

Gauge leptoquark and B-anomalies

“From Flavour to New Physics”

Lyon, IPNL - 18.04.2018

Luca Di Luzio



[Based on:
LDL, Nardecchia 1706.01868
LDL, Greljo, Nardecchia, 1708.08450
LDL, Fuentes-Martin, Greljo, Nardecchia, Renner - work in progress]

Outline

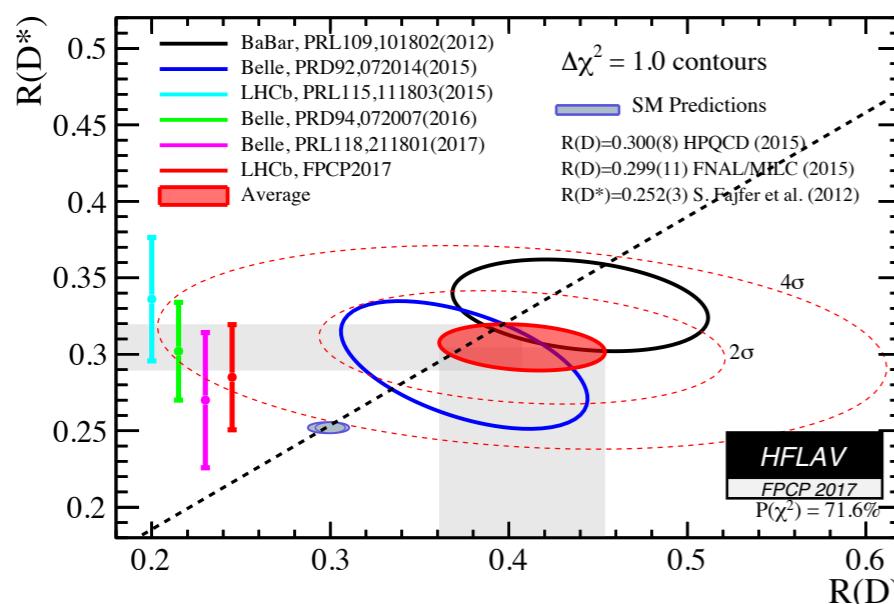
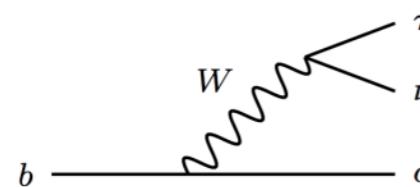
I. Combined explanations of B anomalies

- EFT
- Simplified models
- UV completions → gauge leptoquark ['4321' model]

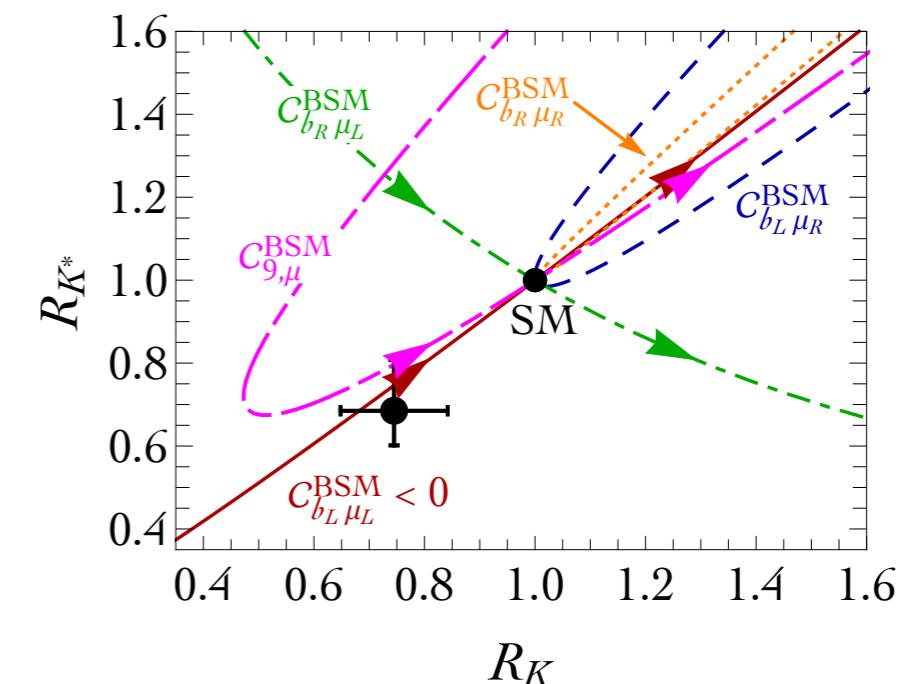
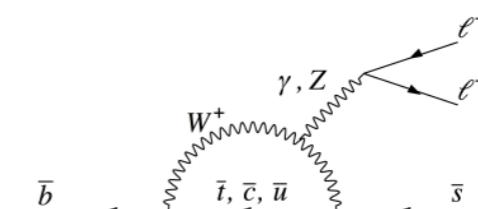
“B-anomalies”

- A seemingly coherent pattern of SM deviations building up since ~ 2012

$b \rightarrow c\tau\nu$



$b \rightarrow s\mu\mu$



+ $R(J/\psi)$

+ $B \rightarrow K^*\mu\mu$ (P'_5) + other BR

EFT [general considerations]

- $SU(2)_L$ triplet operator (combined explanation in SMEFT)

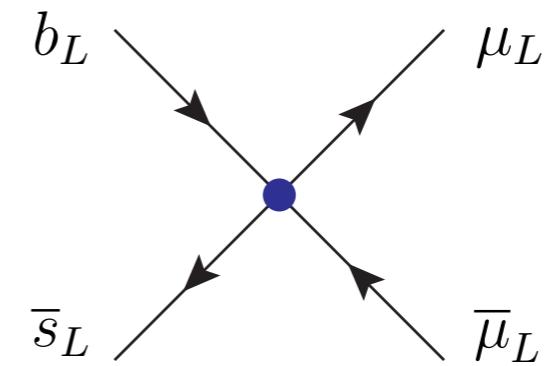
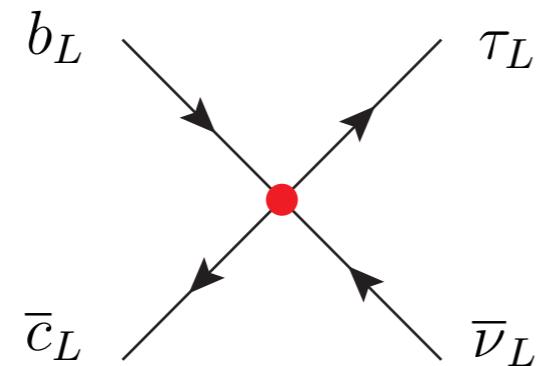
$$\frac{\lambda_{ij}^q \lambda_{kl}^\ell}{\Lambda^2} (\bar{Q}_L^i \sigma^a \gamma_\mu Q_L^j) (\bar{L}_L^k \sigma^a \gamma^\mu L_L^l)$$

[Bhattacharya et al 1412.7164
Alonso, Grinstein, Camalich 1505.05164,
Greljo, Isidori, Marzocca 1506.01705,
Calibbi, Crivellin, Ota 1506.02661, ...]

EFT [general considerations]

- $SU(2)_L$ triplet operator (combined explanation in SMEFT)

$$\frac{\lambda_{ij}^q \lambda_{kl}^\ell}{\Lambda^2} (\bar{Q}_L^i \sigma^a \gamma_\mu Q_L^j) (\bar{L}_L^k \sigma^a \gamma^\mu L_L^l) \supset -\frac{1}{\Lambda_{R_D}^2} 2 \bar{c}_L \gamma^\mu b_L \bar{\tau}_L \gamma_\mu \nu_L + \frac{1}{\Lambda_{R_K}^2} \bar{s}_L \gamma^\mu b_L \bar{\mu}_L \gamma_\mu \mu_L$$



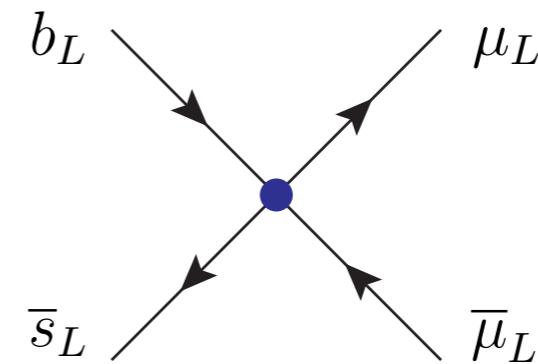
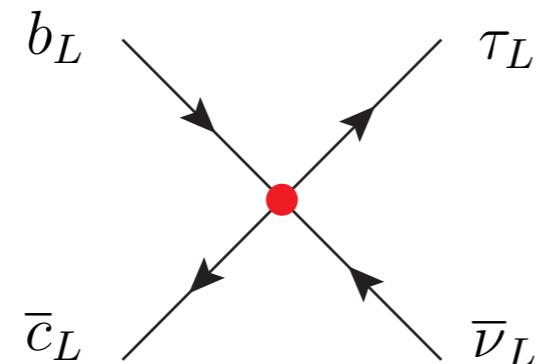
$$\Lambda_{R_D} = 3.4 \text{ TeV} \quad \ll \quad \Lambda_{R_K} = 31 \text{ TeV}$$

what is the scale of new physics ? ($1/\Lambda^2 = g^2/M^2$)

EFT [general considerations]

- $SU(2)_L$ triplet operator (combined explanation in SMEFT)

$$\frac{\lambda_{ij}^q \lambda_{kl}^\ell}{\Lambda^2} (\bar{Q}_L^i \sigma^a \gamma_\mu Q_L^j) (\bar{L}_L^k \sigma^a \gamma^\mu L_L^l) \supset -\frac{1}{\Lambda_{R_D}^2} 2 \bar{c}_L \gamma^\mu b_L \bar{\tau}_L \gamma_\mu \nu_L + \frac{1}{\Lambda_{R_K}^2} \bar{s}_L \gamma^\mu b_L \bar{\mu}_L \gamma_\mu \mu_L$$



$$\Lambda_{R_D} = 3.4 \text{ TeV}$$

\ll

$$\Lambda_{R_K} = 31 \text{ TeV}$$

- Perturbative unitarity bound from $2 \rightarrow 2$ fermion scatterings [↗ see Marco's talk]

$$\sqrt{s}_{R_D} < 9.2 \text{ TeV}$$

$$\sqrt{s}_{R_K} < 84 \text{ TeV}$$



no-loose theorem for HL/HE-LHC ? [LDL, Nardecchia 1706.01868]

EFT [general considerations]

- $SU(2)_L$ triplet operator (combined explanation in SMEFT)

$$\frac{\lambda_{ij}^q \lambda_{kl}^\ell}{\Lambda^2} (\bar{Q}_L^i \sigma^a \gamma_\mu Q_L^j) (\bar{L}_L^k \sigma^a \gamma^\mu L_L^l) \supset -\frac{1}{\Lambda_{R_D}^2} 2 \bar{c}_L \gamma^\mu b_L \bar{\tau}_L \gamma_\mu \nu_L + \frac{1}{\Lambda_{R_K}^2} \bar{s}_L \gamma^\mu b_L \bar{\mu}_L \gamma_\mu \mu_L$$

- Flavour structure:

1. large couplings in taus (SM tree level)
2. sizable couplings in muons (SM one loop)
3. negligible couplings in electrons (well tested, not much room)

$$\lambda_{ij}^{q,\ell} = \delta_{i3}\delta_{j3} + \text{corrections} \quad U(2)_q \times U(2)_\ell \quad \text{approx flavor symmetry}$$

[Barbieri et al | 105.2296, 1512.01560]

$$Q_L^{(3)} \sim q_L^{(b)} = \begin{pmatrix} V_{ib}^* u_L^i \\ b_L \end{pmatrix}$$



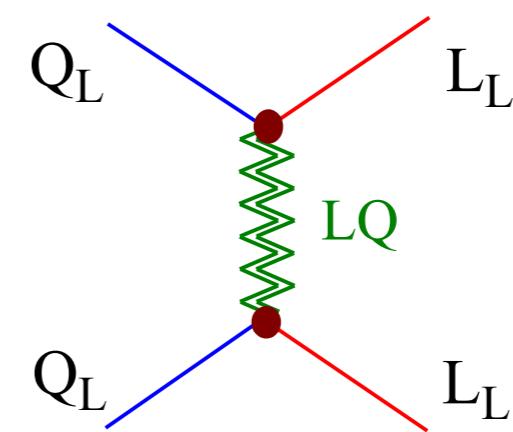
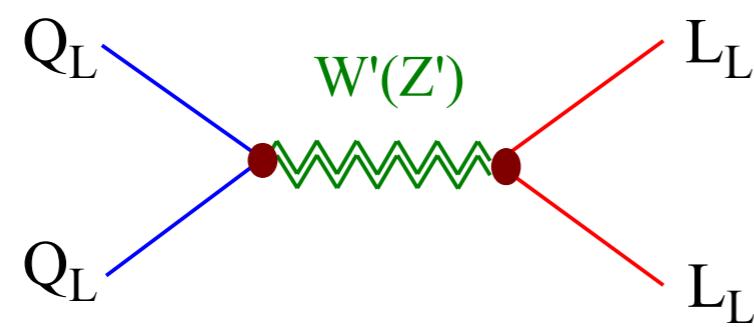
link to SM Yukawa pattern ?

EFT [general considerations]

- $SU(2)_L$ triplet operator (combined explanation in SMEFT)

$$\frac{\lambda_{ij}^q \lambda_{kl}^\ell}{\Lambda^2} (\bar{Q}_L^i \sigma^a \gamma_\mu Q_L^j)(\bar{L}_L^k \sigma^a \gamma^\mu L_L^l)$$

- Tree-level mediators:



Simplified models

- Finite list of tree-level mediators

[Zürich's guide for combined explanations, I706.07808]

[☞ see Admir's talk]

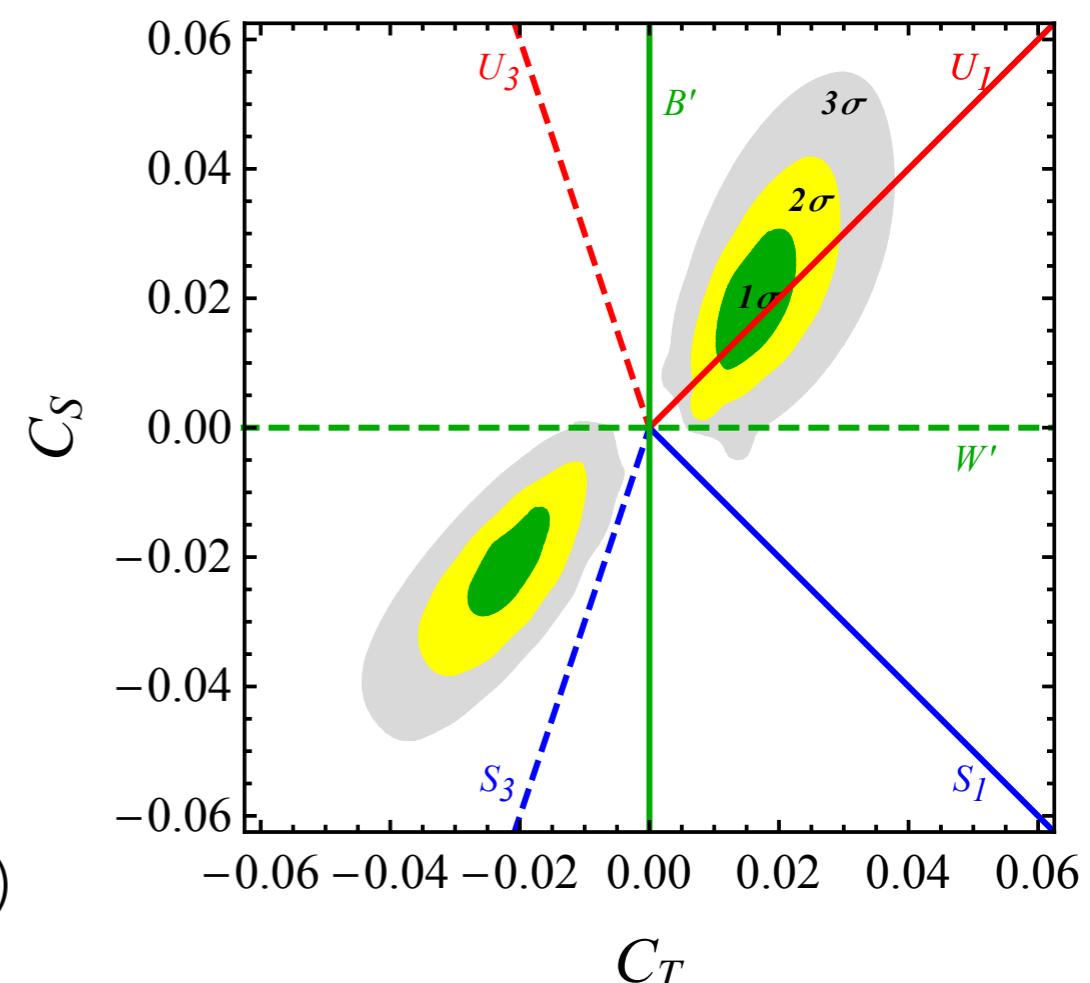
Simplified Model	Spin	SM irrep	C_S/C_T	$R_{D(*)}$	$R_{K(*)}$
Z'	1	(1, 1, 0)	∞	✗	✓
V'	1	(1, 3, 0)	0	✓	✓
S_1	0	($\bar{3}$, 1, 1/3)	-1	✓	✗
S_3	0	($\bar{3}$, 3, 1/3)	3	✓	✓
U_1	1	(3, 1, 2/3)	1	✓	✓
U_3	1	(3, 3, 2/3)	-3	✓	✓

$$\mathcal{B}(B \rightarrow K^* \nu \nu) \propto (C_T - C_S)$$

A clear winner: U_1 → $C_T = C_S$ (at threshold)

Linear combinations also possible → some tuning required

$S_1 + S_3$ [☞ see Svjetlana's talk] or $Z' + V'$ (Bs-mixing constraints) [☞ see Matthew's talk]



UV completion: $U_I \sim (3, 1, 2/3)$

- Massive vectors point to UV dynamics at the TeV scale

composite resonance of
a new strong dynamics

gauge boson of an
extended gauge sector

UV completion: $U_I \sim (3, 1, 2/3)$

- Massive vectors point to UV dynamics at the TeV scale

composite resonance of
a new strong dynamics

$$\frac{G}{H} = \frac{SU(4) \times SO(5) \times U(1)_X}{SU(4) \times SO(4) \times U(1)_X}$$

[Barbieri, Isidori, Pattori, Senia 1502.01560
Barbieri, Murphy, Senia 1611.0493
Buttazzo et al, 1706.07808
Barbieri, Tesi 1712.06844]

- Ambitious program: pNGB Higgs + U_I as composite state of G

- 😊 conceptual link with the naturalness issue of EW scale
- 😢 light LQ lowers the whole resonances' spectrum: issue with direct searches + EWPTs
- 😢 intrinsically non-calculable (e.g. divergent loop observables)

UV completion: $U_1 \sim (3, 1, 2/3)$

- An interesting option: minimal Pati-Salam (PS)

$$G_{PS} = SU(4)_{PS} \times SU(2)_L \times SU(2)_R$$

$$G_{PS}/G_{SM} = U_1 + Z' + W_R$$

gauge boson of an extended gauge sector

😊 hinted by SM chiral structure + everything's calculable

😢 $M_{U_1} \gtrsim 86$ TeV from $K_L^0, B^0, B_s \rightarrow \ell \ell'$ decays (L × R couplings)

[Kutznov et al 1203.0196
+ update from A. D. Smirnov
1801.02895]

mass basis

$$\mathcal{L}_{PS} \supset \frac{g_4}{\sqrt{2}} \left(\bar{d}_L^i \delta_{ij} \gamma_\mu e_L^j + \bar{d}_R^i \delta_{ij} \gamma_\mu e_R^j \right) U_1^\mu \quad \xrightarrow{\text{red arrow}} \quad \frac{g_4}{\sqrt{2}} \left(\bar{d}_L^i \beta_{ij}^L \gamma_\mu e_L^j + \bar{d}_R^i \beta_{ij}^R \gamma_\mu e_R^j \right) U_1^\mu$$

$$\beta^{L,R} = U_{d_{L,R}}^\dagger U_{e_{L,R}} \quad (\text{unitary matrices})$$

UV completion: $U_1 \sim (3, 1, 2/3)$

- An interesting option: minimal Pati-Salam (PS)

$$G_{PS} = SU(4)_{PS} \times SU(2)_L \times SU(2)_R$$

$$G_{PS}/G_{SM} = U_1 + Z' + W_R$$

gauge boson of an extended gauge sector

- 😊 hinted by SM chiral structure + everything's calculable
 - 🙁 $M_{U_1} \gtrsim 86$ TeV from $K_L^0, B^0, B_s \rightarrow \ell \ell'$ decays (L × R couplings)
 - 🙁 Z' direct searches ($M_{U_1} \sim M_{Z'} \sim$ TeV + $\mathcal{O}(g_s)$ Z' couplings to valence quarks)
 - 🙁 neutrino masses also suggest $M_{U_1} \gg$ TeV ($y_{\text{top}} \sim y_{\nu_3-\text{Dirac}}$)
- Minimal PS cannot explain B-anomalies

Beyond minimal PS

- We look for something like

$$\beta^L \sim \begin{pmatrix} \epsilon & \epsilon & \epsilon \\ \epsilon & 0.01 & 0.2 \\ \epsilon & 0.05 & 1 \end{pmatrix} \quad \beta^R \sim \epsilon \quad (\beta^\dagger \beta \neq 1)$$

Beyond minimal PS

- We look for something like

$$\beta^L \sim \begin{pmatrix} \epsilon & \epsilon & \epsilon \\ \epsilon & 0.01 & 0.2 \\ \epsilon & 0.05 & 1 \end{pmatrix} \quad \beta^R \sim \epsilon \quad (\beta^\dagger \beta \neq 1)$$

I): non-minimal **matter** content (mixing with heavy fermions)

[Calibbi, Crivellin, Li 1709.00692]

$$\frac{g_4}{\sqrt{2}} \bar{\mathcal{D}}^A \hat{\beta}_{AB} \gamma_\mu \mathcal{E}^B U_1^\mu \quad \hat{\beta} = \begin{pmatrix} \beta_{LL} & \beta_{LH} \\ \beta_{HL} & \beta_{HH} \end{pmatrix} \quad \hat{\beta}^\dagger \hat{\beta} = 1 \\ \beta_{LL}^\dagger \beta_{LL} \neq 1$$

😢 Z' direct searches

😢 neutrino masses

Beyond minimal PS

- We look for something like

$$\beta^L \sim \begin{pmatrix} \epsilon & \epsilon & \epsilon \\ \epsilon & 0.01 & 0.2 \\ \epsilon & 0.05 & 1 \end{pmatrix} \quad \beta^R \sim \epsilon \quad (\beta^\dagger \beta \neq 1)$$

1): non-minimal **matter** content (mixing with heavy fermions) [Calibbi, Crivellin, Li 1709.00692]

2): non-universal **gauge** interactions [Bordone, Cornella, Fuentes-Martin, Isidori 1712.01368]

$$(422)^3 \sum_{i=1,2,3} \frac{g_4^i}{\sqrt{2}} \bar{Q}^i \gamma^\mu \textcolor{red}{L}^i U_\mu^i \xrightarrow[m_{U_1} \gg m_{U_2} \gg m_{U_3}]{} \beta^{LO} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

😊 flavour hierarchies

😢 neutrino masses [Greljo, Stefanek 1802.04274]

[👉 see Javier's talk]

Beyond minimal PS

- We look for something like

$$\beta^L \sim \begin{pmatrix} \epsilon & \epsilon & \epsilon \\ \epsilon & 0.01 & 0.2 \\ \epsilon & 0.05 & 1 \end{pmatrix} \quad \beta^R \sim \epsilon \quad (\beta^\dagger \beta \neq 1)$$

- 1): non-minimal **matter** content (mixing with heavy fermions) [Calibbi, Crivellin, Li 1709.00692]
- 2): non-universal **gauge** interactions [Bordone, Cornella, Fuentes-Martin, Isidori 1712.01368]
- 3): non-minimal **matter** and **gauge** content (4321 model)

$$G = SU(4) \times SU(3)' \times SU(2)_L \times U(1)' + \text{heavy fermions}$$

[LDL, Greljo, Nardecchia 1708.08450,
See also Diaz, Schmaltz, Zhong 1706.05033,
Georgi, Nakai 1606.05865 for similar constructions]



The '4321' model

$$G = SU(4) \times SU(3)' \times SU(2)_L \times U(1)' \longrightarrow G_{SM} = SU(3)_C \times SU(2)_L \times U(1)_Y$$

$$\langle \Omega_{1,3} \rangle$$

Embedding:

$$SU(3)_C = (SU(3)_4 \times SU(3)')_{diag}$$

$$g_s = \frac{g_4 g_3}{\sqrt{g_4^2 + g_3^2}} \simeq g_3$$

$$U(1)_Y = (U(1)_4 \times U(1)')_{diag}$$

$$g_Y = \frac{g_4 g_1}{\sqrt{g_4^2 + \frac{2}{3}g_1^2}} \simeq g_1$$

Gauge boson spectrum:

$$G/G_{SM} = U + Z' + g'$$

$$M_{g'} \simeq \sqrt{2} M_U \quad M_{Z'} \simeq \frac{1}{\sqrt{2}} M_U$$

→ Structure of gauge symmetry breaking does not allow to decouple g' and Z'

The '4321' model

$$G = SU(4) \times SU(3)' \times SU(2)_L \times U(1)' \longrightarrow G_{SM} = SU(3)_C \times SU(2)_L \times U(1)_Y$$

$$\langle \Omega_{1,3} \rangle$$

Matter content:

Field	$SU(4)$	$SU(3)'$	$SU(2)_L$	$U(1)'$
$q_L'^i$	1	3	2	1/6
$u_R'^i$	1	3	1	2/3
$d_R'^i$	1	3	1	-1/3
$\ell_L'^i$	1	1	2	-1/2
$e_R'^i$	1	1	1	-1
Ψ_L^i	4	1	2	0
Ψ_R^i	4	1	2	0
H	1	1	2	1/2
Ω_3	$\frac{1}{4}$	3	1	1/6
Ω_1	$\frac{1}{4}$	1	1	-1/2

Would-be SM fields

Vector-like fermions (Q'+L')

SSB

} mix after SSB

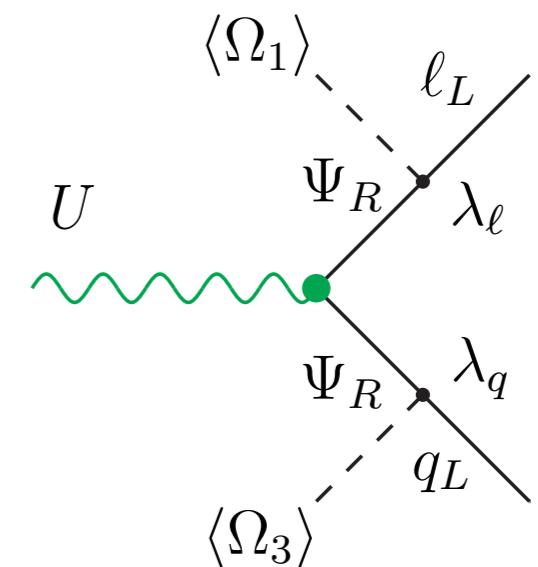
Yukawa sector:

$$\begin{aligned} \mathcal{L}_Y = & -\bar{q}'_L Y_d H d'_R - \bar{q}'_L Y_u \tilde{H} u'_R - \bar{\ell}'_L Y_e H e'_R \\ & - \bar{q}'_L \lambda_q \Omega_3^T \Psi_R - \bar{\ell}'_L \lambda_\ell \Omega_1^T \Psi_R - \bar{\Psi}_L M \Psi_R \end{aligned}$$

Key phenomenological features

I. LQ couples dominantly to 3rd generation LH fields (can satisfy Zürich's EFT criteria)

Field	$SU(4)$	$SU(3)'$	$SU(2)_L$	$U(1)'$
$q_L'^i$	1	3	2	1/6
$u_R'^i$	1	3	1	2/3
$d_R'^i$	1	3	1	-1/3
$\ell_L'^i$	1	1	2	-1/2
$e_R'^i$	1	1	1	-1
Ψ_L^i	4	1	2	0
Ψ_R^i	4	1	2	0
H	1	1	2	1/2
Ω_3	$\frac{1}{4}$	3	1	1/6
Ω_1	$\frac{1}{4}$	1	1	-1/2



$$\begin{aligned} \mathcal{L}_Y = & -\bar{q}'_L Y_d H d'_R - \bar{q}'_L Y_u \tilde{H} u'_R - \bar{\ell}'_L Y_e H e'_R \\ & - \bar{q}'_L \lambda_q \Omega_3^T \Psi_R - \bar{\ell}'_L \lambda_\ell \Omega_1^T \Psi_R - \bar{\Psi}_L M \Psi_R \end{aligned}$$

Key phenomenological features

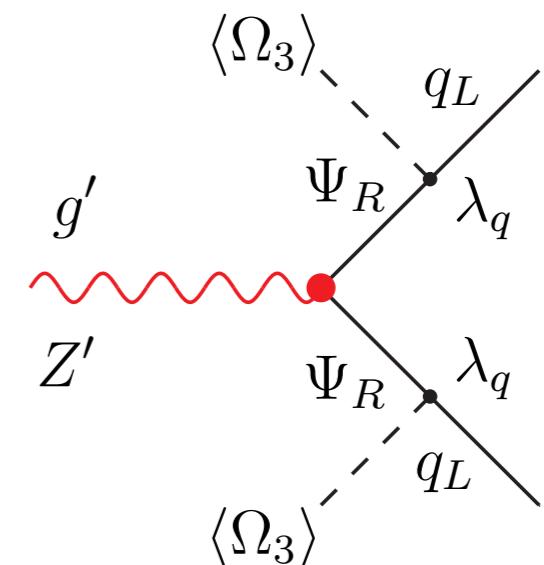
1. LQ couples dominantly to 3rd generation LH fields (can satisfy Zürich's EFT criteria)
2. Down-alignment to avoid tree-level FCNC in the down sector

$$\mathcal{M}_d = \begin{pmatrix} \frac{v}{\sqrt{2}} Y_d^{\text{diag}} & \frac{v_3}{\sqrt{2}} \lambda_q \\ 0 & M^{\text{diag}} \end{pmatrix} \quad \lambda_q = \begin{pmatrix} \lambda_q^d & 0 & 0 \\ 0 & \lambda_q^s & 0 \\ 0 & 0 & \lambda_q^b \end{pmatrix}$$

$$\mathcal{M}_u = \begin{pmatrix} \frac{v}{\sqrt{2}} V^\dagger Y_u^{\text{diag}} & \frac{v_3}{\sqrt{2}} \lambda_q \\ 0 & M^{\text{diag}} \end{pmatrix}$$



CKM-induced
D-mixing



$$\begin{aligned} \mathcal{L}_Y = & - \bar{q}'_L Y_d H d'_R - \bar{q}'_L Y_u \tilde{H} u'_R - \bar{\ell}'_L Y_e H e'_R \\ & - \bar{q}'_L \lambda_q \Omega_3^T \Psi_R - \bar{\ell}'_L \lambda_\ell \Omega_1^T \Psi_R - \bar{\Psi}_L M \Psi_R \end{aligned}$$

Key phenomenological features

1. LQ couples dominantly to 3rd generation LH fields (can satisfy Zürich's EFT criteria)
2. Down-alignment to avoid tree-level FCNC in the down sector
3. Suppressed Z' and g' couplings to light generations

$$\begin{aligned}\mathcal{L}_L = & \frac{g_4}{\sqrt{2}} \bar{Q}'_L \gamma^\mu L'_L U_\mu + \text{h.c.} \\ & + g_s \left(\frac{g_4}{g_3} \bar{Q}'_L \gamma^\mu T^a Q'_L - \frac{g_3}{g_4} \bar{q}'_L \gamma^\mu T^a q'_L \right) g'_\mu^a \\ & + \frac{1}{6} \sqrt{\frac{3}{2}} g_Y \left(\frac{g_4}{g_1} \bar{Q}'_L \gamma^\mu Q'_L - \frac{2}{3} \frac{g_1}{g_4} \bar{q}'_L \gamma^\mu q'_L \right) Z'_\mu \\ & - \frac{1}{2} \sqrt{\frac{3}{2}} g_Y \left(\frac{g_4}{g_1} \bar{L}'_L \gamma^\mu L'_L - \frac{2}{3} \frac{g_1}{g_4} \bar{\ell}'_L \gamma^\mu \ell'_L \right) Z'_\mu\end{aligned}$$

Key phenomenological features

1. LQ couples dominantly to 3rd generation LH fields (can satisfy Zürich's EFT criteria)
2. Down-alignment to avoid tree-level FCNC in the down sector
3. Suppressed Z' and g' couplings to light generations

$$\begin{aligned}\mathcal{L}_L = & \frac{g_4}{\sqrt{2}} \bar{Q}'_L \gamma^\mu L'_L U_\mu + \text{h.c.} \\ & + g_s \left(\frac{g_4}{g_3} \bar{Q}'_L \gamma^\mu T^a Q'_L - \frac{g_3}{g_4} \bar{q}'_L \gamma^\mu T^a q'_L \right) g'_\mu{}^a \\ & + \frac{1}{6} \sqrt{\frac{3}{2}} g_Y \left(\frac{g_4}{g_1} \bar{Q}'_L \gamma^\mu Q'_L - \frac{2}{3} \frac{g_1}{g_4} \bar{q}'_L \gamma^\mu q'_L \right) Z'_\mu \\ & - \frac{1}{2} \sqrt{\frac{3}{2}} g_Y \left(\frac{g_4}{g_1} \bar{L}'_L \gamma^\mu L'_L - \frac{2}{3} \frac{g_1}{g_4} \bar{\ell}'_L \gamma^\mu \ell'_L \right) Z'_\mu\end{aligned}$$



requires the phenomenological limit
 $g_4 \gg g_3 \simeq g_s \gg g_1 \simeq g_Y$

Key phenomenological features

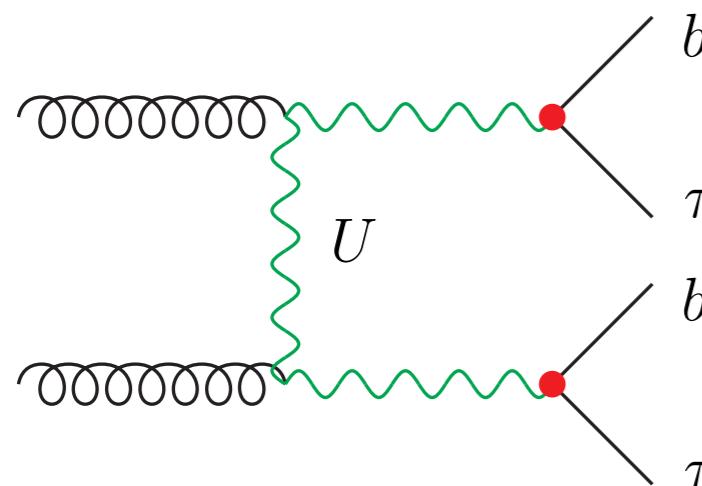
1. LQ couples dominantly to 3rd generation LH fields (can satisfy Zürich's EFT criteria)
2. Down-alignment to avoid tree-level FCNC in the down sector
3. Suppressed Z' and g' couplings to light generations
4. B and L accidental global symmetries as in the SM ($m_\nu = 0$)

$$\mathcal{O}_5 = \frac{1}{\Lambda_L} \ell' \ell' H H \quad \Lambda_L \gg v$$

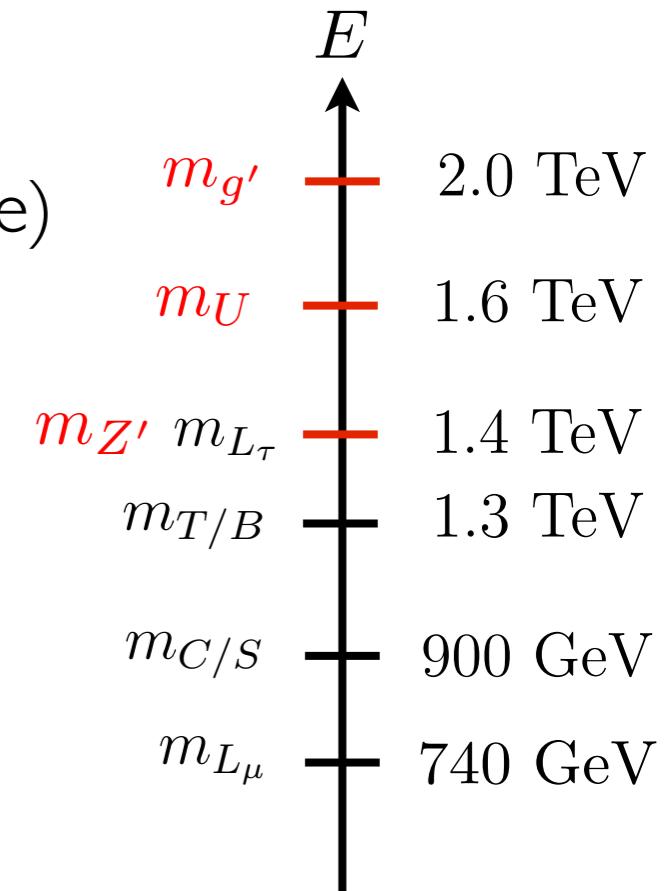
High-pT searches

- LQ pair production via QCD

- 3rd generation final states (fixed by anomaly and $SU(2)_L$ invariance)



$$\begin{cases} U \rightarrow b\tau^+, & \text{BR} = 50 \% \\ U \rightarrow t\bar{\nu}, & \text{BR} = 50 \% \end{cases}$$



[CMS search for spin-0, 1703.03995
recast for spin-1 1706.01868 (see also 1706.05033) + Moriond EW update]

$$m_U \gtrsim 1.5 \text{ TeV}$$

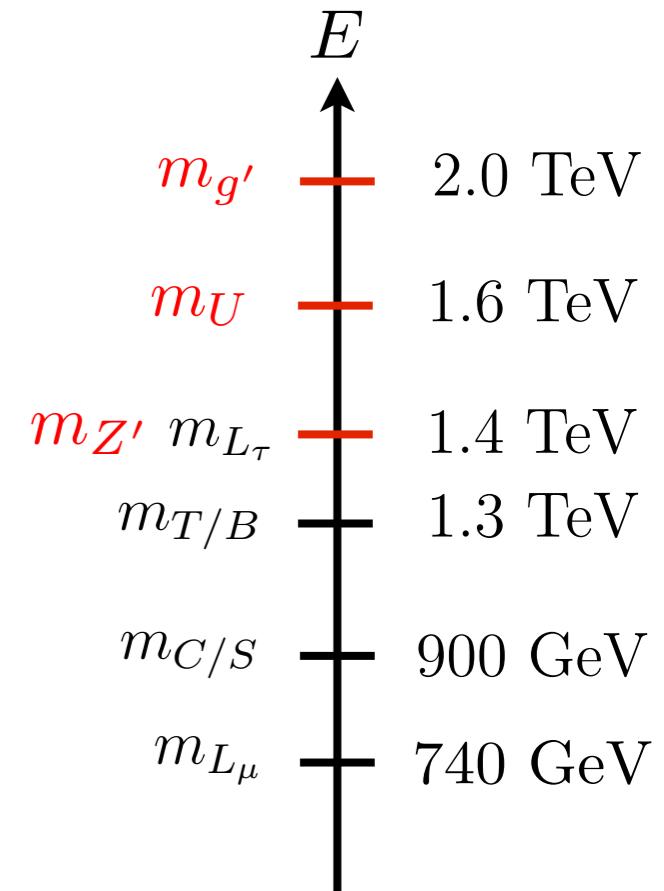


LQ mass sets the overall scale: $M_{g'} \simeq \sqrt{2} M_U$ $M_{Z'} \simeq \frac{1}{\sqrt{2}} M_U$

High-pT searches

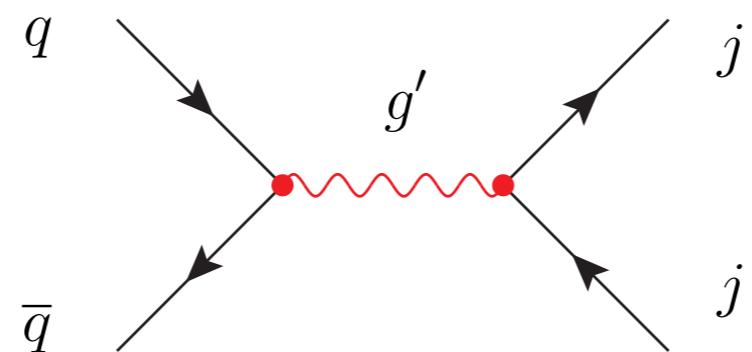
- LQ pair production via QCD
- Z' Drell-Yan production naturally suppressed

$$\sin \theta_{Z'} = \sqrt{\frac{3}{2}} \frac{g_Y}{g_4} \simeq 0.09 \quad \xrightarrow{\hspace{1cm}} \quad \text{requires } g_4 \gtrsim 3$$



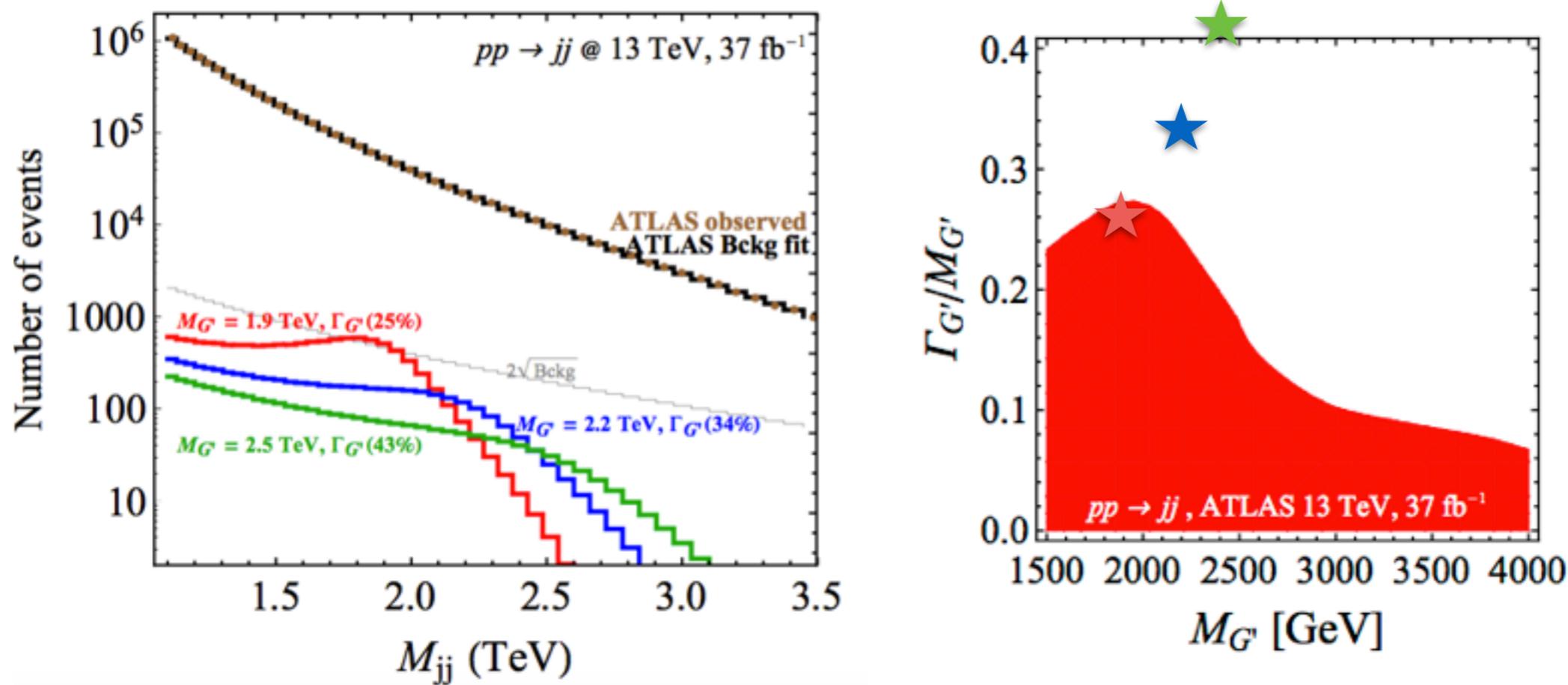
- g' resonant di-jet searches [ATLAS, 1703.09127]

$$\sin \theta_{g'} = \frac{g_s}{g_4} \simeq 0.3 \quad \xrightarrow{\hspace{1cm}} \quad 2 \text{ TeV coloron naively excluded}$$



High-pT searches

[LDL, Fuentes-Martin, Greljo, Nardecchia, Renner (work in progress)]



- However, bump-searches loose in sensitivity for large width/mass

$$\frac{\Gamma}{m} \lesssim 15\% \quad (\text{exp. analysis})$$

$$\frac{\Gamma_{g'}}{m_{g'}} \simeq 25\% \quad (\text{unavoidable in our scenario: large } g_4 + \text{extra channel in VLF})$$

Conclusions

1. We will know much more by ~ 2020 (LHCb + Belle II)
2. Early speculations point to TeV-scale vector leptoquark ($R(D) + R(K)$ explanation)



who ordered that ?

3. In the meantime, lesson from UV complete models



unexpected experimental signatures (coloron, D-mixing, ...) + playground to compute correlations

[More pheno to come: LDL, Fuentes-Martin, Greljo, Nardecchia, Renner (work in progress)]

Backup slides

Down-alignment (flavour symmetry)

- Down-alignment to avoid tree-level FCNC in the down sector
- d'_R, Ψ_L, Ψ_R as triplets of the flavour group $U(3)_{d'_R} \equiv U(3)_{\Psi_L} \equiv U(3)_{\Psi_R}$
- $M \propto$ identity
- Y_d and $\lambda_q \propto$ to the same spurion $(\bar{3}, 3)$ of $U(3)_{q'_L} \times U(3)_{d'_R}$  simultaneously diagonalizable

$$\mathcal{M}_d = \begin{pmatrix} \frac{v}{\sqrt{2}} Y_d^{\text{diag}} & \frac{v_3}{\sqrt{2}} \lambda_q \\ 0 & M^{\text{diag}} \end{pmatrix} \quad \lambda_q = \begin{pmatrix} \lambda_q^d & 0 & 0 \\ 0 & \lambda_q^s & 0 \\ 0 & 0 & \lambda_q^b \end{pmatrix}$$

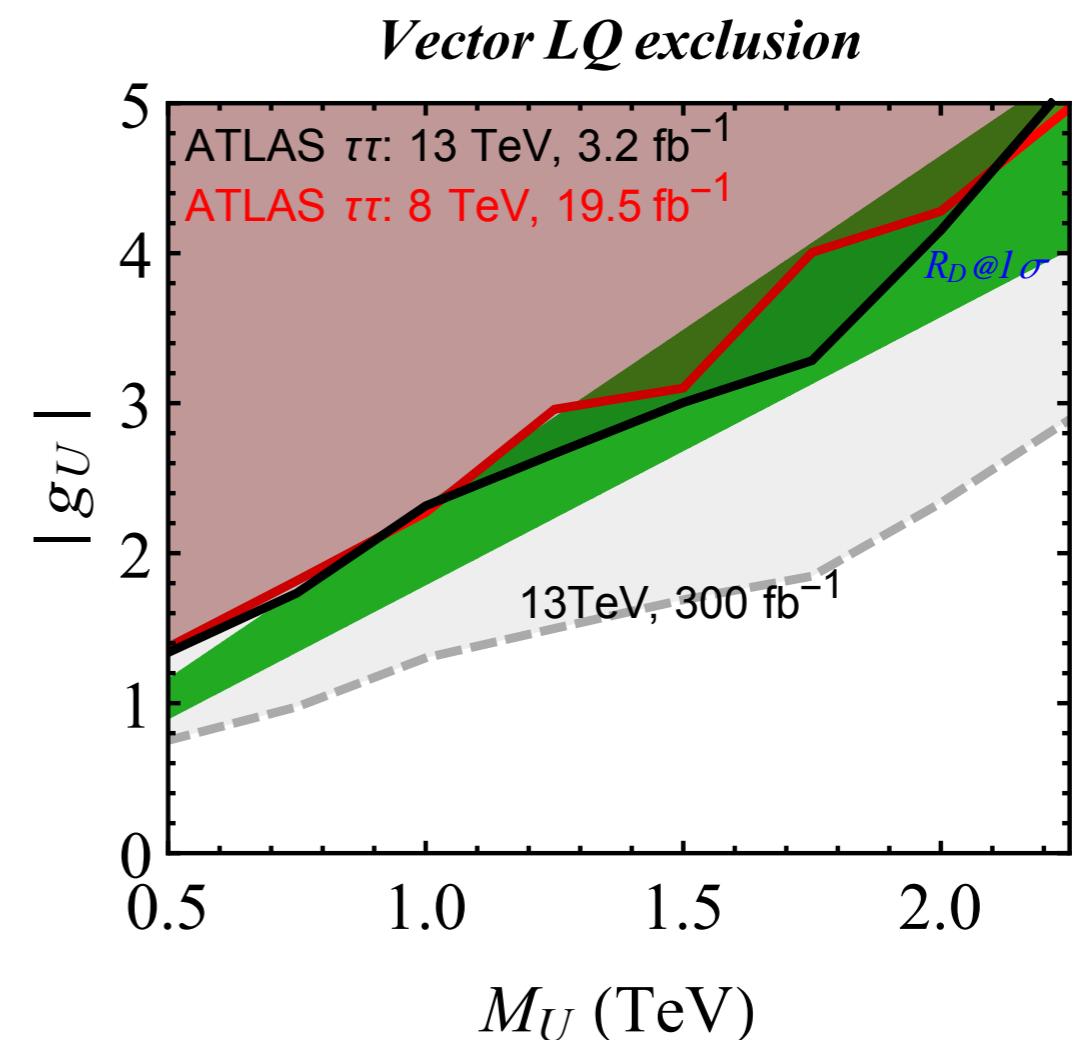
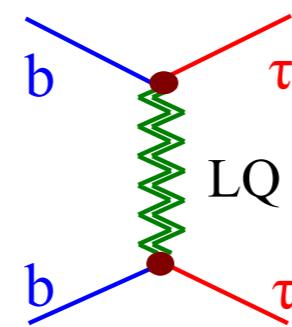
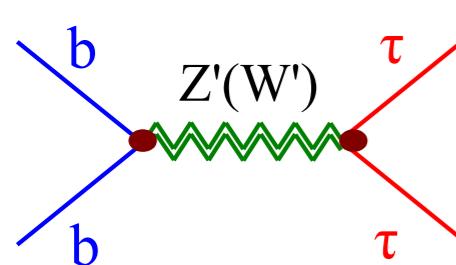
$$\mathcal{M}_u = \begin{pmatrix} \frac{v}{\sqrt{2}} V^\dagger Y_u^{\text{diag}} & \frac{v_3}{\sqrt{2}} \lambda_q \\ 0 & M^{\text{diag}} \end{pmatrix}$$

$$\begin{aligned} \mathcal{L}_Y = & -\bar{q}'_L Y_d H d'_R - \bar{q}'_L Y_u \tilde{H} u'_R - \bar{\ell}'_L Y_e H e'_R \\ & - \bar{q}'_L \lambda_q \Omega_3^T \Psi_R - \bar{\ell}'_L \lambda_\ell \Omega_1^T \Psi_R - \bar{\Psi}_L M \Psi_R \end{aligned}$$

EFT [problems]

- Three main problems mainly driven by $R(D)$

I. High- p_T constraints



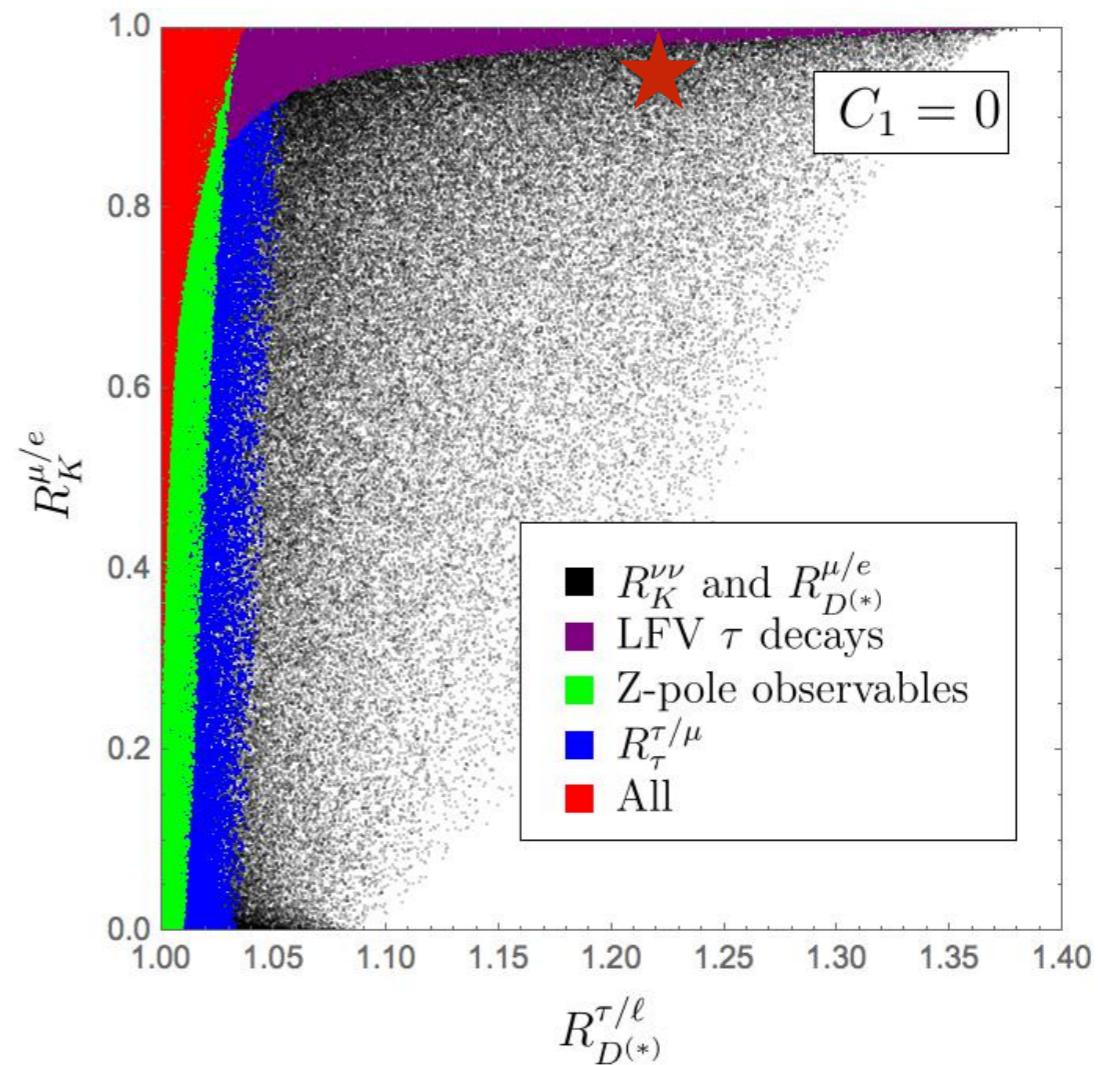
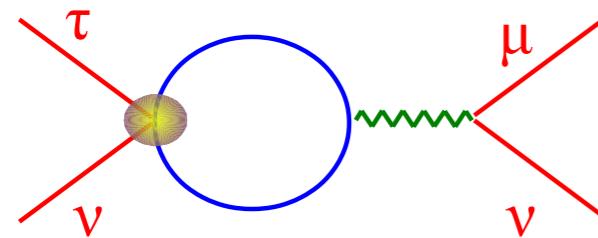
[Faroughy, Greljo, Kamenik | 609.07 | 38]

EFT [problems]

- Three main problems mainly driven by $R(D)$

1. High- p_T constraints

2. Radiative constraints

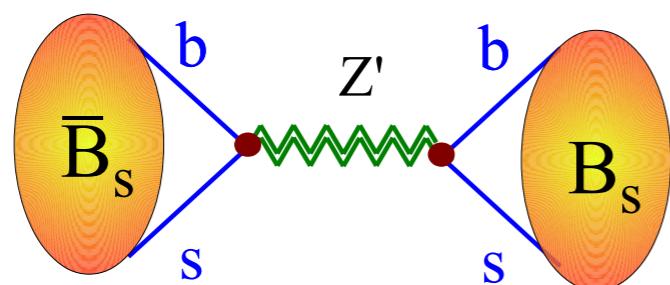


[Feruglio, Paradisi, Pattori | 606.00524, 1705.00929]

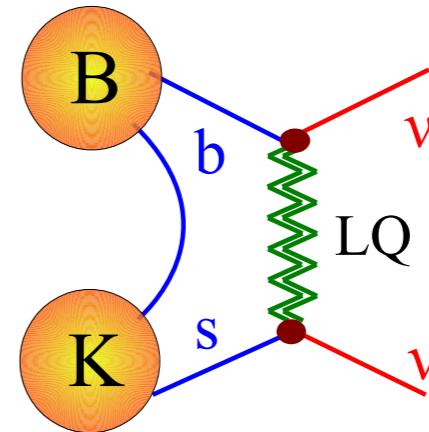
EFT [problems]

- Three main problems mainly driven by R(D)

1. High- p_T constraints
2. Radiative constraints
3. Flavour bounds



(absent at tree-level with LQ)



(consequence of $SU(2)_L$ invariance)

EFT [solutions]

- Tension gets drastically alleviated if [Zürich's guide for combined explanations, 1706.07808]

I. Triplet + Singlet operator (more freedom in $SU(2)_L$ structure)

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} - \frac{1}{v^2} \lambda_{ij}^q \lambda_{\alpha\beta}^\ell \left[C_T (\bar{Q}_L^i \gamma_\mu \sigma^a Q_L^j)(\bar{L}_L^\alpha \gamma^\mu \sigma^a L_L^\beta) + C_S (\bar{Q}_L^i \gamma_\mu Q_L^j)(\bar{L}_L^\alpha \gamma^\mu L_L^\beta) \right]$$

2. Deviation from 'pure-mixing' scenario

$$\bar{Q}^i \lambda_{ij}^q Q^j = \begin{pmatrix} \bar{u}^k V_{ki} & \bar{d}^i \end{pmatrix} \lambda_{ij}^q \begin{pmatrix} V_{jl}^\dagger u^l \\ d^j \end{pmatrix} \supset \bar{c} (V_{cb} \lambda_{bb}^q + V_{cs} \lambda_{sb}^q + \dots) b$$

$$R_{D^{(*)}}^{\tau\ell} \approx 1 + 2C_T \left(1 - \lambda_{sb}^q \frac{V_{tb}^*}{V_{ts}^*} \right) \quad \xrightarrow{\text{red arrow}} \quad \lambda_{sb}^q > \mathcal{O}(V_{cb}) \quad \text{allows for larger NP scale}$$

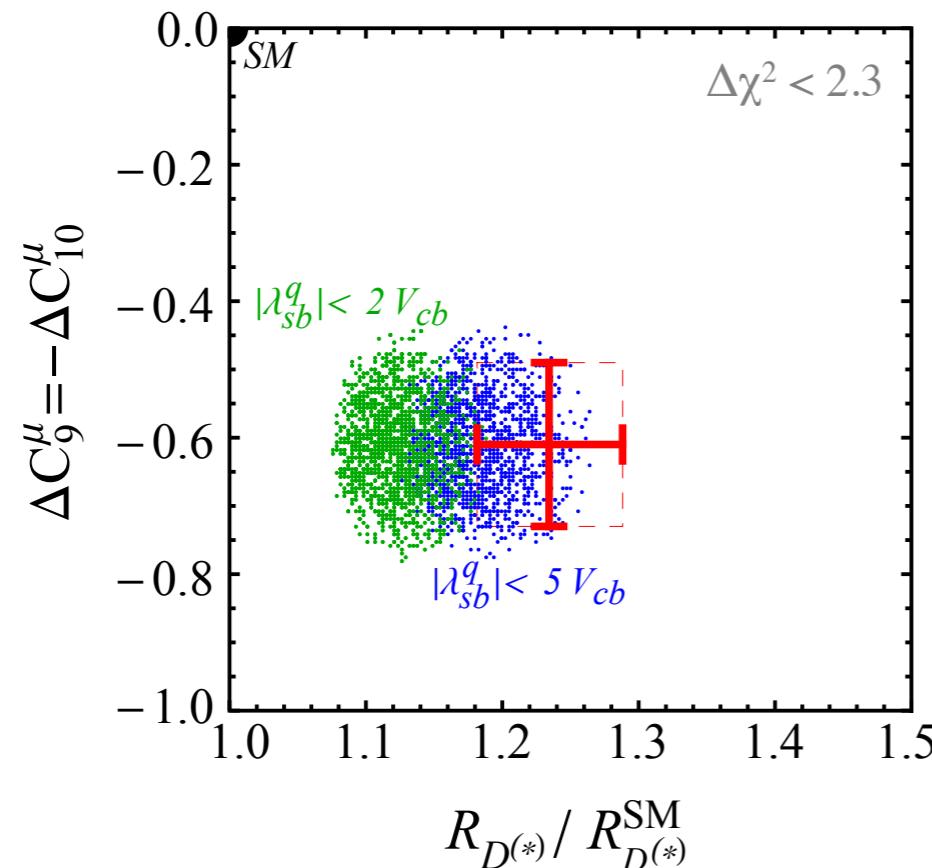
EFT [solutions]

- Tension gets drastically alleviated if [Zürich's guide for combined explanations, 1706.07808]

1. Triplet + Singlet operator (more freedom in $SU(2)_L$ structure)

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} - \frac{1}{v^2} \lambda_{ij}^q \lambda_{\alpha\beta}^\ell \left[C_T (\bar{Q}_L^i \gamma_\mu \sigma^a Q_L^j)(\bar{L}_L^\alpha \gamma^\mu \sigma^a L_L^\beta) + C_S (\bar{Q}_L^i \gamma_\mu Q_L^j)(\bar{L}_L^\alpha \gamma^\mu L_L^\beta) \right]$$

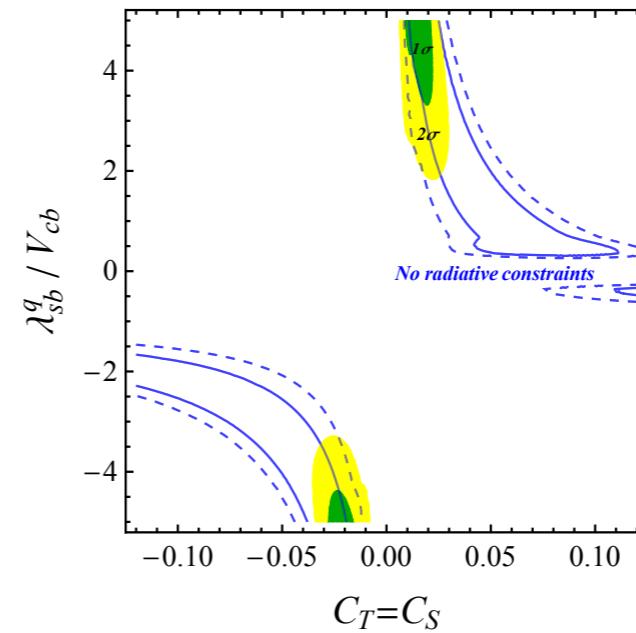
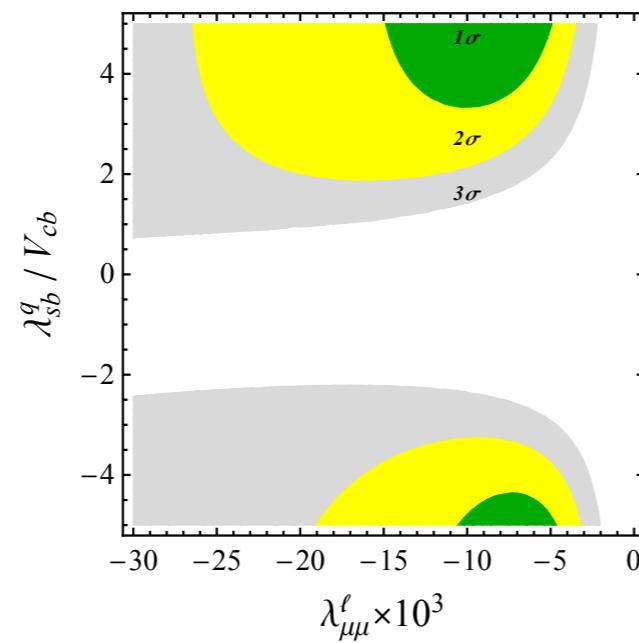
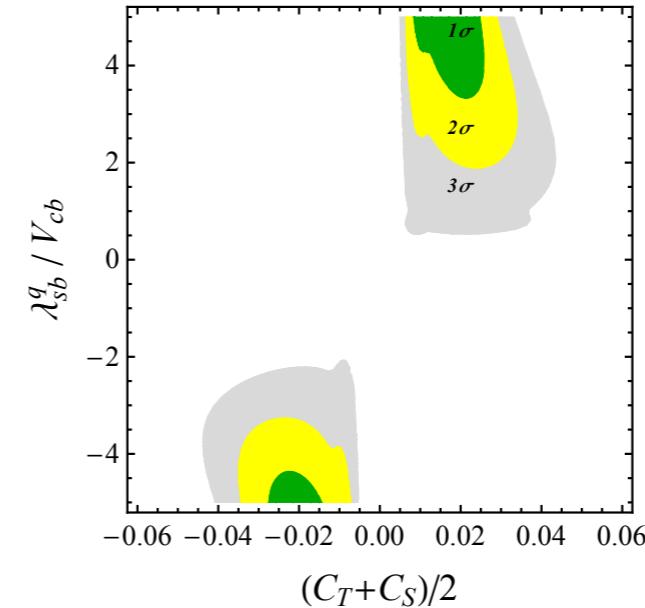
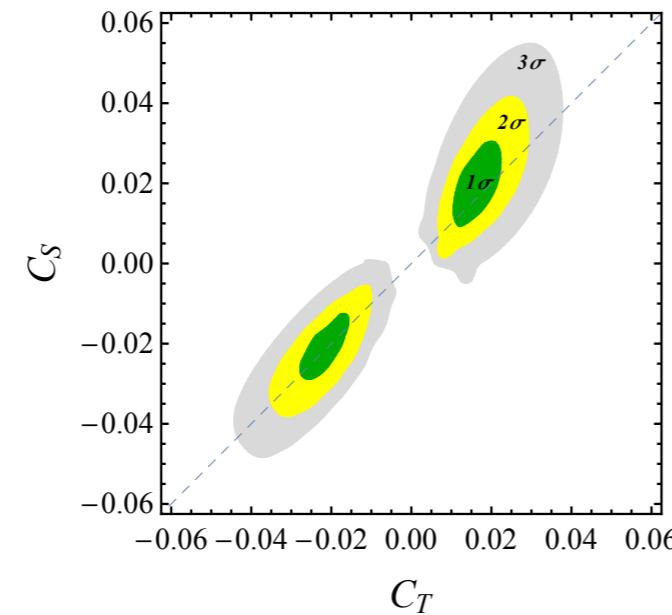
2. Deviation from 'pure-mixing' scenario



$\lambda_{sb}^q > \mathcal{O}(V_{cb})$ allows for larger NP scale

EFT [details fit]

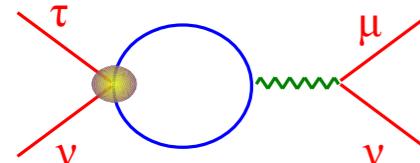
- 4 parameters fit: $C_S, C_T, \lambda_{bs}^q, \lambda_{\mu\mu}^\ell$ ($\lambda_{bb}^q = \lambda_{\tau\tau}^\ell = 1$) [Zürich's guide for combined explanations, 1706.07808]



EFT [details fit]

- 4 parameters fit: $C_S, C_T, \lambda_{bs}^q, \lambda_{\mu\mu}^\ell$ ($\lambda_{bb}^q = \lambda_{\tau\tau}^\ell = 1$) [Zürich's guide for combined explanations, 1706.07808]

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} - \frac{1}{v^2} \lambda_{ij}^q \lambda_{\alpha\beta}^\ell \left[C_T (\bar{Q}_L^i \gamma_\mu \sigma^a Q_L^j) (\bar{L}_L^\alpha \gamma^\mu \sigma^a L_L^\beta) + C_S (\bar{Q}_L^i \gamma_\mu Q_L^j) (\bar{L}_L^\alpha \gamma^\mu L_L^\beta) \right]$$



LH Z - τ - τ coupling

LH Z - ν - ν coupling

LFUV in τ decays

LFV in τ decays

Observable	Experimental bound	Linearised expression
$R_{D^{(*)}}^{\tau\ell}$	1.237 ± 0.053	$1 + 2C_T(1 - \lambda_{sb}^q V_{tb}^*/V_{ts}^*)(1 - \lambda_{\mu\mu}^\ell/2)$
$\Delta C_9^\mu = -\Delta C_{10}^\mu$	-0.61 ± 0.12 [36]	$-\frac{\pi}{\alpha_{\text{em}} V_{tb} V_{ts}^*} \lambda_{\mu\mu}^\ell \lambda_{sb}^q (C_T + C_S)$
$R_{b \rightarrow c}^{\mu e} - 1$	0.00 ± 0.02	$2C_T(1 - \lambda_{sb}^q V_{tb}^*/V_{ts}^*) \lambda_{\mu\mu}^\ell$
$B_{K^{(*)}\nu\bar{\nu}}$	0.0 ± 2.6	$1 + \frac{2}{3} \frac{\pi}{\alpha_{\text{em}} V_{tb} V_{ts}^* C_\nu^{\text{SM}}} (C_T - C_S) \lambda_{sb}^q (1 + \lambda_{\mu\mu}^\ell)$
$\delta g_{\tau_L}^Z$	-0.0002 ± 0.0006	$0.033C_T - 0.043C_S$
$\delta g_{\nu_\tau}^Z$	-0.0040 ± 0.0021	$-0.033C_T - 0.043C_S$
$ g_\tau^W/g_\ell^W $	1.00097 ± 0.00098	$1 - 0.084C_T$
$\mathcal{B}(\tau \rightarrow 3\mu)$	$(0.0 \pm 0.6) \times 10^{-8}$	$2.5 \times 10^{-4} (C_S - C_T)^2 (\lambda_{\tau\mu}^\ell)^2$