

Neutrino masses and leptogenesis from small lepton number violation



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DESY Hamburg

Moriond 2018

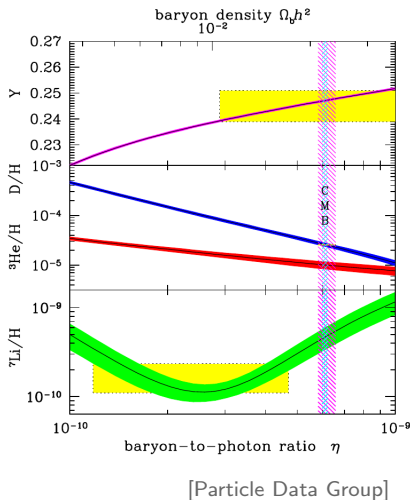
based on

[1507.06215](#), [1709.00414](#)

in collaboration with
A. Abada, G. Arcadi,
and M. Lucente

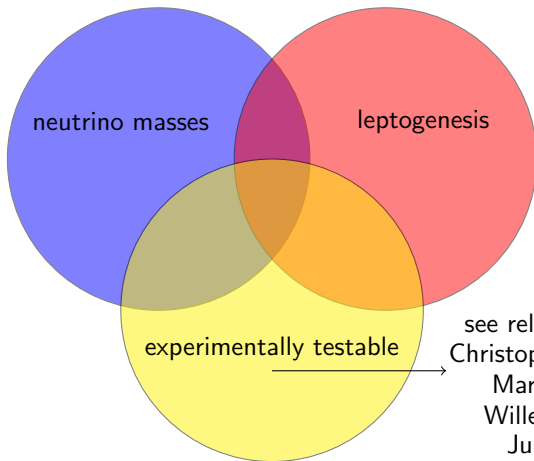


The observed baryon asymmetry of the Universe



- tiny matter-antimatter asymmetry in the early universe
- as particles and anti-particles annihilate, tiny asymmetry results in matter dominated Universe
- Sakharov conditions for dynamical generation [Sakharov '67]
 - violation of B (or L)
 - violation of C and CP
 - departure from thermal equilibrium

Motivation & Outline

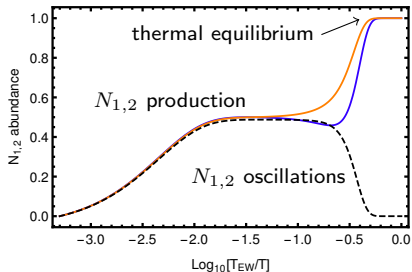


see related talks by:
Christopher Parkinson,
Marco Drewes,
Willem Verbeke,
Juraj Klarič,
Xabier Marciano,
Giulia Negro,
Jacobo Lopez Pávon

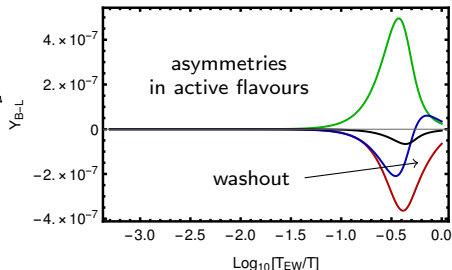
Leptogenesis through neutrino oscillations

neutrino mass spectrum

GeV $\equiv \equiv \equiv N_{1,2}$



eV $\equiv \equiv \equiv \nu_\alpha$



Leptogenesis through neutrino oscillations

[Akhmedov, Rubakov, Smirnov '98]

- nearly degenerate right-handed Majorana neutrinos N with $M_N \sim \text{MeV} - \text{GeV}$ and weak coupling to SM, $h_\nu \sim 10^{-7 \pm 2}$
- starting from a vanishing initial abundance, N s are thermally produced but do **not reach full thermal equilibrium** before the EW phase transition (= sphaleron freeze-out)
- N_R oscillations create a **CP violating** background for the active neutrinos

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- N_R oscillations create a **CP violating** background for the active neutrinos
- defining a generalized lepton number \mathcal{L} under which also the N s are charged, both the active and sterile sector develop a $\Delta\mathcal{L}$ of equal magnitude but opposite sign (for $M_N \gg T$)
- Sphaleron processes only pick up the $\Delta\mathcal{L} = \Delta L$ of the active sector, leading to a net baryon asymmetry

low-scale leptogenesis mechanism \rightarrow experimentally accessible !

Boltzmann equations

Kinetic equations for neutrino density matrix ρ_{AB} , $A, B \in \{N, \bar{N}, \nu, \bar{\nu}\}$

$$i \frac{d\rho}{dt} = [H, \rho] - \frac{i}{2} \{\Gamma, \rho\} + \frac{i}{2} \{\Gamma^p, I - \rho\} \quad [\text{Sigl, Raffelt '93}]$$

Assuming $M_N \gg m_\nu$: active-sterile oscillations suppressed,

kinetic equilibrium: $\rho_{NN}(T, k) = R_N(T) \rho_{NN}^{\text{eq}}(T, k)$,

thermal equilibrium for active ν : $\rho_{\nu, \bar{\nu}} \sim \rho_{\text{eq}}(T, k) \exp[\pm \text{diag}(\mu_{L\alpha})]$:

$$\begin{aligned} \frac{dR_N}{dt} = & -i [\langle H \rangle, R_N] - \frac{1}{2} \langle \gamma^{(0)} \rangle \{F^\dagger F, R_N - I\} - \frac{1}{2} \langle \gamma^{(1b)} \rangle \{F^\dagger \mu_L F, R_N\} \\ & + \langle \gamma^{(1a)} \rangle F^\dagger \mu_L F, \end{aligned}$$

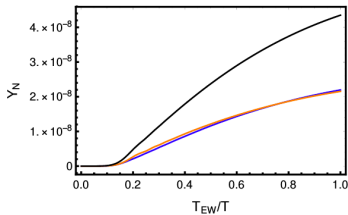
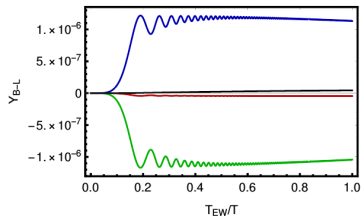
$$\begin{aligned} \frac{d\mu_{\Delta\alpha}}{dt} = & - \frac{9\zeta(3)}{2N_D \pi^2} \left\{ \langle \gamma^{(0)} \rangle (F R_N F^\dagger - F^* R_{\bar{N}} F^T) - 2 \langle \gamma^{(1a)} \rangle \mu_L F F^\dagger + \right. \\ & \left. + \langle \gamma^{(1b)} \rangle \mu_L (F R_N F^\dagger + F^* R_{\bar{N}} F^T) \right\}_{\alpha\alpha}, \quad [\text{Asaka, Shaposhnikov '05}] \end{aligned}$$

w. F = Yukawa coupling in mass basis, H = eff. Hamiltonian, $\langle \cdot \rangle$ = thermal average

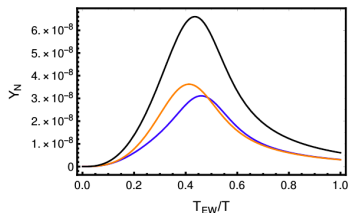
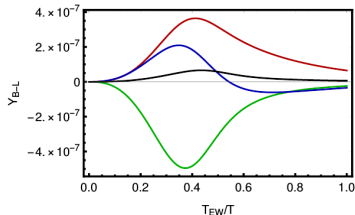
Benchmark solutions

[Abada, Arcadi, VD, Lucente '15, '17]

weak washout regime, $|F| \lesssim 10^{-7}$ (analytical solutions available)



intermediate/strong washout regime, $|F| \sim 10^{-6}$ (perturbative expansion)



Neutrino mass models for ARS leptogenesis

Inverse Seesaw [Wyler, Wolfenstein '83; Mohapatra, Valle '86]
(schematically for 'active', 'RH', 'sterile')

$$M_\nu = \underbrace{\begin{pmatrix} 0 & \frac{1}{\sqrt{2}}h_\nu v & 0 \\ \frac{1}{\sqrt{2}}h_\nu v & 0 & \Lambda \\ 0 & \Lambda & 0 \end{pmatrix}}_{\text{L conserving}} + \xi \underbrace{\begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \Lambda \end{pmatrix}}_{\Delta L=2}$$

$$\Rightarrow m_\nu \simeq \frac{\xi(h_\nu v)^2}{2\Lambda}, \quad \Delta m^2 = M_{N_2}^2 - M_{N_1}^2 \simeq 2\xi\Lambda^2$$

Small neutrino masses and small mass splitting generated by the same small lepton number violation (LNV) parameter

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Linear Seesaw [Barr '04; Malinsky, Romao, Valle '05]

$$M_\nu = \underbrace{\begin{pmatrix} 0 & \frac{1}{\sqrt{2}}h_\nu v & 0 \\ \frac{1}{\sqrt{2}}h_\nu v & 0 & \Lambda \\ 0 & \Lambda & 0 \end{pmatrix}}_{\text{L conserving}} + \epsilon \underbrace{\begin{pmatrix} 0 & 0 & \frac{1}{\sqrt{2}}h'_\nu v \\ 0 & 0 & 0 \\ \frac{1}{\sqrt{2}}h'_\nu v & 0 & 0 \end{pmatrix}}_{\Delta L=2}$$

$$\Rightarrow m_\nu \simeq \frac{\epsilon(h_\nu v)^2}{\Lambda}, \quad \Delta m^2 = M_{N_2}^2 - M_{N_1}^2 \simeq 2\epsilon(h_\nu v)^2$$

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Linear Seesaw [Barr '04; Malinsky, Romao, Valle '05] plus Inverse Seesaw

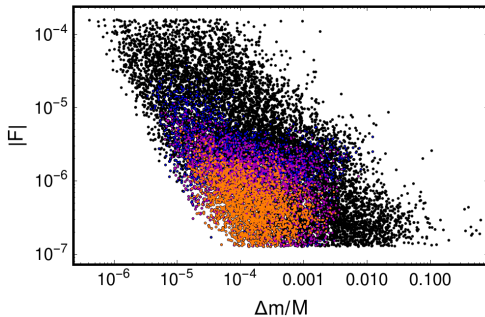
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[Abada, Arcadi, VD, Lucente '15,'17]

- Generate data points which respect all constraints on active neutrino masses/mixings as well as laboratory constraints on sterile neutrinos
⇒ minimal models: LISS(1,1), ISS(2,2), ISS(2,3) [Abada, Lucente '14]
- Calculate baryon asymmetry by solving Boltzmann equations
 - analytically in the weak washout regime
 - perturbative numerical procedure otherwise
- ⇒ requirements for successful leptogenesis:
constraints on masses, Yukawa couplings, mixing angles...
- ⇒ motivates parameter regions to focus experimental searches

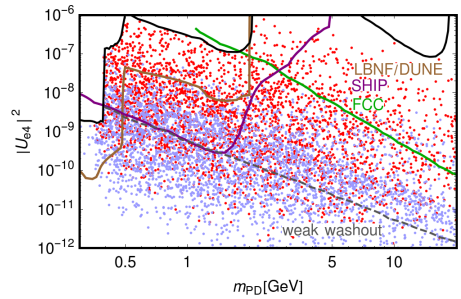
Neutrino mass models for ARS leptogenesis



here: LISS + strong washout

Yukawa coupling and mass splitting.
coloured: $\eta_b \geq \eta_b^{\text{obs}}$

mixing and mass scale
for **NH** and **IH**



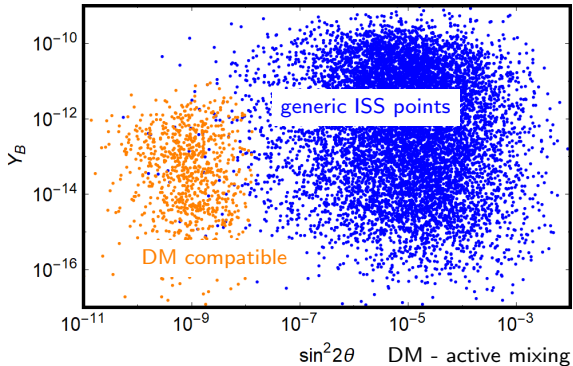
An example: dark matter in the ISS(2,3)

	=====	M_{PD}	—————>	decouple
GeV	=====	m_{PD}	—————>	leptogenesis
	=====	$\Delta m \sim \text{keV}$		
keV	—————	m_{DM}	—————>	dark matter (Dodelson-Widrow)
eV	—————	m_ν	—————>	ν oscillations

with

$$\frac{m_{\text{DM}}}{m_{\text{PD}}} \simeq \frac{\Delta m}{m_{\text{PD}}} \simeq 10^{-3} \left(\frac{10^{-5}}{|F|} \right) \left(\frac{m_{\text{PD}}}{\text{GeV}} \right)^{1/2} \left(\frac{m_\nu}{0.05 \text{ eV}} \right)^{1/2}$$

An example: dark matter in the ISS(2,3)



ISS + ARS leptogenesis too constraining on mixing angles

Conclusion

- observed matter-antimatter asymmetry requires BSM physics
- leptogenesis through neutrino oscillations provides a testable mechanism
 - oscillations of GeV right-handed neutrinos before EW phase transition
 - nonzero lepton number in the active sector
 - transmitted to baryons through sphaleron processes
- Generic predictions (mass range, sizeable active sterile mixing) but also sensitivity to underlying neutrino mass model
 - symmetry inspired vs minimal tuning
 - strong washout vs weak washout
 - minimal particle content vs minimal parameters
 - individual mixing elements experimentally accessible

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Thank you!