Probing the seesaw mechanism and leptogenesis

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based on 1710.03744, with Stefan Antusch, Eros Cazzato, Marco Drewes, Oliver Fischer, Björn Garbrecht and Dario Gueter

Recontres de Moriond EW 2018

La Thuile, March 15th 2018
Some of the missing pieces of the standard model:

The Baryon Asymmetry of the Universe

\[ \frac{n_B}{n_\gamma} = 6.05(7) \times 10^{-10} \]

Neutrino masses
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Neutrino masses

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Where to look for them: Normal hierarchy

- lower limits on the mixing angles from the seesaw mechanism and BBN
- low scale seesaw testable with current experiments
  [see talks by: Parkinson, Verbeke, Drewes, Marcano, Negro]
- leptogenesis within reach of NA62/ LHC
  [see talks by: Domcke, Lopez Pavon]
Where to look for them: Normal hierarchy

- Lower limits on the mixing angles from the seesaw mechanism and BBN
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- Leptogenesis within reach of NA62/LHC [see talks by: Domcke, Lopez Pavon]
Constraints on flavour patterns

- large mixing angles require a **flavour asymmetric washout**, which corresponds to a flavour asymmetric mixing

- together with **seesaw constraints** this imposes constraints on the mixing patterns for **large mixing angles**

\[ \bar{M} = 30 \text{ GeV} \]

[Antusch/Cazzato/Drewes/Fischer/Garbrecht/Gueter/JK 1710.03744]
Measuring the mass splitting

Normal Ordering:

- large range of viable $\Delta M$ consistent with leptogenesis
- extremely small mass splitting can be measured via oscillations $10^{-13}\text{GeV} \sim 1\text{cm}$
- nontrivial same sign/opposite sign dilepton ratios require $\Delta M \sim -\Delta M_{\theta\theta}$
Conclusions

- adding GeV-scale RHNs to the standard model can explain both the observed *neutrino masses* and the *BAU*
- the seesaw mechanism gives constraints on **RHN mixing patterns** (stronger if $\delta$ is measured!)
- testable seesaw and leptogenesis within reach of present (NA62 and LHC) experiments
- large mixing angles $+ \text{leptogenesis} \rightarrow$ even stronger predictions on the flavour patterns, easily falsifiable LG
- while complete determination of the RHN parameters is possible in principle, it requires extreme experimental sensitivity
Leptogenesis via Neutrino Oscillations

SM Thermal Bath

$CP$ - even
RHN states

Vanishing initial abundance

$CP$ - odd
RHN states

Oscillations

Conserved Generalized Lepton Number

Baryon Number
Asymmetry

$\text{number density of RHN} = 1. \times 10^{-4}$

$\text{source} = -3. \times 10^{-9}$

$\text{lepton flavour asymmetry} = 10^{-12}$

$|Y_B| = 10^{-2} 10^{-1} 100$

$z = \frac{TEW}{T}$

<table>
<thead>
<tr>
<th>number density of RHN</th>
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<tr>
<td>0.2. $\times 10^{-3}$</td>
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<table>
<thead>
<tr>
<th>lepton flavour asymmetry</th>
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<tr>
<td>-3. $\times 10^{-9}$</td>
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<tr>
<td>-2. $\times 10^{-9}$</td>
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<tr>
<td>-1. $\times 10^{-9}$</td>
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</tbody>
</table>

$z = \frac{TEW}{T}$

$10^{-2}$ $10^{-1}$ $10^0$
Full Testability?

- full testability requires a complete determination of the RHN parameters
- in principle possible from a measurement of all mixing angles and masses
- “flavoured” mixing angles $U^2_a$ can be measured via flavour ratios
- the two remaining parameters $(\Delta M, \text{Re}(\omega))$ could be probed by:
  - neutrinoless double $\beta$ decay requires small $M$, large $\Delta M$ and $U^2$
    
    [Hernández/Kekic/López-Pavón/Salvado 1606.06719, Eijima/Drewes 1606.06221]
  - $CP$ violation requires $\Delta M$ comparable to the decay width $\Gamma$ obscured by $\Delta M_{\theta\theta}$
  - lepton number violation requires $\Delta M$ comparable to $\Gamma$ obscured by $\Delta M_{\theta\theta}$
  - oscillations in the lab: tiny $\Delta M_{\text{phys}}$
[Antusch/Cazzato/Fischer 1709.03797]
\[
\begin{align*}
U_e^2/U^2 & = 0.451 \\
U_\mu^2/U^2 & = 0.574 \\
\delta & = 261^\circ
\end{align*}
\]
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