Resonant pp → W⁺Zjj events at the LHC from a unitarized study of the EChL

Claudia García-García

Instituto de Física Teórica, Universidad Autónoma de Madrid/CSIC

claudia.garcia@uam.es

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Based on JHEP 1711 (2017) 098, [arXiv:1707.04580]

R. Delgado, A. Dobado, D. Espriu, **CGG**, M.J. Herrero, X. Marcano & J.J. Sanz-Cillero

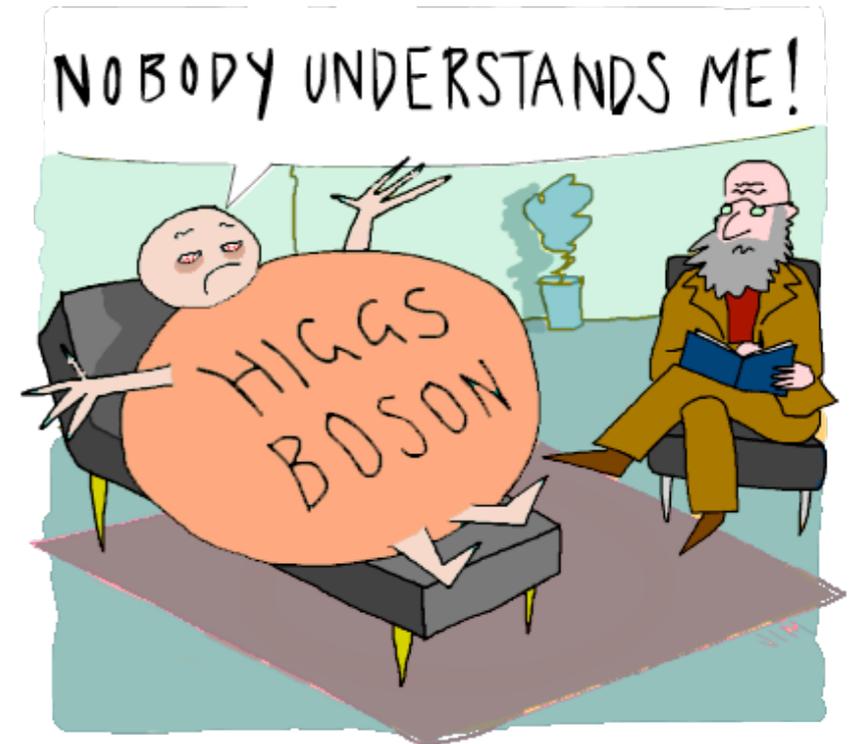53rd Rencontres de Moriond - EW 2018

Motivation

What is the dynamical generation of EWSB?

Effective Theories for EWSB

- Describe **dynamical generation of EWSB**
  - ↳ Strong Dynamics?
    - ↳ **Resonances** predicted!

- **Model independent**

NOBODY UNDERSTANDS ME!
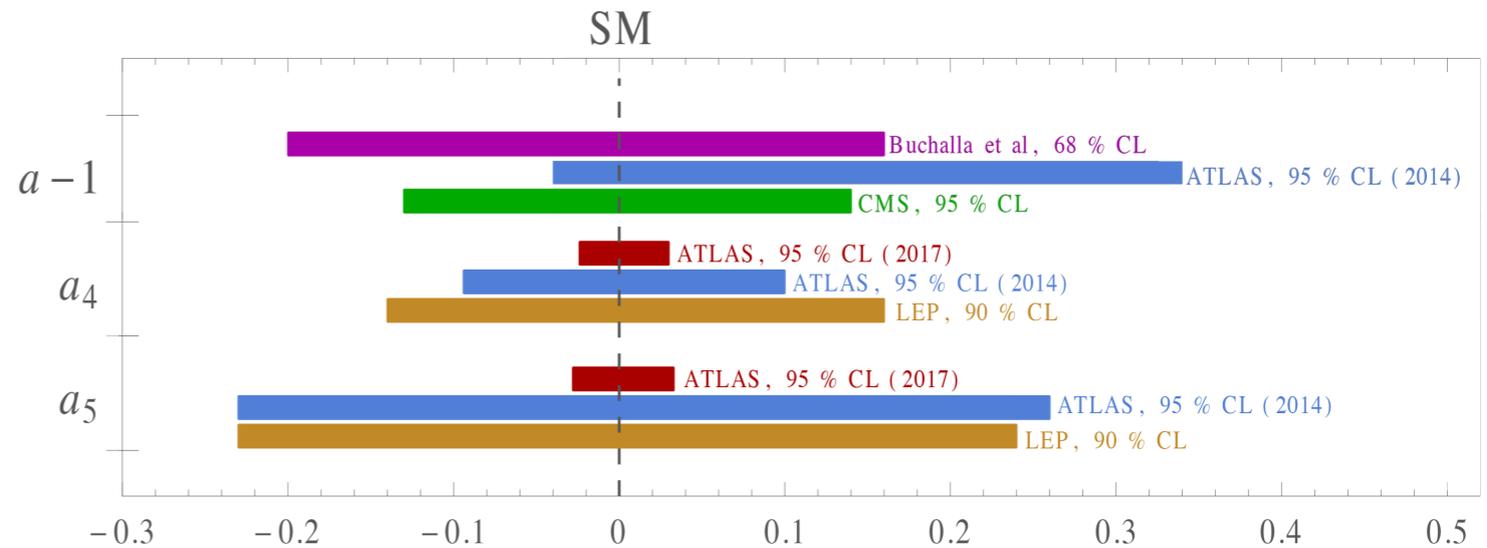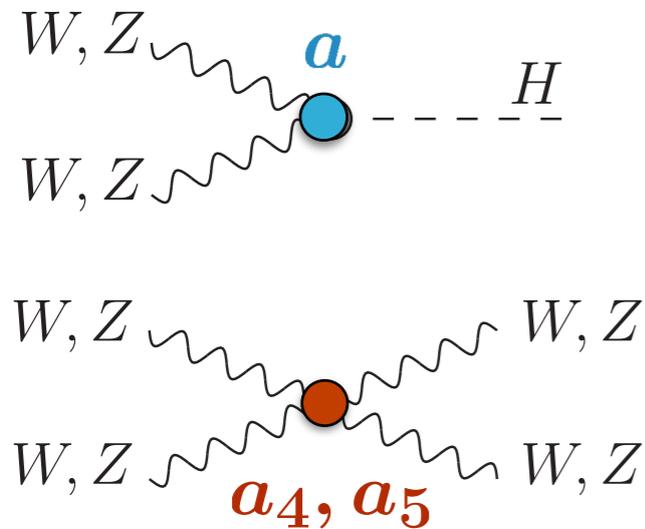
HIGGS BOSON

Vector Boson Scattering @ the LHC

- New Pheno in EWSB sector ⟶ **New Pheno** in EW boson interactions

- **Vector Boson Scattering** (VBS) very sensitive to **New Physics**

- **Searches for VBS** planned @ the **LHC**!

[A. C. Longhitano, 1980; T. Appelquist et al., 1980]   EChL copy of ChPT in QCD

- **EChL** symmetries: Gauge $\longrightarrow$ $SU(2)_L \times U(1)_Y$ and EW Chiral $\longrightarrow$ $SU(2)_L \times SU(2)_R$

- New interactions in EW sector $\longrightarrow$ Controlled by *chiral* parameters a a_4 a_5

[A. C. Longhitano, 1980; T. Appelquist et al., 1980]

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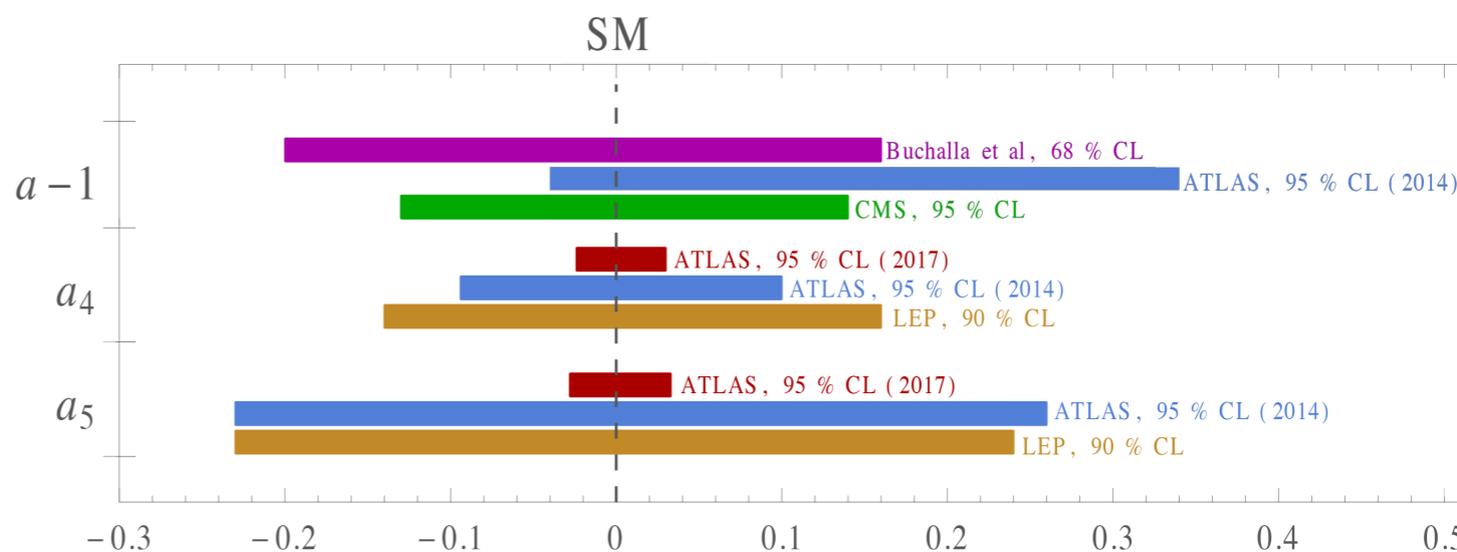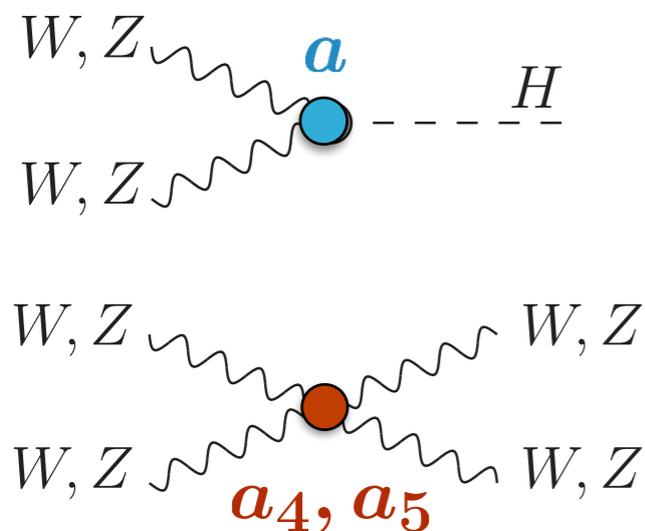- New interactions in EW sector $\longrightarrow$ Controlled by *chiral* parameters a a_4 a_5

- Expansion in powers of p $\xrightarrow{\text{typically}}$ **Violation of Unitarity**

- Unitarization via the **Inverse Amplitude Method** (IAM) $\longrightarrow$ **generates resonances dynamically**

[Tran N. Truong, Phys. Rev. Lett. 61 (1988)2526]

Unitarized Partial Waves

$$a_{IJ}^{\text{IAM}}(s) = \frac{(a_{IJ}^{(0)}(s))^2}{a_{IJ}^{(0)} - a_{IJ}^{(1)}}$$

$M_V,\ \Gamma_V$ depend on a a_4 a_5

We select vector scenarios

[A. C. Longhitano, 1980; T. Appelquist et al., 1980]

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Unitarized **Partial Waves**

$$a_{IJ}^{\mathrm{IAM}}(s) = \frac{(a_{IJ}^{(0)}(s))^2}{a_{IJ}^{(0)} - a_{IJ}^{(1)}}$$

Difficult to Implement in Monte Carlo!

$M_V, \ \Gamma_V$ depend on a a_4 a_5

We select vector scenarios

[A. C. Longhitano, 1980; T. Appelquist et al., 1980]

EChL copy of ChPT in QCD

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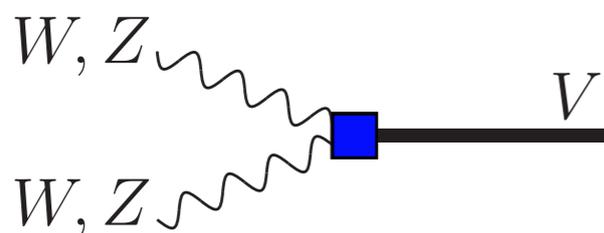
- New interactions in EW sector ⟶ Controlled by *chiral* parameters a a_4 a_5
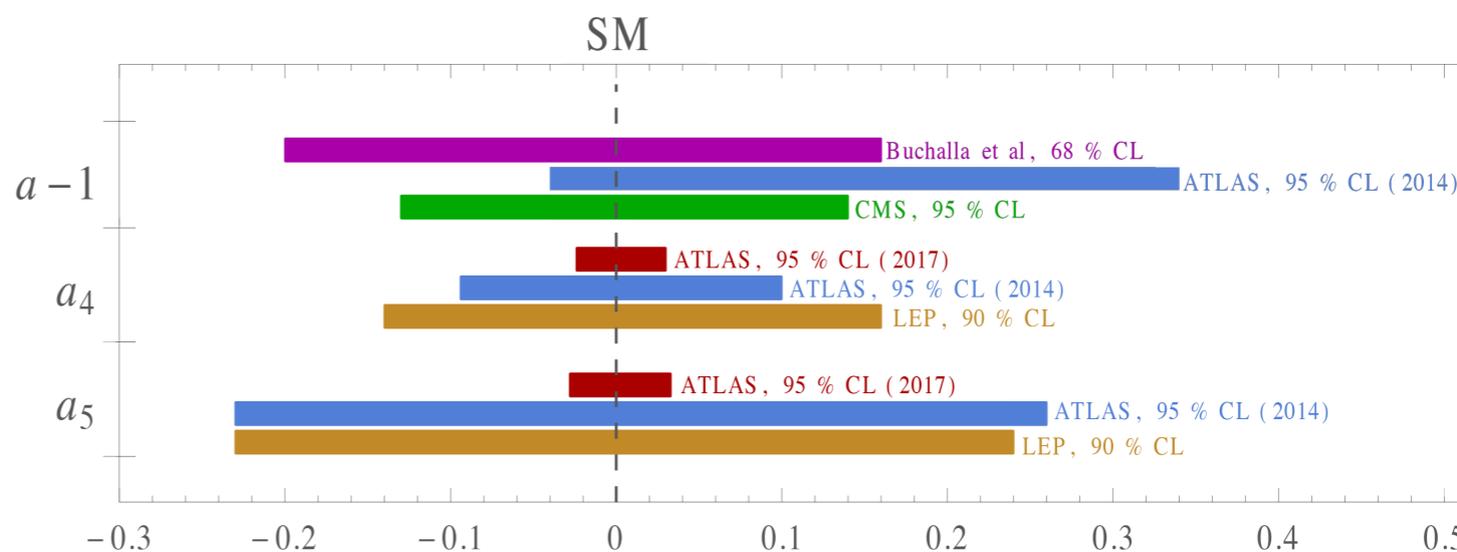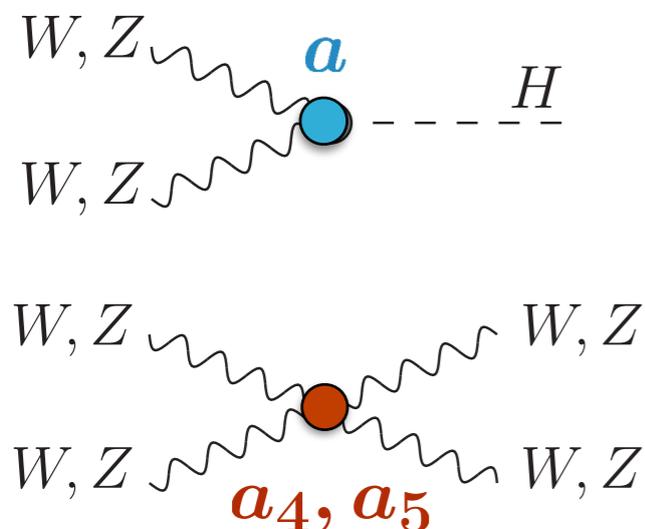
- Expansion in powers of p ⟶ typically ⟶ **Violation of Unitarity**

- Unitarization via the **Inverse Amplitude Method** (IAM) ⟶ generates **resonances** dynamically

Efective vertex to mimic the IAM vector resonances!

introduced in MG5

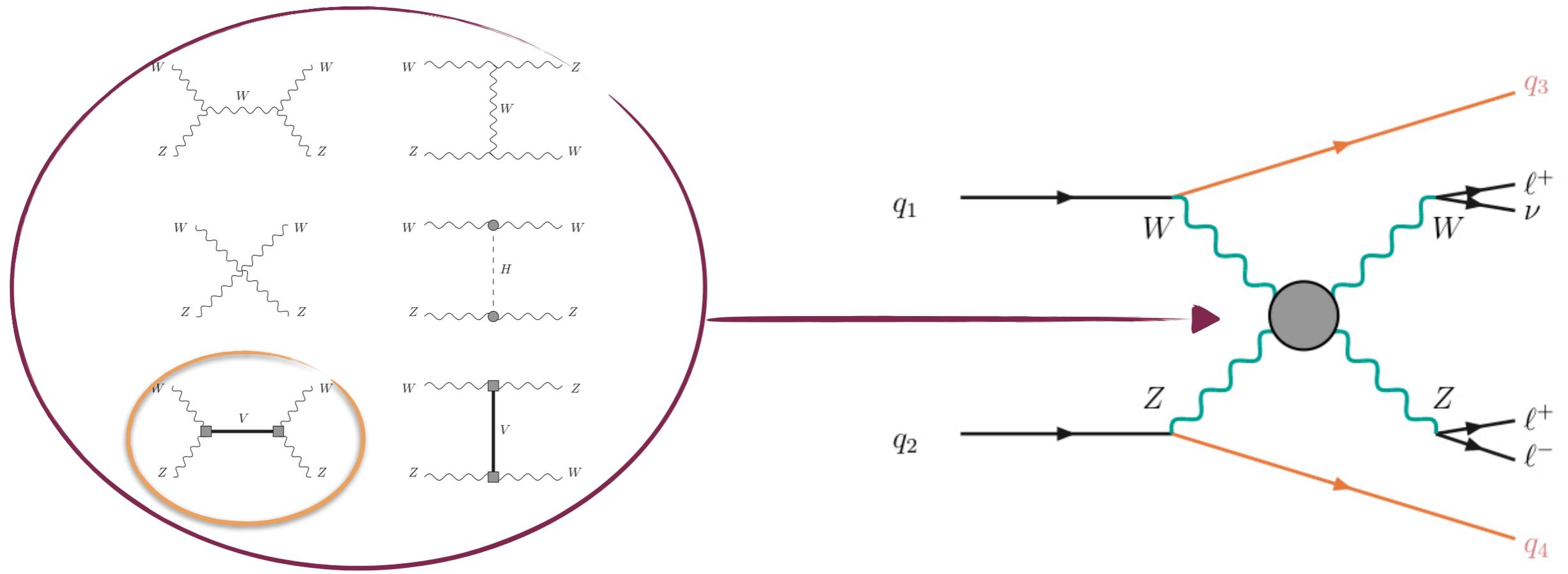$M_V, \ \Gamma_V$ depend on a a_4 a_5

We select vector scenarios

- We study **charged vector resonances**, V, from a triplet, $V^{\pm}, V^0$

- **W⁺Z** channel very promising → Mediated only by V
  → No severe QCD backgrounds
  → Clean leptonic signal

- Signals
  $$pp \rightarrow W^+ Z jj$$
  $$pp \rightarrow \ell_1^+ \ell_1^- \ell_2^+ \slashed{E}_T jj$$
  $$pp \rightarrow JJjj$$

- We study **charged vector resonances**, V, from a triplet, $V^{\pm}, V^0$

- **W⁺Z** channel very promising
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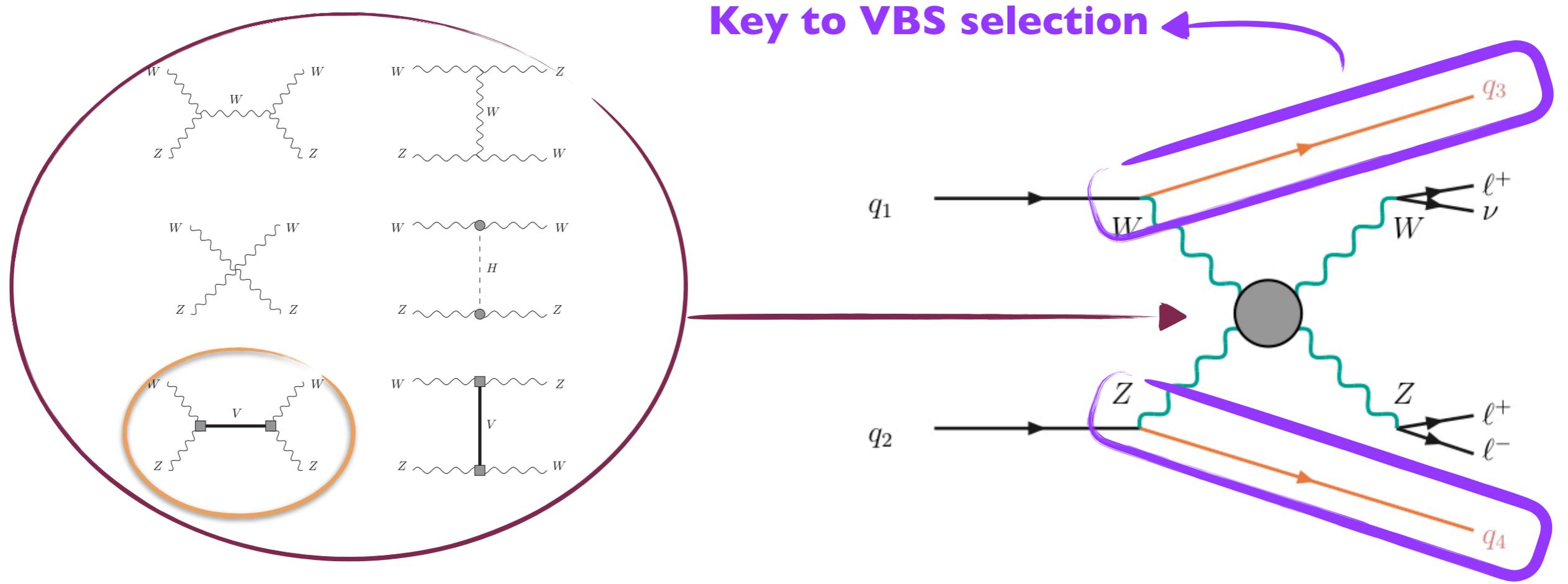 - → No severe QCD backgrounds
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- Signals

$$pp \to W^+ Z jj$$

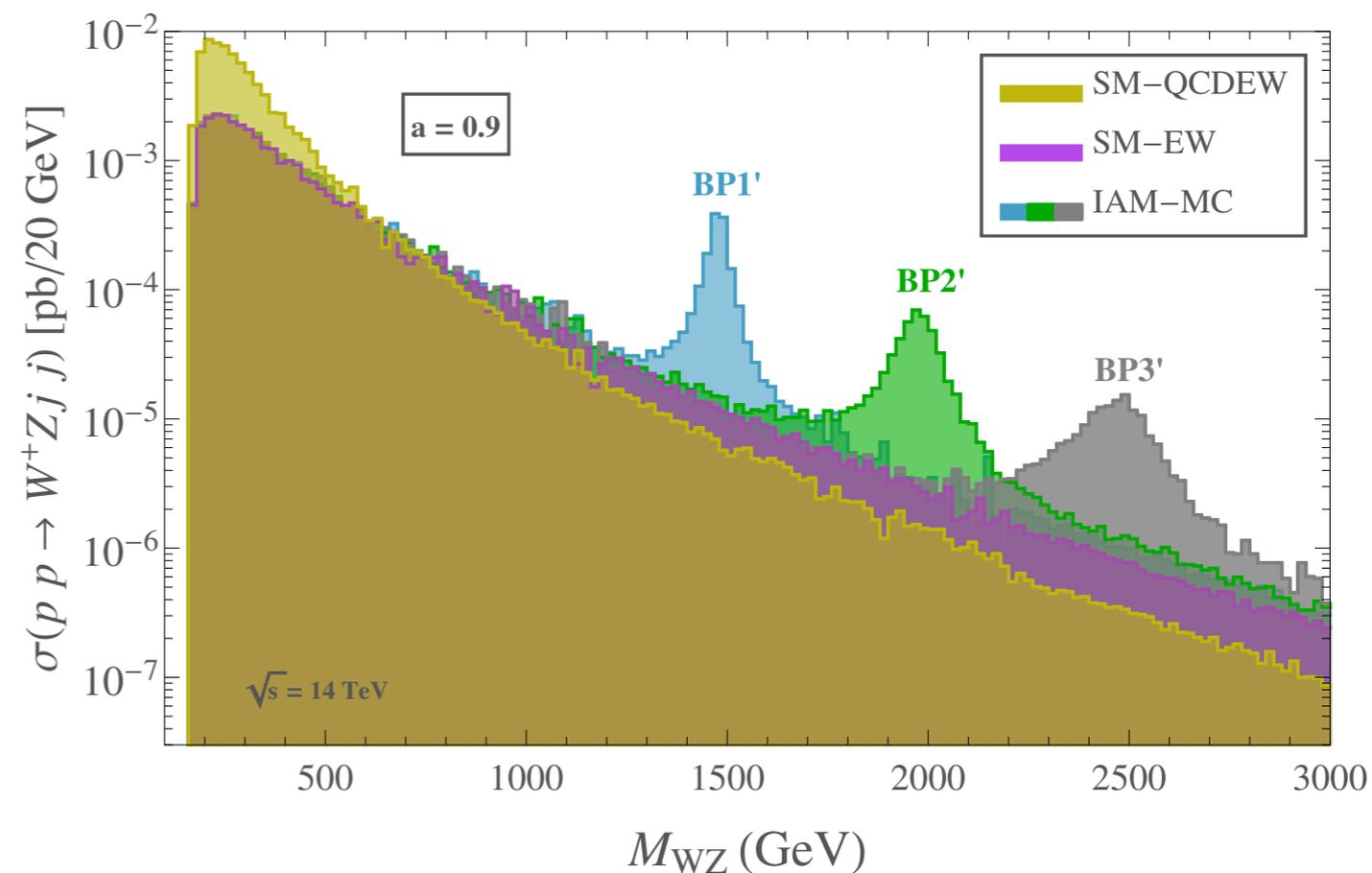$$pp \to \ell_1^+ \ell_1^- \ell_2^+ \not{E}_T jj$$

$$pp \to JJjj$$

Key to VBS selection
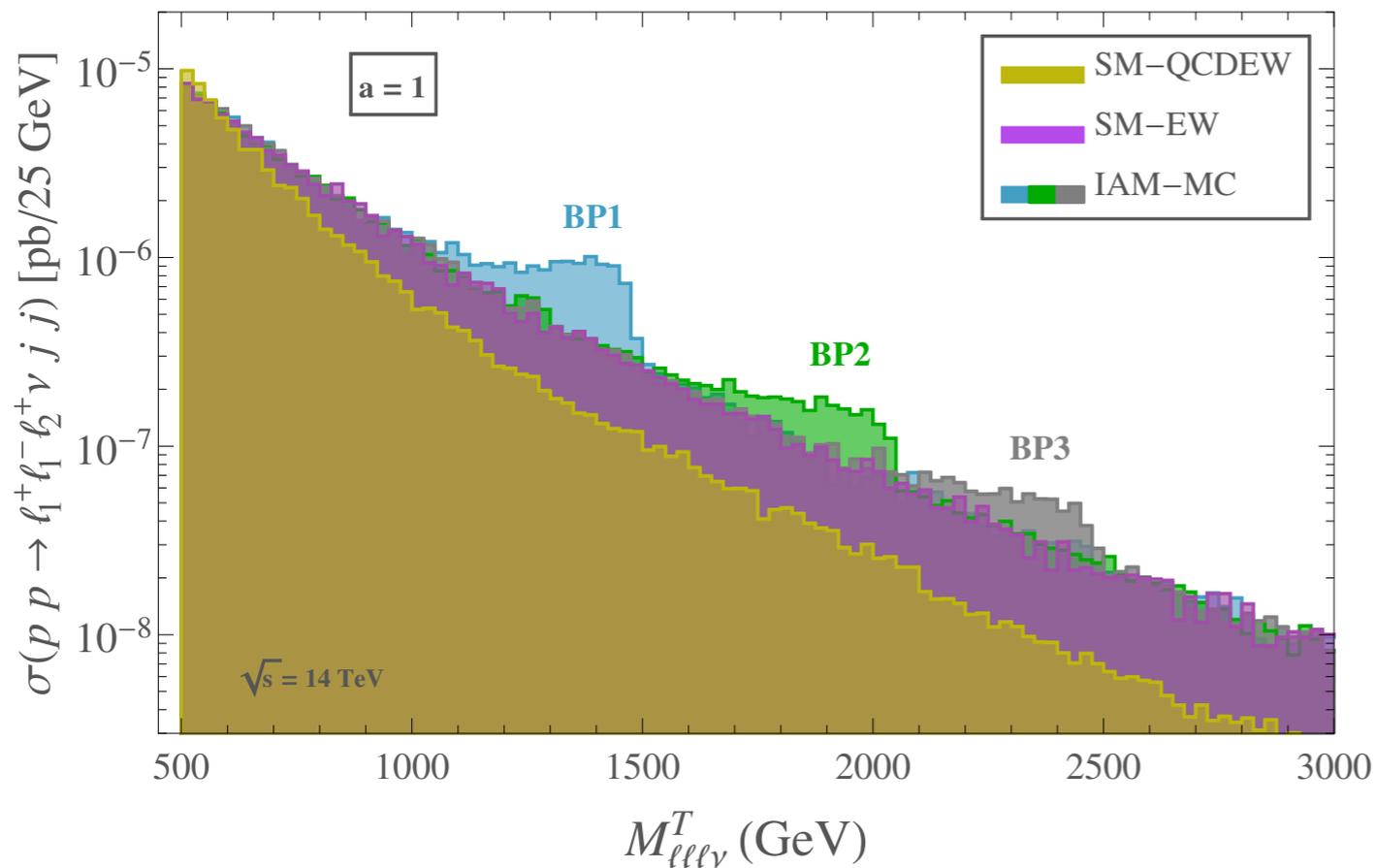
- Very **clear resonant peaks** on top of SM background

- Many **events available** for LHC luminosities if W & Z are well reconstructed

- Some scenarios are **still** very **visible** above SM background

- Very clean signal in the leptonic channel
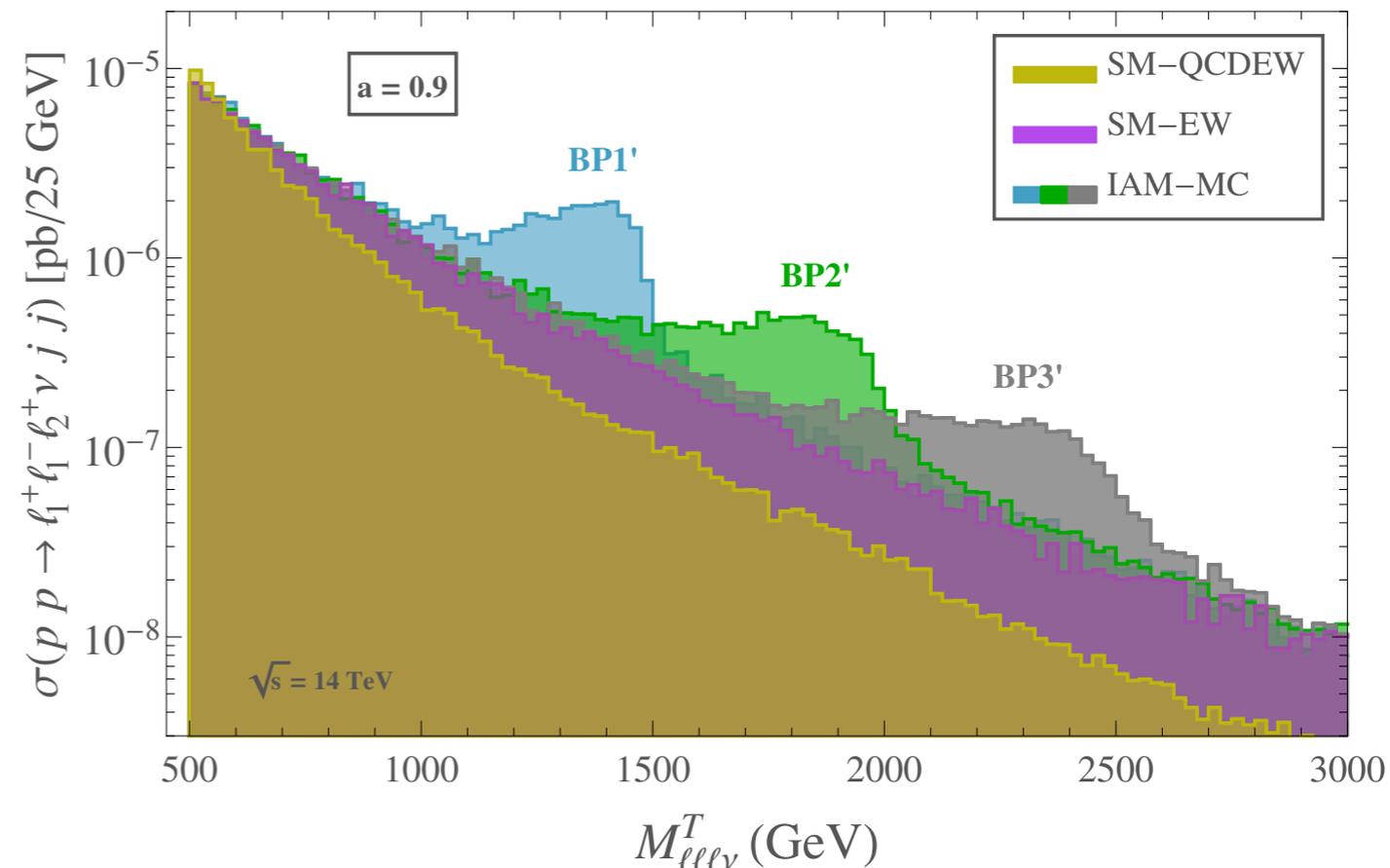
Sensitivity to a in W⁺Zjj final state

- Given M_V:
  significance **increases** with $(a-1)$

- Given $(a-1)$:
  significance **decreases** with M_V

LHC sensitive to $a \in [0.9, 1]$ for $M_V \in [1.5, 2.5]$ TeV and $\mathcal{L} = 300$ fb^{-1}

Signal & Background definitions

$$S = N_{ev}^{IAM-MC} - N_{ev}^{SM(QCD+EW)}$$

$$B = N_{ev}^{SM(QCD+EW)}$$

stat. significance

$$\sigma^{stat} = \frac{S}{\sqrt{B}}$$

Legend:
- M_V = 1.5 TeV
- M_V = 2 TeV
- M_V = 2.5 TeV

$\mathcal{L} = 3000$ fb^{-1}

$M_V \pm 0.5\,\Gamma_V$

$M_V \pm 2\,\Gamma_V$

y-axis: σ_{WZ}^{stat}

x-axis: a
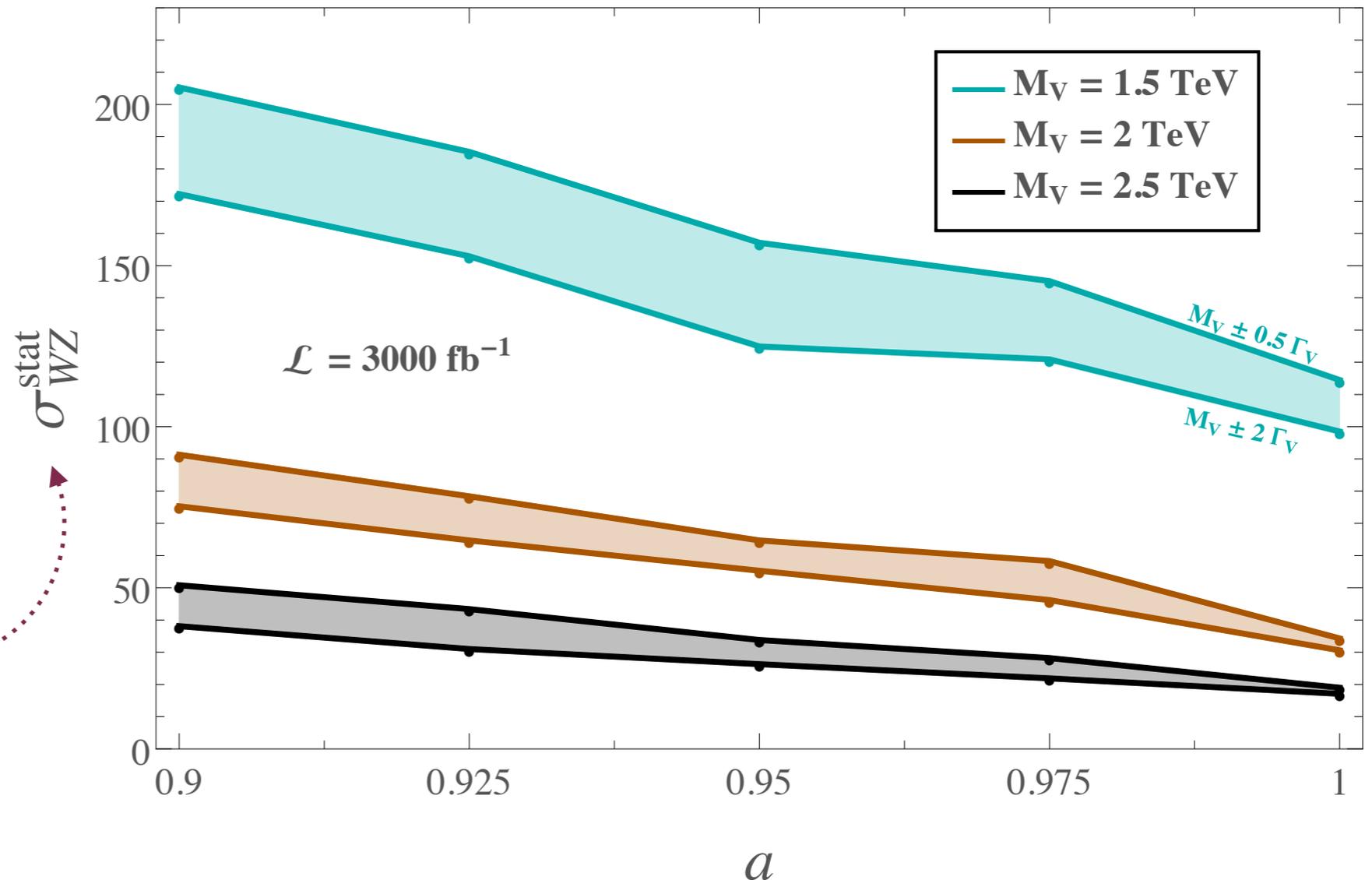
$$pp \to \ell_1^+ \ell_1^- \ell_2^+ \not{E}_T jj \qquad \mathcal{L} = 3000 \text{ fb}^{-1}$$

$a \in [0.9, 1] \longleftrightarrow M_V = 1.5 \text{ TeV}$

$a \in [0.94, 1] \longleftrightarrow M_V = 2 \text{ TeV}$

Poor sensitivity for $M_V = 2.5 \text{ TeV}$
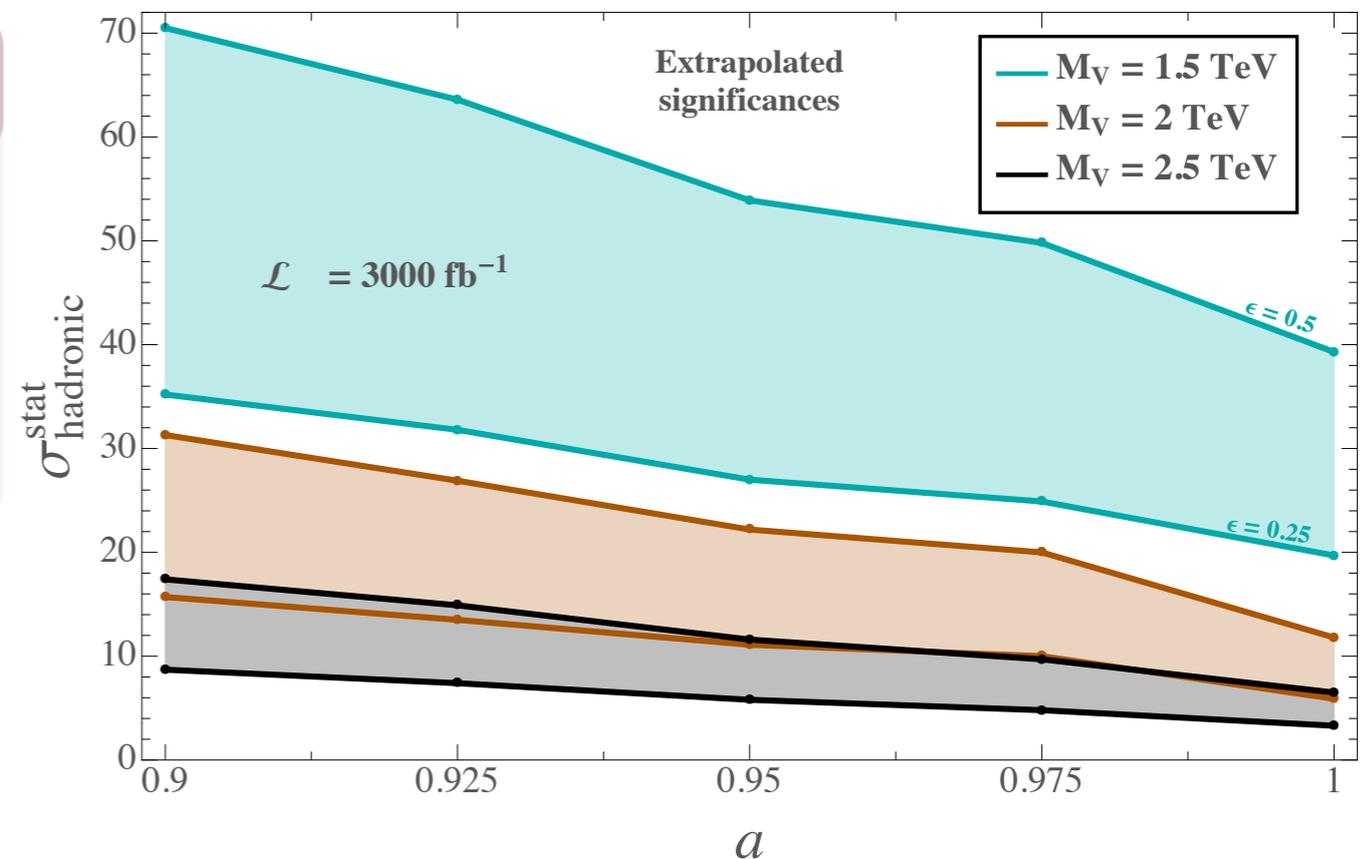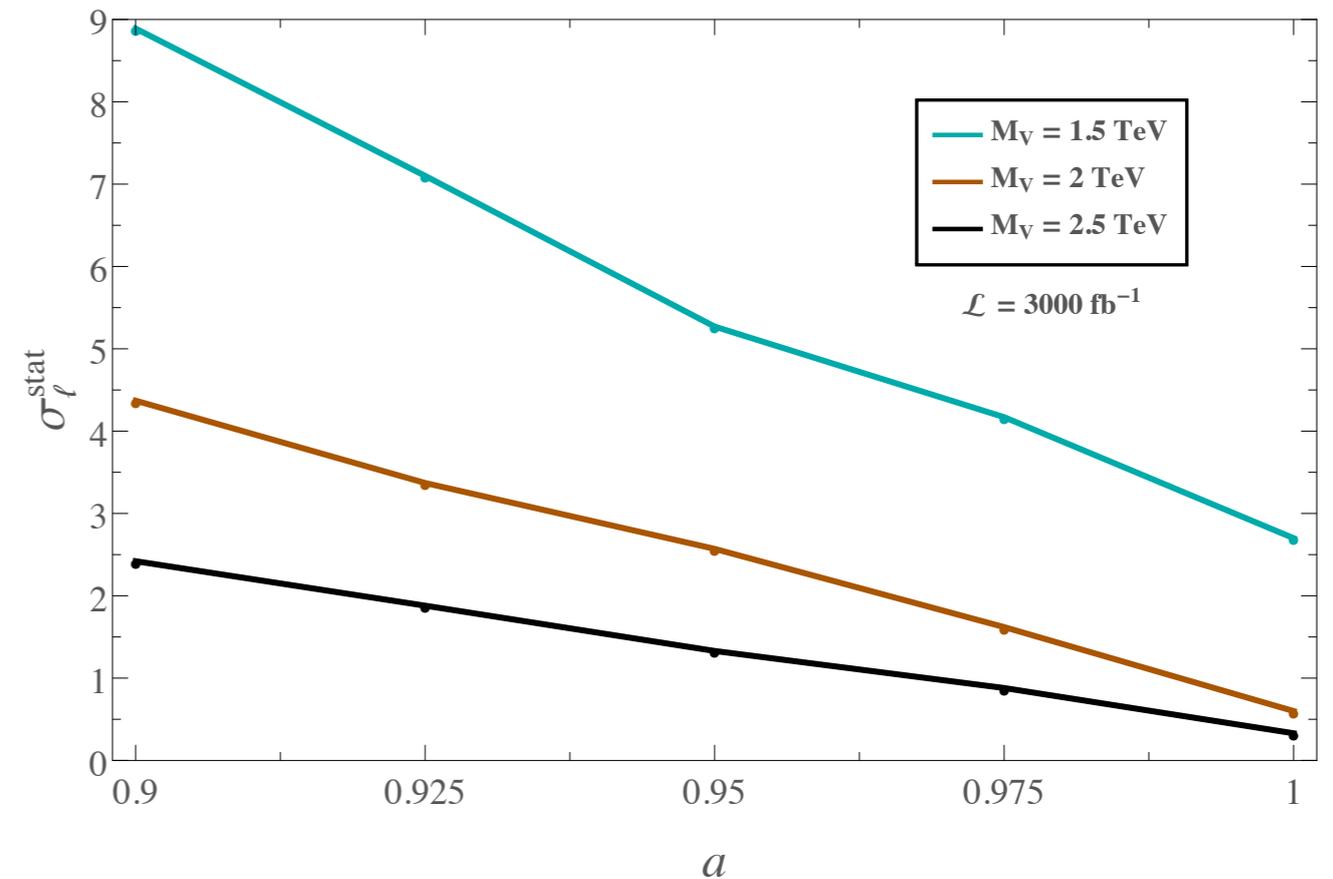
- Clean and controlled
- Requires High Luminosity

$$pp \to JJjj \quad \textbf{ESTIMATES} \quad \mathcal{L} = 300 \text{ fb}^{-1}$$

$a \in [0.9, 1] \longleftrightarrow M_V = 1.5 \text{ TeV}$

$a \in [0.975, 1] \longleftrightarrow M_V = 2 \text{ TeV}$

$a \in [0.925, 1] \longleftrightarrow M_V = 2.5 \text{ TeV}$

- Very high significances!
- Depends on reconstruction efficiency

Conclusions

- **VBS** optimal place to **test deviations** introduced by the **EChL**

- **Dynamically generated vector resonances** from unitarized EChL can be seen **@ LHC**

  → Very high statistical significance in WZ final state

  → Promising results in leptonic channel for some scenarios and HL-LHC

  → High sensitivity in final states with fat jets

- Study Resonances of 1.5 - 2.5 TeV → broad **sensitivity** to **EChL params**

  → Depending on:

  - final state
  - luminosity

- Improving selection cuts + study other distributions might improve the results

Conclusions

- **VBS** optimal place to **test deviations** introduced by the **EChL**

- **Dynamically generated vector resonances** from unitarized EChL can be seen **@ LHC**

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  # THANK YOU!

- Improving selection cuts + study other distributions might improve the results

Back up Slides

The Electroweak Chiral Lagrangian

- Symmetries: **Gauge** $\longrightarrow$ $SU(2)_L \times U(1)_Y$ and **EW Chiral** $\longrightarrow$ $SU(2)_L \times SU(2)_R$
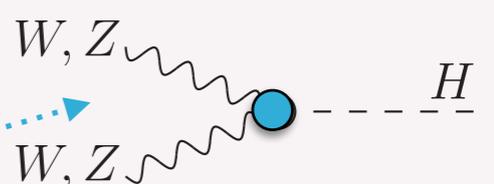
  EChL copy of ChPT in QCD

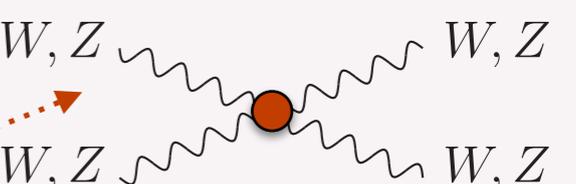- **Light d.o.f. :** $w^\pm, z$ $W^\pm, Z$ h

- Building blocks:

$$D_\mu U = D_\mu \left(e^{\frac{iw^a \tau^a}{v}} \right) \qquad \hat{W}_{\mu\nu}, \hat{B}_{\mu\nu} \qquad \mathcal{F}(h) = 1 + 2a\frac{h}{v} + b\left(\frac{h}{v}\right)^2 \qquad \mathcal{V}_\mu = (D_\mu U)U^\dagger$$

$$\mathscr{L}_{EChL} = \mathscr{L}_2 + \mathscr{L}_4 \qquad \textbf{Relevant terms \& chiral parameters for VBS}$$
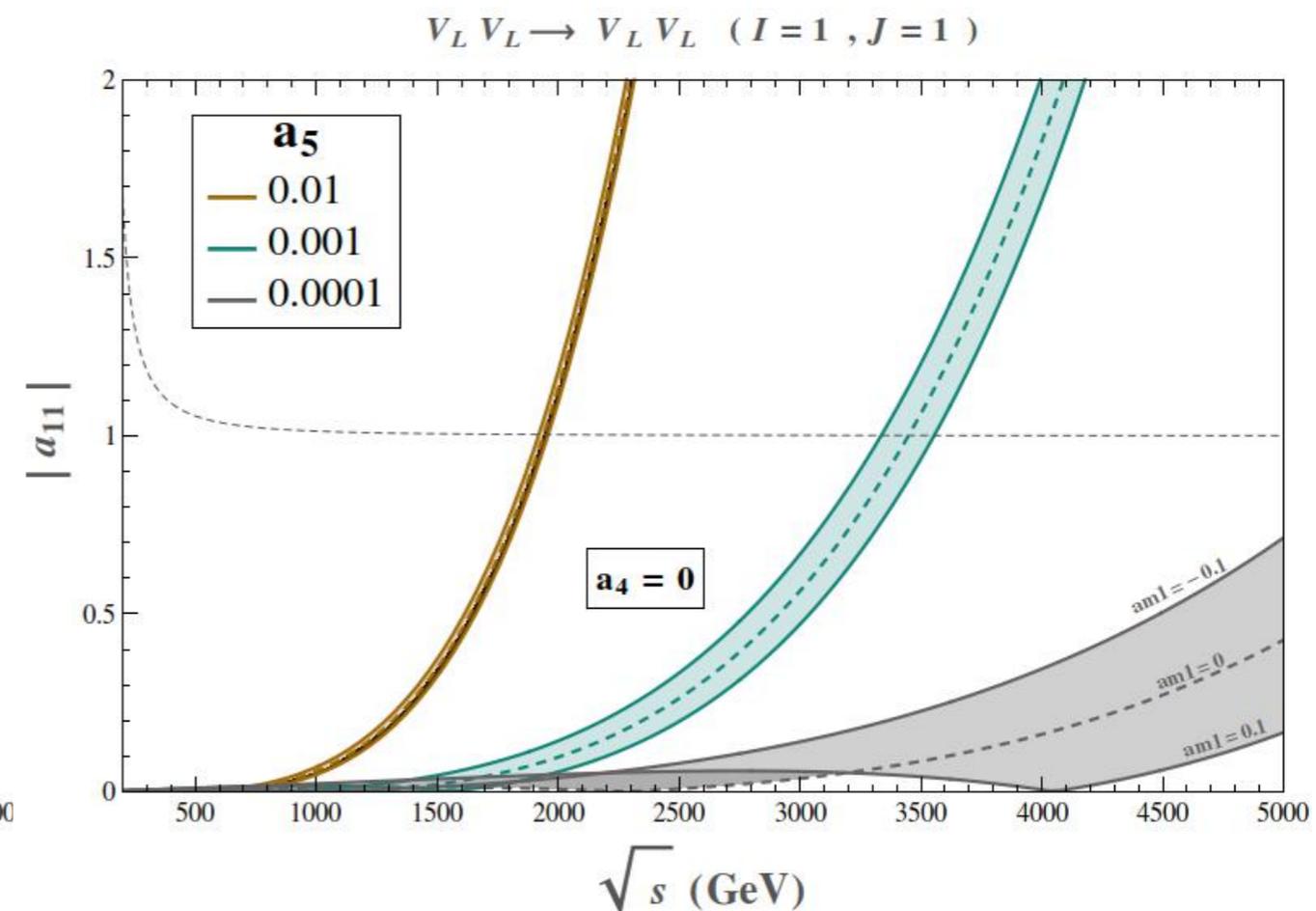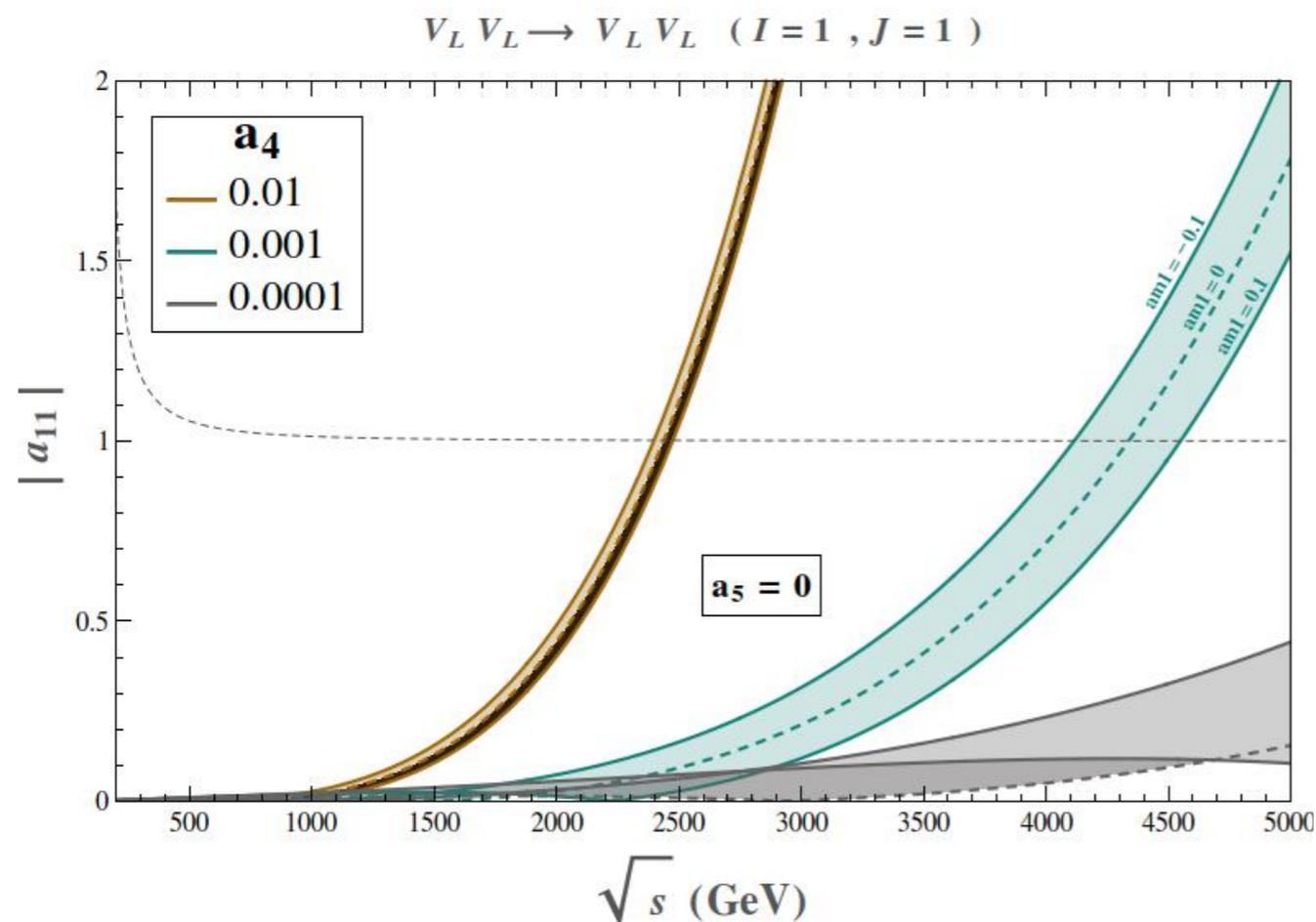
[A. C. Longhitano, 1980; T. Appelquist et al., 1980]

$$b = a^2$$

$$\mathscr{L}_2 = \frac{v^2}{4}\left[1 + 2a\frac{H}{v} + b\frac{H^2}{v^2} \right] \mathrm{Tr}\left(D^\mu U^\dagger D_\mu U \right) + \frac{1}{2}\partial^\mu H\,\partial_\mu H + \dots,$$

W, Z W, Z H

$$\mathscr{L}_4 = a_4\left[\mathrm{Tr}(\mathcal{V}_\mu \mathcal{V}_\nu)\right]\left[\mathrm{Tr}(\mathcal{V}^\mu \mathcal{V}^\nu)\right] + a_5\left[\mathrm{Tr}(\mathcal{V}_\mu \mathcal{V}^\mu)\right]\left[\mathrm{Tr}(\mathcal{V}_\nu \mathcal{V}^\nu)\right]$$

W, Z W, Z W, Z W, Z

Unitarity Violation

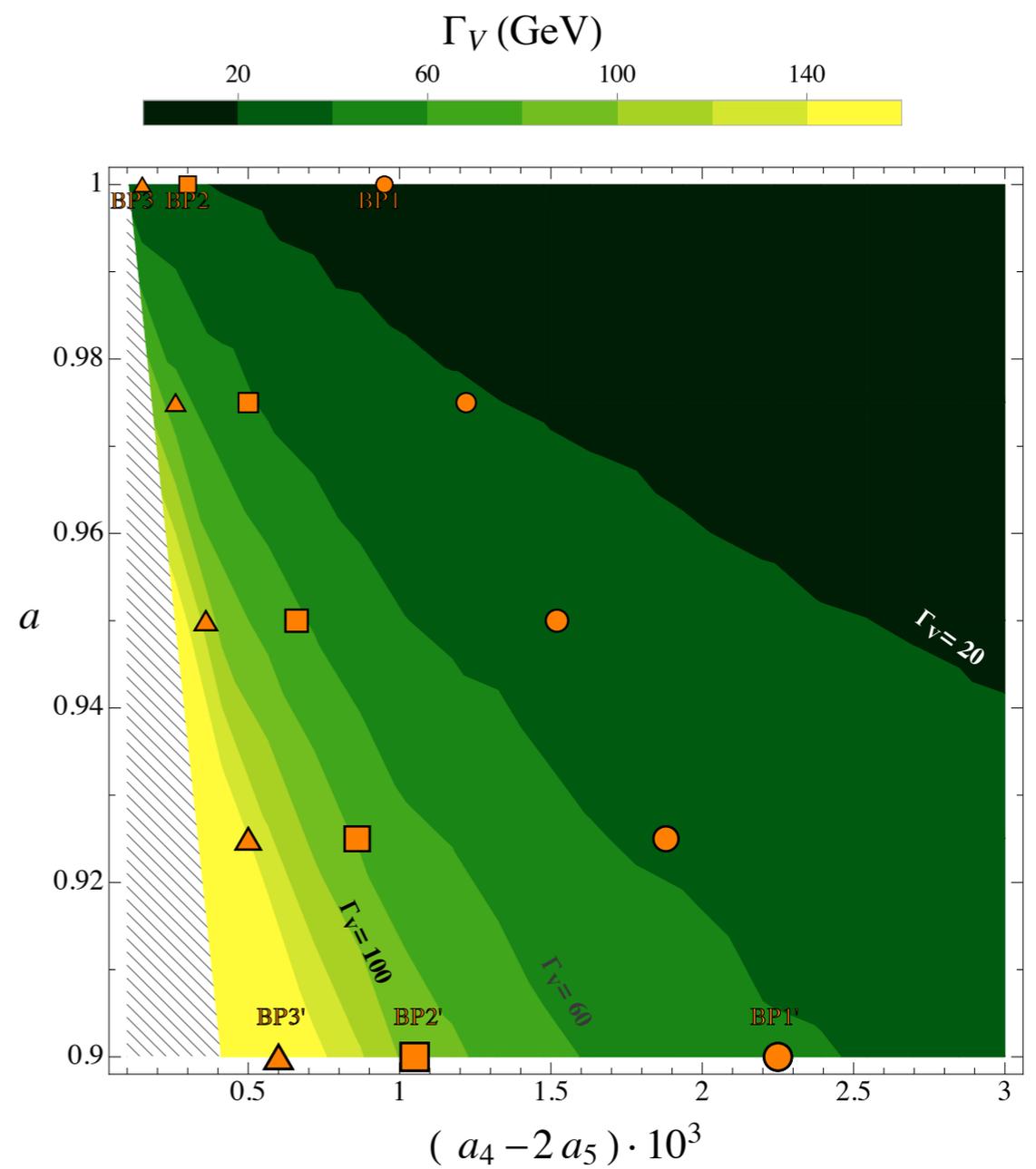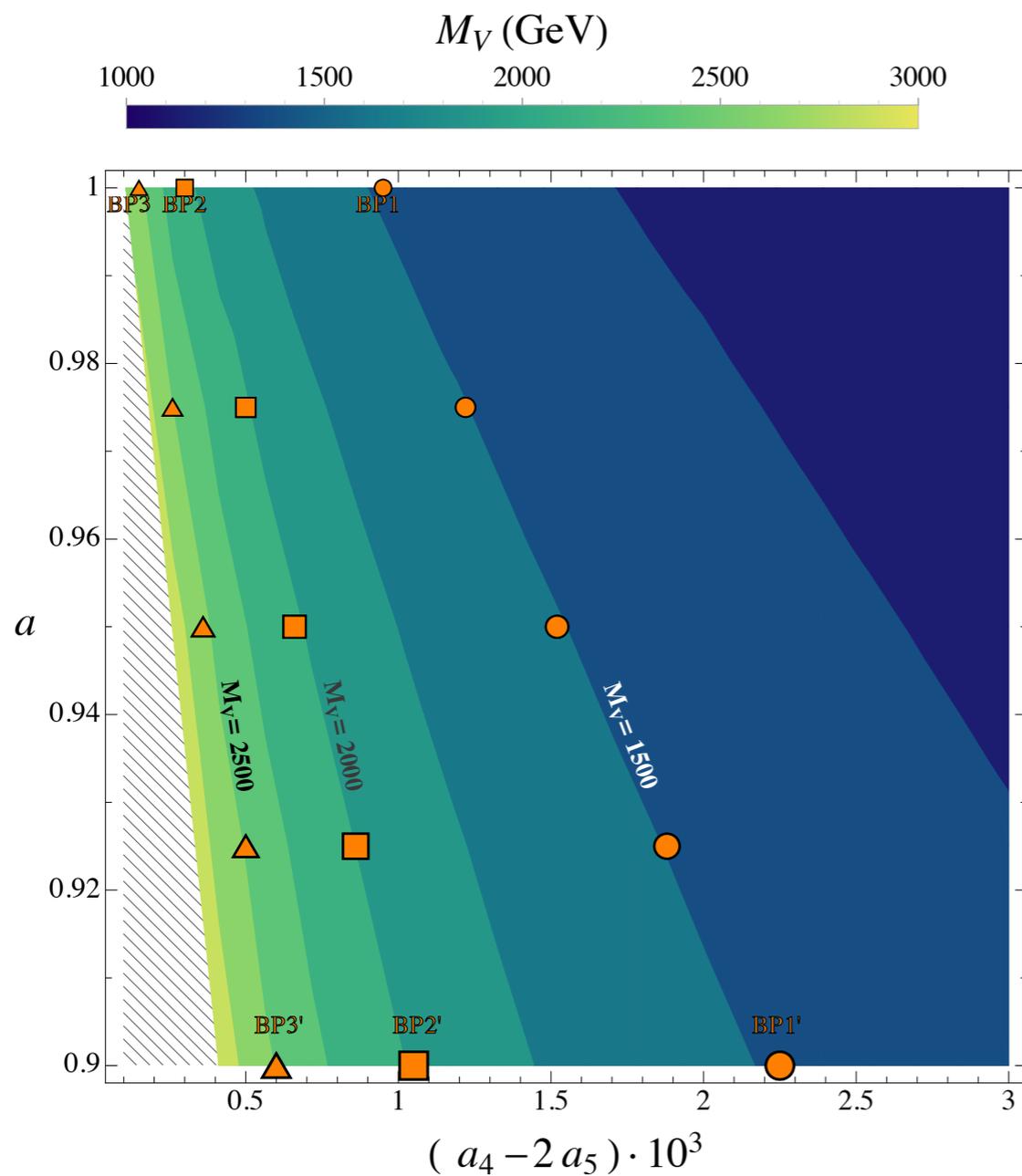- Energy at which unitarity violation occurs is very sensitive to chiral param. values

- At tree level, for the values considered $a \in [0.9, 1]$, $a_4, a_5 \in [10^{-3}, 10^{-4}]$ (allowed by exp.) unitarity violation happens at the few TeV scale

Selection of Scenarios

- 5 different values of a for three lines of fixed mass value M_V

- Resonance properties M_V, Γ_V depend on chiral parameters

[S. Dawson, G. Valencia, Phys. Lett. B 246 (1990) 156]

- Top loop contribution decreases with energy

- Negligible with respect to Goldstone boson loops above 1 TeV

- Subleading effects on resonance properties

- Indirect constraints: $b \in [0, 2]$ [R.L. Delgado, A. Dobado et al., Phys. Rev. Lett. 114 (2015) 221803]
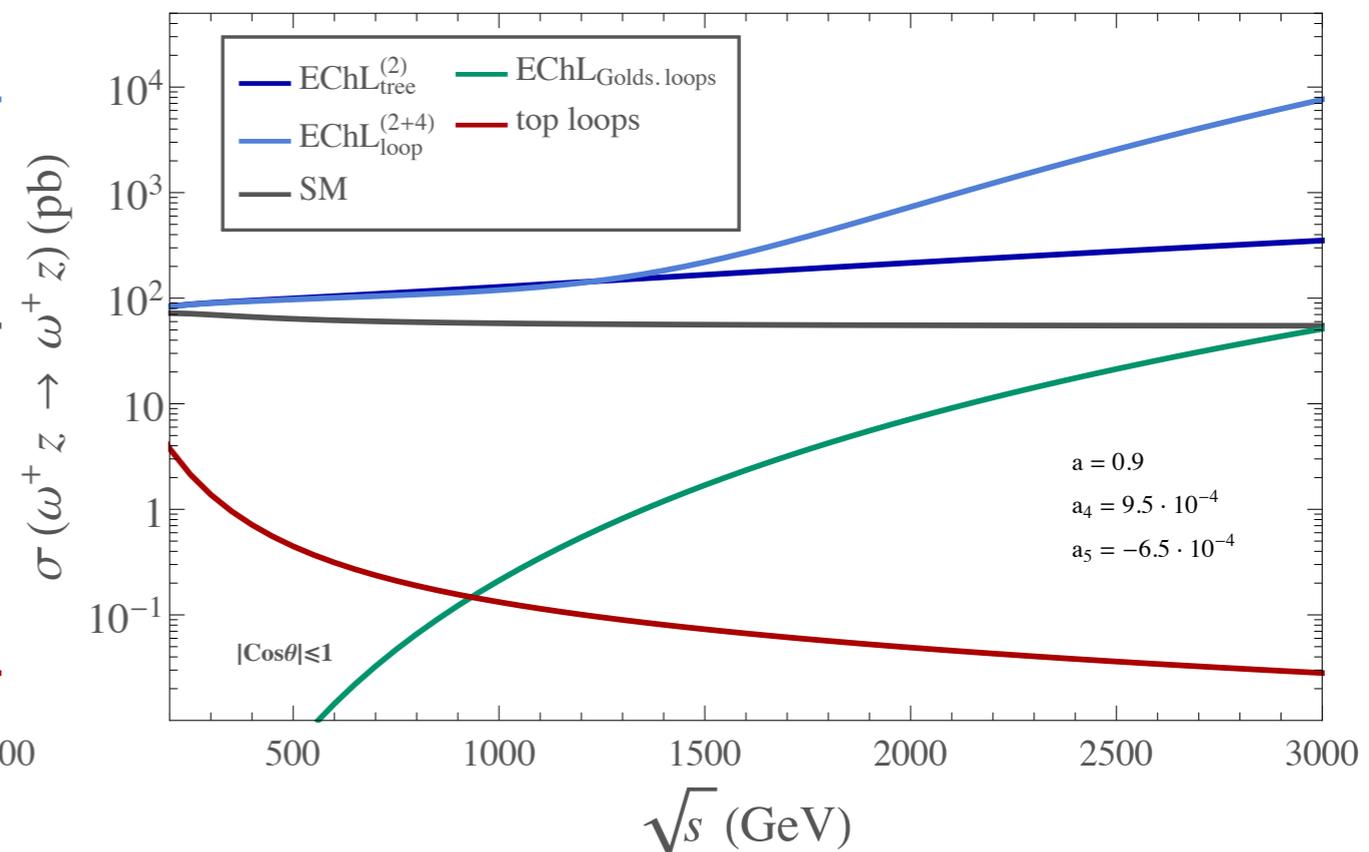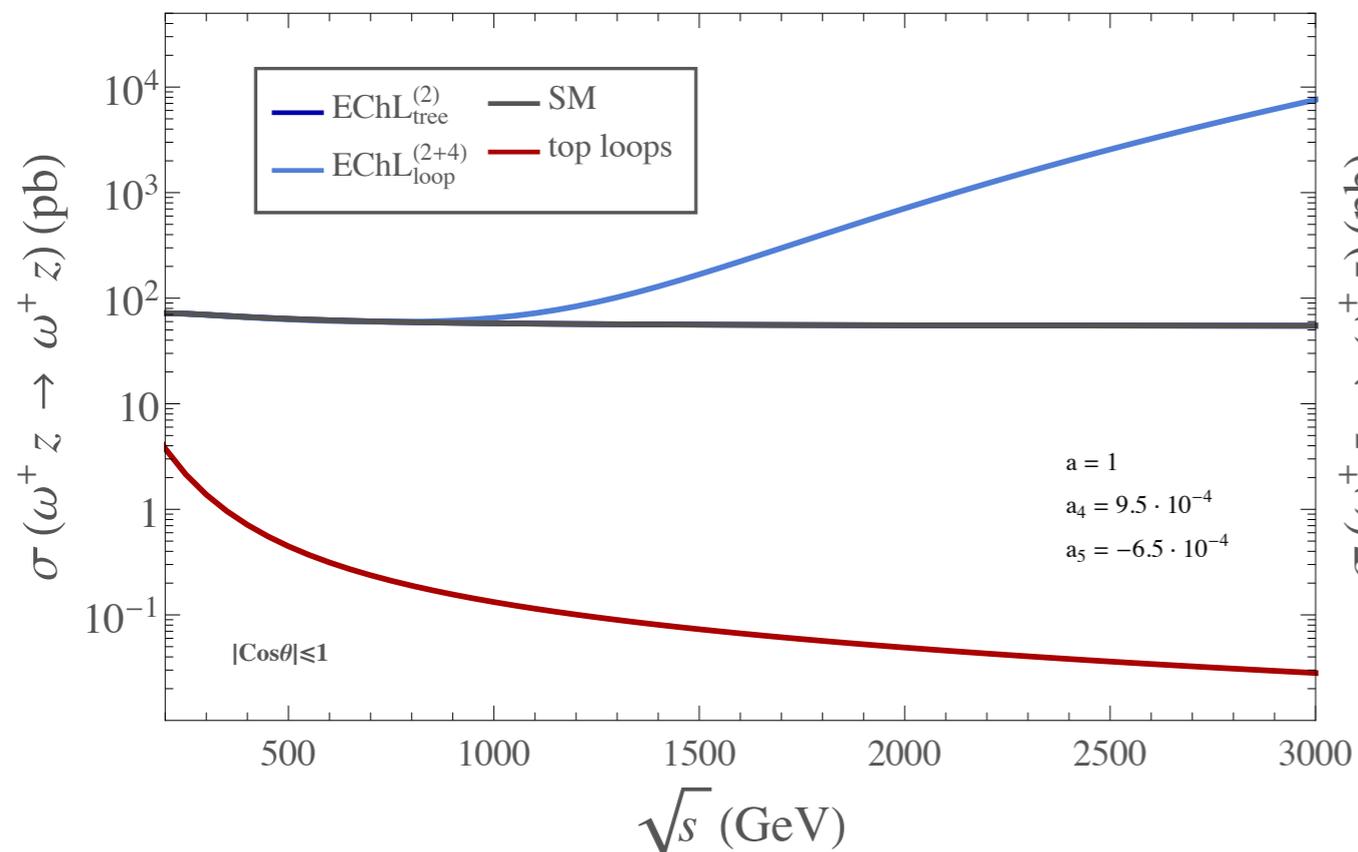
- Relaxing condition $b = a^2$ does not modify properties of resonances

○ MG5 ⟶ Prediction for resonance observables in W⁺Z ⟶ Works with FR as input

Chiral inv. Lagrangian to mimic IAM vector resonances

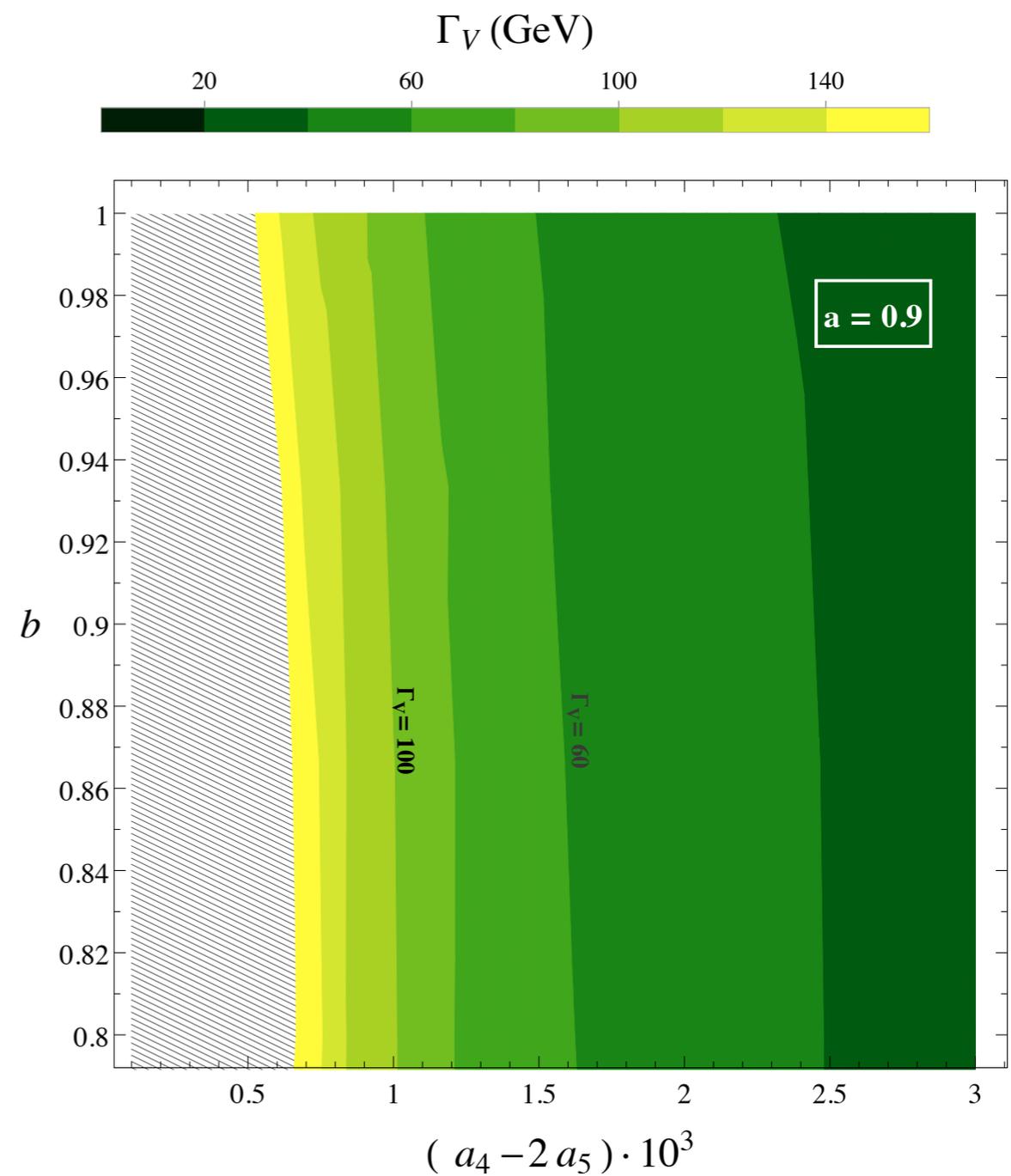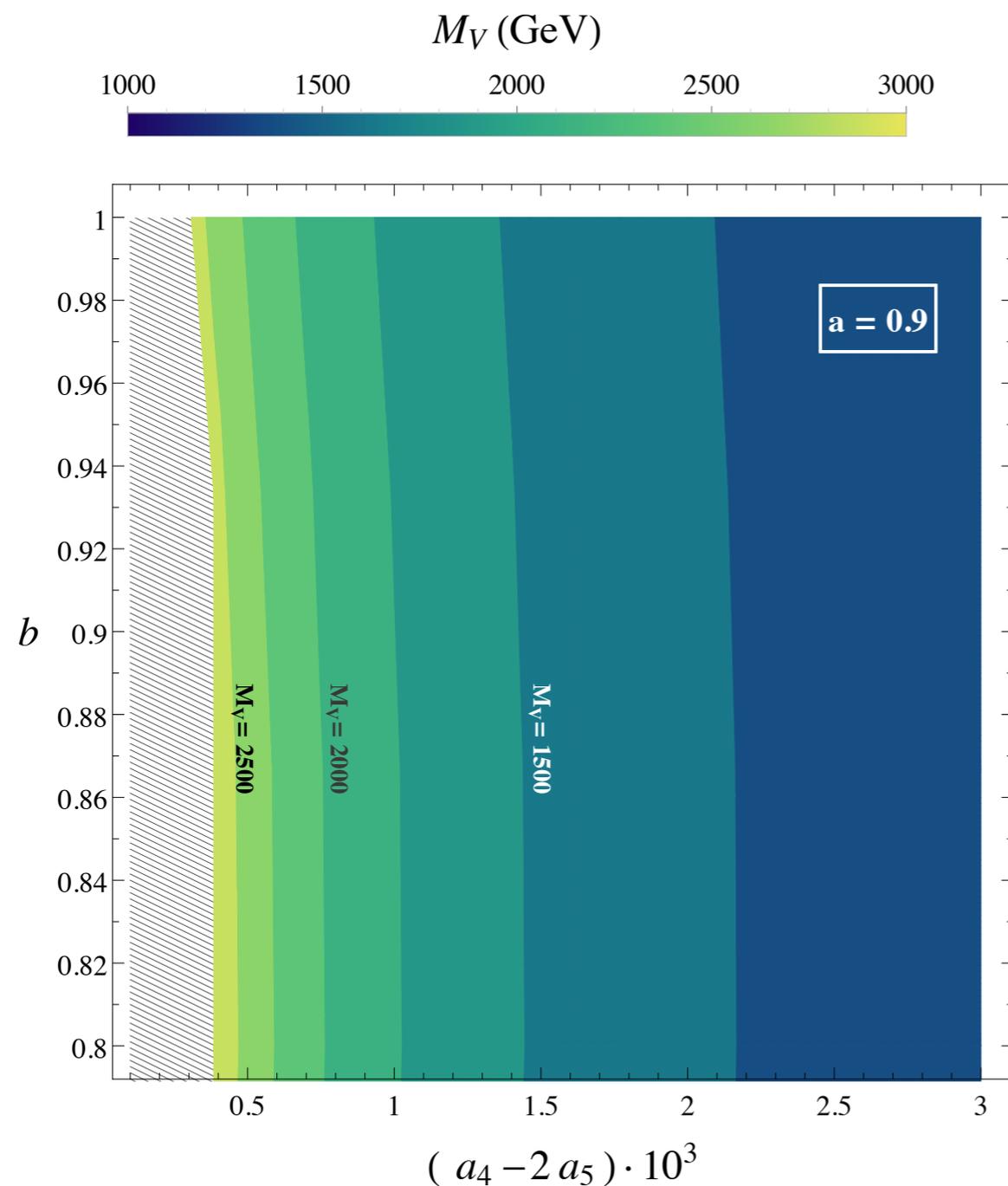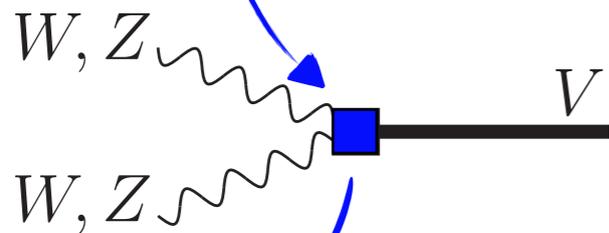$$\mathscr{L}_V = \frac{i2g_V}{v^2}\left[m_W^2\, V^{0\ \mu\nu}W_\mu^+ W_\nu^- + m_W\, m_Z\, V^{+\ \mu\nu}W_\mu^- Z_\nu + m_W\, m_Z\, V^{-\ \mu\nu}Z_\mu W_\nu^+\right] + \dots$$

From Proca Lagrangian [G. Ecker, et al., Phys. Lett. B223, 425 (1989)]

$$\mathscr{L}_{\mathrm{IAM-MC}} = \mathscr{L}_2 + \mathscr{L}_V$$

W, Z

V

W, Z

Agreement with IAM unitarized Partial Waves

g_V ⟷ **Form Factor**

- EChL $\mathscr{L}_4$ recovered for $s < M_V^2$

- $g_V(M_V^2) \equiv g_V(a, b, a_4, a_5) \sim 10^{-2}$

- Froissart unitarity bound for $s > M_V^2$

$g_V^2(z) = g_V^2(M_V^2)\frac{M_V^2}{z}$; $s < M_V^2$

$g_V^2(z) = g_V^2(M_V^2)\frac{M_V^4}{z^2}$; $s > M_V^2$

$g_V(M_V) = 0.058$

Legend:
— ECLh$_{\mathrm{tree}}^{(2+4)}$
— IAM−MC
— SM

$a = 0.9$
$a_4 = 9.5\cdot 10^{-4}$
$a_5 = -6.5\cdot 10^{-4}$

$|\mathrm{Cos}\theta| \lesssim 1$

$M_W = M_Z \neq 0$

$M_V = 1479$ GeV
$\Gamma_V = 42$ GeV
$f_V = 0$

$\sigma(W_L^+ Z_L \to W_L^+ Z_L)$ (pb)

$\sqrt{s}$ (GeV)
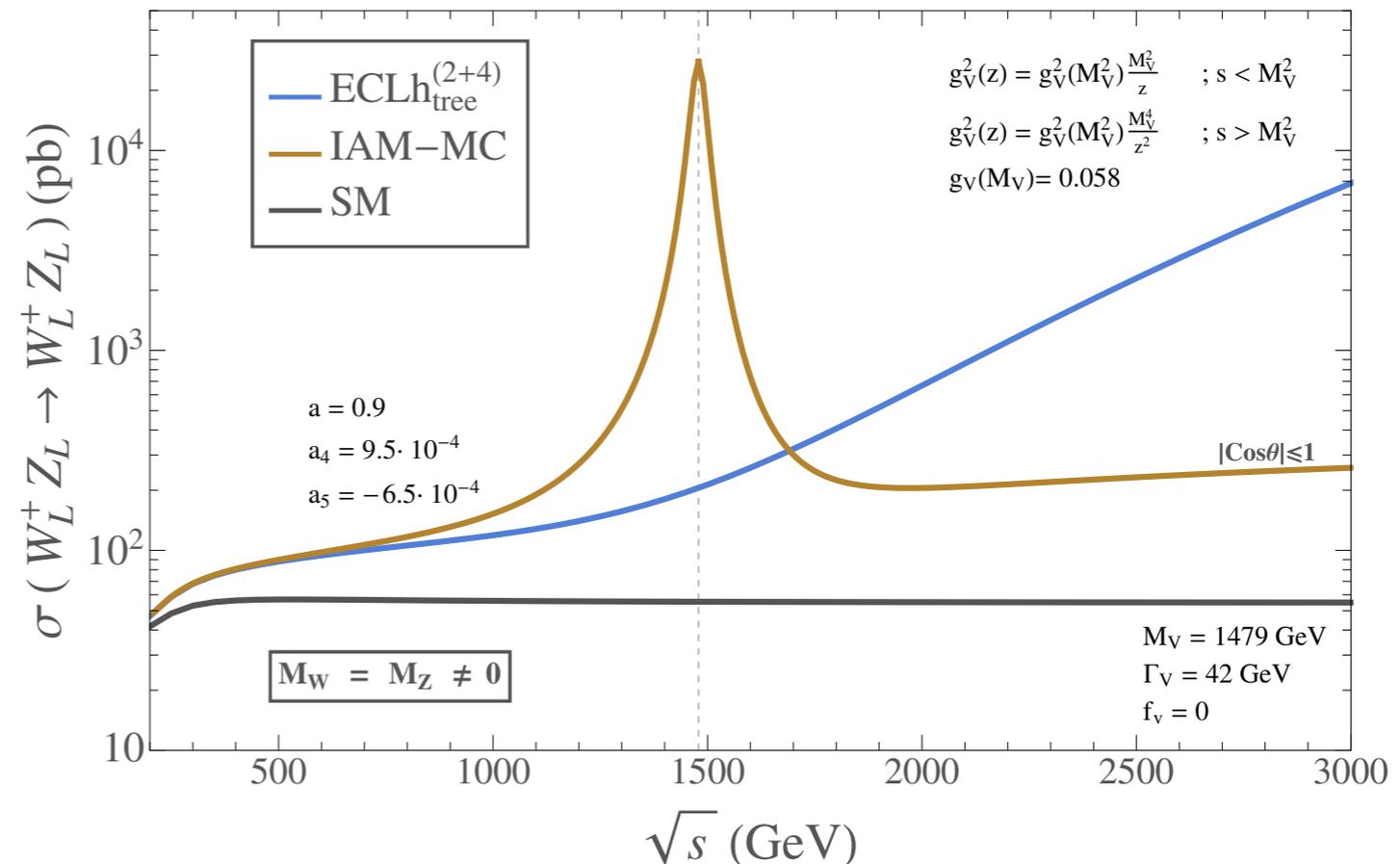
VBS kinematics very characteristic

extra jets key to select VBS efficiently

$|\eta_{W,Z}| < 2$

$2 < |\eta_j| < 5$
$\eta_{j_1} \cdot \eta_{j_2} < 0$
$m_{jj} > 500 \text{ GeV}$
$p_T^j > 20 \text{ GeV}$

leptonic final state

WZ final state

$\eta = 0$

W

$\eta = 2$

j

$\eta = 5$

p

$\eta = -\infty$

$\eta = \infty$

p

j

Z

$M_Z - 10 \text{ GeV} < m_{\ell_Z^+ \ell_Z^-} < M_Z + 10 \text{ GeV}$
$M_{WZ}^T \equiv M_{\ell\ell\nu}^T > 500 \text{ GeV}$
$\not{E}_T > 50 \text{ GeV}$
$p_T^\ell > 40 \text{ GeV}$

Transverse invariant mass

$M_{WZ}^T = \sqrt{\left(\sqrt{M^2(\ell\ell\ell) + p_T^2(\ell\ell\ell)} + |\vec{\not{p}_T}|\right)^2 - (\vec{p_T}(\ell\ell\ell) + \vec{\not{p}_T})^2}$