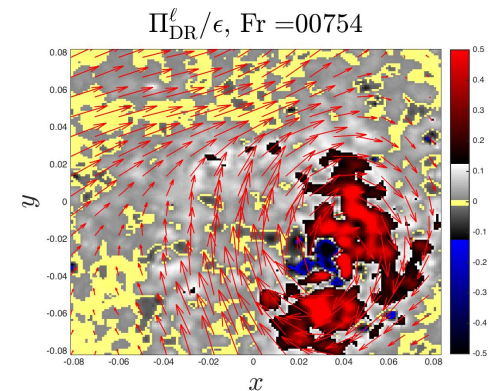
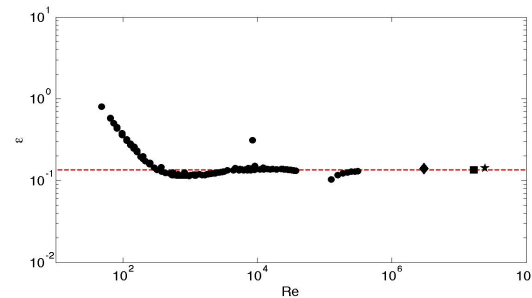
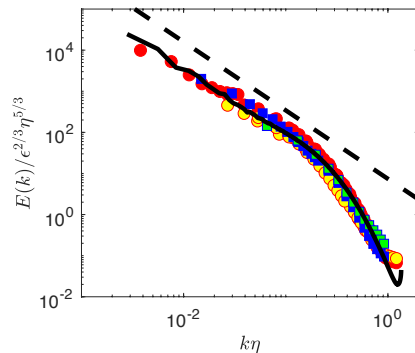


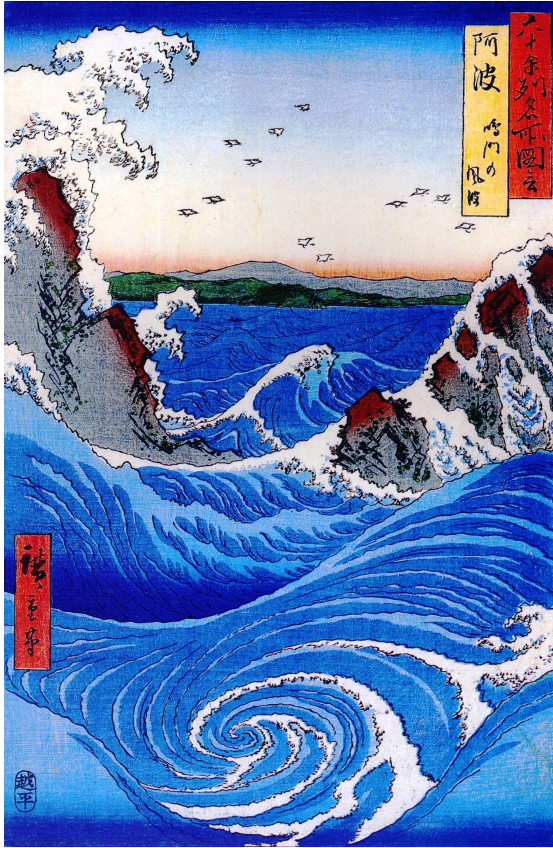
Physique de la turbulence à l'échelle de Kolmogorov: *(Enquête autour de 5 énigmes de la turbulence)*

B. Dubrulle

CEA Saclay/SPEC/SPHYNX
CNRS UMR 3680



What is turbulence?



Hiroshige 1797-1858

Narito vortex



*Définition: **Turbulence** describes the state of a fluid (liquid or gas) in which velocity is in a swirling state.*

$$\begin{aligned}\vec{\nabla} \cdot \vec{u} &= 0 \\ \partial_t \vec{u} + (\vec{u} \cdot \vec{\nabla}) \vec{u} &= -\frac{1}{\rho} \vec{\nabla} p + \nu \Delta \vec{u}\end{aligned}$$

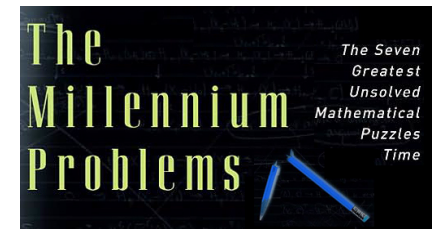
$$\text{Re} = \frac{(u \nabla u)}{\nu \Delta u} = \frac{LU}{\nu}$$

Turbulence: $\text{Re} \gg 1$

Open basic question:

*Navier-Stokes equations are they well-posed?
Are they any singularities?*

Puzzle # 0 of turbulence



Vortices in the Universe...

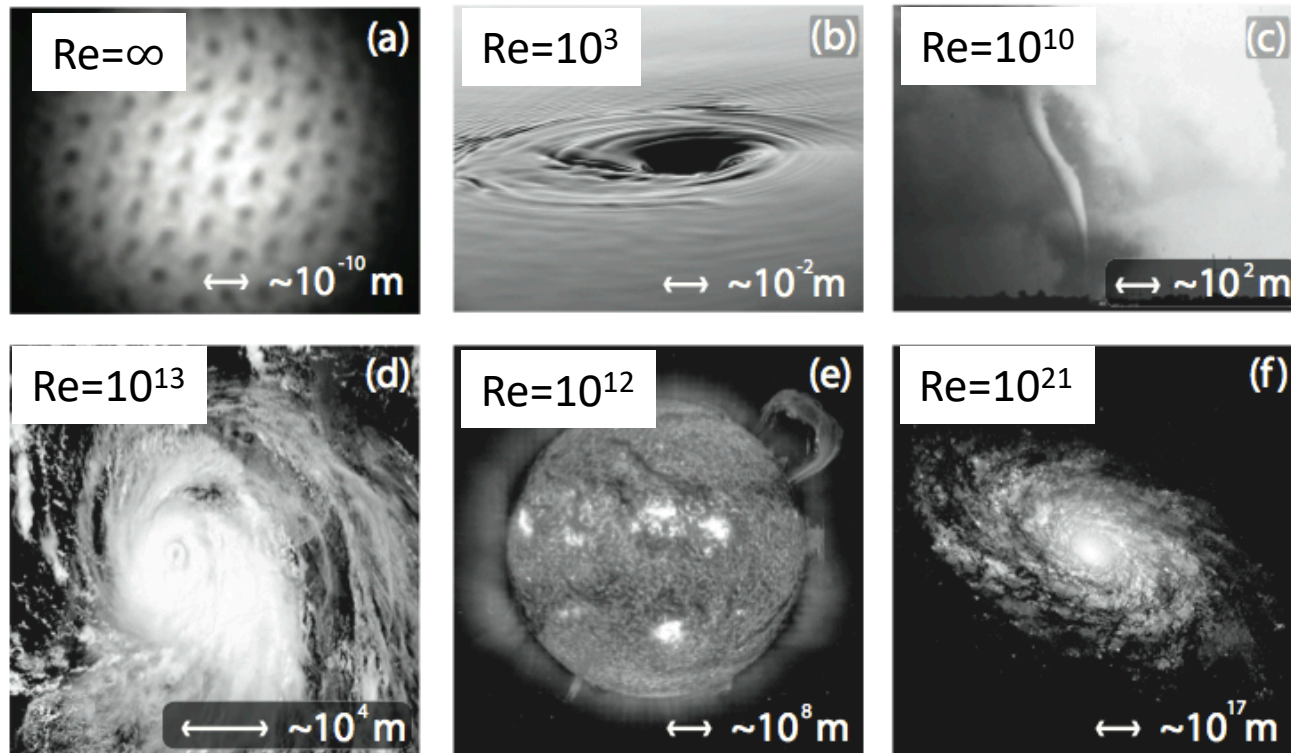
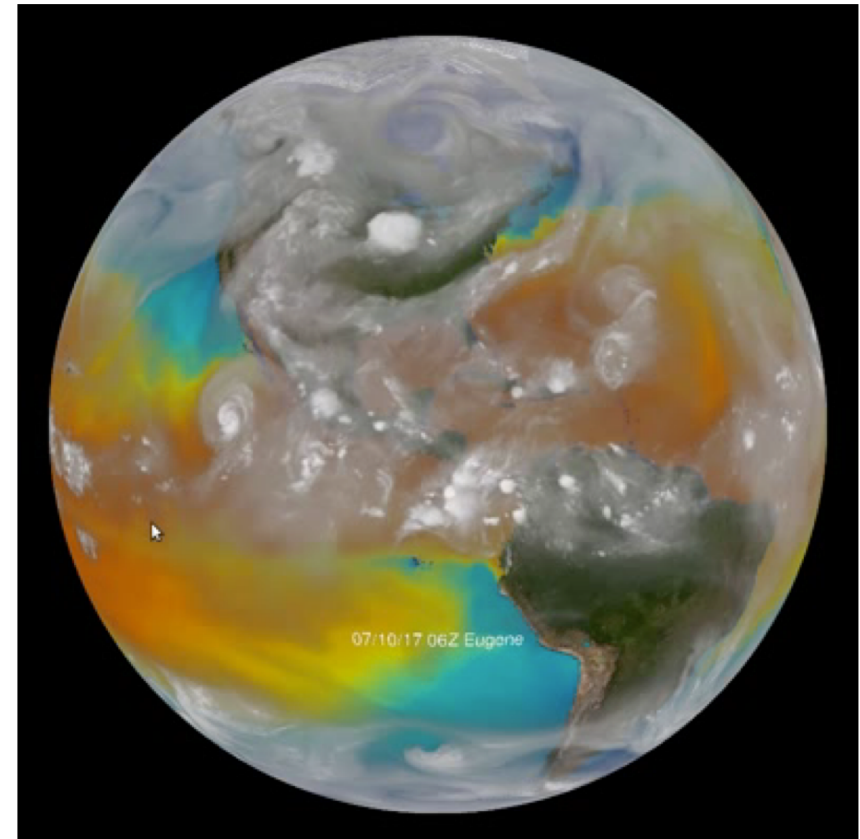
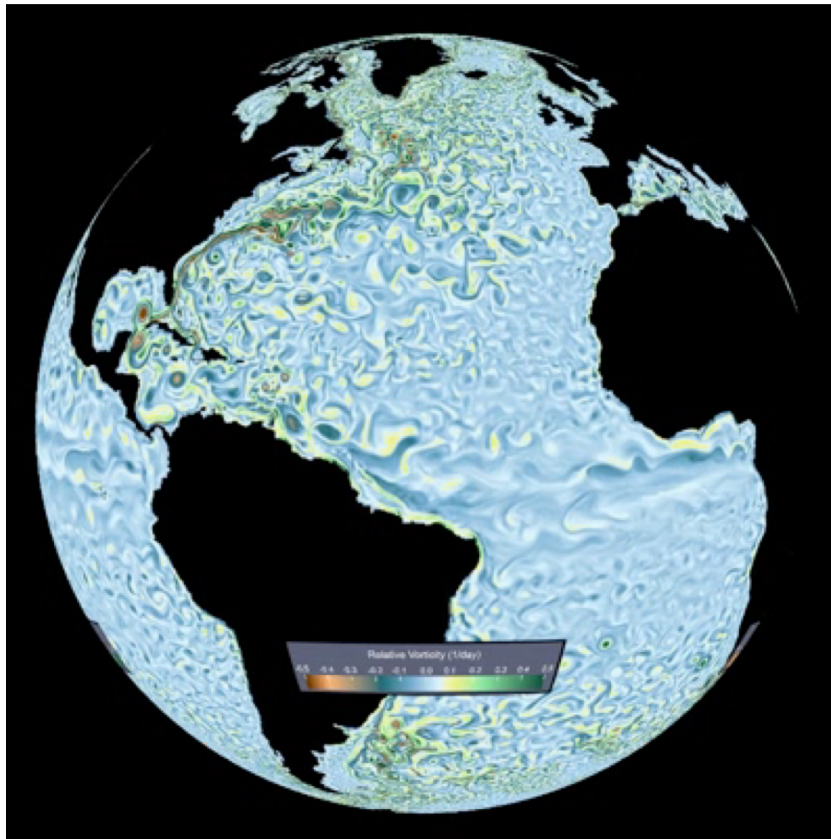
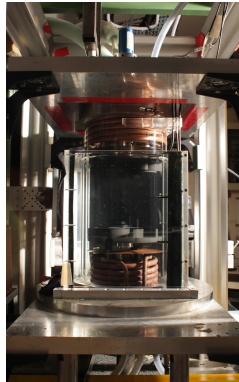
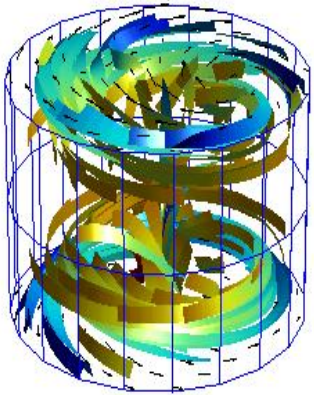


Figure 1.1: Vortices affect fluid behavior on all scales. (a) quantum vortices in a superfluid [130] (b) bathtub vortex [152] (c) tornado [109] (d) hurricane [106] (e) sun spot vortices [110] (f) spiral galaxy [105] (numbers approximate)

Vortices hierarchy on the Earth



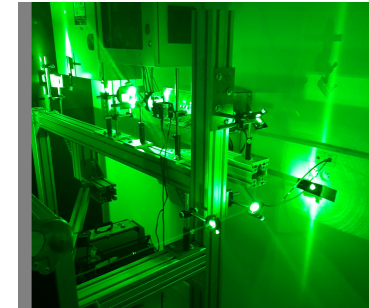
Vorticity hierarchy in the lab



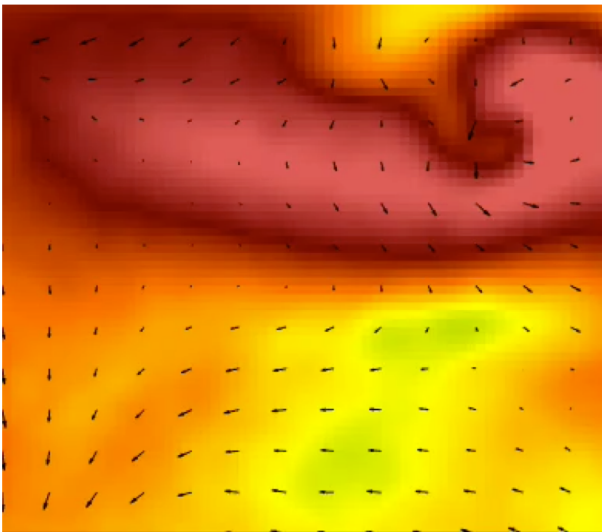
4 fast cameras



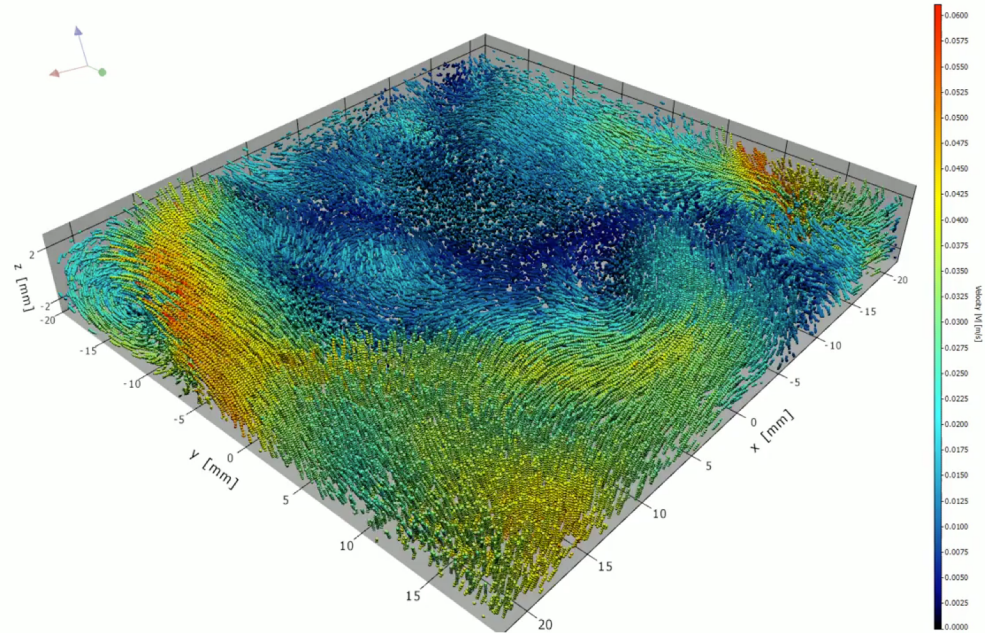
Laser



Turbulent flow in SPIV



Turbulent flow in 4D-PTV



Self-similarity of vortices in turbulence

Seymour Narrows,
Between Vancouver and Quadra Islands

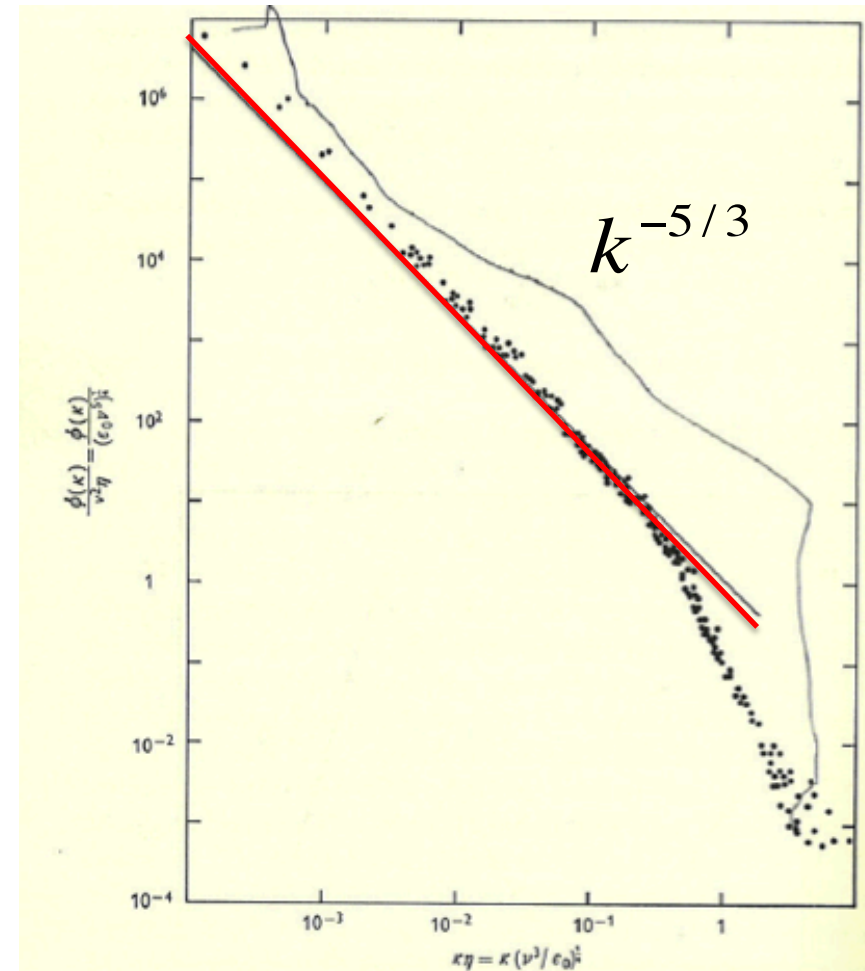


Fig. 6.2. The turbulence spectra, measured by Grant, Stewart and Moilliet (1962) and scaled according to the Kolmogorov parameters. The viscous dissipation rate ϵ_0 varied over a range of values of the order 100. The straight line represents variation as $k^{-5/3}$. The top few points are believed to be rather high on account of the low frequency heaving motions of the ship.

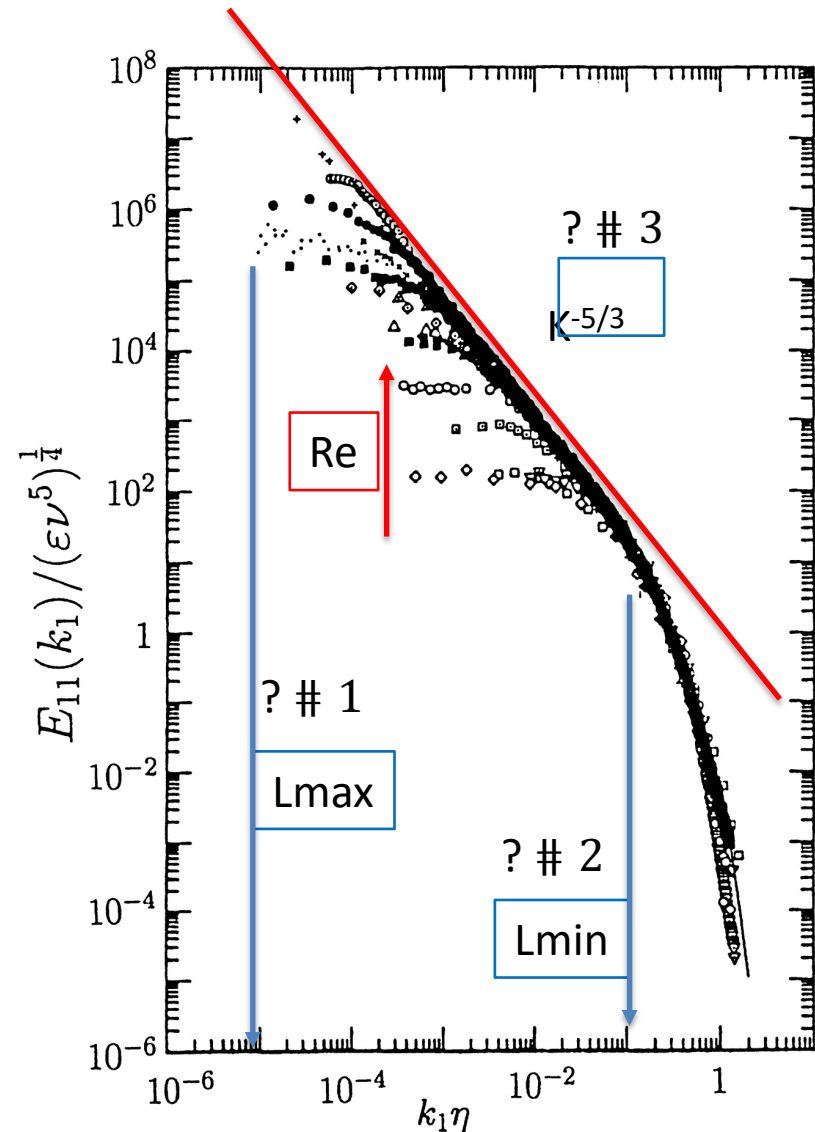
3 puzzles set by the vortices hierarchy



Number of degrees of freedom

$$N = \left(\frac{L_{\max}}{L_{\min}} \right)^3$$

3 more puzzles of turbulence
Practical importance for DNS and theories





Solving # 1 and # 2 puzzles:

How are L_{\max} and L_{\min} selected?

The forcing matters

$$\begin{aligned}\vec{\nabla} \cdot \vec{u} &= 0 \\ \partial_t \vec{u} + (\vec{u} \cdot \vec{\nabla}) \vec{u} &= -\frac{1}{\rho} \vec{\nabla} p + \nu \Delta \vec{u}\end{aligned}$$

Energy dissipation!



Without forcing, unique equilibrium state
 $u=0$

$$\begin{aligned}\vec{\nabla} \cdot \vec{u} &= 0 \\ \partial_t \vec{u} + (\vec{u} \cdot \vec{\nabla}) \vec{u} &= -\frac{1}{\rho} \vec{\nabla} p + \nu \Delta \vec{u} + \vec{F}\end{aligned}$$



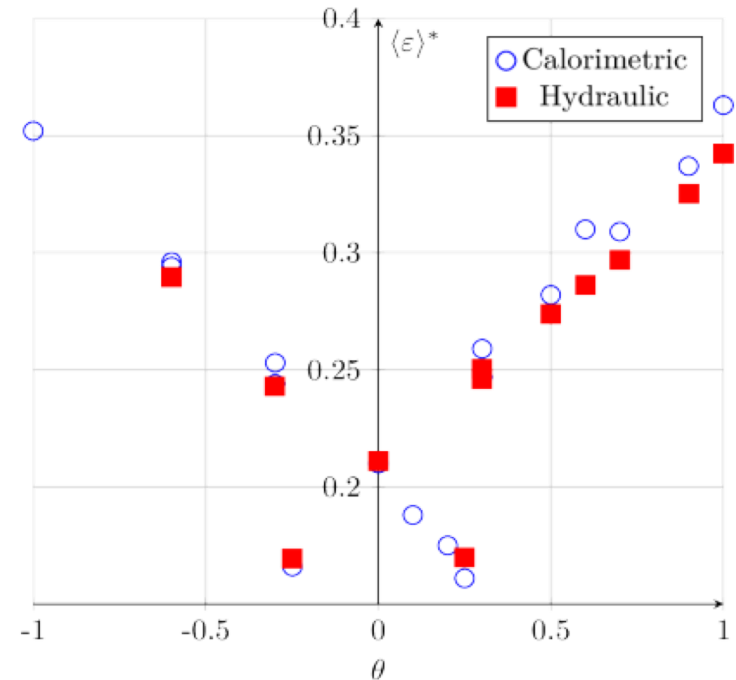
With forcing (possibly multiple) out-of-equilibrium
stationary state

The dissipation matters

$$\partial_t \int \frac{u^2}{2} dx^3 = \int F \cdot u dx^3 - \nu \int (\nabla u)^2 dx^3$$

Energy balance in a turbulent flow

$$dE = dW + dQ$$



$$P_{inj} = P_{diss} \equiv \rho L^3 \varepsilon$$

Cf Joule's experiment:

Work measured by Torques applied at Shafts

=Heat flux measured By keeping T constant

Solution of puzzle # 1

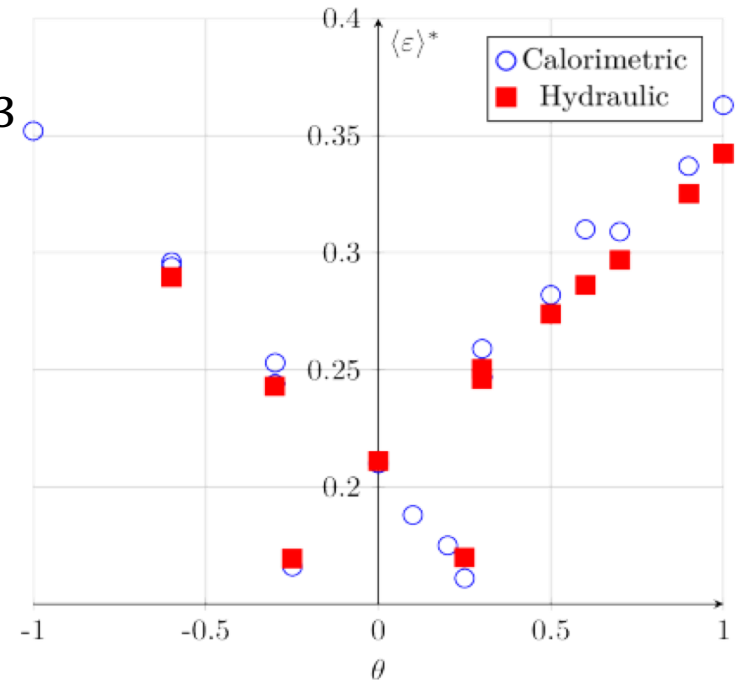
$$\partial_t \int \frac{u^2}{2} dx^3 = \int F \cdot u dx^3 - \nu \int (\nabla u)^2 dx^3$$

Energy balance in a turbulent flow



1st characteristic scale: L: size of the stirrer

Injection scale

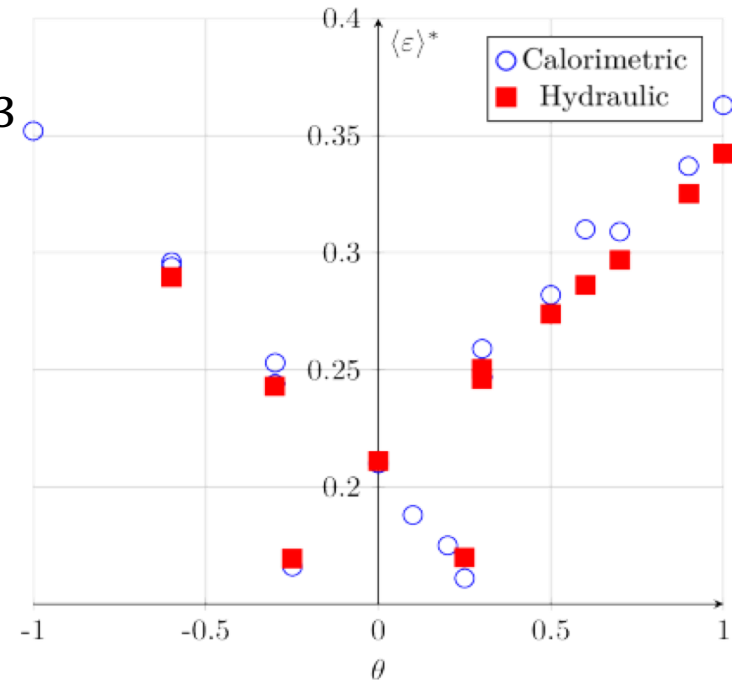


$$P_{inj} = P_{diss} \equiv \rho V \varepsilon$$

Solution of puzzle # 2

$$\partial_t \int \frac{u^2}{2} dx^3 = \int F \cdot u dx^3 - \nu \int (\nabla u)^2 dx^3$$

Energy balance in a turbulent flow



$$P_{inj} = P_{diss} \equiv \rho V \varepsilon$$

Dimensional analysis

$$\begin{aligned} \varepsilon &= [m^2 s^{-3}] \\ \nu &= [m^2 s^{-1}] \end{aligned}$$

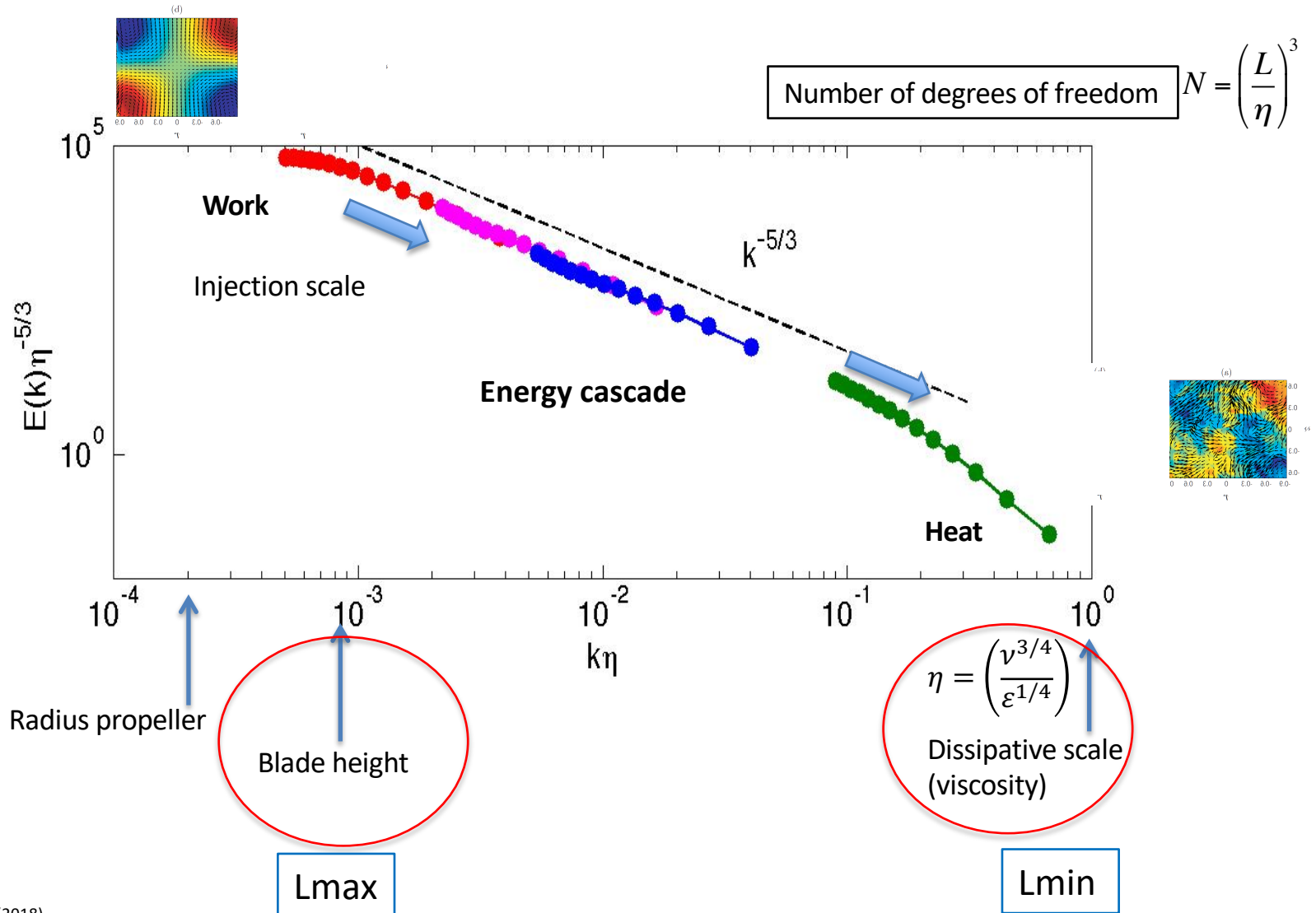


2nd characteristic scale

$$\eta = \left(\frac{\nu^3}{\varepsilon} \right)^{1/4}$$

Kolmogorov scale

Preliminary phenomenological picture



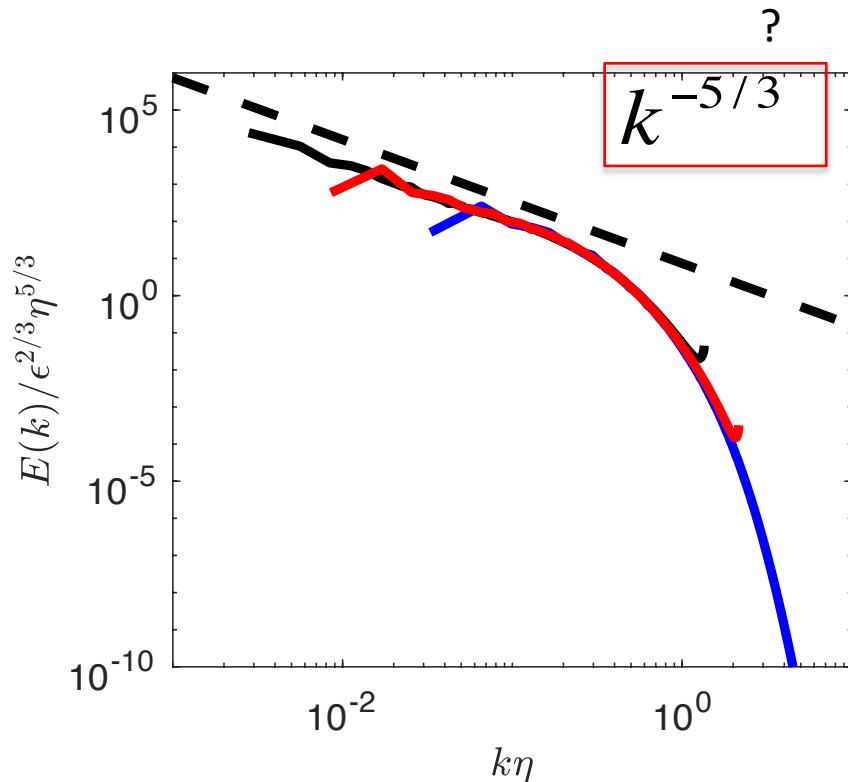


Solving # 3 puzzle:

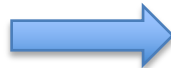
How is spectral exponent selected?

The spectral exponent is very robust but not given by dimensional analysis

Numerical simulations NSE



$$E(k) = [m^3 s^{-2}]$$

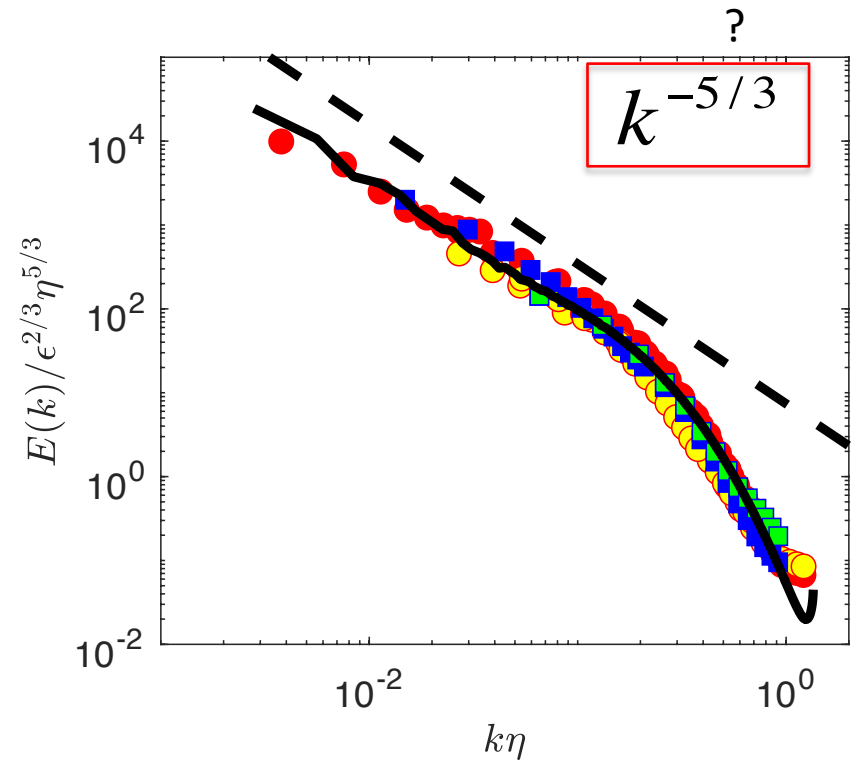


$$E(k) = \varepsilon^{2/3} \eta^{5/3} (k\eta)^{-1-2h}$$

How is the exponent $h=1/3$ selected?

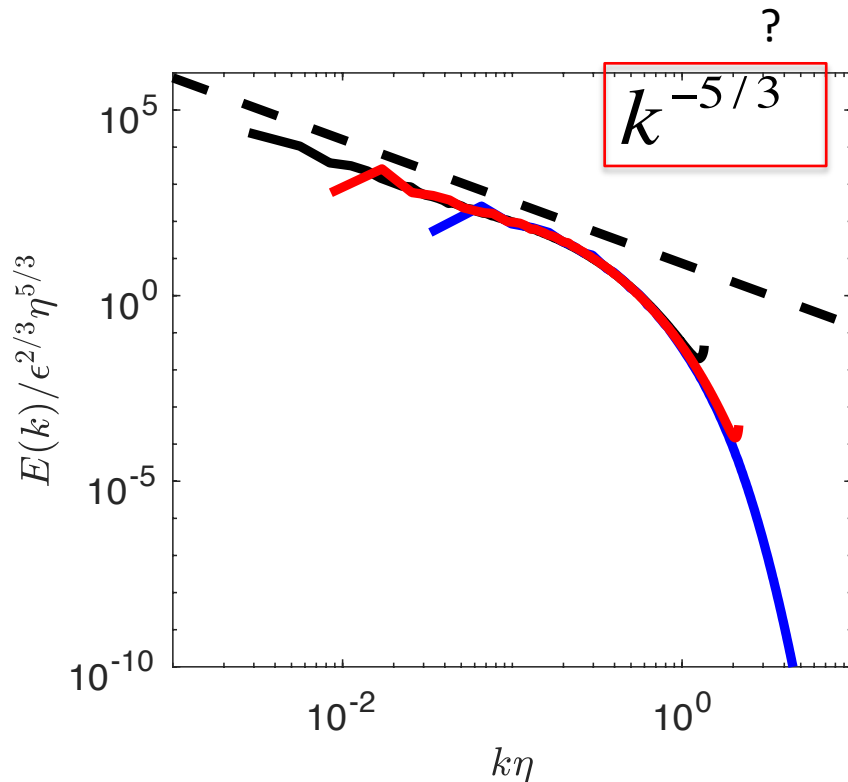
$$v(k) \approx k^{-h}$$

Experiments for a wide class of forcing conditions



The spectral exponent is very robust but not given by dimensional analysis

Numerical simulations NSE



By dimensional analysis

$$E(k) = [m^3 s^{-2}]$$

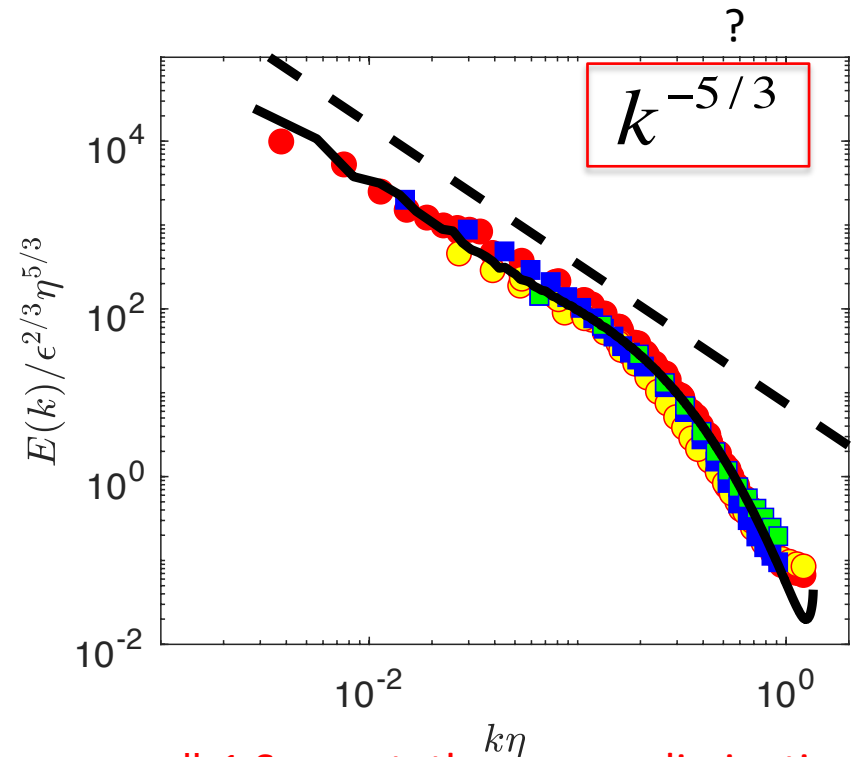


$$E(k) = \epsilon^{2/3} \eta^{5/3} (k\eta)^{-1-2h}$$

How is the exponent $h=1/3$ selected?

$$v(k) \approx k^{-h}$$

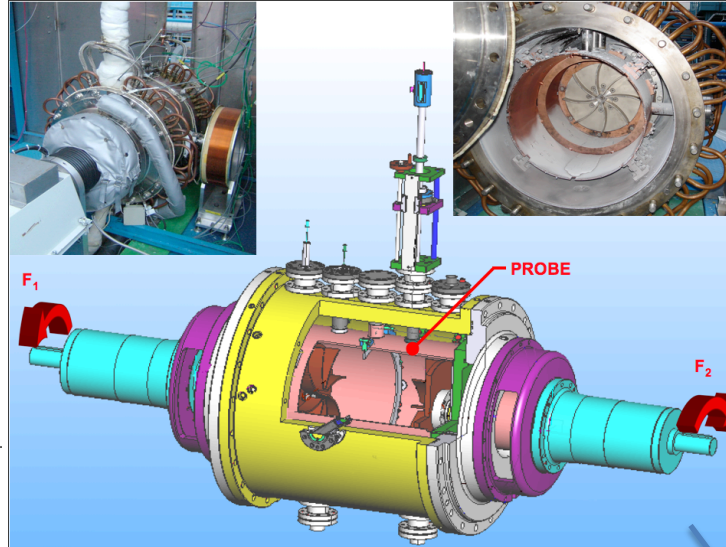
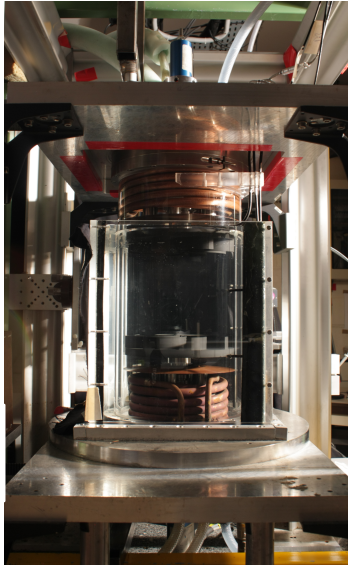
Experiments for a wide class of forcing conditions



1 Suspect: the energy dissipation

A closer look at energy dissipation

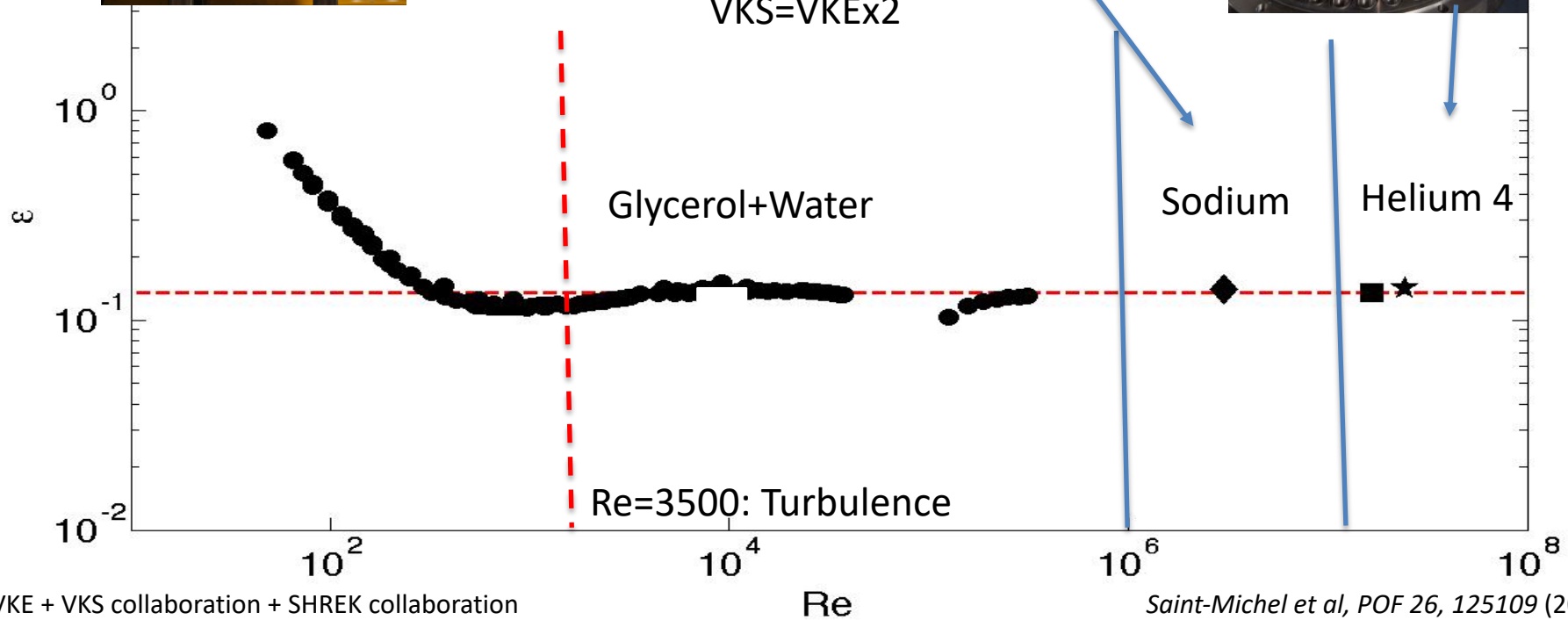
VKE



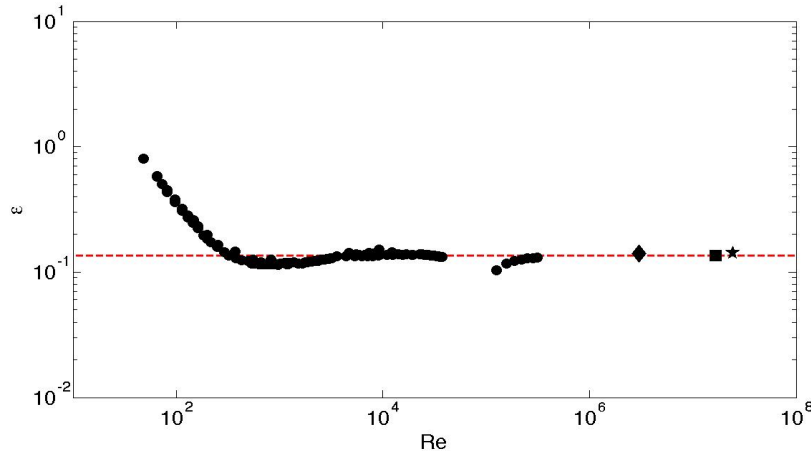
VKS=VKEx2



SHREK=VKEx4

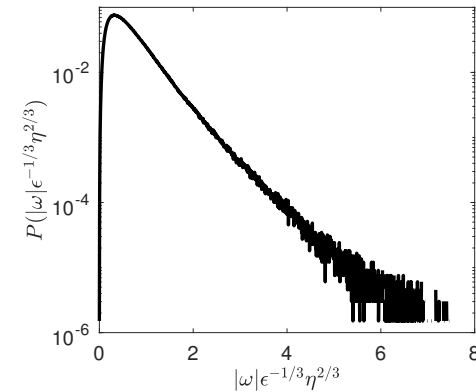
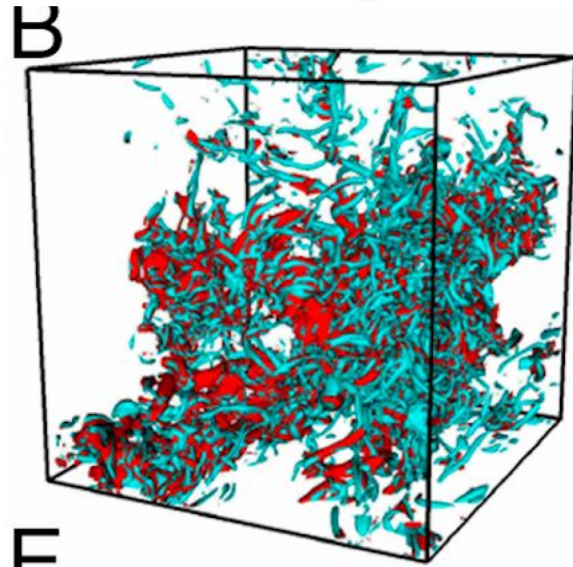


Where we meet again puzzle # 0



Non-dimensional energy dissipation per unit mass
is constant at large Reynolds

Independent of viscosity?

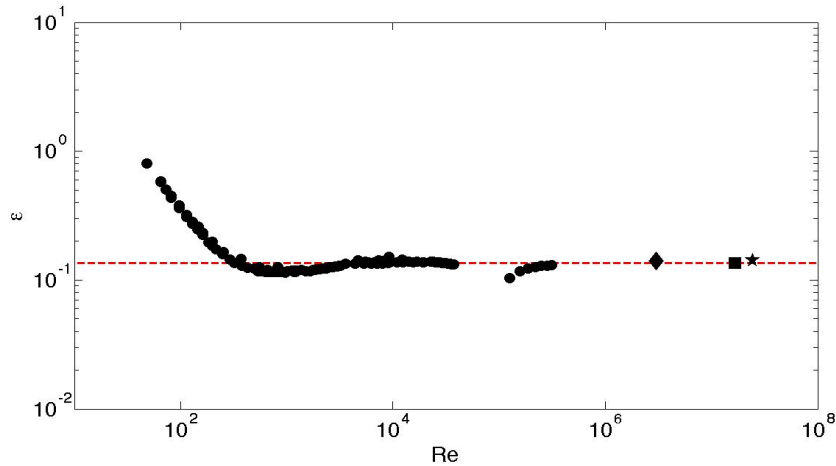


$$\epsilon = \nu \langle (\nabla u)^2 \rangle = \nu \langle \omega^2 \rangle \Rightarrow \langle \omega^2 \rangle \approx \frac{\epsilon}{\nu}$$

$$\lim_{\nu \rightarrow 0} \langle \omega^2 \rangle = \infty$$

Building of very large gradients at small scale... Singularity?
How to measure them/quantify them/understand this?

Local energy budget for irregular fields



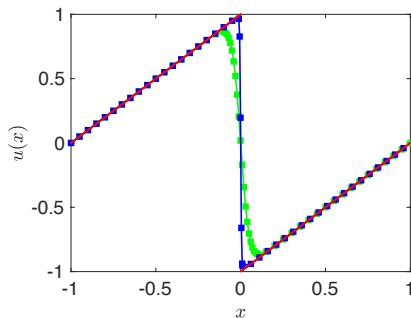
$$\frac{1}{2} \partial_t \mathbf{u}^2 + \operatorname{div} \left(\mathbf{u} \left(\frac{1}{2} \mathbf{u}^2 + p \right) - \nu \nabla \mathbf{u} \right) = D(u) - \nu |\nabla \mathbf{u}|^2$$

Duchon & Robert. Nonlinearity (2000),

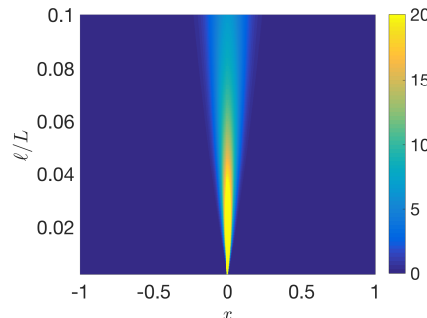
Inertial dissipation= singularity/large gradient detector!

$$D(u) = \lim_{\ell \rightarrow 0} \frac{1}{4} \int_{r \leq \ell} d^3 r \nabla \phi_\ell(r) \cdot \delta u_r |\delta u_r|^2$$

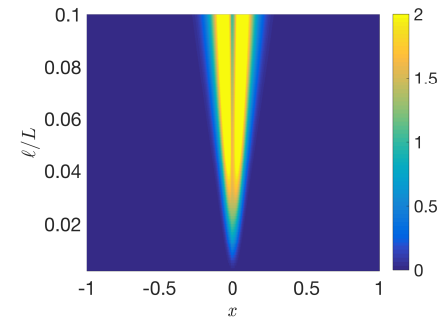
$\delta u = u(x+r) - u(x)$ Velocity increment



Without Viscosity



With Viscosity

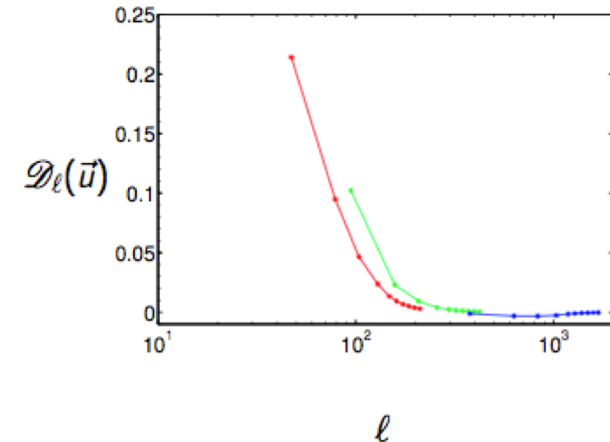
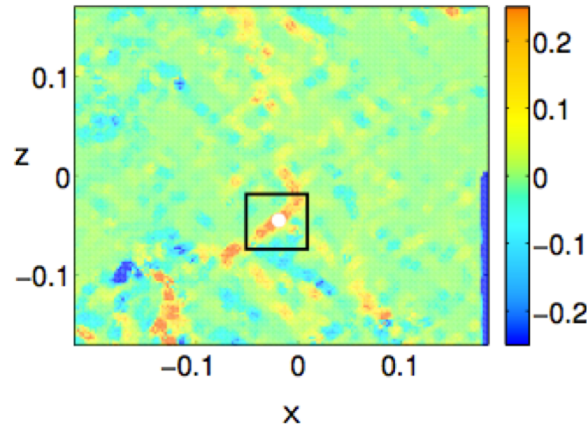
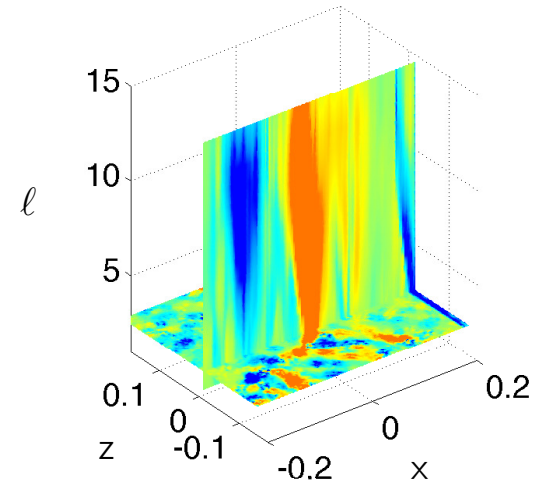
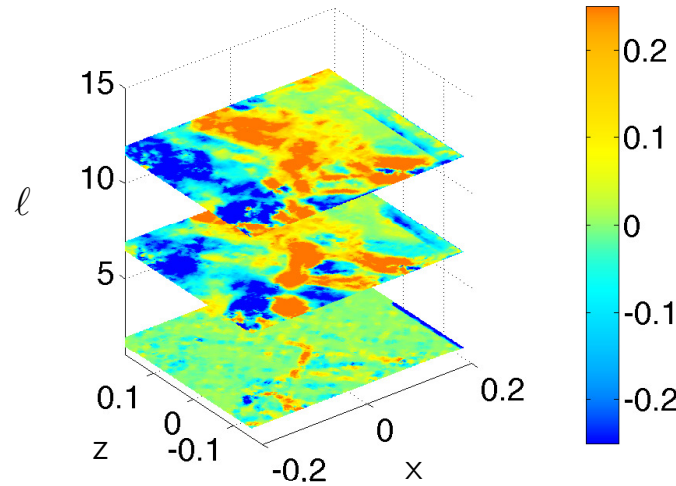


Applying the DR detector in von Karman flow

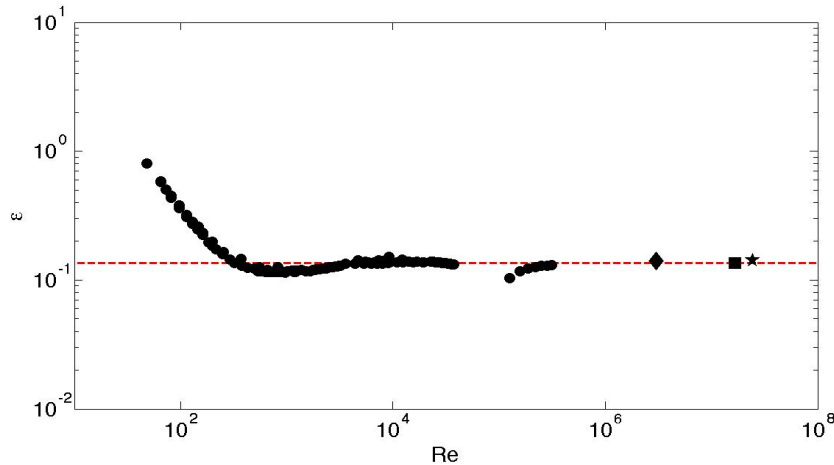
$$D_\ell(\mathbf{u}) = \frac{1}{4} \int_V d^3r (\nabla G_\ell)(\mathbf{r}) \cdot \delta \mathbf{u}(\mathbf{r}) |\delta \mathbf{u}(\mathbf{r})|^2,$$

$$G_\ell(r) = \frac{1}{N} \exp(-1/(1 - (r/2\ell)^2)),$$

-Kuzzay D. et al. (2017),
Nonlinearity



Solving puzzle # 3: Onsager's conjecture



”...in three dimensions a mechanism for complete dissipation of all kinetic energy, even without the aid of viscosity, is available.”

L. Onsager, 1949

$$\frac{1}{2} \partial_t \mathbf{u}^2 + \operatorname{div} \left(\mathbf{u} \left(\frac{1}{2} \mathbf{u}^2 + p \right) - \nu \nabla \mathbf{u} \right) = D(u) - \nu |\nabla \mathbf{u}|^2$$

Duchon & Robert. Nonlinearity (2000),

Inertial dissipation= singularity/large gradient detector!

$$D(u) = \lim_{\ell \rightarrow 0} \frac{1}{4} \int_{r \leq \ell} d^3 r \nabla \phi_\ell(r) \cdot \delta u_r |\delta u_r|^2$$

$$\delta u(\ell) \sim \ell^h \quad \text{In the limit of } \ell \approx 0$$

$$D(u)[x] \propto \lim_{\ell \rightarrow 0} \ell^{3h-1}$$

If $h > 1/3 \rightarrow$ Euler equation conserves energy,
Dissipation in Navier-Stokes by viscosity.

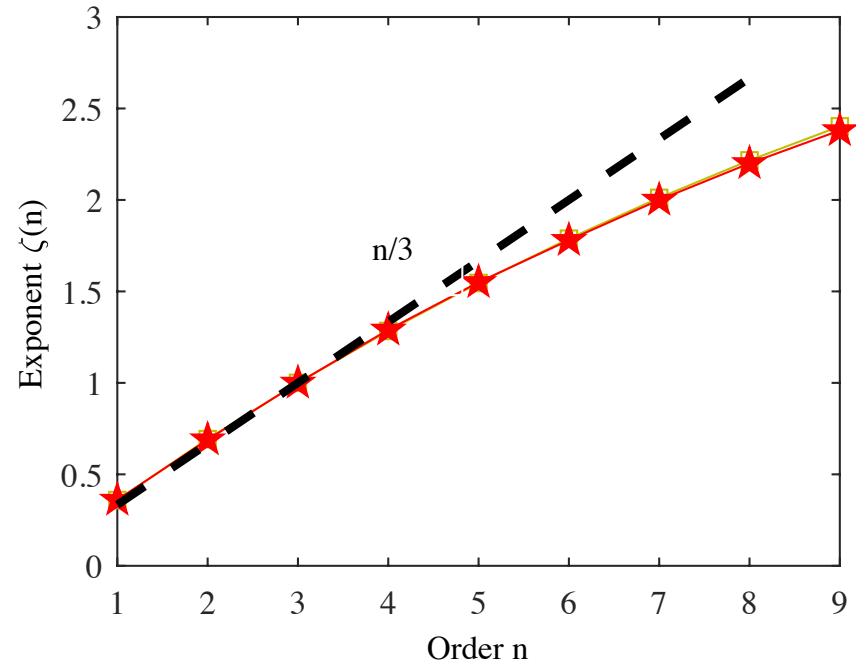
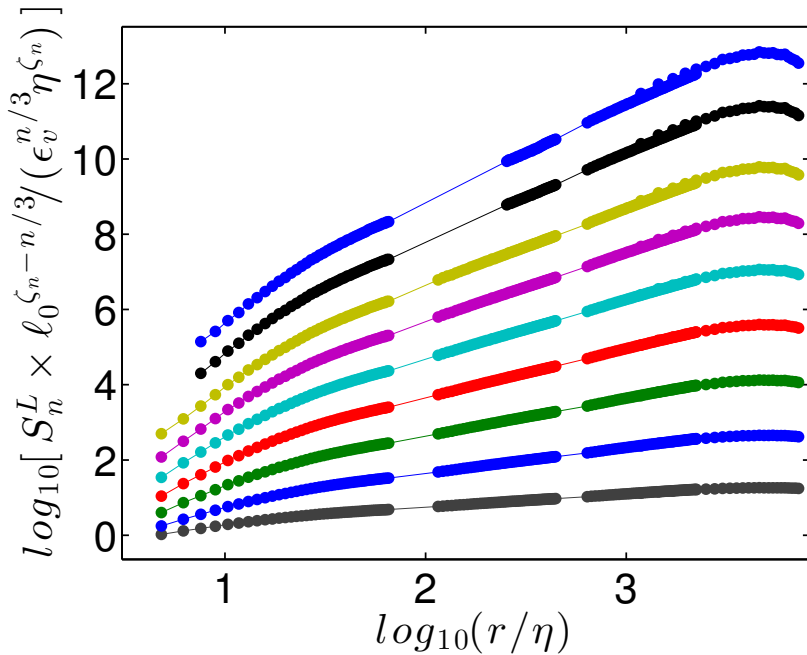
If $h \leq 1/3 \rightarrow$ Dissipation through irregularities (singularities)
Without viscosity !

Kolmogorov: $h=1/3$; $D=\varepsilon cte$

Kolmogorov spectrum exponent is the « critical value » !!!!

Is this the only one? -> look statistics of velocity increments

Puzzle # 4: $h = \frac{1}{3}$ is not unique!



$$\delta u \sim \ell^{1/3} \Rightarrow S_n(\ell) = \langle (\delta u)^n \rangle = \ell^{n/3}$$

Not observed!

Kolmogorov spectrum exponent $h=1/3$ is the « critical value » !!!!
It is not unique...

What is the physics involved?

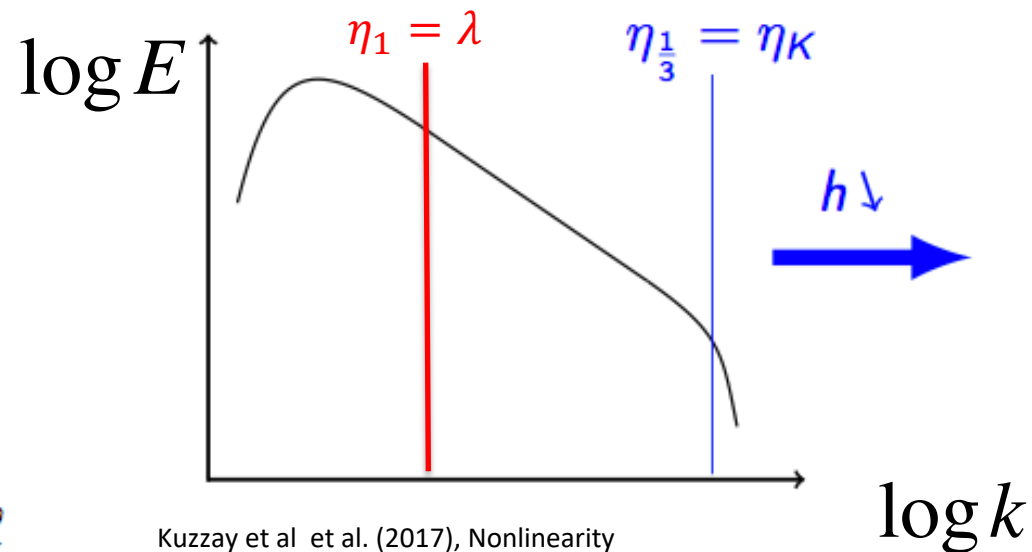
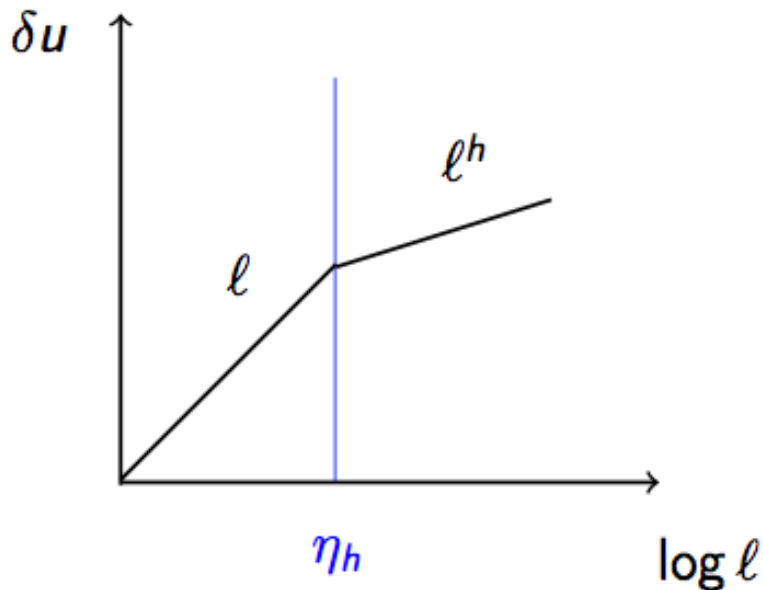
Local energy balance in turbulence

$$\begin{aligned}\partial_t E^\ell + \partial_j J_j^\ell &= -\frac{1}{4} \int \nabla \phi^\ell(\xi) \cdot \delta \mathbf{u} (\delta u)^2 d\xi + \nu \partial^2 E^\ell \\ &\equiv -D_\ell^I - D_\ell^\nu,\end{aligned}$$

$$\delta u \approx \ell^h \Rightarrow D_\ell^I \approx \ell^{3h-1} \text{ and } D_\ell^\nu \approx \nu \ell^{2h-2}$$

$$\frac{D_\ell^I}{D_\ell^\nu} \approx \frac{\ell^{3h-1}}{\nu \ell^{2h-2}} \approx \frac{\ell^{h+1}}{\nu} \Rightarrow \eta_h \approx \text{Re}^{-1/(1+h)}$$

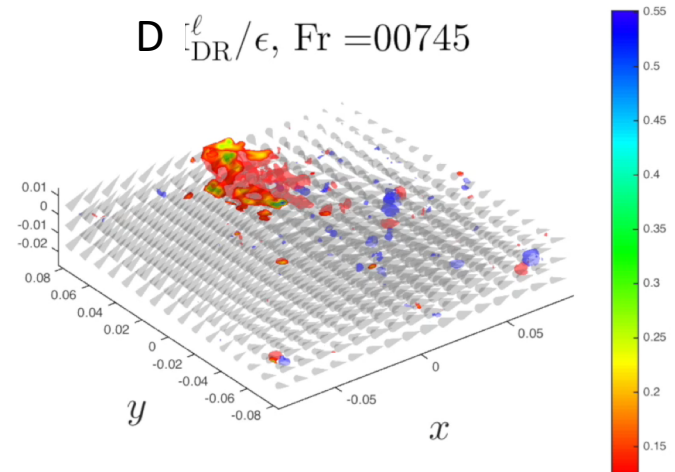
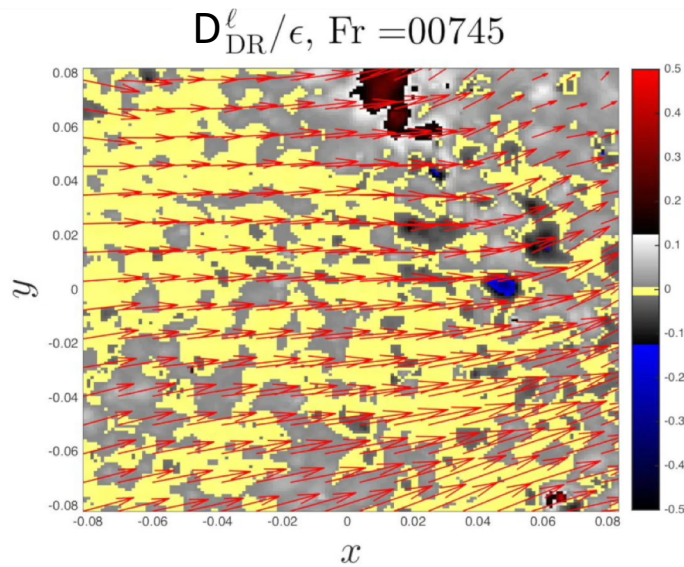
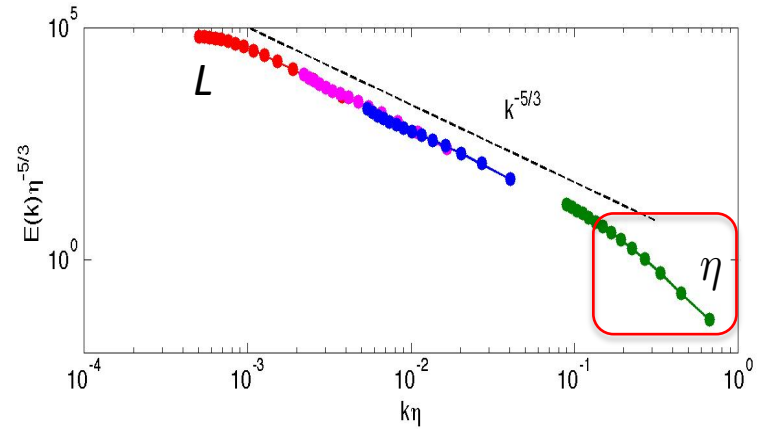
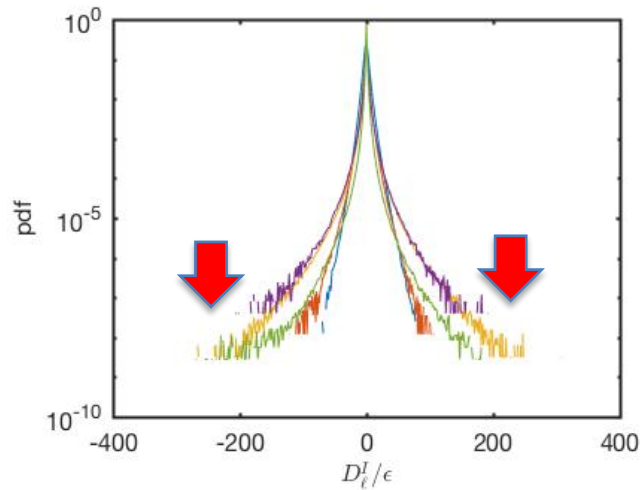
$$\eta_{-1}=0$$





Solving puzzle # 4:
Is $h=1/3$ the only possible exponent?
Look at the physics of turbulence at
Kolmogorov scale

Large events of $D(u)$

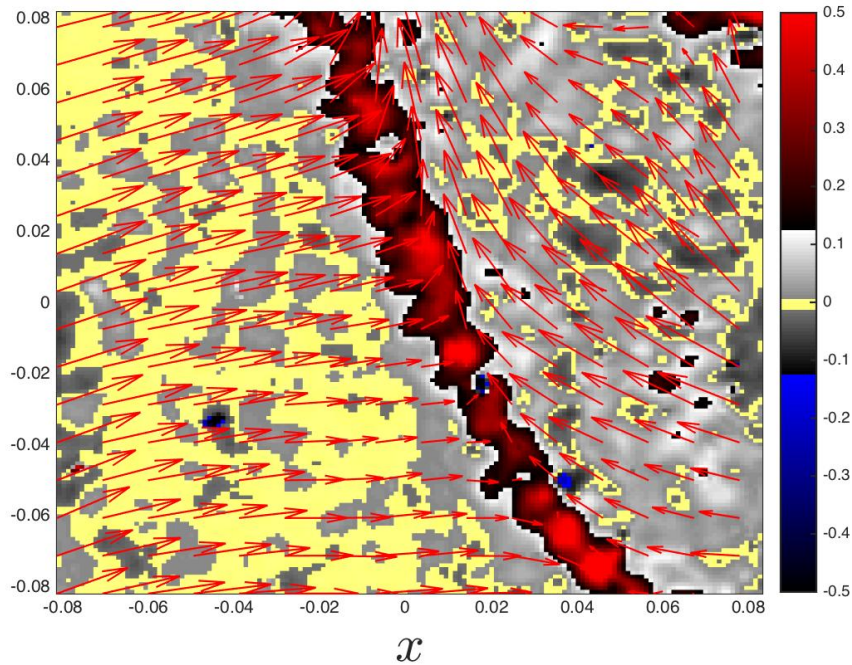


Color map: M. Farge, 1990 *L'Aéronautique et l'Astronautique*, **140**, 24-33

ANR Exploit: collaboration Saclay/LML: Debue, Shukla, Cuvier, Saw, Wiertel, Padilla, Daviaud, Dubrulle, Foucaut, Laval

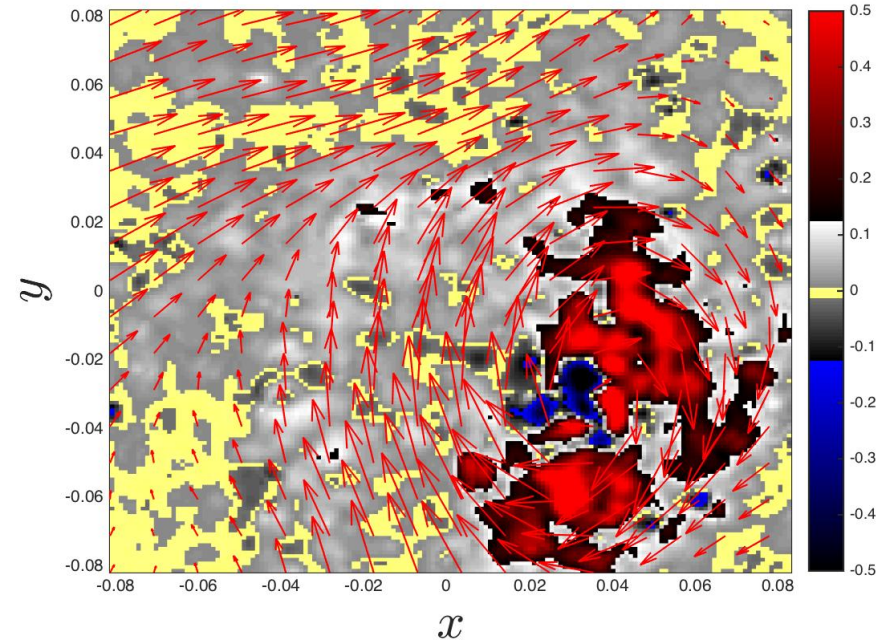
2D Structure of large events of inertial dissipation

$$D_{DR}^{\ell}/\epsilon, Fr = 00207$$



front

$$D_{DR}^{\ell}/\epsilon, Fr = 00754$$

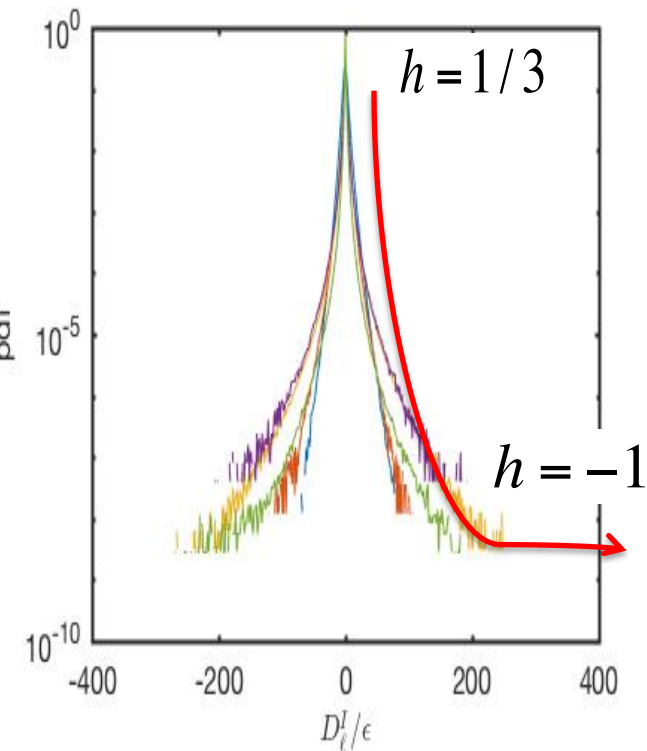


Spiral

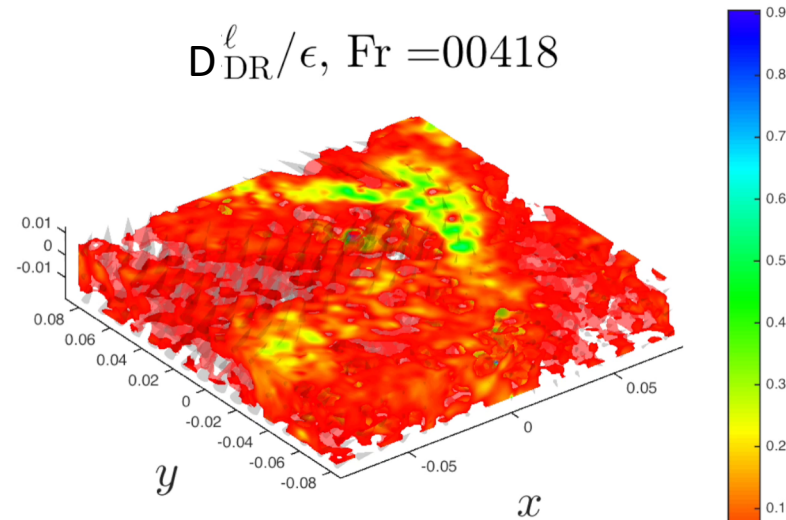
Color map: M. Farge, 1990 *L'Aéronautique et l'Astronautique*, **140**, 24-33

Saw, Kuzzay et al. (2016), *Nature-Comm.* 7

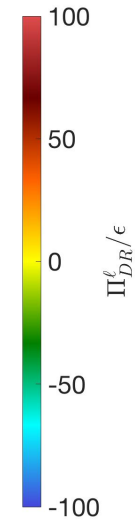
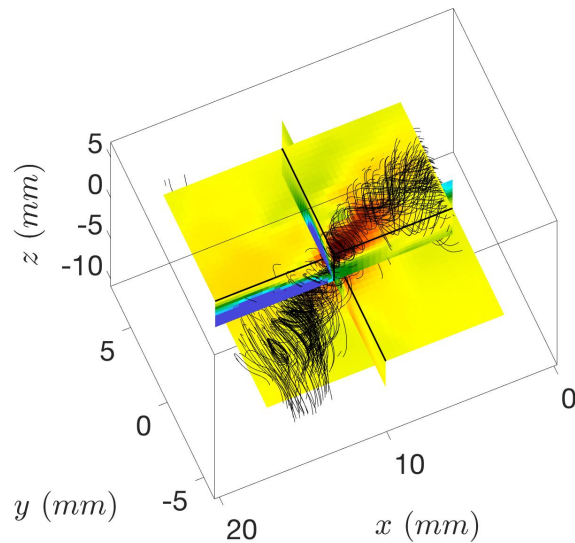
Structure of large events of inertial dissipation



Hierarchy of h !

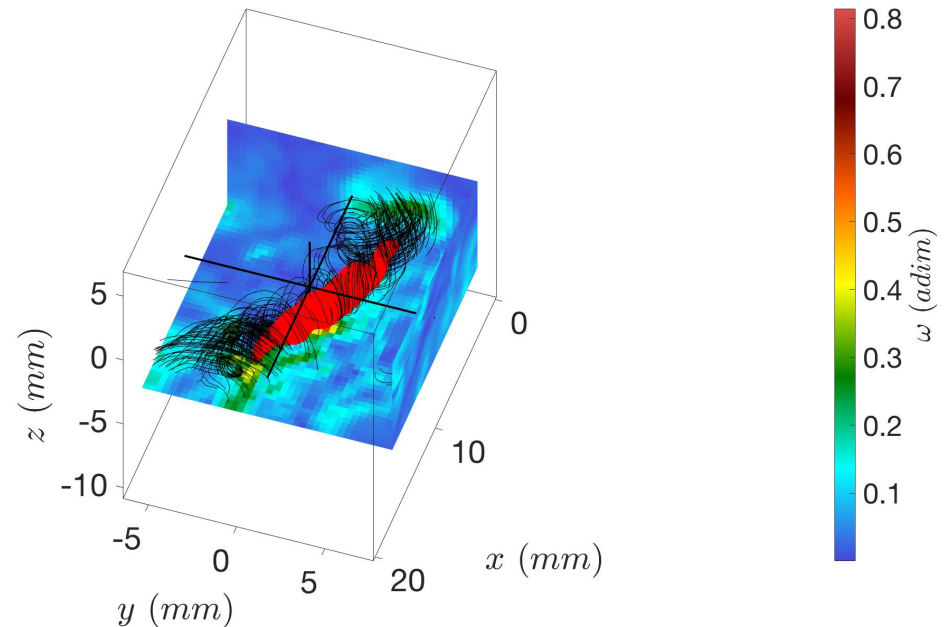


3D Structure of large events of inertial dissipation



Swirling filament

High enstrophy in the
Vicinity of the event



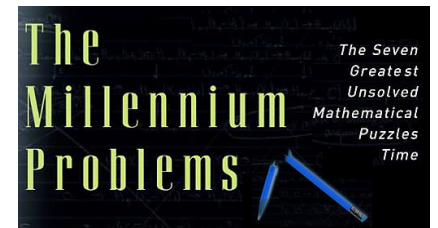


Solving # 0 puzzle:

Are these footprints of singularities?

Open basic question:

*Navier-Stokes equations are they well-posed?
Are there any singularities?*



A possible suspect of NS singularity

Symmetry $(t, x, u) \rightarrow (\gamma^2 t, \gamma x, \gamma^{-1} u) (\gamma \neq 0)$

$u(\gamma^2 t, \gamma x) = \gamma^{-1} u(t, x)$ **homogeneous solutions of NS of degree -1**

Axisymmetric case: classification started by Landau in 1944

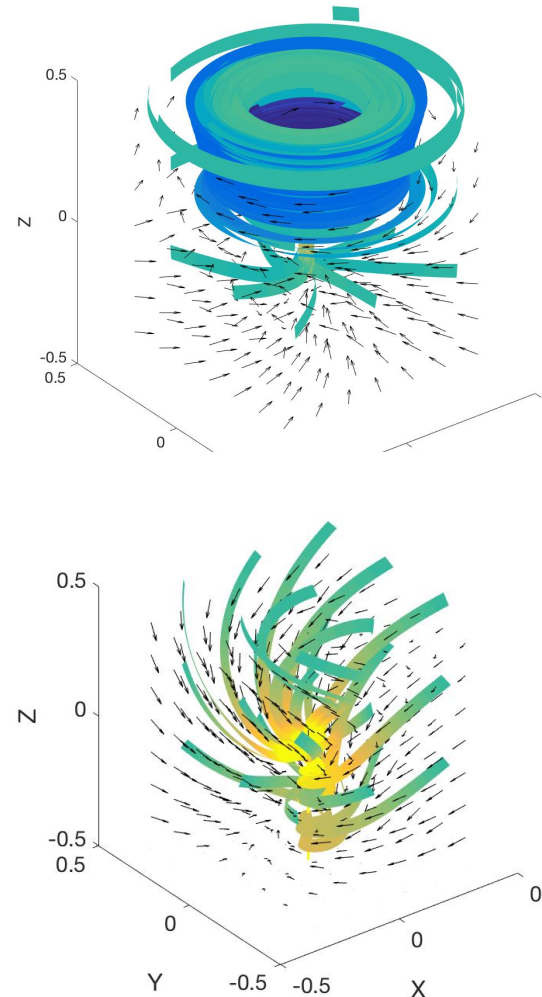
$$U_r = -\frac{1}{r} \left\{ 1 + \frac{2(\gamma + 1)}{(\gamma + 1) \ln\left(\frac{1+\cos\theta}{2}\right) - 2} \right. \quad (1)$$

$$+ \left. \frac{1 - \cos\theta}{1 + \cos\theta} \frac{3(\gamma + 1)}{\left((\gamma + 1) \ln\left(\frac{1+\cos\theta}{2}\right) - 2\right)^2} \right\},$$

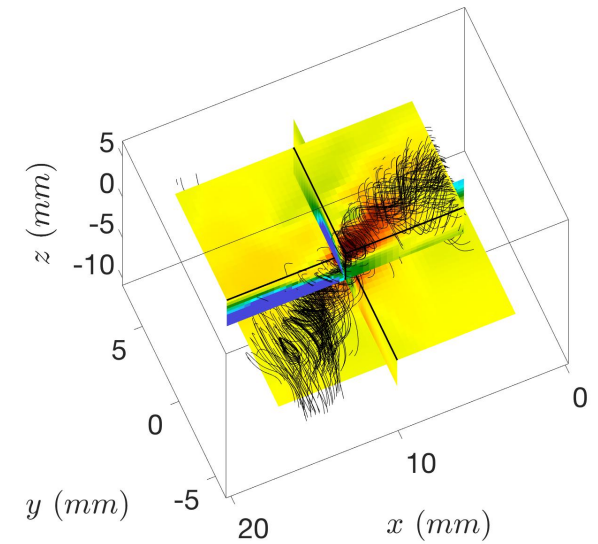
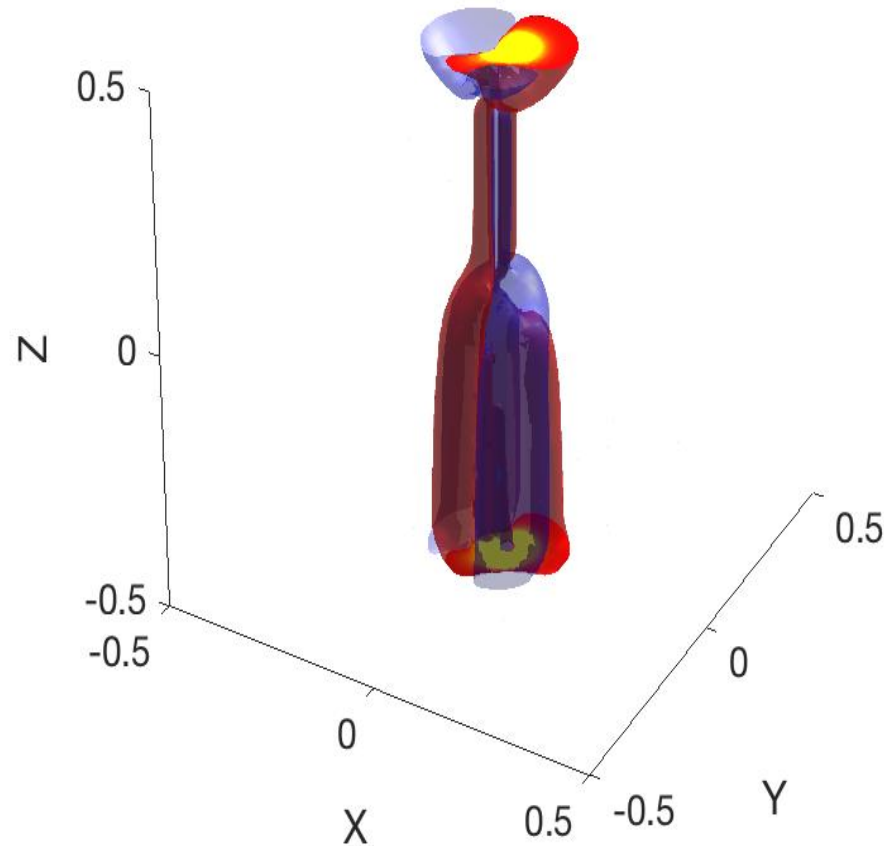
$$U_\theta = \frac{1}{r} \left(\frac{1 - \cos\theta}{\sin\theta} \right) \left(1 + \frac{2(\gamma + 1)}{(\gamma + 1) \ln\left(\frac{1+\cos\theta}{2}\right) - 2} \right)$$

$$U_\phi = \begin{cases} \frac{1}{r \sin(\theta)} (b \ln(1 + \cos\theta) + a) & \text{for } \gamma = -1 \\ \frac{1}{r \sin\theta} \left(\frac{b}{\ln(1+\cos\theta)} + a \right) & \text{for } \gamma > -1 \end{cases}$$

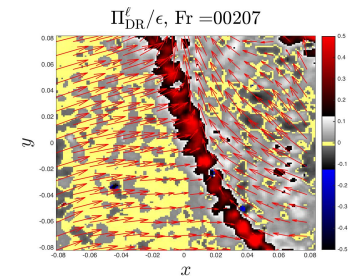
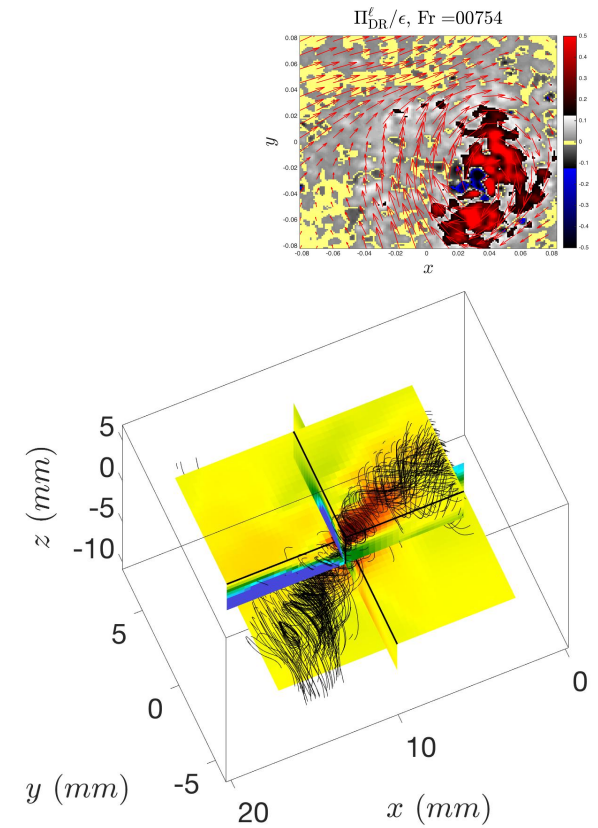
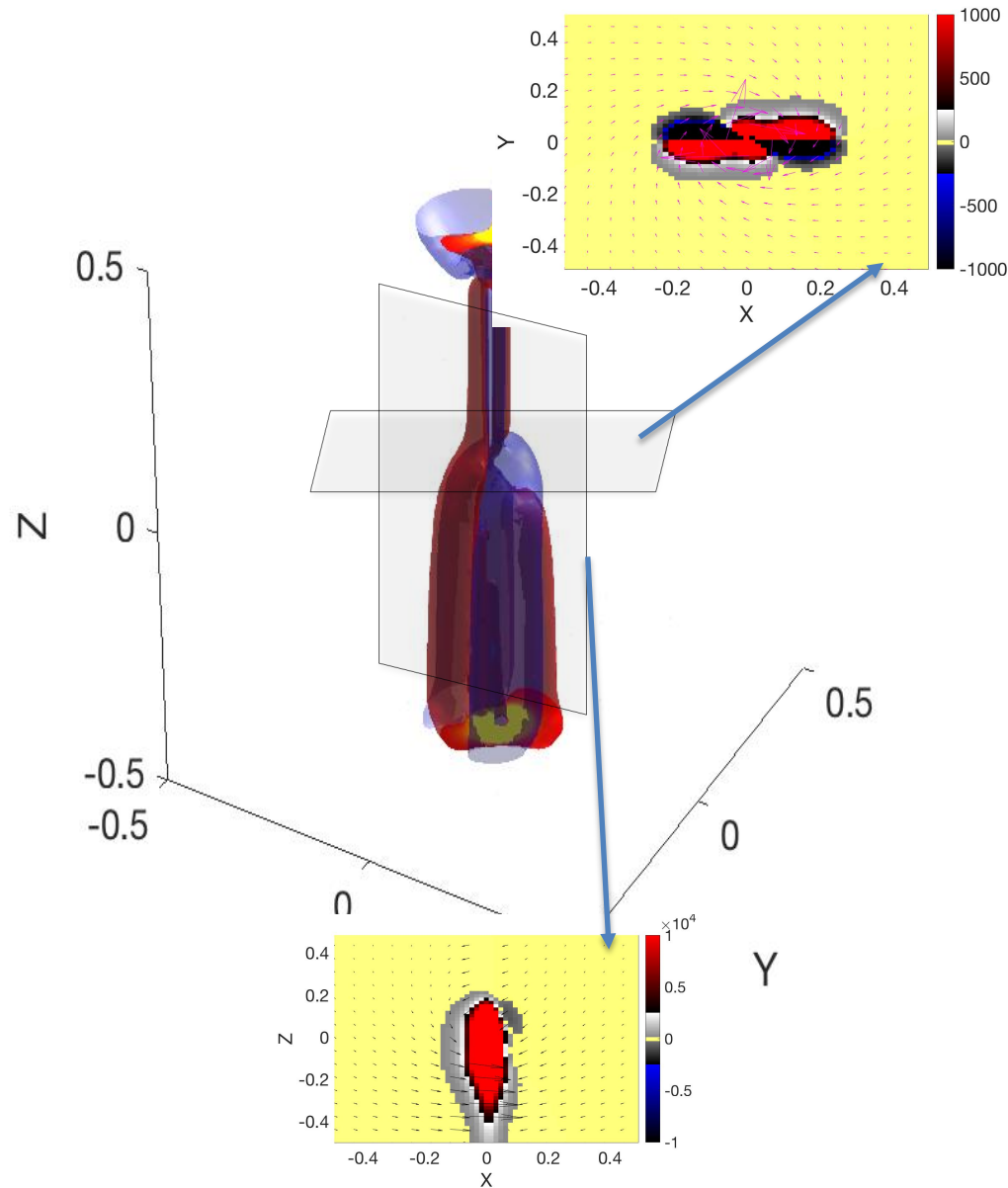
where $\gamma \geq 1$, a , et b are free parameters. Note singular behaviour for $\theta = \pi$.



3D footprints of the -1 degree homogeneous solution



2D footprints of the -1 degree homogeneous solution



Conclusions and perspectives

- Turbulence is plagued with puzzles of fundamental importance. One of them is connected with possible quasi-singularities or singularities of the NSE.
- We have found possible footprints of such quansi singularities. They are linked with coherent structures, living below the Kolmogorov scale,
- -> **Impact on DNS of NSE!!!!**

Impact on DNS of turbulence



$Re=6 \times 10^3$
 $L=10 \text{ cm}$

Resolution of events with exponent h

$$\eta_h = L Re^{-1/(1+h)}$$

$$N = \left(\frac{L}{\eta_h} \right)^3 \sim Re^{3/(1+h)}$$

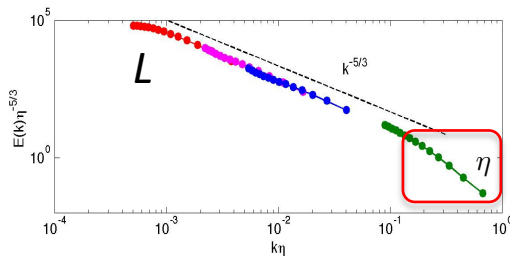
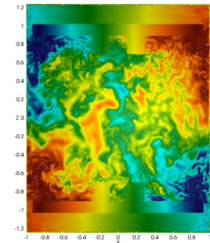
Kolmogorov

$$\eta_{1/3} = L Re^{-3/4}$$

$$N = \left(\frac{L}{\eta_{1/3}} \right)^3 \sim Re^{9/4}$$

$$\eta = 0.4 \text{ mm}$$

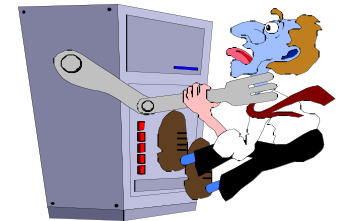
$$N \sim 10^7$$



MFR with $h_{min}=-0.2$

$$\eta_{-1/5} = Re^{-5/4} = \eta Re^{-1/2} = 5 \cdot 10^{-3} \text{ mm}$$

$$N = \left(\frac{L}{\eta_{-1/5}} \right)^3 \sim 8 \cdot 10^{12}$$



Conclusions and perspectives

- Turbulence is plagued with puzzles of fundamental importance. One of them is connected with possible quasi-singularities or singularities of the NSE.
- We have found possible footprints of such quansi singularities. They are linked with coherent structures, living below the Kolmogorov scale,
- -> **Impact on DNS of NSE!!!!**
- **Perspective:**
The study of the dynamics and properties of these structures is underway, using 4D-PTV
We are currently building a larger experiment to explore sub-Kolmogorov regimes, to look for more singular structures).

Prespectives: GVK experiment

20 cm



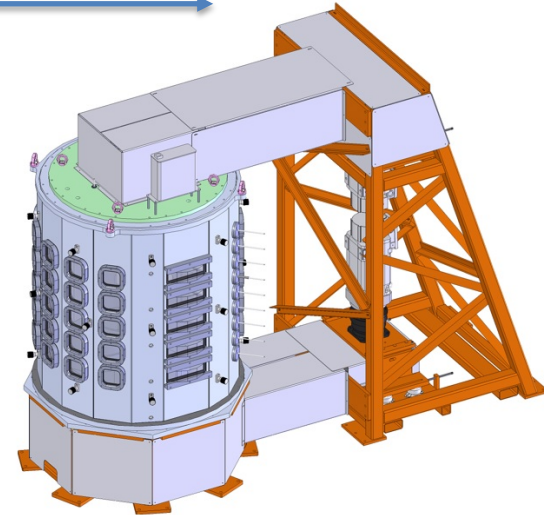
Present experiment

$R=10\text{ cm}$

$Re=10^6$
 $\eta=0,01\text{ mm}$

$Re=6 \times 10^3$
 $\eta=0,4\text{ mm}$
 $\Delta x \sim \eta$

1 m



GVK experiment $R=50\text{ cm}$

$Re=10^6$
 $\eta=0,05\text{ mm}$

$Re=6 \times 10^3$
 $\eta=2\text{ mm}$
 $\Delta x \sim \eta/5$

Possibility to explore sub-Kolmogorov scales

Detection of stronger velocity gradients and quasi-singularities

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VKS collaboration

SHREK collaboration

