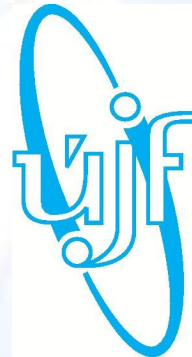


Recent applications of equation of motion phonon method



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**Shapes and Symmetries in Nuclei: from Experiment
to Theory (SSNET'18),
Gif-sur-Yvette, November 2018**

Equation of Motion Phonon Method

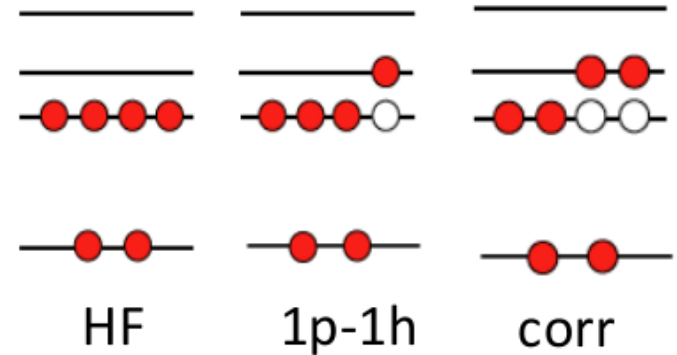
Equation of Motion Phonon Method (EMPM):

Features:

- nuclear **ground state** properties
- energy **spectra**
- **collective** excitations
- **wide-range applicability** (across the nuclear chart)
- exact treatment of **Pauli principle** (unlike the methods based on RPA)
- applicable on any nuclear Hamiltonian but usually **realistic Hamiltonian** is adopted

Applications:

- EMPM first developed for **even-even nuclei** *Phys. Rev. C 85 014313 (2012)*, *Phys. Rev. C 90 014310 (2014)*, *Phys. Rev. C 92 054315 (2015)*
- **quasiparticle** formulation of **EMPM** for open-shell nuclei *Phys. Rev. C 93 044314 (2016)*
- **EMPM** extended to **even-odd** nuclei *Phys. Rev. C 94 061301 (2016)*, *Phys. Rev. C 95 034327 (2017)*
- extension of **EMPM** to **hypernuclei** *in progress...*



EMPM

Hilbert space – divided into subspaces

$$\mathcal{H} = \mathcal{H}_0 \oplus \mathcal{H}_1 \oplus \mathcal{H}_2 \oplus \dots \oplus \mathcal{H}_n$$

HF – Hartree-Fock state (nucleons occupy lowest single-particle levels)

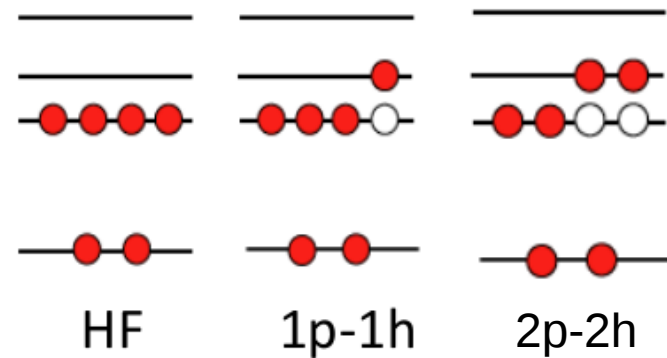
1p-1h = 1particle – 1hole excitation of HF

2p-2h = 2particle – 2hole excitation of HF

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np-nh = nparticle – nhole excitation of HF



Instead of multiple **particle-hole** excitations we can excite multiple **TDA phonons**

Tamm-Dancoff (TDA) phonons

$$O_{\nu}^{\dagger} = \sum_{ph} c_{ph}^{\nu} a_p^{\dagger} a_{\bar{h}}$$

Phonons = linear combination of 1p-1h excitations
can represent **collective modes**

$$\mathcal{H}_0 = \{|HF\rangle\}$$

$$\mathcal{H}_1 = \{O_{\nu_1}^{\dagger} |HF\rangle\}$$

$$\mathcal{H}_2 = \{O_{\nu_1}^{\dagger} O_{\nu_2}^{\dagger} |HF\rangle\}$$

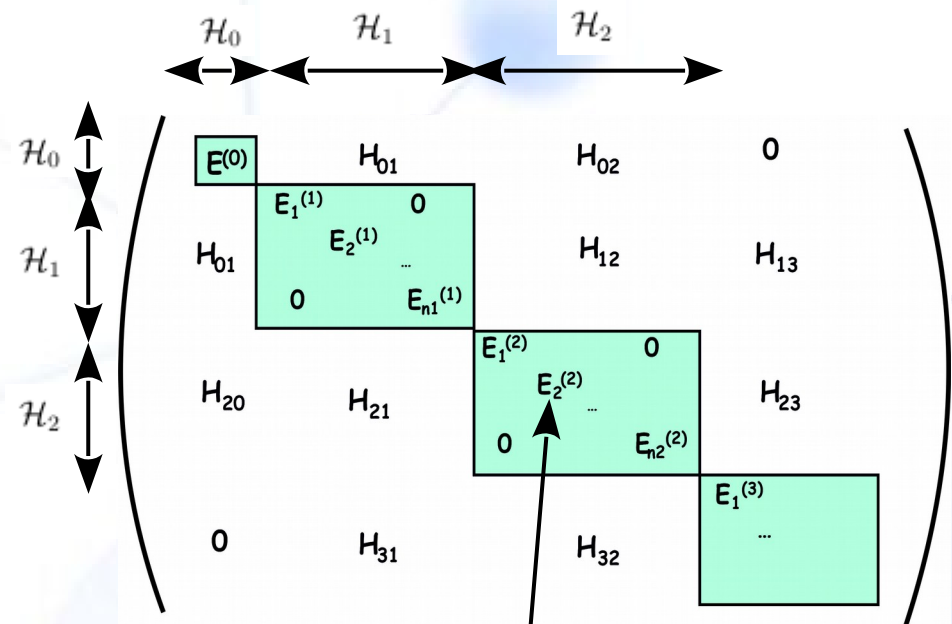
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$$\mathcal{H}_n = \{O_{\nu_1}^{\dagger} O_{\nu_2}^{\dagger} \dots O_{\nu_n}^{\dagger} |HF\rangle\}$$

EMPM

$$\begin{aligned} \mathcal{H}_0 &= \{|HF\rangle\} \\ \mathcal{H}_1 &= \{O_{\nu_1}^\dagger |HF\rangle\} \\ \mathcal{H}_2 &= \{O_{\nu_1}^\dagger O_{\nu_2}^\dagger |HF\rangle\} \\ &\vdots \\ \mathcal{H}_n &= \{O_{\nu_1}^\dagger O_{\nu_2}^\dagger \dots O_{\nu_n}^\dagger |HF\rangle\} \end{aligned}$$



the total **Hamiltonian** mixes configurations from different **Hilbert subspaces**

Equation of Motion (EoM) – recursive eq. to solve **eigen-energies** on each **i-phonon** subspace while knowing the **(i-1)-phonon** solution

$$\langle i, \beta_i | [\hat{H}, O_\nu^\dagger] | i-1, \alpha_{i-1} \rangle = (E_{\beta_i}^i - E_{\alpha_{i-1}}^{i-1}) \langle i, \beta_i | O_\nu^\dagger | i-1, \alpha_{i-1} \rangle$$

non-diagonal blocks of **Hamiltonian** calculated from amplitudes

$$\langle i, \beta_i | O_\nu^\dagger | i-1, \alpha_{i-1} \rangle$$

we diagonalize the total **Hamiltonian**

Ground State Correlations

NN interaction - χ NNLO_{opt}

A. Ekström et al., **PRL 110**, 192502 (2013)

2-phonon correlations in the g.s.

$$|\Psi_{g.s.}\rangle \approx C_{HF}^{g.s.} |HF\rangle + \sum_{\mu_2} C_{\mu_2}^{g.s.} |i=2, \mu_2\rangle$$

G. De Gregorio, J. Herko, F. Knapp, N. Lo Iudice, P. Veselý, **PRC 95**, 024306 (2017)

TABLE I. Binding energies per nucleon. The EMPM value for ⁴⁰Ca was obtained for $N_{\max} = 8$, which is not an extremal point.

^A X	BE/A (MeV)			Exp.
	HF	PT	EMPM	
⁴ He	3.96	7.07	6.67	7.07
¹⁶ O	3.22	8.29	6.77	7.98
⁴⁰ Ca	4.00	9.77	7.02	8.55

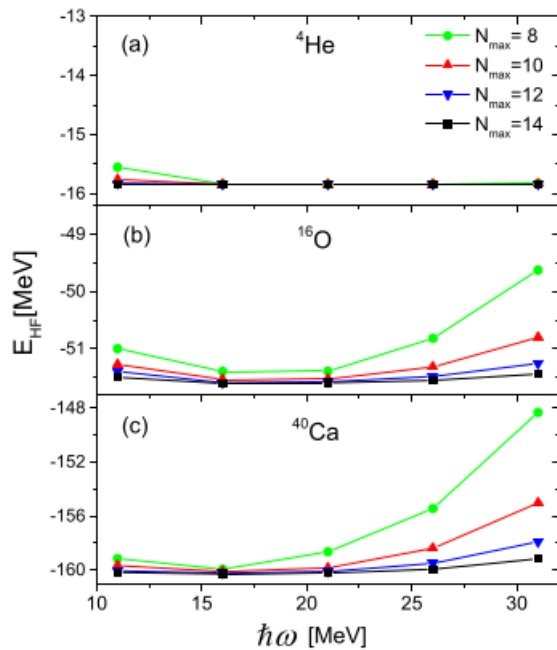


FIG. 1. HF ground-state energy of ⁴He (a), ¹⁶O (b), and ⁴⁰Ca (c) versus the HO frequencies ω for different HO space dimensions N_{\max} .

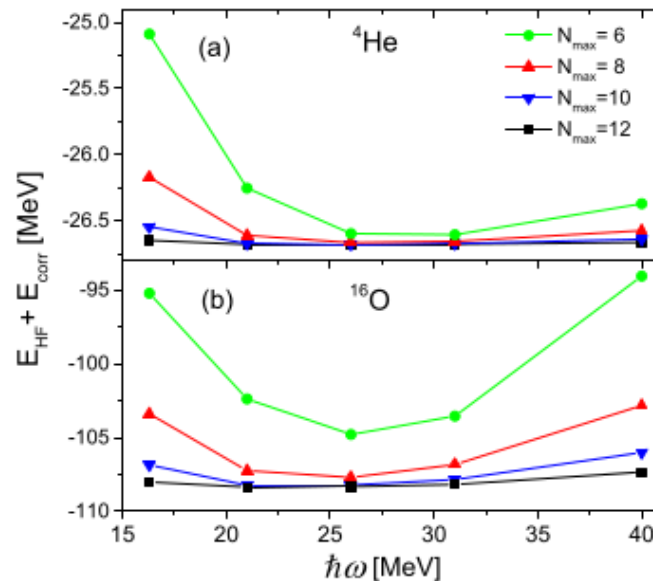


FIG. 4. The EMPM ground-state energy of ⁴He (a) and ¹⁶O (b) versus the HO frequency ω for different N_{\max} .

N_{\max} – maximal osc. shell

$\hbar\omega$ – parameter of basis

Final energy must be converged with respect to N_{\max} and for N_{\max} big enough independent on $\hbar\omega$...

E_{HF}

E_{EMPM}

Ground State Correlations

2-phonon correlations in the g.s.

NN interaction - χ NNLO_{opt}

A. Ekström et al., PRL 110, 192502 (2013)

$$|\Psi_{g.s.}\rangle \approx C_{HF}^{g.s.}|HF\rangle + \sum_{\mu_2} C_{\mu_2}^{g.s.}|i=2, \mu_2\rangle$$

proton point radii

$$\langle r_p^2 \rangle = \langle \Psi_{g.s.} | r_p^2 | \Psi_{g.s.} \rangle = \langle r_p^2 \rangle_{HF} + \langle r_p^2 \rangle_{corr.}$$

G. De Gregorio, J. Herko, F. Knapp, N. Lo Iudice, P. Veselý, PRC 95, 024306 (2017)

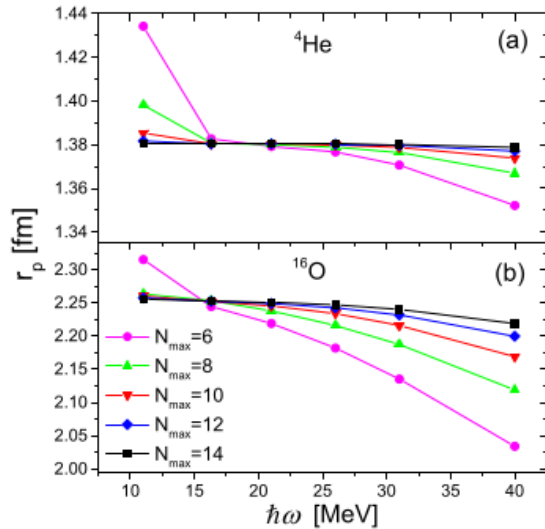


FIG. 5. HF point proton radius versus the HO frequency ω for different N_{max} in ^4He (a) and ^{16}O (b).

small effect of correlations on r_p

NNN forces play important role here!

“plateau” dependence on $\hbar\omega$ – parameter

$\leftarrow (r_p)_{HF}$

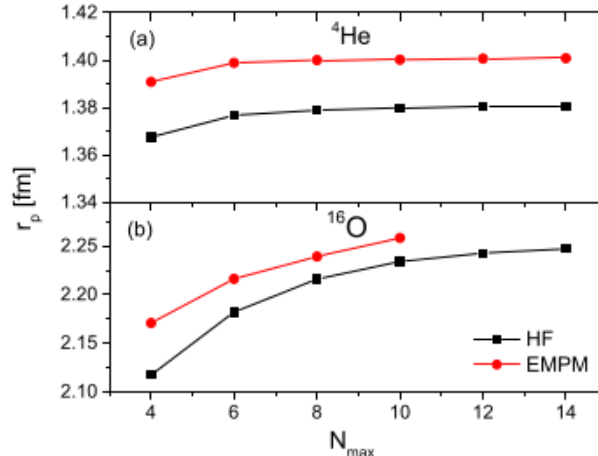


FIG. 7. HF and EMPM point proton radii of ^4He (a) and ^{16}O (b) versus N_{max} for fixed frequency ($\hbar\omega = 26$ MeV).

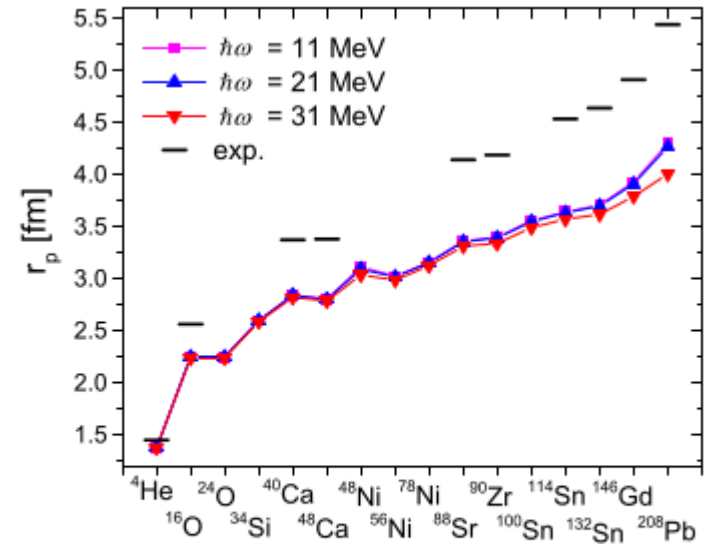


FIG. 6. Systematic of root-mean-square point proton radii computed in HF. The calculations are performed for $N_{\text{max}} = 14$ and different HO frequencies ω . The experimental data are from Ref. [49].

$^A X$	HF	r_p (fm) EMPM	Exp.
^4He	1.38	1.40	1.46
^{16}O	2.25	2.26	2.57

Ground State - NNN Force

NN+NNN interaction - χ NNLO_{sat} (Ekström et al. *Phys. Rev. C* 91 (2015) 051301R)

HO basis

$$N = (2n + l)$$

$\hbar\omega = 16$ MeV

N_{\max} up to 12

charged radii

Table 1: The charge radii $r_{\text{ch}} = \sqrt{\langle r_{\text{ch}}^2 \rangle}$ [fm] of ^{16}O and ^{40}Ca calculated with NN and NN+NNN forces are compared with the experimental data (exp) [23].

$^A X$	NN	NN+NNN	exp
^{16}O	2.19	2.77	2.70
^{40}Ca	2.58	3.54	3.48

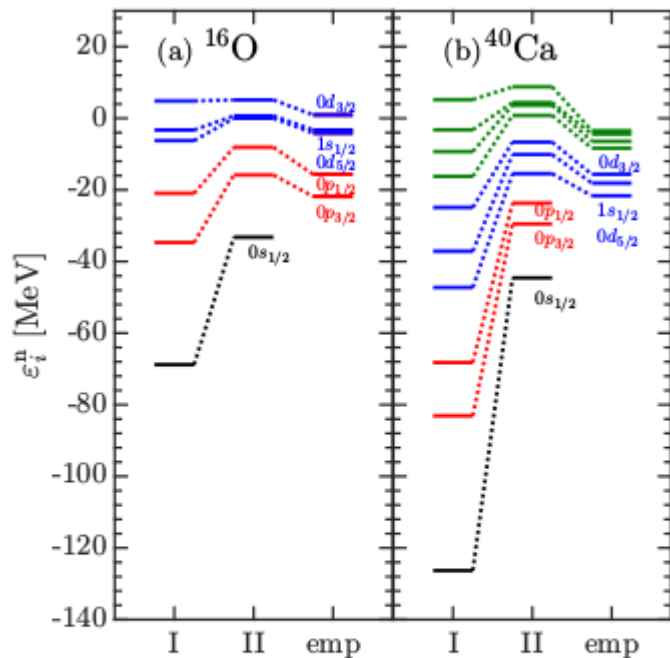


Fig. 3: The neutron single-particle energies ϵ_i^n of ^{16}O (a) and ^{40}Ca (b) calculated with NN (I) and NN+NNN (II) interactions. The empirical data (emp) [24] are shown for comparison.

HF energy

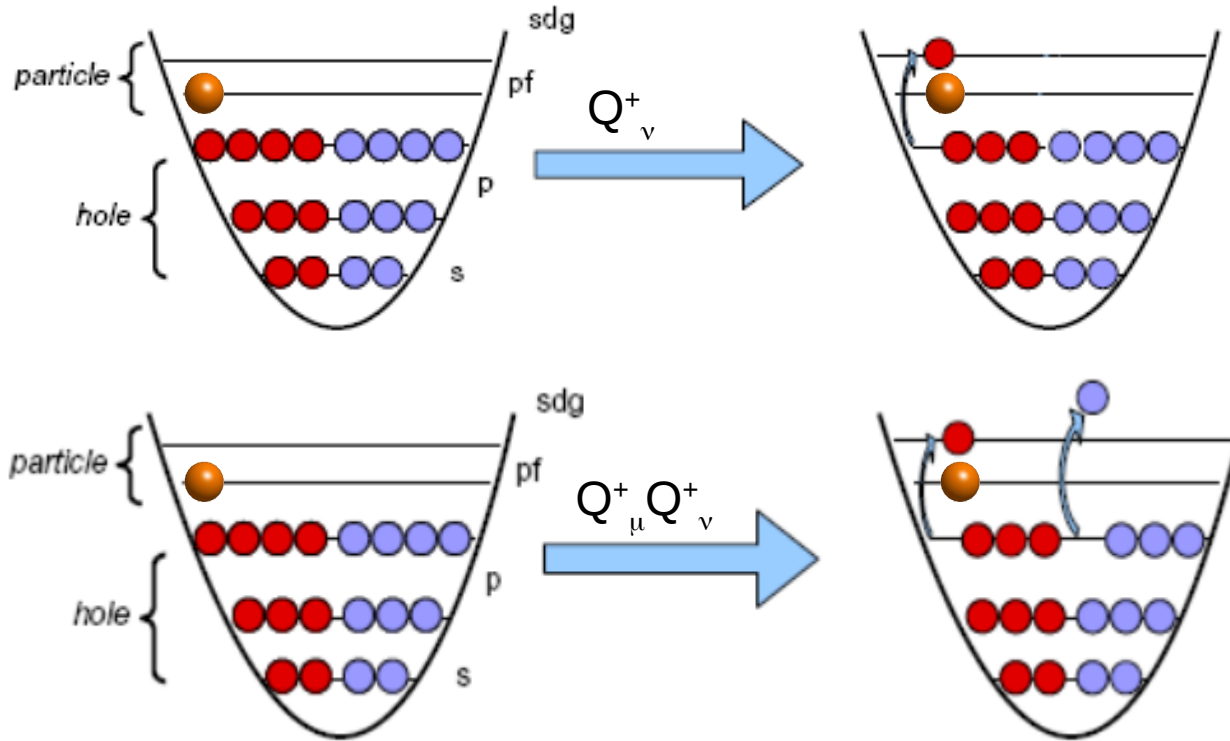
Table 2: Binding energies per nucleon BE/A [MeV] calculated with NN and with NN+NNN forces in ^{16}O and ^{40}Ca compared to the experimental values (exp).

$^A X$	NN	NN+NNN	exp
^{16}O	7.36	2.66	7.98
^{40}Ca	11.65	2.31	8.55

HF underestimates g.s. energy
(correlations necessary)

However NNN force improves significantly radii & single-particle energies already at the mean-field level

EMPM for Even-Odd Nuclei



Hilbert space – divided into separate **n-phonon** subspaces

$$\mathcal{H} = \mathcal{H}_0 \oplus \mathcal{H}_1 \oplus \dots \oplus \mathcal{H}_n$$

$$\mathcal{H}_0 = \{c_m^\dagger |HF\rangle\},$$

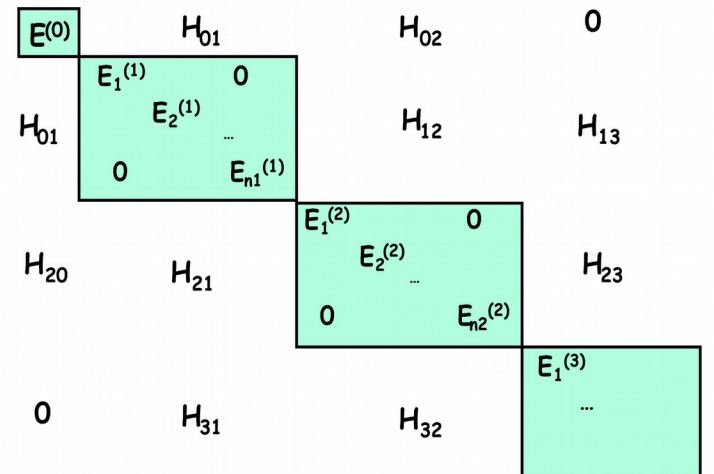
$$\mathcal{H}_1 = \{c_m^\dagger Q_\nu^\dagger |HF\rangle\},$$

$$\vdots$$

$$\mathcal{H}_n = \{c_m^\dagger Q_{\nu_1}^\dagger \dots Q_{\nu_n}^\dagger |HF\rangle\}.$$

explicit **coupling** of valence **nucleon** to general **excitations** of the **nuclear core**

then diagonalization of complete Hamiltonian



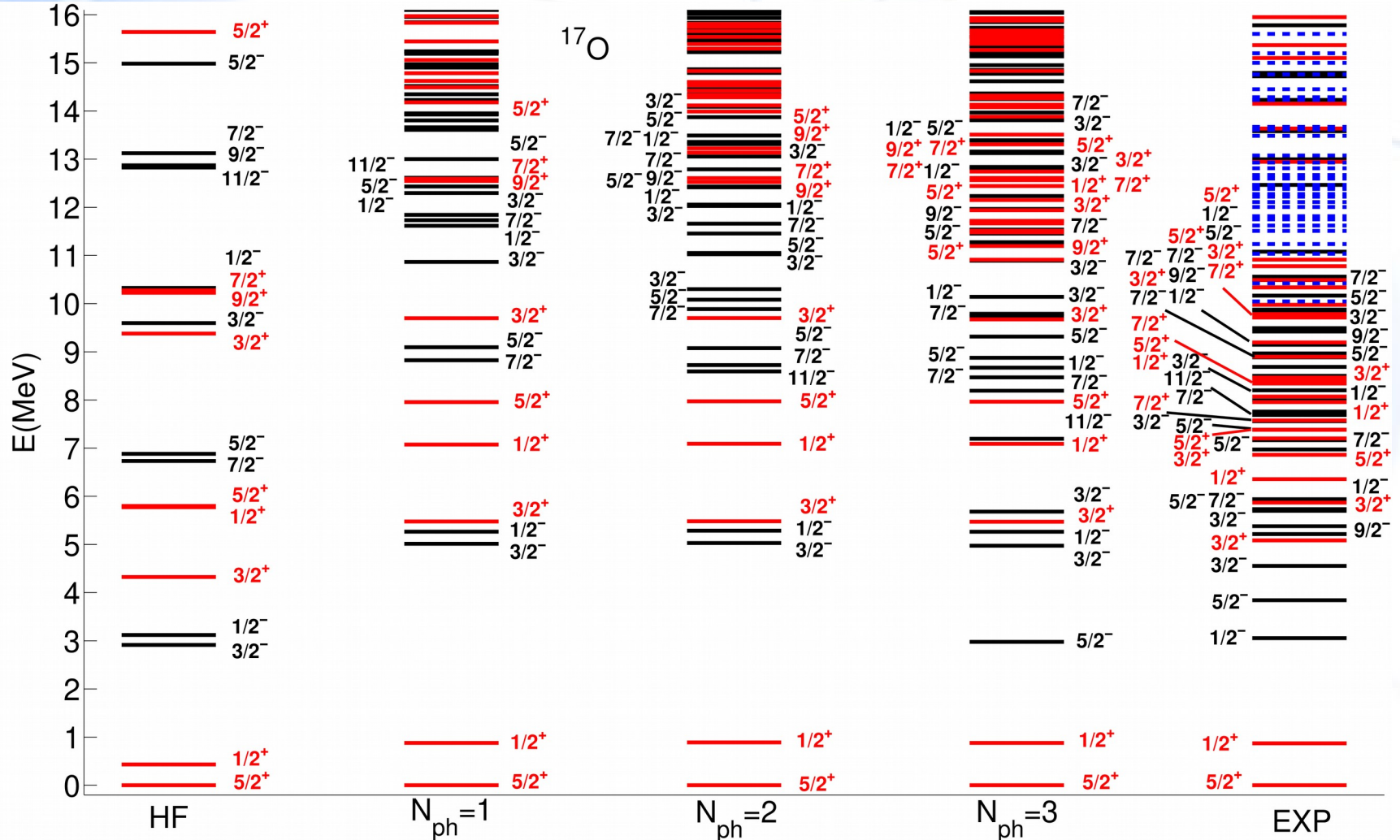
EMPM for Even-Odd Nuclei

^{17}O - particle coupled to (multi)phonon excitations

NN interaction - χ NNLO_{opt}

First application of **EMPM** on the **odd** nuclear systems:

G. De Gregorio, F. Knapp, N. Lo Iudice, P. Veselý, *Phys. Rev. C* **94**, 061301(R) (2016)



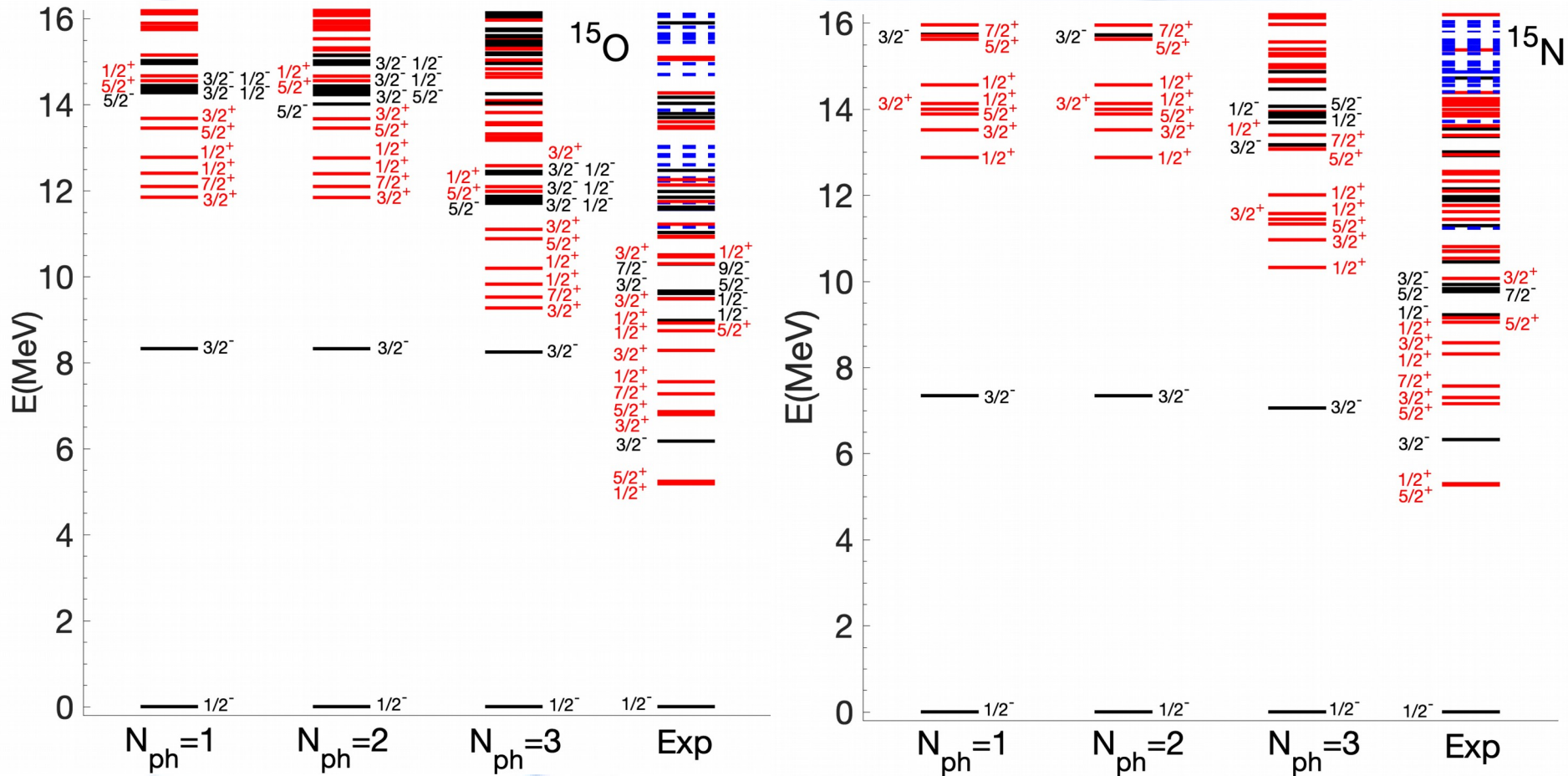
EMPM for Even-Odd Nuclei

^{15}O , ^{15}N , ^{21}O , ^{21}N - hole coupled to (multi)phonon excitations NN interaction - χ NNLO_{opt}

G. De Gregorio, F. Knapp, N. Lo Iudice, P. Veselý, sent to **Phys. Rev. C** (2018)

Lowest states – predominantly from hole-1phonon configurations

For better description, stronger coupling to more-phonon configs. needed

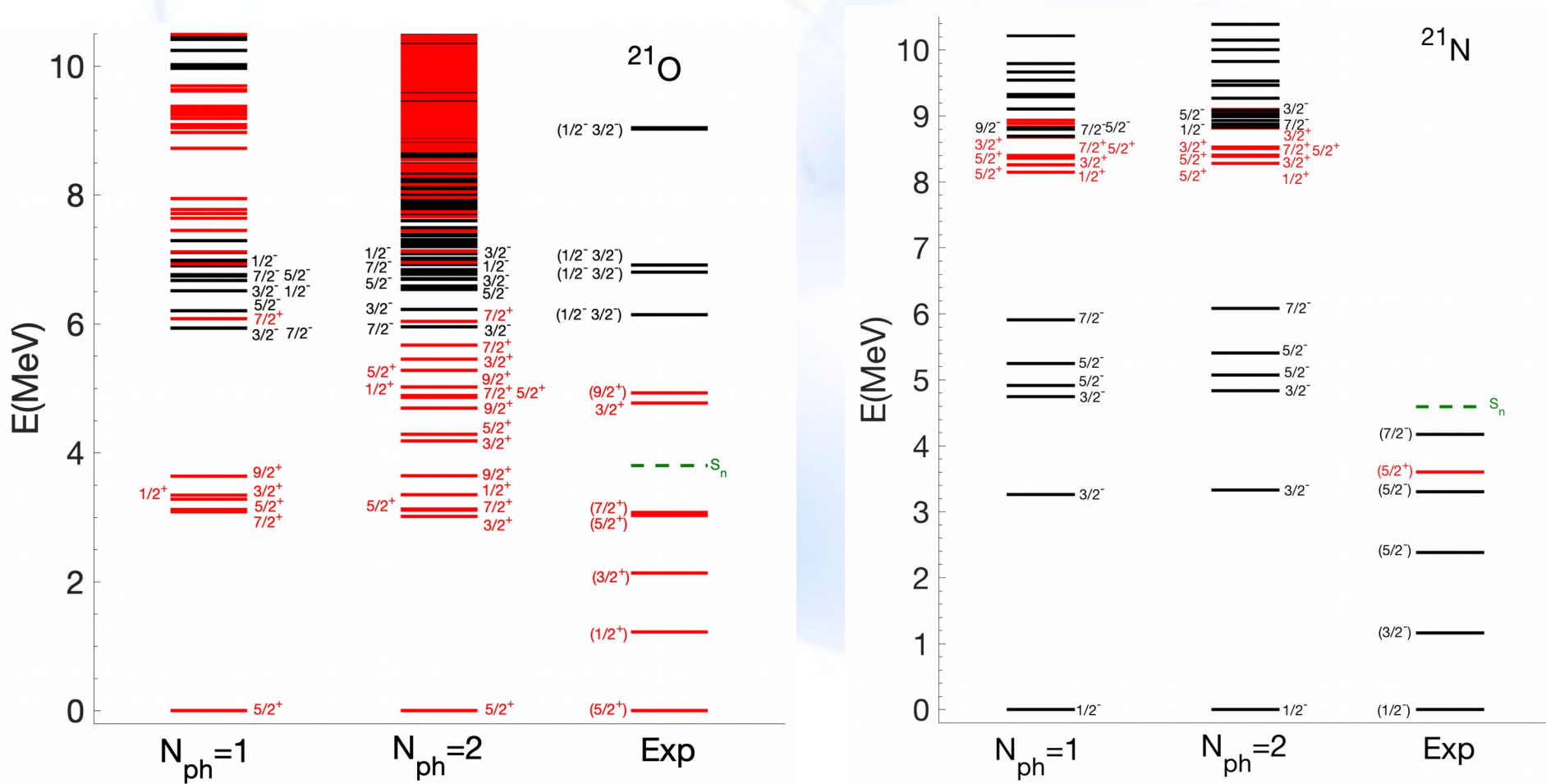


EMPM for Even-Odd Nuclei

^{15}O , ^{15}N , ^{21}O , ^{21}N - hole coupled to (multi)phonon excitations NN interaction - χ NNLO_{opt}

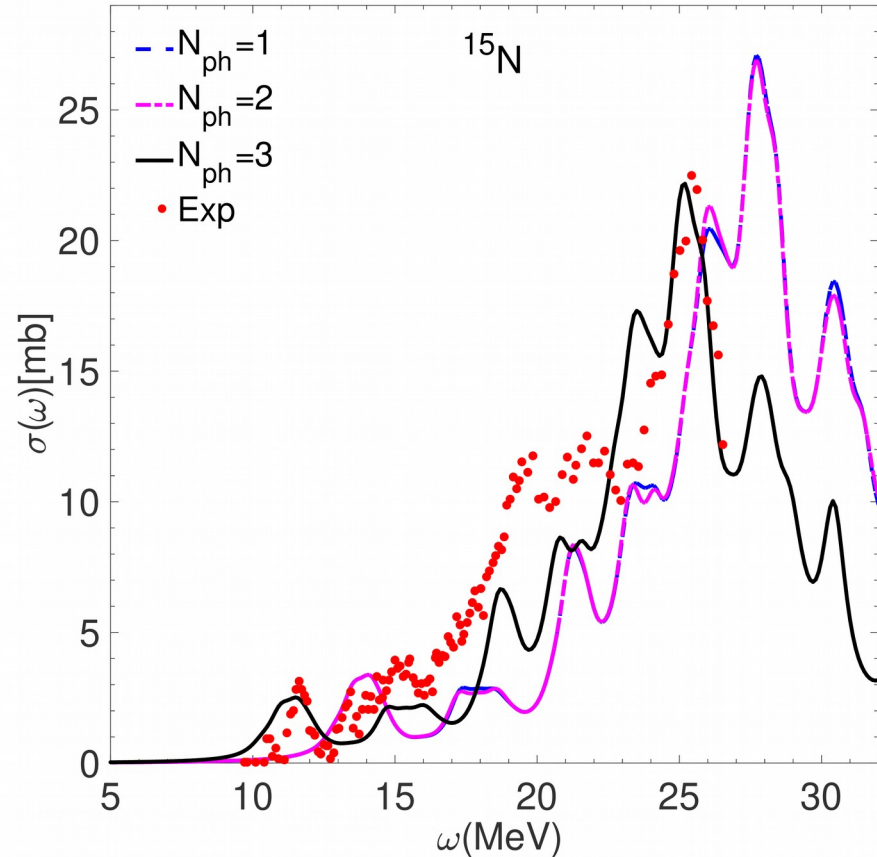
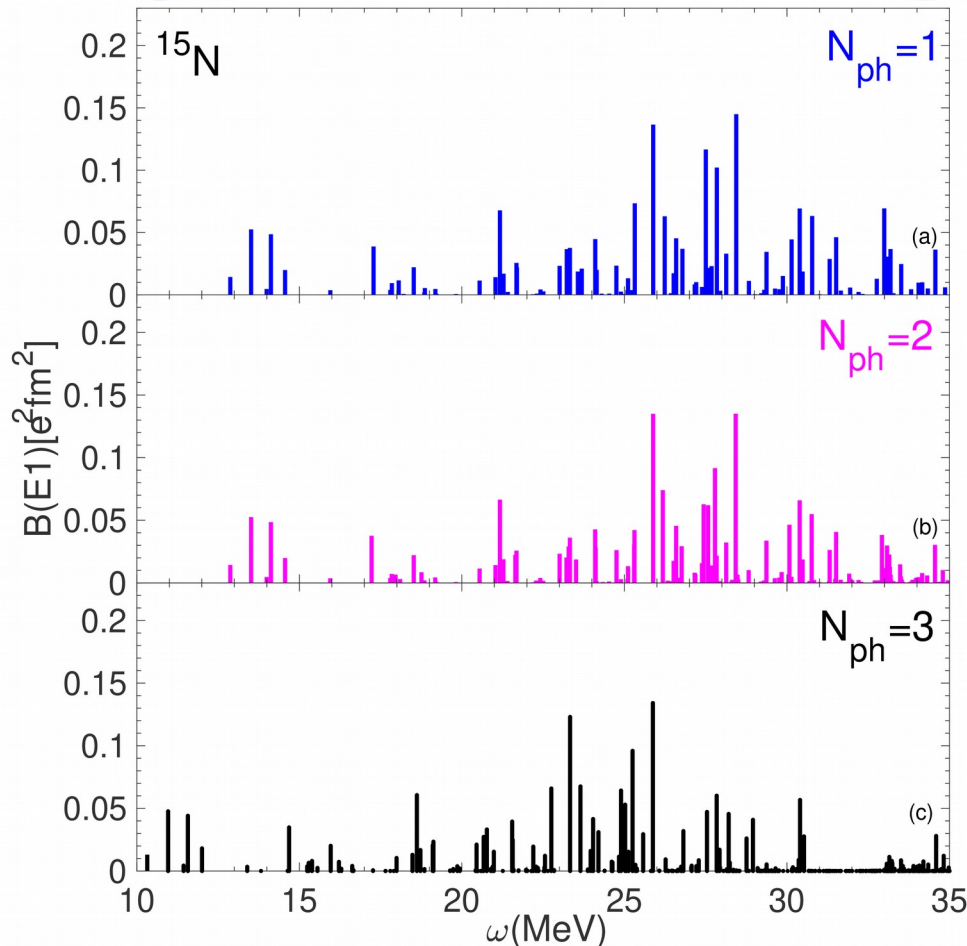
G. De Gregorio, F. Knapp, N. Lo Iudice, P. Veselý, sent to **Phys. Rev. C** (2018)

ground states – predominantly the hole-states nature



EMPM for Even-Odd Nuclei

role of **hole-1phon, hole-2phon, hole-3phon** configurations on **E1** transitions in ^{15}N



~ **112 %** of Thomas-Reiche-Kuhn sum rule up to **40 MeV**

~ **50 %** of TRK sum rule up to **26.5 MeV** (experiment ~ **58 %**)

EMPM for Hypernuclei

$$\hat{H} = \hat{T}_N + \hat{T}_\Lambda + \hat{V}^{NN} + \hat{V}^{NNN} + \hat{V}^{\Lambda N} + \cancel{\hat{V}^{\Lambda NN}} - \hat{T}_{CM}$$

NN+NNN interaction - χ **NNLO_{sat}** (Ekström et al. **Phys. Rev. C** 91 (2015) 051301R)

ΛN part of **YN** interaction - χ **LO** (H. Polinder, J. Haidenbauer, U. Meissner, **Nucl. Phys. A** 779 (2006) 244) **cut-off** $\lambda = 550$ MeV

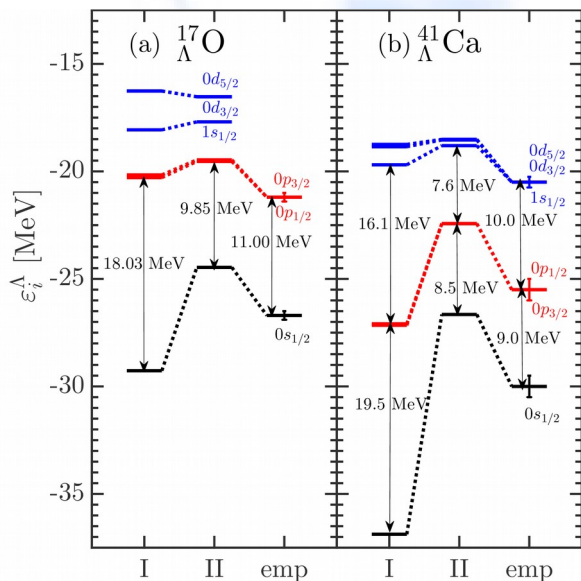
so far implemented:

extension of HF+TDA formalism on hypernuclei \rightarrow **proton-neutron- Λ HF + ΛN TDA**

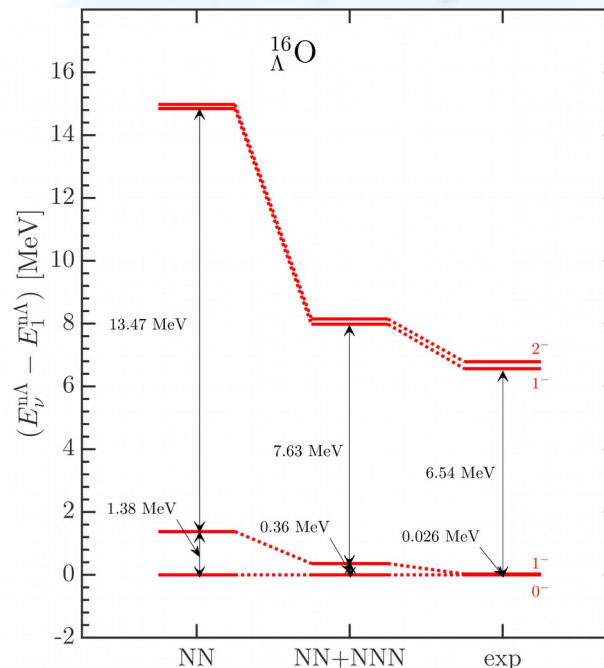
(replacement of the nucleon by Λ)

work in progress: - adding **Λ - Σ coupling** and **ΛNN SRG induced** force into the formalism

- **coupling to (multi)phonon configurations**



single-particle energies of Λ



important role of **NNN** part of **NNLO_{sat}**

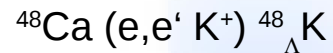
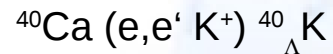
strong dependence on **cut-off** parameter of **YN** int.

in future testing the quality of **YN** interaction (**spin-dependent** part)

Outlook

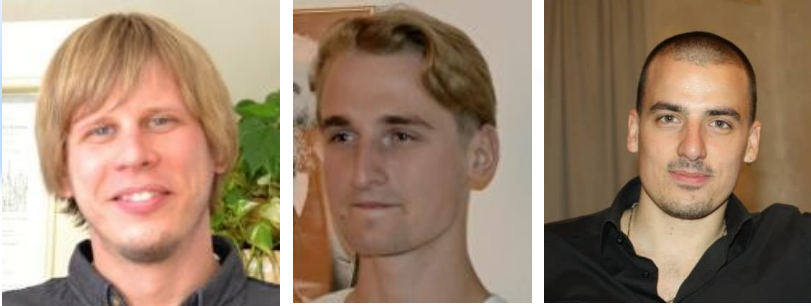
next goals:

- study of the role of **NNN** interaction in nuclear ground state properties
- more systematic studies of **odd nuclei** – heavier systems
- further extensions of **EMPM** formalism – **odd-odd nuclei, hypernuclei, ...**
- transitions in nuclei – **GDR, M1, GMR, β decay (2β decay) ...**
- possibly calculations of electro**production of hypernuclei**



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Thank you!!