



Universität
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A model for **B** anomalies and **SM** flavor hierarchies

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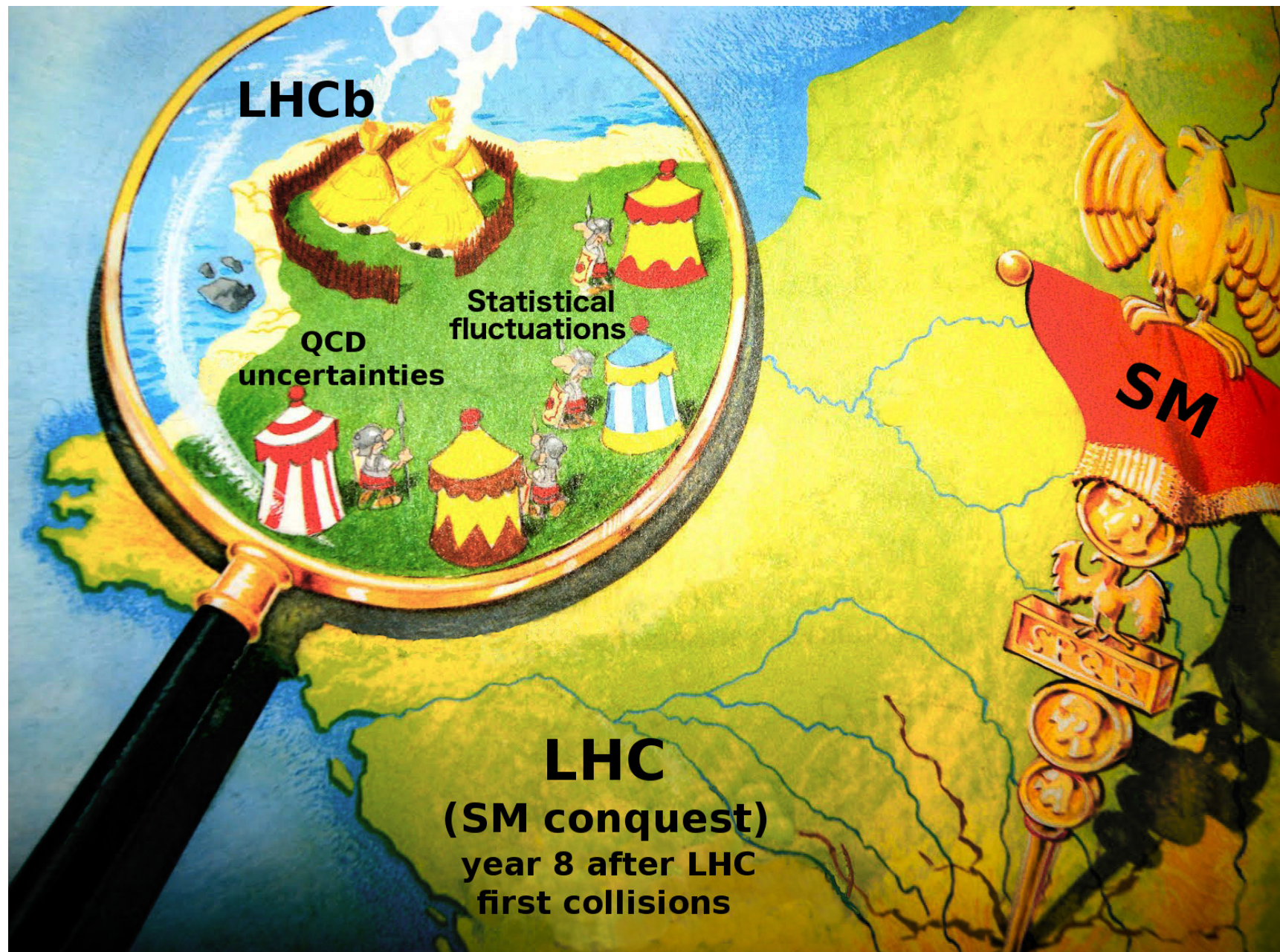
Based on **arXiv:1712.01368** and ongoing work

In collaboration with **M. Bordone**, **C. Cornella** and **G. Isidori**

LIO international conference on Flavour Physics: “From Flavour to New Physics observables”

The LHC landscape

The year is 8 after the LHC first collisions. Experimental data is entirely SM-like. Well, not entirely! The LHCb Collaboration still holds out against the SM. And life is not easy for the SM there...



The challenge: combined explanation of the anomalies

If taken together, the B anomalies are possibly the **largest coherent set** of deviations from the SM we have ever seen...

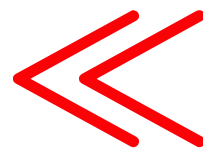
... So let me (aggressively) assume that the anomalies (both!) are genuine hints of NP. Can we make some sense out of them?

Intriguingly, they follow a very peculiar structure

$$b \rightarrow s \ell \ell$$

$$3_Q \rightarrow 2_Q 2_L 2_L$$

~25% of a SM **loop** effect



$$b \rightarrow c \tau \nu$$

$$3_Q \rightarrow 2_Q 3_L 3_L$$

~20% of a SM **tree-level** effect

The only source of **lepton flavor universality violation** in the SM (Yukawas) follow a similar trend: $y_e \ll y_\mu \ll y_\tau$ Are the anomalies connected to them?

The U(2) flavor symmetry

The SM Yukawas respect an approximate U(2) symmetry

Barbieri et al. 1105.2296

$$M_{u,d} \sim \begin{array}{|c|c|} \hline \text{light} & \text{light} \\ \hline \text{light} & \text{heavy} \\ \hline \end{array}$$

$$V_{\text{CKM}} \sim \begin{array}{|c|c|c|} \hline \text{heavy} & \text{light} & \text{light} \\ \hline \text{light} & \text{heavy} & \text{light} \\ \hline \text{light} & \text{light} & \text{heavy} \\ \hline \end{array}$$

$$\text{U}(2)_{\text{q}} \times \text{U}(2)_{\text{u}} \times \text{U}(2)_{\text{d}}$$

$$\psi = (\psi_1 \ \psi_2 \ \psi_3)$$

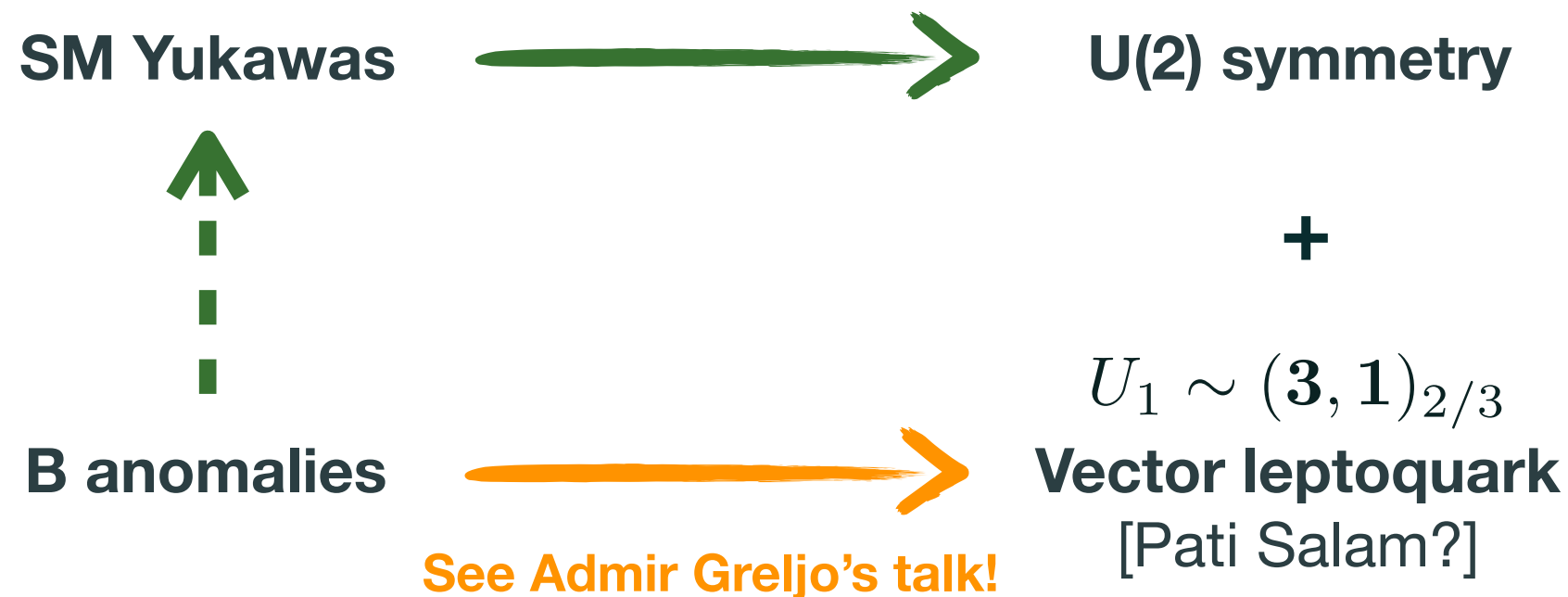
$$Y_{u,d} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \xrightarrow{\text{green}} \begin{pmatrix} 0 & V \\ 0 & 1 \end{pmatrix} \xrightarrow{\text{orange}} \begin{pmatrix} \Delta & V \\ 0 & 1 \end{pmatrix} \quad \begin{array}{l} |V| \sim |V_{ts}| \\ |\Delta| \sim y_c \end{array}$$

Unbroken symmetry

Leading breaking

Final breaking

- Assuming a single leading breaking ensures an effective protection of FCNCs
[**SM-like mixing among light & 3rd generations** \longrightarrow consistent with CKM fits]
- Large NP effects in 3rd generation, light-generation effects controlled by the breaking



Analysis in terms of dynamical (simplified) models select the **vector LQ as the only single-mediator possibility** for a combined explanation of the anomalies
[but of course other solutions with more mediators are available]

Gauged vector LQ: low-scale unification?

The vector-leptoquark solution points to Pati-Salam unification

$$\mathbf{PS} \equiv \mathbf{SU}(4) \times \mathbf{SU}(2)_L \times \mathbf{SU}(2)_R \qquad \Psi_{L,R} = \begin{pmatrix} Q_{L,R}^1 \\ Q_{L,R}^2 \\ Q_{L,R}^3 \\ L_{L,R} \end{pmatrix}$$

Pati, Salam, Phys. Rev. D10 (1974) 275

[**Lepton number as the 4th “color”**]

✓ $\mathbf{SU}(4)$ is the smallest group containing the required vector LQ [$U_1 \sim (\mathbf{3}, \mathbf{1})_{2/3}$]

✓ No proton decay (protected by symmetry)

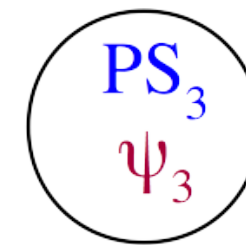
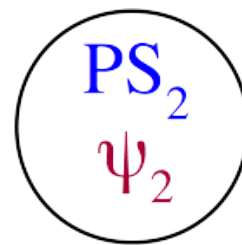
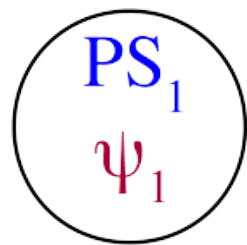
✗ The (flavor blind) Pati-Salam model cannot work

→ The bounds from $K_L \rightarrow \mu e$ and $D - \bar{D}$ lift the LQ mass to 100 TeV

A very interesting proposal [[Di Luzio et al., 1708.08450](#)] uses vector-like fermions to circumvent this problem [[See talks by Admir Greljo and Luca Di Luzio](#)]

At high energies each family is charged under an independent gauge group
(**gauge bosons carry flavor!**)

$$\mathbf{PS}^3 \equiv [\mathbf{SU}(4) \times \mathbf{SU}(2)_L \times \mathbf{SU}(2)_R]^3$$



Unification of quarks and leptons
[natural explanation for $U(1)_Y$ charges]

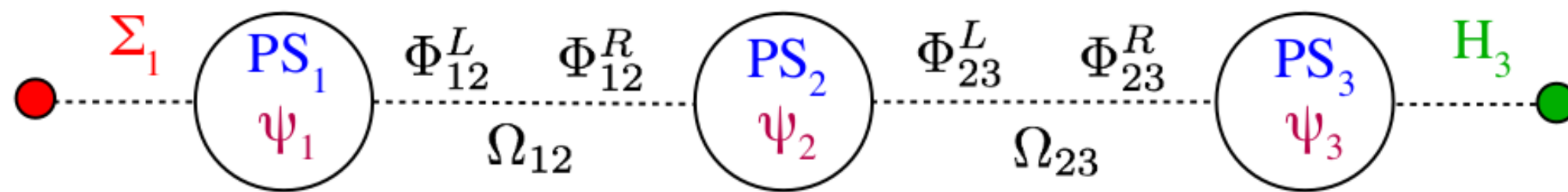
“De-unification” of the gauge symmetry
(= flavor deconstruction)



Available now at low scales!

At high energies each family is charged under an independent gauge group
(**gauge bosons carry flavor!**)

$$\text{PS}^3 \equiv [\text{SU}(4) \times \text{SU}(2)_L \times \text{SU}(2)_R]^3$$



High-scale [$\sim 10^3$ TeV]

“vertical breaking”:

[$\text{PS}_1 \rightarrow \text{SM}_1$]

Link fields



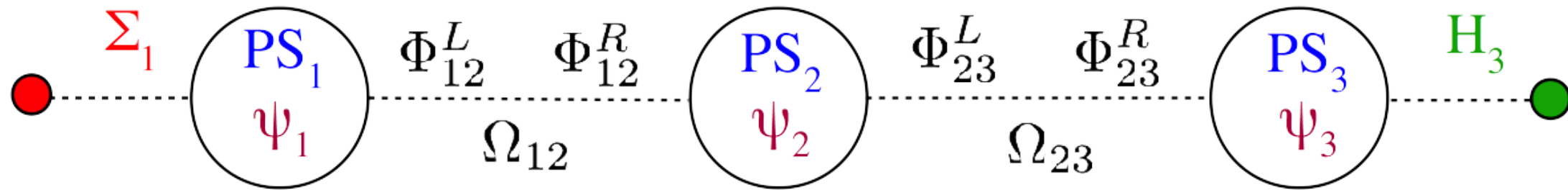
Low-scale [EW]

“vertical breaking”:

[$\text{SM}_3 \rightarrow \text{QED}_3$]

Breaking to the diagonal SM group via “link” fields [$\Omega_{ij}, \Phi_{ij}^{L,R}$],
also responsible of the Yukawa hierarchies

Symmetry breaking pattern

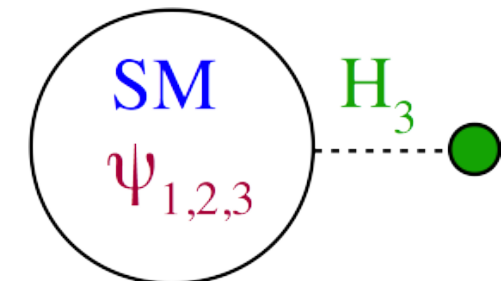
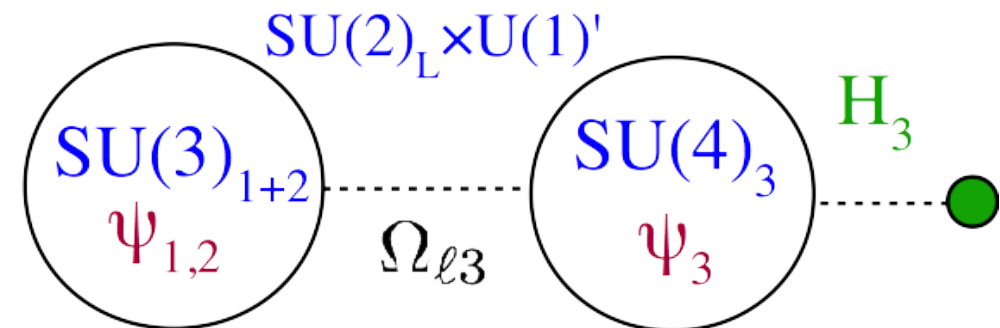
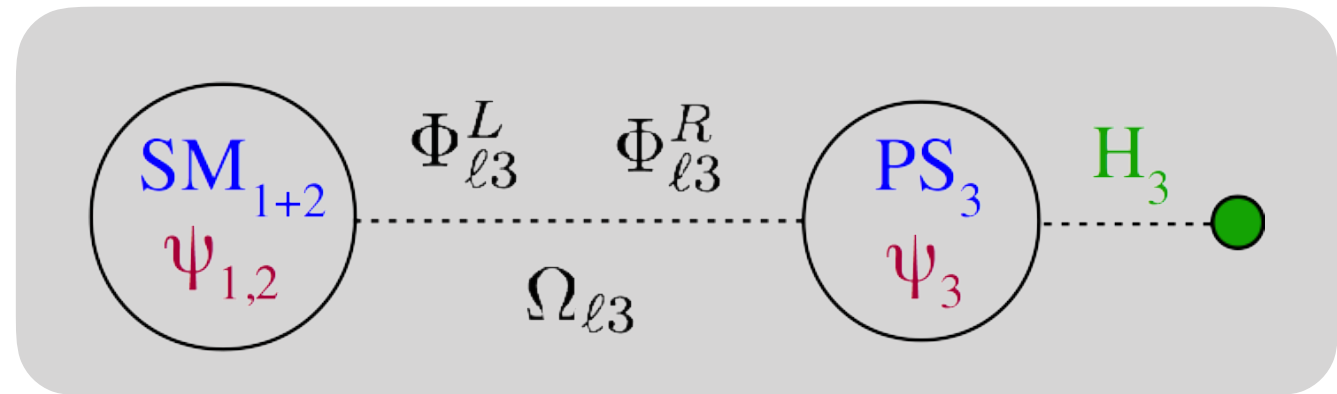


Accidental $U(2)^5$ symmetry

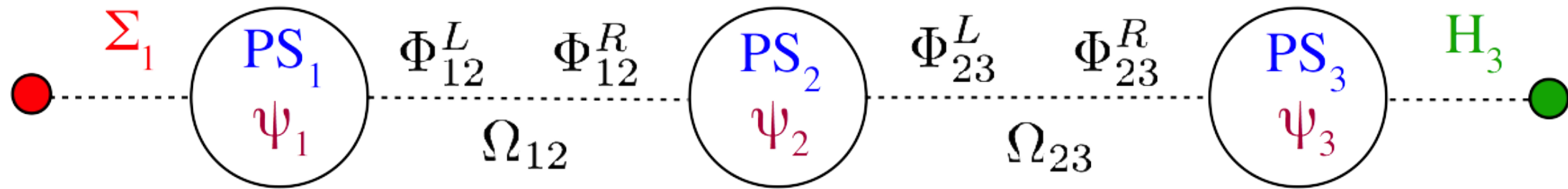


$$Y_f = \begin{bmatrix} \mathbf{y}_{11} & \mathbf{y}_{13} \\ \mathbf{y}_{33} \end{bmatrix} \frac{\langle \Omega_{\ell 3} \rangle}{\Lambda_{23}} \frac{\langle \Phi_{\ell 3}^L \rangle \langle \Phi_{\ell 3}^R \rangle}{\Lambda_{23}^2}$$

Higher dimensional operators act as spurions (i.e. small breakings) of the $U(2)$ symmetry



Symmetry breaking pattern

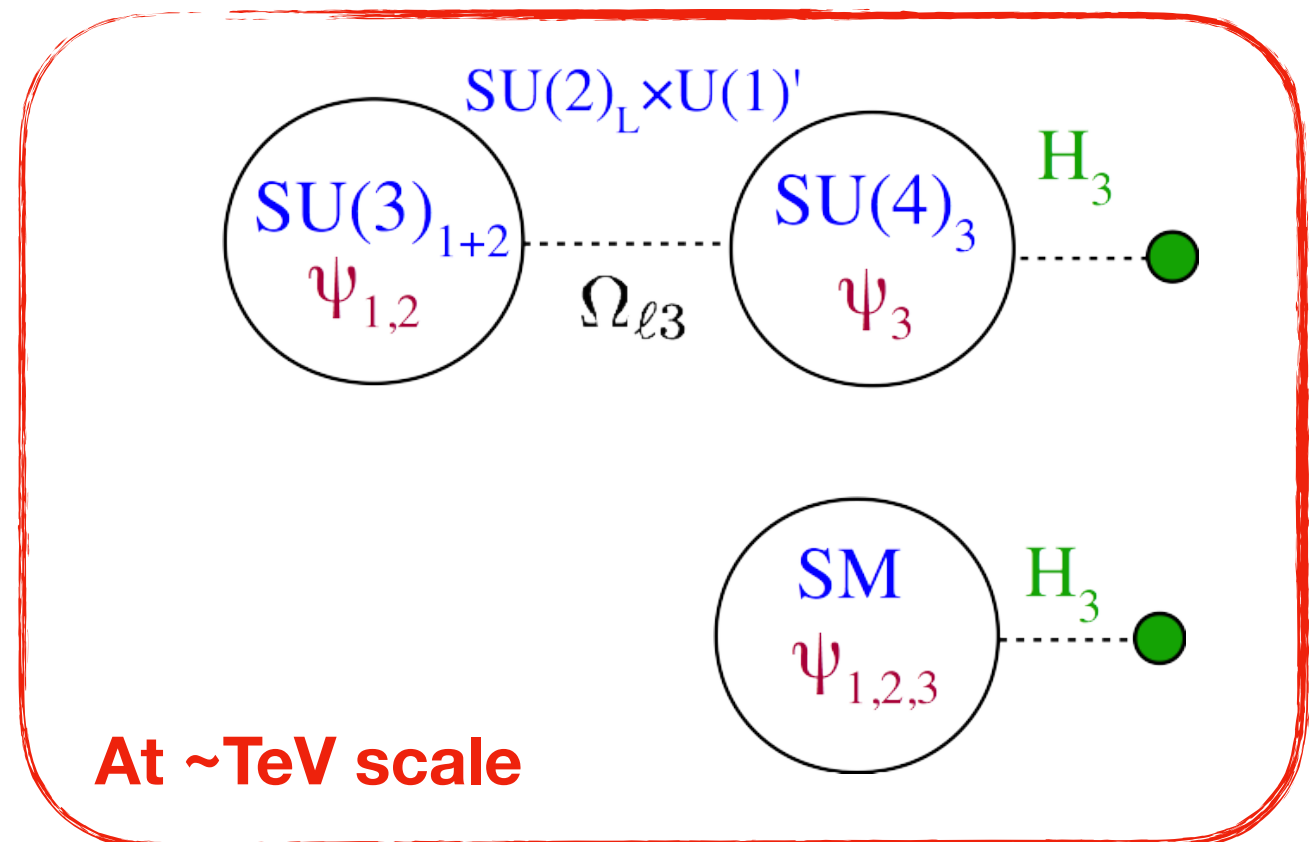
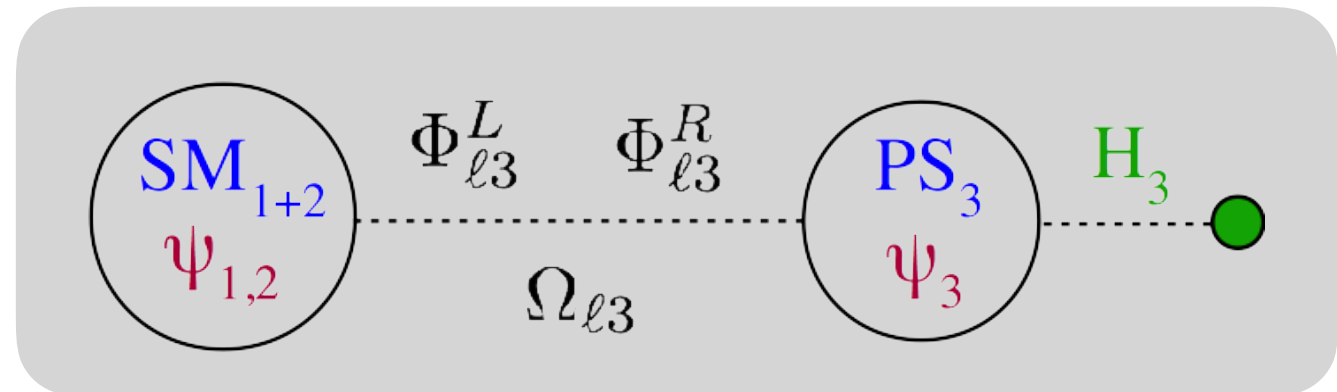


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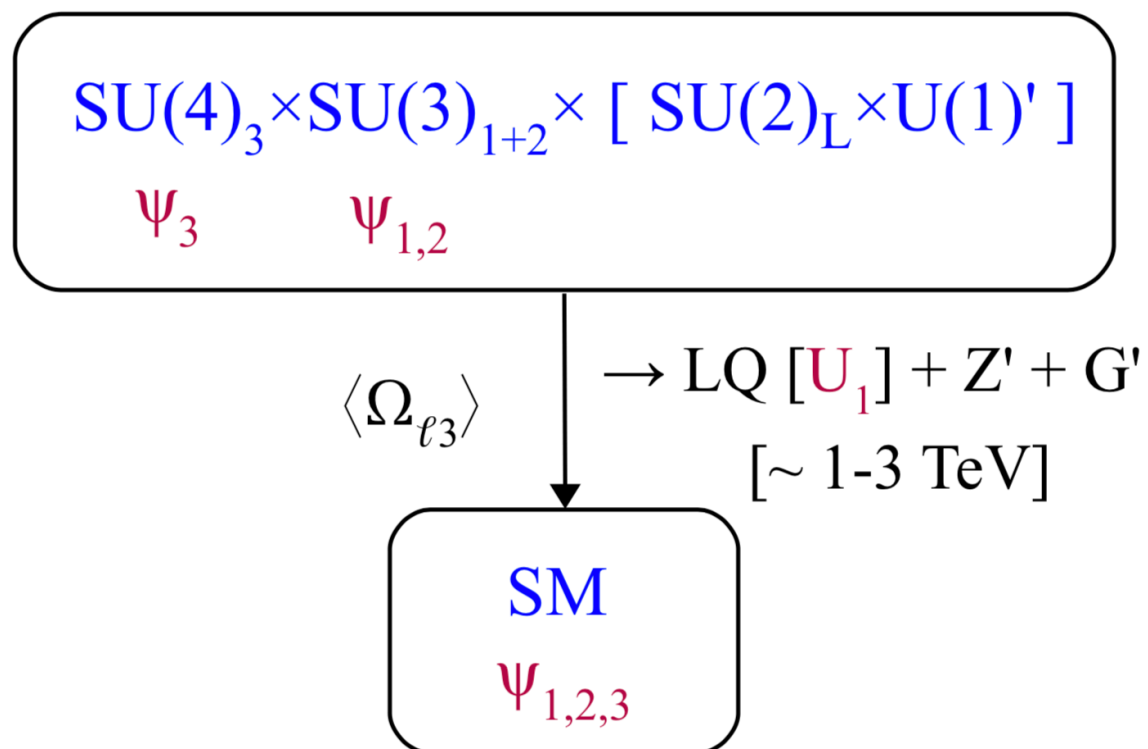


At $\sim \text{TeV}$ scale

Low-energy remnants

Similar constructions in:

- Di Luzio, Greljo, Nardecchia, 1708.08450
- Diaz, Schmaltz, Zhong, 1706.05033
- Georgi, Nakai, 1606.05865



Despite the apparent complexity, the model is highly constrained at $\sim \text{TeV}$ scale

- Uninvited guests: unavoidable Z' and G' , with masses close to U_1
- Key difference to other existing models
 \rightarrow unsuppressed $b_R - \tau_R$ LQ couplings
[Very important pheno implications!]

Flavor structure

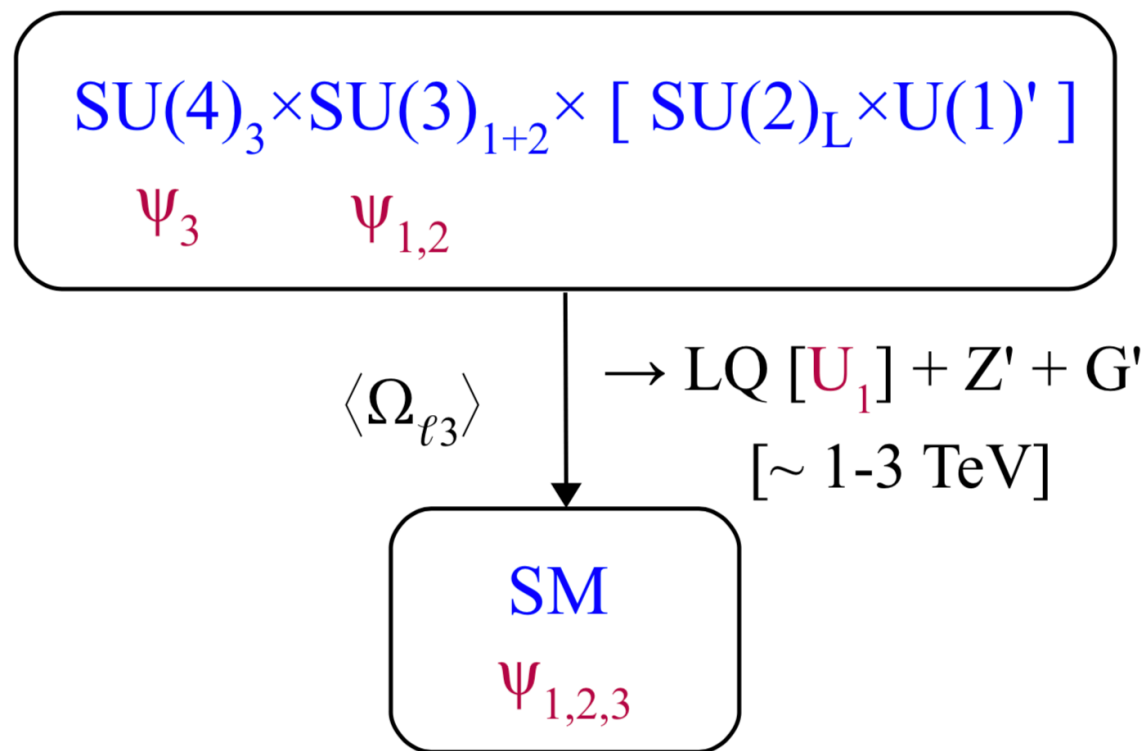
[$g_4 \simeq \mathbf{3}$ enhances NP in $R(D^{(*)})$ and suppresses dangerous couplings]

$$U_1: g_4 \text{diag}(\epsilon'_U, \epsilon_U, 1) \quad [\epsilon_i^{(')} \text{ from d=6 op.}]$$

$$G': g_c \text{diag} \left(-\frac{g_3}{g_4} + \epsilon_q, -\frac{g_3}{g_4} + \epsilon_q, \frac{g_4}{g_3} \right) \quad + \quad \text{(small) flavor rotations controlled by the } U(2) \text{ symmetry}$$

$$Z': g_Y \text{diag} \left(-\frac{g_1}{g_4} + \epsilon_{q,\ell'}, -\frac{g_1}{g_4} + \epsilon_{q,\ell}, \frac{g_4}{g_1} \right)$$

To the arena!... Low-energy phenomenology



- Uninvited guests: unavoidable Z' and G' , with masses close to U_1
- Unsuppressed $b_R - \tau_R$ LQ couplings
- Flavor structure controlled by the $U(2)$ symmetry [and d=6 ops. ($U(2)$ spurions)]

Any candidate to simultaneously explain the anomalies should “fight in two arenas”:

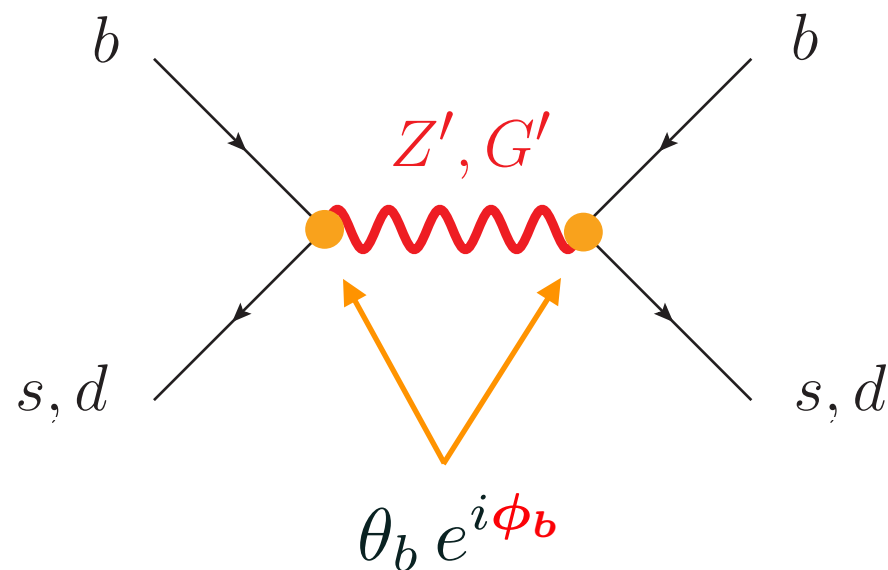


- **Other low energy data**
 - $\Delta F = 2$ transitions
 - τ decays (LFUV and LFV)
 - Other semileptonic processes
- **High- p_T searches at LHC**

and life is not easy for NP models there....

$\Delta F = 2$: one phase to save them all?

Current lattice data hint to a deficit in the experiment w.r.t. SM prediction in $B_{s,d} - \bar{B}_{s,d}$ [Fermilab/MILC 2016 [1602.03560]: SM prediction $1.8 \sigma (B_d)$ and $1.1 \sigma (B_s)$ **above** experiment]



CP violating NP can account for the deficit!
[Di Luzio et al., 1712.06572]

Current data

$$\phi_b \simeq \pi/2$$

$$|\theta_b| = \mathcal{O}(10\%) |V_{ts}|$$

U(2) symmetry

$$\phi_b \text{ free}$$

$$\theta_b = \mathcal{O}(V_{ts})$$

Still early to draw conclusions but it is interesting that the model can “naturally” explain the deficit

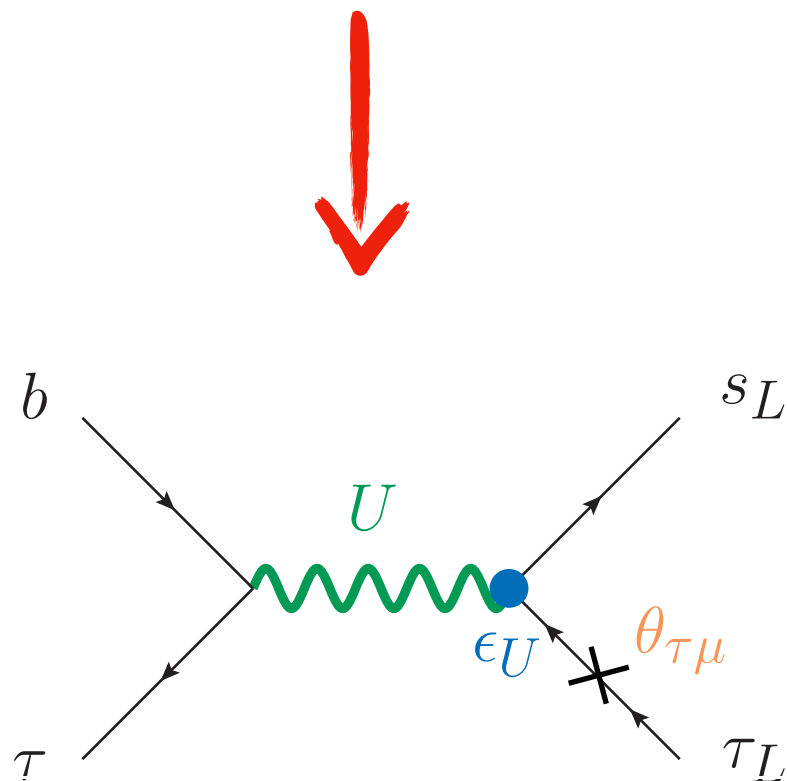
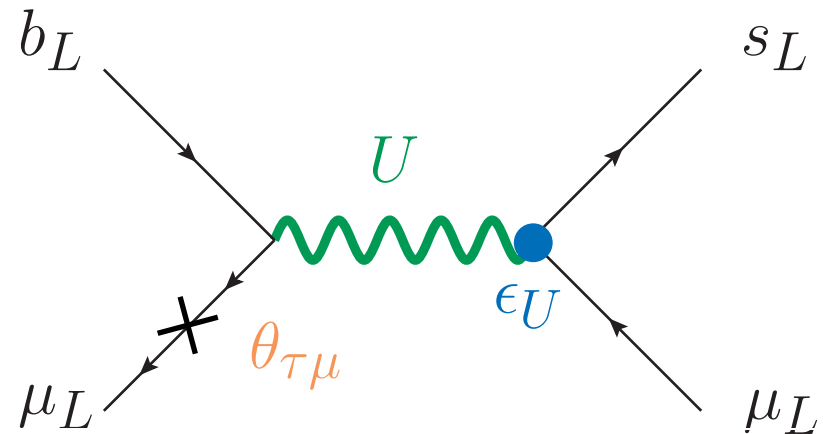
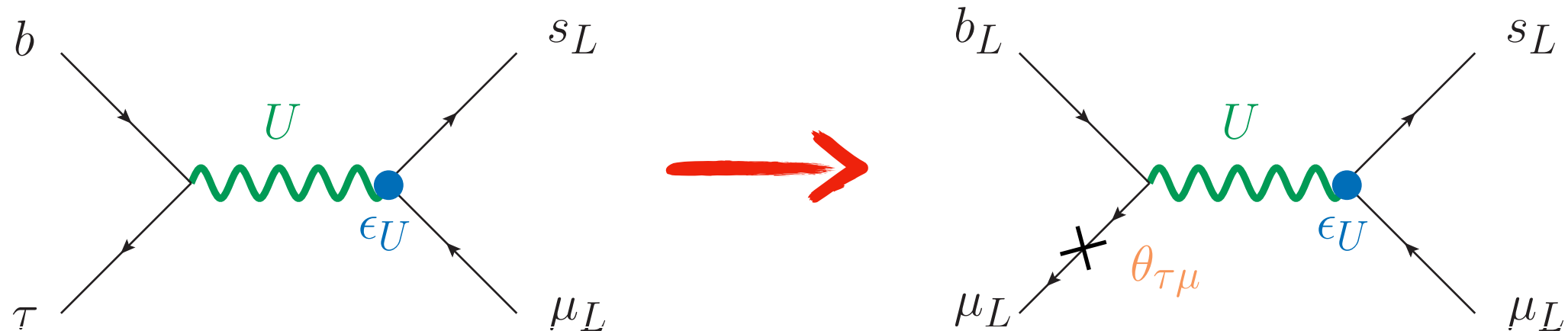


Possible **CP violation** effects in $b \rightarrow s, d$ transitions!
[Currently under investigation]

Other $\Delta F = 2$ transitions: $K - \bar{K}$, $D - \bar{D}$ also **under control!**

Explanation of the $b \rightarrow s\ell\ell$ anomalies

The dominant contribution to $b \rightarrow s\ell\ell$ transitions is mediated by the leptoquark



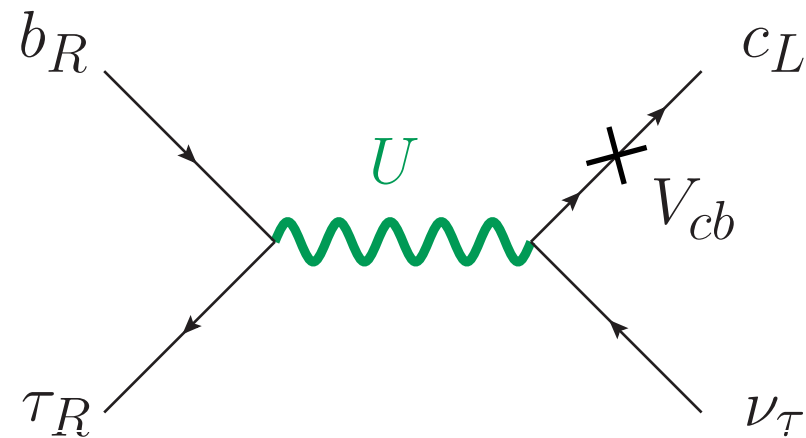
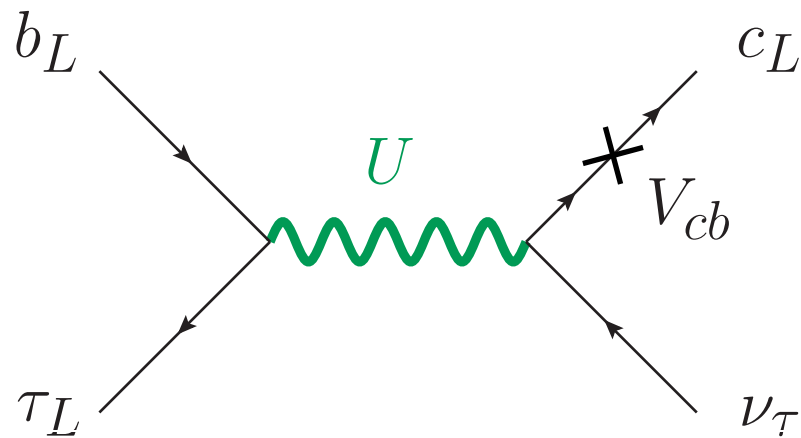
$$\theta_{\tau\mu} \sim 0.2$$

$$b \rightarrow s\tau\tau \sim b \rightarrow s\mu\mu$$

$$b \rightarrow s\tau\mu \sim \mathcal{O}(5) \times b \rightarrow s\mu\mu$$

★ ★ Key difference w.r.t the “standard” solutions

Explanation of the $b \rightarrow c\tau\nu$ anomalies



$C_{\nu edu}^{V,LL}$

$C_{\nu edu}^{S,RL}$



Key difference w.r.t the
“standard” solutions

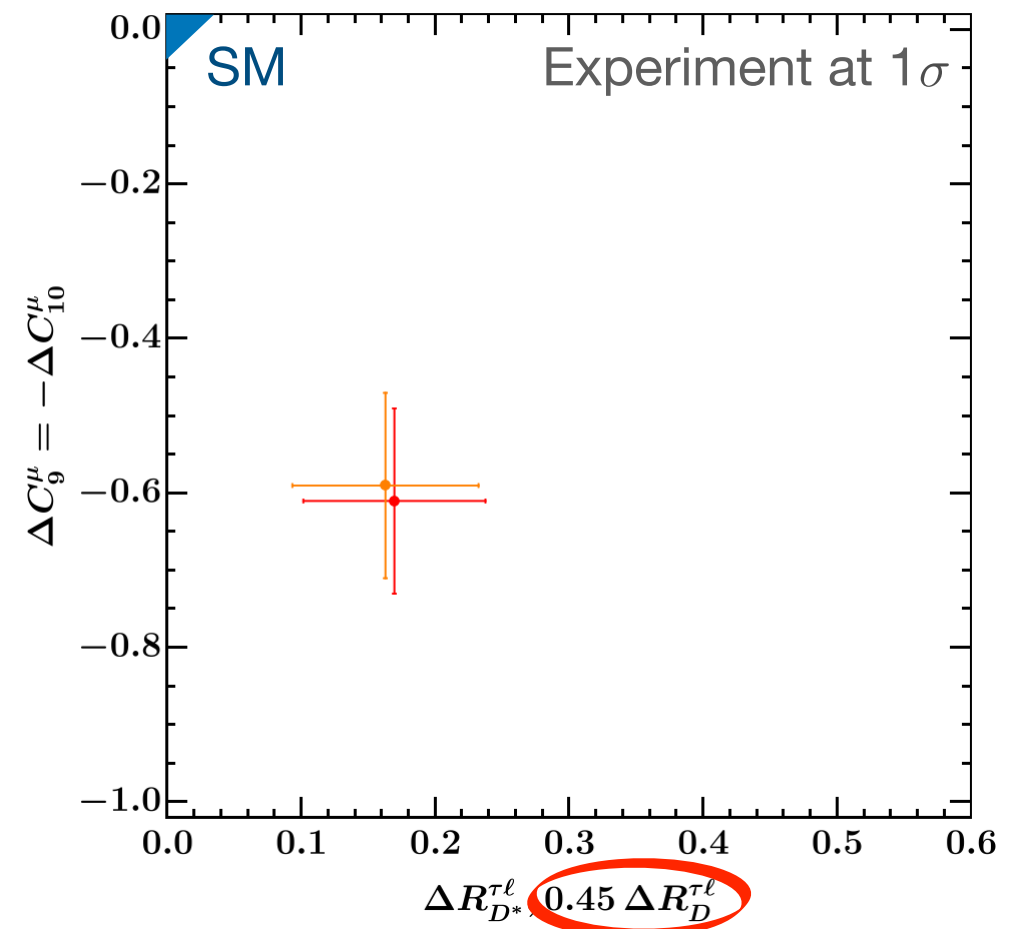
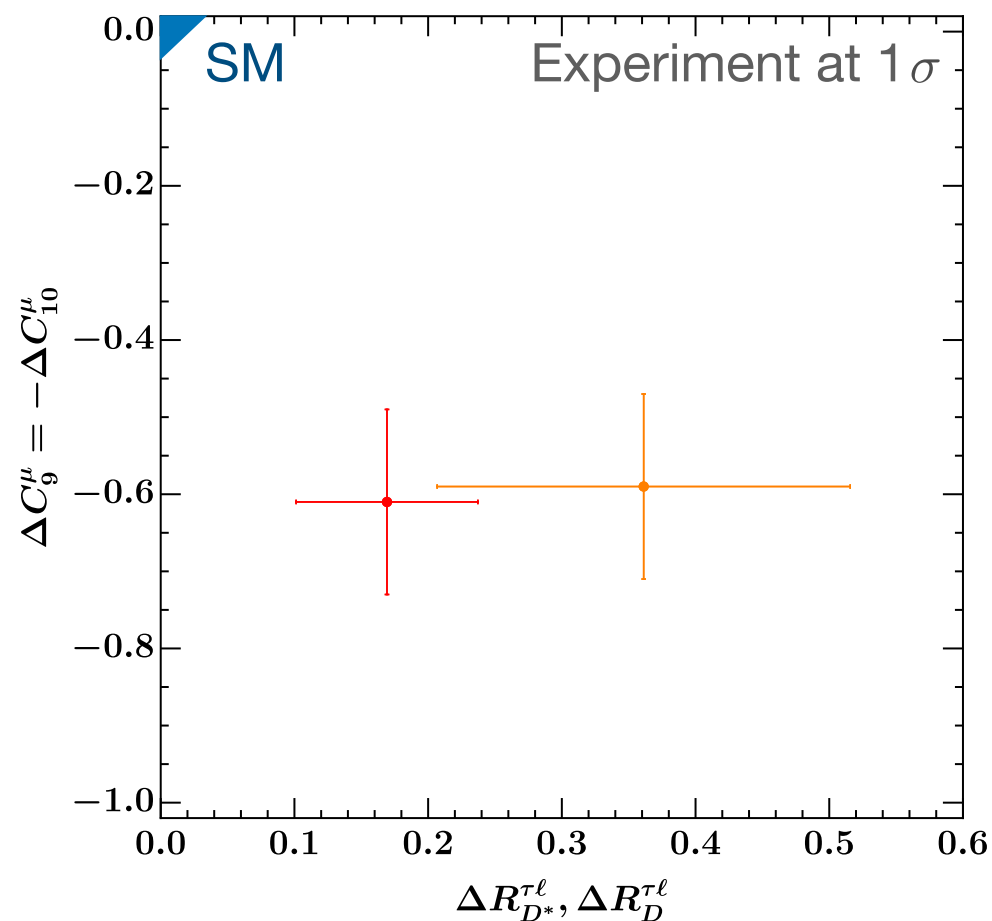
$$\Delta R_{D^*} = R_{D^*}/R_{D^*}^{\text{SM}} - 1 \approx 2(1 + 0.12)C_U$$

$$\Delta R_D = R_D/R_D^{\text{SM}} - 1 \approx 2(1 + 1.5)C_U$$

$$C_U = \frac{g_U^2 v^2}{4M_U^2}$$

Contributions from the **vector/scalar** operator in **green/blue**

The “0.45 rule” for $R(D)/R(D^*)$

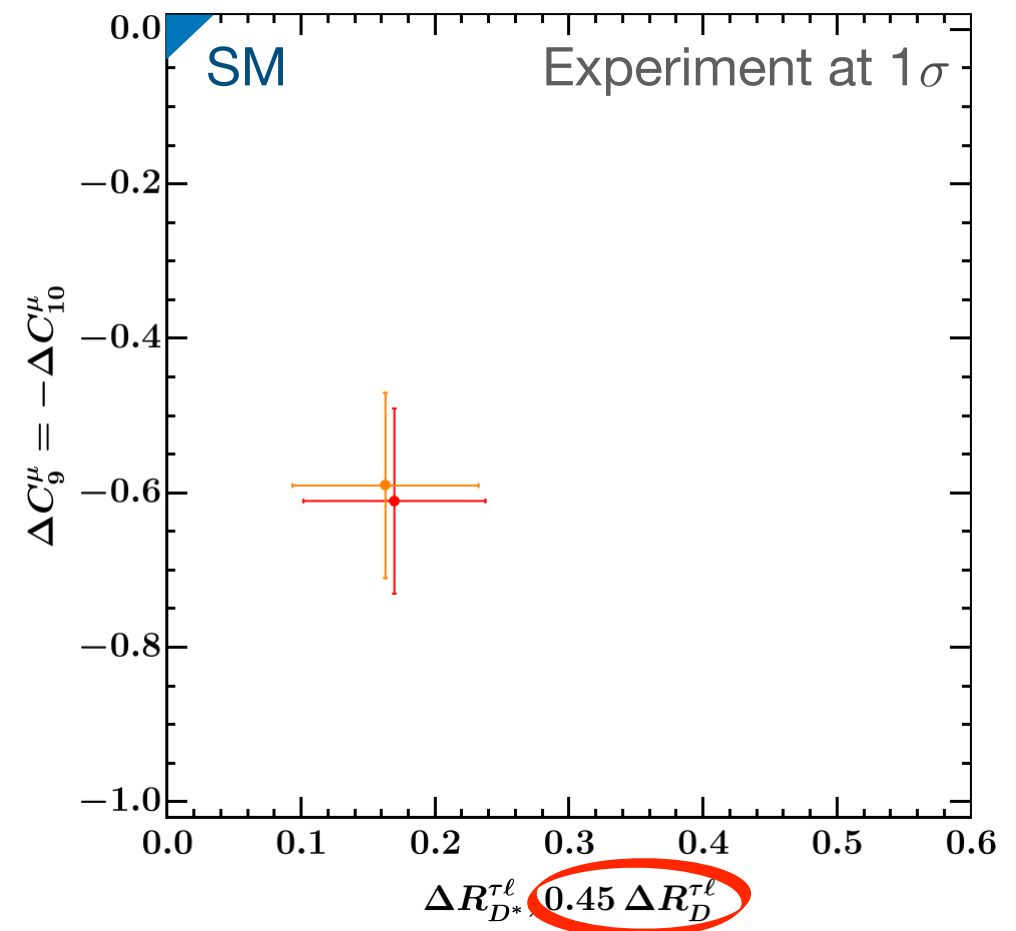
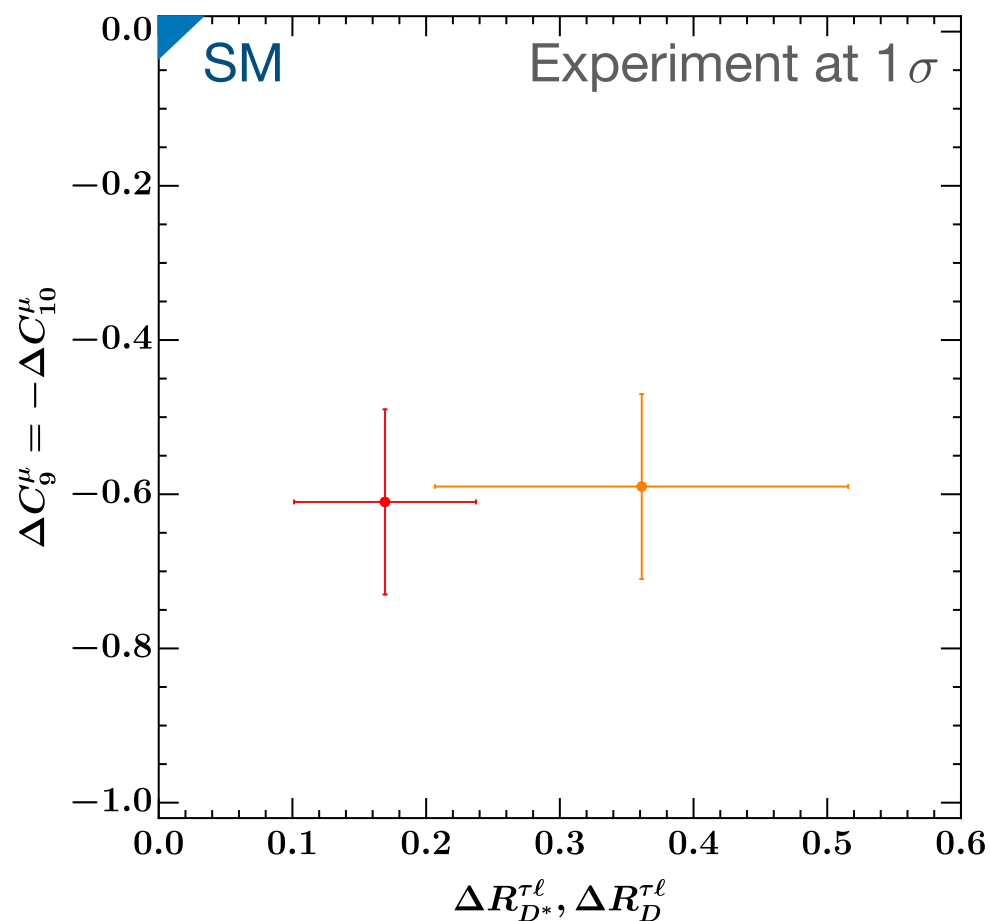


$$\left. \begin{aligned} \Delta R_{D^*} &= R_{D^*}/R_{D^*}^{\text{SM}} - 1 \approx 2(1 + 0.12)C_U \\ \Delta R_D &= R_D/R_D^{\text{SM}} - 1 \approx 2(1 + 1.5)C_U \end{aligned} \right\} \Rightarrow \boxed{\frac{\Delta R_{D^*}}{\Delta R_D} \simeq 0.45} \quad C_U = \frac{g_U^2 v^2}{4M_U^2}$$

The required NP scale can be higher!!

★ ★ Key difference w.r.t other solutions

The “0.45 rule” for $R(D)/R(D^*)$

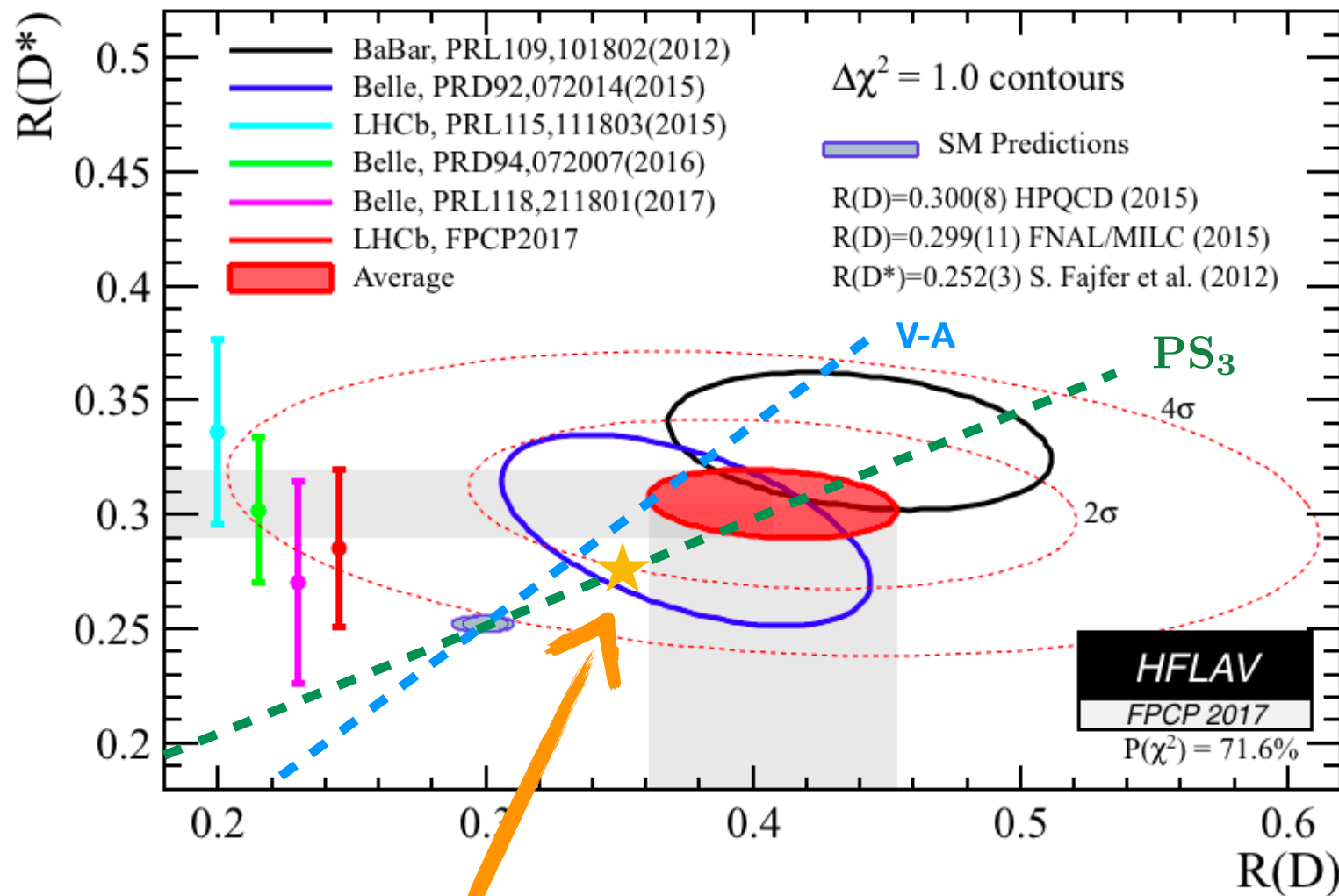


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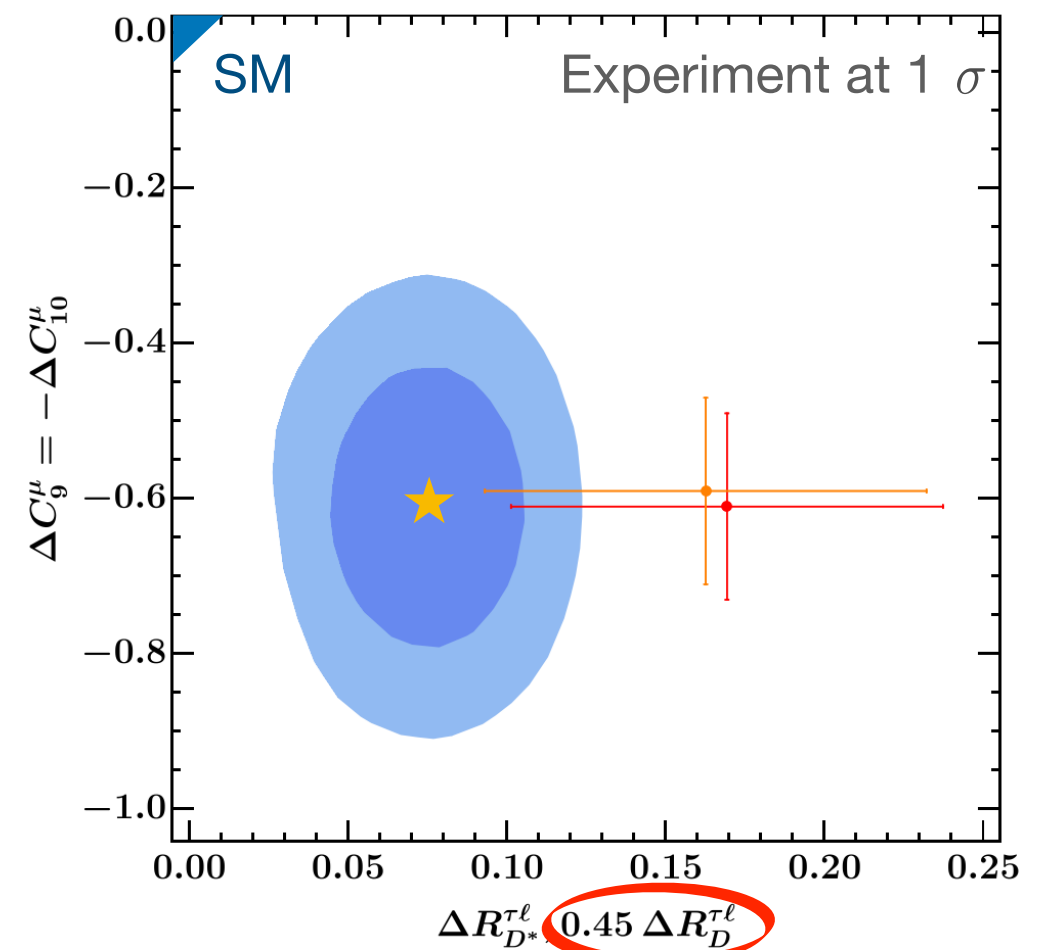
Possible large enhancements in $B \rightarrow \tau\nu$
[eagerly waiting for Belle II]

★ ★ Key difference w.r.t
other solutions

Low-energy fit results



Fit to (mostly) B and τ physics



Dominated by 5 parameters
(NP scale, θ_b , ϕ_b , $\theta_{\tau\mu}$, ϵ_U)

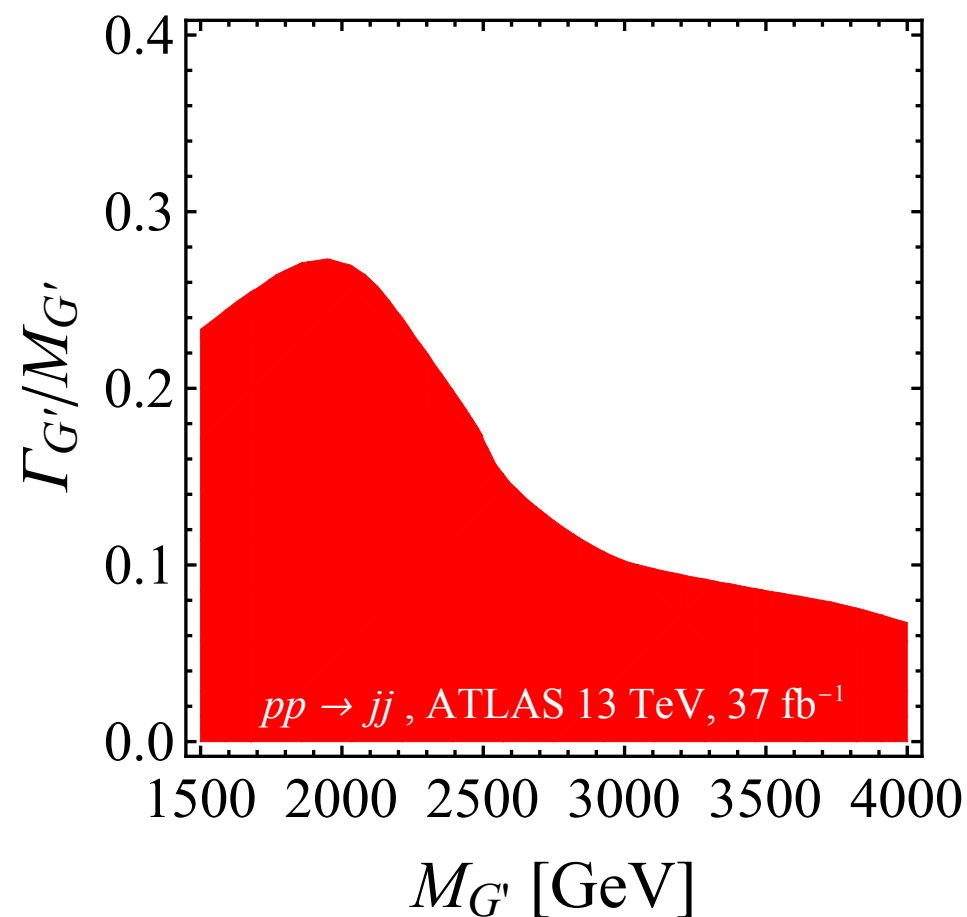
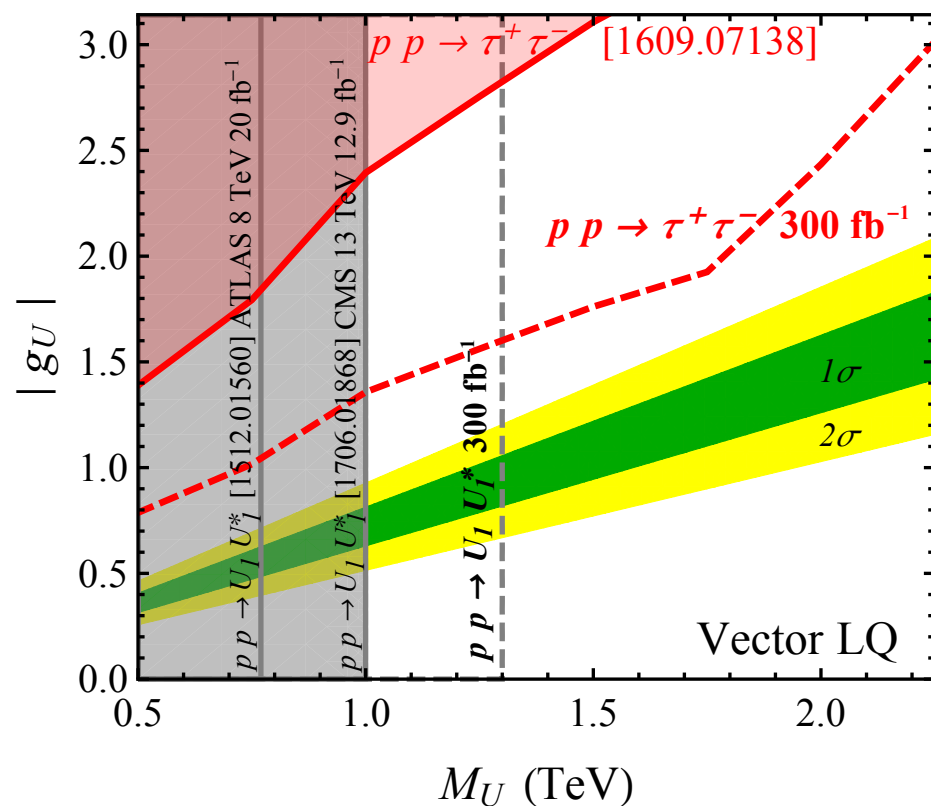
The fit to low-energy data is very good!
[although slightly smaller NP effects in $R(D)$ - $R(D^*)$]

High-pT searches

General characteristics of the new exotic states:

- Flavor non-universal couplings, with much **stronger couplings to the 3rd generation**
- **Large widths** (specially for G' and Z')

Coloron exclusion for the “4321” model in [Di Luzio et al., 1708.08450] (Similar conclusions apply for PS^3)



See talk by Luca di Luzio!

[Di Luzio, JF, Greljo, Nardecchia, Renner, in preparation]

Conclusions

It is possible to find **well-motivated** combined explanations of the B anomalies, compatible with current data in both low- and high-energies

The model I presented, a **Pati-Salam flavor deconstruction**, predicts several characteristic smoking-gun signatures that differentiate it from other solutions

- $\Delta R_{D^*} / \Delta R_D \simeq 0.45$
- NP in $b \rightarrow s\tau\tau \sim b \rightarrow s\mu\mu$ and $b \rightarrow s\tau\mu \sim \mathcal{O}(5) \times b \rightarrow s\mu\mu$
- Possibility of a deficit in $\Delta B = 2$ & CP violation involving the 3rd quark family
- Possible large enhancements in $B \rightarrow \tau\nu$

If the anomalies are really pointing to NP, **new experimental indications** (both in high- p_T and at low energies) will show up soon in several observables

Exciting times ahead of us. Let's have fun



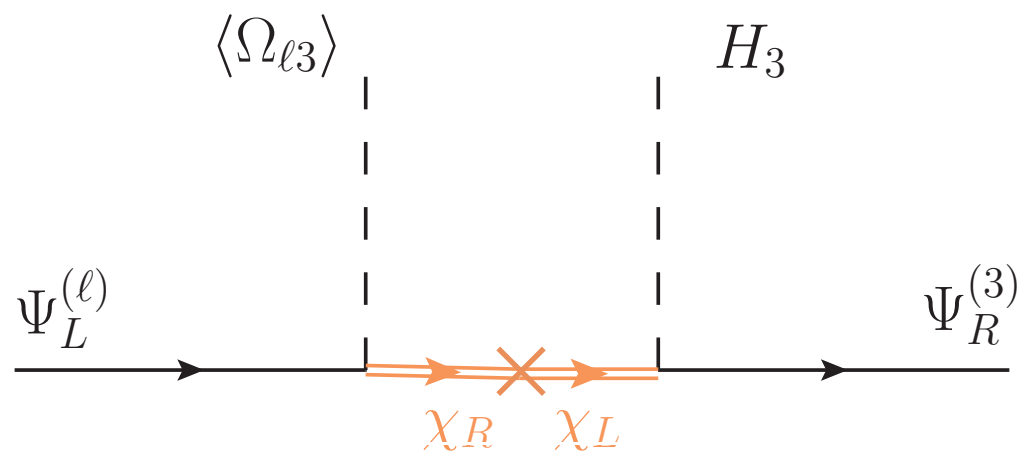
The end

Backup slides

Generation of the effective Yukawas

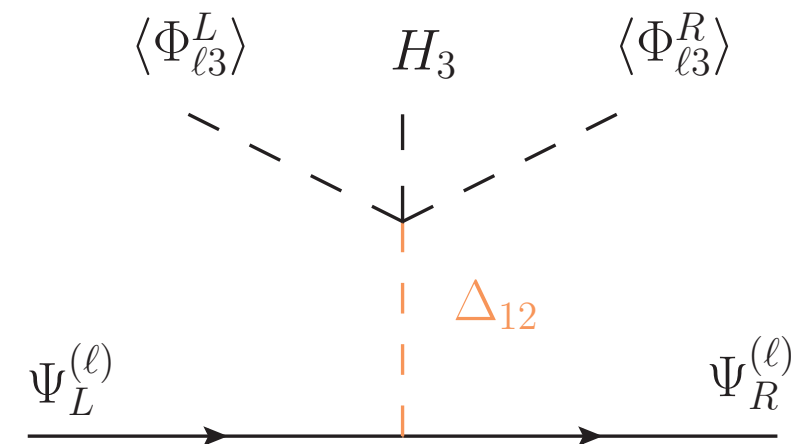
d=5

Like in type-I seesaw
 $[\chi_{L,R} \sim (4, 2, 1)_3]$



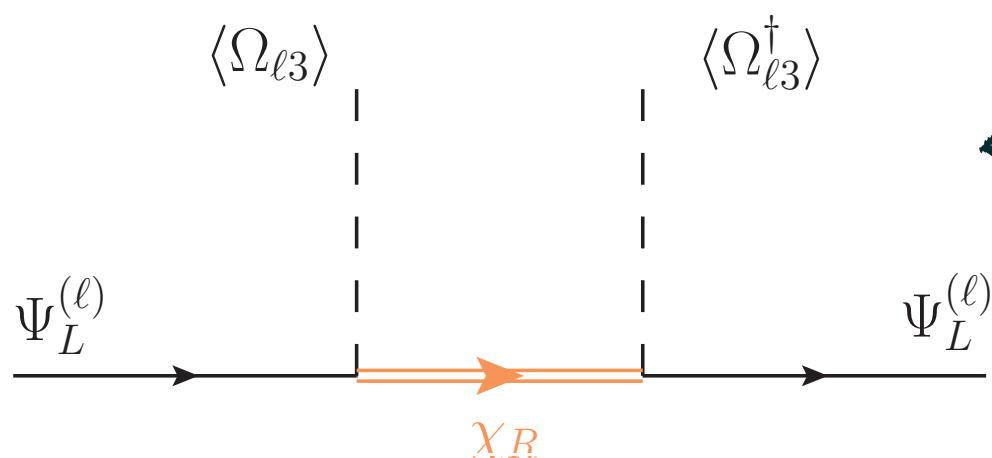
d=6

Like in type-II seesaw
 $[\Delta_{12} \sim (4, 2, 1)_1 \times (\bar{4}, 1, \bar{2})_2]$



“For free”

Leading breaking of the (accidental) U(2) symmetry



Induces a small effective coupling among
 the new vectors and the light fermions
(very important for $b \rightarrow sll$!)

Very similar setup in Greljo, Stefanek, 1802.04274