

Dark matter and long-lived particles at the LHC

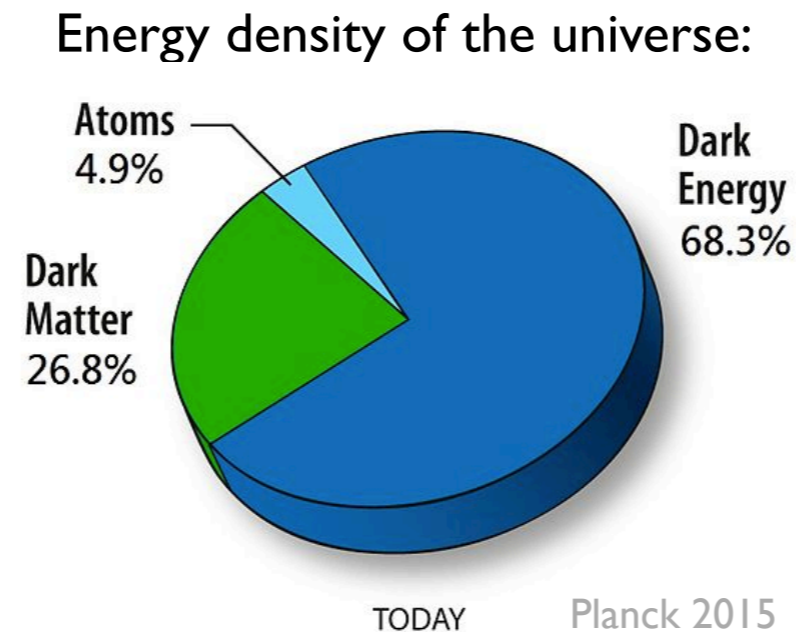
Jan Heisig (RWTH Aachen University)



53rd Rencontres de Moriond session on
Electroweak Interactions and Unified Theories
La Thuile, Aosta Valley, Italy, March 10–17, 2018

Among key scientific goals of LHC:

- Pinpoint the nature of dark matter!

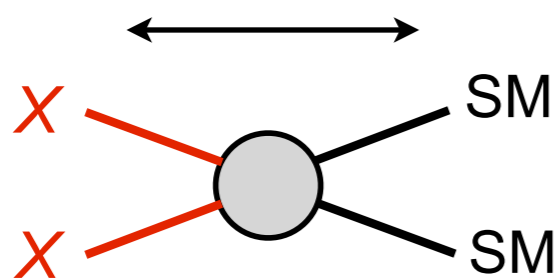


Needed: Predictions for possible signatures of dark matter models

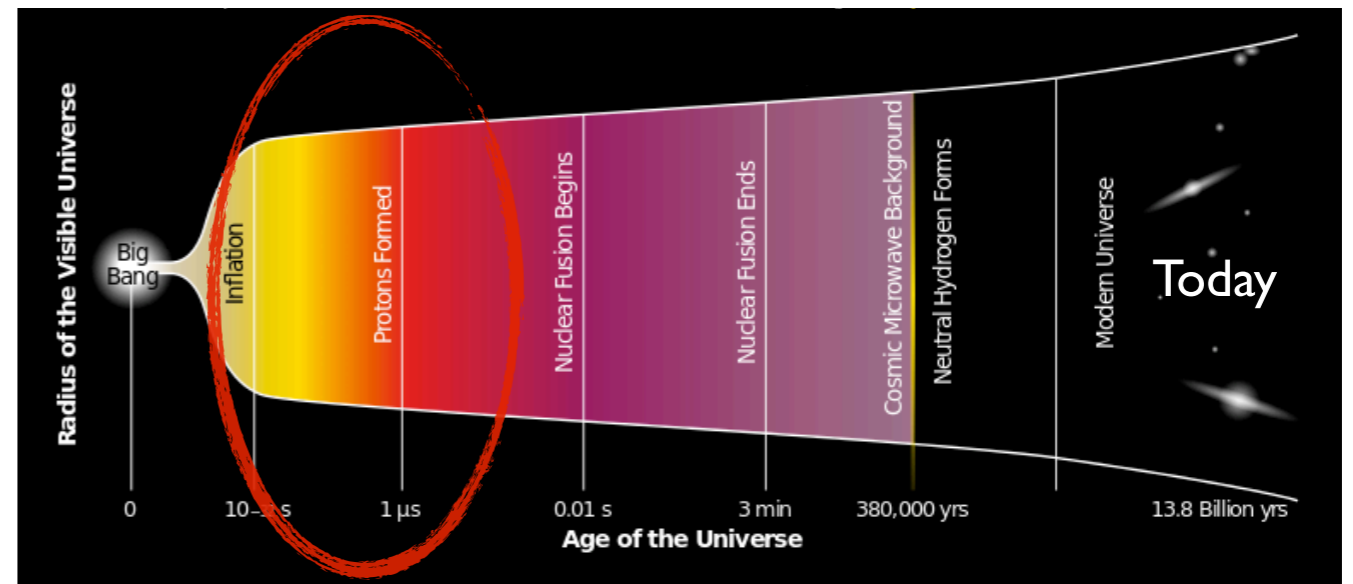
Decisive property: Dark matter thermalized in expanding early Universe?

$$\Gamma_{\text{ann}} \gtrsim H$$

Rate of thermalizing processes:



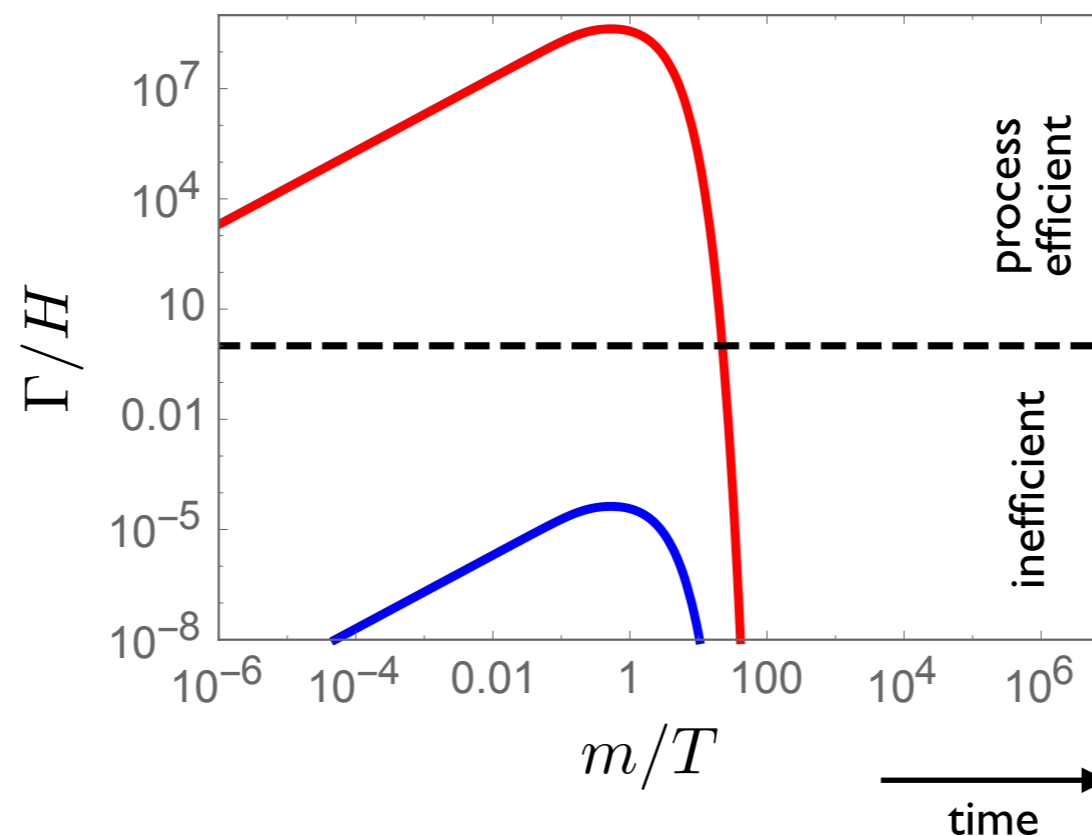
Expansion with Hubble rate H



Decisive property: Dark matter thermalized in expanding early Universe?

$$H \propto T^2, \quad \Gamma_{\text{ann}} \propto \begin{cases} T & , T \gg m \text{ (ultra-rel.)} \\ e^{-m/T} & , T < m \text{ (non-rel.)} \end{cases}$$

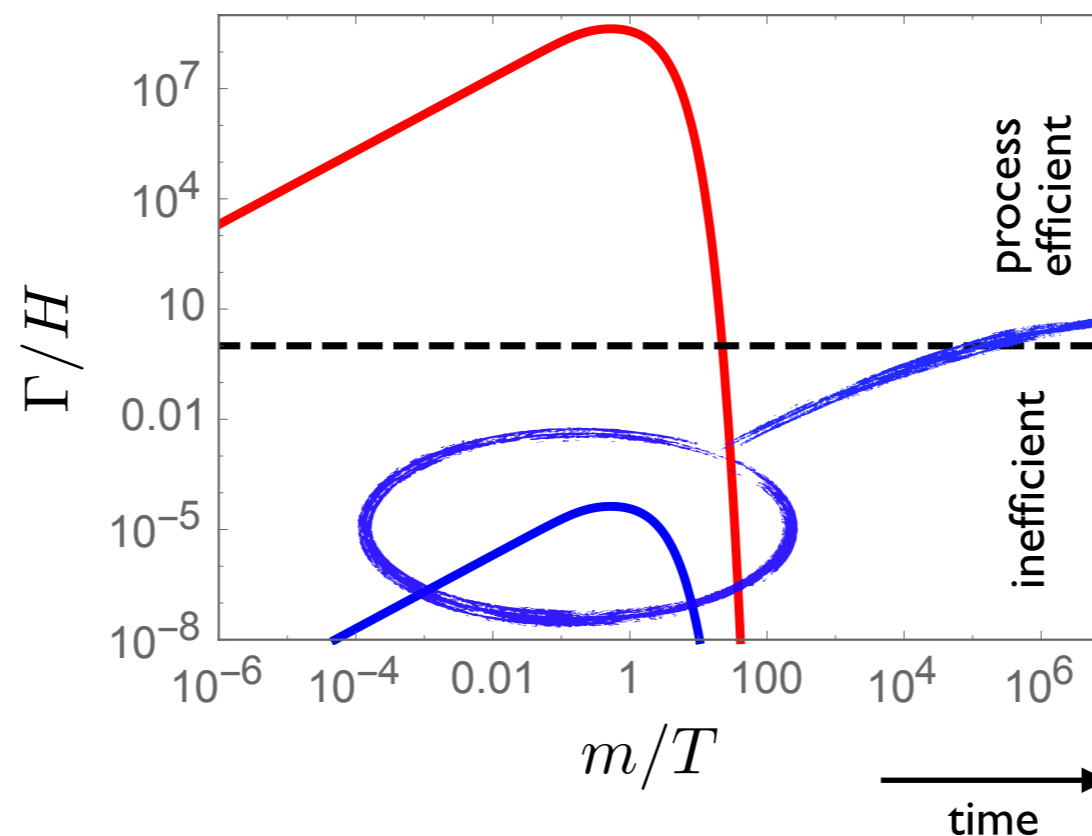
- $2 \rightarrow 2$ processes maximum, thermalization processes can remain inefficient:



Decisive property: Dark matter thermalized in expanding early Universe?

$$H \propto T^2, \quad \Gamma_{\text{ann}} \propto \begin{cases} T & , T \gg m \text{ (ultra-rel.)} \\ e^{-m/T} & , T < m \text{ (non-rel.)} \end{cases}$$

- $2 \rightarrow 2$ processes maximum, thermalization processes can remain inefficient:

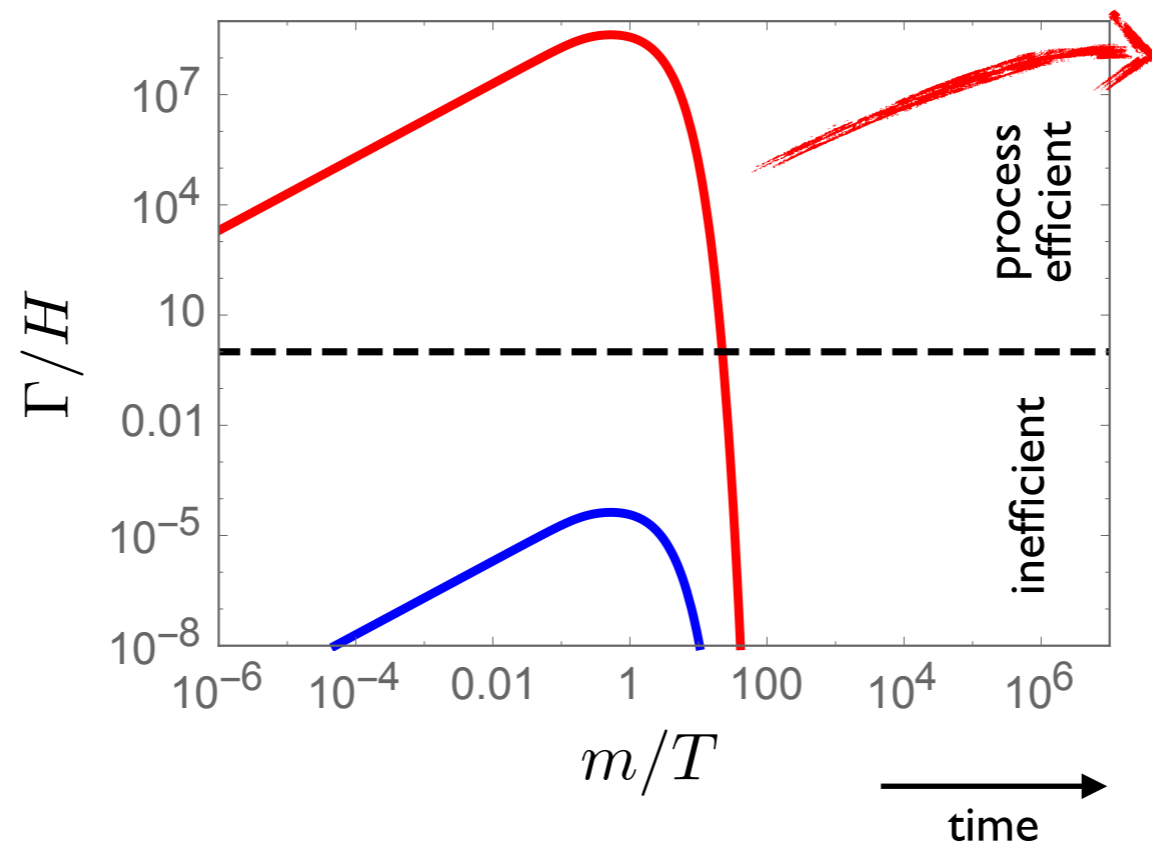


Non-th. dark sector:
No LHC signal!
[see Kahlhoefer
1801.07621]

Decisive property: Dark matter thermalized in expanding early Universe?

$$H \propto T^2, \quad \Gamma_{\text{ann}} \propto \begin{cases} T & , T \gg m \text{ (ultra-rel.)} \\ e^{-m/T} & , T < m \text{ (non-rel.)} \end{cases}$$

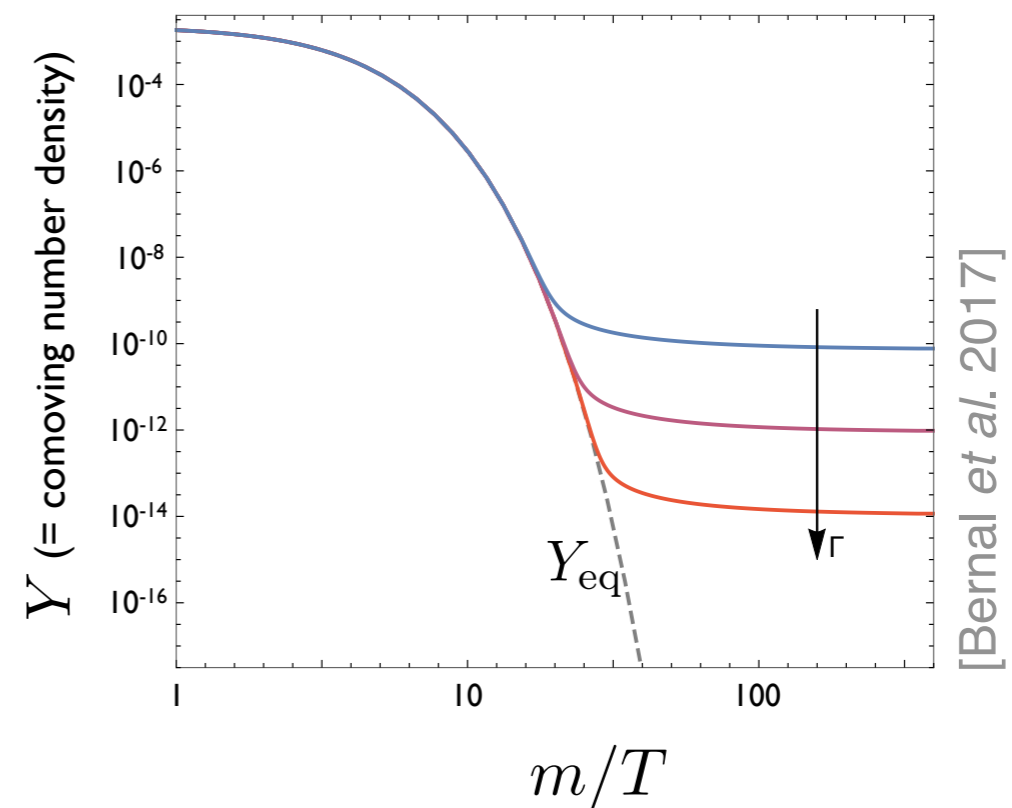
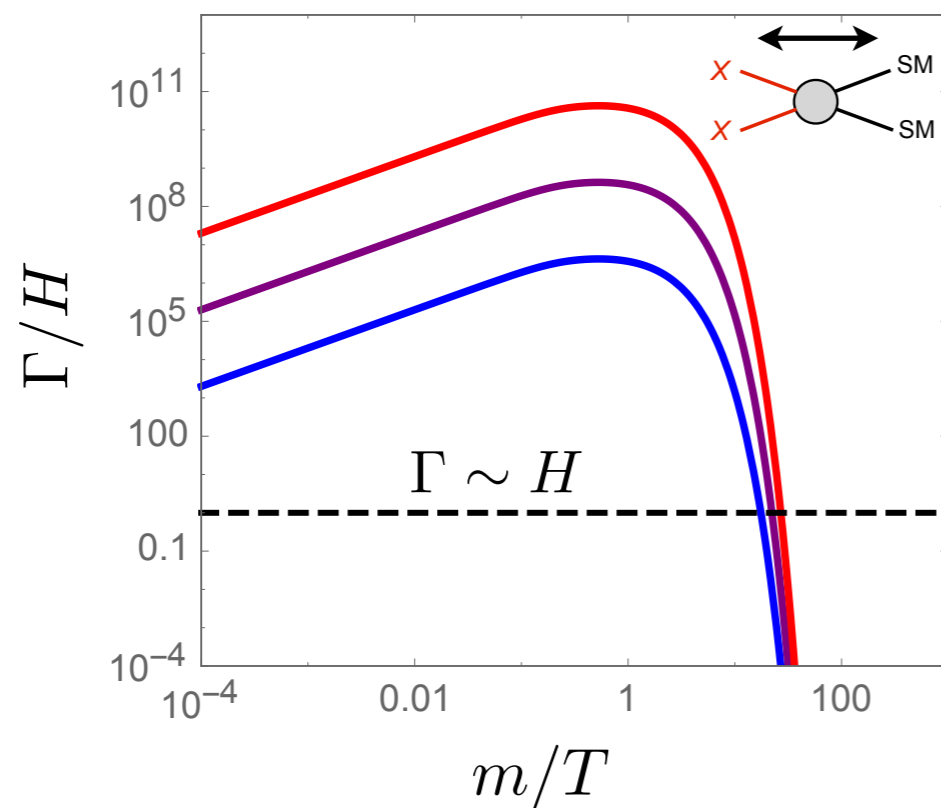
- $2 \rightarrow 2$ processes maximum, thermalization processes can remain inefficient:



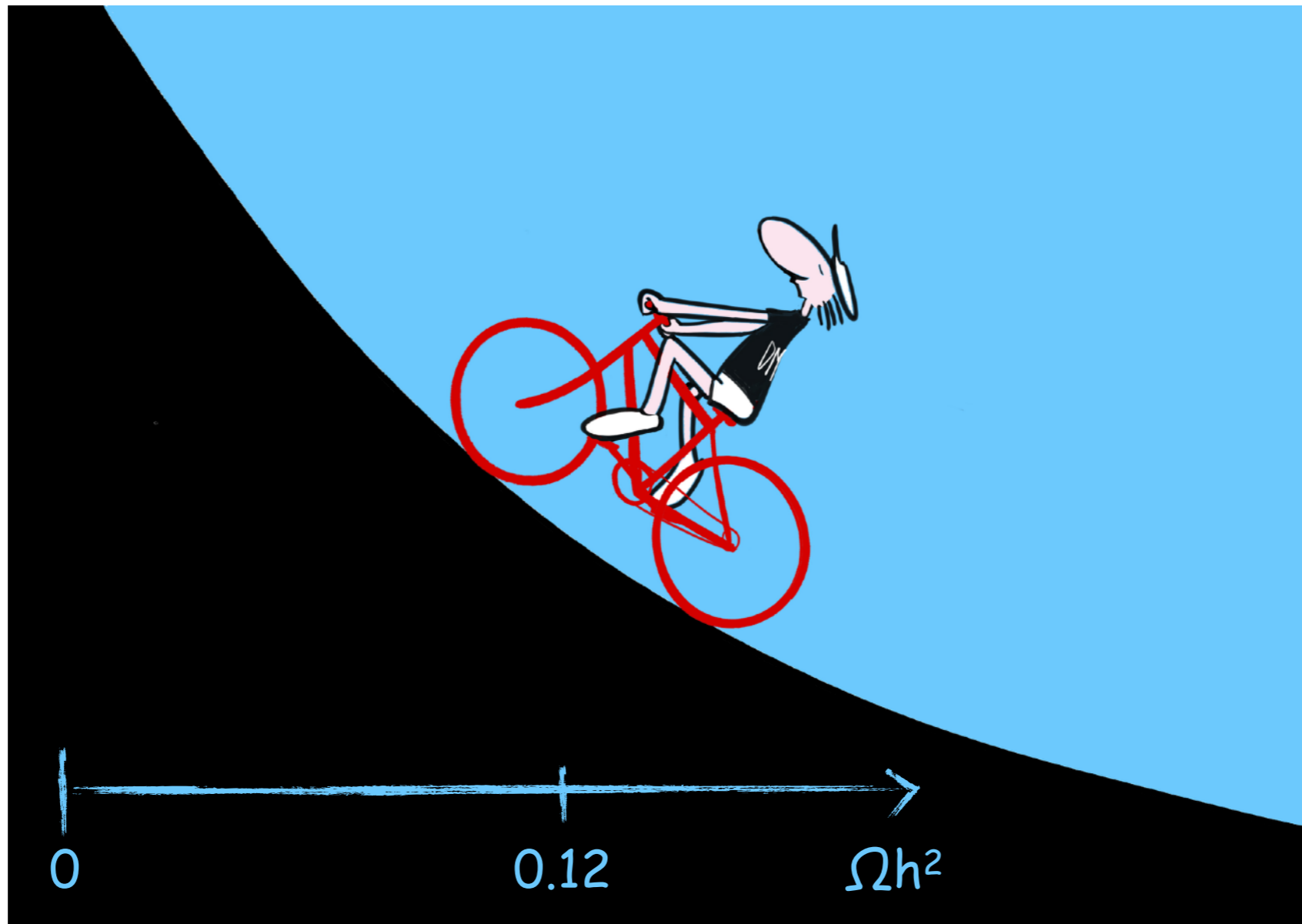
At least part of dark sector thermalized:
Of interest for LHC

Vanilla WIMP

- Weak couplings: well thermalized! → independent of initial cond.
- Well-known freeze-out picture:
Leaves equilibrium density when $\Gamma \sim H$ → cold DM



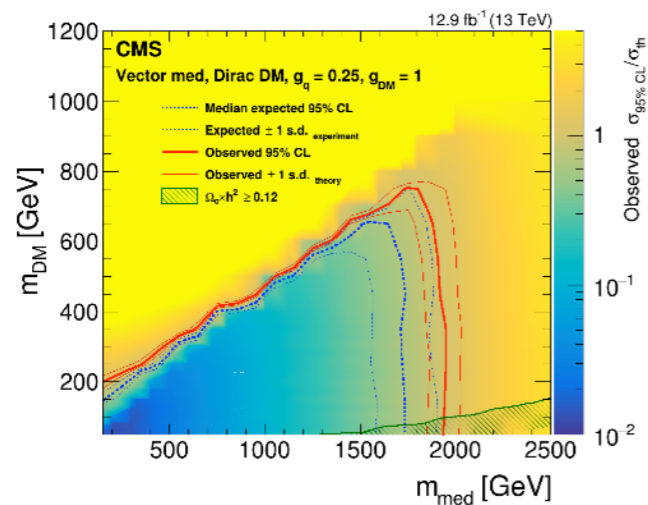
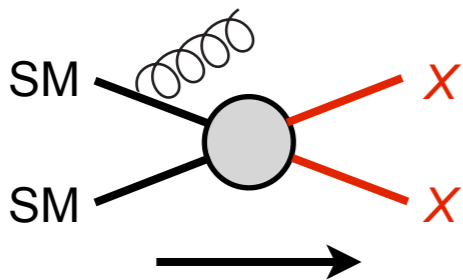
Vanilla WIMP



Dark Matter searches

WIMPs (weakly interacting massive particles)

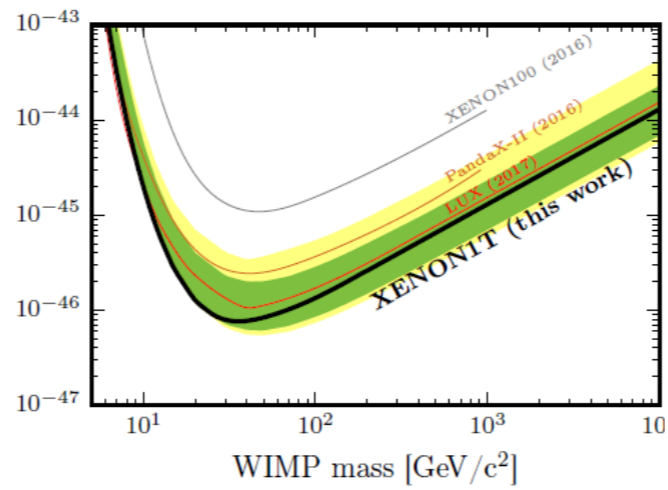
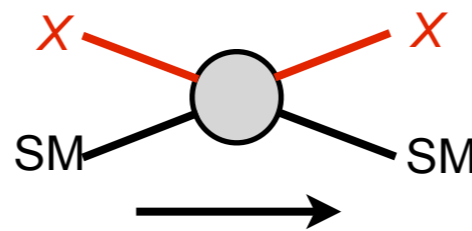
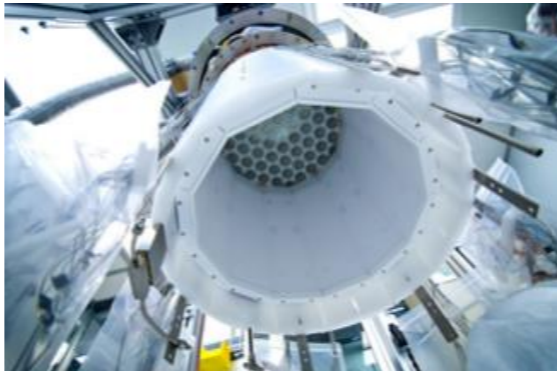
Direct production



[CMS EXO-16-039]

→ talk of Francesco

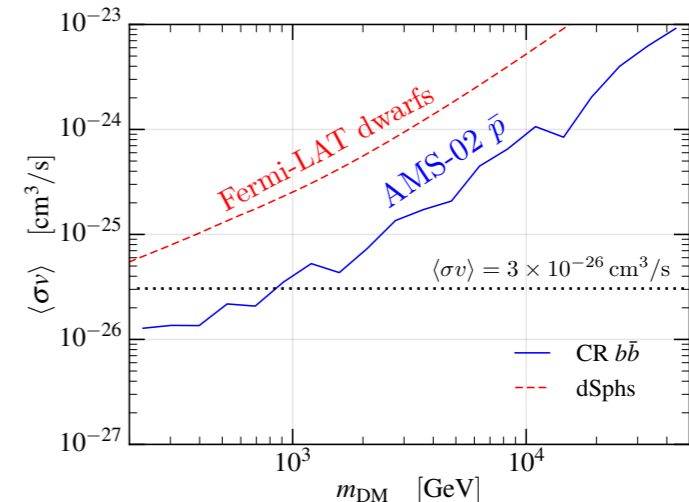
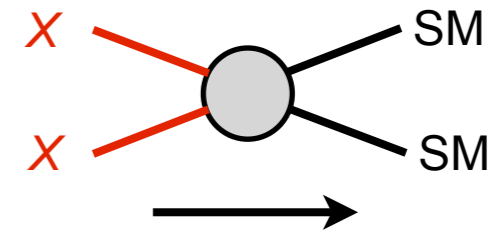
Direct detection



[XenonIT 1705.06655]

→ talk of Daniel

Indirect detection



[Cuoco, JH, Korsmeier, Krämer 2017]

→ talk of Michael (YSF)

Dark Matter searches

WIMPs (weakly interacting massive particles)

Direct production



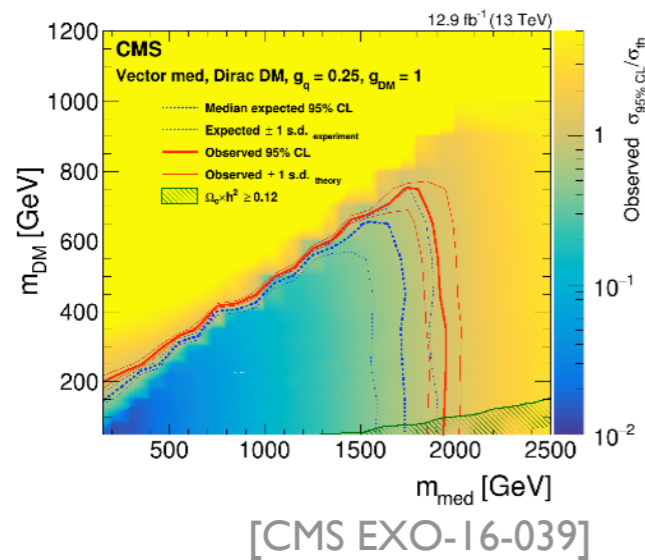
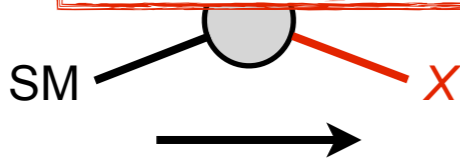
Direct detection



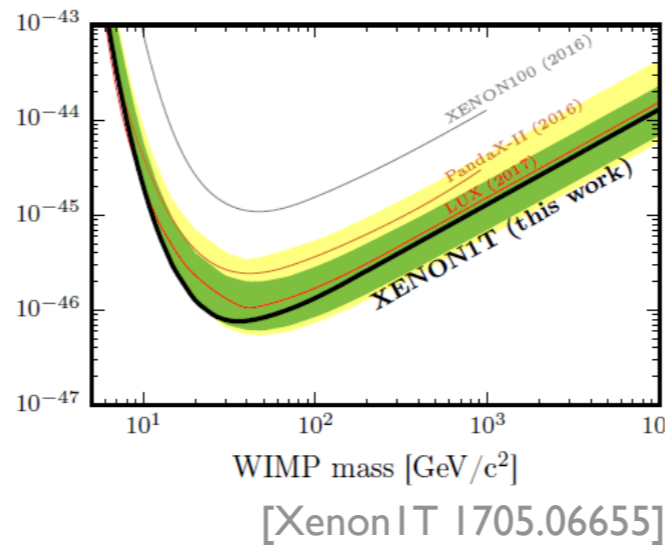
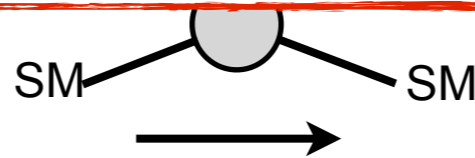
Indirect detection



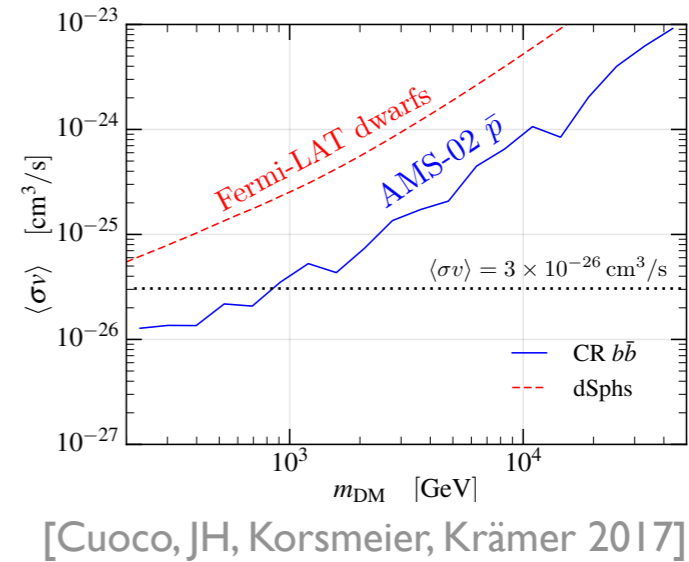
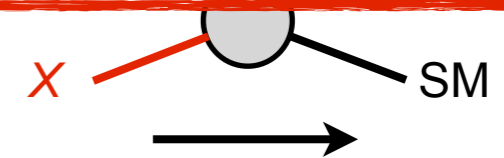
WIMPs strongly constrained by current searches
 Scenarios with other signatures?



→ talk of Francesco



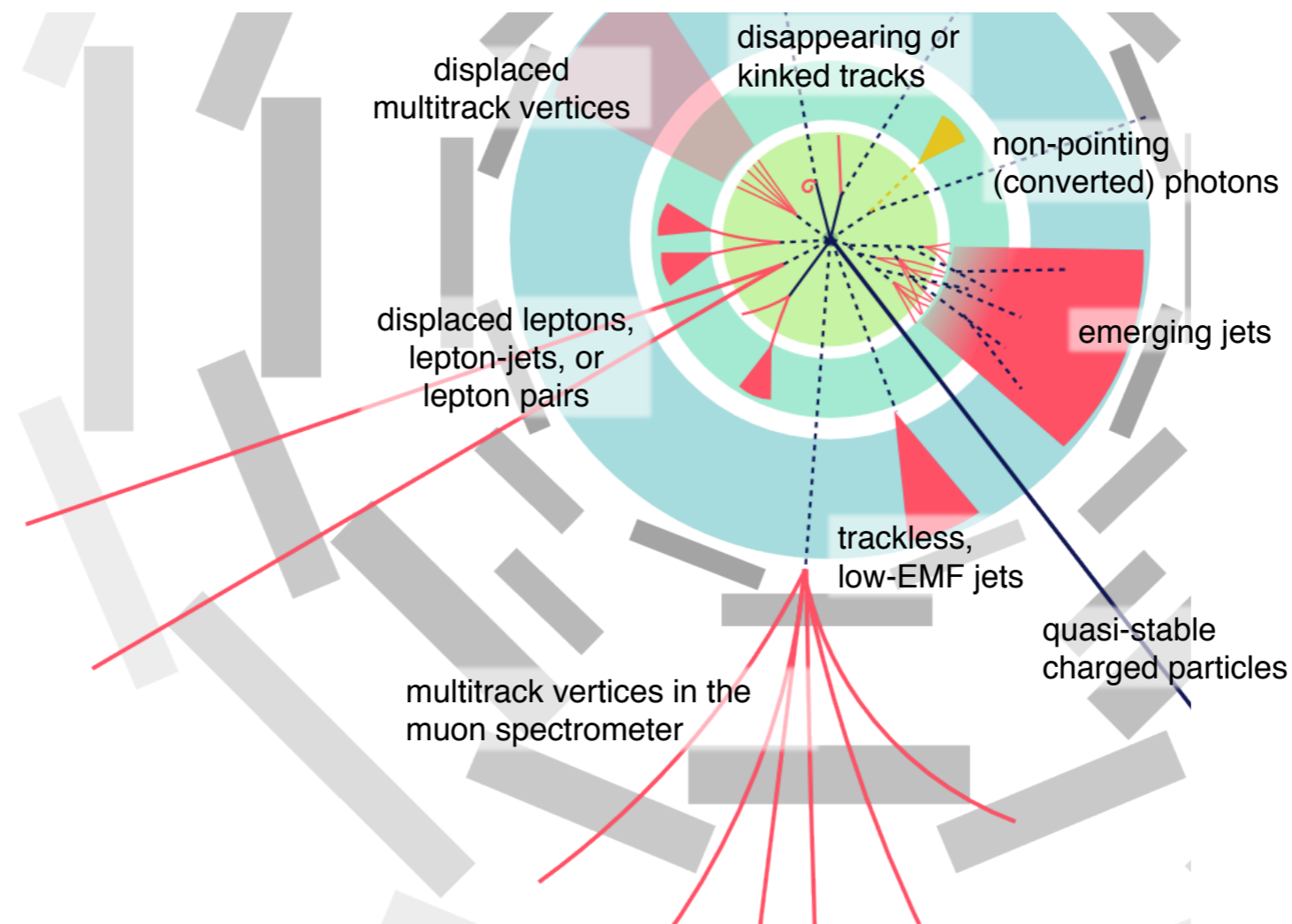
→ talk of Daniel



→ talk of Michael (YSF)

Long-lived particles (LLPs) at the LHC

- Metastable BSM particles (except neutral detector-stable particles \rightarrow MET)
- Wide range of possible signatures:



[Figure from Heather Russell]

- Many LLP signatures low background \rightarrow high sensitivity

LLP searches at the LHC: challenges

- **Cover spectrum of possible LLP signatures and kinematic ranges**
so far mostly SUSY related scenarios
- **Hard to trigger**
often additional signature needed for triggering
- **Interpretation of results**
applicability for a large set of models required
- **Recasting difficult**
exchange of additional information needed

LLP searches at the LHC: challenges

- Cover spectrum of possible LLP signatures and kinematic ranges
so far mostly SUSY related scenarios
- Hard to trigger
often additional signature needed for triggering
- Interpretation of results
applicability for a large set of models required
- Recasting difficult
exchange of additional information needed

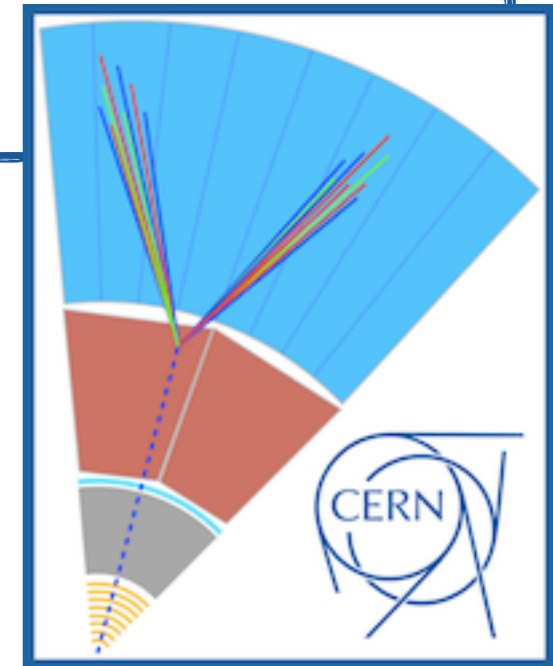
LHC LLP Community Workshops:

April 2017: <https://indico.cern.ch/event/607314/>

October 2017: <https://indico.cern.ch/event/649760/>

Next Workshop: 16-18 May 2018 at CERN

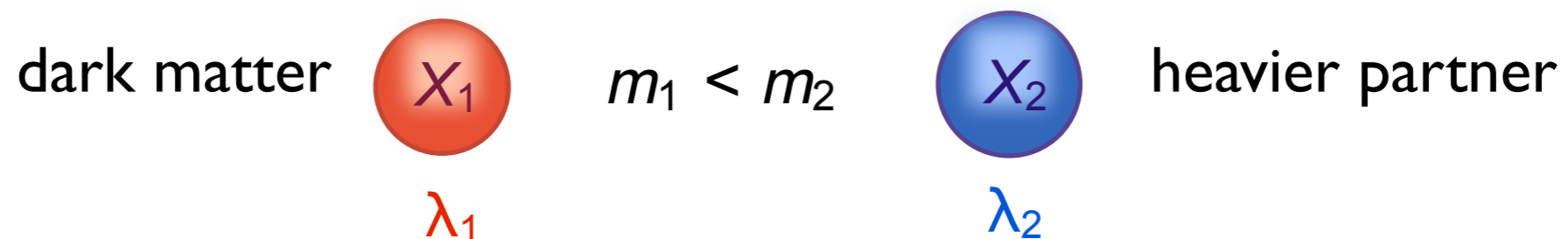
Community white paper in preparation



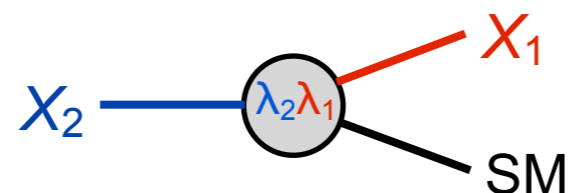
LLPs and dark matter

LLPs and dark matter

- Consider Z_2 -odd dark sector

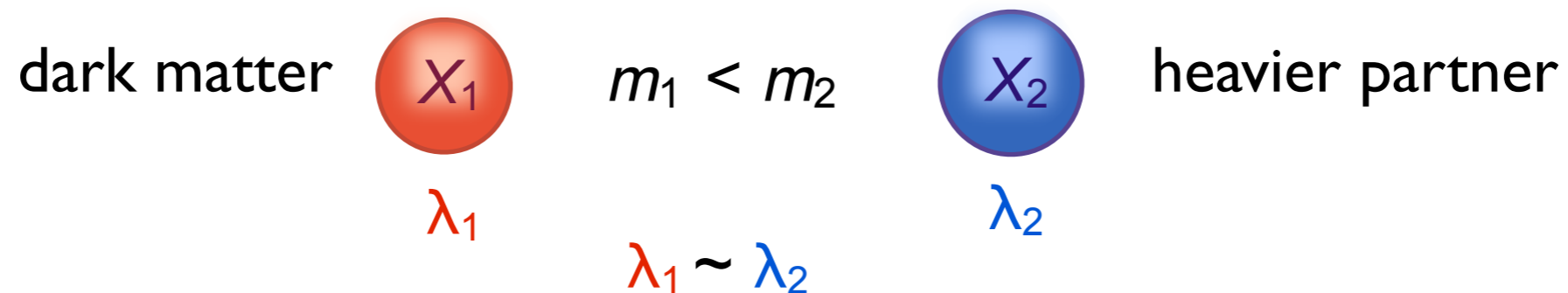


- Decay only in Z_2 -odd states:



- Consider: next heavier state X_2 involved in freeze-out!

Coannihilation [Griest, Seckel 1991; Edsjo, Gondolo 1997]



Number densities during freeze-out:

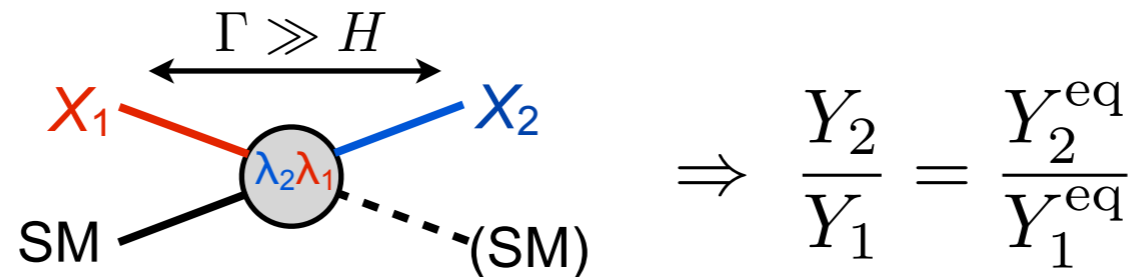
$$\frac{Y_2^{\text{eq}}}{Y_1^{\text{eq}}} \propto e^{-\Delta m/T_f} \sim e^{-25 \Delta m/m}$$

$$\Delta m/m \lesssim 10\% \quad [\text{for an interesting exception see D'Agnolo et al. 1803.02901}]$$

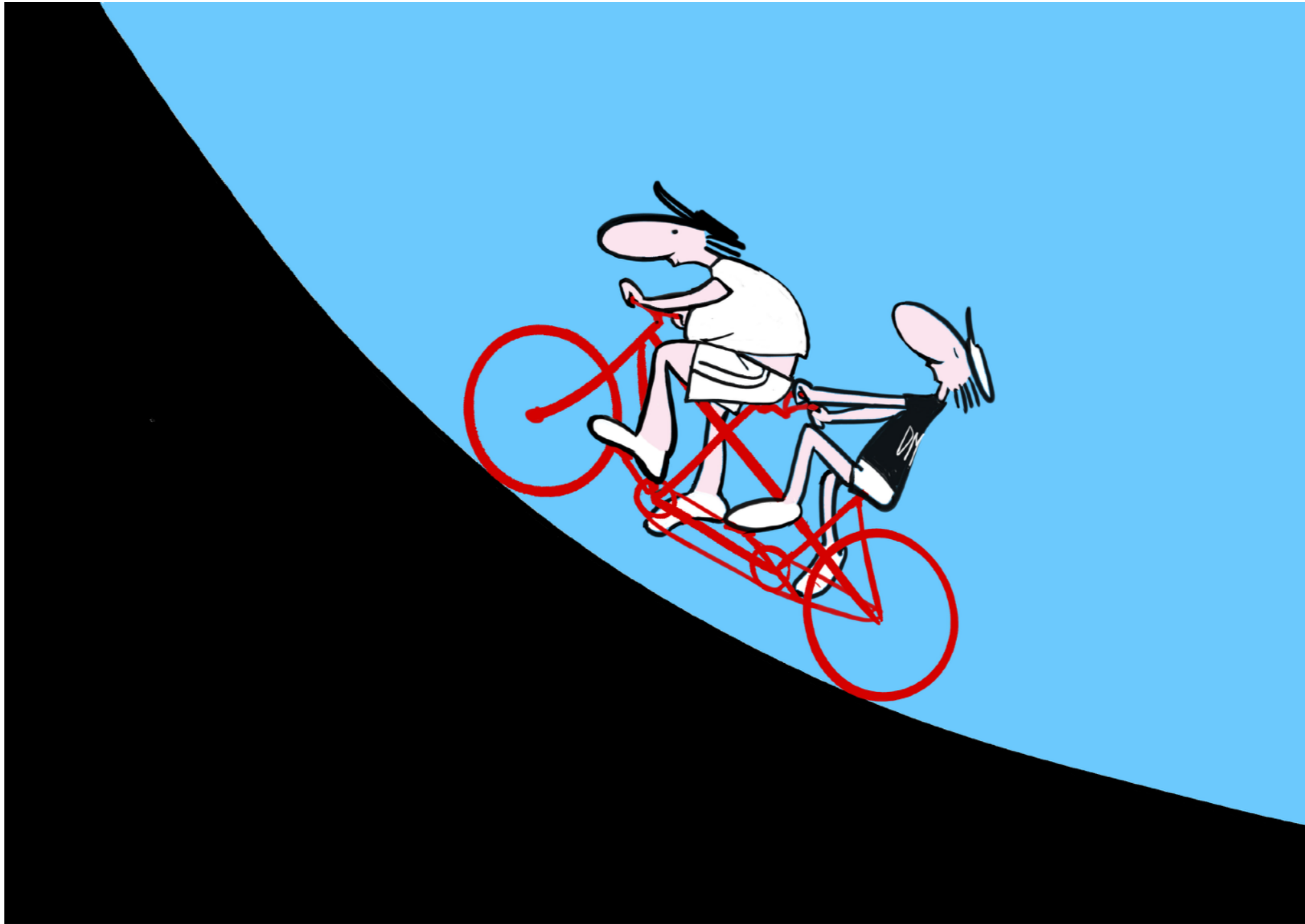
→ X_2 still "around" during freeze-out

Coannihilation

Number densities "tied" together through efficient conversions:

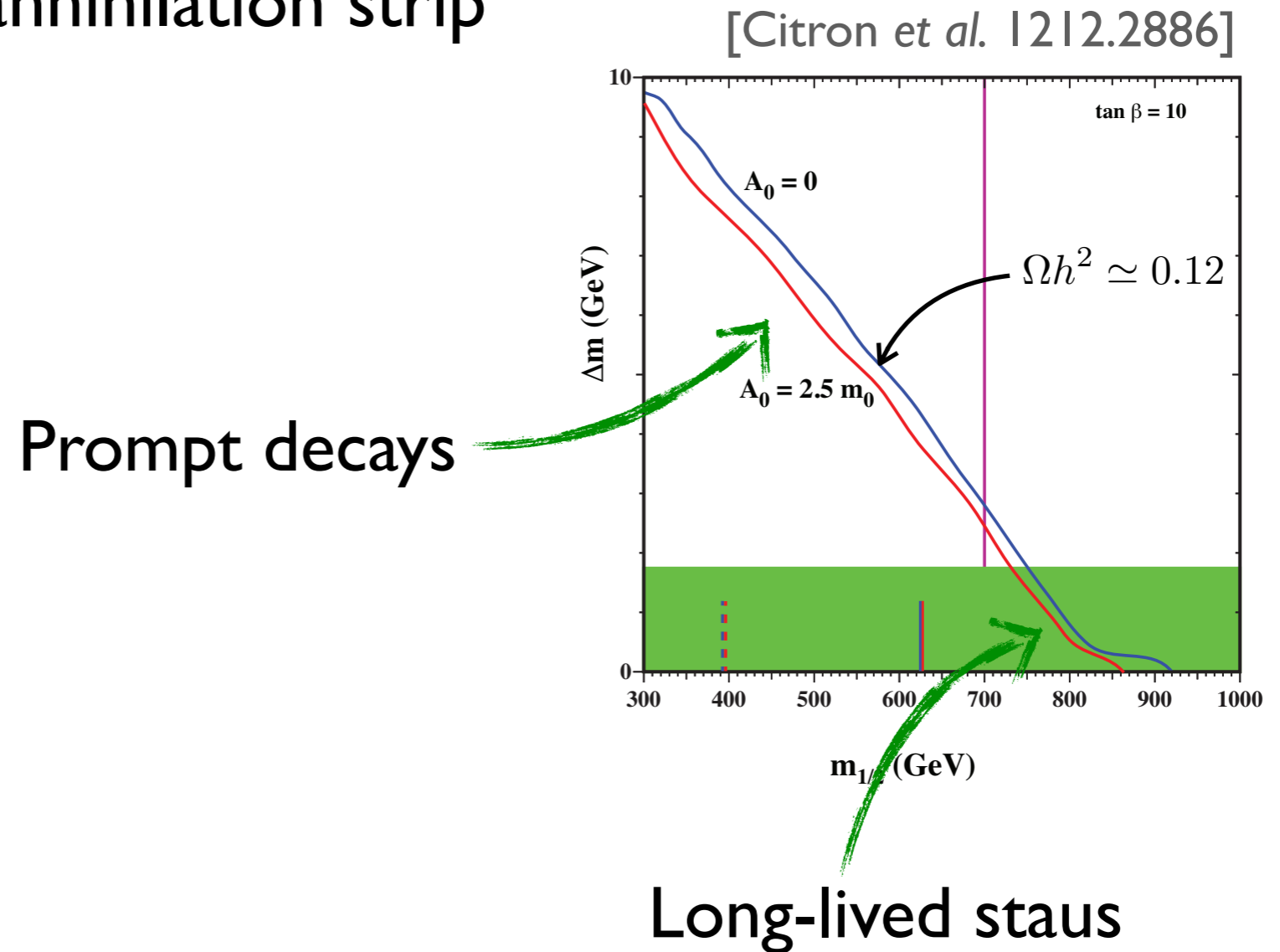


Coannihilation



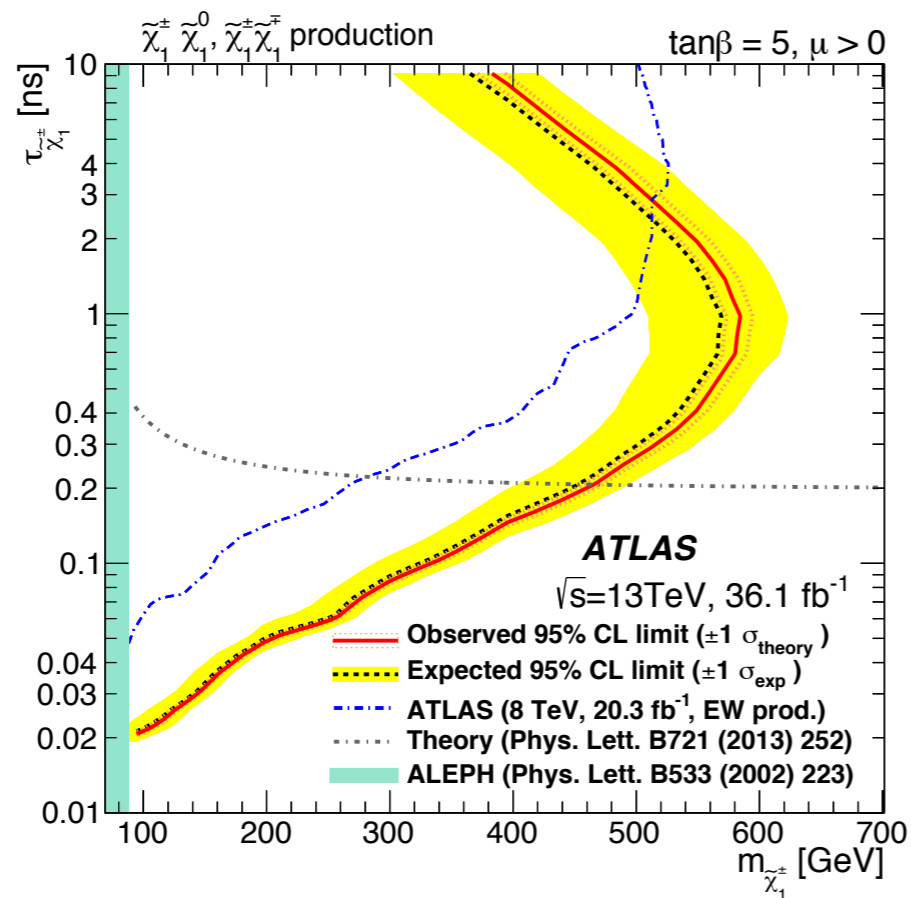
Coannihilation and LLPs

- SUSY: stau coannihilation strip

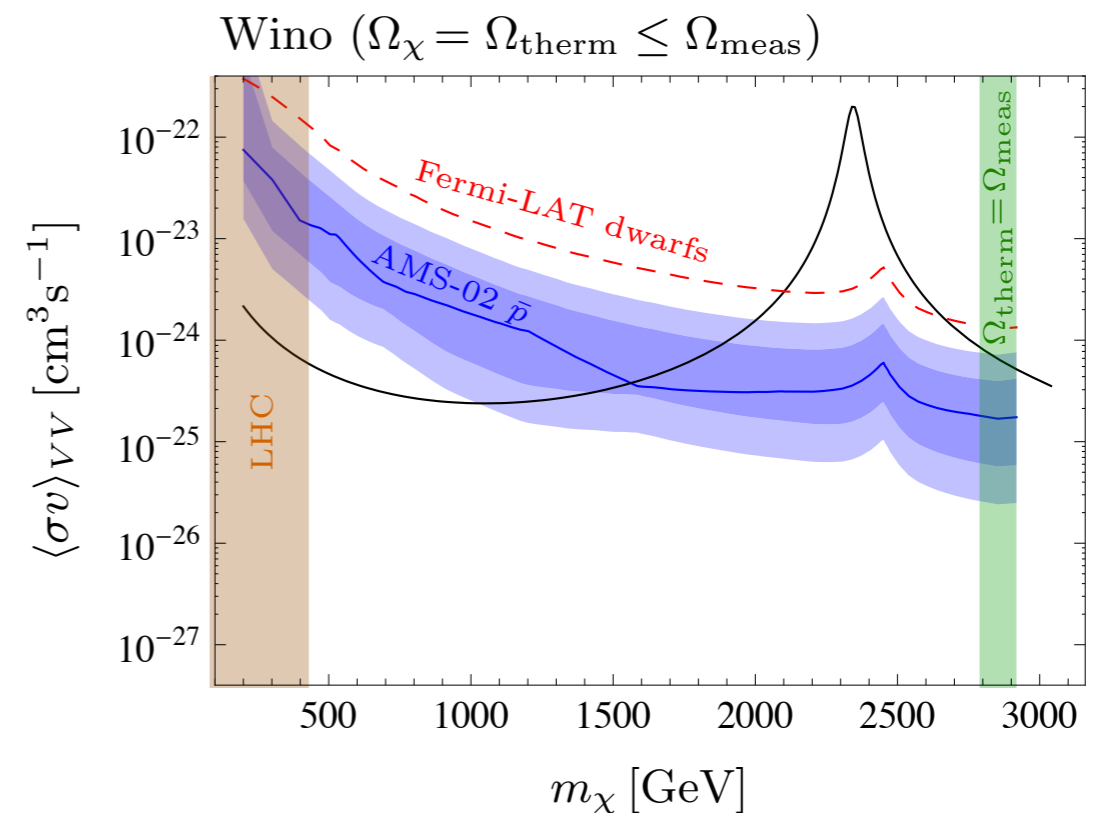


Coannihilation and LLPs

- SUSY: stau coannihilation strip
- Minimal dark matter, e.g. pure wino:
[Cirelli *et al.* hep-ph/0512090]



[Cuoco, JH, Korsmeier, Krämer 1711.05274]



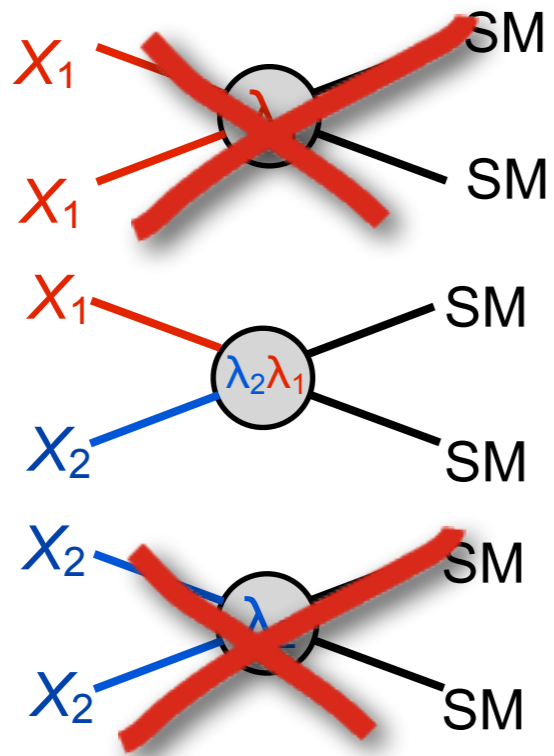
Coannihilation and LLPs

- SUSY: stau coannihilation strip
- Minimal dark matter e.g. pure wino

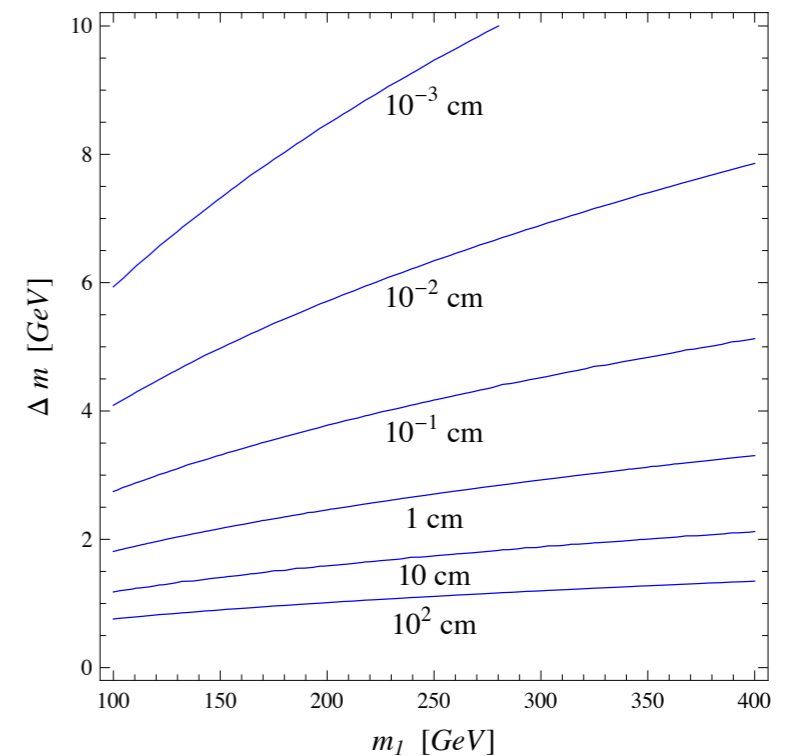
[Cirelli *et al.* hep-ph/0512090]

- Pseudo Dirac dark matter

[De Simone, Sanz, Sato 2010; Davolia, De Simone, Jacquesa, Sanz 2017]



→ evade direct and indirect detection bounds
→ predicts LLPs

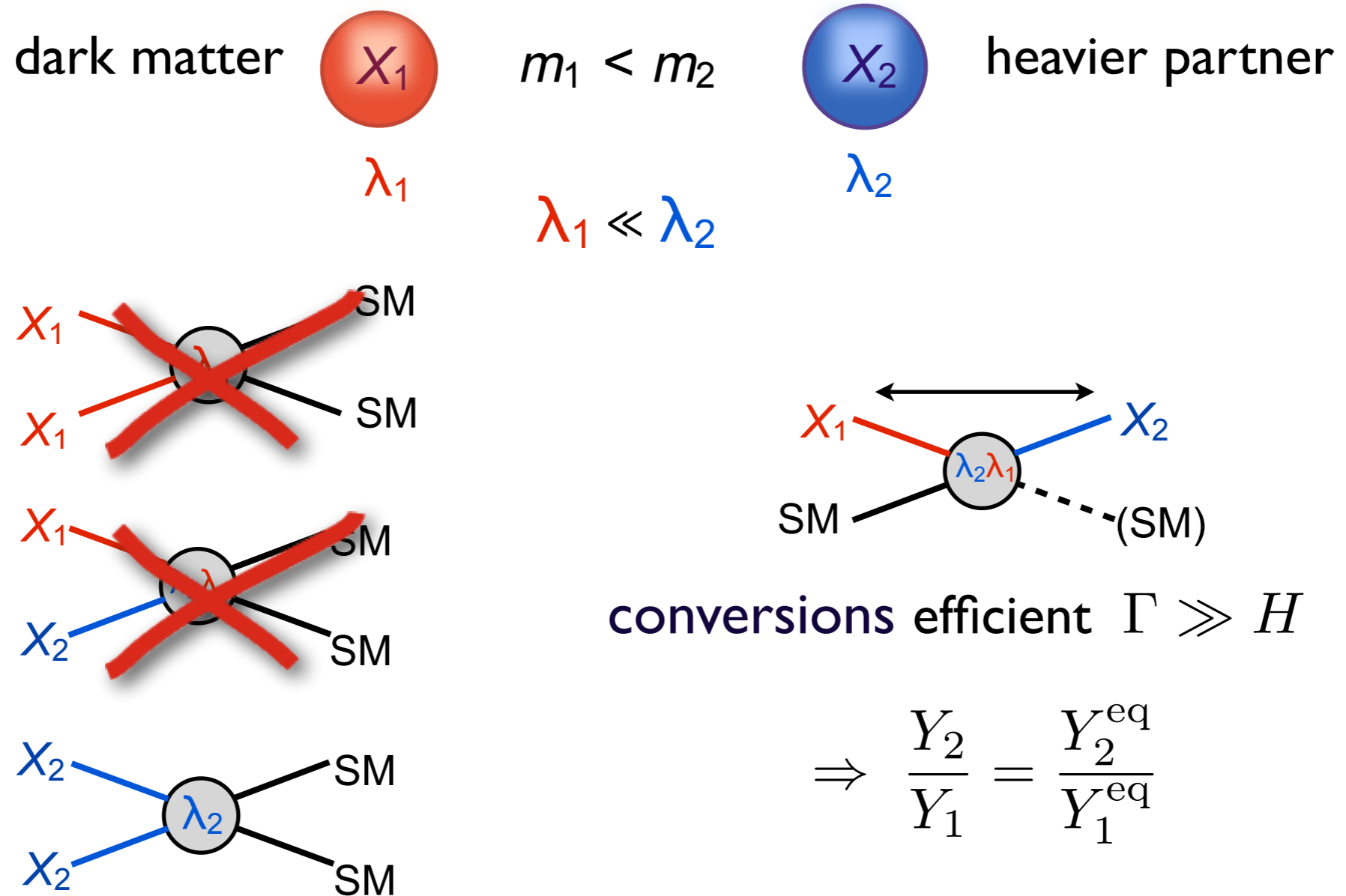


Coannihilation and LLPs

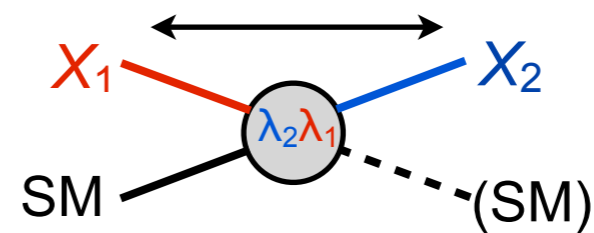
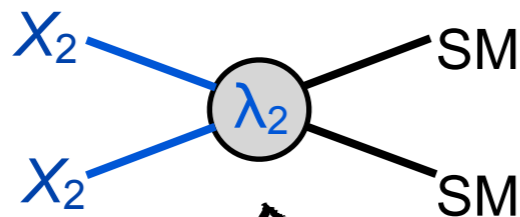
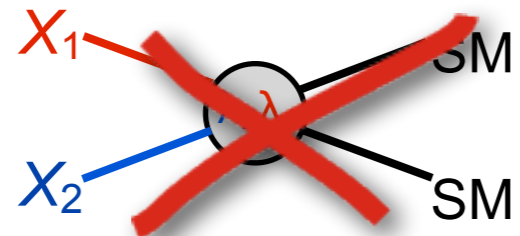
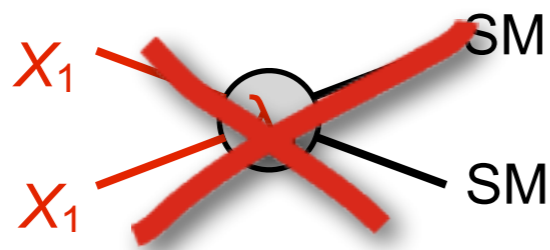
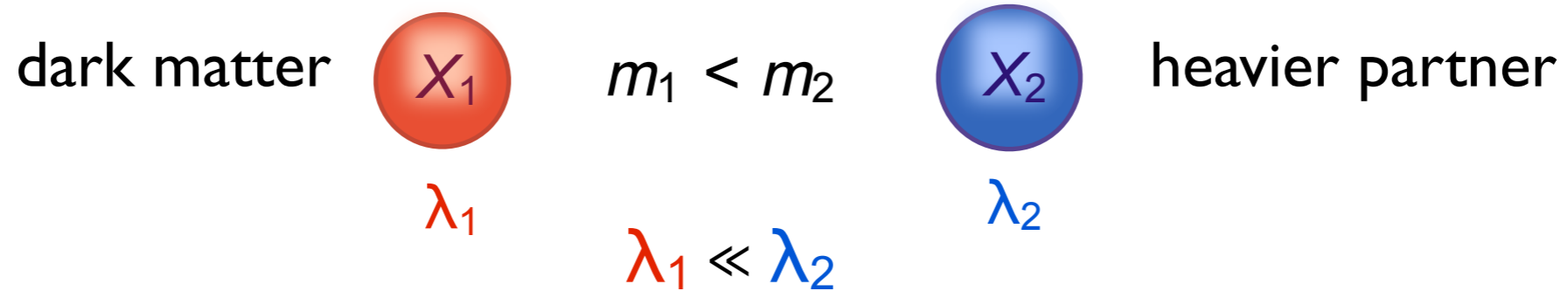
- SUSY: stau coannihilation strip
- Minimal dark matter e.g. pure wino
[Cirelli *et al.* hep-ph/0512090]
- Pseudo Dirac dark matter
[De Simone, Sanz, Sato 2010; Davolia, De Simone, Jacquesa, Sanz 2017]
- ...

LLP due to small mass splittings

Small dark matter couplings



Small dark matter couplings

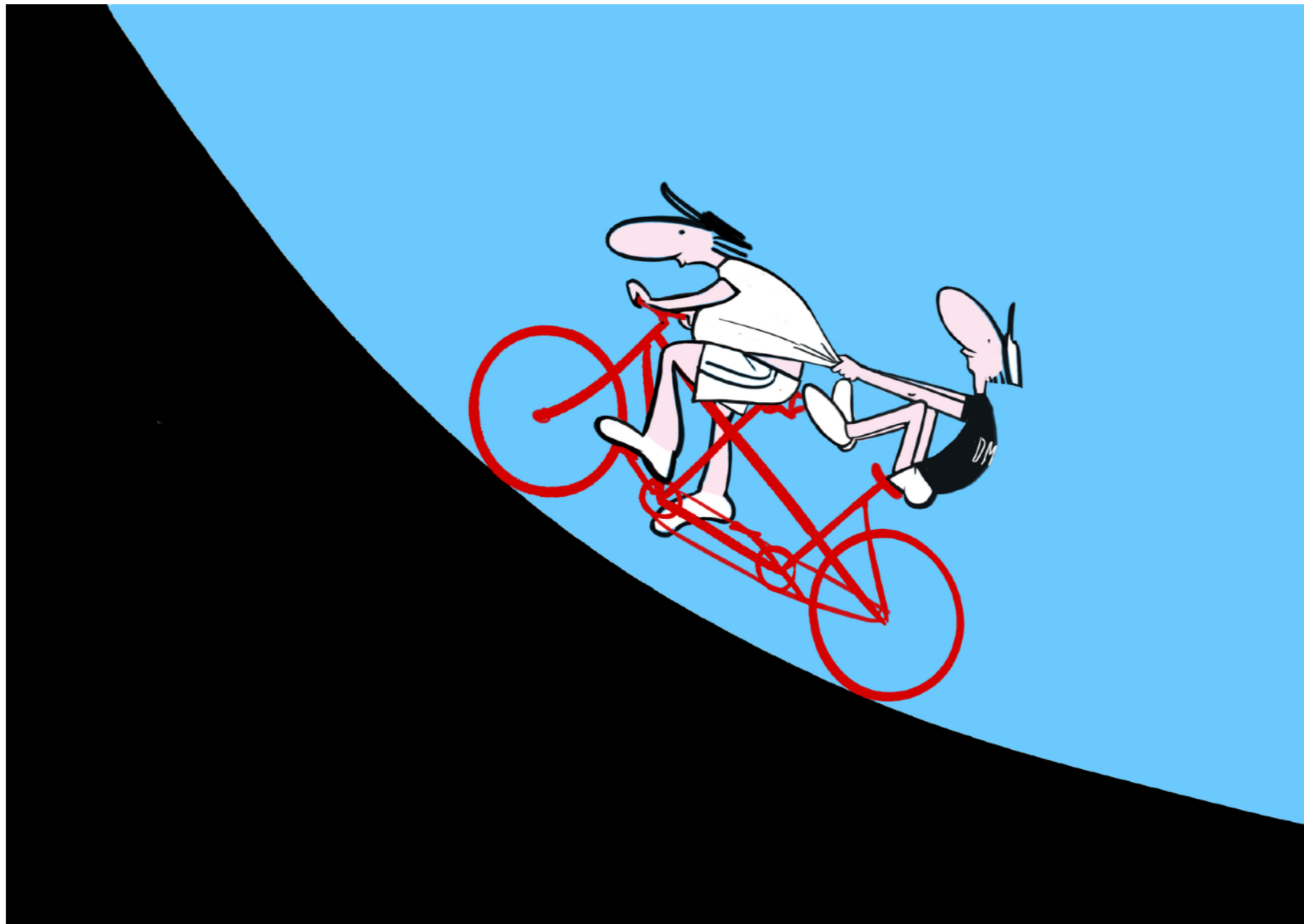


conversions efficient $\Gamma \gg H$

$$\Rightarrow \frac{Y_2}{Y_1} = \frac{Y_2^{\text{eq}}}{Y_1^{\text{eq}}}$$

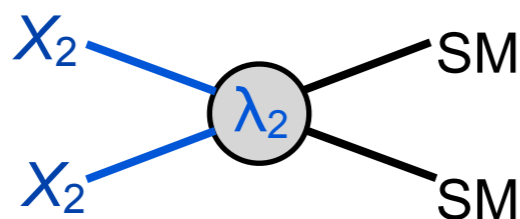
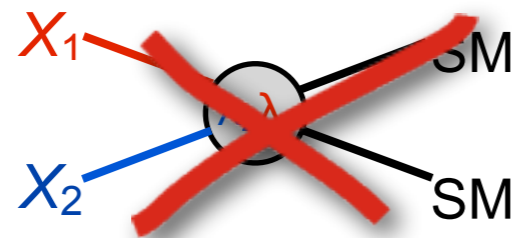
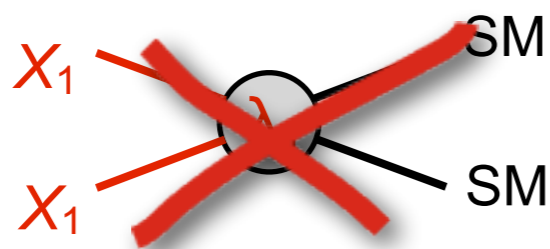
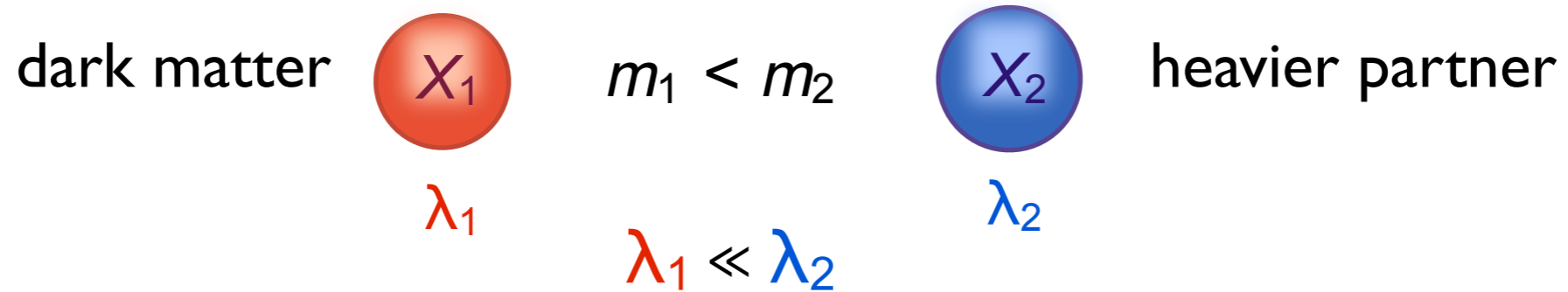
annihilation of X_2 alone sets relic density

Freeloader's freeze-out

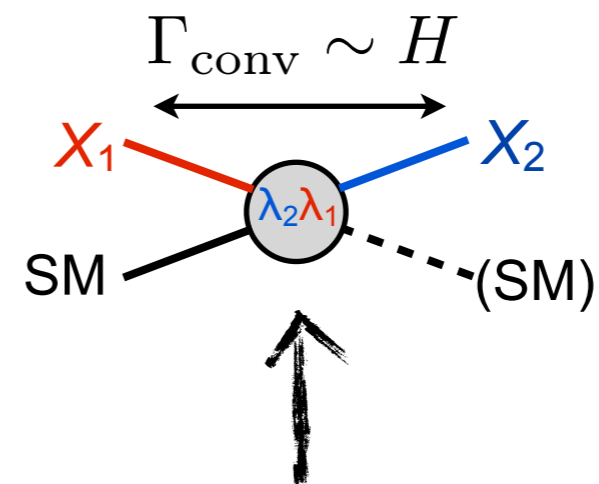


Even smaller dark matter couplings

[Garny, JH, Lülf, Vogl 1705.09292; D'Agnolo, Pappadopulo, Ruderman 1705.08450; Garny, JH, Hufnagel, Lülf 1802.00814]

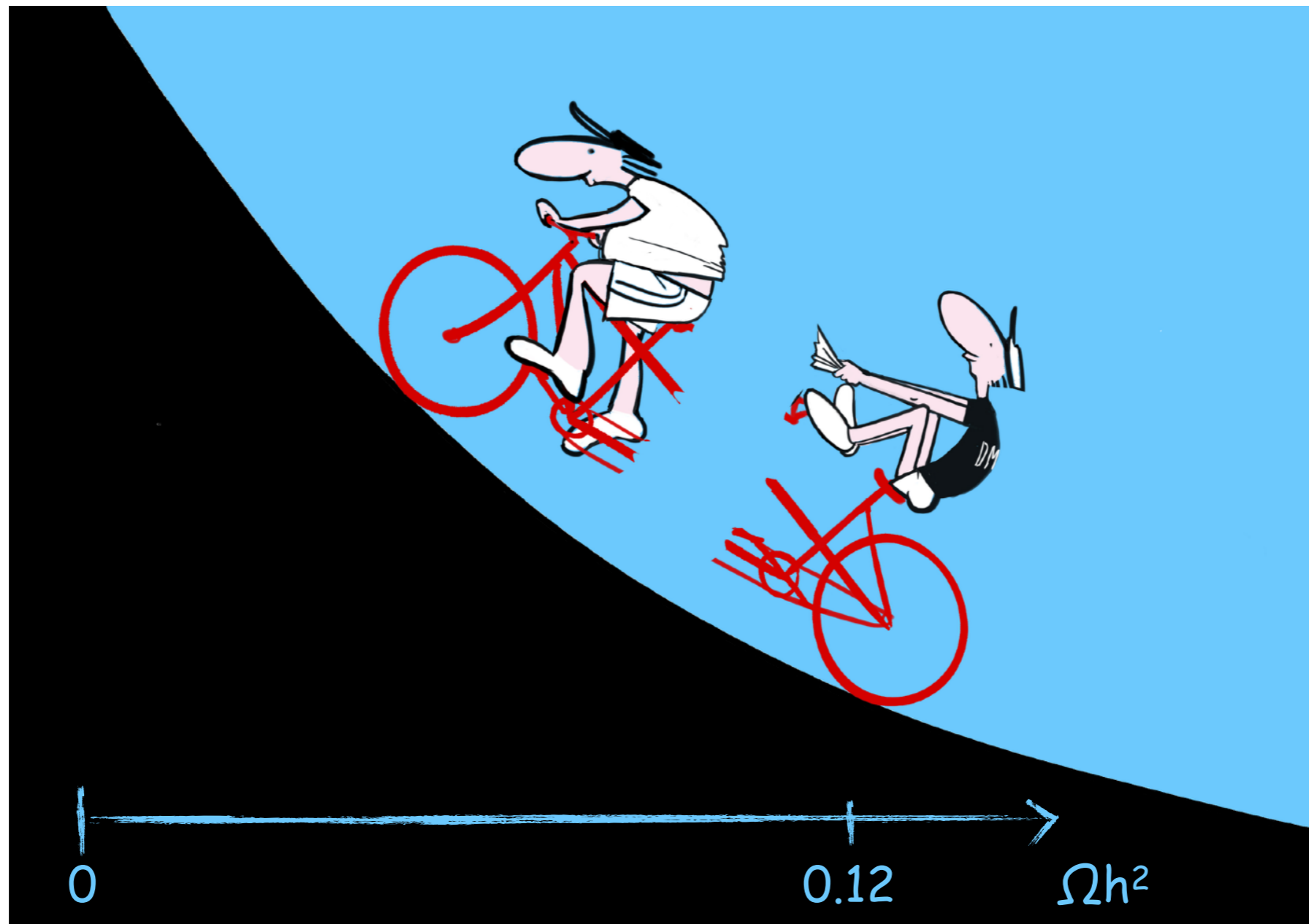


conversions become inefficient



conversion sets relic density

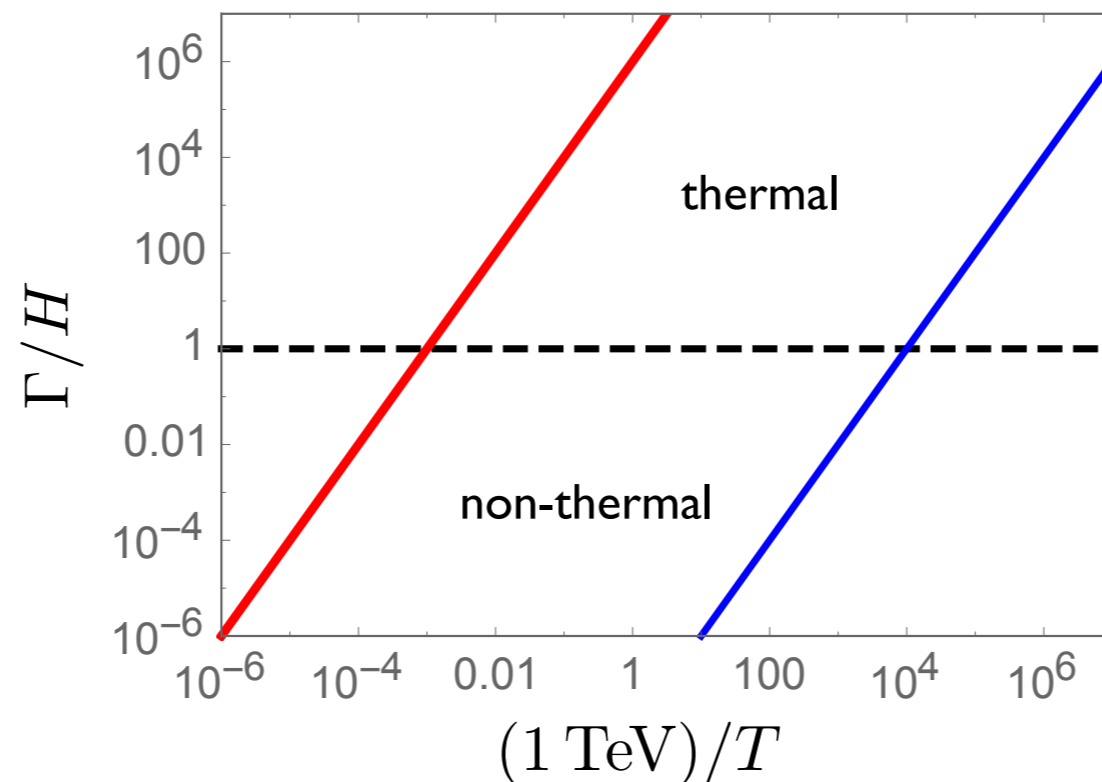
Conversion-driven freeze-out



Pheno of conversion-driven freeze-out

- Consider decay:

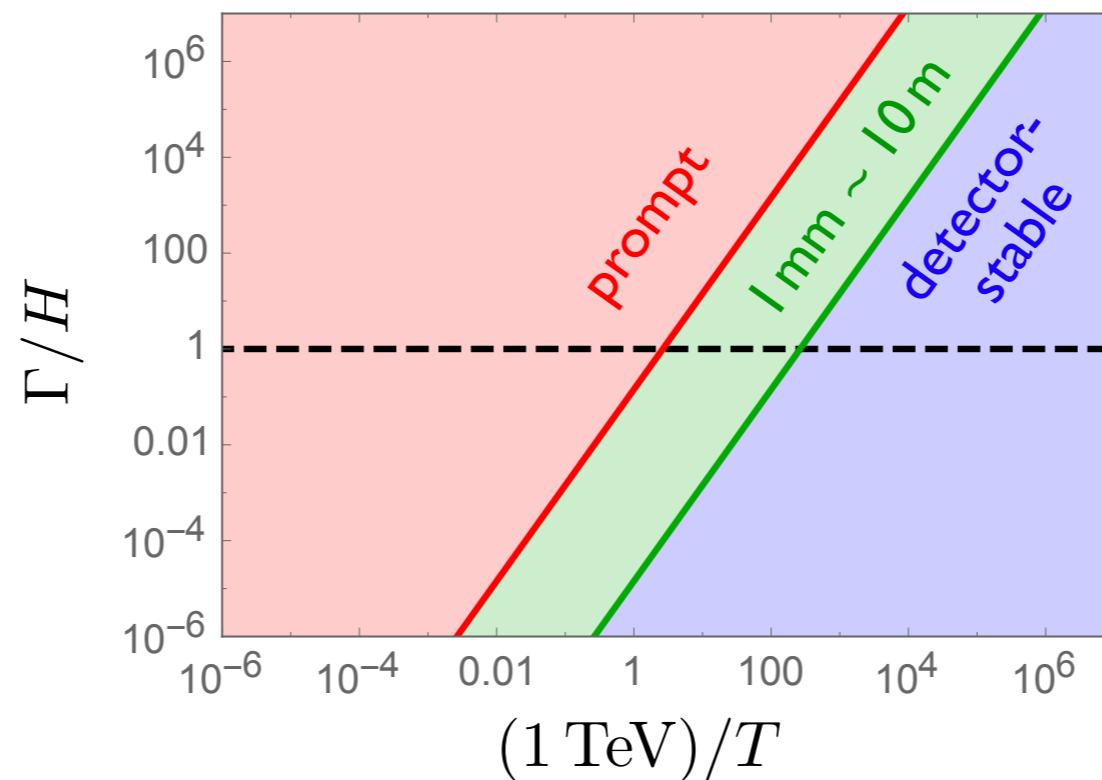
$$H \propto T^2, \quad \Gamma_{\text{dec}} \sim \text{const.}$$



Pheno of conversion-driven freeze-out

- Consider decay:

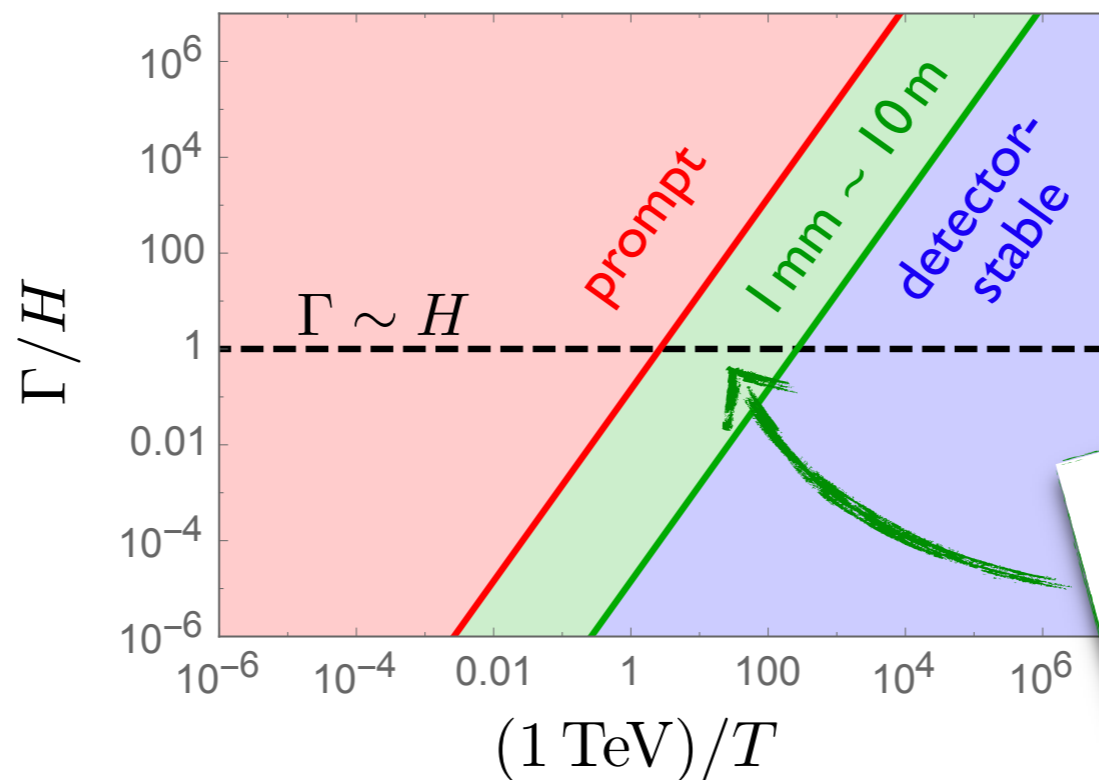
$$H \propto T^2, \quad \Gamma_{\text{dec}} \sim \text{const.}$$



Pheno of conversion-driven freeze-out

- Consider decay:

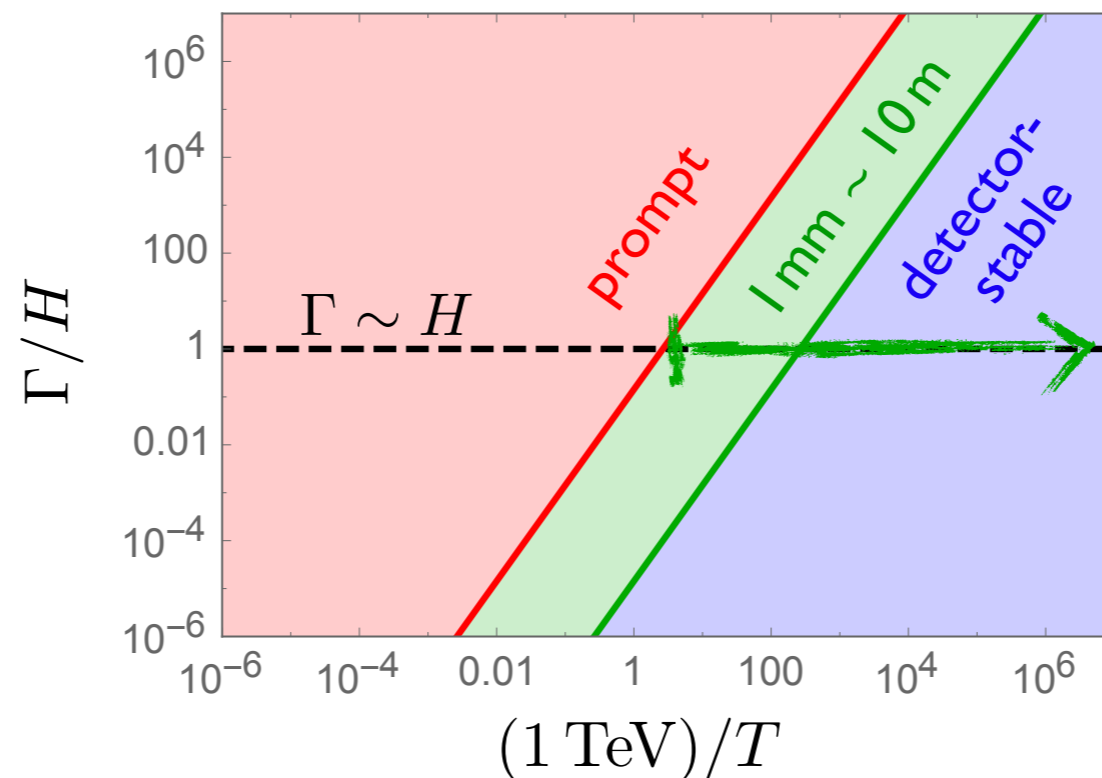
$$H \propto T^2, \quad \Gamma_{\text{dec}} \sim \text{const.}$$



weak scale ~
detector size
LLP miracle

Pheno of conversion-driven freeze-out

Neither decay nor $2 \rightarrow 2$ scattering can be efficient around freeze-out: Decay non-prompt!

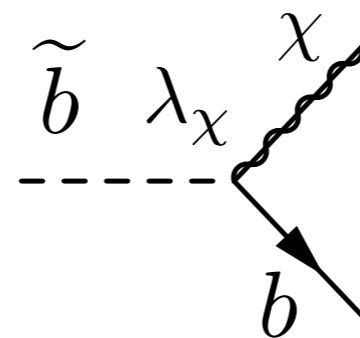


A concrete example

- Specific model: $\mathcal{L}_{\text{int}} = |D_\mu \tilde{q}|^2 - \lambda_\chi \tilde{q} \bar{q} \frac{1 - \gamma_5}{2} \chi + \text{h.c.}$
- SUSY-inspired simplified model:
Choose Majorana DM and scalar bottom-partner



- Yukawa-type interaction:

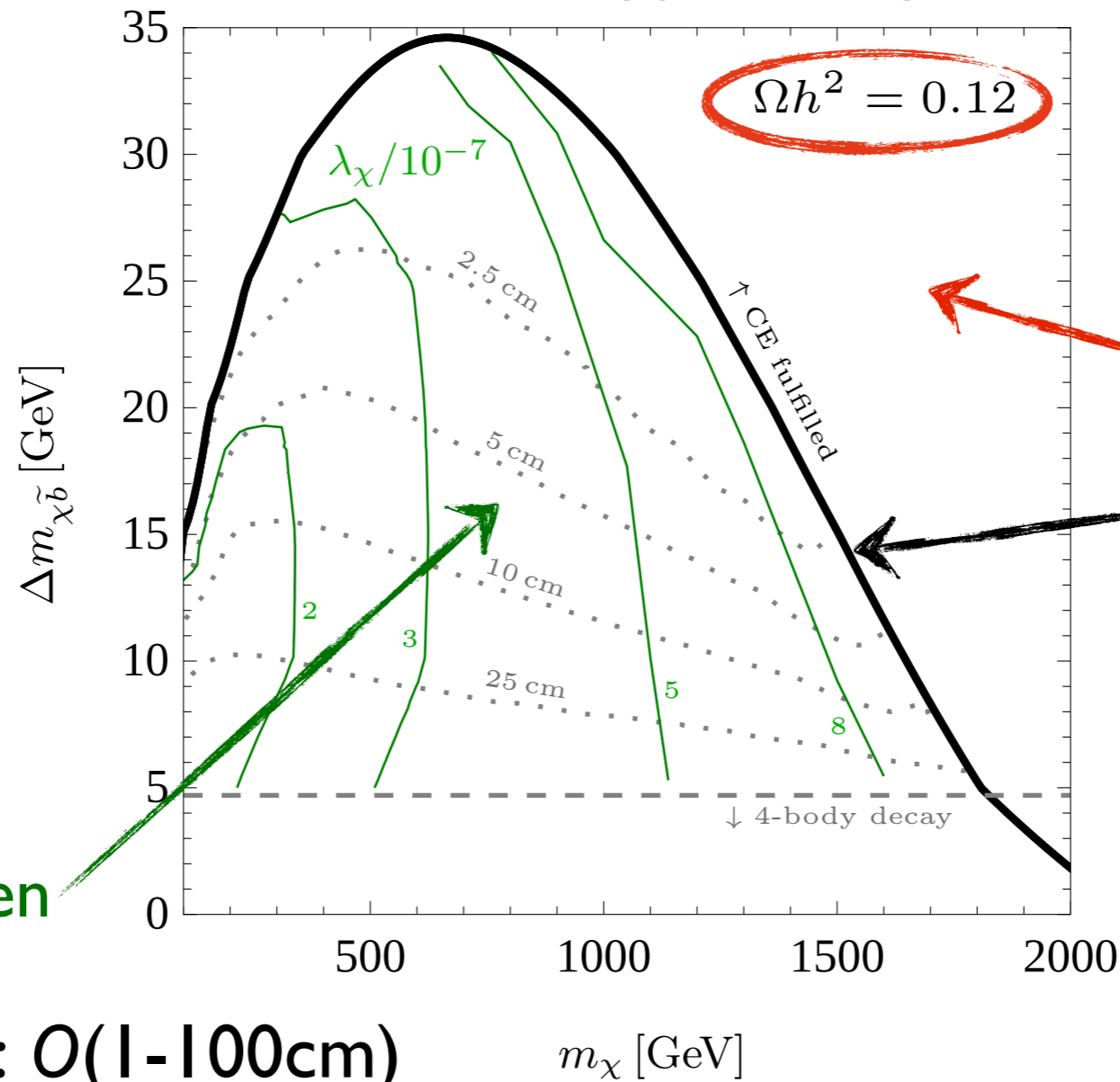


λ_χ is a free parameter here [see Ibarra et al. 2009 for SUSY realization]

Allowed parameter space

- Require Planck relic density

[Garny, JH, Lüpf, Vogl 2017]



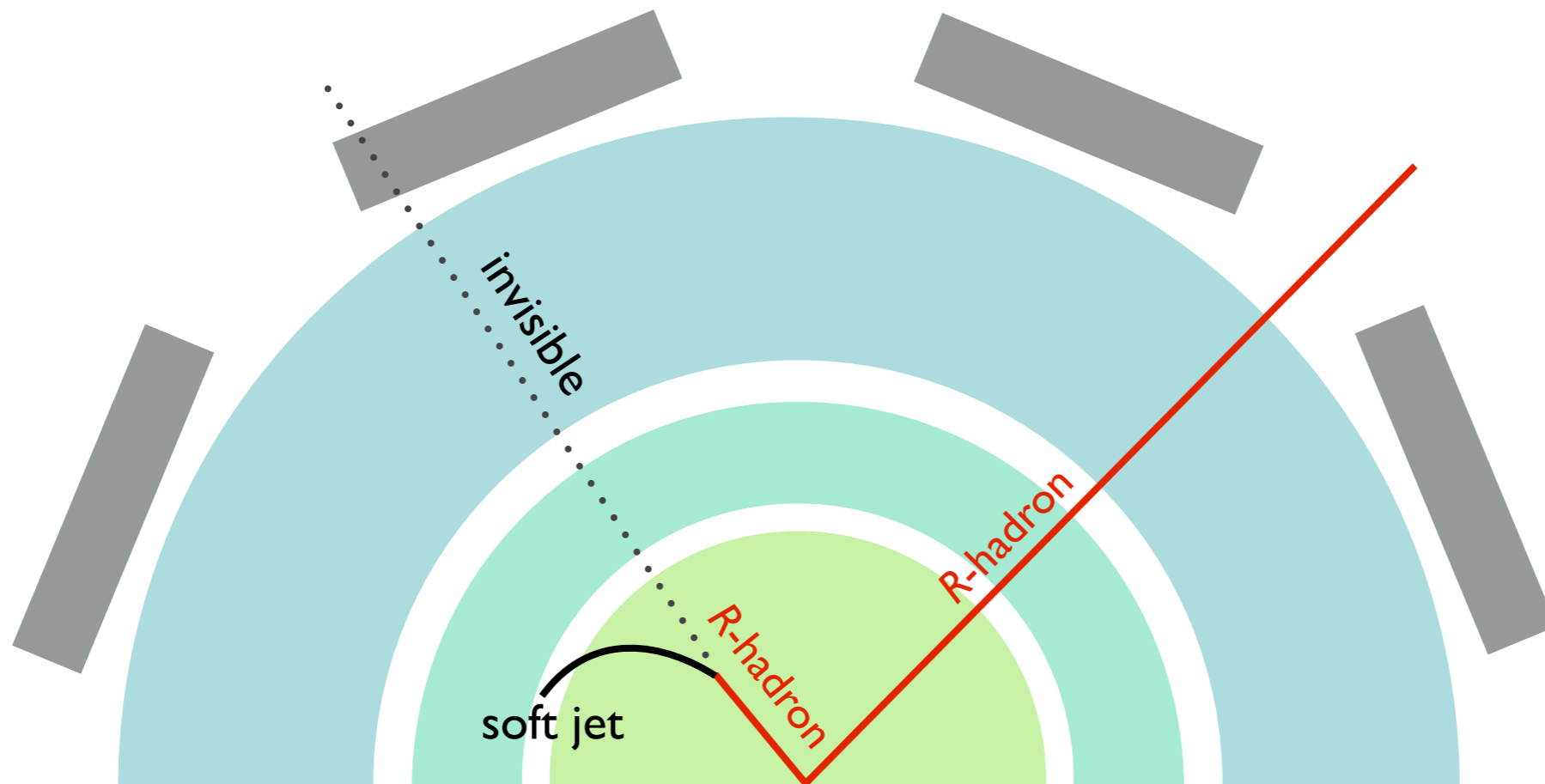
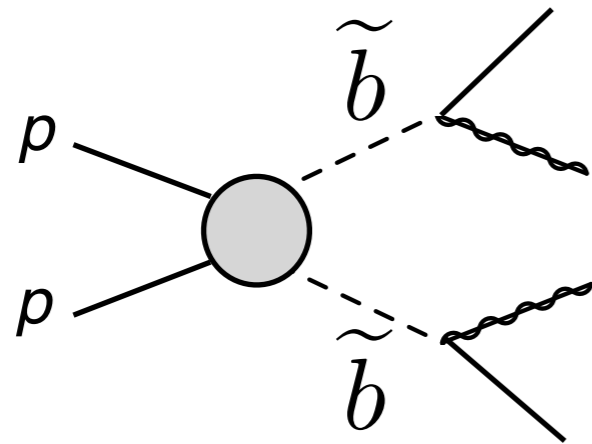
Conversion-driven
freeze-out

\Rightarrow Decay length: $O(1-100\text{cm})$

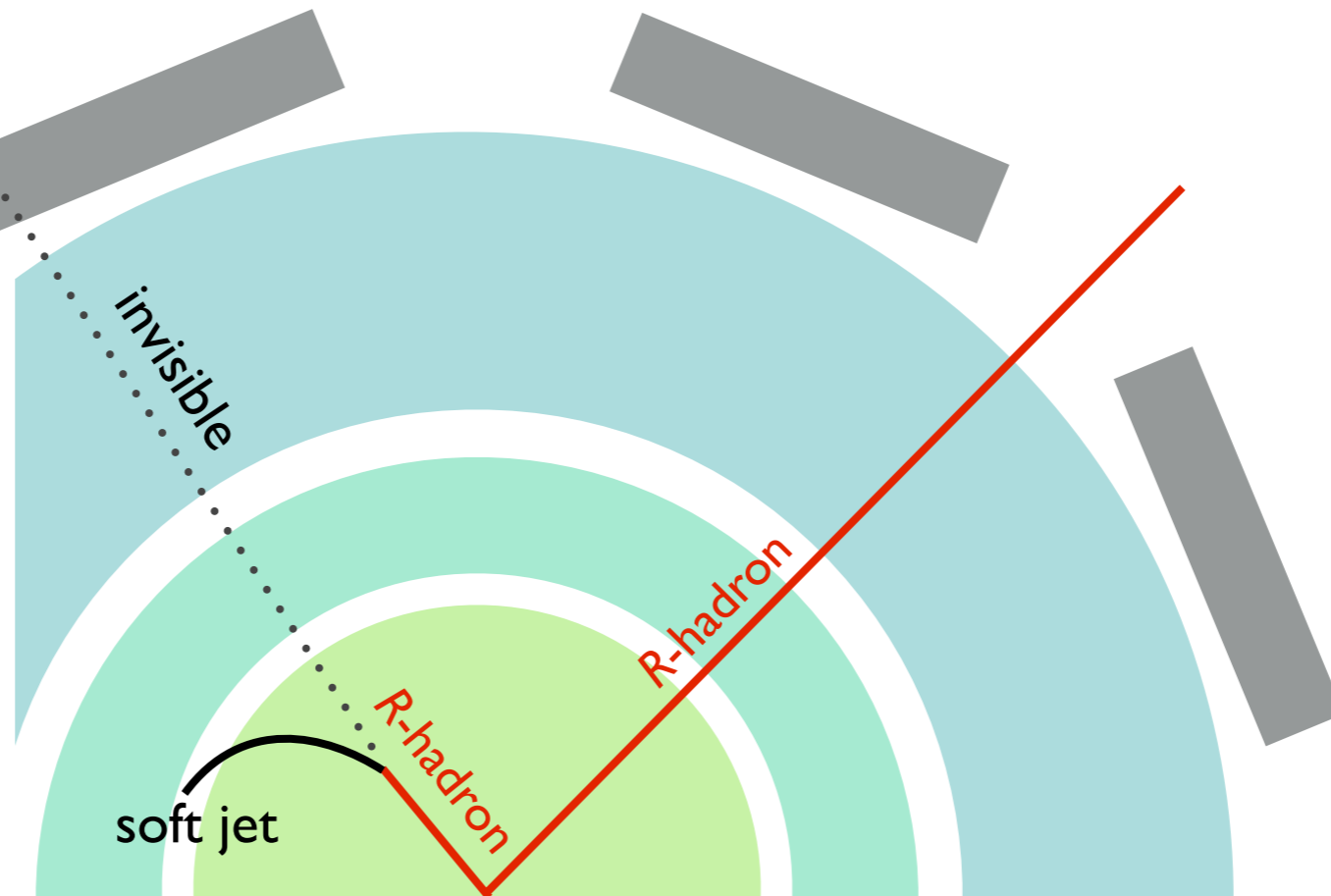
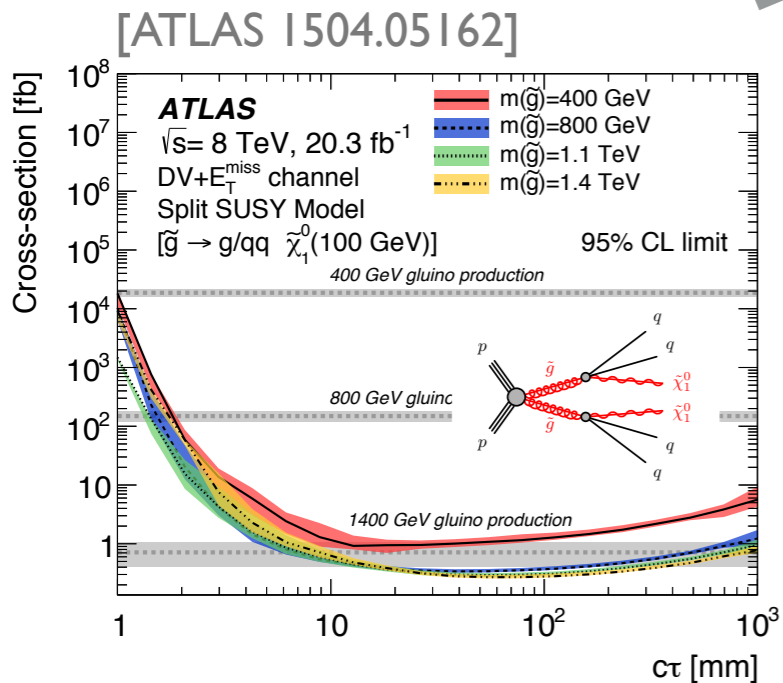
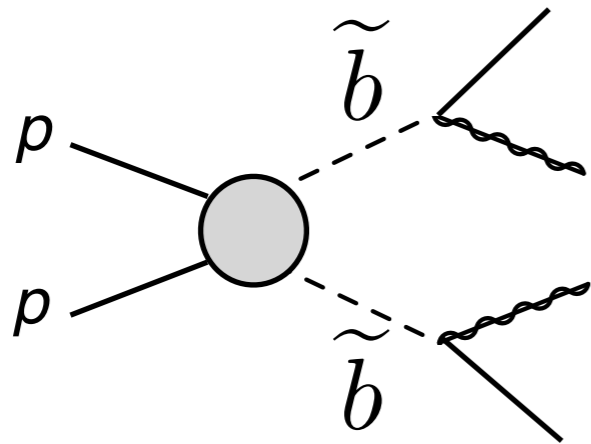
WIMP region

"Freeloader's
freeze-out"

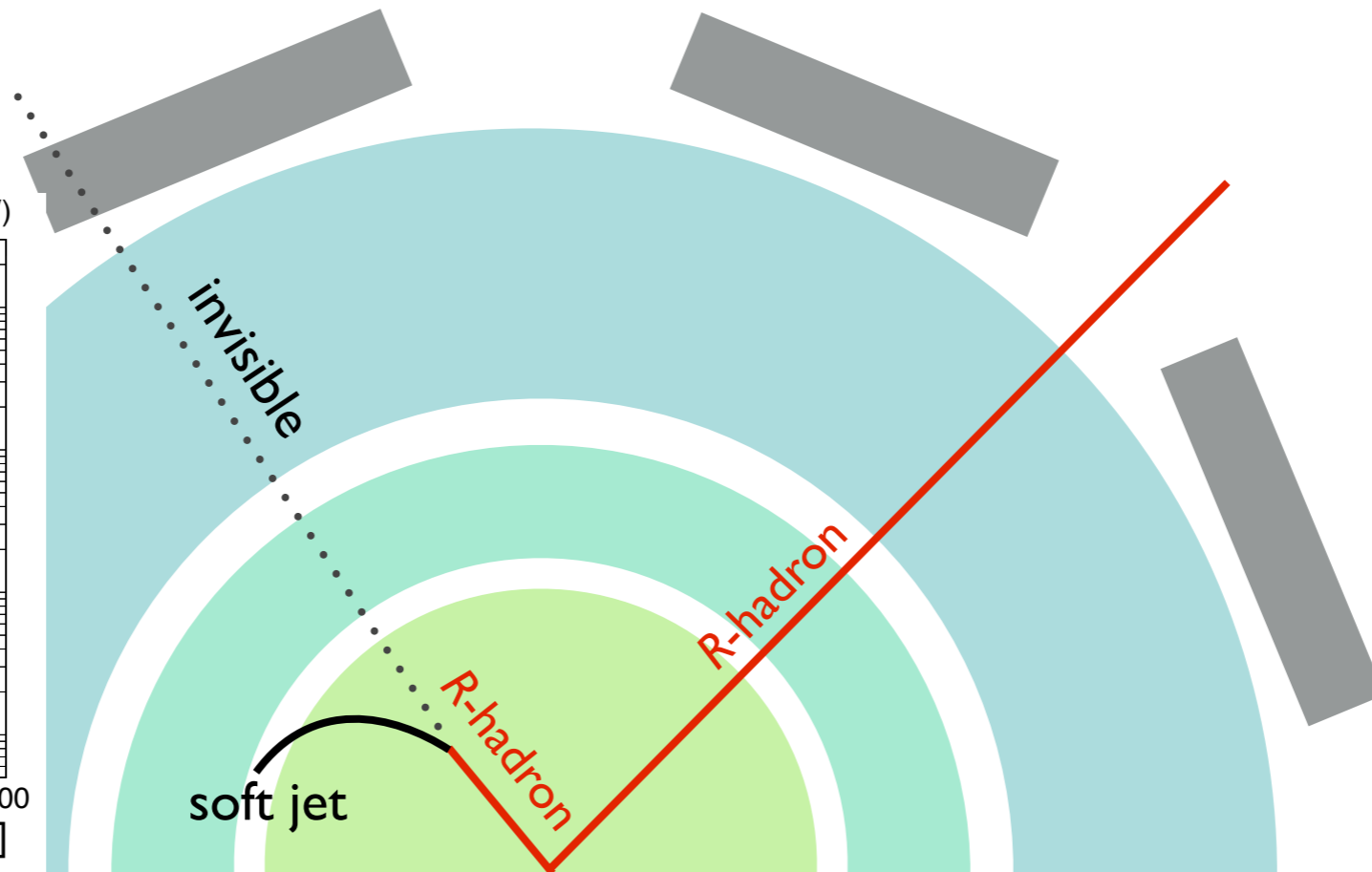
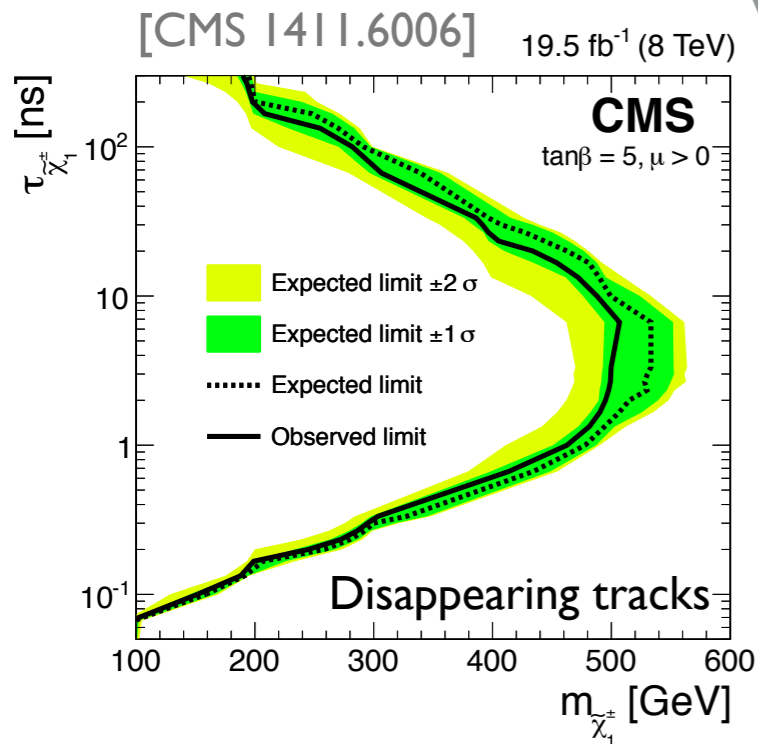
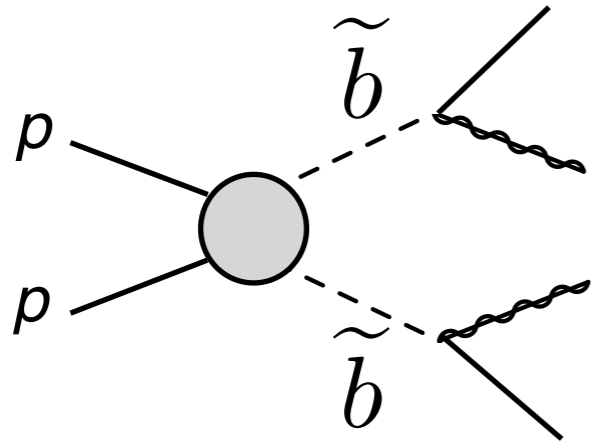
LHC constraints



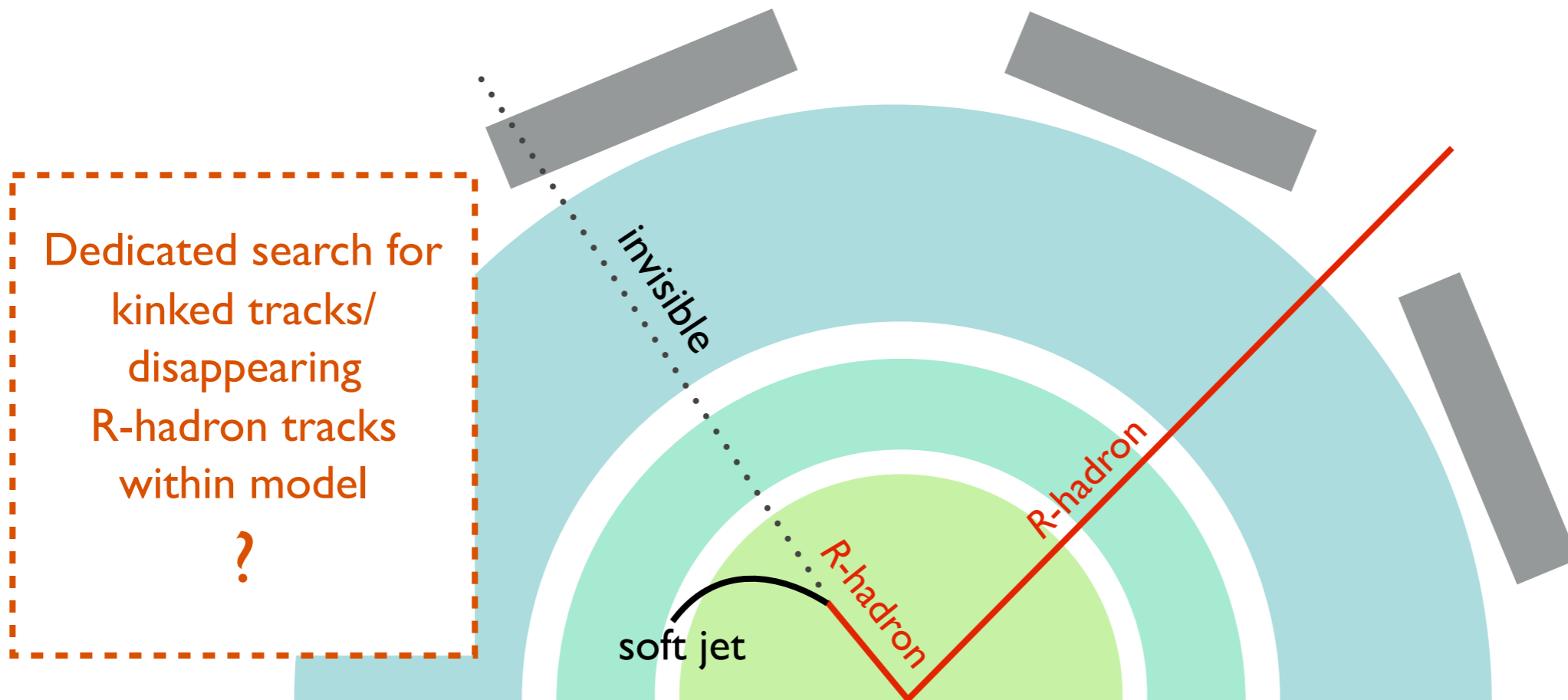
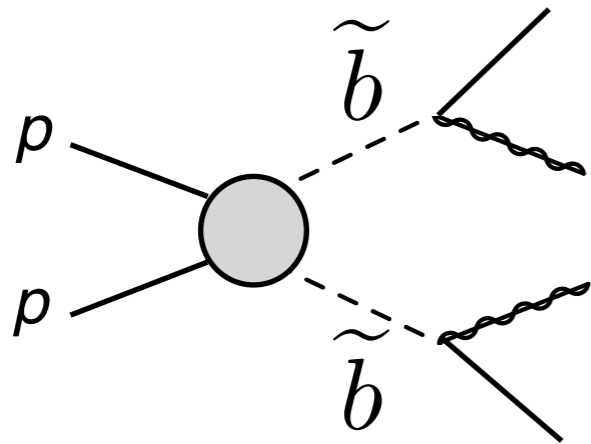
LHC constraints



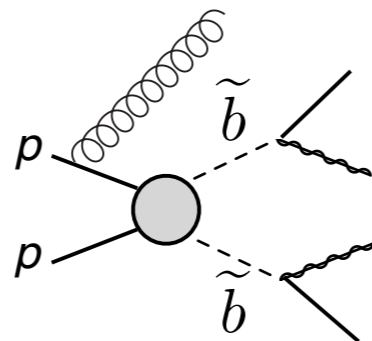
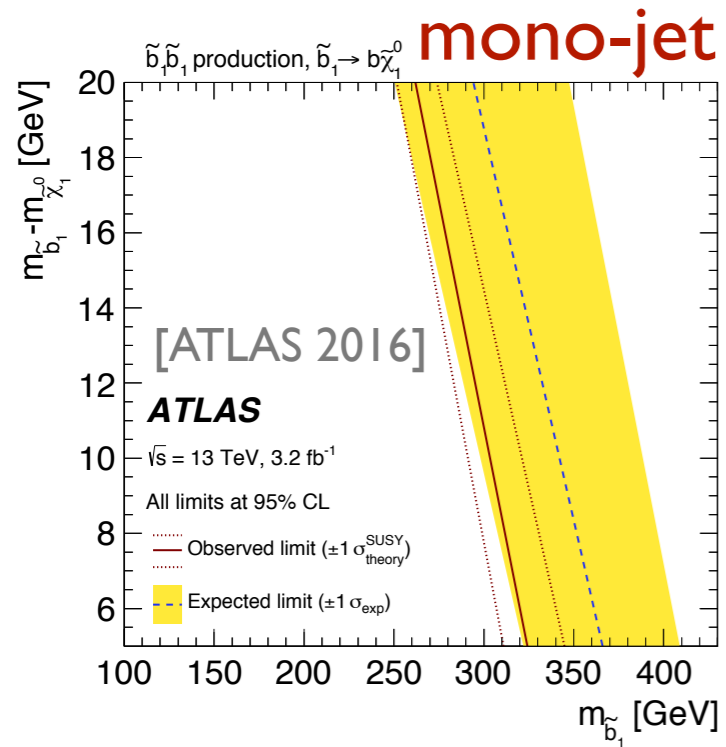
LHC constraints



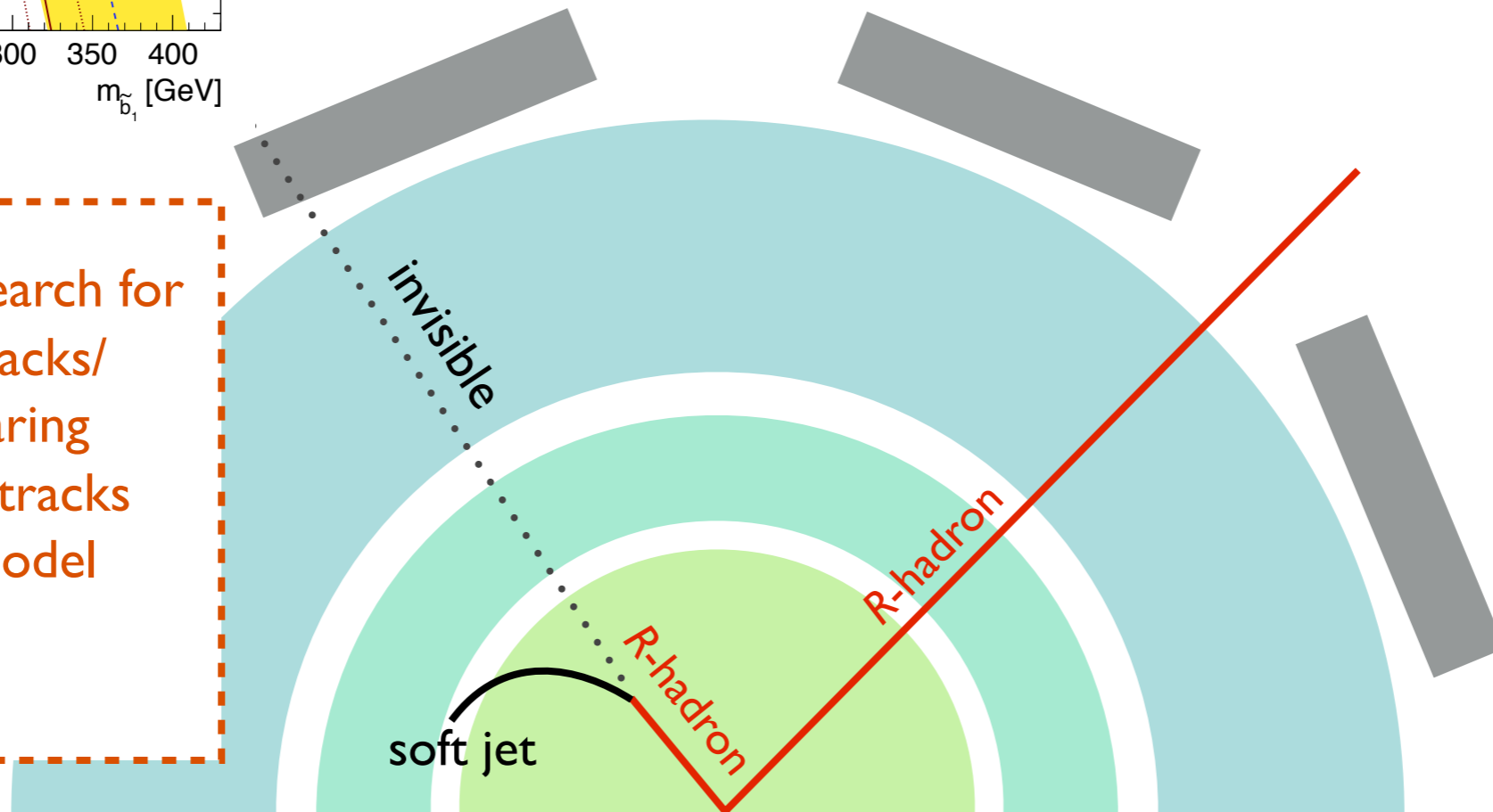
LHC constraints



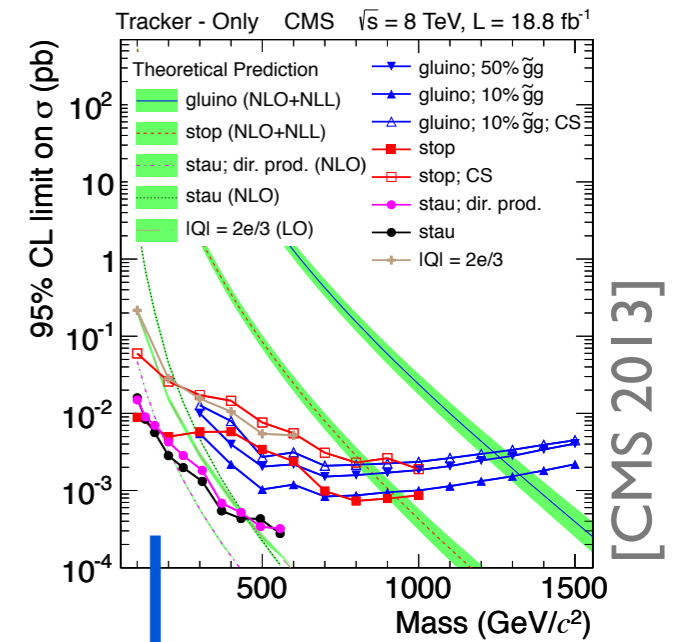
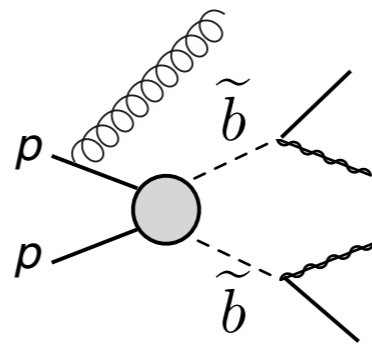
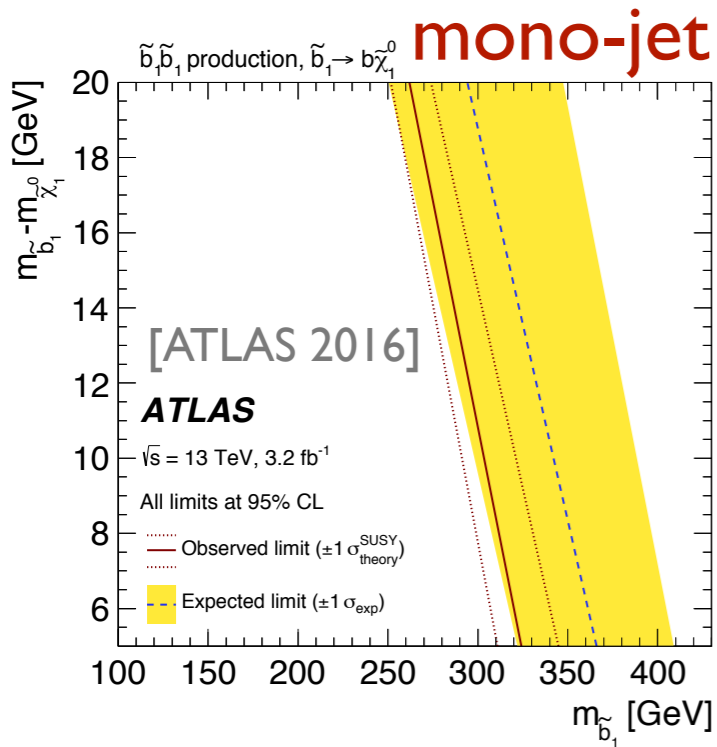
LHC constraints



Dedicated search for
 kinked tracks/
 disappearing
 R-hadron tracks
 within model
 ?

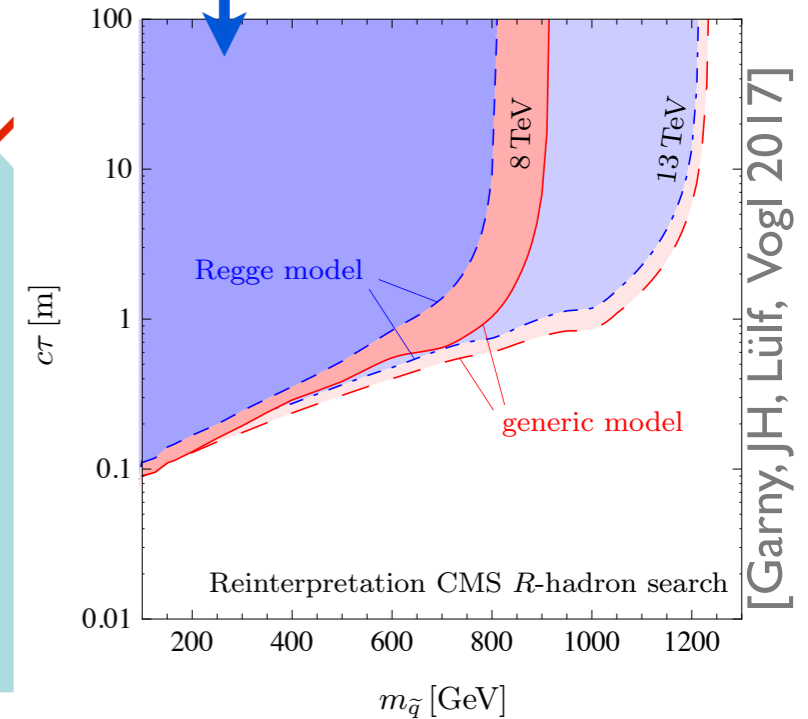
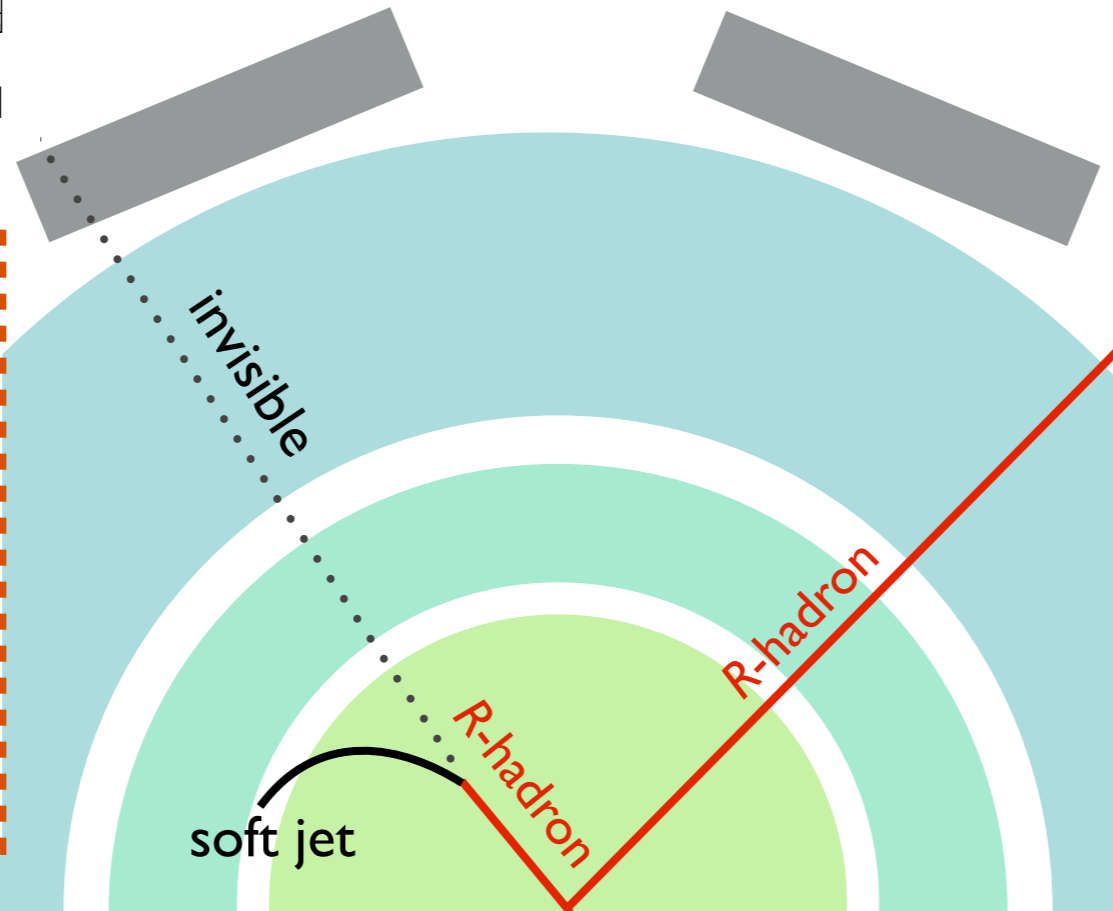


LHC constraints

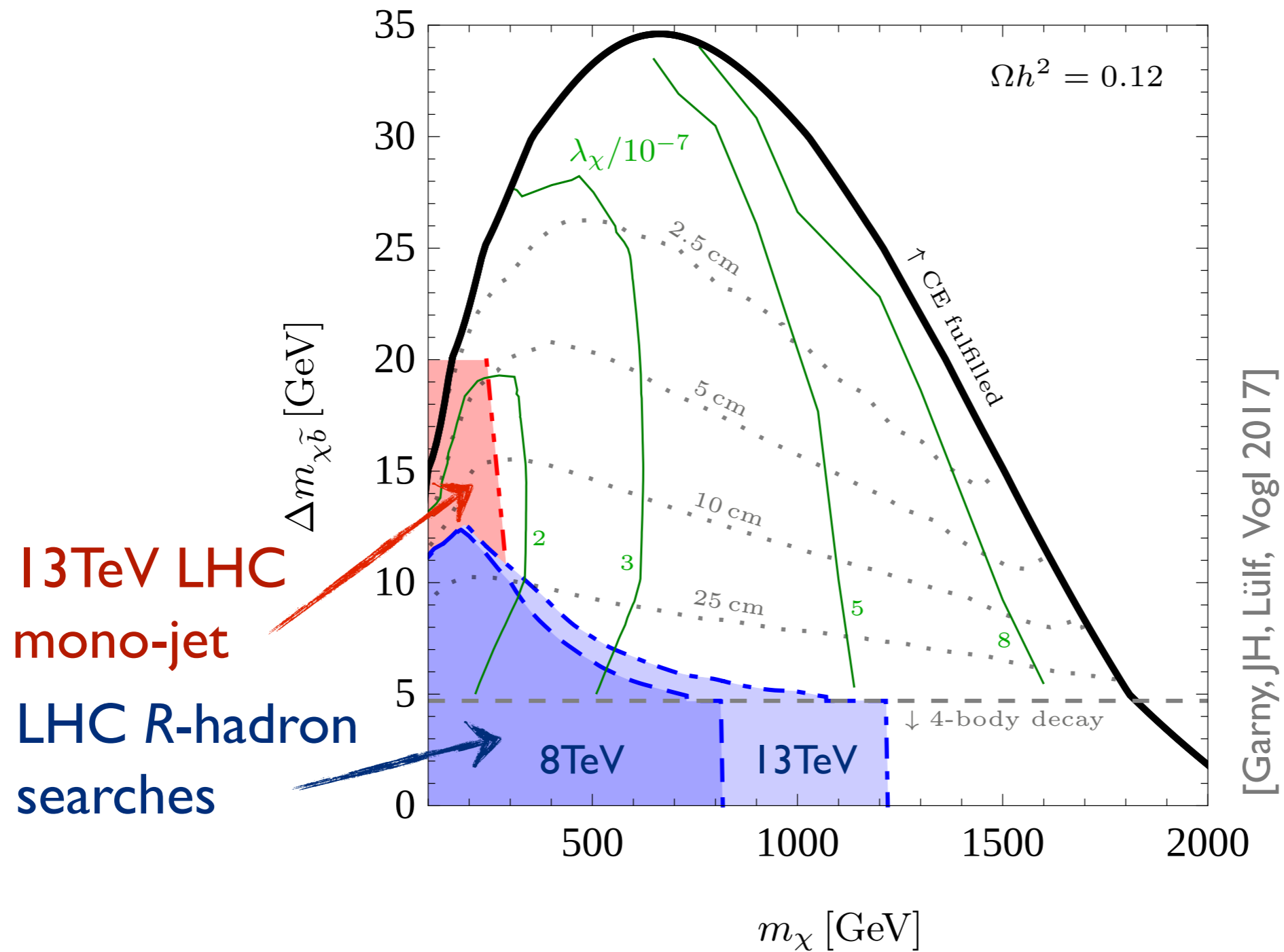


Reinterpretation of R-hadron searches for finite lifetimes

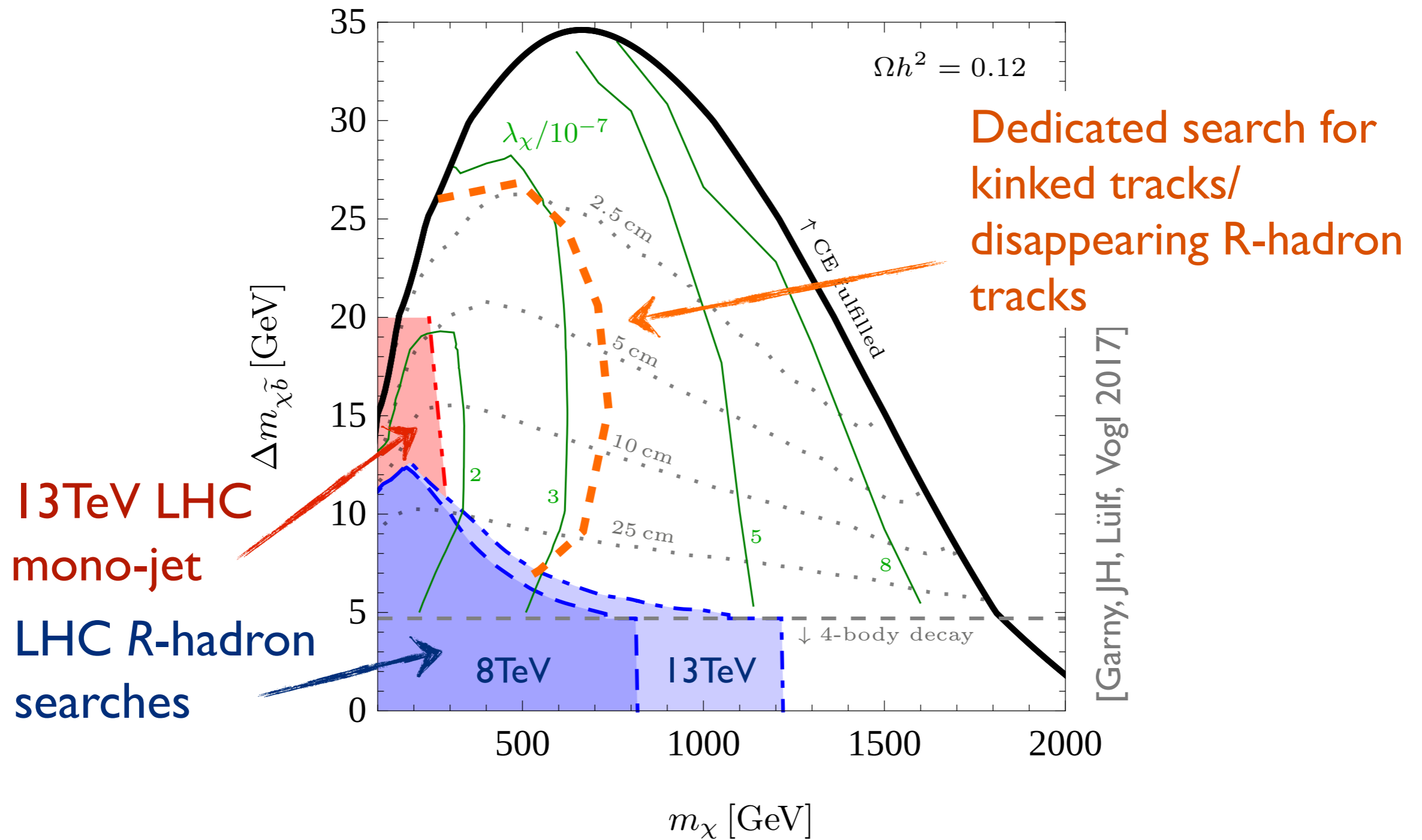
Dedicated search for kinked tracks/ disappearing R-hadron tracks within model ?



Allowed parameter space

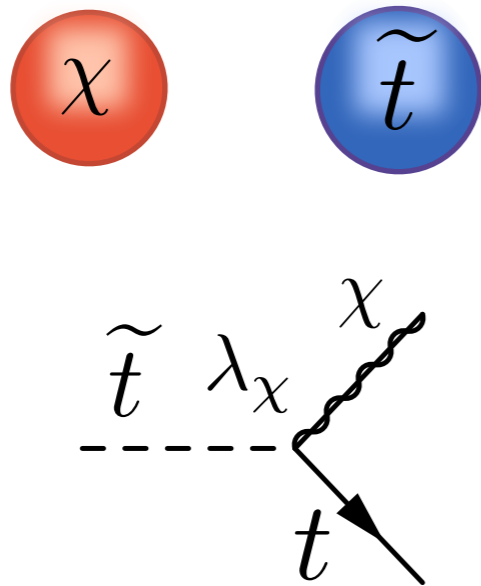


Allowed parameter space



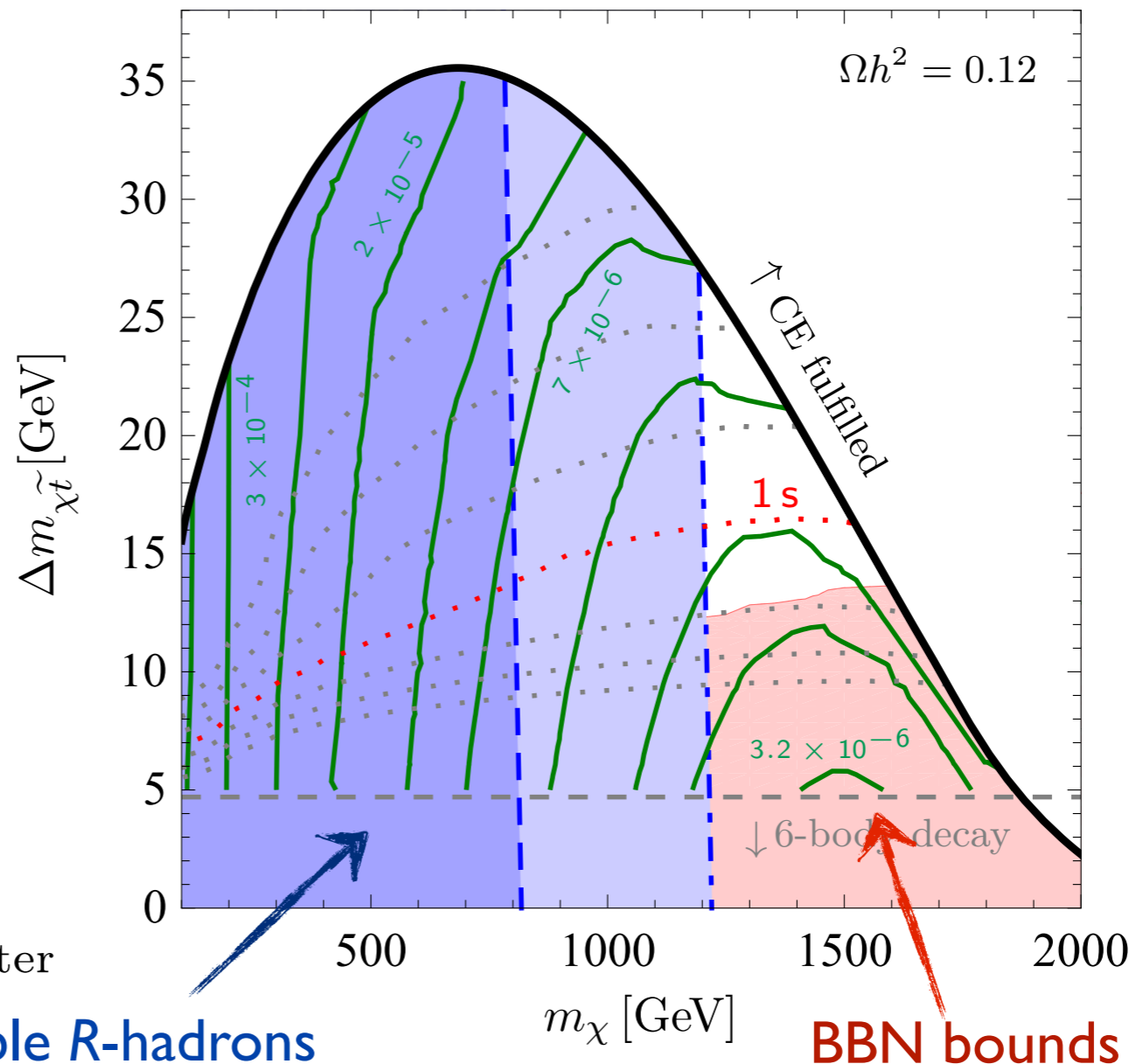
Allowed parameter space: top-partner model

[Garny, JH, Hufnagel, Lulf 1802.00814]



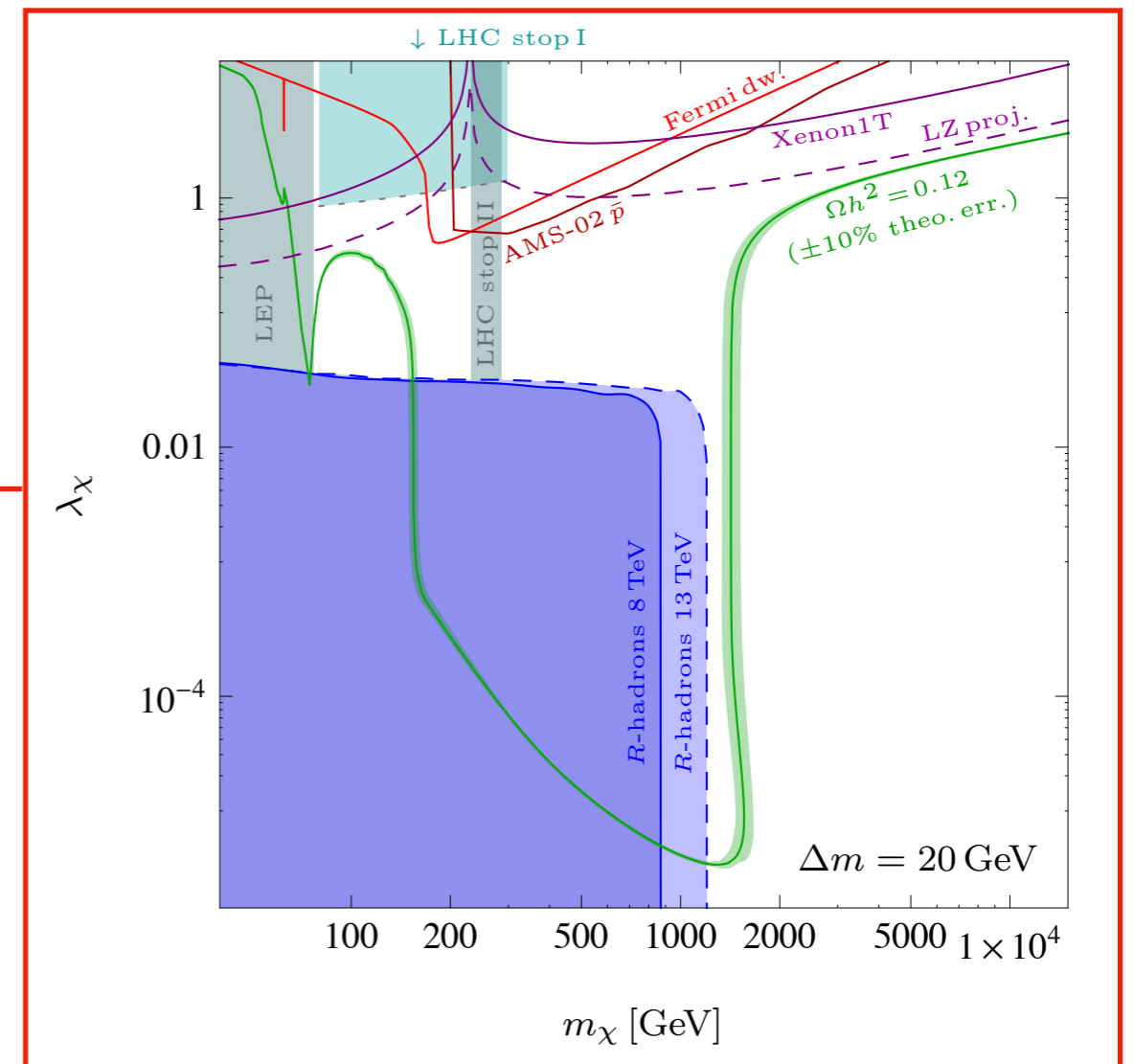
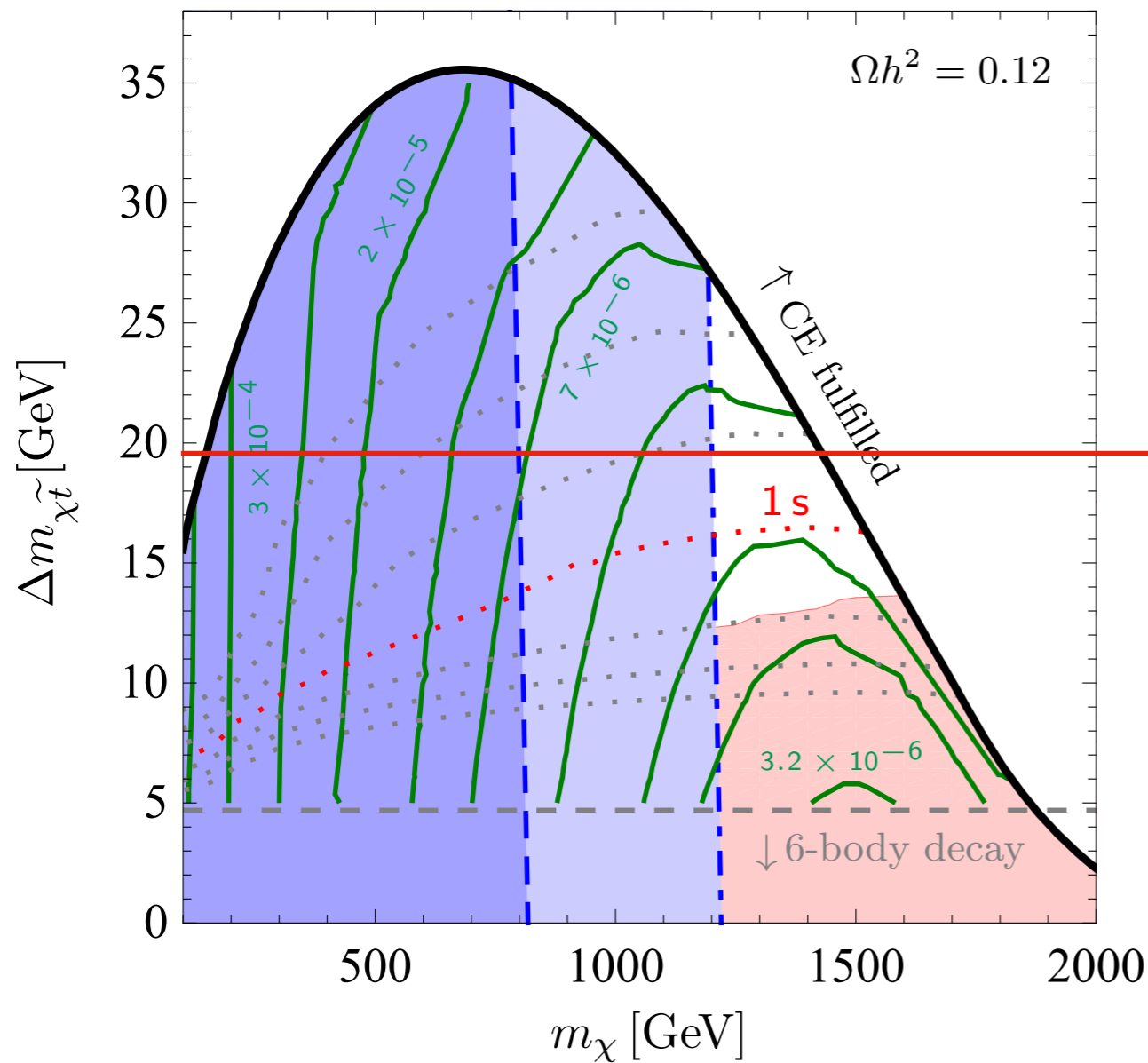
$$\Gamma_{\text{decay}} \ll \Gamma_{\text{scatter}}$$

\Rightarrow **Detector-stable R-hadrons**



Allowed parameter space: top-partner model

[Garny, JH, Hufnagel, Lulf 1802.00814]



Summary

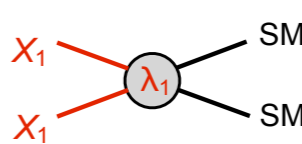
- Vanilla WIMP strongly constrained: Watch out for new avenues beyond WIMPs
 - Variety long-lived particle signatures: Exploit LHC potential
 - Coannihilation with small mass splitting or couplings
 - Coincidence: weak scale decays \sim detector size
 - Conversion-driven freeze-out:
 - Shares nice features of WIMPs!
 - Accommodates null-results from WIMP-searches
 - Dedicated searches needed
-

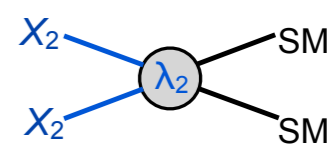
Thanks for your attention!

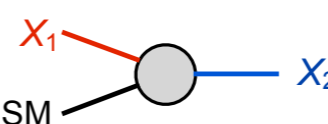
Backup slides

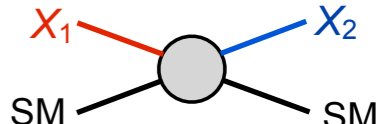
Coupled set of Boltzmann equations

$$\begin{aligned}
 \frac{dn_i}{dt} + 3Hn_i = & - \sum_{j=1}^N \langle \sigma_{ij} v_{ij} \rangle (n_i n_j - n_i^{\text{eq}} n_j^{\text{eq}}) \text{ annihilations} \\
 & - \sum_{j \neq i} [\langle \sigma'_{Xij} v_{ij} \rangle (n_i n_X - n_i^{\text{eq}} n_X^{\text{eq}}) - (i \leftrightarrow j)] \text{ conversions (scattering)} \\
 & - \sum_{j \neq i} [\Gamma_{ij} (n_i - n_i^{\text{eq}}) - (i \leftrightarrow j)] \text{ conversions (decay)}
 \end{aligned}$$





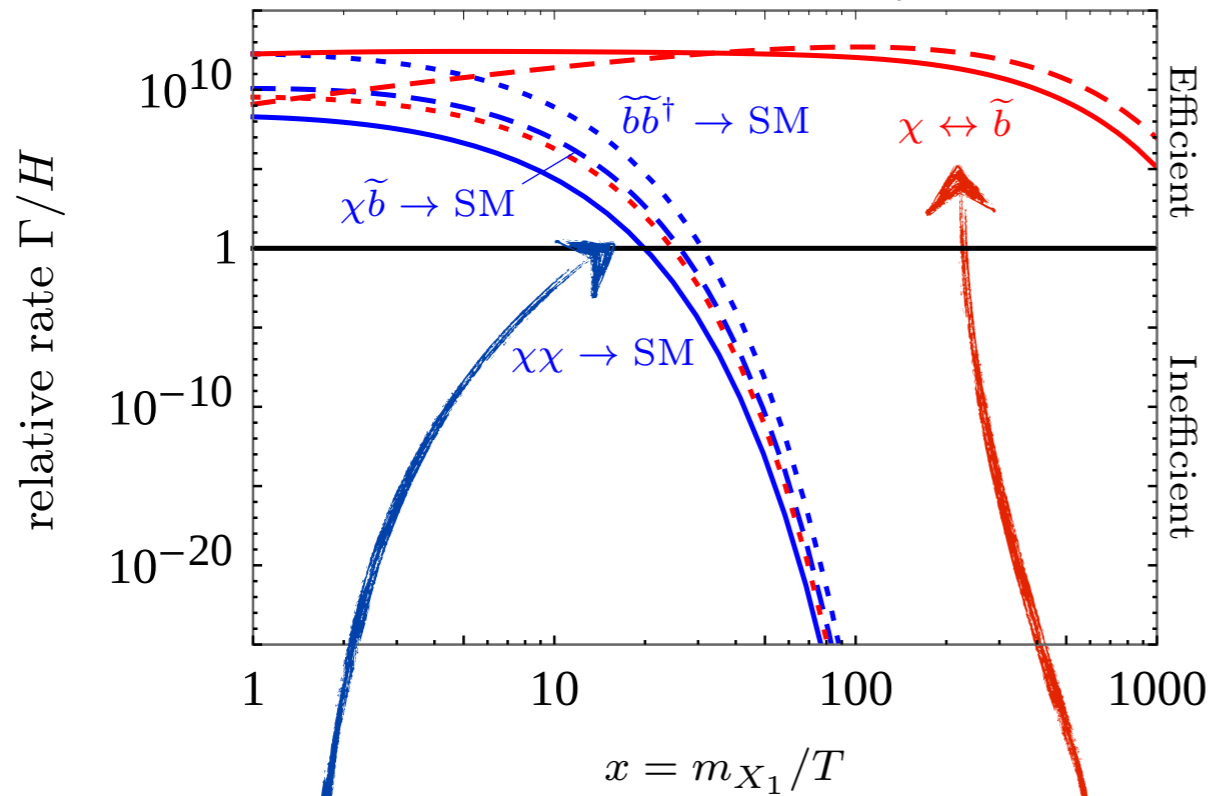




Numerical solution of full coupled system

- SUSY coupling $\lambda_\chi \simeq 0.17$:

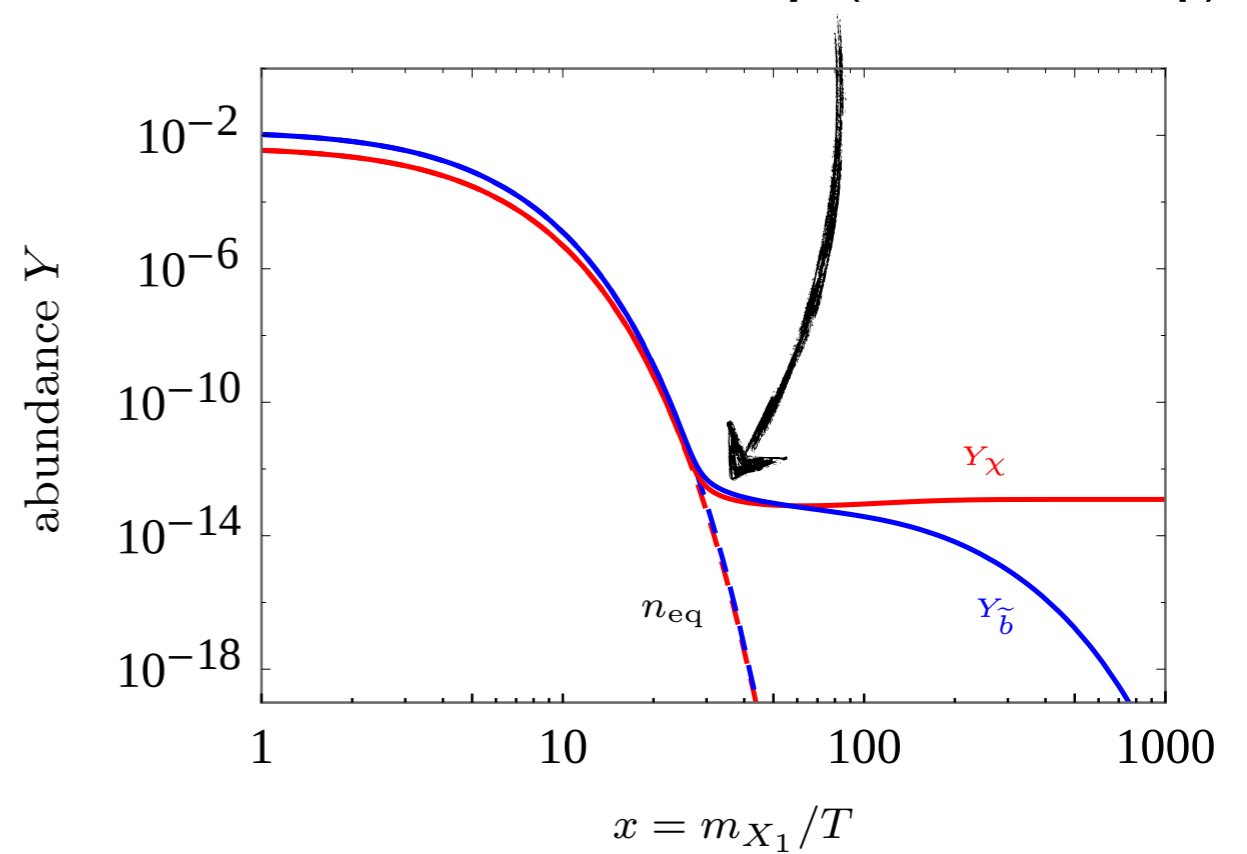
$$m_\chi = 500 \text{ GeV}, m_{\tilde{b}} = 510 \text{ GeV}$$



all annihilations
contribute

conversion thoroughly efficient

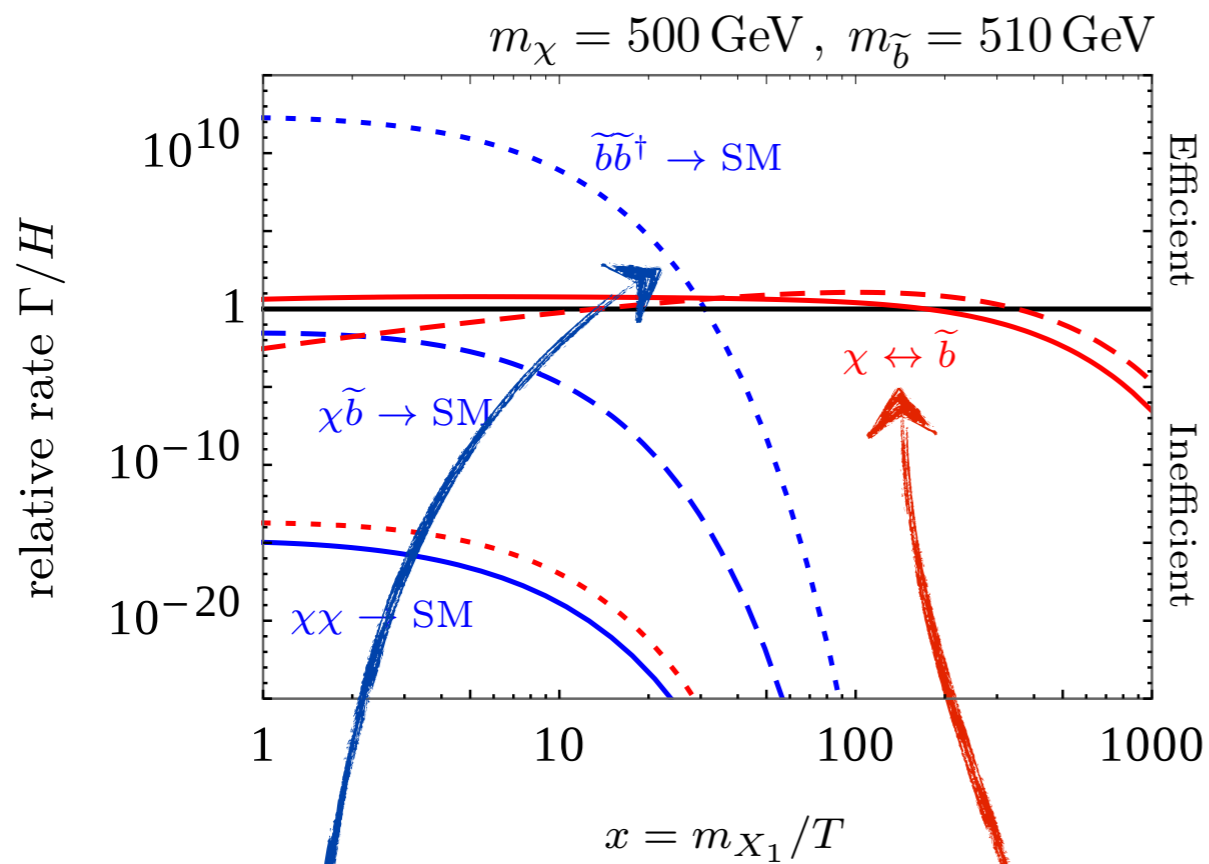
DM and mediator freeze-out
simultaneously (chemical eq.)



Numerical solution of full coupled system

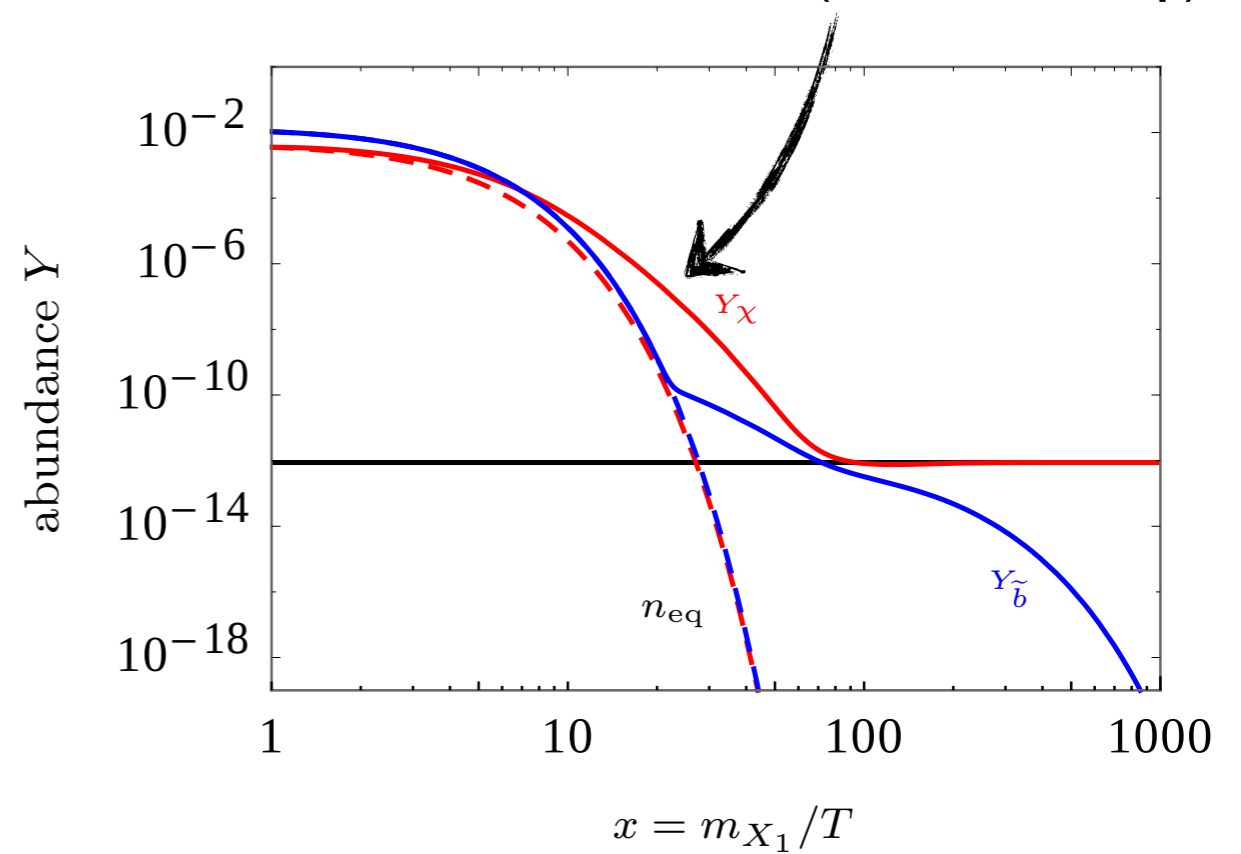
- Very small coupling $\lambda_\chi \simeq 2.6 \times 10^{-7}$:

DM and mediator freeze-out at different x (no chem. eq.)



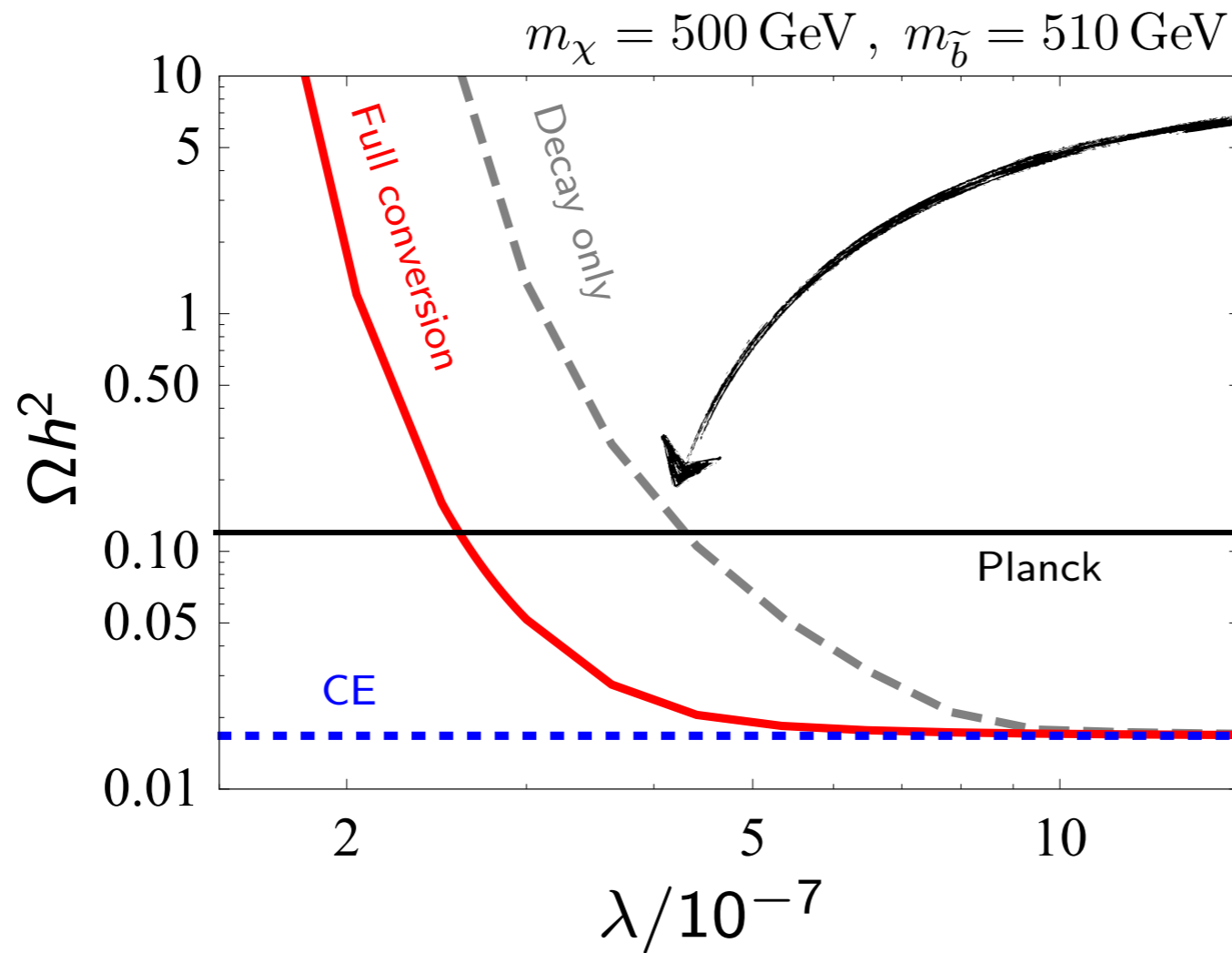
mediator-annihilation
contributes only

conversion on the edge
of being efficient



Numerical solution of full coupled system

- Scan of the coupling:



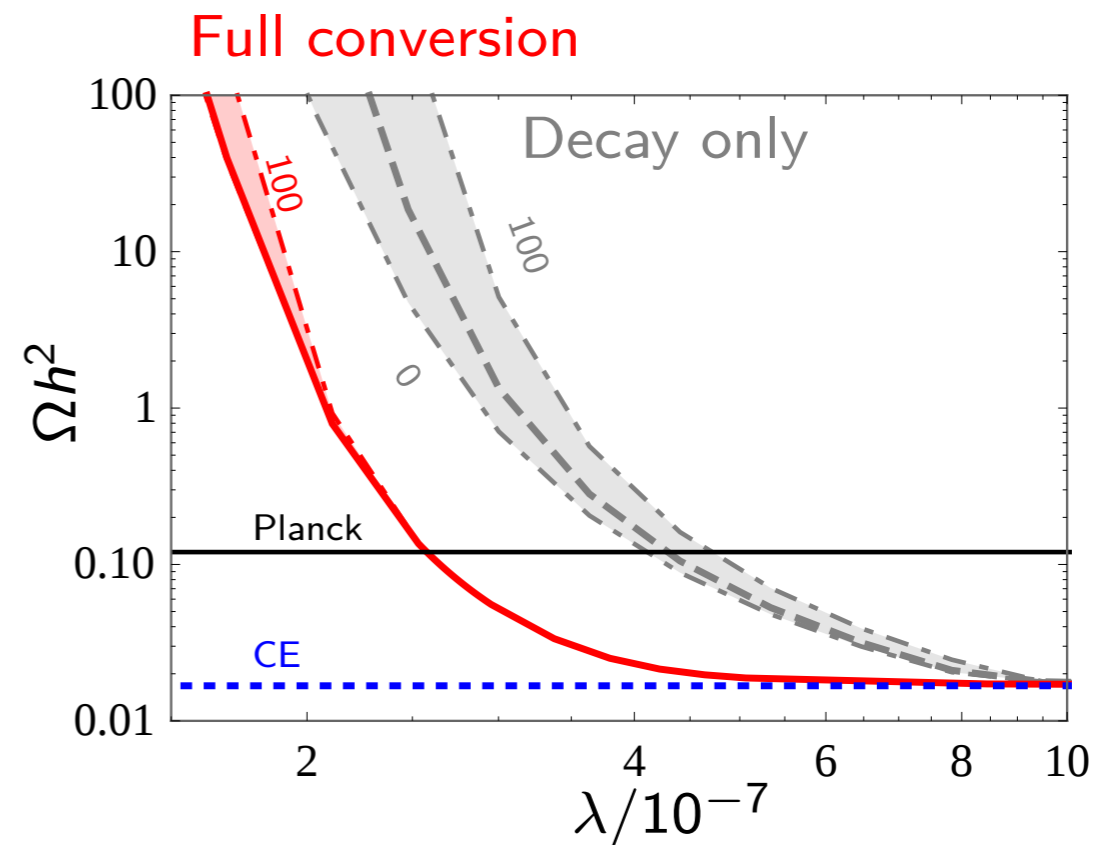
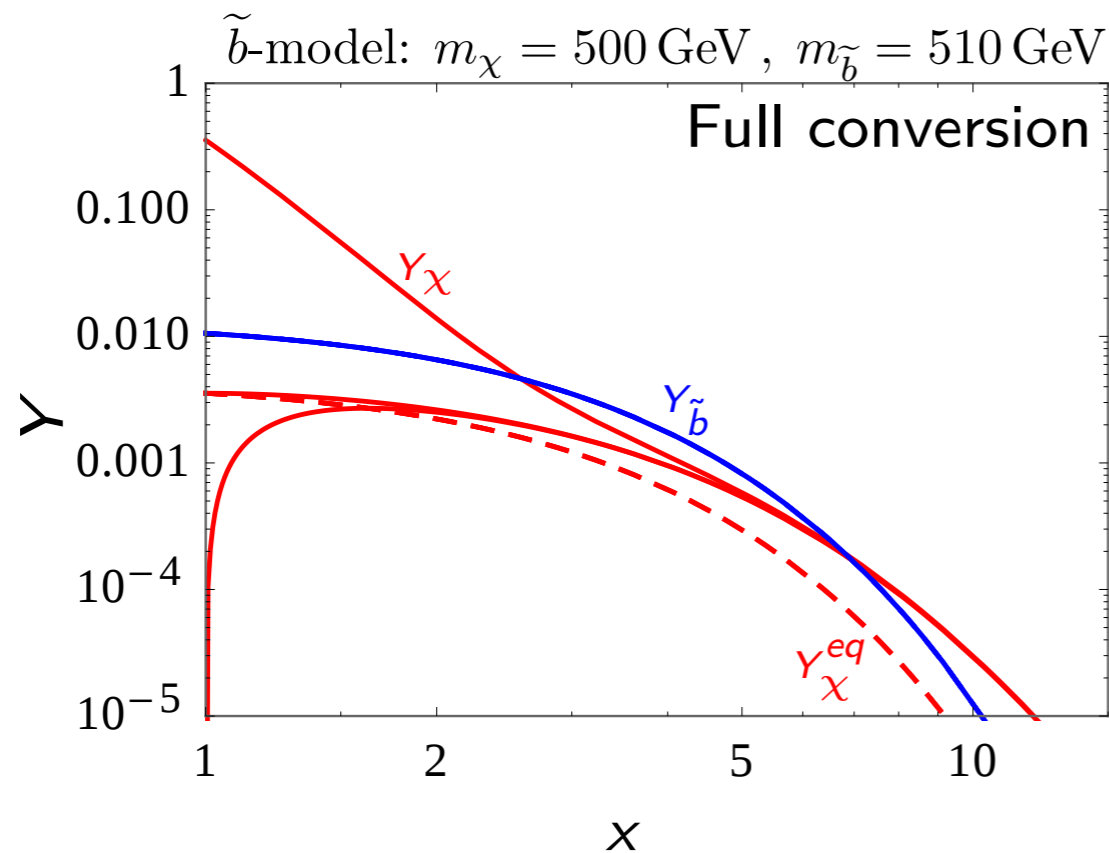
Decay only similar
 $\Gamma_{\text{decay}} \sim \Gamma_{\text{scatter}}$

Scrutinizing some assumptions



Dependence on Initial Conditions

- So far equilibrium density at $x=1$ assumed
- Does DM thermalize?



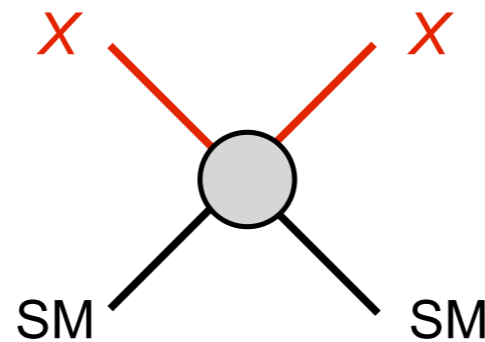
- Insensitive in range $Y_\chi(1) = (0-100) \times Y_\chi^{eq}(1)$
- \Rightarrow Independent of thermal history prior to freeze-out!

Kinetic equilibrium

- Assumption of thermal distributions (via kinetic equilibrium)

$$f_{\chi}(t, p) = f^{\text{eq}}(t, p) \frac{n(t)}{n^{\text{eq}}(t)}$$

- WIMPs: kinetic equilibrium established through efficient elastic scatterings with SM particles:



(kinetic decoupling
takes place well after
freeze-out)

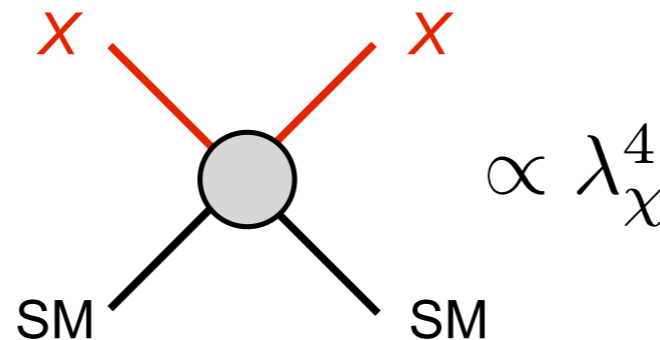
[cf. Chen, Kamionkowski, Zhang 2001,
Bringmann, Hofmann 2006;
Borzumati, Bringmann, Ullio 2007]

Kinetic equilibrium

- Assumption of thermal distributions (via kinetic equilibrium)

$$f_{\chi}(t, p) = f^{\text{eq}}(t, p) \frac{n(t)}{n^{\text{eq}}(t)}$$

- WIMPs: kinetic equilibrium established through efficient elastic scatterings with SM particles:




- Inefficient for DM in conversion-driven freeze-out!
- Mediator is in kinetic equilibrium

Boltzmann equations for particle densities

[Lee, Weinberg 1977; Binetruy, Girardi, Salati 1984; Bernstein, Brown, Feinberg 1985; Srednicki, Watkins, Olive 1988; Kolb, Turner 1990; Griest, Seckel 1991; Gondolo, Gelmini 1991; Edsjo, Gondolo 1997]

DM distribution functions


$$\underline{E_\chi (\partial_t - H p \partial_p) f_\chi(p, t)} = \underline{C [f_\chi]}$$

Relativistic Liouville operator for
homogeneous, isotropic Universe

Collision operator



Cosmology



Particle Physics

Unintegrated Boltzmann equation

- Consider unintegrated Boltzmann equation for χ :

$$Hx\partial_x f_\chi(q, x) = \tilde{C}(q, x) \left(f_\chi^{\text{eq}} \frac{Y_{\tilde{b}}}{Y_{\tilde{b}}^{\text{eq}}} - f_\chi \right)$$

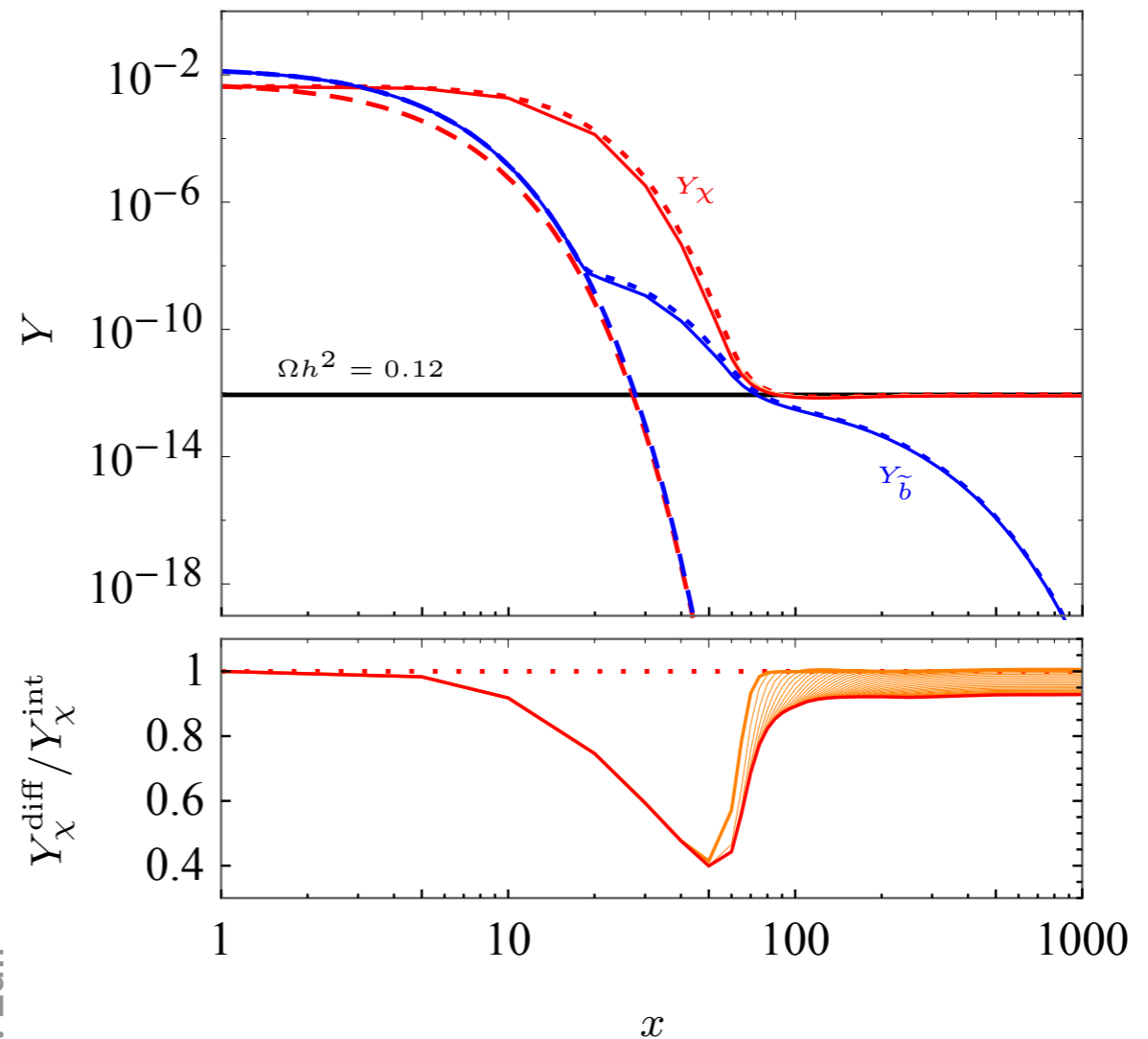
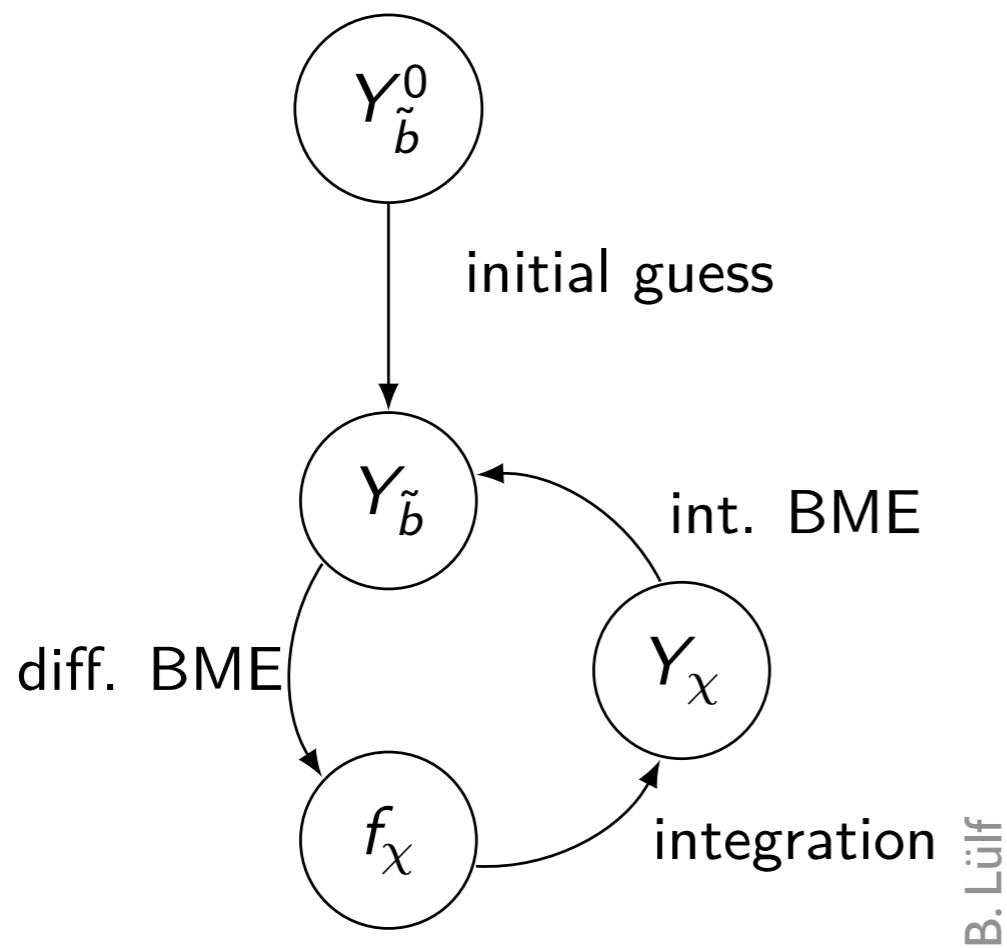
- Conversion only: linear in f_χ
- Can be solved by separation of variables and variation of constants:

$$f_\chi(q, x) = f_\chi^{\text{eq}}(q, x) \frac{Y_{\tilde{b}}}{Y_{\tilde{b}}^{\text{eq}}} - \int_{x_0}^x \frac{d(f_\chi^{\text{eq}}(q, y) Y_{\tilde{b}}(y) / Y_{\tilde{b}}^{\text{eq}}(y))}{dy} \times \exp \left(- \int_y^x \frac{\tilde{C}(q, z)}{zH(z)} dz \right) dy$$

Involves $Y_{\tilde{b}}$ → still coupled system

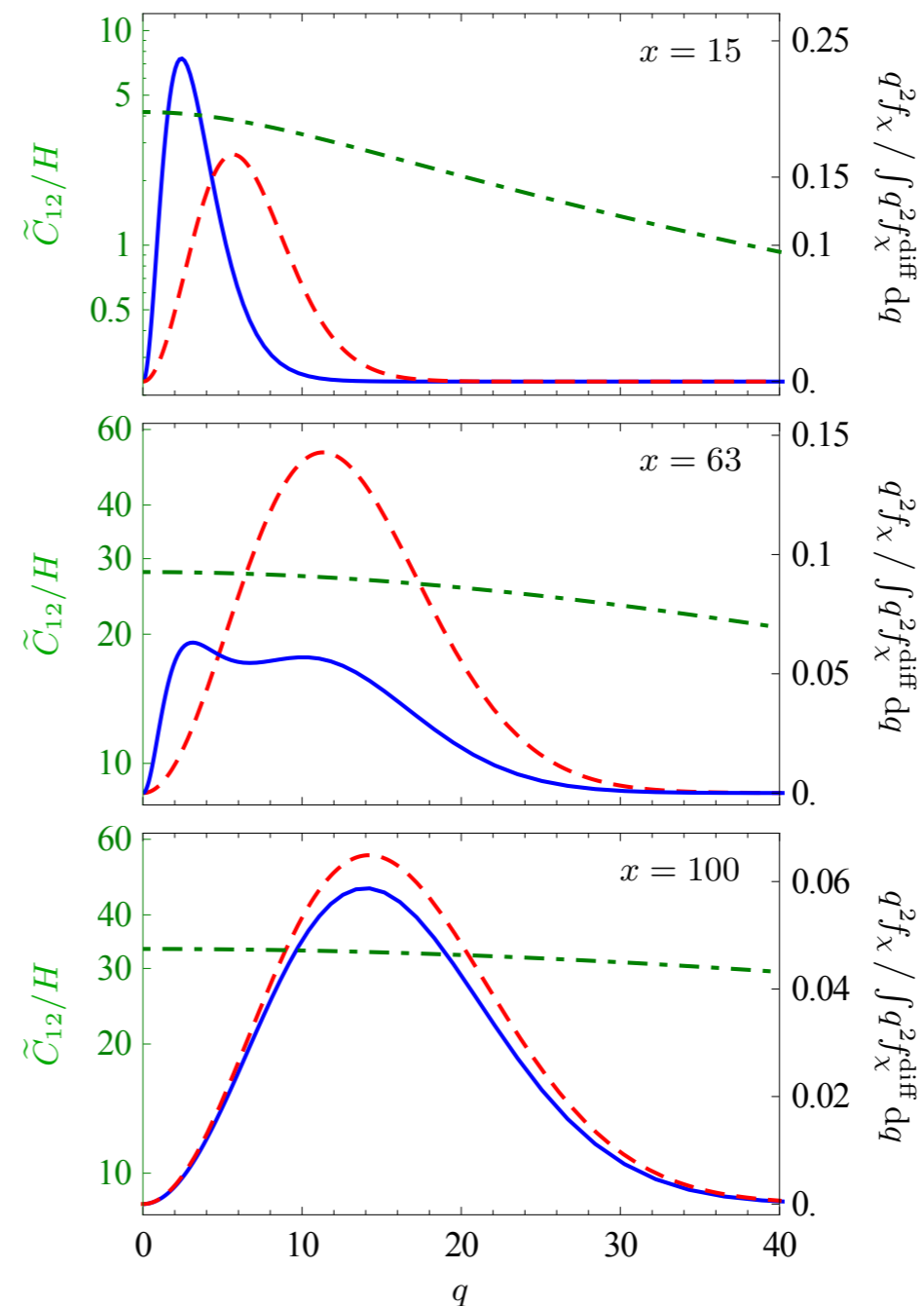
Iterative solution

- Do not solve coupled system at once but iteratively
- Start with "guess" for $Y_{\tilde{b}}$: solution of integrated equations



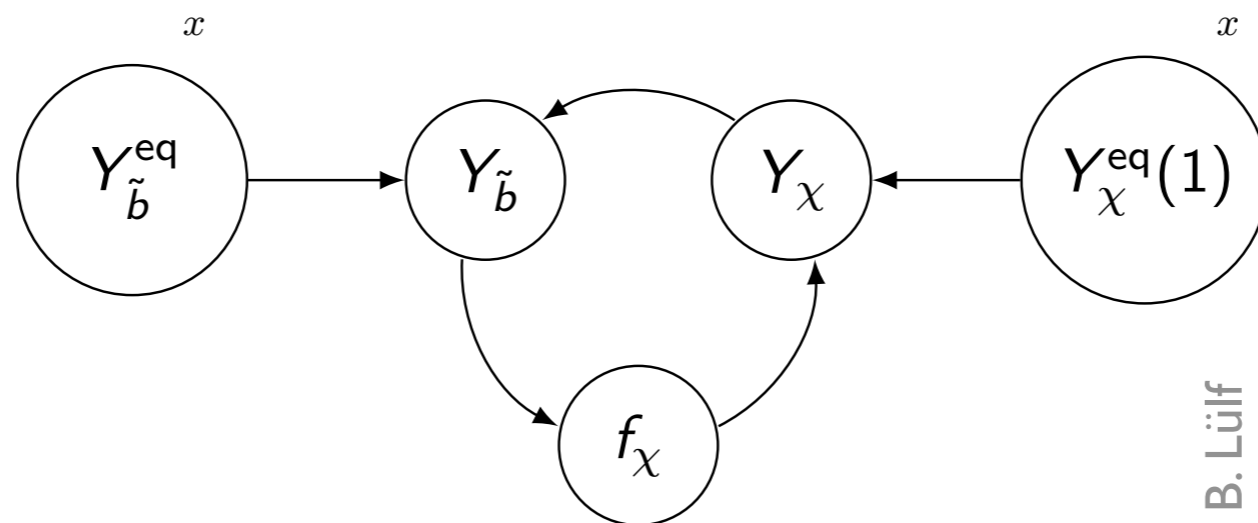
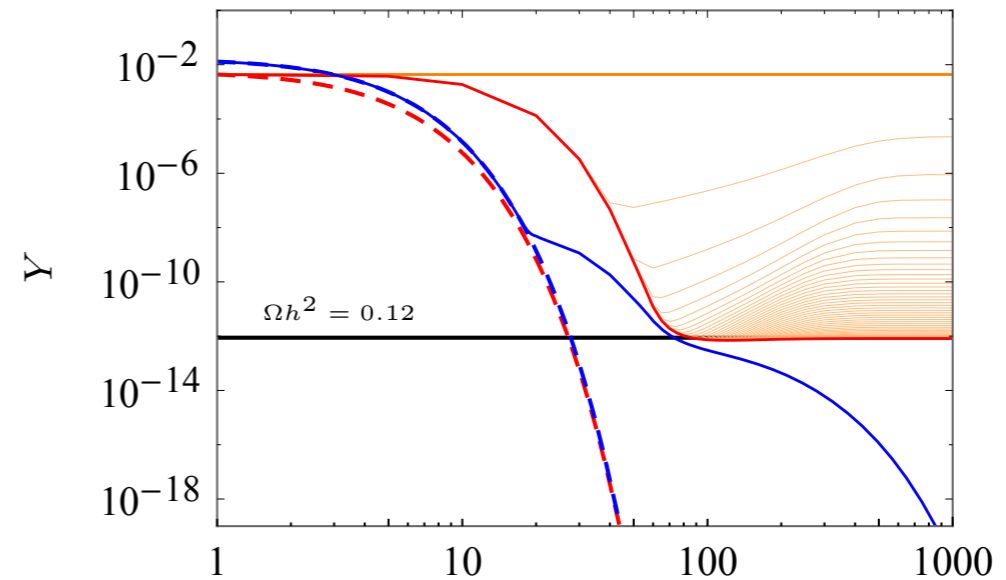
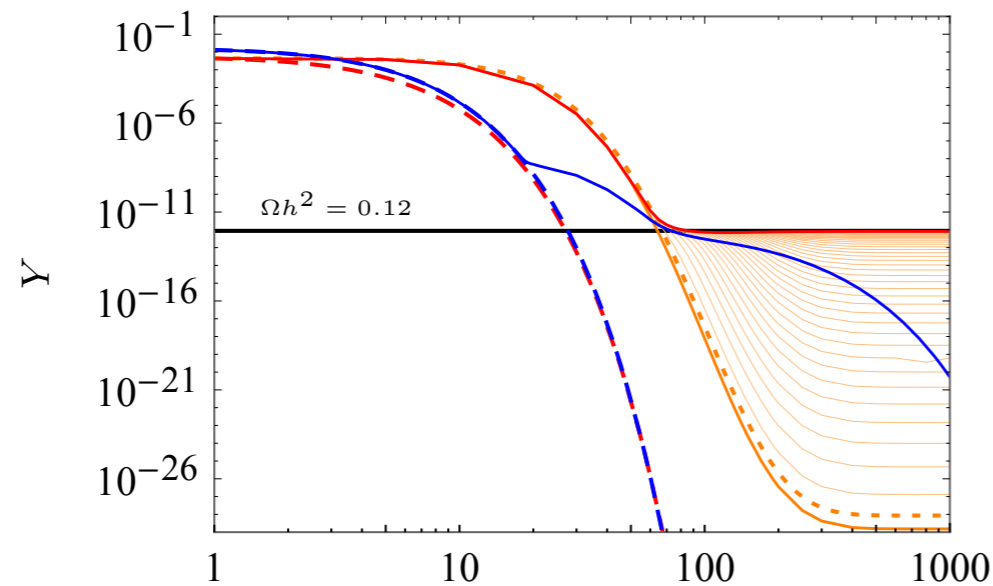
Deviation from thermal distribution

- small x : redshift only
- Conversion inset: thermalization starts
- Close-to-thermal distribution



Testing initial guess

- Extreme cases for initial evolutions of abundances
- Converge to same solution:

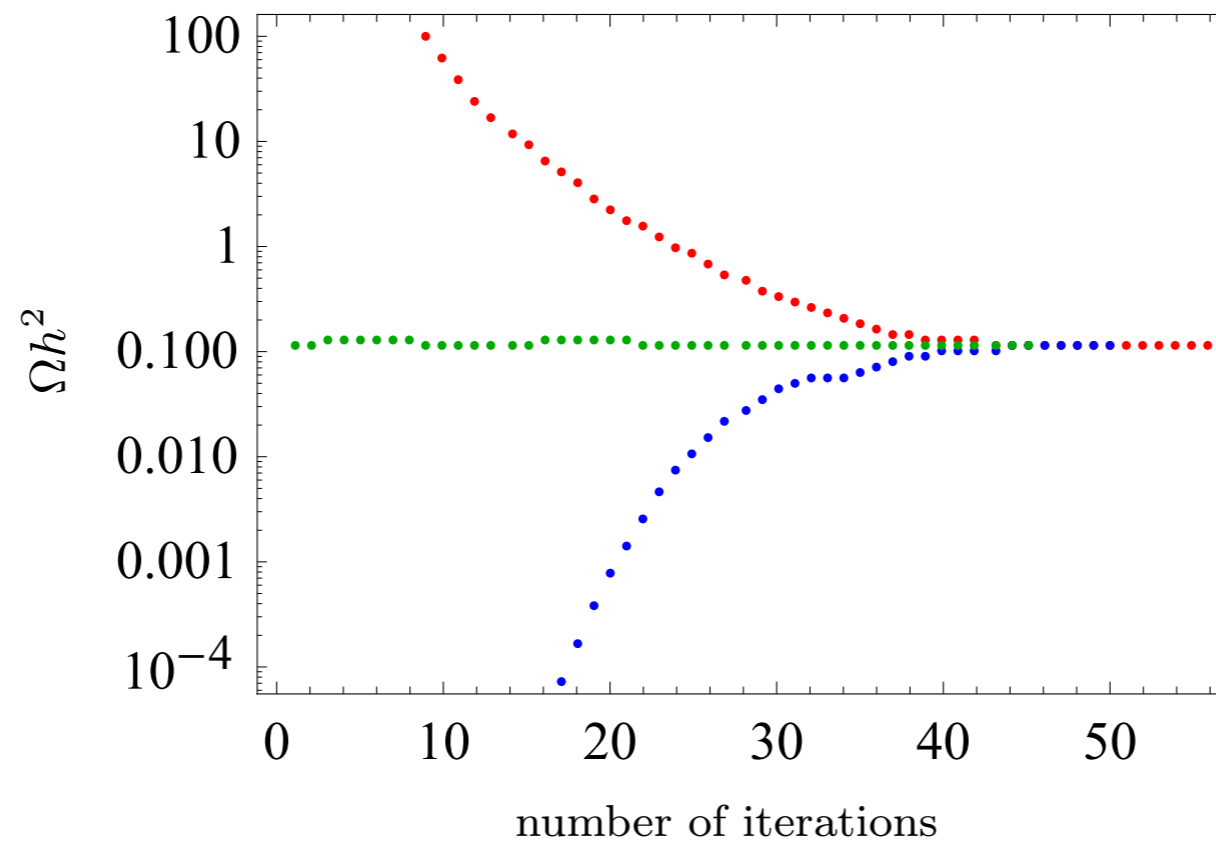


B. Lulf

Iterative solution

- All initial guesses converge to the same solution
- Difference to integrated treatment below 10%
- Solution of coupled system more important

[cf. D'Agnolo, Pappadopulo, Ruderman, 2017]



superWIMP scenario

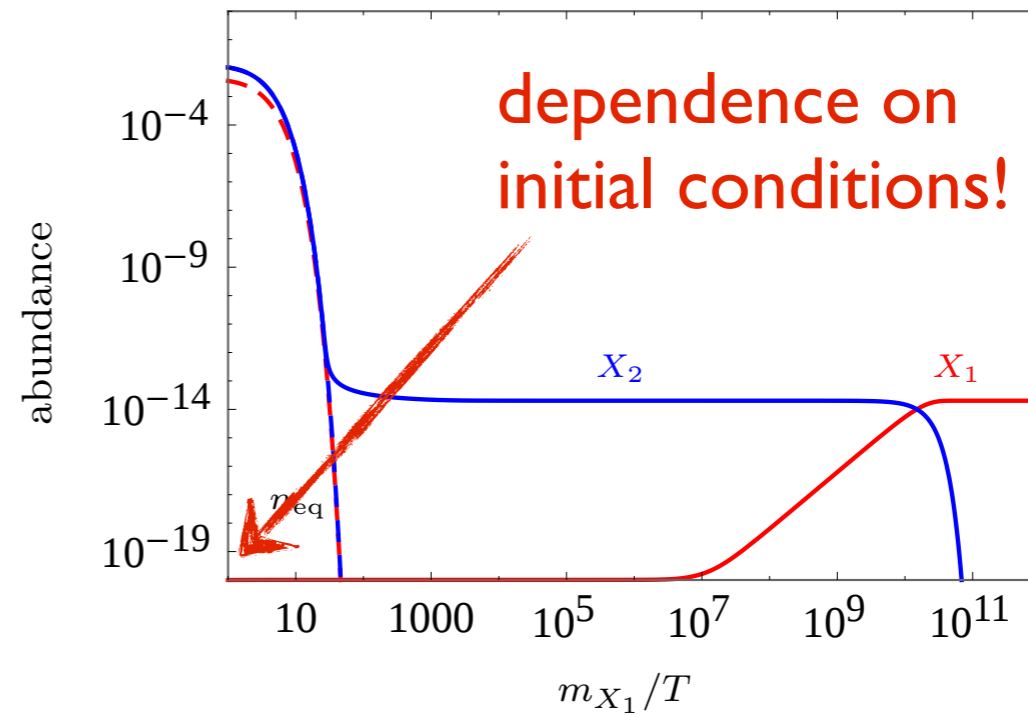
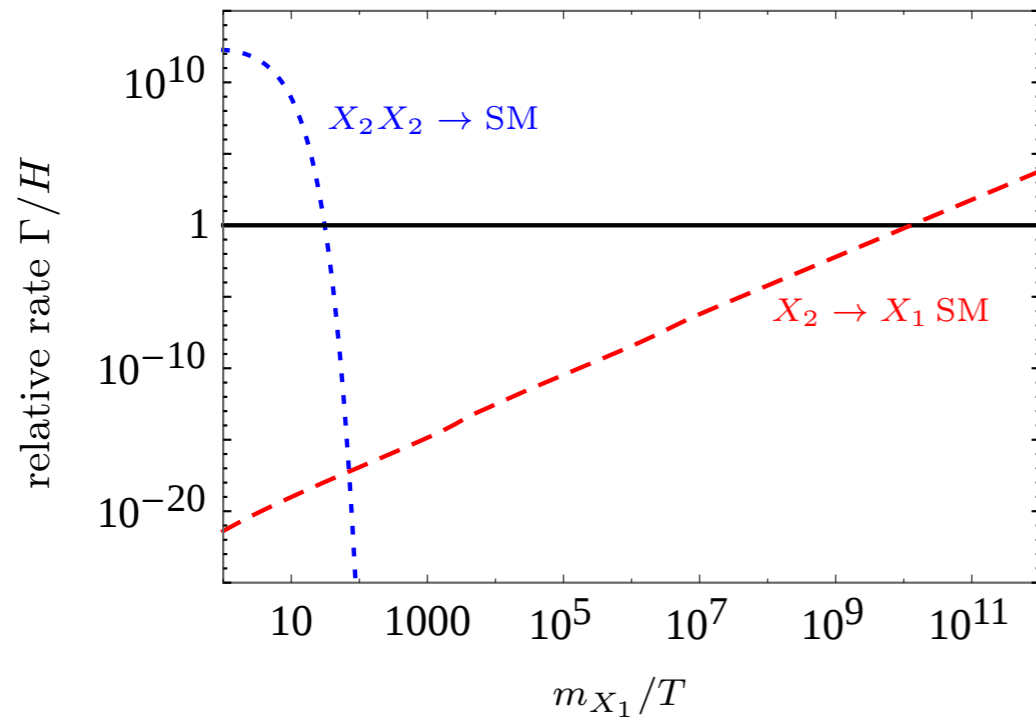
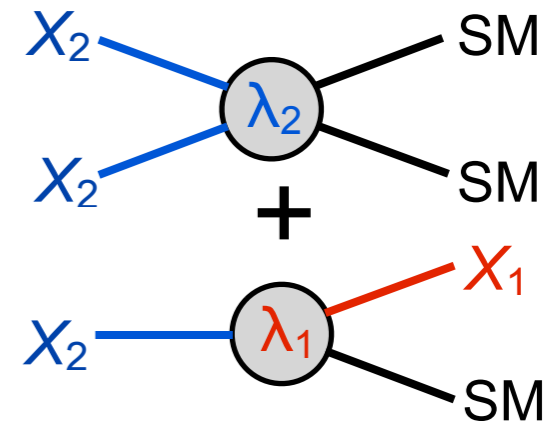
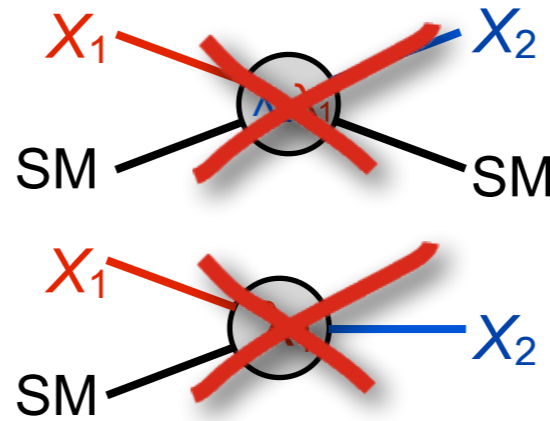
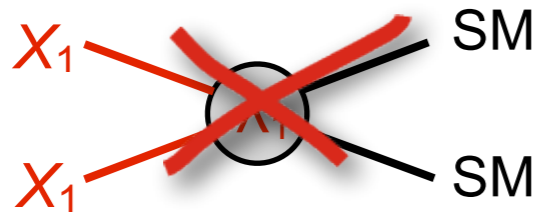
Annihilation:

$$\lambda_1 \lll \lambda_2$$

Conversion:

+

Co-annihilation:



superWIMP scenario

