



Colored Dark Matter

based on: 1801.01135 with de Luca, Redi, Smirnov, Strumia



SCUOLA
NORMALE
SUPERIORE



The model

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \bar{Q} (i\not{D} - M_Q) Q$$

$$Q = (8, 1)_0$$

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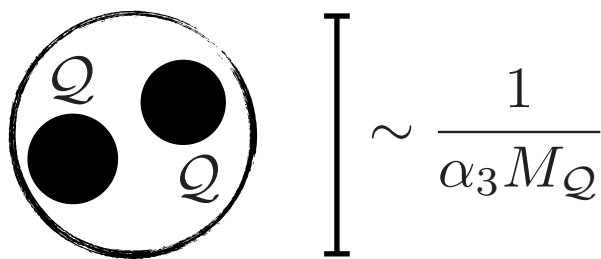
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Q -onlyum



$$E_B \sim \alpha_3^2 M_Q$$

DM candidate

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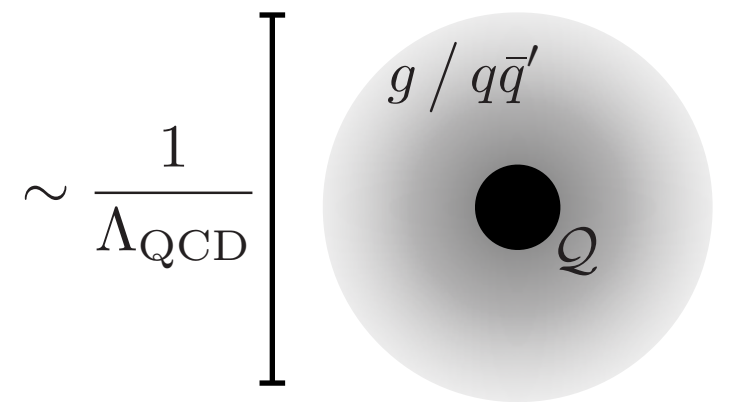
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Hybrids

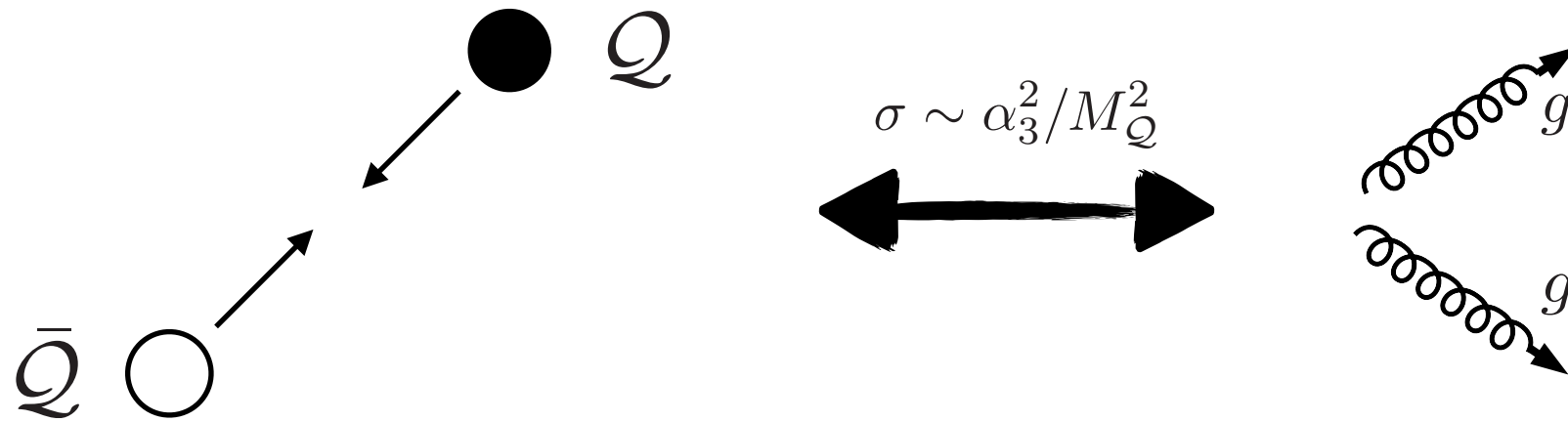


$$E_B \sim \Lambda_{\text{QCD}} \quad \& \quad \sigma \sim \Lambda_{\text{QCD}}^{-2}$$

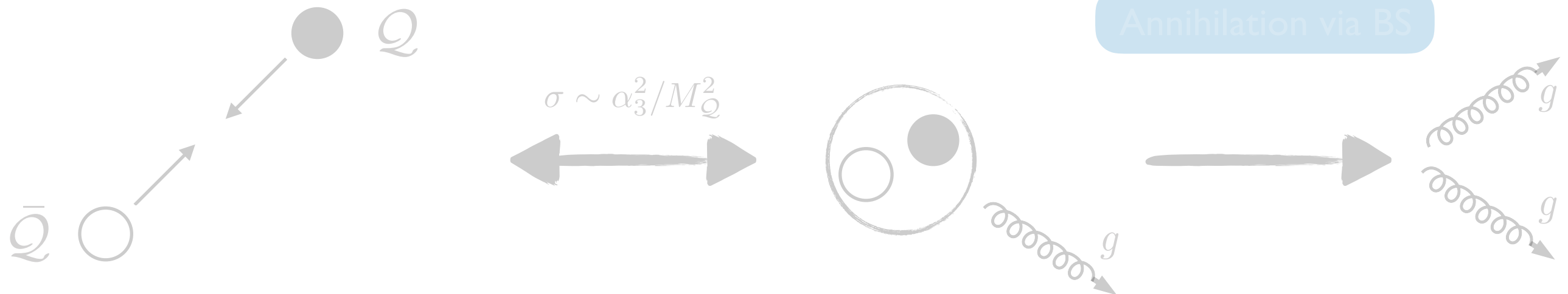
Dangerous

(Pre-confinement) Cosmological evolution

Direct annihilation

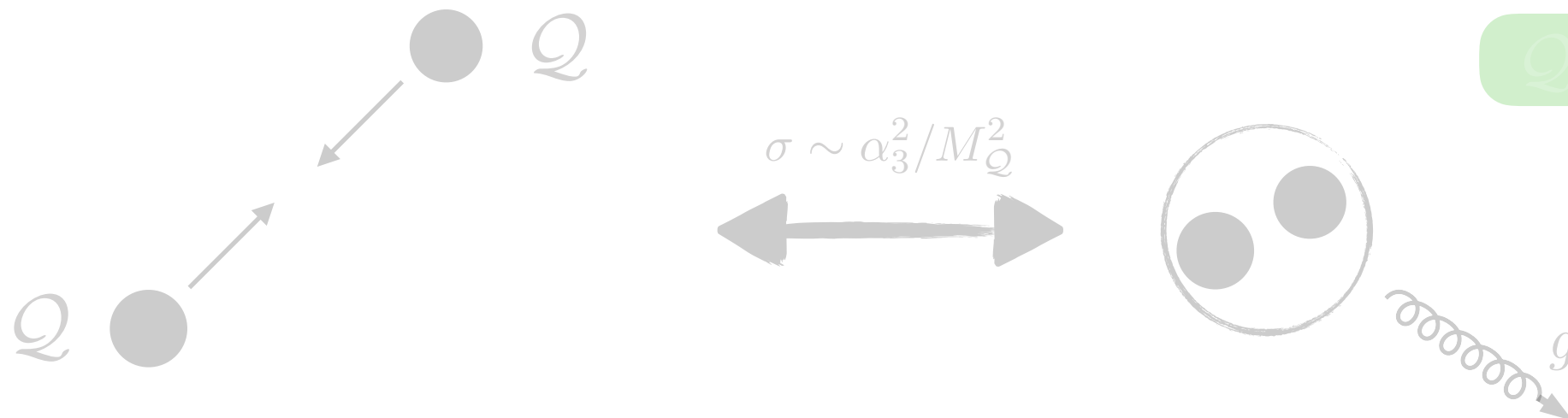


Annihilation via BS



AM, Redi, Smirnov, Strumia '17

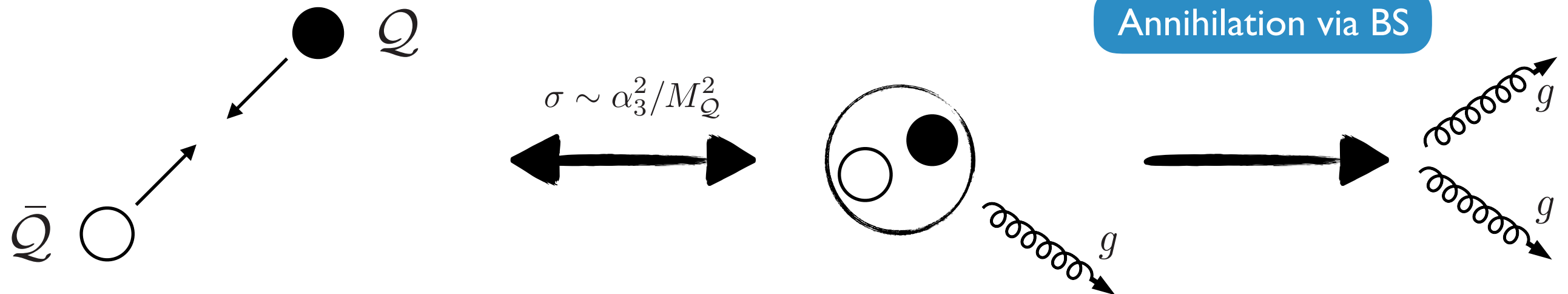
Q -onlyum production



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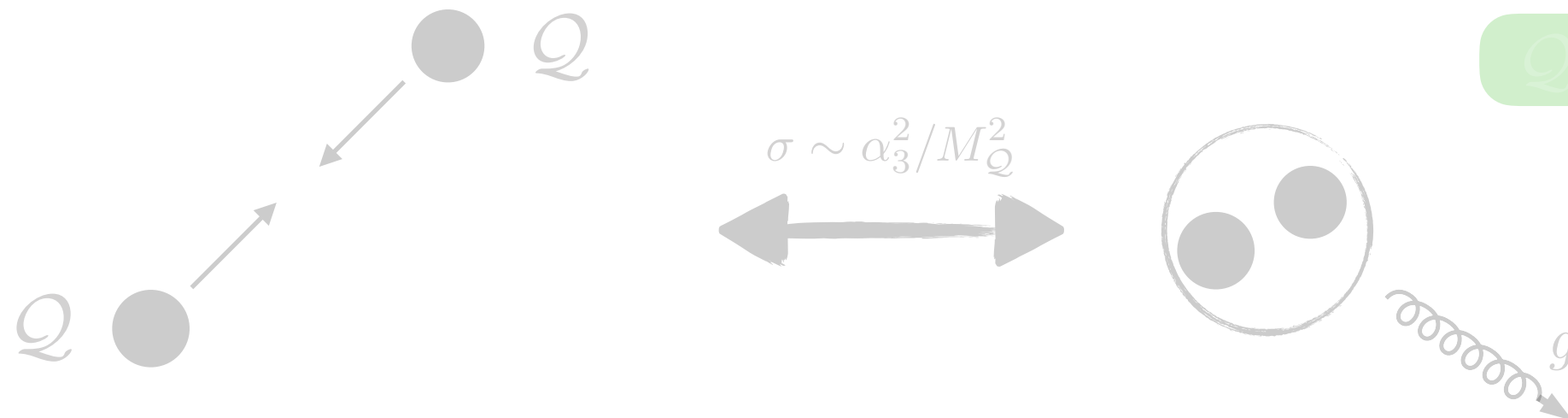


Direct annihilation



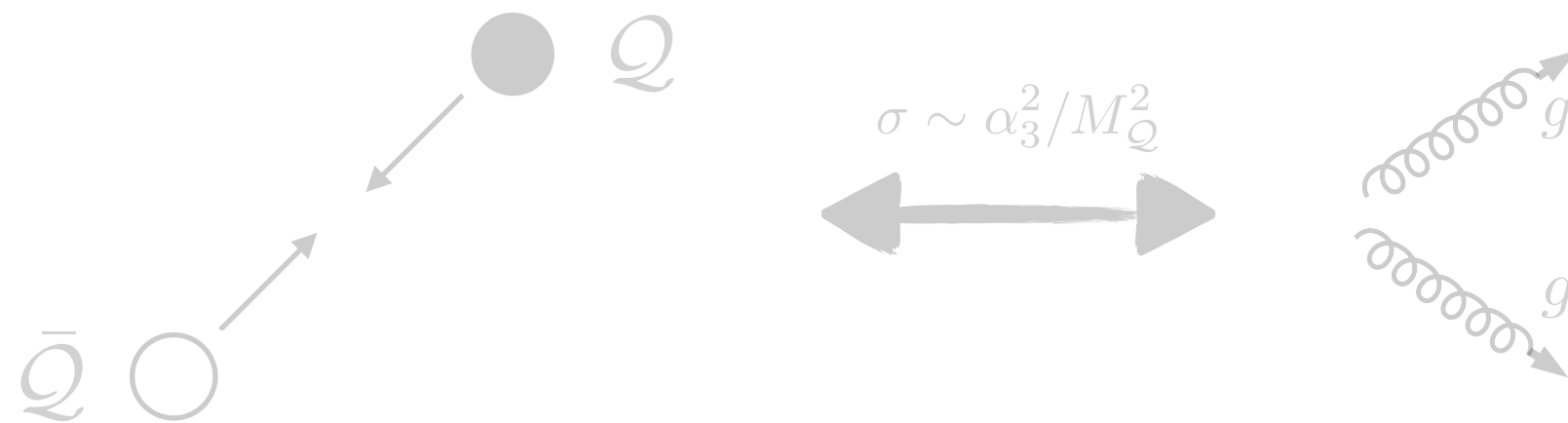
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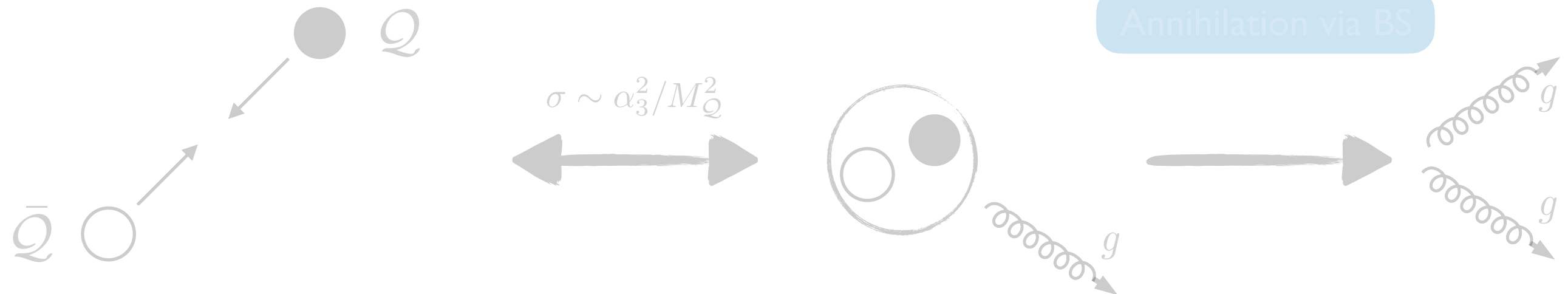


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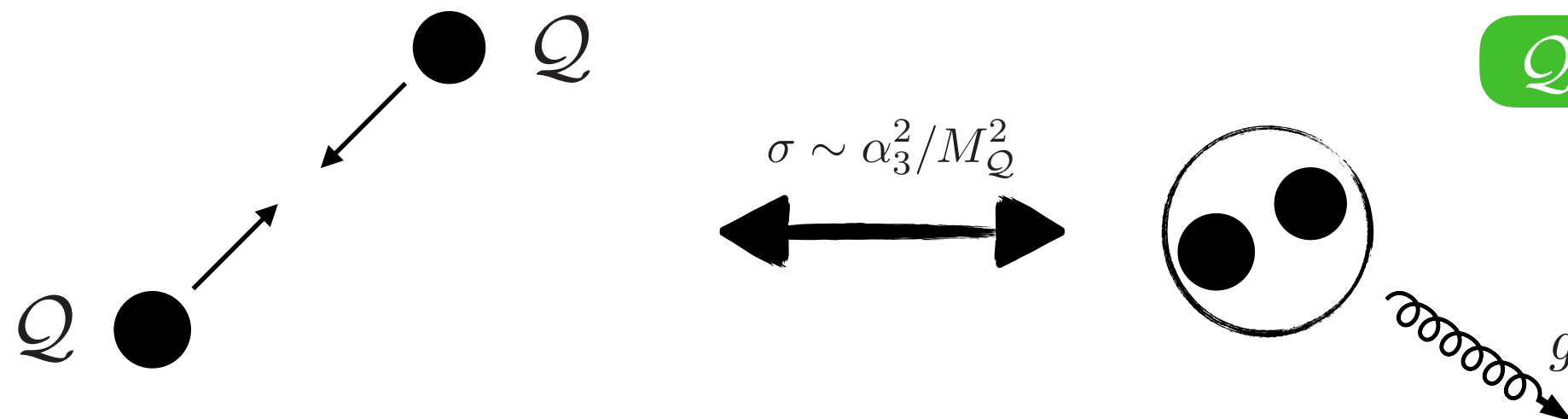
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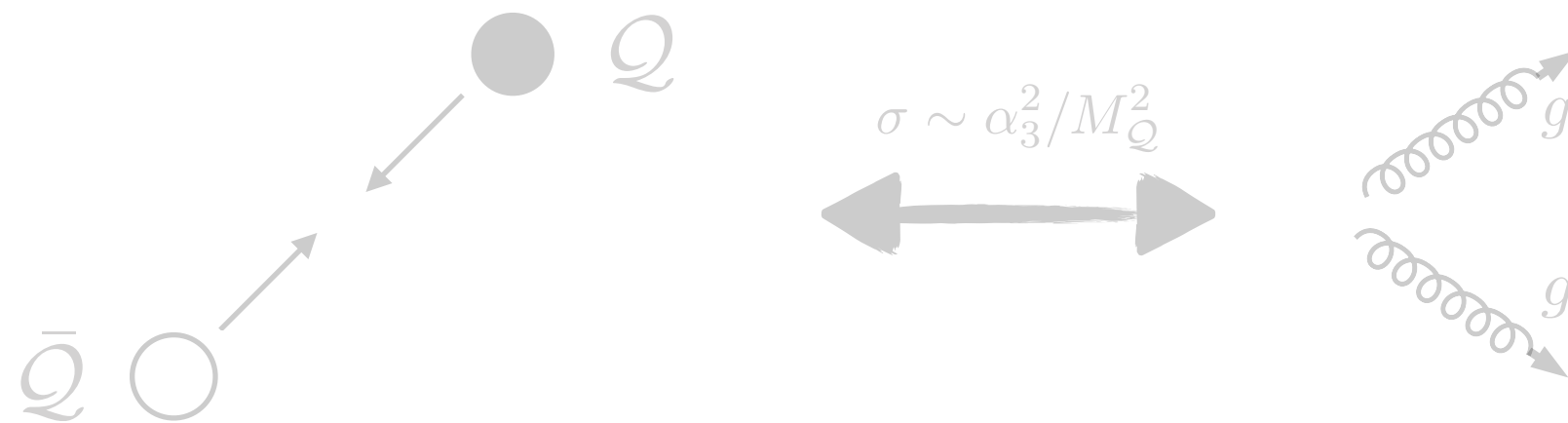
Annihilation via BS



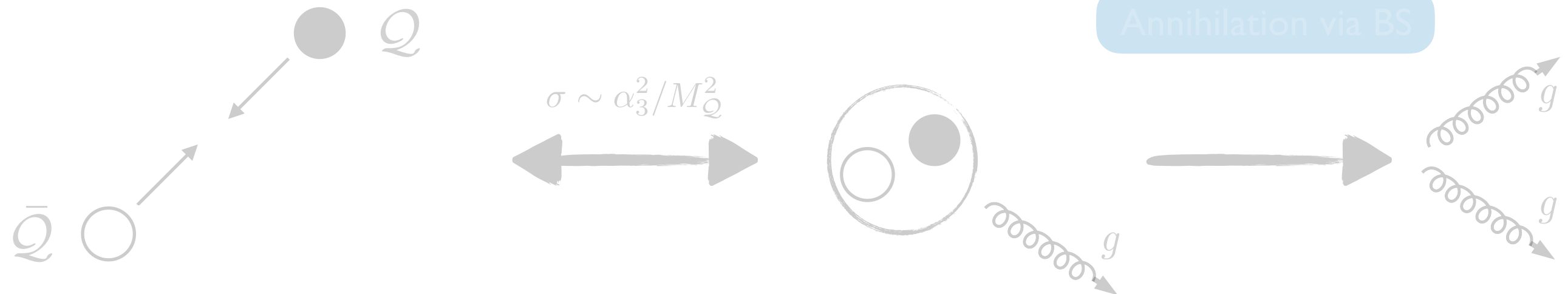
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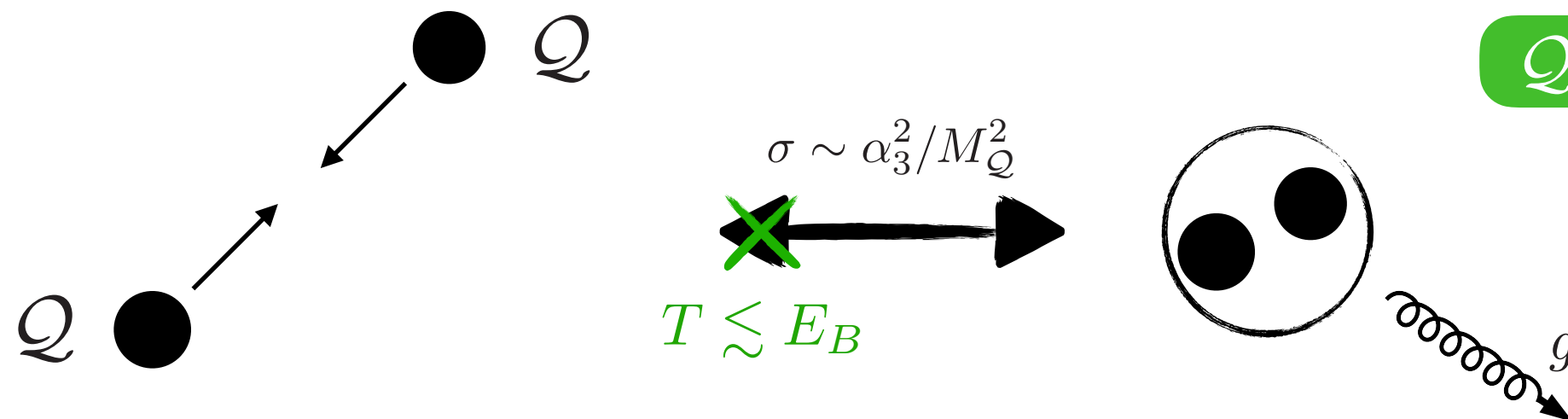
(Pre-confinement) Cosmological evolution



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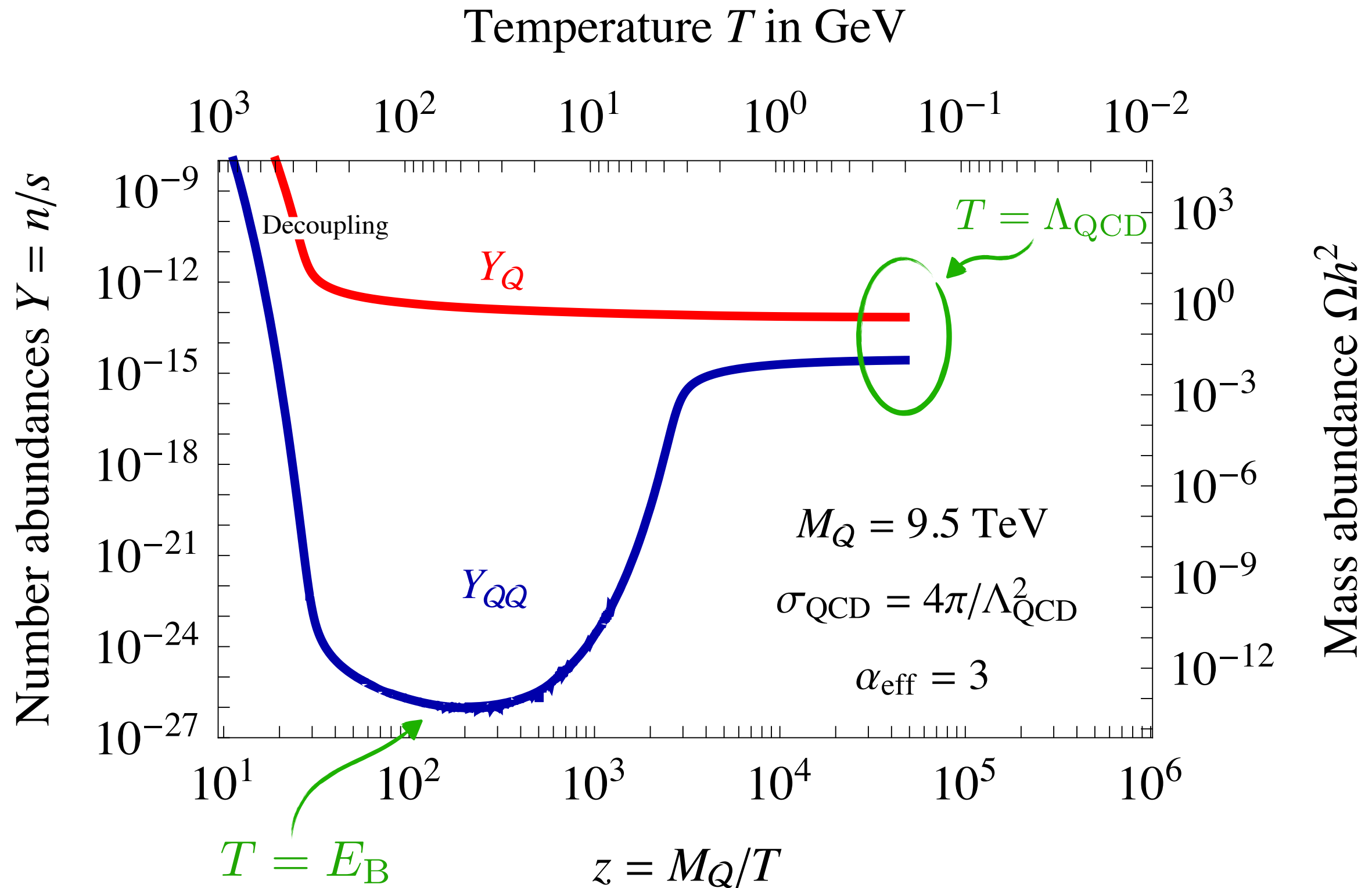
Annihilation via BS



Q -onlyum production

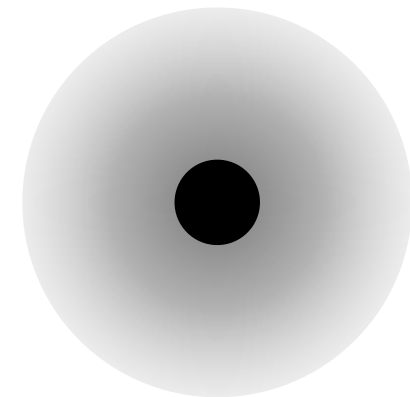
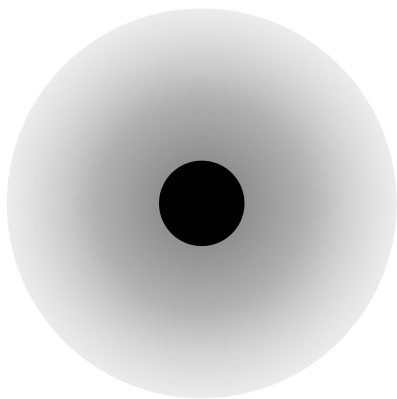
AM, Redi, Smirnov, Strumia '17

(Pre-confinement) Cosmological evolution



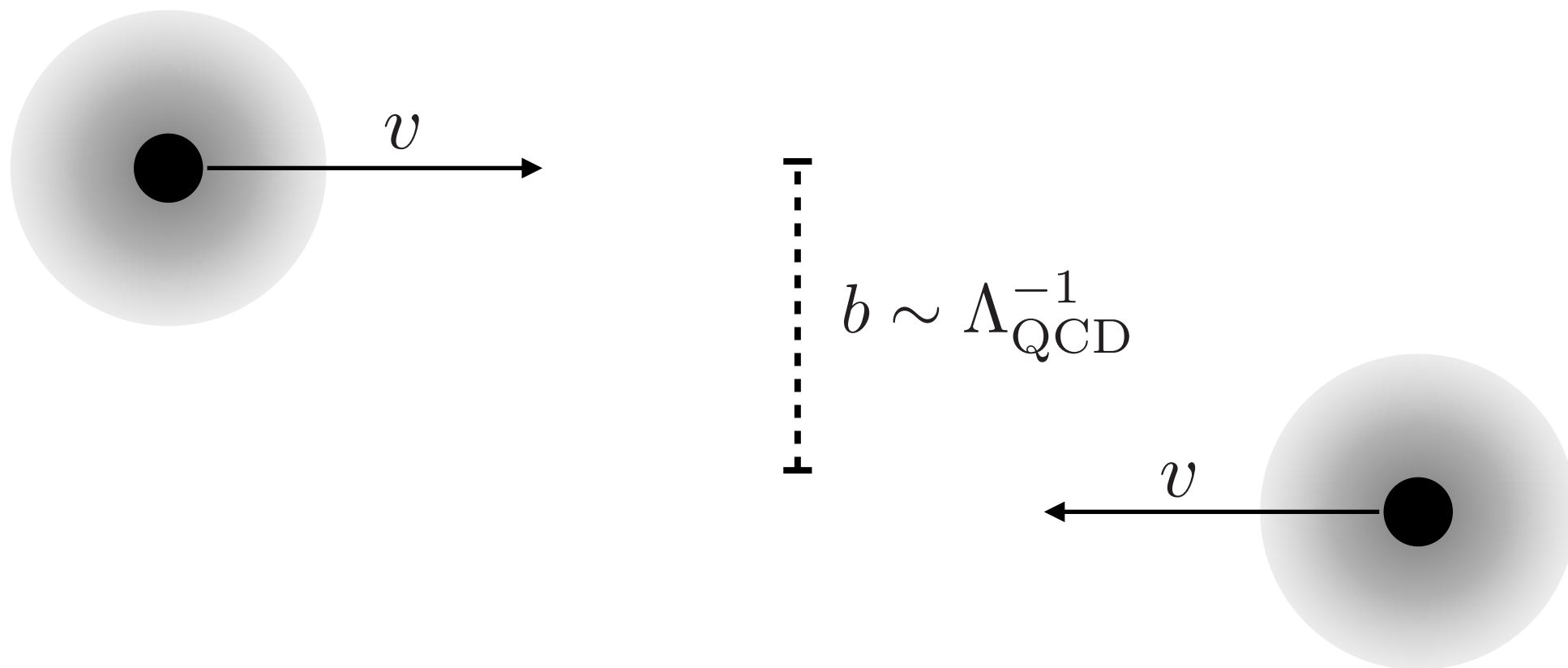
(Post-confinement) **Cosmological evolution**

@ confinement all free Q hadronize into hybrids



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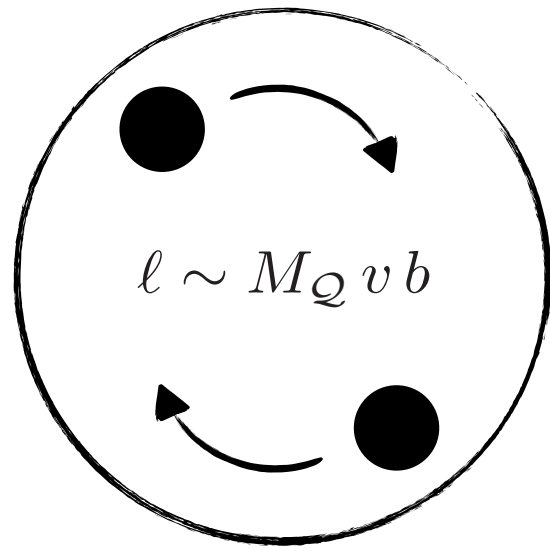


hybrids can collide with large impact parameters, giving rise to self scattering cross-sections of typical QCD size:

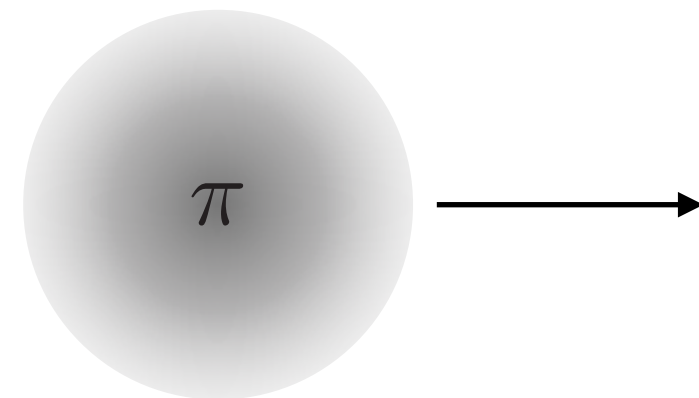
$$\sigma \sim \Lambda_{\text{QCD}}^{-2} \gg \alpha_3^2 / M_Q^2$$

(Post-confinement) **Cosmological evolution**

in hybrids collisions \mathcal{Q} -onlyum hadrons or unstable $\mathcal{Q}\bar{\mathcal{Q}}$ can be formed



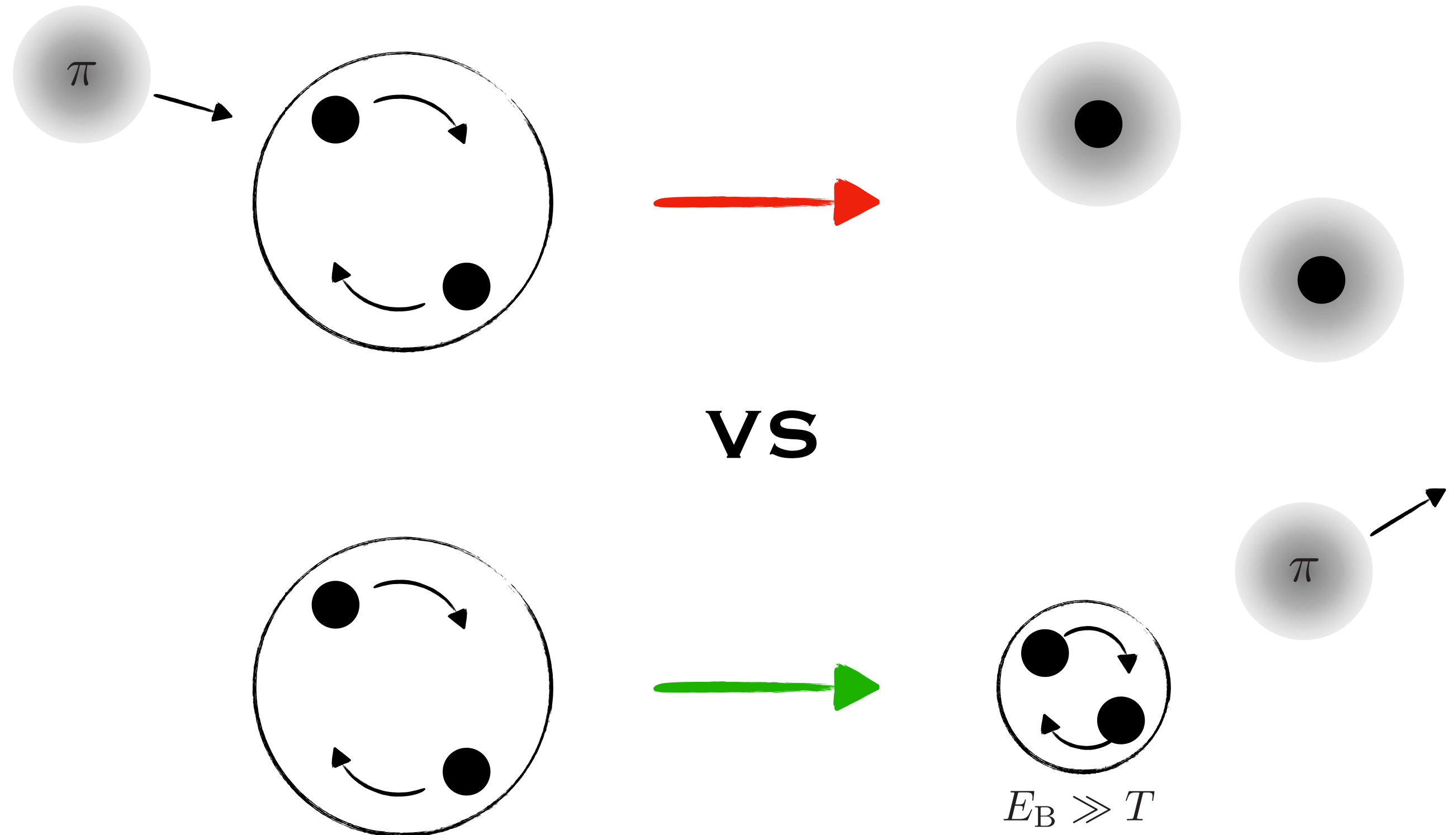
$$E_{\text{B}}^* \ll E_{\text{B},0} \sim \alpha_3^2 M_{\mathcal{Q}}$$



\mathcal{Q} -onlyum are typically formed in states with large angular momenta

(Post-confinement) Cosmological evolution

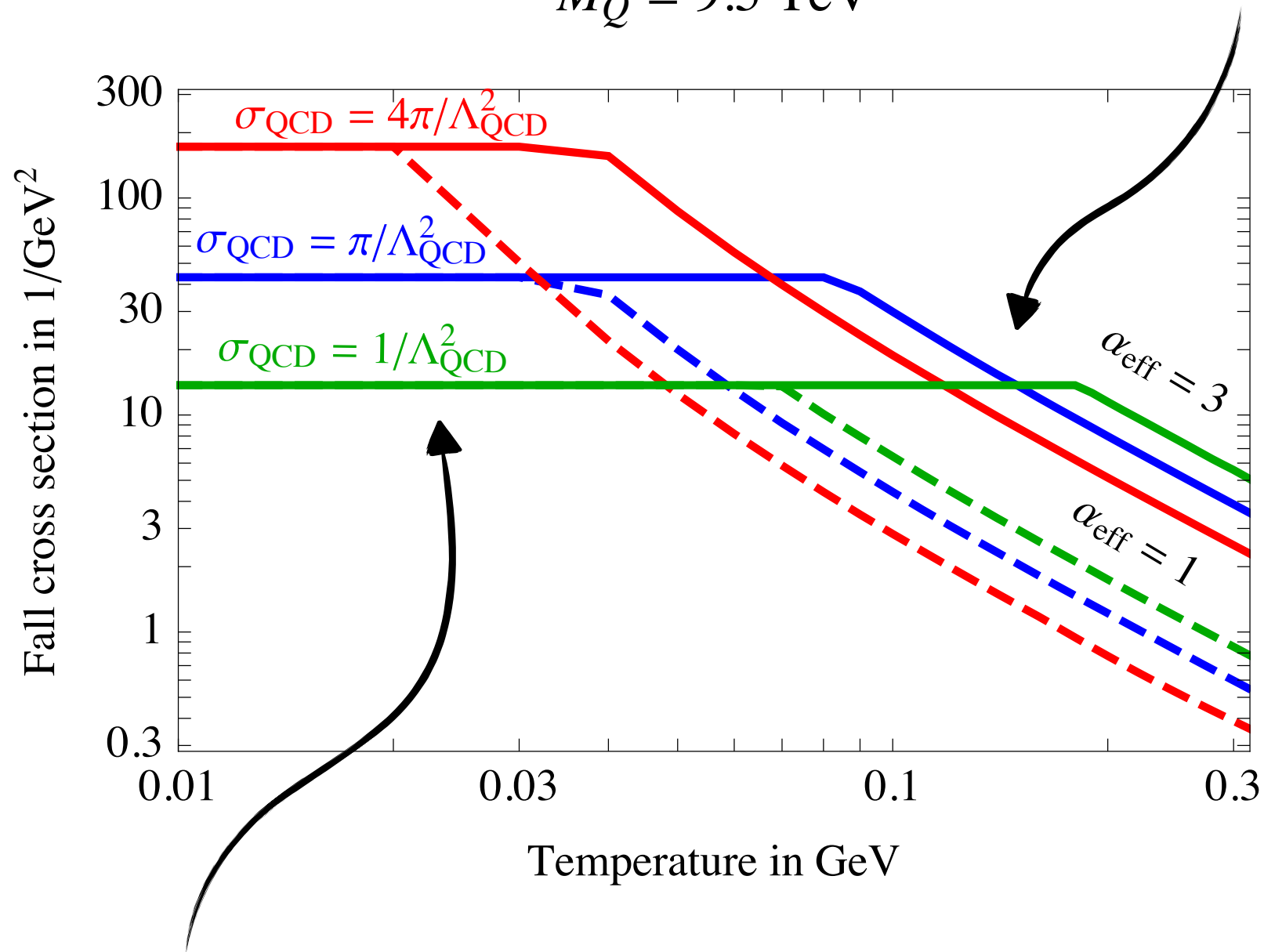
after formation $Q\bar{Q}$ states can **break** or **fall** to an unbreakable (deep enough) level



(Post-confinement) Cosmological evolution

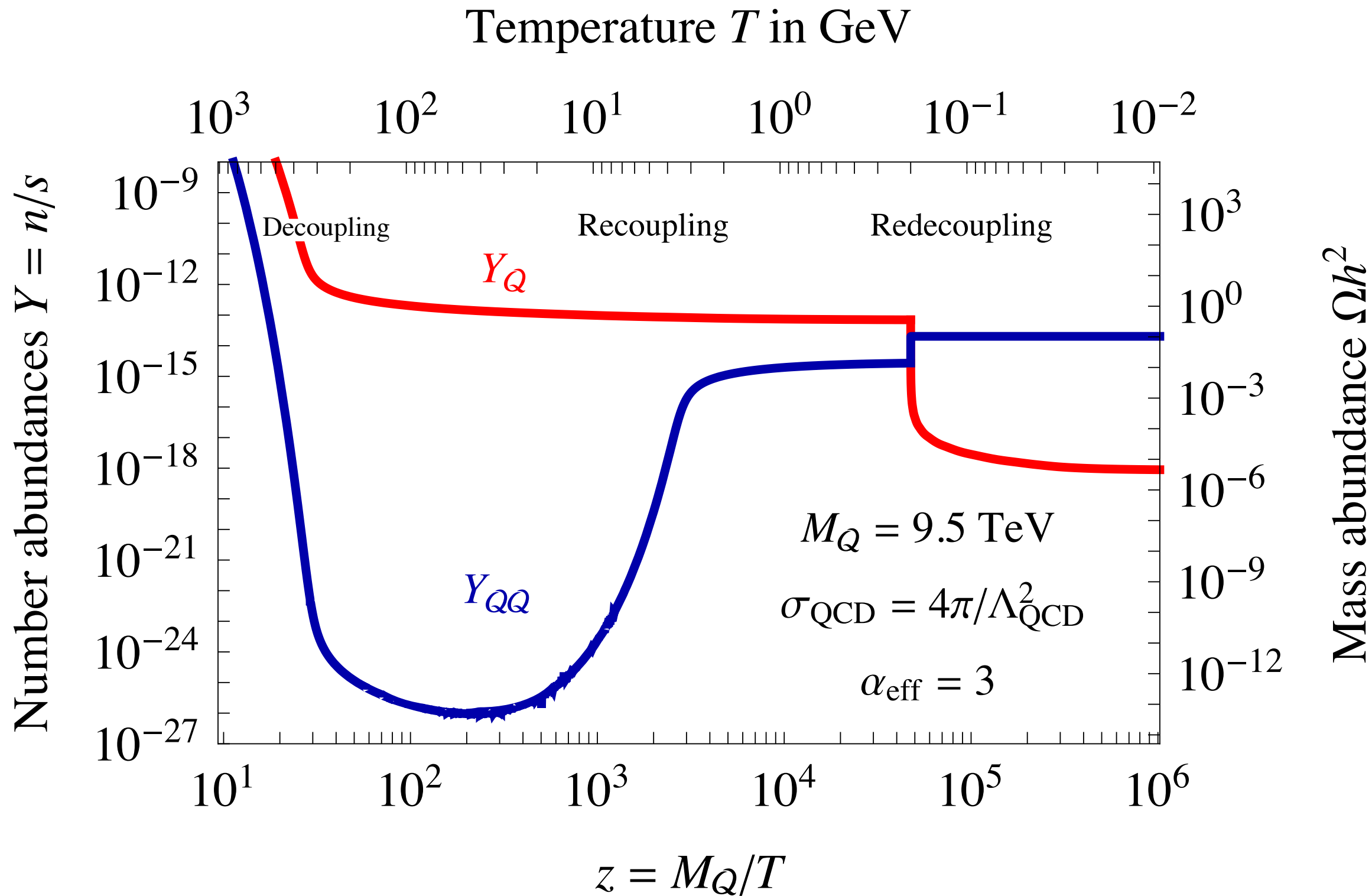
for $T \gg m_\pi$ the effective cross section to form an unbreakable Q -onlyum is reduced compared to the geometric one

$$M_Q = 9.5 \text{ TeV}$$



at low temperatures all the Q -onlyum manage to reach an unbreakable state

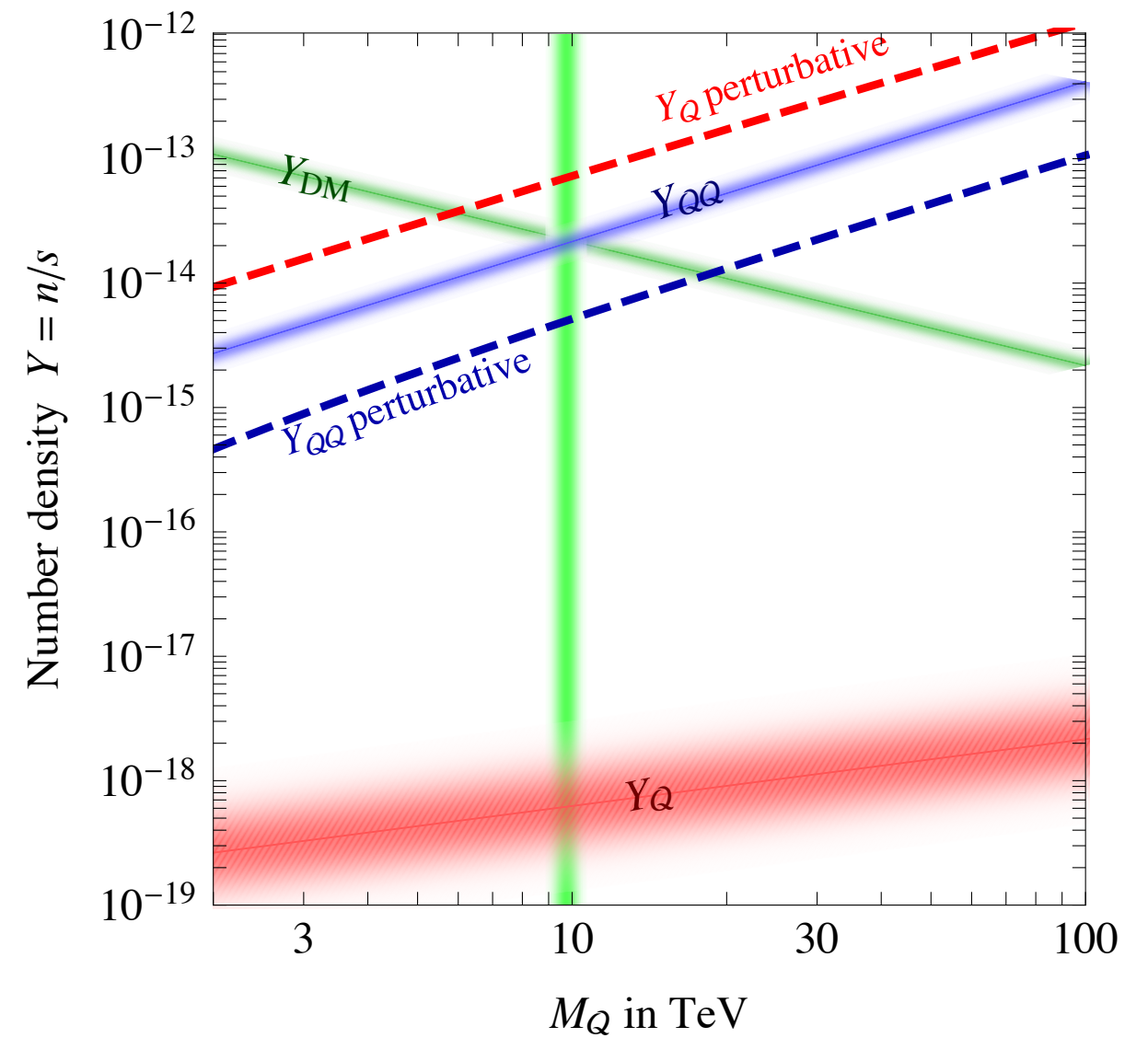
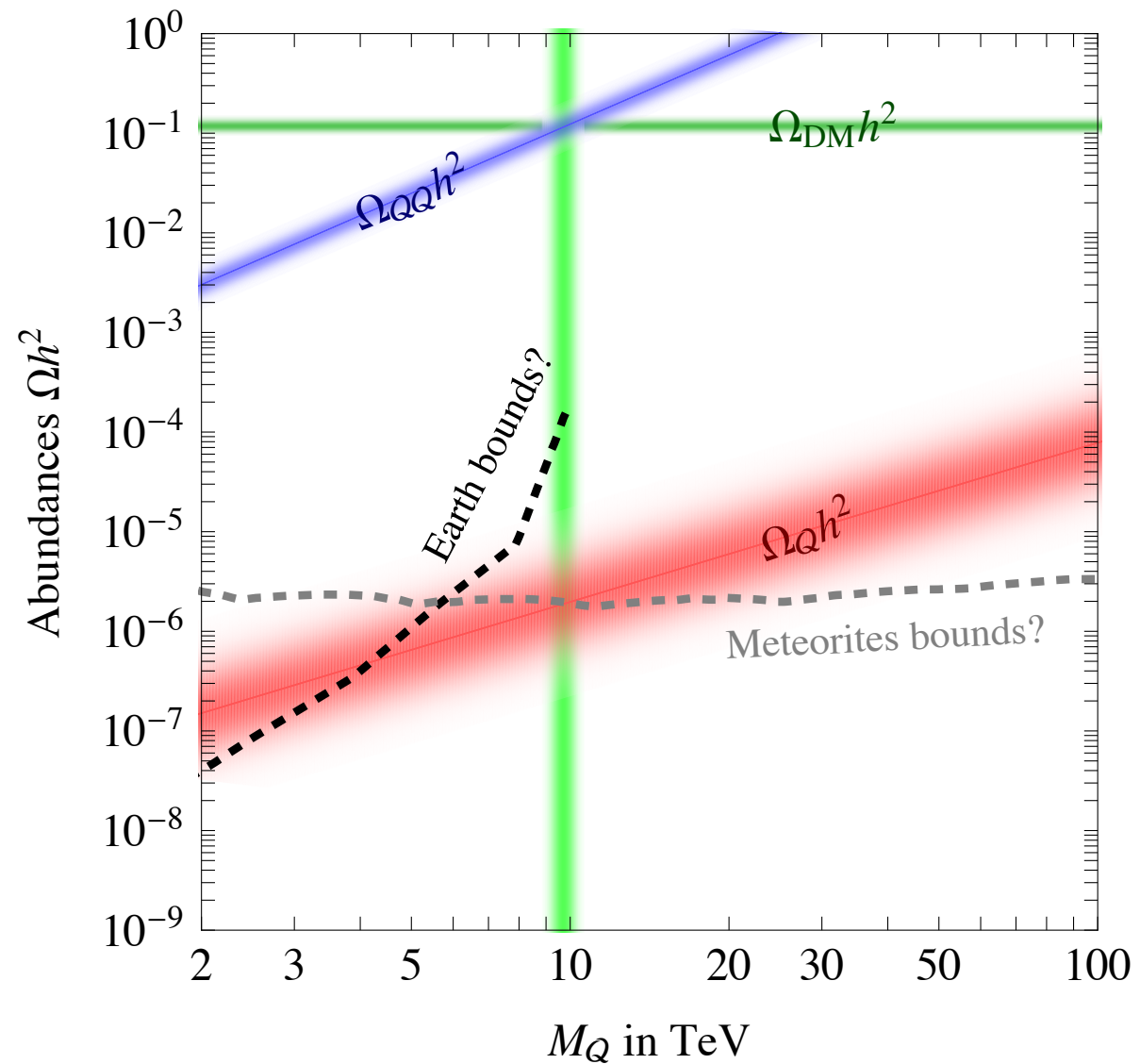
(Post-confinement) Cosmological evolution



Relic abundances

$$M_Q \simeq 9.5 \text{ TeV}$$

$$\Omega_{\text{hyb}} \sim 10^{-5} \Omega_{\text{DM}}$$



Dark Matter signals

DM direct detection

QQ interacts with **gluons** through **induced** chromo-dipole moments analogous to Rayleigh scattering hydrogen/light

$$\mathcal{L}_{\text{eff}} = c_E M_{\text{DM}} \bar{B} B \vec{E}^a{}^2$$

the polarizability coefficient is:

$$c_E = \pi \alpha_3 \langle B | \vec{r} \frac{1}{H_8 - E_{10}} \vec{r} | B \rangle \simeq 1.5 \pi a^3$$

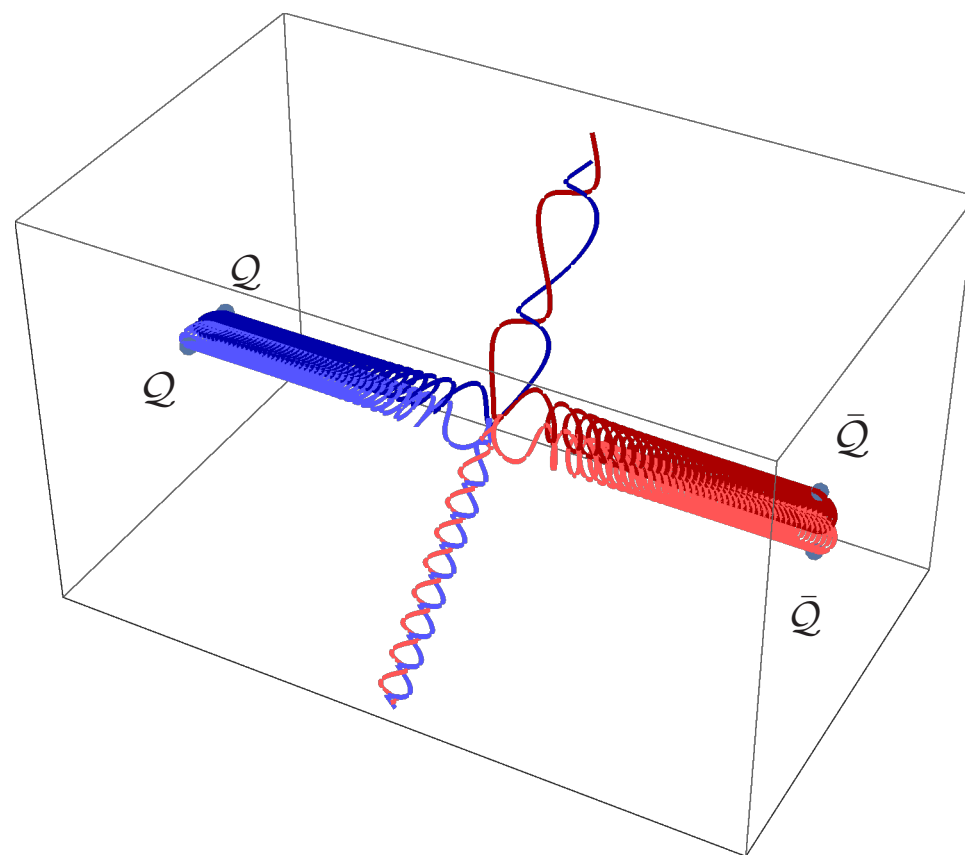
this gives a SI cross section slightly below XENON1T bounds

$$\sigma_{\text{SI}} = 2.3 \cdot 10^{-45} \text{ cm}^2 \times \left(\frac{20 \text{ TeV}}{M_{\text{DM}}} \right)^6 \left(\frac{0.1}{\alpha_3} \right)^8 \left(\frac{c_E}{1.5 \pi a^3} \right)^2$$

DM indirect detection

DM annihilation is dominated by **recombination** followed by a **decay** into SM particles

$$(Q Q) + (\bar{Q} \bar{Q}) \xrightarrow{\text{blue}} (Q \bar{Q}) + (Q \bar{Q}) \xrightarrow{\text{red}} \text{SM}$$

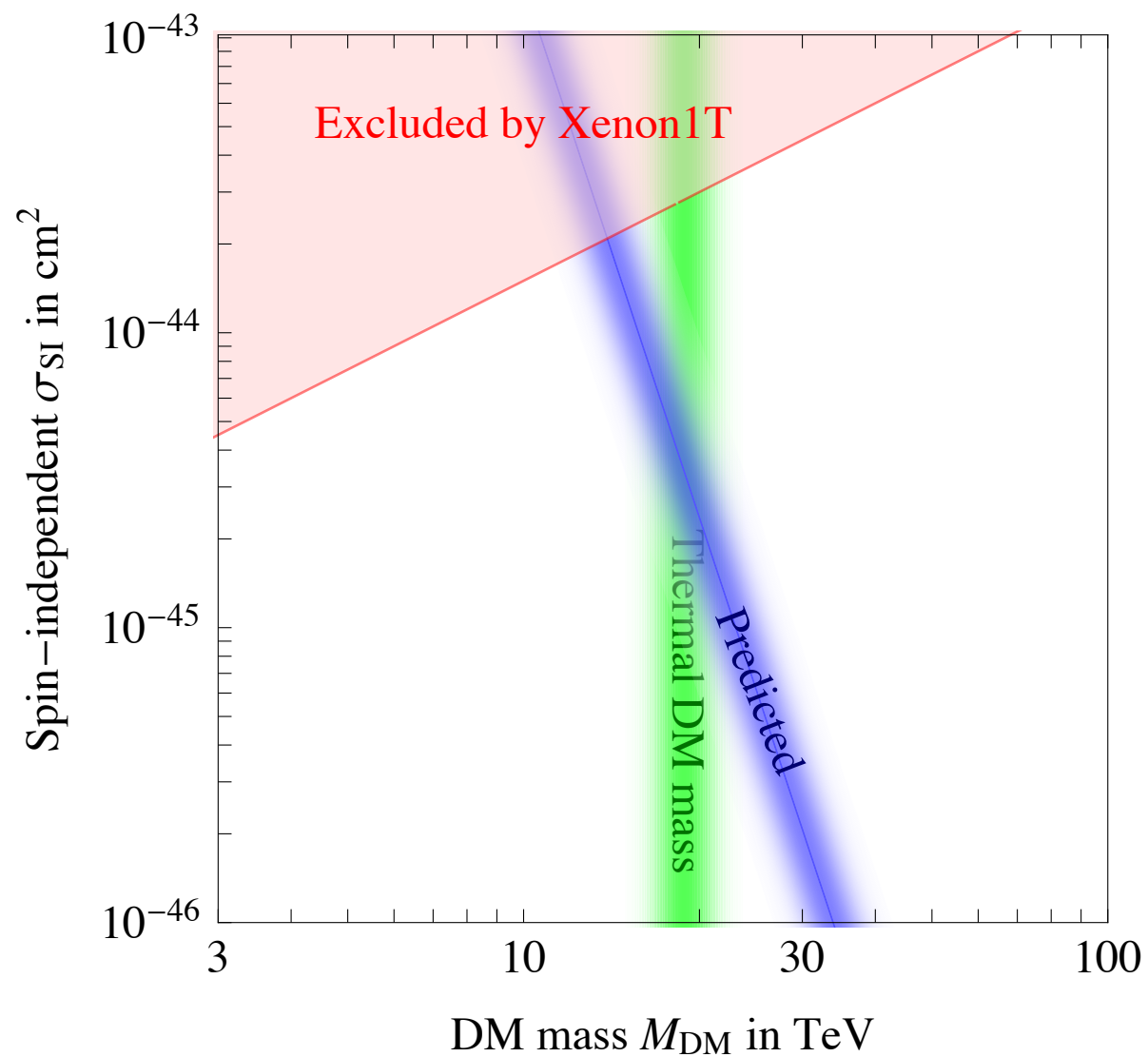


the recombination cross section is $\sim \alpha_3^3$ bigger than the one for direct annihilation

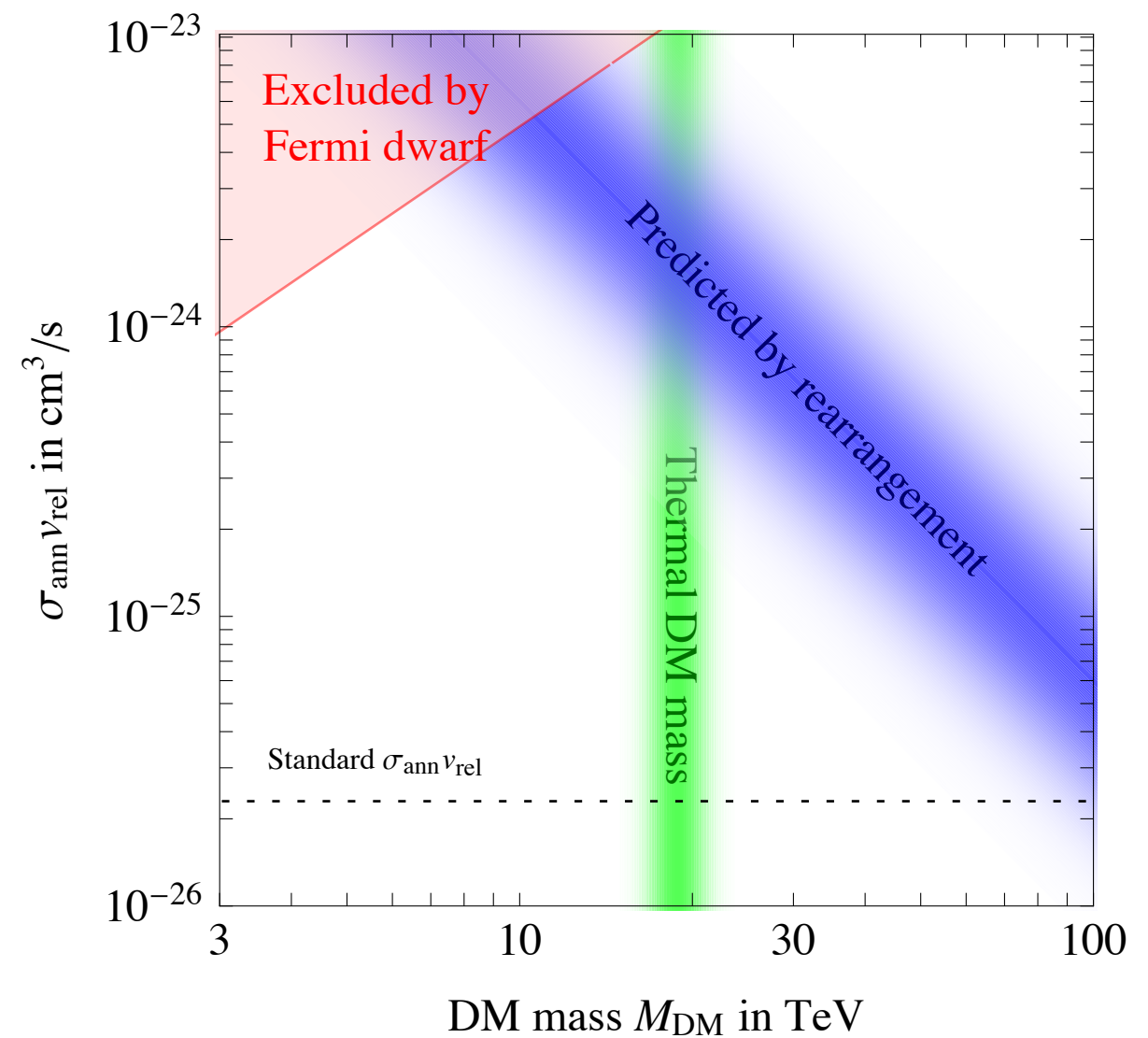
$$\sigma_{\text{ann}} v_{\text{rel}} \sim \frac{\pi a^2 v_{\text{rel}} / 2}{\sqrt{E_{\text{kin}} / E_B}} = \frac{\sqrt{2} \pi}{3 M_Q^2 \alpha_3} = 1.5 \cdot 10^{-24} \frac{\text{cm}^3}{\text{sec}}$$

DM direct & indirect detection bounds

Direct detection



Indirect detection



Collider detection of \mathcal{Q}

\mathcal{Q} can be pair produced through QCD interactions

$$p + p \rightarrow \mathcal{Q} + \bar{\mathcal{Q}}$$

and then hadronize into stable or long lived hadrons which give rise to (**charged**) tracks

$$\mathcal{Q} \rightarrow \mathcal{Q}g \qquad \text{or} \qquad \mathcal{Q} \rightarrow \mathcal{Q}q\bar{q}'$$

$\sqrt{s} \sim 65 \text{ TeV}$ is needed to discover \mathcal{Q} with $M_{\mathcal{Q}} \sim 9.5\text{TeV}$

$\sqrt{s} \sim 100 \text{ TeV}$ would be sensitive up to $M_{\mathcal{Q}} \lesssim 15\text{TeV}$

Hybrids signals/bounds

Direct detection of hybrids

hybrids in the galactic halo hit the Earth with energies $E_0 = M_{\mathcal{Q}} v^2 / 2 \sim \text{MeV}$
but then they lose energy due to interaction with matter

Starkman et al. 90'

$$E = E_0 \exp(-x/x_0) \qquad x_0 = \frac{\rho N_A \sigma_A}{M_{\mathcal{Q}}} \sim 100 \text{ m} \ll \text{atmosphere thickness}$$

hybrids reach underground detectors with energies below typical thresholds $\mathcal{O}(\text{keV})$

Mack, Beacom, Bertone 07'

hybrids are still excluded by balloon searches or Earth over-heating if $\Omega_{\text{hyb}} = \Omega_{\text{DM}}$
but not if $\Omega_{\text{hyb}} = 10^{-5} \Omega_{\text{DM}}$

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Searches for heavy nuclei in terrestrial samples

Hemmick et al. 89'

Norman, Gazes, Bennett 86'

$$\frac{N_{\text{hyb}}}{N} < \begin{cases} 10^{-14} & \text{Oxygen} \\ 10^{-16} & \text{Enriched C} \end{cases}$$

CAUTION

not clear whether or not hybrids bind to nuclei

for iso-spin triplet $Qq\bar{q}'$ pions can mediate long ranges forces.
So we expect they do

for iso-spin singlet Qg pions cannot mediate long ranges forces.
They do not bind? They bind only to big nuclei?

the following bounds apply only if they do bind

$$N_n \left| \frac{M_Q}{M_{\text{Earth}}} \right|$$

Searches for heavy nuclei in terrestrial samples

Hemmick et al. 89'

$$\frac{N_{\text{hyb}}}{N_n} < \begin{cases} 10^{-14} & \text{Oxygen} \\ 10^{-16} & \text{Enriched C} \\ 10^{-12} & \text{Iron} \end{cases}$$

Norman, Gazes, Bennett 86'

hybrids get captured and thermalized in the upper atmosphere

$$M \sim \rho_{\text{hyb}} v_{\text{rel}} \times \pi R_E^2 \times \Delta t \sim 2.5 \cdot 10^{10} \text{ kg}$$

then collisions in the Earth atmosphere could make hybrid N, O, He kept in the crust by electromagnetic binding

$$\left. \frac{N_{\text{hyb}}}{N_n} \right|_{\text{Earth}} = \frac{M}{M_Q} \frac{m_N}{M_{\text{Earth}}} \approx 4 \cdot 10^{-19}$$

Searches for heavy nuclei in meteorites

limits by meteorites samples give

$$\frac{N_{\text{hyb}}}{N_n} < 4 \times 10^{-14}$$

Polikanov et al. 91'

meteorites capture hybrids only if they bind to nuclei

$$\left. \frac{N_{\text{SIMP}}}{N_n} \right|_{\text{meteorite}} = n_{\text{SIMP}} \sigma_{\text{capture}} v_{\text{rel}} \Delta t \approx 10^{-14} \frac{\sigma_{\text{capture}}}{0.01/\Lambda_{\text{QCD}}^2}$$

estimated scaling typical neutron-capture cross sections

Conclusions

DM can be made of two color octet with mass around 9.5 TeV

Direct detection signals within the reach of future experiments

Indirect detection signals enacted by recombination

Manifest as stable tracks at collider, 65 TeV needed

Look for heavy isotopes

Backup slides

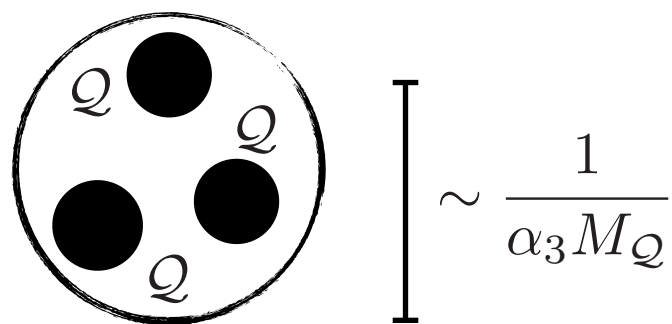
Why adjoint representation?

We could have chosen

$$Q = (3, 1)_0$$

in this case the new hadron would have been

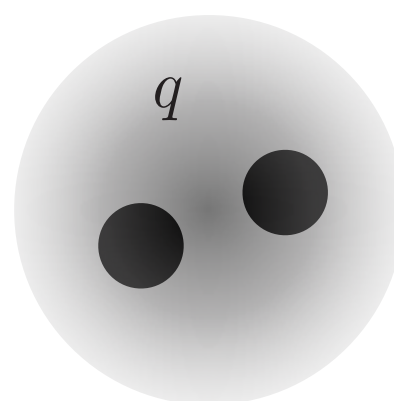
Q -onlyum



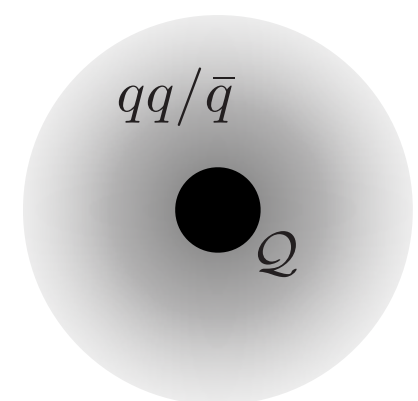
$$E_B \sim \alpha_3^2 M_Q$$

DM candidate

Hybrids



$$\sim \frac{1}{\Lambda_{\text{QCD}}}$$



$$E_B \sim \Lambda_{\text{QCD}} \quad \& \quad \sigma \sim \Lambda_{\text{QCD}}^{-2}$$

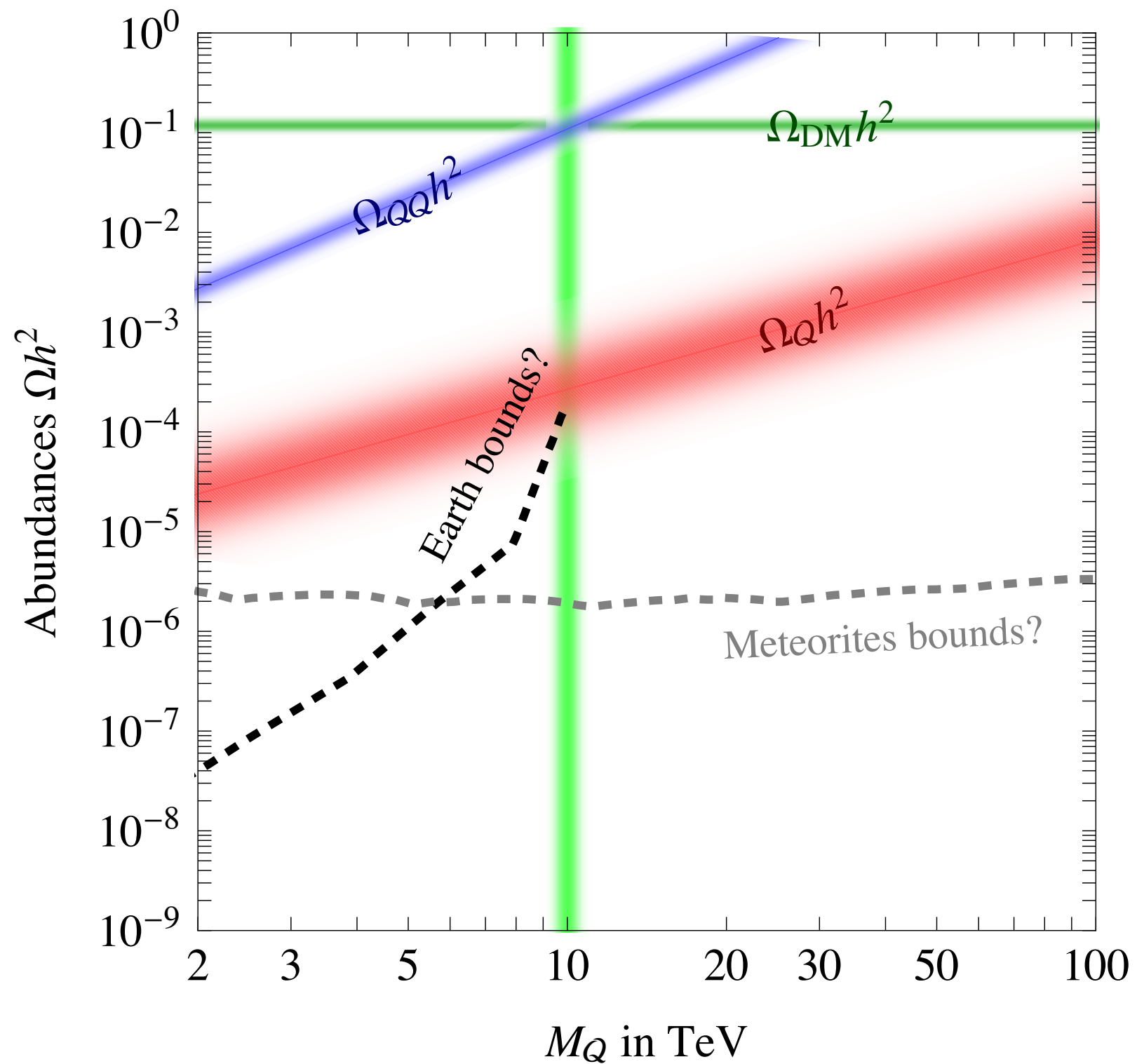


fractional charge



Too Dangerous

The worst case scenario



Smooth re-coupling

