Closing talk: Flavor anomalies, model independent and new physics considerations

From flavor to new physics.

Lyon, France
Amarjit Soni, BNL-HET
04/20/2018



Clarification

- Closing talkNOT A SUMMARY
- CLOSING TALK ~ LAST TALK of this very nice meeting in a really nice flavory town on flavor physics
- Of course, in the past several years, I have been giving considerable attention to these anomalies.....So will make some observations based on presentations here and elsewhere and come up with important take home messages from my perspective
- Bound to be subjective.....Sorry

Outline

- several looming deviations from SM ...i.e.
 "anomalies"
- For each case:
- briefly mention reservations for expt & for theory/comments
- Model independent collider implications
- Assuming NP is a source: array of BSMs
- An illustration, of one minimal, well motivated setup for a BSM origin
- Summary, Outlook & emphasize 4 taking home

SET OF KEY ANOMALIES

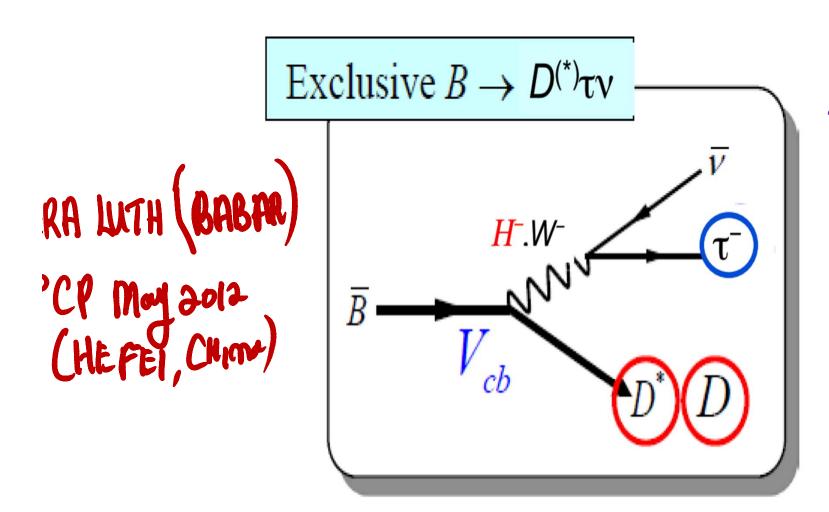
Anomalies galore!

- RD(*) ~ 46(?)
 RK(*): 2.66(A_K);
- g-2...BNL'06 => FNAL expty 3.66 myn lattie progress y
- E': a personal obsession....for a long^3 time=>'cause of the strong conviction that it is super-sensitive to NP/EVER

216[PRL 2015] => ~1200 now => ~1400

[2.1 σ (2.9 σ Buras; Nierste) => ??]few more months to new Inesultang BSM scenarios it is important to keep all these [INCLUDING E] + Higgs nadiative stability in mind

RD(*)



MANUEL FRANCE SEVILLA PLD Thesis

Independent of Vcb!

To test the SM Prediction, we measure

$$R(D) = \frac{\Gamma(\overline{B} \to D\tau\nu)}{\Gamma(\overline{B} \to D\ell\nu)} \qquad R(D^*) = \frac{\Gamma(\overline{B} \to D^*\tau\nu)}{\Gamma(\overline{B} \to D^*\ell\nu)}$$

Leptonic τ decays only

Several experimental and theoretical uncertainties cancel in the ratio!

DD avanta and fully na sametovicted.

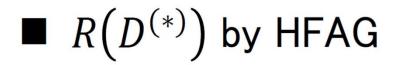
Improving constraints on $\tan \beta/m_H$ using $B \rightarrow D \tau \overline{\nu}$

Ken Kiers* and Amarjit Soni[†]

Department of Physics, Brookhaven National Laboratory, Upton, New York 11973-5000
(Received 12 June 1997)

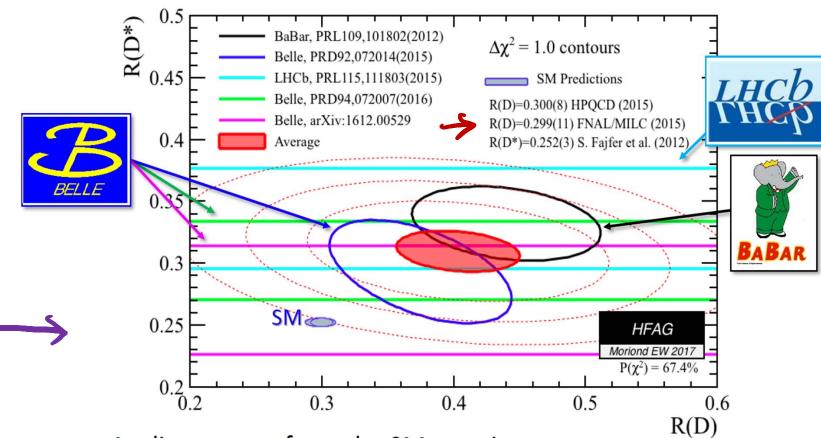
We study the q^2 dependence of the exclusive decay mode $B \to D \tau \overline{\nu}$ in type-II two Higgs doublet models (2HDM's) and show that this mode may be used to put stringent bounds on $\tan \beta/m_H$. There are currently rather large theoretical uncertainties in the q^2 distribution, but these may be significantly reduced by future measurements of the analogous distribution for $B \to D(e,\mu)\overline{\nu}$. We estimate that this reduction in the theoretical uncertainties would eventually (i.e., with sufficient data) allow one to push the upper bound on $\tan \beta/m_H$ down to about 0.06 GeV⁻¹. This would represent an improvement on the current bound by about a factor of 7. We

=) Follower my Vienste et ali fajfen et al 12



Hirose [BELLE]@EW MORIOND Mar. 2017





- \sim 4 σ discrepancy from the SM remains
 - All the experiments show the larger $Rig(D^{(*)}ig)$ than the SM
- More precise measurements at Belle II and LHCb are essential

Lepton universality tests

In the SM, ratios

LHCb introduced such v ule



$$R_{\rm K} = \frac{\int d\Gamma[B^+ \to K^+ \mu^+ \mu^-]/dq^2 \cdot dq^2}{\int d\Gamma[B^+ \to K^+ e^+ e^-]/dq^2 \cdot dq^2}$$

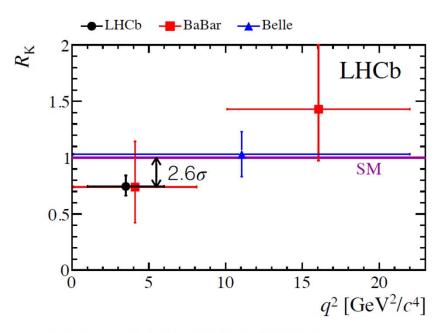


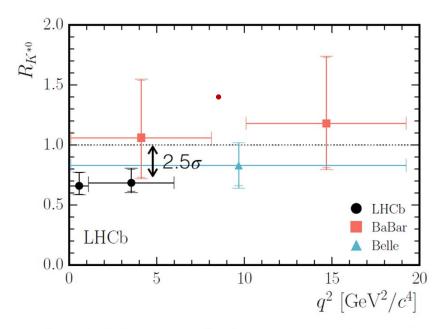
only differ from unity by phase space — the dominant SM processes couple equally to the different lepton flavours.

- Theoretically clean since hadronic uncertainties cancel in the ratio.
- Experimentally challenging due to differences in muon/electron reconstruction (in particular Bremsstrahlung from the electrons).
 - → Take double ratios with $B \rightarrow J/\psi X$ decays to cancel possible sources of systematic uncertainty.
 - → Correct for migration of events in q² due to FSR/Bremsstrahlung using MC (with PHOTOS).

Lepton universality tests

 We have interesting hints of non-universal lepton couplings in LHCb run 1 dataset:





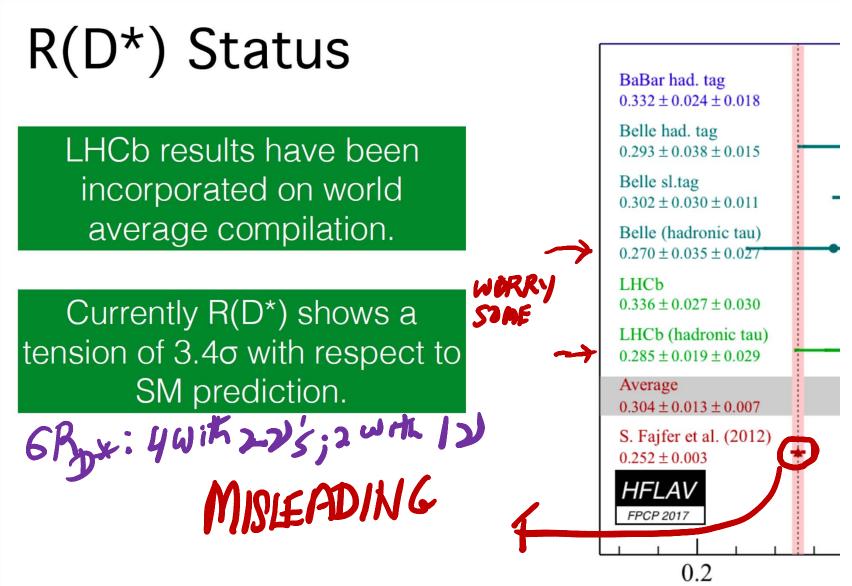
[LHCb , PRL113 (2014) 151601] [LHCb, LHCb-PAPER-2017-013] [BaBar, PRD 86 (2012) 032012] [Belle, PRL 103 (2009) 171801]

NB $R_{\rm K} \simeq 0.8$ is a prediction of one class of model explaining the $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ angular observables, see $L\mu$ - $L\tau$ models W. Altmannshofer et al. [PRD 89 (2014) 095033]

RECENT UPDATES FROM EXPERIMENT AND FROM THEORY



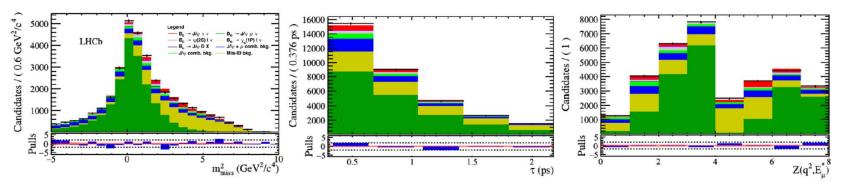
LIO International Conference on Flavour Physics 2018



 $B_c \to J/\psi \tau \nu$ 2 PM 2 on 30 18 Greg Ciezarek,



on behalf of the LHCb collaboration

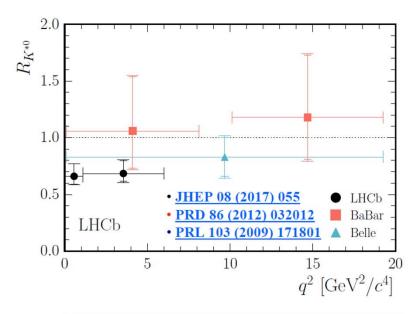


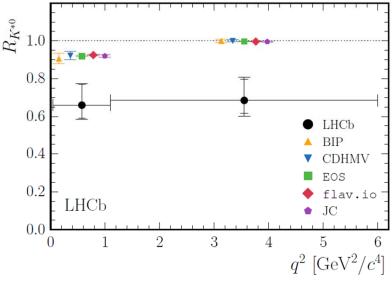
- $R_{J/\psi} \equiv B_c \rightarrow J/\psi \tau \nu/B_c \rightarrow J/\psi \mu \nu$
- Measured using very similar techniques to $\mathcal{R}(D^*)$, on run 1 data
- $R_{J/\psi} = 0.71 \pm 0.17 \pm 0.18$
 - $\sim 2\sigma$ from SM
 - But nearly as far from consistency with $\mathcal{R}(D^*)$
- LHCb-PAPER-2017-035(Run 1 data)



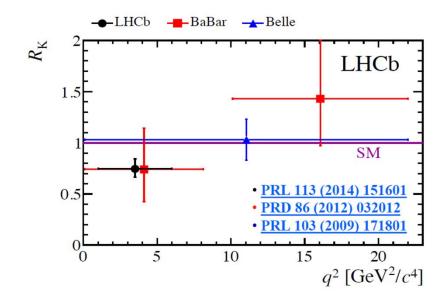


Lepton Universality results





RKX Shows similar results as RK



 R_K : Central- q^2 : 2.6 σ from SM

 R_{K^*} : Low- q^2 : 2.1-2.3 σ from SM

 R_{K^*} : Central- q^2 : 2.4-2.5 σ from SM

Rusa Mandal, PhD Thesis [IMSchemai]

• Another hint of deviation (at a level of more than 3σ), for a particular neutral-current decay mode is evinced by $B_s \to \phi \mu \mu$ [8,62,63].

$$\Phi = \frac{d}{dq^2} BR(B_s \to \phi \mu \mu) \Big|_{q^2 \in [1:6] \text{ GeV}^2} = \begin{cases} \left(2.58^{+0.33}_{-0.31} \pm 0.08 \pm 0.19\right) \times 10^{-8} \text{ GeV}^{-2} \text{ (exp.)} \\ (4.81 \pm 0.56) \times 10^{-8} \text{ GeV}^{-2} \end{cases}$$
(SM).

where $q^2 = m_{\mu\mu}^2$. Intriguingly, the q^2 region where this measurement has relatively low error (and data is quoted) is virtually the same as that for R_K and $R_{K^*}^{\text{central}}$. This

M. Patel C Lo: LHC may be able to give Rp.

That word be great

Flavor anomalies; Lyon; A Soni(BNL-HET)

Concerns on SM-theory

- Good news is that lattice[FERMIL-MILC] study largely confirms pheno calculations for R_D [our RBC-UKQCD, Witzel et al needs bit more time]
- For B=>D* no complete lattice study so far; 4 rather than 2 FF, so, from the lattice perspective, anticipate appreciably larger errors than for B=>D
- Therefore, O(1%) errors in RD* (and in fact smaller than in RD) are difficult to understand; lattice results should come in some months
- HFLAG should update the SM-theory with more realistic errors otherwise their fig is bit misleading
- Meantime recent phenomenological study of Bernlochner, Ligeti, Papucci and Robinson, 1703.05330 [and even more recently...is/are very timely and greatly appreciated.
- For now, for RD*, keeping these recent calculations and other reservations in mind best (conservative) guess is RD* ~ 0.258 +-0.020 [based on FERMIL-MILC error for RD]

REMARKABLY: FOR RD* CENTRAL VALUE OF BEST THEORY ESTIMATE APPEARS BIT HIGHER THAN ALL 6 MEASUREMENTS!



Comments on Theory

- For RD and RD*, non-lattice pheno efforts are on good, theo. grounds based on HQS....unlikely to be off by much
- Need however precise lattice results esp for B=>D*...see next p.
- For FCNC B=>K(*) II, LUV tests theory is essentially irrelevant
- FCNC, B=>K(*) II, absolute measured rates vs SM, theory is not reliable because of serious LD, non-perturbative contaminations

RBC-UKQCD [WITZEL, JUTTNER, TSANG, FLYNN, LEHNER, IZUBUCHI + AS] work in progress

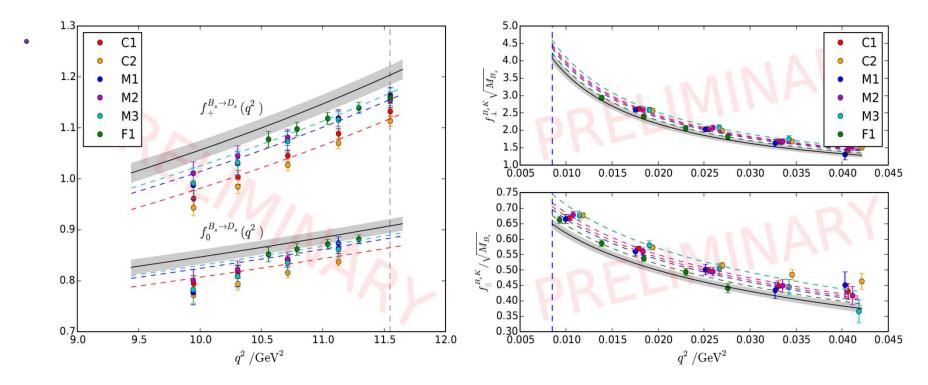


Figure 4. Chiral-continuum extrapolation for semi-leptonic form factors for $B_s \to D_s \ell \nu$ (left) and $B_s \to K \ell \nu$ (right). Performing a simple pole-ansatz for $B_s \to D_s$ we directly fit the phenomenological form factors f_+ and f_0 . For $B_s \to K$ we use heavy meson chiral perturbation theory and show the fit to the "lattice" form factors f_{\parallel} and f_{\perp} . The colored data points show results for our lattice calculations obtained at three values of the lattice spacing, whereas the black lines with the gray error band shows the chiral-continuum extrapolation. Only the statistical uncertainties are shown and no kinematical constraints are imposed.

Comments/Reservations pros & cons on Expts p1 of 2

- For RD(*), B=> D(*) τ v; most experimental results
 are with τ => μ v vi.e 2 v'sso D** potential
 contamination is a serious problem, in my view, as I
 have been stressing for past few years
- With that perspective in mind, it is worrysome that 2 new results, one from Belle and the other from LHCb with τ => hadrons + ν [i.e. 1 ν] is basically in agreement with SM though central value is mildly higher [i.e. roughly within ~ 1- σ]

Comments/Reservations pros & cons on Expts p2 of 2

HFLAV: Its high time they update their SM theory blob. It is true that combining theoretical predictions is bit difficult but it isnot good that they continue showing only predictions based on 2012 when recent ones with larger errors exist. This can be esp. misleading since using these recent calculations [not my work] will for sure reduce the tension with the SM.

See, as an illustration how ADS' handles this a bit later.

Reg. RK(*) mu/e UV

- Needless to say its of profound importance, if true
- If true not just B=>K, B=>K* but also Bs=> phi,
- B-baryon decays should show it
- Current statistics is marginal; more final states are needed and even more important other experiments esp. BELLE (II) confirmation is essential
- This can take years as Br are O(10⁻⁶) so not easy even for Belle-II;
- however, Belle-II will be able to do RXs...inclusive and that will likely have more sensitivity for them; See P Urquijo @ LIO
- OTOH, LHCB will have Bs and B-baryons; See M. Patel
 @ LIO

SURVEY OF BSMS

Array of BSM approaches to address flavor anomalies

Model ind eff L: Greljo; Camalich; Mandal;

LQ: Fajfer; Luzio; Neubert; Silva; Becirevic;

[Partial]Composite/warped: Stangl; Ahmady; Barbieri;

Panico; Blanke ... Natural set up for flavor Non-Universality

- SUSY-like: Hiller; ADS'=> RPV3 Natural flavor Non-Universality
- More BSMs: Guadagnoli; Grinstein; Jung; Ricciardi; Fuentes-Martin; Neshatpour; Crivelin
- New approaches: Valli; Camalich;
- Status + Outlook: Matias....Nice update ...some diff of opinion

EXPT: Status + (near) future outlook

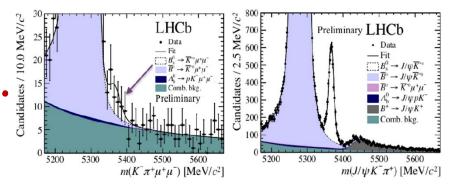
- Martinez; Lima; Bertolini @LHCb
- Urquijo; Goi; @ Belle-II
- Patel @ LHCb-UG

 $B_s^0
ightarrow ar{K^{*0}} \mu^+ \mu^-$ [LHCb-PAPER-2018-004]

CONGRATS CONGRATS LHCGTIC FANTASTIC TREMELY IMPTUS COCALHCLUPACES

Run1 + part Run2 (2015 and 2016) data (4.6 fb $^{-1}$)

 $N(B_s^0 o ar K^{*0} \mu^+ \mu^-) = 38 \pm 12$, 3.4σ above bkg-only hypothesis (first evidence) $\mathcal{B}(B_s^0 o ar K^{*0} \mu^+ \mu^-) = [2.9 \pm 1.0 (\mathrm{stat}) \pm 0.2 (\mathrm{syst}) \pm 0.3 (\mathrm{norm})] \times 10^{-8}$ (first measurement)

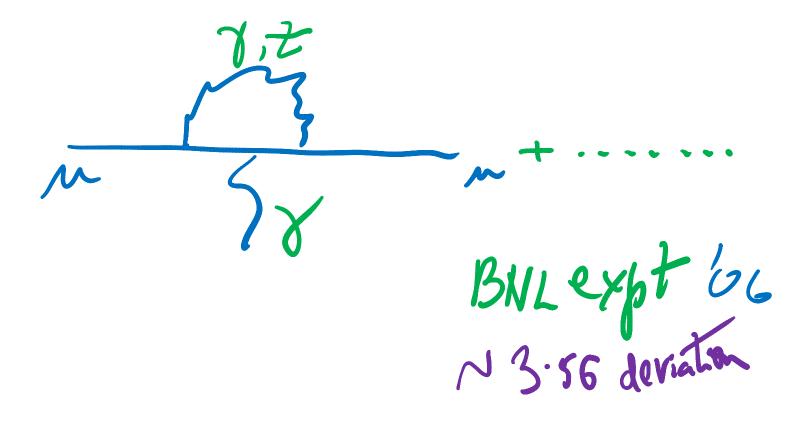


Menu...more

- SM/CKM....Hurth; Rusov
- Model independent tests@collider: Isidori;
 Kamenik;Hou; Nardecchia; ADS';
- Model Independent @ IF: Mandal; Fajfer; Neubert;

Kaons [SM/NP]: Buras; D'Ambrosio; Pich

NOW FEW WORDS ON MUON G-2



POSSIBLE CONNECTION OF G-2 TO OTHER FLAVOR ANOMALIES

MUON MAY NOT BE JUST A HEAVY ELECTRON: KILE, KOBACH AND AS

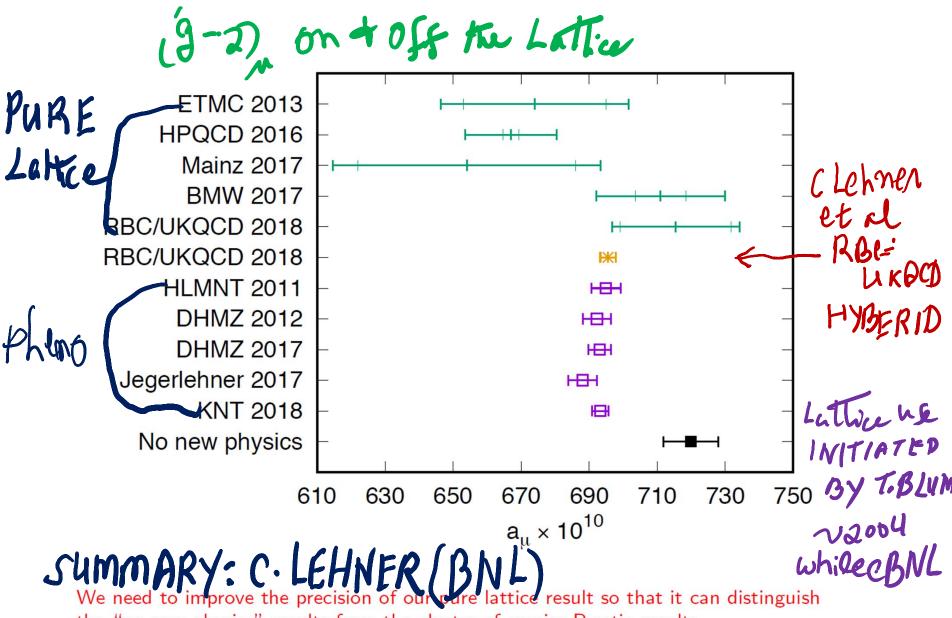


Table 1

Constraints on lepton-flavor violating and conserving processes. For the last four observables, the experimental null results are given in terms of a dimension-6 operator, suppressed by two orders of Λ , which can be interpreted as the nominal scale of new physics.

	Observable	Limit	PKD2015
	$Br(\mu \rightarrow 3e)$	$< 1.0 \times 10^{-12} [1]$	11/
	$Br(\mu \to e \gamma)$	$< 5.7 \times 10^{-13} [1]$	
Istgennot	Br(au o 3e)	$< 2.7 \times 10^{-8} [1]$	Spaultia
エリー	$Br(\tau \to e^- \mu^+ \mu^-)$	$< 2.7 \times 10^{-8}$ [1]	
a ative D	$Br(\tau \to e^+\mu^-\mu^-)$	$< 1.7 \times 10^{-8}$ [1]	
sensitive to	$Br(\tau \to \mu^- e^+ e^-)$	$< 1.8 \times 10^{-8}$ [1]	Mayla 1st
	$Br(\tau \to \mu^+ e^- e^-)$	$< 1.5 \times 10^{-8} [1]$	I I I I W WE I S'
NP	$Br(\tau \to 3\mu)$	$< 2.1 \times 10^{-8} [1]$	~
منائد ا	$Br(\tau \to \mu \gamma)$	$< 4.4 \times 10^{-8}$ [1]	
	$Br(\tau \to e\gamma)$	$< 3.3 \times 10^{-8} [1]$	gen. is
(0.2)	μ –e conversion	$\Lambda \gtrsim 10^3 \text{ TeV } [5]$	
(g-2)m	$e^+e^- \rightarrow e^+e^-$		Cuntombata
	$e^+e^- \rightarrow e^+e^-$ $e^+e^- \rightarrow \mu^+\mu^-$	$\Lambda \gtrsim 5 \text{ TeV } [3]$ $\Lambda \gtrsim 5 \text{ TeV } [3]$	
	$e^+e^- \rightarrow \tau^+\tau^-$	$\Lambda \gtrsim 3 \text{ TeV } [3]$ $\Lambda \gtrsim 4 \text{ TeV } [3]$	its protecte
		~	
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		- 11-20 A1	VEK F. LIVIV

QUICK UPDATE ON G-2



the "no new physics" results from the cluster of precise R-ratio results.

LUMCh Seminor 030118

Personal take on g-2

- If you take pheno estimate of hadronic VP contributions via use of R-ratio method deviation for BNL-expt ~3.6 σ so likely culprit is under-estimate error on theory of around ½%; though recently RBC-UKQCD lattice hybrid method finds support for this pheno estimate
- Need to wait on pure lattice result after another factor of 4-5 reduction in error, may take another 2-3 years
- By that time improved experimental results should also become available
- Final verdict may need another 2-3 years

Bottom line

- NP or not depends critically not just on precise experiment but also reliable SM prediction from the lattice become mandatory
- Experiment + Lattice M.E. has the last word....[of course should be stressed that the lattice calculations often require sophisticated and demanding and essential input from perturbation theory]
- Experimental results often attained at huge cost can be used effectively, iff commensurate theory predictions are available......mantra for past several decades

A.S. in Proceedings of Lattice '85 (FSU)..1st Lattice meeting ever attended

The matrix elements of some penguin operators control in the standard model another CP violation parameter, namely ϵ'/ϵ . $^{6,8)}$ Indeed efforts are now underway for an improved measurement of this important parameter. 10) In the absence of a reliable calculation for these parameters, the experimental measurements, often achieved at tremendous effort, cannot be used effectively for constraining the theory. It is therefore clearly important to see how far one can go with MC techniques in alleviating this old but very difficult

With C. Bernard [UCLA]

malies; Lyon; A Soni(BNL-HET)

• Already heard pheno perspective on K=> $\pi\pi$ and ϵ' in nice talks from Buras and also from Pich

So I skip exptal route for ε' and such

LATTICE QUEST FOR EPSILON'

$$Re\left(\frac{\epsilon'}{\epsilon}\right) = \frac{\omega}{\sqrt{2}|\epsilon|} \left[\frac{\operatorname{Im}(A_2)}{\operatorname{Re}(A_2)} - \frac{\operatorname{Im}(A_0)}{\operatorname{Re}(A_0)}\right], \quad \omega = \frac{R_0 A_0}{R_0 A_0}.$$

$$ReA_0 = \frac{\delta}{\sqrt{2}|\epsilon|} \left[\frac{\operatorname{Im}(A_2)}{\operatorname{Re}(A_2)} - \frac{\operatorname{Im}(A_0)}{\operatorname{Re}(A_0)}\right], \quad \omega = \frac{R_0 A_0}{R_0 A_0}.$$

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$$ReA_0 = \frac{\operatorname{Im}(A_0)}{\operatorname{Im}(A_0)} - \frac{\operatorname{Im}(A_0)}{\operatorname{Im}(A_0)}$$

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$$ReA_0 = \frac{\operatorname{Im}(A_0)}{\operatorname{Im}(A_0)} - \frac{\operatorname{Im}(A_0)}{\operatorname{$$

ton Simplicity: 1St Strategy via ChPT

PHYSICAL REVIEW D

VOLUME 32, NUMBER 9

1 NOVEMBER 1985

Application of chiral perturbation theory to $K \rightarrow 2\pi$ decays

Claude Bernard, Terrence Draper,* and A. Soni
Department of Physics, University of California, Los Angeles, California 90024

H. David Politzer and Mark B. Wise

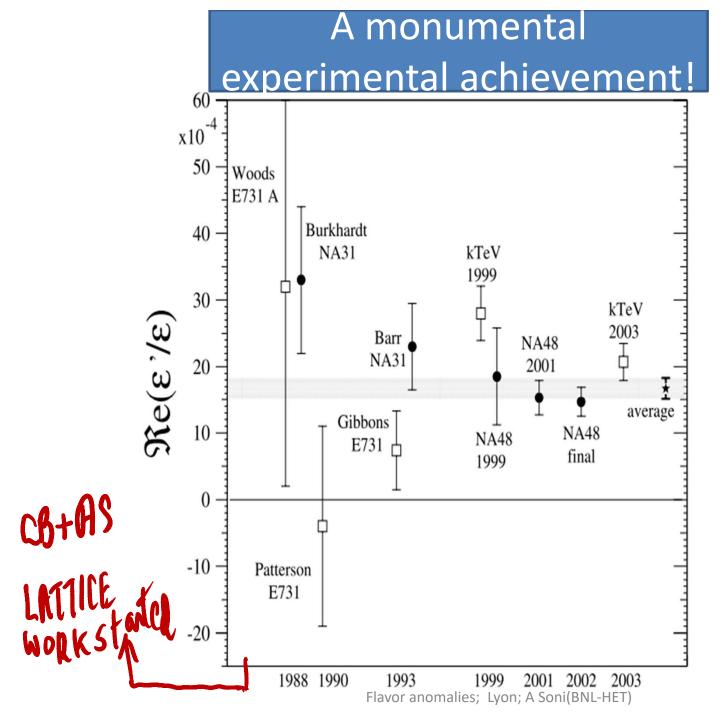
Department of Physics, California Institute of Technology, Pasadena, California 91125 (Received 3 December 1984)

Chiral perturbation theory is applied to the decay $K \rightarrow 2\pi$. It is shown that, to quadratic order in meson masses, the amplitude for $K \rightarrow 2\pi$ can be written in terms of the unphysical amplitudes $K \rightarrow \pi$ and $K \rightarrow 0$, where 0 is the vacuum. One may then hope to calculate these two simpler amplitudes with lattice Monte Carlo techniques, and thereby gain understanding of the $\Delta I = \frac{1}{2}$ rule in K decay. The reason for the presence of the $K \rightarrow 0$ amplitude is explained: it serves to cancel off unwanted renormalization contributions to $K \rightarrow \pi$. We make a rough test of the practicability of these ideas in Monte Carlo studies. We also describe a method for evaluating meson decay constants which does not require a determination of the quark masses.

C. Bernard, A. Soni / Weak matrix elements on the lattice 0.60 \Box a $\triangle d$ 0.45 1.2 * f $(GeV)^4$ PATHEOLUM 0.8 $\Box \alpha$ 0.15 4.0 0 C 0.00 0.12 0.25 0.37 0.50 0.62 0.75 0.87 1.00 m² (GeV)² 0.0 $^{0.4}$ $^{0.6}$ $^{0.6}$ $^{0.6}$ 0.0 0.2 8.0

Flavor anomalies; Lyon; A Soni(BNL-H

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Komad Kleinknecht "" "UncabiyC?V

16.6(2.3) X10 PDG InspiredI.P. by papers of Shamir [+Furman] + discussions with Creutz

QCD with domain wall quarks

T. Blum* and A. Soni[†]

Department of Physics, Brookhaven National Laboratory, Upton, New York 11973

(Received 27 November 1996)

1St Simulation With DUQ

DNO proserve CHIRAL SYM We present lattice calculations in QCD using <u>Shamir's variant of Kaplan fermions</u> which retain the continuum $SU(N)_L \times SU(N)_R$ chiral symmetry on the lattice in the limit of an infinite extra dimension. In particular, we show that the pion mass and the four quark matrix element related to $K_0 - \overline{K_0}$ mixing have the expected behavior in the chiral limit, even on lattices with modest extent in the extra dimension, e.g.,

expected behavior in the chiral limit, even on lattices with modest extent in the extra dimension, e.g., $N_c = 10. [S0556-2821(97)00113-6]$ even $M_{LL}/|\langle K|\overline{s}\gamma_5 d|0\rangle|^2$ 0.6 0.0 0.20 0.0 0.00 0.10 0.15 0.05 MAJOR BREAKTHROUGHLFOIK-ITT Letter Could K->2T ChPT

With DWG in Quench

PHYSICAL REVIEW D 68, 114506 (2003)

Kaon matrix elements and CP violation from quenched lattice QCD: The 3-flavor case

T. Blum, P. Chen, N. Christ, C. Cristian, C. Dawson, G. Fleming, R. Mawhinney, S. Ohta, G. Siegert, A. Soni, P. Vranas, M. Wingate, L. Wu, and Y. Zhestkov

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 ⁵IBM Research, Yorktown Heights, New York 10598, USA

(Received 19 July 2002; published 30 December 2003)

We report the results of a calculation of the $K \rightarrow \pi\pi$ matrix elements relevant for the $\Delta I = 1/2$ rule and ϵ'/ϵ in quenched lattice QCD using domain wall fermions at a fixed lattice spacing a^{-1} ~2 GeV. Working in the three-quark effective theory, where only the u, d, and s quarks enter and which is known perturbatively to next-to-leading order, we calculate the lattice $K \rightarrow \pi$ and $K \rightarrow |0\rangle$ matrix elements of dimension six, fourfermion operators. Through lowest order chiral perturbation theory these yield $K \rightarrow \pi\pi$ matrix elements, which we then normalize to continuum values through a nonperturbative renormalization technique. For the ratio of isospin amplitudes $|A_0|/|A_2|$ we find a value of 25.3 \pm 1.8 (statistical anorom), empared to the value of 22.2, with individual isospin amplitudes 10%-20% below the experimental values. For ϵ'/ϵ , using known central values for standard model parameters, we calculate \ 4.0\pm 2.3\)\times 10\pm (statistical error only compared to the current experimental average of (17.2±1.8)×10⁻⁴. Because we find a large cancellation between the I=0 and I=2 contributions to ϵ'/ϵ , the result may be very sensitive to the approximations employed. Among these are the use of quenched QCD, lowest order chiral perturbation theory, and continuum perturbation theory below 1.3 GeV. We also calculate the kaon B parameter B_K and find $B_{K,MS}(2 \text{ GeV})$ =0.532(11). Although currently unable to give a reliable systematic error, we have control over statistical errors and more simulations will yield information about the effects of the approximations on this firstprinciples determination of these important quantities,

RBC Collabonation

x 17F

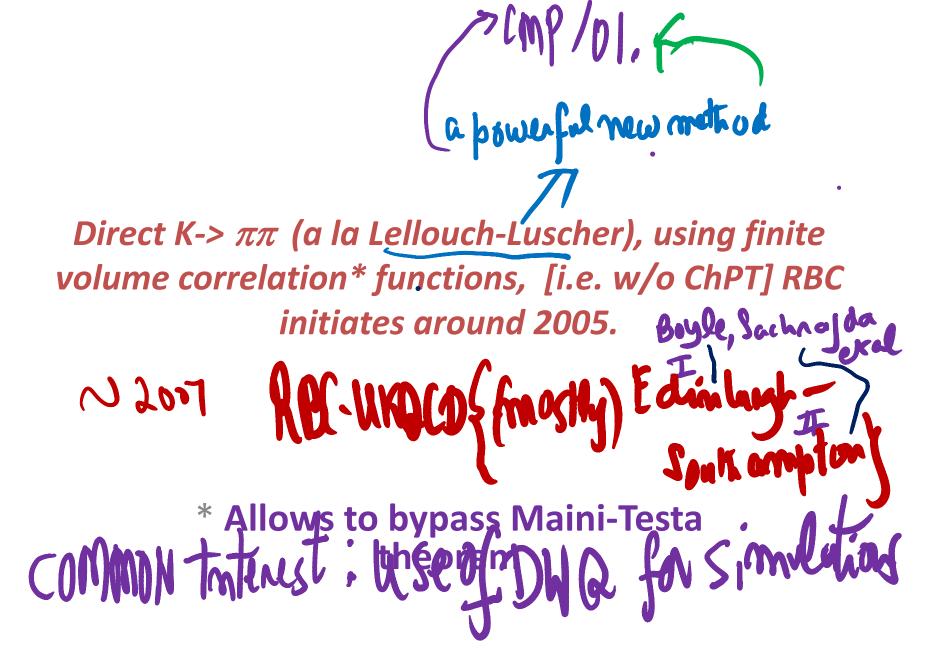
Flaship Project

Now 120 yrs!

Ist laye Sale Simbolia With DUO

PRO103

45



$$\Delta S = 1 \text{ He KLD}$$

$$Buchalla, Bunas, Lautenleider Guchimiert$$

$$H_W = \frac{G_F}{\sqrt{2}} V_{us}^* V_{ud} \sum_{i=1}^{10} \left[z_i(\mu) + \tau y_i(\mu) \right] Q_i(\mu).$$

$$V_{ReAD}$$

$$V_{ReAD}$$

$$M_{i} = \langle K|Q_{i}|\pi i \rangle \text{ from the } T = -V_{i}^{*}V_{i}/V_{i}^{*}V_{i}$$
Needed eather

$$Q_1 = (\bar{s}_{\alpha} d_{\alpha})_L (\bar{u}_{\beta} u_{\beta})_L,$$

$$Q_2 = (\bar{s}_{\alpha} d_{\beta})_L (\bar{u}_{\beta} u_{\alpha})_L,$$

$$Q_2 = (\bar{s}_{\alpha} d_{\beta})_L (\bar{u}_{\beta} u_{\alpha})_L,$$

$$Q_3 = (\bar{s}_{\alpha} d_{\alpha})_L \sum_{q=u,d,s} (\bar{q}_{\beta} q_{\beta})_L,$$

$$Q_4 = (\bar{s}_{\alpha} d_{\beta})_L \sum_{q=u,d,s} (\bar{q}_{\beta} q_{\alpha})_{L_2}$$

$$Q_{3} = (\bar{s}_{\alpha}d_{\alpha})_{L} \sum_{q=u,d,s} (\bar{q}_{\beta}q_{\beta})_{L},$$

$$Q_{4} = (\bar{s}_{\alpha}d_{\beta})_{L} \sum_{q=u,d,s} (\bar{q}_{\beta}q_{\alpha})_{L},$$

$$Q_{5} = (\bar{s}_{\alpha}d_{\alpha})_{L} \sum_{q=u,d,s} (\bar{q}_{\beta}q_{\beta})_{R},$$

$$Q_{6} = (\bar{s}_{\alpha}d_{\beta})_{L} \sum_{q=u,d,s} (\bar{q}_{\beta}q_{\alpha})_{R},$$

$$Q_6 = (\bar{s}_{\alpha} d_{\beta})_L \sum_{\alpha} (\bar{q}_{\beta} q_{\alpha})_R,$$

$$Q_7 = \frac{3}{2} (\bar{s}_{\alpha} d_{\alpha})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\beta})_R,$$

$$Q_7 = \frac{3}{2} (\bar{s}_{\alpha} d_{\alpha})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\beta})_R,$$

$$Q_8 = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\alpha})_R,$$

$$Q_9 = \frac{3}{2} (\bar{s}_{\alpha} d_{\alpha})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\beta})_L,$$

$$3 (\bar{s}_{\alpha} d_{\alpha})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\beta})_L,$$

$$Q_9 = \frac{3}{2} (\bar{s}_{\alpha} d_{\alpha})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\beta})_L,$$

$$Q_{10} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\alpha})_L,$$

Flavor anomalies; Lyon; A Soni(BNL-HET)

Acknowledge many significant contributions off & on the lattice

- While focus is on lattice calculations of K=>ππ primarily by our RBC-UKQCD Collab
- Over the years many important contributors, in particular:
- (Mary K)Gaillard, (Ben) Lee; Altarelli, Maiani;
 Shifman, Vainshtein, Zhakrov; Gilman + Wise;
 Buras & Co; Martinelli & Co; (Claude)
 Bernard; de Rafael; Pich, Bijnens.....

DIRFCT K-> TITE

[No (hf 7) · Using Re(A)

Im(A₀) and I

Results for ε'

• Using Re(A) and Re(A) from experiment $Im(A_0)$ and $Im(A_2)$ and the phase shifts,

and our lattice values for

USING 216 independent measurements

Ľ15

 $\left(\frac{\varepsilon'}{\varepsilon}\right) = \operatorname{Re}\left\{\frac{i\omega e^{i(\delta_2 - \delta_0)}}{\sqrt{2}\varepsilon}\right\}$

 ReA_2

 $-\frac{\mathrm{Im}A_0}{\mathrm{Re}A_0}$

 $\overline{\operatorname{Re} A_0}$ (80-85%) (this work)

RBC-UKQCD PRL'15 EDITOR'S CHOICE

 $= 1.38(5.15)(4.43) \times 10^{-4},$ $16.6(2.3) \times 10^{-4}$

(experiment)

Factorionality of all emony

LARGE CANCELLATION!!

Bearing in mind the largish errors in this first calculation, we interpret that our result are consistent with experiment at $\sim 2\sigma$ level

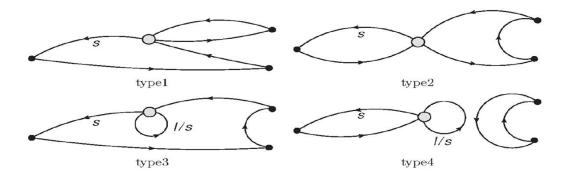
w= Re Az ~ 0.145

Computed ReA2 excellent agreement with expt
Computed ReA0 good agreement with expt
Offered an "explanation" of the Delta I=1/2
Enhancement [c later]

Scalars 2017: HET-BNL: son

12

BURAS et al use ow IMF => effect is ~ 2.95 Niente: Flavor anomalies; Lyon; A Soni(BNL-HET)



RBC-UKQCD PRL 2015

FIG. 1. Examples of the four types of diagram contributing to the $\Delta I=1/2,\,K\to\pi\pi$ decay. Lines labeled ℓ or s represent light or strange quarks. Unlabeled lines are light quarks.

4.66(0.96)(0.27)

Tot

 $\operatorname{Im}(A_0)(\operatorname{GeV})$ $Re(A_0)(GeV)$ $1.02(0.20)(0.07) \times 10^{-7}$ 0 $3.63(0.91)(0.28) \times 10^{-7}$ $-1.19(1.58)(1.12) \times 10^{-10}$ $1.54(2.04)(1.45) \times 10^{-12}$ $-1.86(0.63)(0.33) \times 10^{-9}$ $1.82(0.62)(0.32) \times 10^{-11}$ $-8.72(2.17)(1.80) \times 10^{-10}$ $1.57(0.39)(0.32) \times 10^{-12}$ $3.33(0.85)(0.22) \times 10^{-9}$ $-3.57(0.91)(0.24) \times 10^{-11}$ $2.40(0.41)(0.00) \times 10^{-11}$ $8.55(1.45)(0.00) \times 10^{-14}$ $-1.33(0.04)(0.00) \times 10^{-10}$ $-1.71(0.05)(0.00) \times 10^{-12}$ $-7.12(1.90)(0.46) \times 10^{-12}$ $-2.43(0.65)(0.16) \times 10^{-12}$ $7.57(2.72)(0.71) \times 10^{-12}$ $-4.74(1.70)(0.44) \times 10^{-13}$ 10

-1.90(1.19)(0.32)

TABLE I. Contributions to A_0 from the ten continuum, $\overline{\text{MS}}$ operators $Q_i(\mu)$, for $\mu = 1.53$ GeV. Two statistical errors are shown: one from the lattice matrix element (left) and one from the lattice to $\overline{\text{MS}}$ conversion (right). See the Supplemental Material at [*URL to be inserted*] for tables of the separate matrix elements in the lattice, RI/SMOM and $\overline{\text{MS}}$ schemes as well as the renormalization matrices which relate them.

While ReA0 and ReA2 and δ 2 agree well with expt a possible difficulty: δ0

The continuum and our lattice determinations of strong phase

difference differs at the ~2
$$\sigma$$
 level:
$$\phi_{\varepsilon'} = \delta_2 - \delta_0 + \frac{\pi}{2} = \begin{cases} (42.3 \pm 1.5)^{\circ} \text{ PDG} & \text{Colongila stal} \\ (54.6 \pm 5.8)^{\circ} & \text{RBC-UKQC} \end{cases}$$
Fortunately, due to the central value of the combination $\delta_2 - \delta_0 + \pi/2 - \phi_{\varepsilon}$ and to the large uncertainties in the determination of the various matrix elements, these two

Fortunately, due to the central value of the combination $\delta_2 - \delta_0 + \pi/2 - \phi_\varepsilon$ and to the large uncertainties in the determination of the various matrix elements, these two choices yield almost identical results; for the large uncertainties in the determination of the various matrix elements, these two choices yield almost identical results; for the large uncertainties in the determination of the various matrix elements, these two choices yield almost identical results; for the large uncertainties in the determination of the various matrix elements, these two choices yield almost identical results; for the large uncertainties in the determination of the various matrix elements, these two choices yield almost identical results; for the large uncertainties in the determination of the various matrix elements, these two choices yield almost identical results; for the large uncertainties in the determination of the various matrix elements, the large uncertainties in the determination of the various matrix elements.

Sensitivity of ε' to strong phase(s)

Significant effort is being put now to calculate and understand pi pi rescattering phases from diff pts of view

Generation of New gaye empigs For past 7, 3 years

SUPERCOMPUTERS OVER 3 CONTINENTS!

Progress in the calculation of arepsilon' on the lattice

C. Relly

Resource	Million BG/Q equiv core-hours	Independent cfgs.
USQCD (BNL 512 BG/Q nodes)	50	220
RBRC/BNL (BNL 512 BG/Q nodes)	17	50
UKQCD (DiRAC 512 BG/Q nodes)	17	50
NCSA (Blue Waters)	108	380
KEK (KEKSC 512 BG/Q nodes)	74	296
Total	266	996

Table 1: A breakdown of the various resources we intend to utilize. Note that we require 4 molecular dynamics time units per independent configuration.

4 hilf statems

By mw ~ 1200

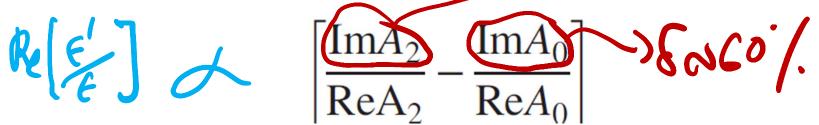
mershame 54

Guess estimate of reduction of errors

- δ (Im A0) from 65% => 20-25%
- δ (ReA0) from 35% => 15-20% [don't use for ϵ ' for now]
- Uncorrelated fits (due to lack of stat) =>

Very good chance we'll be able to correlated fits now Systematic error from ~ 27% => ~ 20%

- Effect on ε' unclear: 'cause of hefty cancellations



Additional Improvements/checks in lattice E' determination underway for past ~2 years

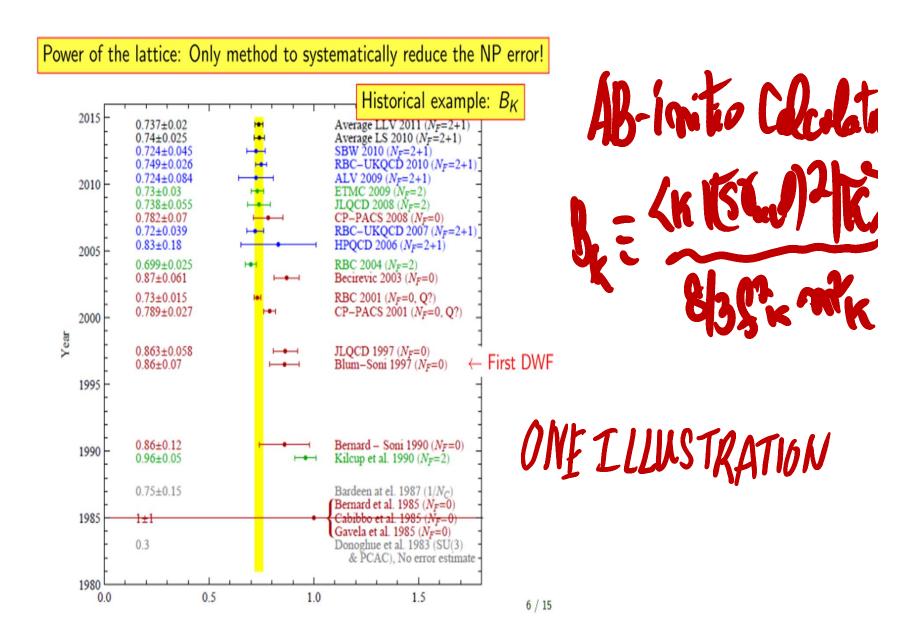
• EM+ isospinfforder extensive study Ciniglians, Neufoly Educat Pich 13 U Com PhD Student Dante Completely diff method(s) • 1) excited pipi stateph 1 (Im Muyly ChlT was!!
• 11) Revisit ChPT well for physical 11, K moss of e NLO!

Reexamine BDSPW 84, Lailot + AS'04 for E

underlying method is systematically improvable =>multitude of successful demonstrations by now

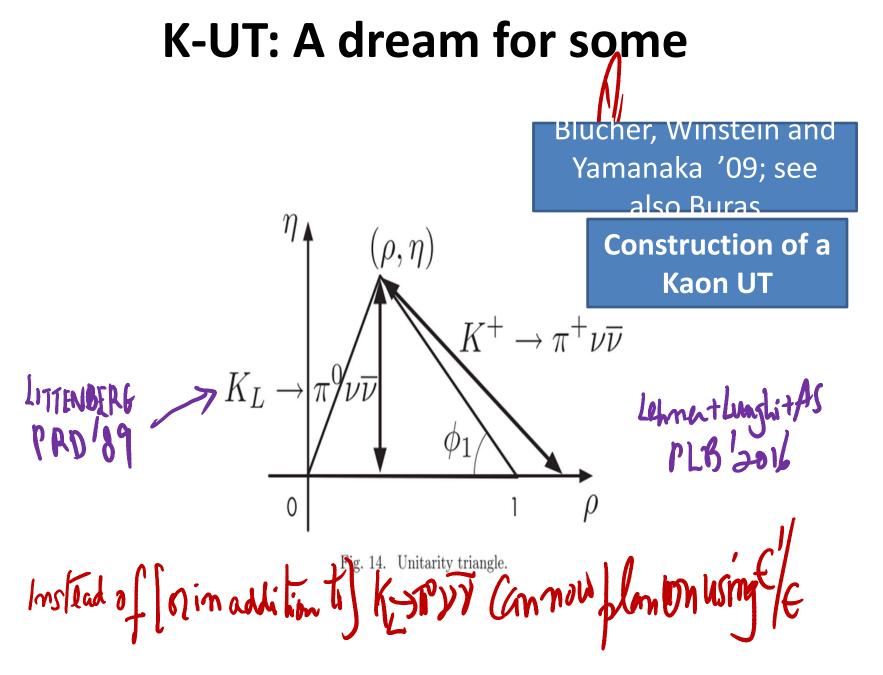
- BK in full QCD with DWF '07 RBC-UKQCD error O(7%)
- Since ~2012 many discretizations, WA error O(1-2%)
- Re A2 from ~25% around 2012 to now ~10% (now no longer due to lattice but only only due to perturbation theory error upto NLO!)
- KI3, A2, fB's , BB's......
- Quark masses; in particular ms no longer anywhere around ~150
 MeV [used to be PDG value] but now ~100 MeV.
- No doubt that A0 and ε' will also go that way for quite sometime to come......to ~10% total in another ~ 3-5years!.

After that EM& isospin effects need to be ascertained quantitatively; WIP

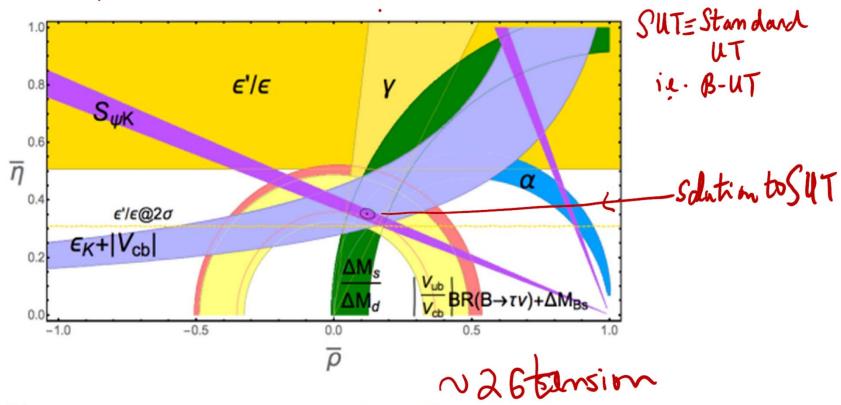


MANY IMPLICATIONS, AS ONE EXAMPLE:K-UT

See Lehner, lingli + ASP18/6







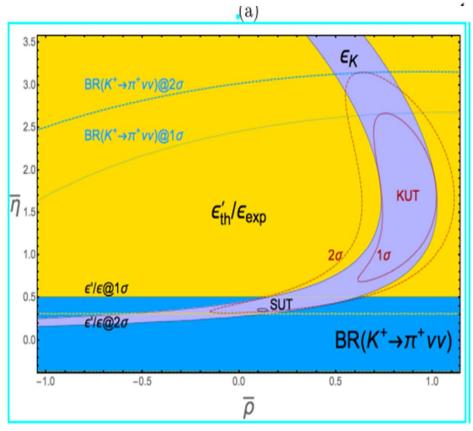
9/18/2017

LmC2017; SIEGEN; HET-BNL;soni

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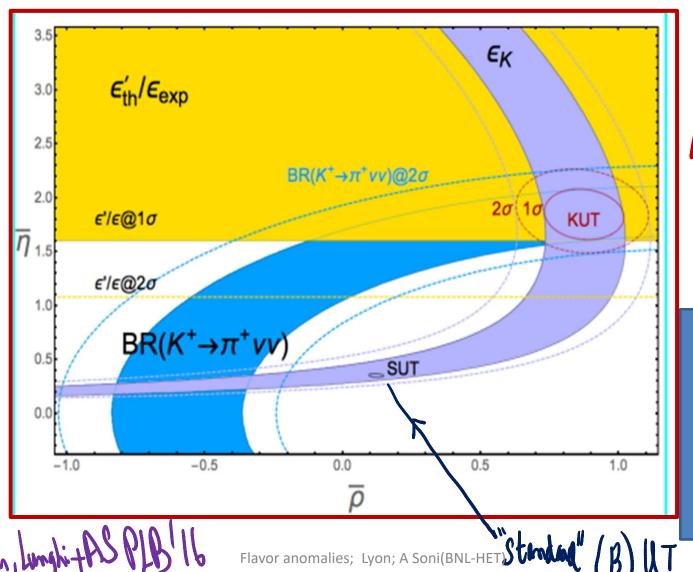
Sketch of an emerging K-W: 3 keykamic imputs $\mathcal{L} = \begin{cases}
(8.64 \pm 0.60) \times 10^{-11} \text{ SM} \\
(17.3^{+11.5}_{-10.5}) \times 10^{-11} \text{ E949}
\end{cases}$

$$BR(K^{+} \to \pi^{+} \nu \bar{\nu}) = \begin{cases} (8.64 \pm 0.60) \times 10^{-11} \text{ SM} \\ (17.3^{+11.5}_{-10.5}) \times 10^{-11} \text{ E949} \end{cases}$$



$$\operatorname{Re}\left(\frac{\varepsilon'}{\varepsilon}\right)_{K} = \begin{cases} (16.7 \pm 1.6) \times 10^{-4} \\ (1.36 \pm 5.21_{\text{stat}} \pm 4.49_{\text{syst}}) \times 10^{-4} \end{cases}$$

Possible KUT Circa 2020: DUE NA62+ RBG-UKBGD



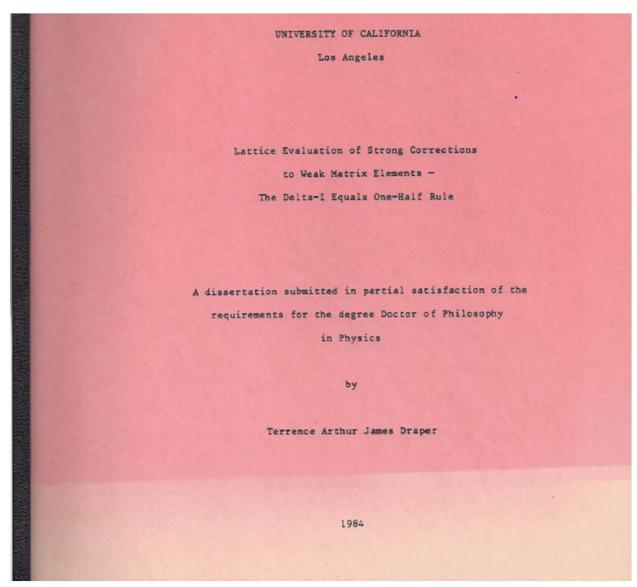
Assumed: NA62, 100 events with ~7% error RBC-UKQCD, δ(ImA0)~18% [current ~60%]

Lehner, Lunghi+ASPLB/16

ITS MY ~36TH YEAR ON K=>ΠΠ & ε'.....WHY?



The 1st PhD Nesio



Grew from

End of year

Been Pandy

A June 20, 1882.

[ULLA]

MOTHER of all (lattice) calculations to date: A Personal Perspective

- Calculation K=> $\pi\pi$ & ϵ' were the reasons I went into lattice over 1/3 of a century ago!
- 9 + (3 new) PhD thesis: Terry Draper (UCLA'84), George Hockney(UCLA'86), Cristian Calin (Columbia=CU'01), Jack Laiho(Princeton'04), Sam Li(CU'06), Matthew Lightman(CU'09), Elaine Goode(Southampton'10), Qi Liu(CU'12), Daiqian Zhang(CU'15)+ [new ones starting from CU, U Conn and Southampton] + many PD's & junior facs.. obstacles & challenges (and of course "mistakes"!) ad infinitum.....

WHY FOCUS with SUCH intense DETERMINATION All these many many years?

UNDERLYING REALIZATION E': MOST LIKELY A GEM IN SEARCH OF NEW PHENOMENA

Contrarian/Complementary view

- flavor physics is actually hanging by perhaps the weakest link i.e. a single CP-phase endowed by the 3g –SM.
- In many ways this is a contrarian (or complementary)
 point of view, in sharp contrast to the overwhelming
 majority following the naturalness lamp post via Higgs
 radiative stability.

 E' due to its miniscule value, esp because it results from unnatural large cancellations seemed clearly highly vulnerable...The mantra being followed for a very very long time

Its presumed importance:

- lies in its very small size => Perhaps new phenomena has a better chance of showing up
- Smallness also renders it exceedingly sensitive monitor of flavor –alignment
- Simple naturalness arguments strongly suggest ε' very sensitive to BSM – CP odd phases
- In many ways, (superficially) ε' is rather analogous to nedm......both being very sensitive to BSM-CP phases; however, key diff for (now) nedm expt is the key, theory has marginal role, in sharp contrast to ε'
- Understanding ε', nedm are extremely important for uncovering new physics and/or learning how naturalness really works in nature



IF YOU BUILD IT THEY WILL COME

If there is new physics around below ~5 TeV, there is an excellent chance that ε' will find it!

[of course requires accurate theory calculation... RBC-UKQCD plans for X5 in stat and appreciable improvements in systematic in ~2 years]

9/18/2017 LmC2017; SIEGEN; HET-BNL;soni

Soni's ε' Rule

If ε' ends up being due SM then:

- I. Other anomalies, RD(*), RK(*), g-2 will ALL go away

 9fso, that
- 2. It will give us a vitally important new info on naturalness

MODEL INDEPENDENT IMPLICATIONS OF RD(*) ANOMALIES FOR [LHC] COLLIDER EXPERIMENTS: SAMPLE ILLUSTRATIONS

Implications of Anomalies

For Collider:From Isidori; Kamenik; ADS';
 Luzio; Hou;

 FOR IF: Guadagnoli; Fajfer; Neubert; Mandal; Fuentes-Martin

Conclusions

FUENTES-MARTIN

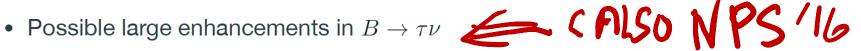
It is possible to find well-motivated combined explanations of the B anomalies, compatible with current data in both low- and high-energies ALSO Mande

The model I presented, a Pati-Salam flavor deconstruction, predicts several characteristic smoking-gun signatures that differentiate it from other solutions

- $\Delta R_{D^*}/\Delta R_D \simeq 0.45$
- NP in $b \to s \tau \tau \sim b \to s \mu \mu$ and $b \to s \tau \mu \sim \mathcal{O}(5) \times b \to s \mu \mu$



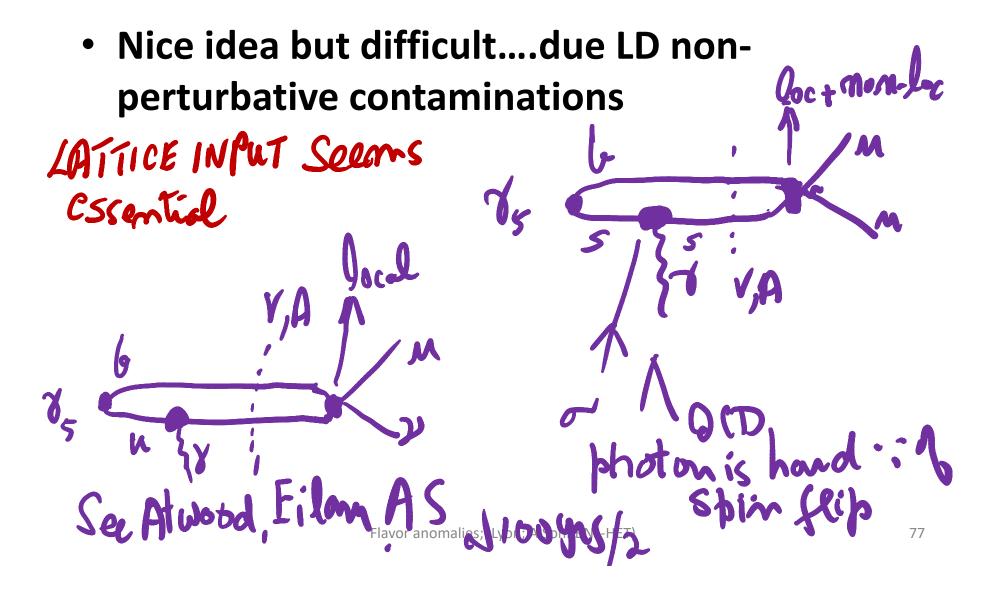
- Possibility of a deficit in $\Delta B = 2$ & CP violation involving the 3rd quark family



If the anomalies are really pointing to NP, new experimental indications (both in high- p_T and at low energies) will show up soon in several observables

Exciting times ahead of us. Let's have fun

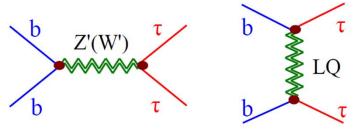
Guadagnoli B(s)=>mu mu gamma



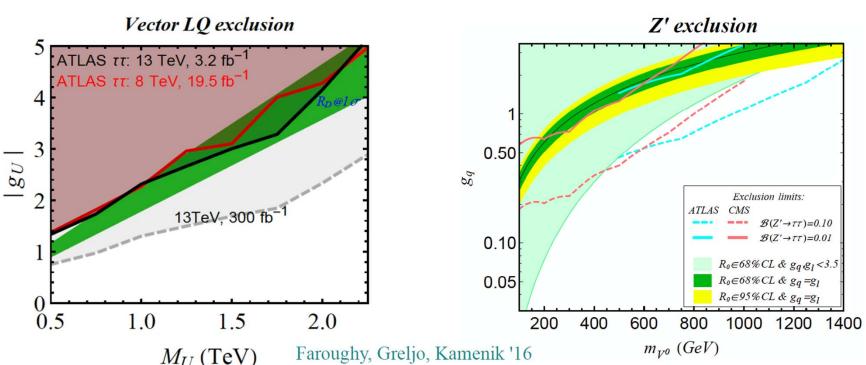
EFT-type considerations [The main problems]

Three main problems identified in the recent literature (*driven mainly by* R_D ...):

I. <u>high</u>-p_T <u>constraints</u>

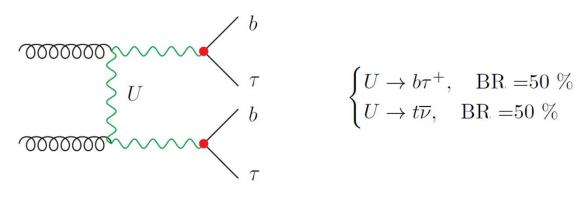


Naïve EFT scale [from R_D - setting g, $\lambda \to 1$]: $\Lambda \sim 700 \text{ GeV}$

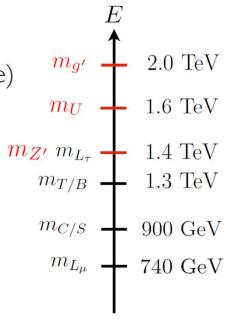


High-p_T searches

- LQ pair production via QCD
- 3rd generation final states (fixed by anomaly and SU(2)_L invariance)



$$\begin{cases} U \to b\tau^+, & \text{BR} = 50 \% \\ U \to t\overline{\nu}, & \text{BR} = 50 \% \end{cases}$$



[CMS search for spin-0, 1703.03995 recast for spin-1 1706.01868 (see also 1706.05033) + Moriond EW update]

 $m_U \gtrsim 1.5 \text{ TeV}$



LQ mass sets the overall scale: $M_{g'} \simeq \sqrt{2} M_U \quad M_{Z'} \simeq \frac{1}{\sqrt{2}} M_U$

ALTMANNSHOFFR, Devtas 1704.06659 + Sey WIP

MODEL INDEPENDENT IMPLICATIONS OF RD(*) ANOMALIES FOR [LHC] COLLIDER EXPERIMENTS

- In a nut-shell B-experiments seem to find anomalous behavior in the underlying b=>c tau nu
- This necessarily [by XSym] implies there should be analogous anomaly in g + c => b tau nu...=>pp => b tau nu

 Thus it immediately leads to inescapable search channels for possible NP at the high energy frontier for ATLAS & CMS and these are urgently urged

Xsymm implications of anomalies for colliders



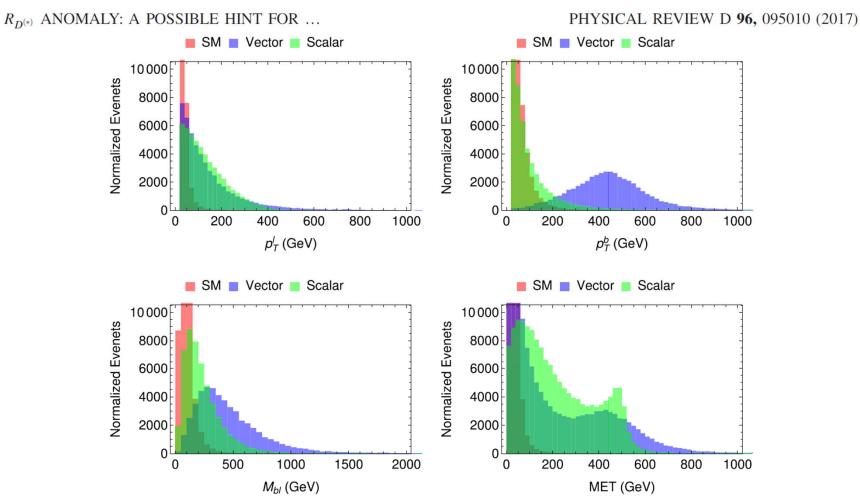


FIG. 1. Normalized kinematic distributions for the $pp \to b\tau\nu \to b\ell + \not\!\!E_T$ signal and background.

EXPECT DISTINCTIVE NP CONTRIBUTIONS IN COLLIDERS

ANOMALY: POSSIBLY A HINT FOR (NATURAL) SUSY-WITH RPV

- ASSUMING the anomaly is REAL & HERE TO STAY [BIG ASSUMPTION due to caveats mentioned]
- Anomaly involves simple tree-level semi-leptonic decays
- Also b => tau (3rd family)
- Speculate: May be related to Higgs naturalness
- Seek minimal solution: perhaps 3rd family super-partners(a lot) lighter than other 2 gens > proton decay concerns may not be relevant=> RPV ["natural" SUSY]
- RPV natural setting for LUV ...can accommodate g-2 and eps' if needs be
- Collider signals tend to get a lot harder than (usual-RPC) SUSY

RPV3 preserves gange coupling unification i mespecture of ## of elsection gens. 1, 2 on 3.

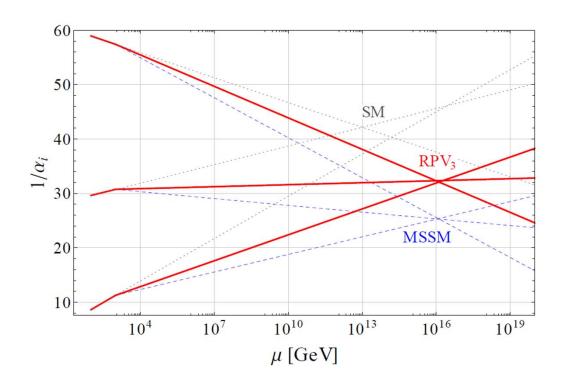


FIG. 2. RG evolution of the gauge couplings in the SM, MSSM and with partial supersymmetrization.

Unification scale astoys some, only value of couplings hifts

For phono relayant tems:

ADS'PRD 2017

$$\mathcal{L} = \lambda'_{ijk} \left[\tilde{\nu}_{iL} \bar{d}_{kR} d_{jL} + \tilde{d}_{jL} \bar{d}_{kR} \nu_{iL} + \tilde{d}_{kR}^* \bar{\nu}_{iL}^c d_{jL} \right]$$
$$-\tilde{e}_{iL} \bar{d}_{kR} u_{jL} - \tilde{u}_{jL} \bar{d}_{kR} e_{iL} - \tilde{d}_{kR}^* \bar{e}_{iL}^c u_{jL} \right] + \text{H.c.}$$

$$\mathcal{L}_{\text{eff}} \supset \frac{\lambda'_{ijk} \lambda'^*_{mnk}}{2m_{\tilde{d}_{kR}}^2} \left[\bar{\nu}_{mL} \gamma^{\mu} \nu_{iL} \bar{d}_{nL} \gamma_{\mu} d_{jL} - \nu_{mL} \gamma^{\mu} e_{iL} \bar{d}_{nL} \gamma_{\mu} \left(V^{\dagger}_{\text{CKM}} u_L \right)_j + \text{h.c.} \right] - \frac{\lambda'_{ijk} \lambda'^*_{mjn}}{2m_{\tilde{u}_{jL}}^2} \bar{e}_{mL} \gamma^{\mu} e_{iL} \bar{d}_{kR} \gamma_{\mu} d_{nR} ,$$

For addressing RK(*) in RPV, see e.g. Das et al , 1705.09188

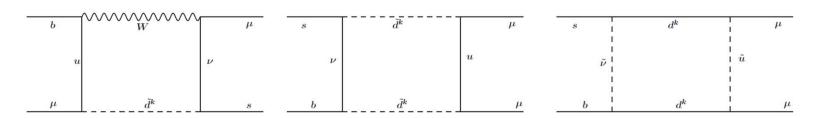


FIG. 1: Representative diagrams for $b \to s\mu^+\mu^-$ transition in R-parity violating interactions.

g-2 with RPV has a long history, see, e.g.Kim, Kyae and Lee, PLB 2001

We [ALTHANNISHOFER+DEV+AS] acexamining+up date in light of current flavor amomalies WORK IN Progress

CONSTRAINTS: TIGHTENING EXPT'S NOOSE AGAINST SPECIFIC MODELS

Table 13-6. Model-dependent effects of new physics in various processes.

	CP Violation			D^0 – \overline{D}^0
Model	$B_d^0 - \overline{B}_d^0$ Mixing	Decay Ampl.	Rare Decays	Mixing
MSSM	O(20%) SM	No Effect	$B \to X_s \gamma$ – yes	No Effect
	Same Phase		$B o X_s l^+ l^-$ – no	
SUSY – Alignment	O(20%) SM	$\mathcal{O}(1)$	Small Effect	Big Effect
	New Phases			
SUSY -	$\mathcal{O}(20\%)$ SM	$\mathcal{O}(1)$	No Effect	No Effect
Approx. Universality	New Phases			
R-Parity Violation	Can Do	Everything	Except Make	Coffee
MHDM	\sim SM/New Phases	Suppressed	$B \to X_s \gamma, B \to X_s \tau \tau$	Big Effect
2HDM	\sim SM/Same Phase	Suppressed	$B \to X_s \gamma$	No Effect
Quark Singlets	Yes/New Phases	Yes	Saturates Limits	Q = 2/3
Fourth Generation	~ SM/New Phases	Yes	Saturates Limits	Big Effect
$LRM - V_L = V_R$	No Effect	No Effect	$B \to X_s \gamma, B \to X_s l^+ l^-$	No Effect
$-V_L \neq V_R$	Big/New Phases	Yes	$B \to X_s \gamma, B \to X_s l^+ l^-$	No Effect
DEWSB	Big/Same Phase	No Effect	$B \to X_s \ell \ell, B \to X - s \nu \overline{\nu}$	Big Effect

though in many cases further data may limit the available parameter space. In the more exciting eventuality that the results are not consistent with Standard Model predictions, the full pattern of the discrepancies both in rare decays and in CP-violating effects will help point to the preferred extension, and possibly rule out others. In either case there is much to be learned.

constraints

• Direct searches via $pp \to \tilde{b} \tilde{b} \to \tau^+ \tau^- t \bar{t}$

Indirect constraints considered due $B=>\tau v$; $\pi \tau v$; $\pi(K) v v....$ Also $B_C =>\tau v....$

To a/c (within 1σ) of expt for RD(*) needs largish $\lambda'333~1-2~$ range with quite heavy sbottoms but such large couplings develop landau pole below GUT scale.We require couplings stay perturbative below GUT so with $\lambda'333<^{1}$,

- ⇒ TAKE HOME: This version of RPV is actually (surprisingly) well constrained
- ⇒ With improved measurements RD(*) in RPV3 may be difficult

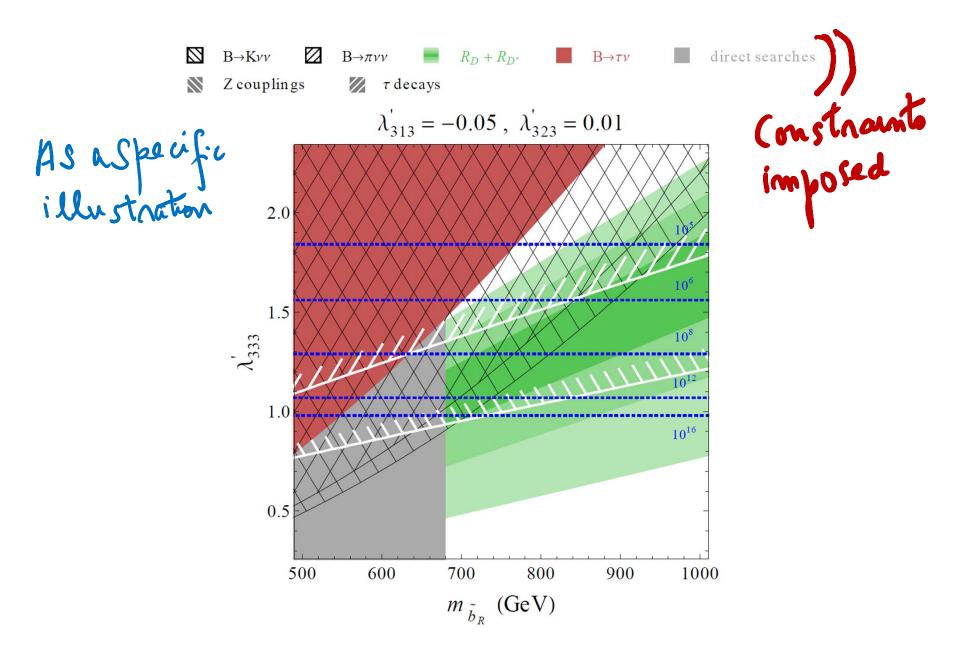


FIG. 3. RPV parameter space satisfying the $R_{D^{(*)}}$ anomaly and other relevant constraints. Lyon; A Soni(BNL-HET)





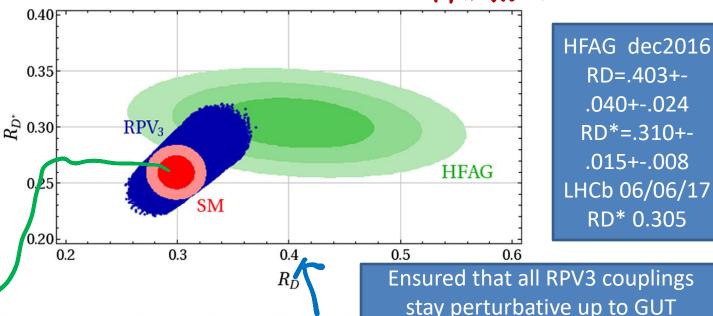


FIG. 4. The SM predictions (red), experimental world average (green), and accessible values in our RPV-SUSY scenario (blue) in the R_D vs. R_{D^*} plane. For the SM, bearing in mind recent works [17,20,22] we are taking $(R_D^{\rm SM}, R_{D^*}^{\rm SM}) = (0.299 \pm 0.011, 0.260 \pm 0.010)$.

all constraints......RPV(blue) region obtained by scanning with sbottom mass 680-1000Gev, 0<λ333<2;|λ323|<0.1;|λ313|<0.3

••••

IN CLOSING

Contrarian/Complementary view

- flavor physics is actually hanging by perhaps the weakest link i.e. a single CP-phase endowed by the 3g –SM.
- In many ways this is a contrarian (or complementary)
 point of view, in sharp contrast to the overwhelming
 majority following the naturalness lamp post via Higgs
 radiative stability.

In this context it is useful to stress

Importance of the "IF": score card

- Beta decay => Gf => W....
- Huge suppression of KL => mu mu; miniscule ΔmK=> charm
- KL =>2 pi but very rarely; mostly to 3pi =>CP violation
 => 3 families
- Largish Bd –mixing => large top mass
- etc.....
- => extremely unwise to put all eggs in HEF
- Complementary info from IF can be a crucial guide for pointing to new thresholds as well as provide important clues to the nature of the signals there from

Summary and Outlook..[p1 of 2]

- Hints of LUV [from 3 B experiments] in sl B decays are interesting though not yet
 compelling despite the claimed 4.1 sigma; FCNC mu/e LUV indications in B=K (and K*) l l
 are also extremely interesting and may well be related to LUV in sl B....If so similar mu/e
 LUV should be seen in FCNC channels [e.g] in baryonic B decays as well as a multitude of
 other channels.
- Though not yet compelling; OTOH no good reason that experimental hints are wrong.
 While we are accustomed to respecting LU the sensitivities of the current experiments are better than the past.
- This is not a violation of SR or GR, but only of accepted dogma; experimental results on LUV therefore deserve serious attention.
- That is my interest, try get to the bottom of it all AFAP using on & off the lattice tools
- Accurate SM predictions are usually very challenging and demanding but are extremely important; BSM scenarios are easy to cook up but 1/N^N' of being correct with N, N' >1000
- However, in the context of BSM note also the truly impressive constraints on models that past decade or so studies have resulted [as illustrated by RPV3].
- More data from LHCb from Run 2 < 1 year should help and further ~2 years down Belle II should start to help more....On the theory side, of critical importance is FF for B=>D*...should start to get results in ~ 6 months to 1 year RBC/UKQCD as well as others
- FCNC mu/e LUV so far only LHCb data shows the effect...confirmation from other experiments would be extremely helpful.

Summary and Outlook ...p.2

- It may well be that BNL's observed g-2 signals of possible NP were just a precursor to these observations in B decays.
- Lattice progress in g-2 by RBC-UKQCD as well as global efforts are impressive ...But needs to reduce errors further by ~X4...Expect next reduction X2 in a year or so
- ε': RBC-UKQCD should be able to appreciably improve their 2015 result of ~2.1 sigma tension, in <6 months
- Personally, this is the ~36th year of trying to tame this really wild beast; so it'd be welcome indeed.
- There is now an exciting and may be even a revolutionary possibility that one or more of these avenues will show significant departure from SM in ~ 1-2 years

XTRAS

Lattice computation of the decay constants of B and D mesons

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PHYSICAL REVIEW D

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Semileptonic decays on the lattice: The exclusive 0⁻ to 0⁻ case

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Lattice study of semileptonic decays of charm mesons into vector mesons

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We present our lattice calculation of the semileptonic form factors for the decays $D \to K^*$, $D_s \to \phi$, and $D \to \rho$ using Wilson fermions on a $24^3 \times 39$ lattice at $\beta = 6.0$ with 8 quenched configurations. For $D \to K^*$, we find for the ratio of axial form factors $A_2(0)/A_1(0) = 0.70 \pm 0.16 \pm 0.00$. Results for other form factors and ratios are also given.

Department of Player anomalies; Lyon, A sound process.

SU(3) flavor breaking in hadronic matrix elements for B-B oscillations

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