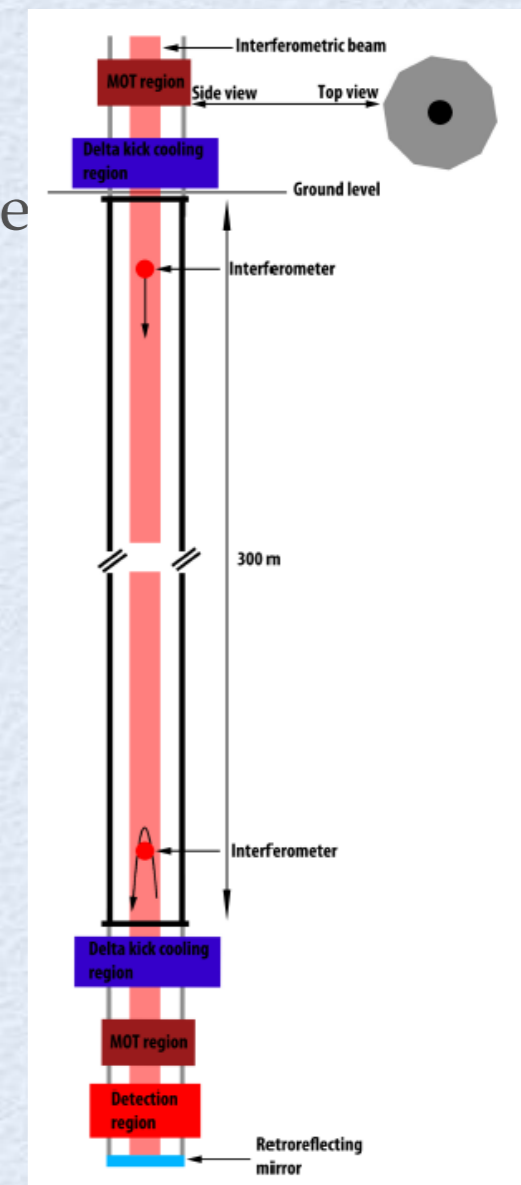
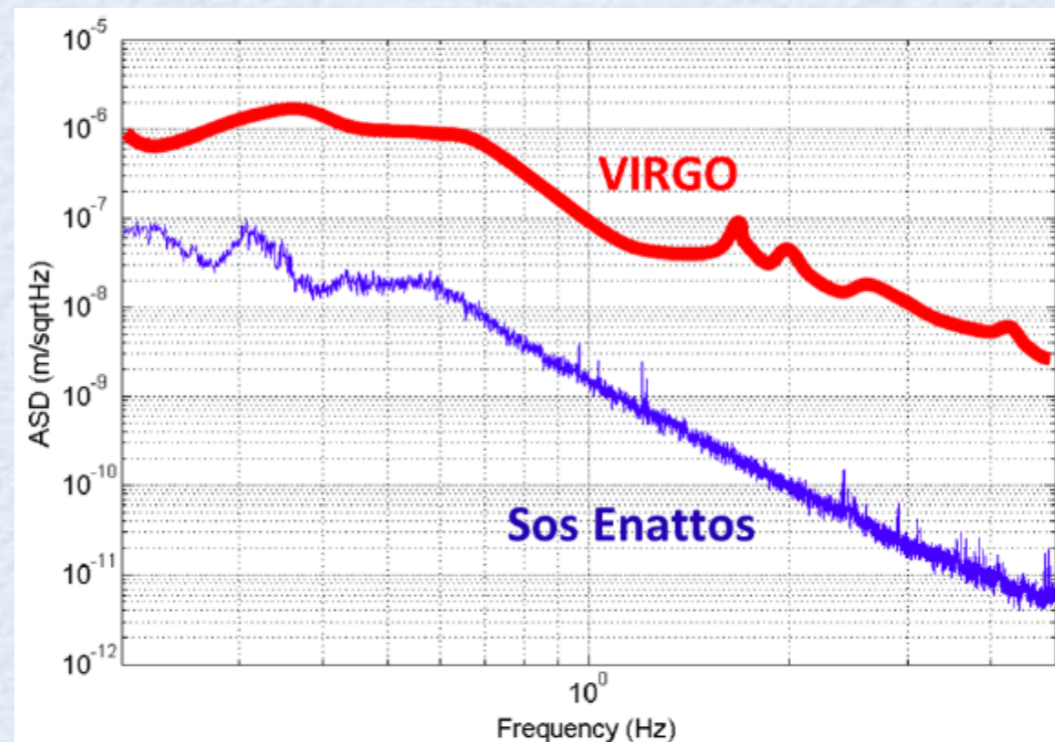
/01/20



Atom interferometry gradiometers

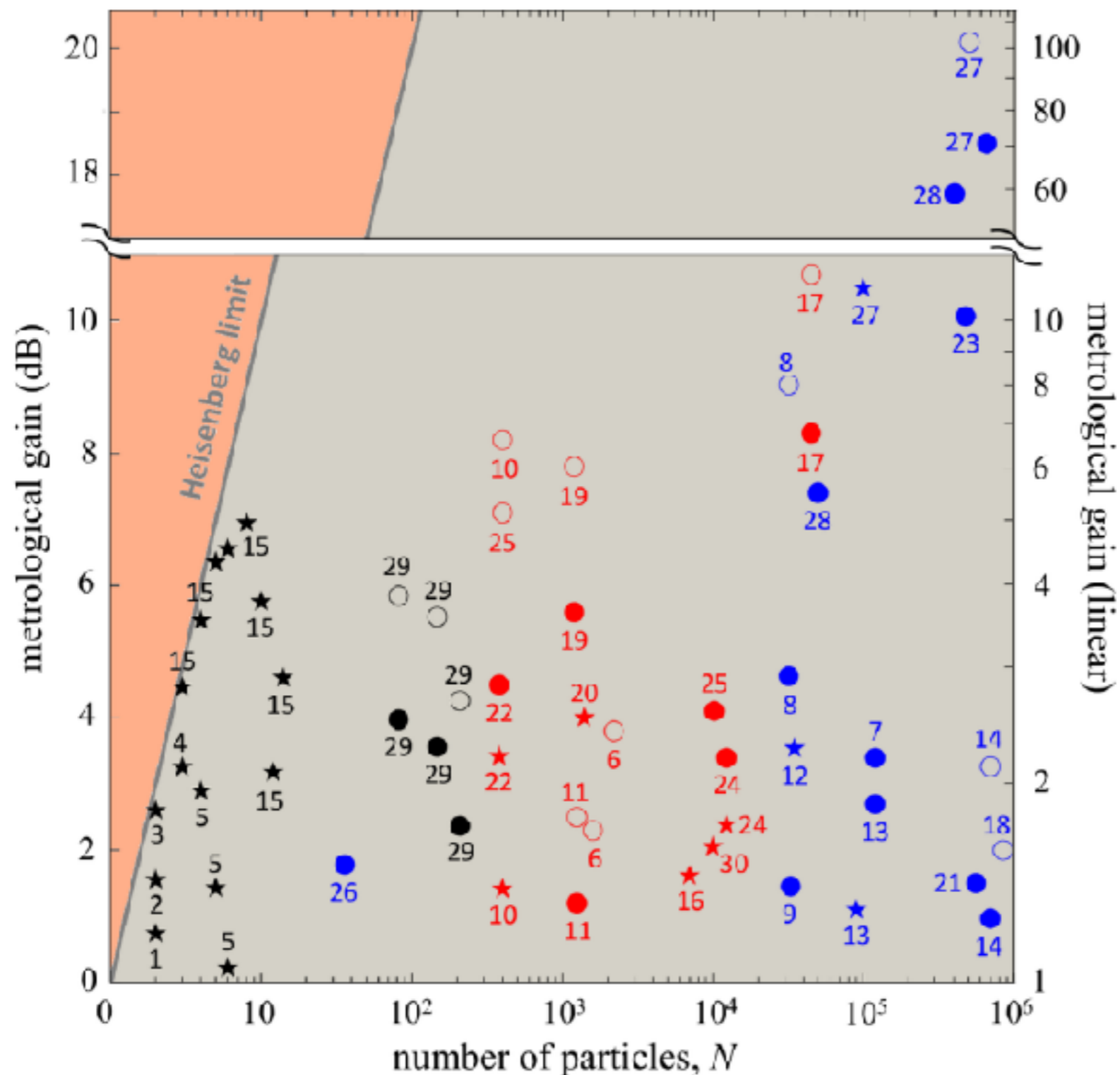
# Long baseline

- CMRR for seismic noise will be maintained when increasing the baseline  $L$
- proportional improvement on gradiometer sensitivity, different projects to reach  $L \sim 1$  km
  - long horizontal baseline  $L$  (MIGA, see talk by R. Geiger)
  - long vertical baseline: large  $L$  and large  $T$ 
    - MAGIA Adv: plan to install a differential atom interferometer in a decommissioned carbon mine in Sardinia
    - vertical shafts available up to 500 m
    - microseism in Sardinia is extremely low





# Squeezing



## TRAPPED IONS

- [1] Sackett *et al.*, 2000
- [2] Meyer *et al.*, 2001
- [3] Leibfried *et al.*, 2003b
- [4] Leibfried *et al.*, 2004
- [5] Leibfried *et al.*, 2005
- [15] Monz *et al.*, 2011
- [29] Bohnet *et al.*, 2016

## BOSE-EINSTEIN CONDENSATES

- [6] Estève *et al.*, 2008
- [10] Gross *et al.*, 2010
- [11] Riedel *et al.*, 2010
- [16] Lücke *et al.*, 2011
- [17] Hamley *et al.*, 2012
- [19] Berrada *et al.*, 2013
- [20] Ockeloen *et al.*, 2013
- [22] Strobel *et al.*, 2014
- [24] Muessel *et al.*, 2014
- [25] Muessel *et al.*, 2015
- [30] Kruse *et al.*, 2016

## COLD THERMAL ATOMS

- [7] Appel *et al.*, 2009
- [8] Leroux *et al.*, 2010a
- [9] Schleier-Smith *et al.*, 2010b
- [12] Leroux *et al.*, 2010b
- [13] Louchet-Chauvet *et al.*, 2010
- [14] Chen *et al.*, 2011
- [18] Sewell *et al.*, 2012
- [21] Sewell *et al.*, 2014
- [23] Bohnet *et al.*, 2014
- [26] Barontini *et al.*, 2015
- [27] Hosten *et al.*, 2016a
- [28] Cox *et al.*, 2016

L. Pezzé and A. Smerzi, arXiv:1609.01609v2

# Summary on AI gradiometer sensitivity



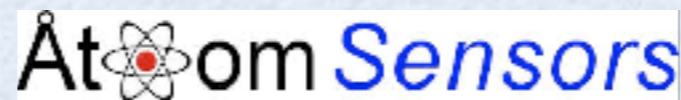
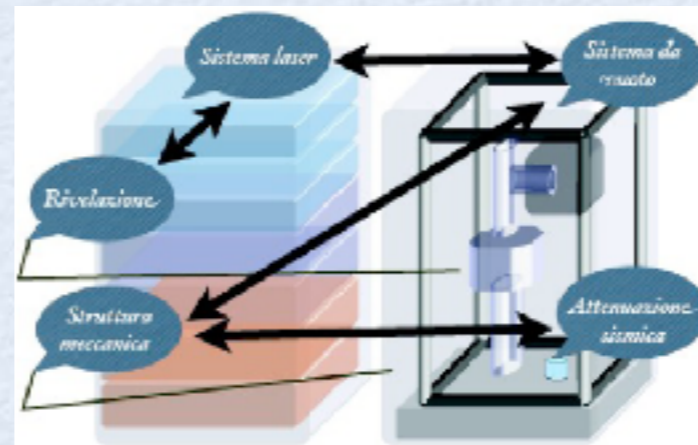
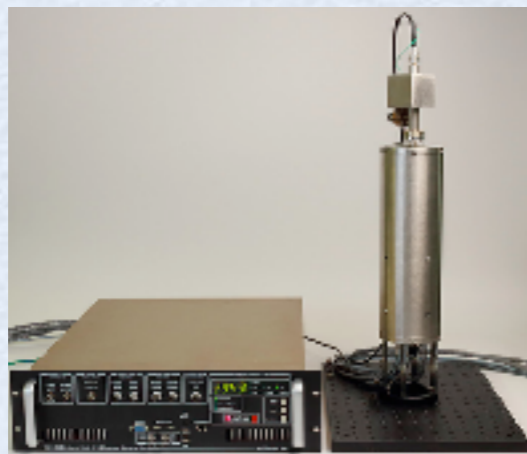
- Best demonstrated performance of a gravity gradiometer:
  - small scale instrument (MAGIA, 1 m fountain, 30 cm baseline):  $\sim 30 \text{ E @ } 20 \text{ s}$  [*F. Sorrentino et al., Phys. Rev. A 89, 023607 (2014)*]
  - medium scale instrument (Stanford 10 m tower, 30 cm baseline):  $\sim 3 \text{ E @ } 20 \text{ s}$  [*P. Asenbaum et al., Phys. Rev. Lett. 118, 183602 (2017)*]
- Differential acceleration sensitivity of a large vertical gradiometer
  - interaction time  $T$
  - $n$  photon recoils
  - $10^6$  atoms at QPN  $\rightarrow$  phase noise 1 mrad/shot
  - contrast  $C$
  - distance between atomic clouds  $L$
- Realistic projection for
  - medium baseline, long  $T$ :  $L=10 \text{ m}$ ,  $n=30$ ,  $T=1.5 \text{ s}$ ,  $C=0.6 \rightarrow 0.2 \text{ mE @ } 10 \text{ s}$
  - long baseline, long  $T$ :  $L=300 \text{ m}$ ,  $n=100$ ,  $T=3 \text{ s}$ ,  $C=0.6 \rightarrow 0.6 \mu\text{E @ } 10 \text{ s}$
- Substantial improvement expected with squeezing (long term plan)

# Technology readiness



- Commercial atomic gravimeters & gradiometers

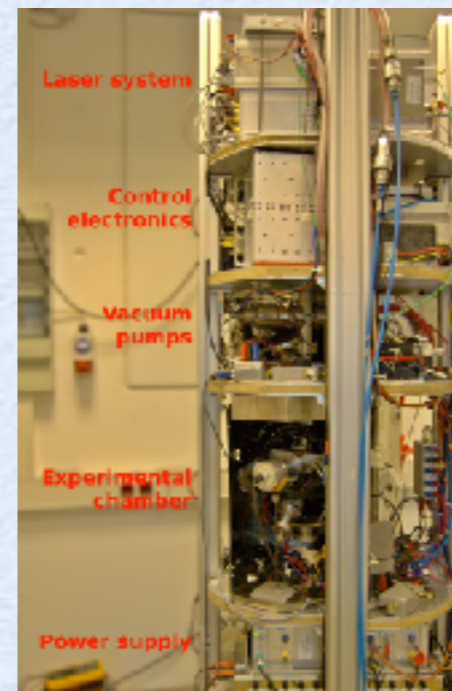
- MuQuans (France)
- AOSense (USA)
- AtomSensors (Italy)



- Prototypes for microgravity

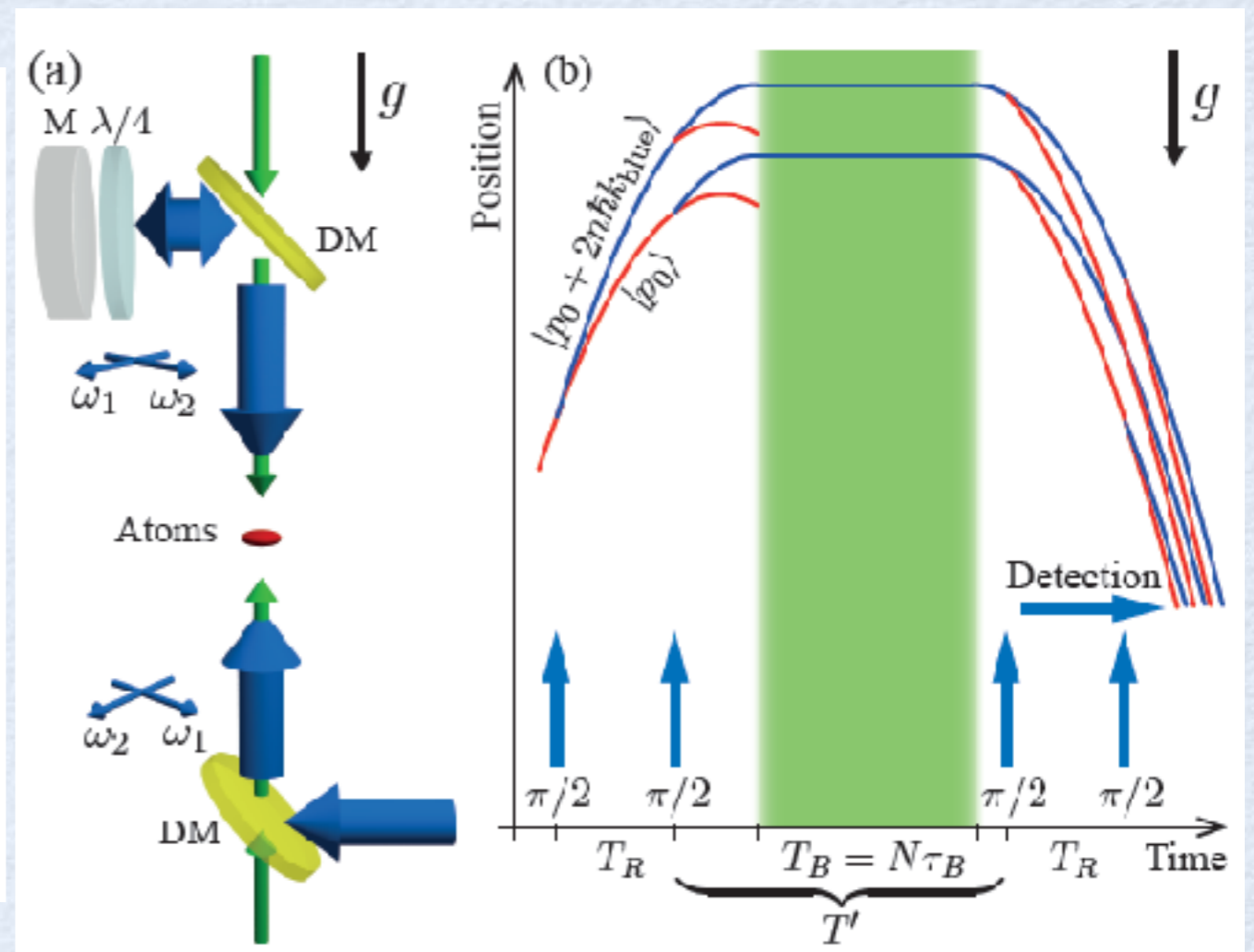
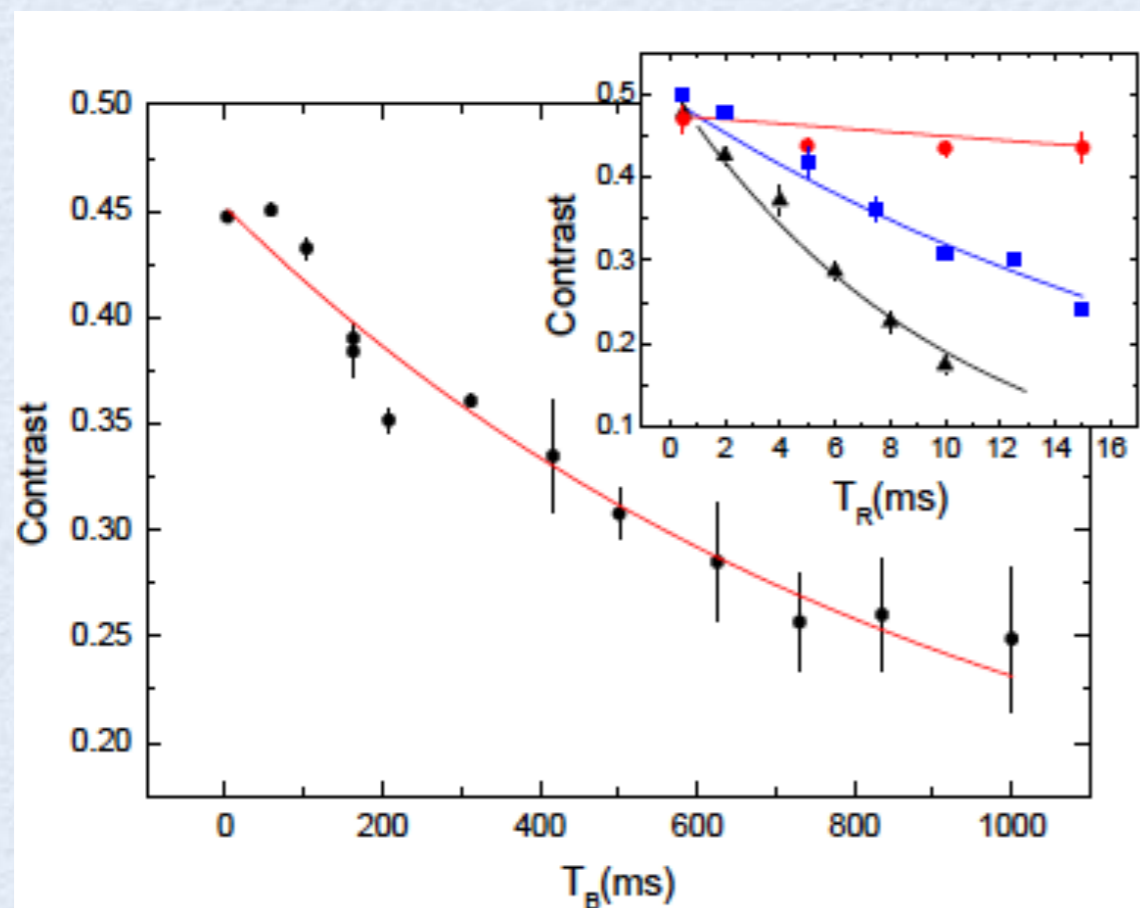
- I.C.E. (CNES)
- QUANTUS/MAIUS (DLR)
- S.A.I. (ESA)

G. Tino et al., Precision Gravity Tests with Atom Interferometry in Space, Nuclear Physics B 243, 203 (2013)



# Combination of trapping & free fall

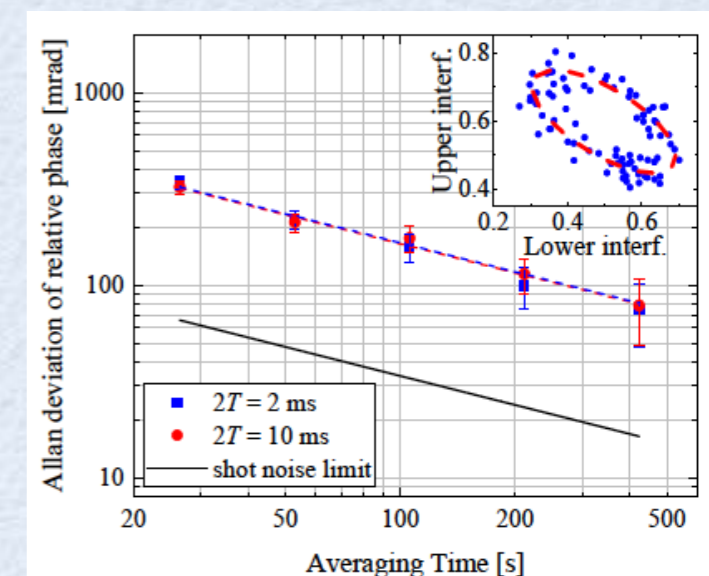
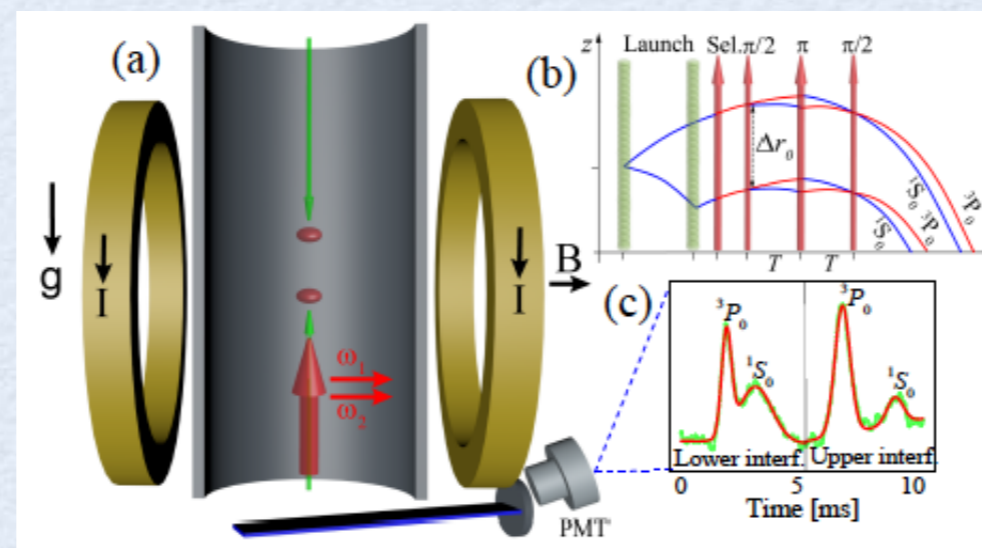
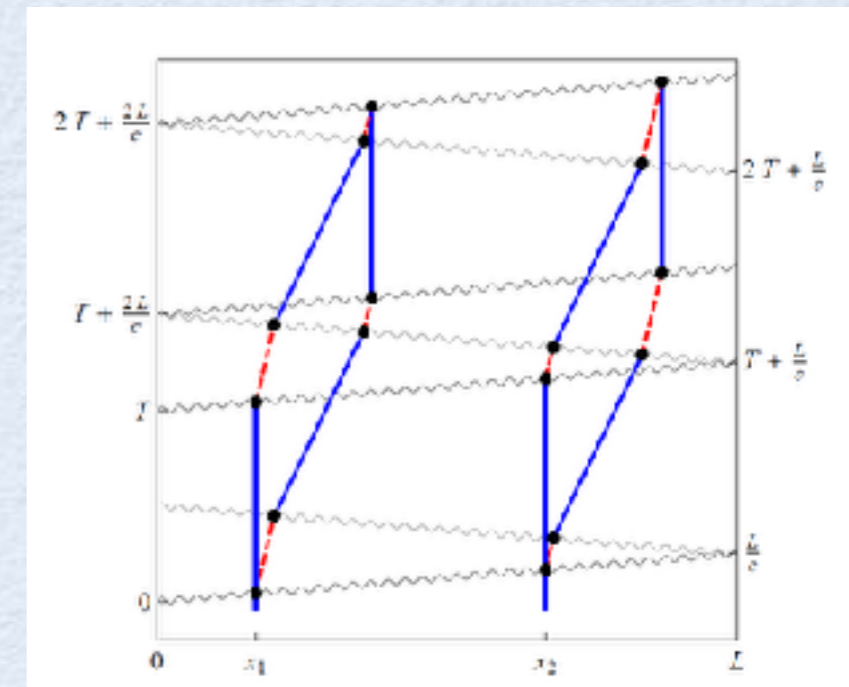
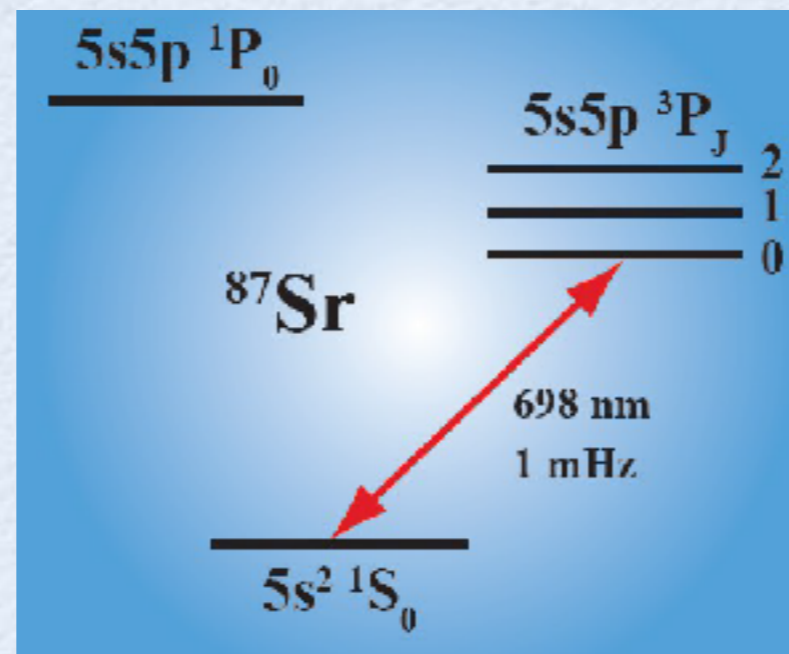
- LMT splitter + free fall to spatially separate interferometer arms
- Proof of principle demonstration with Rb
  - limited to  $T \sim 0.1$  s due to decoherence
  - R. Charrière et al., PRA 85, 013639 (2012)
- Up to  $T = 1$  s using Sr
  - X. Zhang et al., PRA 94, 043608-1 (2016)



# Single photon interferometry with Sr

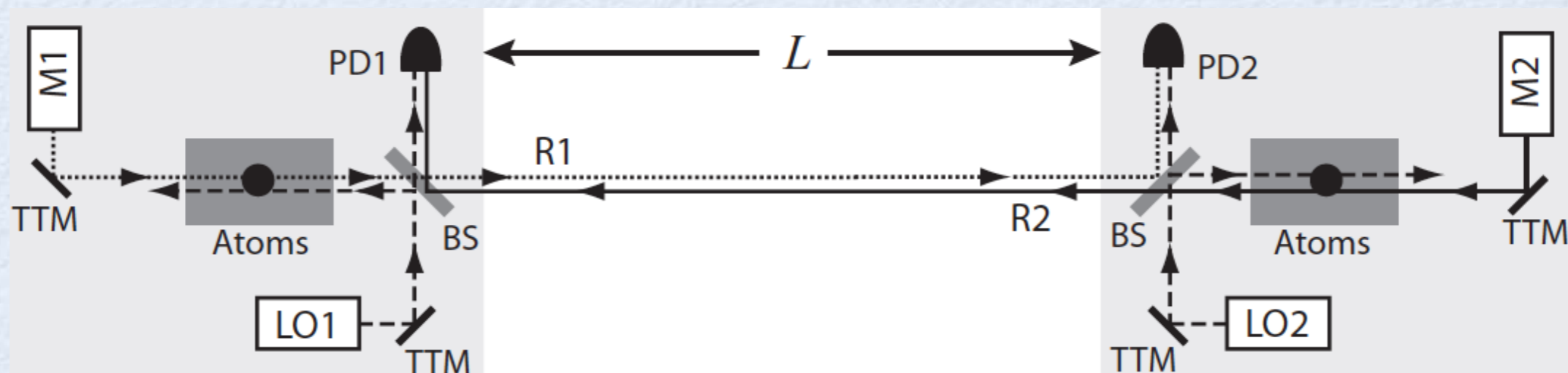


- Cancellation of frequency noise in long-baseline interferometry
- Proposed for GW detection in space
  - *P. W. Graham et al., PRL 110, 171102 (2013)*
- Experimentally demonstrated on small baseline
  - *L. Hu et al., PRL 119, 263601 (2017)*



# Fiber network

- Fiber Linked Atom-Interferometry Gravity Sensors (FLAGS)
  - *F. Sorrentino et al., to be published*
- gradiometer via coherent link between distant AI gravimeters
  - already proposed for GW detection with free-space links [*J. M. Hogan and M. A. Kasevich, Phys. Rev. A 94, 033632 (2016)*]
- two-way phase noise cancellation is routinely applied to optical frequency dissemination with residual phase noise at the level of  $\sim 0.1$  rad rms over  $\sim 100$  km
- trade off phase noise for much longer gradiometer arm
- in principle sub-mE/ $\sqrt{\text{Hz}}$  at reach with medium / small AI setups using state of the art technologies
- scalable to a network of many coherent AI gravity sensors



# Conclusions



- State of the art of AI gravity gradiometers
  - demonstrated sensitivity at  $\sim 10 \text{ E} / \sqrt{\text{Hz}}$  level
  - compact (mobile) sensors already available
- Further improvements expected in near future
  - long baseline
  - long T
  - squeezing
- realistic projection at
  - $\sim \text{mE} / \sqrt{\text{Hz}}$  level with mid-size sensors (10 m fountain)
  - $\sim \mu\text{E} / \sqrt{\text{Hz}}$  level with large scale antennas (1 km baseline)
- promising schemes for 10÷100 improvement over current sensitivity with small-scale sensors