

A clockwork model of flavor

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Outline

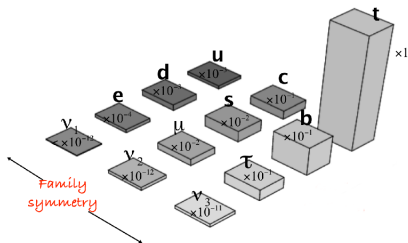
- 1 Flavor puzzle
- 2 The clockwork mechanism
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- 4 Phenomenology
 - The clockwork GIM mechanism
 - Low-energy bounds
 - Landau poles of α_s
- 5 Comparison with Froggatt-Nielsen
 - Collider signatures

The flavor puzzle of the SM

- The matter content of the SM comes in **3 almost identical generations**
 - Symmetry broken by Yukawa couplings

$$-\mathcal{L}_Y = \bar{q}_L Y_d d_R H + \bar{q}_L Y_u u_R \tilde{H} + \bar{\ell}_L Y_e e_R H + h.c.$$

- Masses** $M_u = L_u Y_u R_u^\dagger$, $M_d = L_d Y_d R_d^\dagger$ and **mixing**, $V_{CKM} = L_u L_d^\dagger$ are **hierarchical!**



CKM

$$|V| = \begin{matrix} & \begin{matrix} d & s & b \end{matrix} \\ \begin{matrix} u \\ c \\ t \end{matrix} & \begin{bmatrix} \text{orange} & \text{green} & \cdot \\ \text{green} & \text{orange} & \text{blue} \\ \cdot & \text{blue} & \text{orange} \end{bmatrix} \end{matrix}$$

The flavor puzzle of the SM

- The matter content of the SM comes in **3 almost identical generations**
 - ▶ Symmetry broken by Yukawa couplings

$$-\mathcal{L}_Y = \bar{q}_L Y_d d_R H + \bar{q}_L Y_u u_R \tilde{H} + \bar{\ell}_L Y_e e_R H + h.c.$$

- **Rotations** are “not physical” in the SM ,only **masses** and **V_{CKM}**


Where does flavor come from?

- ▶ **Geometrical:** *à la* Warped x-dimensions $\Lambda_{NP} \sim \Lambda_{\text{Natural}}$ [Randall&Sundrum'98](#)
- ▶ **Dynamical:** Horizontal symmetries $\Lambda_{NP} \gg \Lambda_{\text{Natural}}$ [Froggatt&Nielsen '78](#)
- ▶ ...

The clockwork mechanism

- Generic mechanism to produce exponential hierarchies

Choi & Im '15, Kaplan & Rattazzi '15, Giudice & McCullough '16



g $g_{\text{eff}} = g \times q^{-N}$

$\frac{1}{q} \times \frac{1}{q} \times \dots \times \frac{1}{q}$ N - times

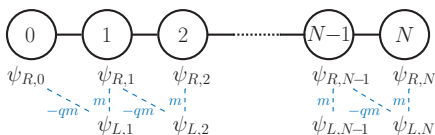
Kamenik's at PLANCK'17

- **Effective scale:** $\Lambda_{\text{eff}} = \Lambda/g_{\text{eff}}$

**Solve the EW hierarchy problem
by making $v_{\text{EW}} = M_P \times q^{-N}$**

Clockworking a single fermion

- **Chiral fermion** $\psi_R \implies$ Chain of “mirror” N vector-like fermions Giudice & McCullough'16



$$\mathcal{L}_{\psi_R} = i \sum_{j=0}^N \bar{\psi}_{R,j} \not{D} \psi_{R,j} + i \sum_{j=1}^N \bar{\psi}_{L,j} \not{D} \psi_{L,j} - m \sum_{j=1}^N (\bar{\psi}_{L,j} \psi_{R,j} - q \bar{\psi}_{L,j} \psi_{R,j-1}) + \text{h.c.},$$

- **Asymmetric** “off-site” interaction

One chiral fermion remains

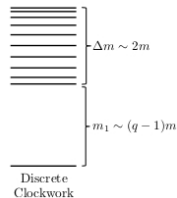
$N + (N + 1)$ χ symmetries $- 2 N$ χ -breaking parameters

- The N -vector-like fermions $m_k \sim m$: **Gears**

The clockwork physical basis

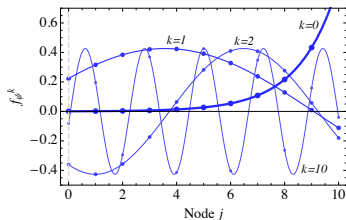
● Spectrum

- ▶ **Mass gap:** $M_1 \simeq m(q-1)$
 - ★ **Limit** $q \rightarrow 1 \implies$ A light gear
- ▶ **Spectrum:** Within a band $\sim 2m$
 - ★ **Limit** $q \gg 1 \implies m_k \sim qm$: **compressed**



McCullough&Giudice arXiv: 1610.07962

● Mixings



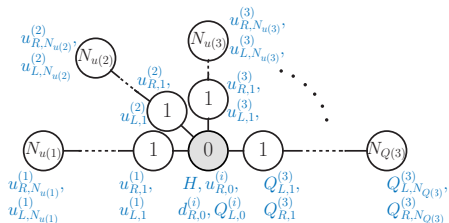
- ▶ **Overlaps** f_{ψ}^k : Mixings with the 0-th “node”

$$\psi'_{R,k} = \sum_{j=0}^N V_{jk}^R \psi_{R,j}, \quad f_{\psi}^k = V_{0k}^R$$

Clockwork mechanism

$$f_{\psi}^0 \sim 1/q^N \text{ for } q > 1$$

A clockwork model of flavor



Simple recipe

- 1 Each SM field ψ in a family i has a clockwork chain $(q_{\psi^{(i)}} , N^{\psi^{(i)}})$
- 2 Identify the massless fermions with the SM fields
- 3 Make the SM-Higgs doublet interact to all chains at the “0-node”

$$\mathcal{L}_{c'k} = \sum \left(\mathcal{L}_{u_R^{(i)}} + \mathcal{L}_{d_R^{(i)}} + \mathcal{L}_{Q_L^{(i)}} \right) - \sum \left[(\gamma_D)_{ij} \bar{Q}_{L,0}^{(i)} H d_{R,0}^{(j)} + (\gamma_U)_{ij} \bar{Q}_{L,0}^{(i)} \tilde{H} u_{R,0}^{(j)} + \text{h.c.} \right],$$

- 4 Adjust $(q_{\psi^{(i)}} , N^{\psi^{(i)}})$ to reproduce the exponential flavor hierarchies

The flavor clockwork in action

- Start with **anarchic** Yukawa matrices Y_U, Y_D

$$\left(Y_U^{\text{SM}}\right)_{ij} = f_{Q(i)} (Y_U)_{ij} f_{U(j)}, \quad \left(Y_D^{\text{SM}}\right)_{ij} = f_{Q(i)} (Y_D)_{ij} f_{d(j)}$$

- Zero-mode overlaps:** $f_{\psi(i)} \sim q_{\psi(i)}^{-N_{\psi(i)}}$

- This produces . . . [Froggatt & Nielsen '78](#)

- CKM mixings**

$$(V_{\text{CKM}})_{ij} = \frac{f_{Q(i)}}{f_{Q(j)}}, \quad i < j$$

Fix doublet overlaps

- Masses of SM quarks**

$$m_{u(i)} \sim v f_{Q(i)} f_{u(i)}, \quad m_{d(i)} \sim v f_{Q(i)} f_{d(i)}$$

Fix singlet overlaps

The flavor clockwork in action

- **Phenomenological scalings**

$$\begin{aligned} f_{Q(1)} &\sim \lambda^3, & f_{Q(2)} &\sim \lambda^2, & f_{Q(3)} &\sim 1, \\ f_{u(1)} &\sim \lambda^5, & f_{u(2)} &\sim \lambda^2, & f_{u(3)} &\sim 1, \\ f_{d(1)} &\sim \lambda^5, & f_{d(2)} &\sim \lambda^3, & f_{d(3)} &\sim \lambda^3. \end{aligned}$$

- **Illustrative benchmark for proto-Yukawas**

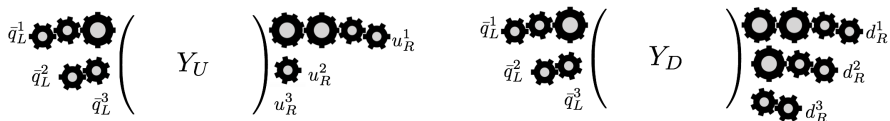
$$Y_U = \begin{pmatrix} 2.68 - 1.14 i & 1.15 + 0.34 i & -1.28 + 0.63 i \\ -0.35 - 1.00 i & 0.53 + 0.28 i & 0.28 + 1.14 i \\ -1.13 + 0.85 i & 0.19 + 0.31 i & 0.76 - 0.12 i \end{pmatrix},$$
$$Y_D = \begin{pmatrix} -1.76 + 0.57 i & 0.57 - 0.09 i & 0.30 + 0.10 i \\ 0.23 + 0.52 i & 0.04 - 0.91 i & 0.73 - 0.67 i \\ 0.23 + 0.43 i & -0.63 - 0.16 i & 0.04 - 0.97 i \end{pmatrix},$$

- ▶ **Lead to the right quark masses and CKM!**

Similar to solutions of flavor puzzle in Randall-Sundrum or Froggatt-Nielsen

Variations of the flavor-clockwork model

- “Randall-Sundrum” variation: **Same** $N_{\psi(i)} = N$ for all ψ and i , **different** $q_{\psi(i)}$
 - ▶ Large number of mirror fermions: $9 \times N$
 - ▶ Unsuppressed $Q(3)$ and $u(3) \implies q \sim 1 \implies$ **Light gears**
- “Froggatt-Nielsen” variation: **Same** $1/q = \lambda$ for all ψ and i , **different** $N_{\psi(i)}$
 - ▶ Number of mirror fermions can be minimized
 - ▶ $q \gg 1 \implies$ **Compressed gear spectrum** $m_k \simeq qm$
 - ▶ **Artist’s impression of the FN setup**



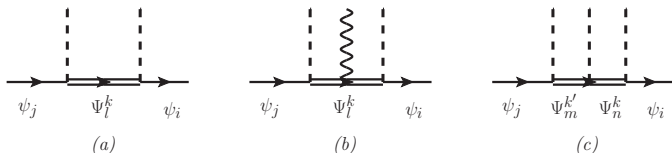
- ▶ This setup requires **5 doublet gears**, **7 u -type gears** and **11 d -type gears**

Low-energy consequences: The clockwork GIM mechanism

- **Gears** (VLQs) can have disruptive effects in low-energy flavor observables!

Bobeth *et al.* JHEP 1704 (2017) 079

- ▶ Contributions to FCNCs can push the **flavor-mass scale to PeV!**
- ▶ **Example:** Contributions to Higgs and weak boson couplings



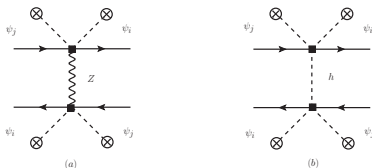
Anomalous $d_{L(i)} d_{L(j)}$ Z couplings: $[\delta g_L]_{ij}^{Z_d} = \frac{v^2}{4} f_{Q(i)} f_{Q(j)} [Y_D M_d^{-2} Y_D^\dagger]_{ij}$

Clockwork GIM mechanism

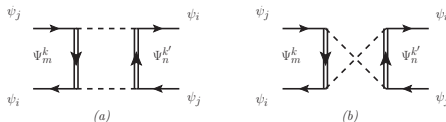
SM-gear mixing via Higgs \Rightarrow **Overlap** $\sim (f_{\psi(i)})^2$ suppression!

Low-energy consequences: Neutral-meson mixing

- **Tree-level** at $\mathcal{O}(v^4/m^4)$



- **Loop-level** at $\mathcal{O}(v^2/m^2)$



- **However** these contributions are very suppressed by overlaps
 - ▶ **For example**, $K-\bar{K}$ mixing the “leading” contribution is of the type

$$(f_{Q(1)} f_{Q(2)})^2 (\bar{s}_L \gamma^\mu d_L)^2 \sim \lambda^{10} (\bar{s}_L \gamma^\mu d_L)^2$$

- ★ Same parametric suppression (and structure) as the top-box in the SM!

Low-energy systematically: SMEFT

- Contributions of the gears in the SMEFT if $m \gg v_{EW}$ Buchmuller & Wyler'88, ...

$$\mathcal{L}_{\text{SMEFT}}^{(6)} = \sum w_i \mathcal{O}_i$$

$\mathcal{O}_{HQ}^{(1)}$	$(H^\dagger_i \overleftrightarrow{D}_\mu H) \bar{Q}_L \gamma^\mu Q_L$	$\mathcal{O}_{QQ}^{(1)}$	$(\bar{Q}_L \gamma^\mu Q_L)(\bar{Q}_L \gamma_\mu Q_L)$
$\mathcal{O}_{HQ}^{(3)}$	$(H^\dagger_i \overleftrightarrow{D}_\mu^I H) \bar{Q}_L \gamma^\mu \tau^I Q_L$	$\mathcal{O}_{QQ}^{(3)}$	$(\bar{Q}_L \gamma^\mu \tau^I Q_L)(\bar{Q}_L \gamma_\mu \tau^I Q_L)$
\mathcal{O}_{Hu}	$(H^\dagger_i \overleftrightarrow{D}_\mu H) \bar{u}_R \gamma^\mu u_R$	$\mathcal{O}_{Qd}^{(1)}$	$(\bar{Q}_L \gamma^\mu Q_L)(\bar{d}_R \gamma_\mu d_R)$
\mathcal{O}_{Hd}	$(H^\dagger_i \overleftrightarrow{D}_\mu H) \bar{d}_R \gamma^\mu d_R$	$\mathcal{O}_{Qu}^{(1)}$	$(\bar{Q}_L \gamma^\mu Q_L)(\bar{u}_R \gamma_\mu u_R)$
\mathcal{O}_{Hud}	$(\tilde{H}^\dagger_i D_\mu H) \bar{u}_R \gamma^\mu d_R$	\mathcal{O}_{dd}	$(\bar{d}_R \gamma^\mu d_R)(\bar{d}_R \gamma_\mu d_R)$
\mathcal{O}_{uH}	$(H^\dagger H) \bar{Q}_L \tilde{H} u_R$	\mathcal{O}_{uu}	$(\bar{u}_R \gamma^\mu u_R)(\bar{u}_R \gamma_\mu u_R)$
\mathcal{O}_{dH}	$(H^\dagger H) \bar{Q}_L H d_R$		

- Match neutral-meson mixing to one-loop accuracy at $\mathcal{O}(v^2/m^2)$

- ▶ $\psi^2 H^2 D \rightarrow$ tree-level matching
- ▶ ψ^4 operators \rightarrow one loop matching
- ▶ EW RGE of $\psi^2 H^2 D$ at one-loop
- ▶ One-loop matching between SMEFT and Low-Energy-EFT

Low-energy observables

- **EWPO:** Weak boson LH and RH couplings to the quarks Falkowski, Gonzalez-Alonso, Mimouni:17

$$M_d \gtrsim 5 \text{ TeV}, \quad M_u \gtrsim 0.5 \text{ TeV}, \quad M_Q \gtrsim 0.1 \text{ TeV}$$

- ▶ One of strongest bounds from $Z \rightarrow b_L b_L$: **It is not suppressed by overlaps!**

- **Rare decays:** $\Delta S = 1, \Delta B = 1, \Delta C = 1, \dots$

$$M_d \gtrsim 5 \text{ TeV}, \quad M_u \gtrsim 0.1 \text{ TeV}$$

- **Neutral-meson mixing:** $K-\bar{K}, B-\bar{B}, D-\bar{D}, \dots$ M. Bona et al.:16, UTfit collaboration

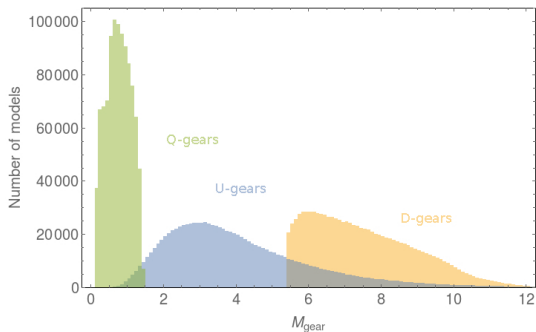
- ▶ The contributions of the gears is complicated ...

Process	U	D	Q	UD	UQ	DQ
$B_s-\bar{B}_s$	λ^4, \square^*	λ^4, \times and \square				
$B_d-\bar{B}_d$	λ^6, \square^*	λ^6, \times and \square				
$K-\bar{K}$	λ^{10}, \square^*	λ^{10}, \times and \square	$\lambda^{13}, \square^{*\dagger}$		$\lambda^{13}, \square^\dagger$	$\lambda^{12}, \times^\dagger$
$D-\bar{D}$	λ^{10}, \times and \square	λ^{10}, \square			$\lambda^{10}, \times^\dagger$	

$\dagger = \chi$ -enhancement; $*$ =logarithmic enhancement

Bounds scanning over models

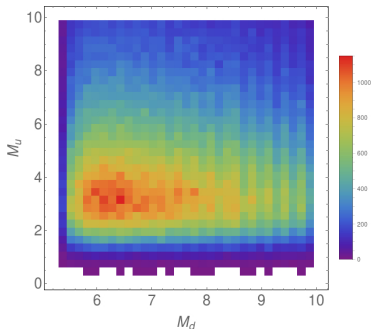
- **Strategy:** Scan over **anarchic** proto-Yukawas leading to suitable $Y_{U,D}^{\text{SM}}$
 - ▶ For each model, derive the stringest bounds on M_d, M_u, M_Q from all observables
- **Scan 1:** Assume only 1 type of gear active (decouple the other two)



- ▶ **U- and D-gears can be very light** $M_{u,d} \lesssim 10$ TeV
- ▶ **Q-gears unconstrained!**

Bounds scanning over models

- **Strategy:** Scan over anarchic proto-Yukawas leading to suitable $Y_{U,D}^{\text{SM}}$
 - ▶ For each model, derive the stringest bounds on M_d , M_u , M_Q from all observables
- **Scan 2:** Fix $M_Q = 2$ TeV and put bounds on M_u and M_d
 - ▶ **Rationale:** Dangerous contributions to meson mixing from QU and QD diagrams



Most of our clockwork models can live at the TeV scale!

Bounds scanning over models

- **Strategy:** Scan over **anarchic** proto-Yukawas leading to suitable $Y_{U,D}^{\text{SM}}$
 - ▶ For each model, derive the stringest bounds on M_d, M_u, M_Q from all observables

- **Remember our benchmark model...**

$$Y_U = \begin{pmatrix} 2.68 - 1.14 i & 1.15 + 0.34 i & -1.28 + 0.63 i \\ -0.35 - 1.00 i & 0.53 + 0.28 i & 0.28 + 1.14 i \\ -1.13 + 0.85 i & 0.19 + 0.31 i & 0.76 - 0.12 i \end{pmatrix},$$
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- We obtain ...

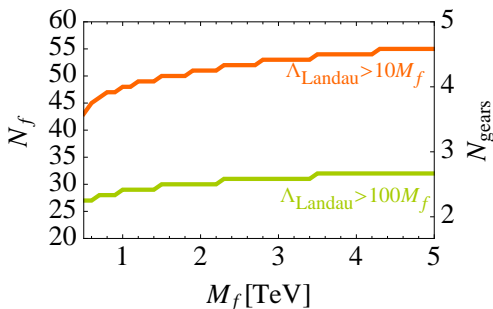
$$M_u \gtrsim 3.8 \text{ TeV}, \quad M_u \gtrsim 5.9 \text{ TeV}, \quad (M_Q = 2 \text{ TeV})$$

Bounds from Landau poles in α_s

- New colored states change the β -function of QCD

$$\frac{d\alpha_s}{d\ln\mu} = -2\beta_0 \frac{\alpha_s^2}{4\pi}, \quad \beta_0 = \frac{11N_c - 2N_f}{3},$$

- ▶ For $N_f \gtrsim 16$ QCD develops a UV Landau pole!
- Demanding $\Lambda_{\text{QCD}} > 100 \times M_{\text{gear}}$ we set a bound on N_f



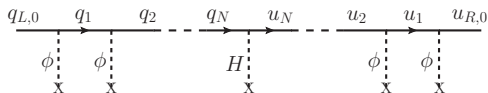
- In our benchmark $N_f = 34 \implies \Lambda_{\text{QCD}} \sim 100 \text{ TeV}$

Comparison with Froggatt-Nielsen mechanism

- Clockwork model of flavor looks similar to **Froggatt-Nielsen**

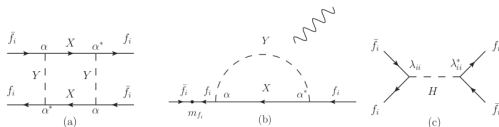
- ▶ **Horizontal symmetry** $\mathcal{H} (U(1))$

- ★ Different charges to fields in different generations
- ★ Add mirror VLQs and scalars (**flavon**) charged under \mathcal{H}



$$\mathcal{L}_Y = (Y_D)_{ij} \left(\frac{\langle \phi \rangle}{M} \right)^{H(\bar{Q}_i) + H(d_j)} \bar{Q}_i H d_j + (Y_U)_{ij} \left(\frac{\langle \phi \rangle}{M} \right)^{H(\bar{Q}_i) + H(u_j)} \bar{Q}_i \tilde{H} u_j$$

- Hierarchies in **small parameter** $\langle \phi \rangle / M$ given **dynamically** by the charges!
- **Richer phenomenology** \implies Flavons, gauged horizontal bosons, etc

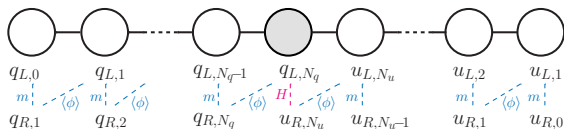


- ▶ **Bounds typically in the order of 50 TeV - 1 PeV** Calibbi, Lalak, Pokorski, Ziegler arXiv: 1204.1275

Pointing to a UV completion of the clockwork flavor model

- Clockwork mechanism is nondynamical as is
 - ▶ **Where does the q spurion comes from?**
- It can be UV-completed with an \mathcal{H} à la FN!

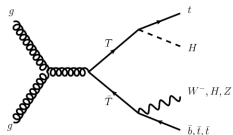
$$q M = \langle \phi \rangle \gg M$$



- **“Inverted” Froggatt-Nielsen mechanism**
- Does it help reducing flavon contributions?

Collider signatures from the clockwork gears

- **Pair-produce** lightest gears of Q -type at the **LHC** via gg fusion

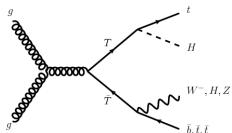


- ▶ Heaviest gears decay down the clockwork-chain
- ▶ Eventually they decay into t_L , b_L and h , W and Z

- **LHC** searches of VLQs set bounds typically of $\mathcal{O}(1)$ TeV
- Need more sophisticated searches ...

Collider signatures from the clockwork gears

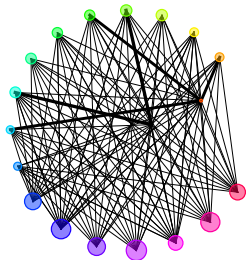
- **Pair-produce** lightest gears of Q -type at the **LHC** via gg fusion



- ▶ Heaviest gears decay down the clockwork-chain
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- **LHC** searches of VLQs set bounds typically of $\mathcal{O}(1)$ TeV
- Need more sophisticated searches ...

- **Psychodelic collider pheno (in progress) ...**



▶ Decay patterns of gears

- ★ **Central node:** SM fermion
 - ★ **Colored bulbs:** Gears classified by mass (order) and width (size)
 - ★ **Lines:** Decays with thickness representing BR
- ▶ **Simulating this in MG for promising signals!**

Conclusions

- **Clockwork mechanism:** Solution to the EW hierarchy problem by reducing M_P
- **Can we solve the flavor puzzle within this paradigm with low mass scale?**

YES! ...

- ... If we forget about “*Where does the clockwork comes from?*”
- **If not**, we are in a **Froggatt-Nielsen** setup with $\langle \phi \rangle \gg M$...