





Momentum-space atom correlations in a Mott insulator







0.02

200

1.8

1.2

0

 $g^{(2)}(\delta k)$

• x-axis (1,0,0)

y-axis (0,1,0)

z-axis (0,0,1)

0.04

 $\delta k \left[k_a \right]$



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Congrès général de la SFP 2019





A small detour in quantum optics: HBT experiment



R. Hanbury-Brown and R. Q. Twiss: correlations in the light intensity Nature 177, 27-29 (1956)

$$g^{(2)}(\vec{r_1}, \vec{r_2}, \tau) = \frac{\langle i(\vec{r_1}, t)i(\vec{r_2}, t+\tau) \rangle}{\langle i(\vec{r_1}, t) \rangle \langle i(\vec{r_2}, t+\tau) \rangle}$$



A small detour in quantum optics: HBT experiment





Nature 178, 1046-1048 (1956)

Nature 177, 27-29 (1956)

Correlation at zero distance equals 2:

$$g^{(2)}(d=0) = 2$$

Correlation length l_c provides measure of the **angular diameter** α **of the star**:

$$l_c \propto \lambda/lpha$$

The HBT stellar interferometer: a classical explanation

A star is an incoherent source of light: modelled by <u>many random emitters</u>



The detected field is sum of many independent random variables:

$$\mathcal{E}(P_1, t) = \sum_{\text{star}} a_j e^{i[\phi_j - \omega_j t - k.r_j]}$$

<u>Central limit theorem:</u> $\mathcal{E}(P_1, t)$ is a **Gaussian random variable**

$$g^{(2)}(d) = 1 + |g^{(1)}(d)|^2$$

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distribution of the **incoherent** emitters: <u>diameter of the star</u>

(identical to Michelson interferometry)

The HBT stellar interferometer: a quantum explanation

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Thermal light = many modes of ideal bosons (photons) populated:

$$H = \sum_{k} N(k)\hbar\omega(k) \ a(k)^{\dagger}a(k)$$

Described by a Gaussian density operator

Wick theorem applies:

$$\langle a_1^{\dagger} a_2^{\dagger} a_3 a_4 \rangle = \langle a_1^{\dagger} a_3 \rangle \langle a_2^{\dagger} a_4 \rangle + \langle a_1^{\dagger} a_4 \rangle \langle a_2^{\dagger} a_3 \rangle$$

$$\mathbf{g}^{(2)}(d) = 1 + |g^{(1)}(d)|^2$$

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<u>Characteristics of HBT experiments with</u> incoherent source of non-interacting bosons:

 \longrightarrow $g^{(2)}(d=0) = 2$: source with Gaussian statistics

 $\mathcal{E}(P_1, t) = \sum a_j \, \mathrm{e}^{i[\phi_j - \omega_j t - k.r_j]}$

star

<u>Central limit theorem:</u> $\mathcal{E}(P_1, t)$

is a Gaussian random variable

 $g^{(2)}(d) = 1 + |g^{(1)}(d)|^2$

 \longrightarrow Shape of $g^{(2)}(d)$ is set by the spatial distribution of the incoherent emitters

Correlations in a Mott insulator deep in the insulating regime

Bosons loaded in the lowest band of a lattice described by the **Bose-Hubbard Hamiltonian**:



$$H = -J\sum_{j,j'} b_j^{\dagger} b_{j'} + \frac{U}{2} \sum_j n_j (n_j - 1)$$

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In the regime $U \gg J$ the atoms are pinned in the lattice site with unity occupation: **Mott insulator**

Deep in the Mott regime:

- negligible fluctuations of atom number per lattice site
- negligible site-to-site phase coherence



Interaction term is negligible:
$$H \simeq \sum_{k} \epsilon(k) a^{\dagger}(k) a(k)$$

Similar to ideal bosons with no coherence!

Detection of lattice gases with the Helium detector

Adiabatic loading of Bose-Einstein Condensate into a 3D square optical lattice (wave-length 1550 nm)



3 pairs of counter-propagating laser beams forming standing waves

Bose-Einstein condensates of metastable Helium-4



Phys. Rev. A **90** 063407 (2014) Phys. Rev. A **91** 061402(R) (2015)

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Bose-Einstein condensate when released from a 3D lattice



3D distribution of single atoms





Mott insulator with U/J = 100

Momentum density is featureless (due to absence of phase coherence)



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Two-body correlations in momentum-space:



$$g^{(2)}(\vec{k},\vec{k'}) = \frac{\langle a^{\dagger}(\vec{k})a^{\dagger}(\vec{k'})a(\vec{k})a(\vec{k'})\rangle}{\rho(\vec{k})\rho(\vec{k'})}$$

1D cuts along a given axis:
$$ec{k} - ec{k'} = \delta k.ec{u}$$

perfectly contrasted bosonic bunching



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Pioneering experiment using noise correlation on atomic densities:



Bunching observed in the expanded density!

Nature **434**, 481 (2005)



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Two-body correlations in momentum-space:



Two-body correlation length



Mott insulator with U/J = 100

Momentum density is featureless (due to absence of phase coherence)

in-trap size L





Two-body correlation length

Excellent agreement with no adjustable parameters!

(incompressible Mott state)





dashed line: Gutzwiller solution for the experiment (3D Mott state and $\langle n_j \rangle = 1$)



Mott insulator with U/J = 100

Momentum density is featureless (due to absence of phase coherence)

From **recorded full distribution of atoms** in 3D, we can extract higher-order correlations:



Two-body correlations in momentum-space:



Momentum-space correlations deep in the Mott regime are that of a Gaussian many-body ground-state

Depletion of strongly-interacting condensate

Momentum distribution of lattice condensates:



Conclusion



Single-atom detection of lattice gases in 3D momentum space

H. Cayla et al., Phys. Rev. A 97 061609(R) (2018)

Many-body correlations in the deep Mott regime





C. Carcy et al., arXiv:1904.10995 (2019)

Non-Gaussian correlations in interacting depleted bosons

Quantum-Monte Carlo:



(Flatiron Institute,

New York)



T. Roscilde (ENS Lyon)



M. Mancini C. Carcy H. Cayla



A. Tenart



A. Aspect