# Shape coexistence and collective low-spin states in <sup>112,114</sup>Sn (Lifetime measurements with SONIC@HORUS)

M. Spieker et al., PRC 97, 054319 (2018)

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supported by DFG (ZI 510/7-1)

MS supported in part by the National Science Foundation under Contract No. PHY-1565546 (NSCL)

### Sn – transitional nuclei and collectivity?



<sup>[</sup>Figures: NNDC, NuDat (2018)]

Sn isotopes considered as one of the prime examples for spherical nuclei

- → Well-developed Z = 50 magic shell closure (no significant contribution of protons to low-lying excitation spectrum)
- $\rightarrow$  Little variation of  $E_{2^+_1}$ , i.e. no strong p-n interaction and no onset of deformation



## The $B(E2; 0_1^+ \rightarrow 2_1^+)$ puzzle understood?

#### Novel Shape Evolution in Sn Isotopes from Magic Numbers 50 to 82

Tomoaki Togashi,<sup>1</sup> Yusuke Tsunoda,<sup>1</sup> Takaharu Otsuka,<sup>2,1,3,4,5,\*</sup> Noritaka Shimizu,<sup>1</sup> and Michio Honma<sup>6</sup>
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 <sup>5</sup>National Superconducting Cyclotron Laboratory, Michigan State University, East Lansing, Michigan 48824, USA
 <sup>6</sup>Center for Mathematical Sciences, University of Aizu, Ikki-machi, Aizu-Wakamatsu, Fukushima 965-8580, Japan



- Activating protons in the 1g<sub>9/2</sub> shown to provide a possible explanation (proton holes)
  - $\rightarrow$  Breaking of Z = 50 core!
  - → Modest prolate deformation in
     N = 50 to 64 isotopes
  - → 2<sup>nd</sup>-order phase transition from modestly deformed to pairing phase in Sn nuclei (N = 66)!

[T. Togashi, Y. Tsunoda, T. Otsuka, et al., PRL 121, 062501 (2018)]

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### Shape coexistence in Sn isotopes



[P.E. Garrett, J. Phys. G 43, 084002 (2016); D. Rowe and J.L. Wood, Models of Nuclear Structure: Foundational Models (World Scientific, Singapore, 2010)]

- Excited 0<sup>+</sup> state strongly excited in two-proton transfer reaction
  - $\rightarrow$  2p-2h proton structure?
- Parabolic evolution of additional structure with minimum at midshell
  - $\rightarrow$  One of the key signatures of intruder configurations
- Large  $\rho^2(E0; 0_3^+ \rightarrow 0_2^+)$  suggests strong mixing between excited 0<sup>+</sup> states in <sup>116</sup>Sn
  - $\rightarrow$  Which 0<sup>+</sup> state is the bandhead of the intruder structure?
- Collective intraband E2 transitions



### Shape coexistence in Sn isotopes



[J.L. Pore et al., EPJA 52, 27 (2017); D. Rowe and J.L. Wood, Models of Nuclear Structure: Foundational Models (World Scientific, Singapore, 2010)]

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#### **Collective structures in Sn isotopes?**





#### **Collective structures in Sn isotopes?**

#### Quasi-rotational structure of Cd isotopes and mixing between the different configurations

(are there "true" vibrational states at all?)









#### **Collective structures in Sn isotopes?**



#### **Experimental requirements:**

- Selective probe needed to excite those low-spin states (non-Yrast)
- Small γ-decay branching ratios need to be detected
- Lifetimes and multipole-mixing ratios need to be measured for the determination of reduced transition strengths

### **Proton-γ coindicences with SONIC@HORUS**



#### **10 MV FN Tandem ion accelerator**

- Three ion sources available
  - $\rightarrow$  typically ions up to Z = 30
- Terminal voltages from 1 MV to 10 MV
- Current on target up to 1 μA (protons)
- Experimental setups: new AMS beamline, Plunger setup, Orange spectrometer, LYCCA, HORUS spectrometer

#### The SONIC@HORUS setup

#### HORUS for $\gamma$ -ray detection

- Up to 14 HPGe detectors ( $\epsilon_{\text{FEP}} \sim 2\%$  @ 1.3 MeV)
  - → Six BGO shields and two Clover HPGe detectors available

#### SONIC for particle detection

- 7-12 silicon detectors (thickness  $\leq$  1.5 mm)
- Particle-γ coincidences (Lifetimes, branching ratios, angular correlations)



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#### **Proton-γ coindicences with SONIC@HORUS**

py-coincidence matrix



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### (p,p' $\gamma$ ) DSA coincidence technique



**(p,p'γ) DSA coincidence technique:** A. Hennig *et al.*, NIM **794**, 171 (2015) **SONIC@HORUS (UoC, Germany):** S.G. Pickstone *et al.*, NIM **875**, 104 (2017)



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### (p,p' $\gamma$ ) DSA coincidence technique

#### Our method is very similar to the (n,n' $\gamma$ ) technique used at the University of Kentucky (USA)



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#### Lifetimes of intruder states in <sup>112,114</sup>Sn



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#### $\gamma$ -decay behavior of the states of interest

<sup>112</sup> Sn					<sup>114</sup> Sn				
$E_x$ [keV]	$J_i^{\pi}$	$J_f^\pi$	$E_{\gamma}$ [keV]	<i>Ι<sub>γ</sub></i> [%]	$E_x$ [keV]	$J_i^{\pi}$	$J_f^\pi$	$E_{\gamma}$ [keV]	$I_{\gamma}$ [%]
1256.5(2)	$2^{+}_{1}$	$0_{1}^{+}$	1256.5(2)	100	1299.7(2)	$2^{+}_{1}$	$0_{1}^{+}$	1299.7(2)	100
2150.5(3)	$2^{+}_{2}$	$0_{1}^{+}$	2150.5(2)	20(3)	1952.9(2)	$0^{+}_{2}$	$2^{+}_{1}$	653.2(2)	100
	$2^{+}_{2}$	$2^{+}_{1}$	893.9(2)	100	2155.9(2)	$0_{3}^{+}$	$2^{+}_{1}$	856.2(2)	100
					2187.3(3)	$4_{1}^{+}$	$2^{+}_{1}$	887.6(2)	100
2190.5(2)	$0^{+}_{2}$	$2^+_1$	934.0(2)	100	2238.6(2)	$2^{+}_{2}$	$0_{1}^{+}$	2238.5(2)	100
2247.0(3)	$4_{1}^{+}$	$2^+_1$	990.47(10)	100		$2^{+}_{2}$	$2^{+}_{1}$	938.9(2)	81(12)
2353.7(2)	$3^{-}_{1}$	$2^+_1$	1097.2(2)	100		$2^{\frac{2}{+}}_{2}$	$0^{+}_{2}$	286.5(10)	0.9(3)
2475.5(2)	$2^+_3$	$0_{1}^{+}$	2475.5(2)	100	2274.5(2)	$3\frac{1}{1}$	$2^{2}_{1}$	974.8(2)	100
	$2^{+}_{3}$	$2^{+}_{1}$	1218.9(2)	36(5)	2420.5(2)	$0^{+}_{4}$	$2^{\frac{1}{+}}$	1120.8(2)	100
	$2^+_3$	$0^{+}_{2}$	284.9(2)	0.70(10)	2453.8(2)	$2^{4}_{2}$	$0_{1}^{+}$	2453.7(2)	28(4)
2520.5(2)	$4^+_2$	$2^+_1$	1264.0(2)	100	(_)	$2^{+}_{2}$	$2^{+}_{1}$	1154.0(2)	100
		:	-			$2^{+}_{2}$	$2^{+}_{2}$	215.4(4)	1.3(3)
		Ŀ			2514.4(2)	$3^{+}_{1}$	$4^{+}_{1}$	327.1(2)	100
2945.0(7)	4+	$2^{+}_{1}$	1688.5(2)	100	2613.7(4)	$4^{+}_{2}$	$2^{+}_{1}$	1314.5(2)	100
	4+	$2^{+}_{2}$	794.2(2)	5.4(10)		$4^{+}_{2}$	$4^{+}_{1}$	426.0(4)	1.6(6)
	4+	$4_1^{\tilde{+}}$	697.9(2)*	<1.5		4+	2+	375.2(3)	1.8(6)
	4+	$2^{+}_{3}$	469.5(2)	18(3)		•2	-2	0,012(0)	110(0)
	4+	$4^{+}_{2}$	424.6(3)*	4.9(9)					
	4+	$6_{1}^{+}$	396.4(4)*	2.3(5)					
	4+	4+	$161.4(2)^*$	9(2)					



#### $\gamma$ -decay behavior of the states of interest

<sup>112</sup> Sn					<sup>114</sup> Sn				
$E_x$ [keV]	$J_i^{\pi}$	$J_f^\pi$	$E_{\gamma}$ [keV]	<i>Ι<sub>γ</sub></i> [%]	$E_x$ [keV]	$J_i^{\pi}$	$J_f^\pi$	$E_{\gamma}$ [keV]	$I_{\gamma}$ [%]
1256.5(2)	$2^{+}_{1}$	$0_{1}^{+}$	1256.5(2)	100	1299.7(2)	$2^{+}_{1}$	$0_{1}^{+}$	1299.7(2)	100
2150.5(3)	$2^{+}_{2}$	$0_{1}^{+}$	2150.5(2)	20(3)	1952.9(2)	$0^+_2$	$2^{+}_{1}$	653.2(2)	100
	$2^{+}_{2}$	$2^+_1$	893.9(2)	100	2155.9(2)	$-\frac{1}{0_{3}^{+}}$	$2^{+}_{1}$	856.2(2)	100
· — — — ·					2187.3(3)	$4_{1}^{+}$	$2^{+}_{1}$	887.6(2)	100
2190.5(2)	$0^+_2$	$2_1^+$	934.0(2)	100	2238.6(2)	$2^{+}_{2}$	$0^+_1$	2238.5(2)	-100
2247.0(3)	$-4_1^+$	$2_{1}^{+}$	990.47(10)	100		$2^{+}_{2}$	$2^{+}_{1}$	938.9(2)	81(12)
2353.7(2)	$3_{1}^{-}$	$2^+_1$	1097.2(2)	100		$2^{+}_{2}$	$0^{+}_{2}$	286.5(10)	0.9(3)
2475.5(2)	$2^+_3$	$0^+_1$	2475.5(2)	100	2274.5(2)	$-\frac{2}{3_1^-}$	$-\frac{2}{2_1^+}$	974.8(2)	-100
1	$2^+_3$	$2^+_1$	1218.9(2)	36(5)	2420.5(2)	$0_{4}^{+}$	$2^{+}_{1}$	1120.8(2)	100
	$-\frac{2_3}{4+}$	$-\frac{0_2}{2^+}$	$-\frac{284.9(2)}{1264.9(2)}$	0.70(10)	2453.8(2)	$2^{+}_{3}$	$0_{1}^{+}$	2453.7(2)	28(4)
2520.5(2)	42	$2_{1}^{+}$	1264.0(2)	100		$2^{+}_{3}$	$2^{+}_{1}$	1154.0(2)	100
		:	-			$2^{+}_{3}$	$2^{+}_{2}$	215.4(4)	1.3(3)
					2514.4(2)	$3_{1}^{+}$	$-4^{\tilde{+}}_{1}$	327.1(2)	100
2945.0(7)	4+	$2_{1}^{+}$	1688.5(2)	100	2613.7(4)	$4^{+}_{2}$	$2^{+}_{1}$	1314.5(2)	100
1	4+	$2^{+}_{2}$	794.2(2)	5.4(10)		$4^{+}_{2}$	$4^{+}_{1}$	426.0(4)	1.6(6)
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* New γ-deo	cay branch	ing			L				

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<sup>112</sup> Sn					<sup>114</sup> Sn				
$E_x$ [keV]	$J_i^{\pi}$	$J_f^\pi$	$E_{\gamma}$ [keV]	$I_{\gamma}$ [%]	$E_x$ [keV]	$J_i^{\pi}$	$J_f^\pi$	$E_{\gamma}$ [keV]	$I_{\gamma}  [\%]$
1256.5(2)	$2_{1}^{+}$	$0_{1}^{+}$	1256.5(2)	100	1299.7(2)	$2^{+}_{1}$	$0^{+}_{1}$	1299.7(2)	100
2150.5(3)	$2^{+}_{2}$	$0_{1}^{+}$	2150.5(2)	20(3)	1952.9(2)	$0^+_2$	$2^+_1$	653.2(2)	100
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2475.5(2)	$2^+_3$	$0^+_1$	2475.5(2)	100	2274.5(2)	$-\frac{2}{3_1}$	$\frac{1}{2_1^+}$	974.8(2)	100
-	$2^+_3$	$2_{1}^{+}$	1218.9(2)	36(5)	2420.5(2)	$0_{4}^{+}$	$2^{+}_{1}$	1120.8(2)	100
2520.5(2)	$-\frac{2_3}{4^+}$	$-\frac{0_2}{2^+}$ -	$\frac{284.9(2)}{1264.0(2)}$	0.70(10)	2453.8(2)	$2^{+}_{3}$	$0_{1}^{+}$	2453.7(2)	28(4)
2520.5(2)	42	$Z_1$	1264.0(2)	100		$2^{+}_{3}$	$2^{+}_{1}$	1154.0(2)	100
0+@2	2617 keV	:	-			$2^{+}_{3}$	$2^{+}_{2}$	215.4(4)	1.3(3)
					2514.4(2)	$3_{1}^{+}$	41	327.1(2)	100
2945.0(7)	4+	$2^{+}_{1}$	1688.5(2)	100	2613.7(4)	$4^{+}_{2}$	$2^{+}_{1}$	1314.5(2)	100
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	4+	4+	161.4(2)*	9(2)		"intr	uder	" states	
* New γ-deo	cay branchir	ng							
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"Quasi-rotational structure" already existent at higher energies in Sn isotopes?























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Influence of underlying single-particle structure or overall structure change?

Different influence of neutron single-particle states?



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### Changing shell structure from <sup>114</sup>Sn to <sup>112</sup>Sn?



## **Comparison to IBM-2 mixing calculations** (assuming <sup>110</sup>Pd to cause intruder structure)

$J_i^{\pi}$	$E_x$	$E_{x,\mathrm{IBM}}$	$J_f^{\pi}$	$B(E2)_{\mathrm{exp.}}\downarrow$	$B(E2)_{\rm IBM}\downarrow$			
	[MeV]	[MeV]	-	[W.u.]	[W.u.]			
		norm	al co	nfiguration				
$2_{1}^{+}$	1.30	1.30	$0_{1}^{+}$	11.1(7)	11			
$4_{1}^{+}$	2.19	2.28	$2_{1}^{+}$	5.9(5)	19			
$0_{2}^{+}$	1.95	1.99	$2_{1}^{+}$	23.2(8)	21			
$2^{+}_{3}$	2.45	2.54	$0_{1}^{+}$	0.023(9)	0.004			
			$2_{1}^{+}$	3(2)	17			
			$2^{+}_{2}$	-	8			
	intruder configuration							
$0^{+}_{3}$	2.16	2.15	$2_{1}^{+}$	$\leq 5$	2			
$2^{+}_{2}$	2.24	2.46	$0_{1}^{+}$	$\leq 0.12$	0.04			
			$2_{1}^{+}$	$\leq 8$	2			
			$0^{+}_{2}$	$\leq 44$	31			
			$0^{+}_{3}$	-	27			
$4_{2}^{+}$	2.61	3.00	$2_{1}^{+}$	6.6(10)	0.2			
			$4_{1}^{+}$	1.6(10)	0.06			
			$2^{+}_{2}$	62(25)	85			
$6^{+}$	3.19	3.63	$4_{1}^{+}$	1.68(9)	1.5			
			$4_{2}^{+}$	97(5)	93			
			$4_{3}^{+}$	18.9(12)	0.7			





## **Comparison to IBM-2 mixing calculations** (assuming <sup>110</sup>Pd to cause intruder structure)

$J_i^{\pi}$	$E_x$	$E_{x,\mathrm{IBM}}$	$J_f^{\pi}$	$B(E2)_{\mathrm{exp.}}\downarrow$	$B(E2)_{\rm IBM}\downarrow$
	[MeV]	[MeV]		[W.u.]	[W.u.]
		norm	nal con	nfiguration	
$2_{1}^{+}$	1.30	1.30	$0_{1}^{+}$	11.1(7)	11
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$2^{+}_{3}$	2.45	2.54	$0_{1}^{+}$	0.023(9)	0.004
			$2_{1}^{+}$	3(2)	17
			$2^{+}_{2}$	-	8
		intru	der co	nfiguration	
$0^{+}_{3}$	2.16	2.15		= 0.9(3) %	(Fxp.)
$2^{+}_{2}$	2.24	2.46	·γ,2 Ι <sub>γ,3</sub>	≈ 0.002 %	6 (IBM)
			$0^{+}_{2}$	$\leq 44$	31
			$0^{+}_{3}$	-	27
$4_2^+$	2.61	<sup>3</sup> B(I	E2) =	= 55.5(9) V	V.u. in <sup>110</sup> Pd
			$2^{+}_{2}$	62(25)	85
$6^{+}$	3.19	3.63	$4_{1}^{+}$	1.68(9)	1.5
			$4^{+}_{2}$	97(5)	93
			$4_{3}^{+}$	18.9(12)	0.7





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#### **Comparison to IBM-2 mixing calculations** (assuming <sup>110</sup>Pd to cause intruder structure)

**S** 

						2p 2ii build
$J_i^{\pi}$	$E_x$ [MeV]	$E_{x,\text{IBM}}$ [MeV]	$J_f^{\pi}$	$\begin{array}{c} B(E2)_{\text{exp.}} \downarrow \\ \text{[W.u.]} \end{array}$	$B(E2)_{\text{IBM}}\downarrow$ [W.u.]	B(E2) = 44(3) W.u.
		norm	al cor	figuration		$_{\gamma,2}^{641} = 5.16(14) \%$
$2^{+}_{1}$	1.30	1.30	$0_{1}^{+}$	11.1(7)	11	$4_{1}^{+}$ 2391 / 138 304
$4_{1}^{+}$	2.19	2.28	$2_{1}^{+}$	5.9(5)	19	279 - 417 - 2225 - 2
$0^{+}_{2}$	1.95	1.99	$2_{1}^{+}$	23.2(8)	21	$\frac{85}{1226}$ $\frac{85}{0327}$ $0^{+}_{-}355$ $469$
$2^{+}_{3}$	2.45	2.54	$0_{1}^{+}$	0.023(9)	0.004	1097 $1230 / 2027 / 1757 0 + 932 / 932 /$
			$2_{1}^{+}$	3(2)	17	$\begin{pmatrix} 319 \\ 1 \end{pmatrix} \begin{pmatrix} 734 \\ 463 \end{pmatrix}$
			$2^{+}_{2}$	-	8	
		intru	der co	nfiguration		2112  B(E2) = 100(8) M/H
$0^{+}_{3}$	2.16	2.15		= 0.9(3) %	(Exp.)	D(L2) = 100(8)  V.d.
$2^{+}_{2}$	2.24	2.46	·γ,2	$\sim 0.002 %$	$(\Box \Lambda \Lambda)$	$I_{\gamma,3} = 0.0091(6) \%$
			γ,3	$\approx 0.002$ /6		
			$0^+_2$	$\leq 44$	31	EPJA 52, 27 (2017)
			$0^+_3$	-	27	$0^+$ $0$ $+$
$4_{2}^{+}$	2.61	<mark>3</mark> В(I	E2) =	= 55.5(9) W	<i>l</i> .u. in <sup>110</sup> Pd	B(E2) = 40(7) W.u. in <sup>112</sup> Pc
			$2^{+}_{2}$	62(25)	85	
6+	3.19	3.63	$4_1^+$ $4_2^+$	1.68(9) 97(5)	How co	uld this large B(E2) in <sup>116</sup> Sn be explained?
			$4_{3}^{2}$	18.9(12)	0.7	
			-			

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<sup>116</sup>Sn – Is 3<sup>rd</sup> 0<sup>+</sup> bandhead?

2n-2h Band

#### **Comparison to IBM-2 mixing calculations** (assuming <sup>110</sup>Pd to cause intruder structure)

S

						2p 2h Build
$J_i^{\pi}$	$E_x$ [MeV]	$E_{x,\text{IBM}}$ [MeV]	$J_f^{\pi}$	$\begin{array}{c} B(E2)_{\text{exp.}} \downarrow \\ [\text{W.u.}] \end{array}$	$B(E2)_{\text{IBM}}\downarrow$ [W.u.]	B(E2) = 44(3) W.u
		norn	nal cor	nfiguration		$_{641}^{641}_{2520}$ $I_{\gamma,2}^{503} = 5.16(14) \%$
$2_{1}^{+}$	1.30	1.30	$0_{1}^{+}$	11.1(7)	11	$4_{1}^{+}$ 2391 $138$ 304
$4_{1}^{+}$	2.19	2.28	$2_{1}^{+}$	5.9(5)	19	279 - 417 - 2025 - 2225 - 2112 - 21
$0^{+}_{2}$	1.95	1.99	$2_{1}^{+}$	23.2(8)	21	$\frac{85}{1226}$ $\frac{85}{2227}$ $\frac{0}{3355}$ $\frac{469}{469}$
$2^{+}_{3}$	2.45	2.54	$0_{1}^{+}$	0.023(9)	0.004	$1097$ $1230   2027   1757 0^+_2  932   1097   1097   1230   1097   100$
			$2_{1}^{+}$	3(2)	17	
			$2^{+}_{2}$	-	8	
		intru	der co	nfiguration		$P(E_2) = 100(8) M/11$
$0^{+}_{3}$	2.16	2.15		= 0.9(3) %	(Exp)	B(LZ) = 100(8) VV.U.
$2^{+}_{2}$	2.24	2.46	γ,2			$I_{\gamma,3} = 0.0091(6) \%$
			<b>γ</b> ,3	≈ 0.002 %		
			$0^{+}_{2}$	$\leq 44$	31	EPJA 52. 27 (2017)
			$0^+_3$	-	27	$O_1^+$ $O$
$4_{2}^{+}$	2.61	<sup>3</sup> B(	E2) =	= 55.5(9) W	<i>l</i> .u. in <sup>110</sup> Pd	B(E2) = 40(7) W.u. in <sup>112</sup> Pc
			$2^{+}_{2}$	62(25)	85	
6+	3.19	3.63	$4_1^+$ $4_2^+$	1.68(9) 97(5)	How co	uld this large B(E2) in <sup>116</sup> Sn be explained?
			$4^{+}_{3}$	18.9(12)	0.7	B(E2) = 101(5) W.u. in <sup>120</sup> X

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<sup>116</sup>Sn – Is 3<sup>rd</sup> 0<sup>+</sup> bandhead?

2n-2h Band

in <sup>120</sup>Xe



[C. Petrache, J.-M. Régis, C. Andreoiu, MS, et al., submitted for publication]

- Strong mixing between normal and intruder configuration
  - $\rightarrow$  Observation of two collective decay branches with "band" structure
- IBM-2 can provide reasonable to good description of experimental data
  - → Discrepancies can be explained by imposed selection rules

(further adjustment would be possible; maybe mix O(6)-like with O(6)-like)



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#### Summary

 SONIC@HORUS to determine lifetimes and γ-decay behavior of low-spin states via (p,p'γ) DSA coincidence technique

[A. Hennig *et al.*, NIM **794**, 171 (2015) ]
[S.G. Pickstone *et al.*, NIM **875**, 104 (2017)]

- Collectivity of low-spin "intruder" states studied in <sup>112,114</sup>Sn
- Mixing hypothesis between normal and intruder configuration tested via schematic IBM-2 mixing calculations
- No clear hints at quadrupole multiphonon structures in <sup>112,114</sup>Sn





 $\frac{5 \times 12}{5 \times 17}$  M. Spieker – Shape coexistence and collective low-spin states in <sup>112,114</sup>Sn