

# SIMP Dark Matter

Based on:

NB, C. Garcia-Cely & R. Rosenfeld: arXiv:1501.01973 - JCAP 1504 (2015) 04, 012

NB, X. Chu, C. Garcia-Cely, T. Hambye & B. Zaldivar: arXiv:1510.08063 - JCAP 1603 (2016) 03, 018

NB & X. Chu: arXiv:1510.08527 - JCAP 1601 (2016) 01, 006

NB, J. Pradler & X. Chu: arXiv:1702.04906 - Phys.Rev. D95 (2017) 11, 115023

**Nicolás BERNAL**



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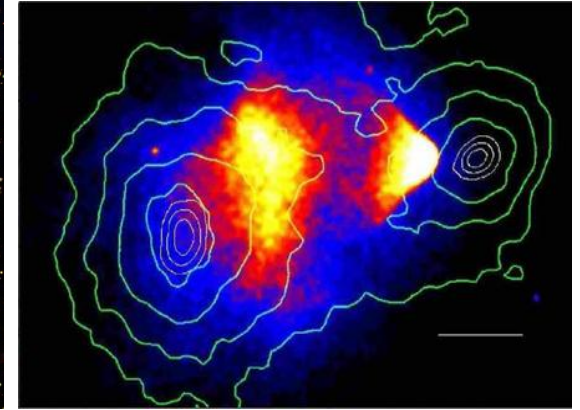
October 17<sup>th</sup>, 2017



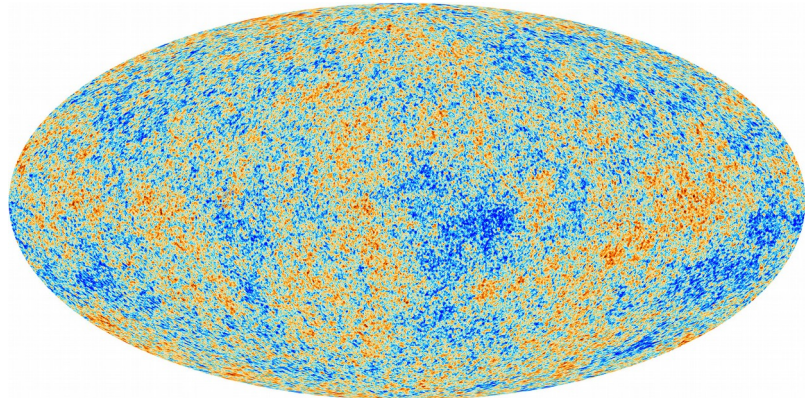
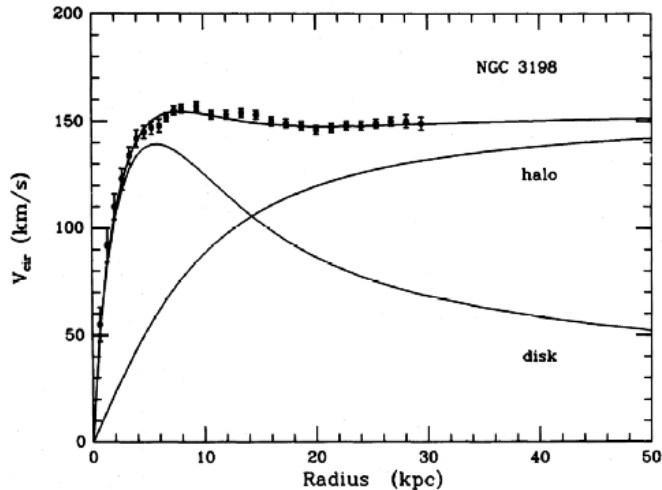
# Evidences for Dark Matter

Several observations indicate the existence of non-luminous Dark Matter (missing force) at very different scales!

- \* Galactic rotation curves
- \* RC in Clusters of galaxies
- \* Clusters of galaxies
- \* CMB anisotropies



DISTRIBUTION OF DARK MATTER IN NGC 3198



# Evidences for Dark Matter

Several observations indicate the existence of non-luminous Dark Matter (missing force) at very different scales!

- \* Galactic rotation curves
- \* RC in Clusters of galaxies
- \* Clusters of galaxies
- \* CMB anisotropies

**Dark Matter is there! :-)**

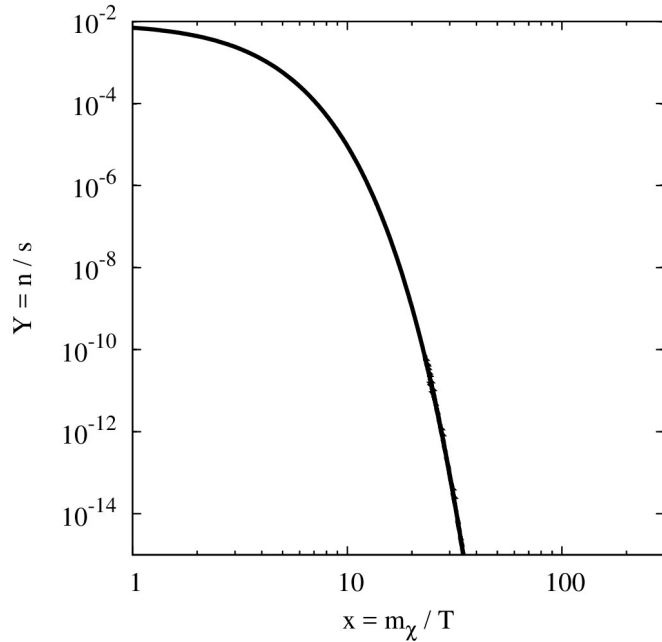
But what is it? :-/

- \* Neutral
- \* Massive enough
- \* 'Weak' interactions
- \* Stable or long-lived

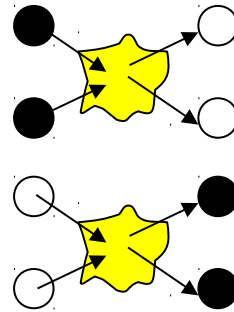
*Dark Matter needs  
New Physics beyond the Standard Model!*

# **Thermal Collisionless Cold Dark Matter**

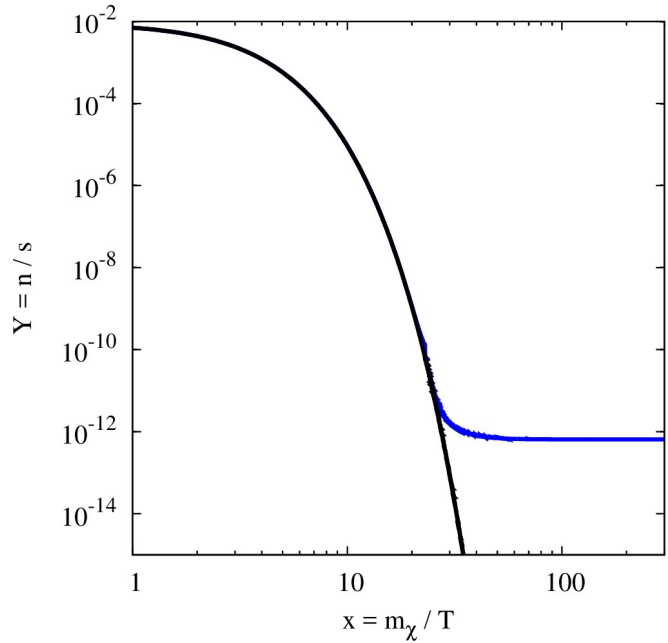
# Vanilla WIMP Dark Matter



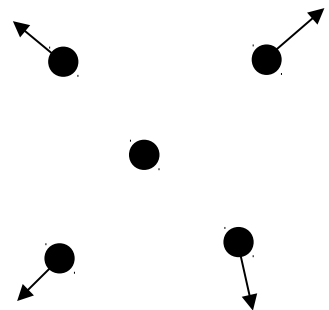
Early Universe:  
DM in thermal equilibrium  
with the Standard Model.



# Vanilla WIMP Dark Matter

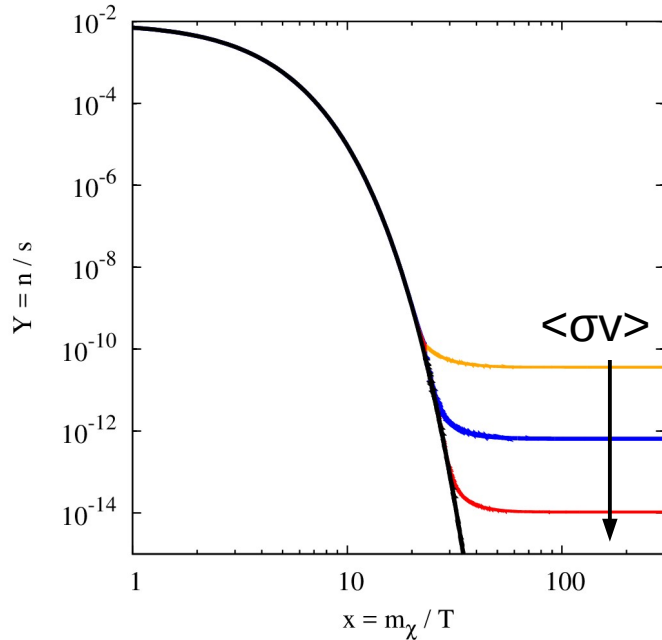


Due to the expansion of the Universe DM particles fall out of equilibrium and cannot annihilate any more.



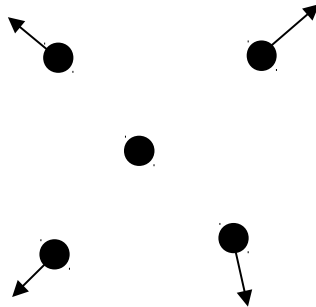
A Relic Density of DM is obtained which remains constant.

# Vanilla WIMP Dark Matter



$$\langle \sigma v \rangle \sim 3 \times 10^{-26} \text{ cm}^3/\text{s}$$

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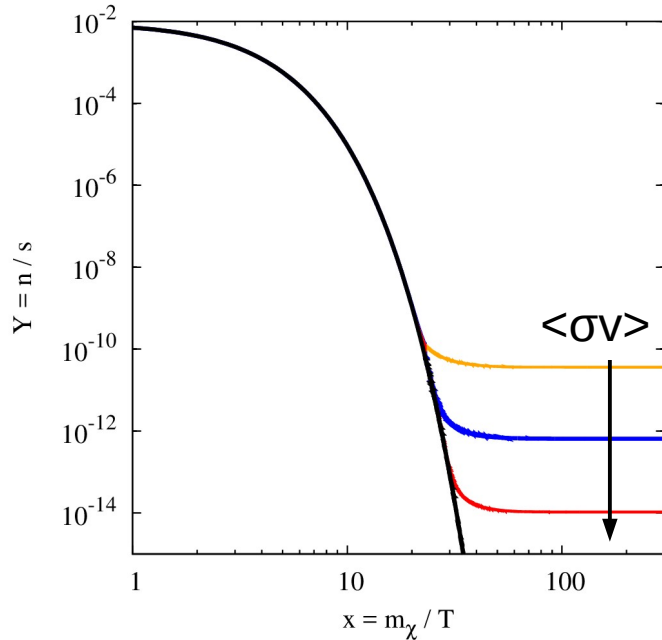
A particle with very weak interactions decouples earlier, having a larger relic density.

A Relic Density of DM is obtained which remains constant.

A particle with stronger interactions keeps in equilibrium for longer, and is more diluted.

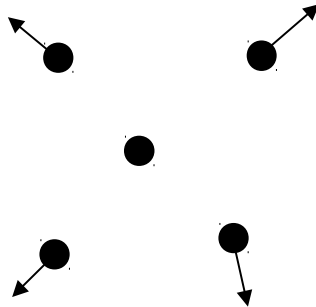


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# **Collisionless Dark Matter in Trouble?!**

# Collisionless cold Dark Matter in Trouble

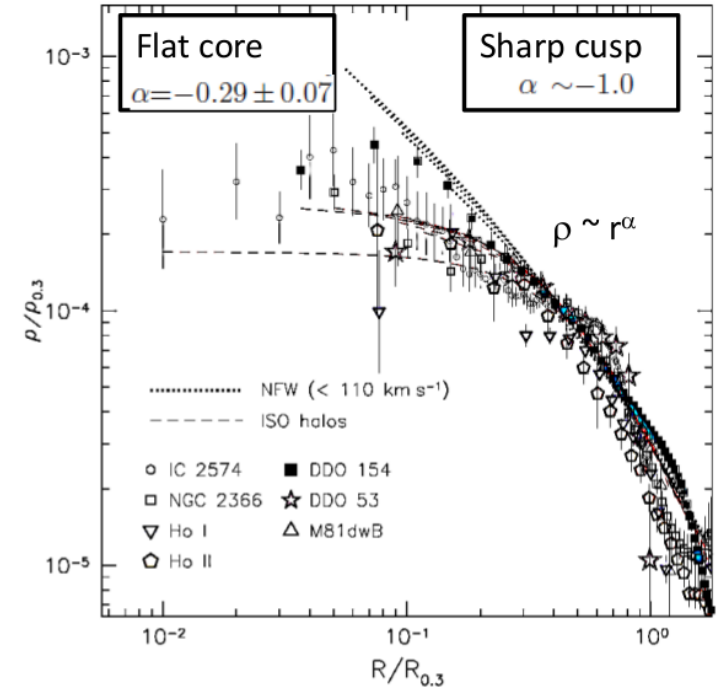
## \* Core-vs-cusp problem

- Central densities of halos exhibit cores
- N-body simulations  $\rho \sim r^{-1}$

Moore '94, Flores & Primack '94

THINGS (dwarf galaxy survey) - Oh et al. (2011)

- \* Field dwarfs
- \* Satellite dwarfs galaxies
- \* Low surface brightness galaxies (LSBs)
- \* Clusters

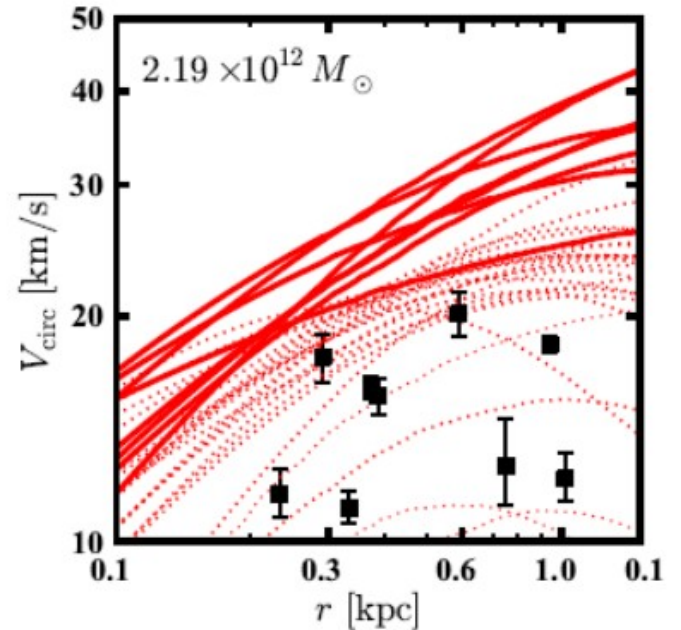
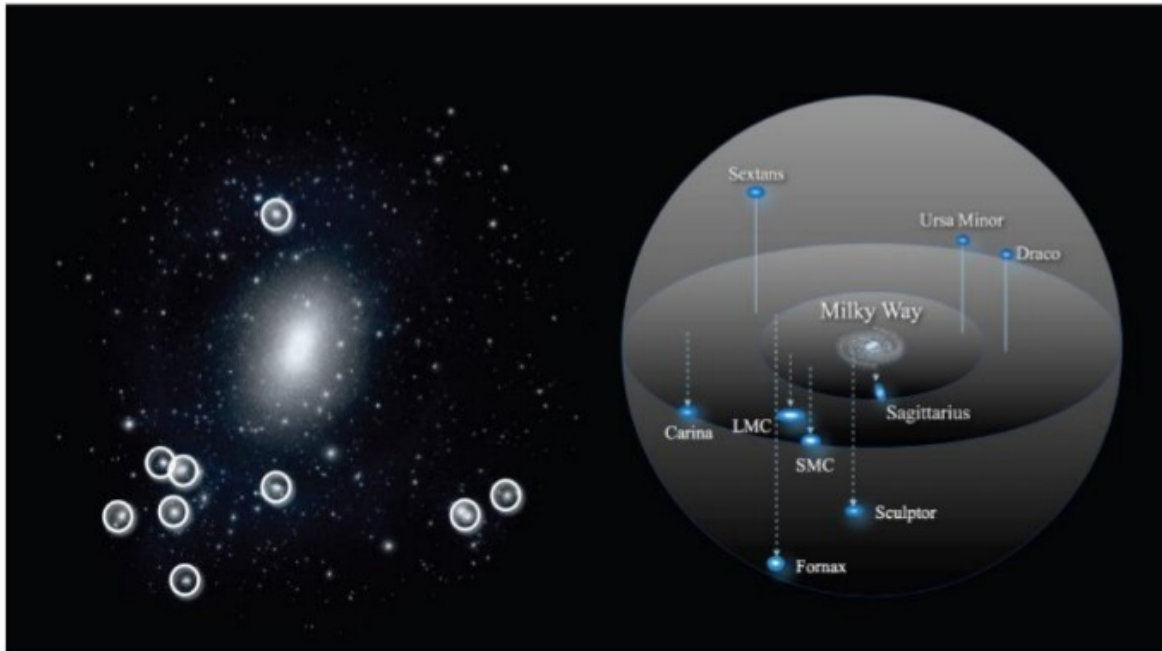


# Collisionless cold Dark Matter in Trouble

## \* Too-big-to-fail problem

Boylan-Kolchin, Bullock, Kaplinghat ('11 + '12)

MW galaxy should have  $\sim 10$  satellite galaxies which are more massive than the most massive dwarf spheroidals



# Small-scale problems → Self-interacting DM

## Small-scale problems:

- \* Core-vs-cusp
- \* Too-big-to-fail

## Possible solutions:

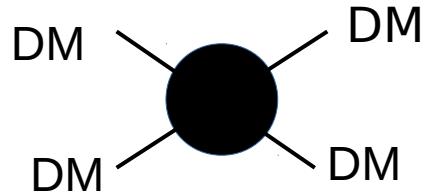
- \* Baryonic physics
  - Can't use DM-only simulations to model real DM+baryon Universe
  - Astrophysical observations not being modeled correctly  
(Suppressed gas cooling efficiency, low star-formation efficiency, supernova feedback, large velocity anisotropy...)
- \* *Dark Matter*
  - DM may not be collisionless*

# Self-interacting Dark Matter

CDM structure problems could be solved if dark matter is **self-interacting**.

Dark matter particles in halos elastically scatter with other dark matter particles.

Spergel & Steinhardt (2000)

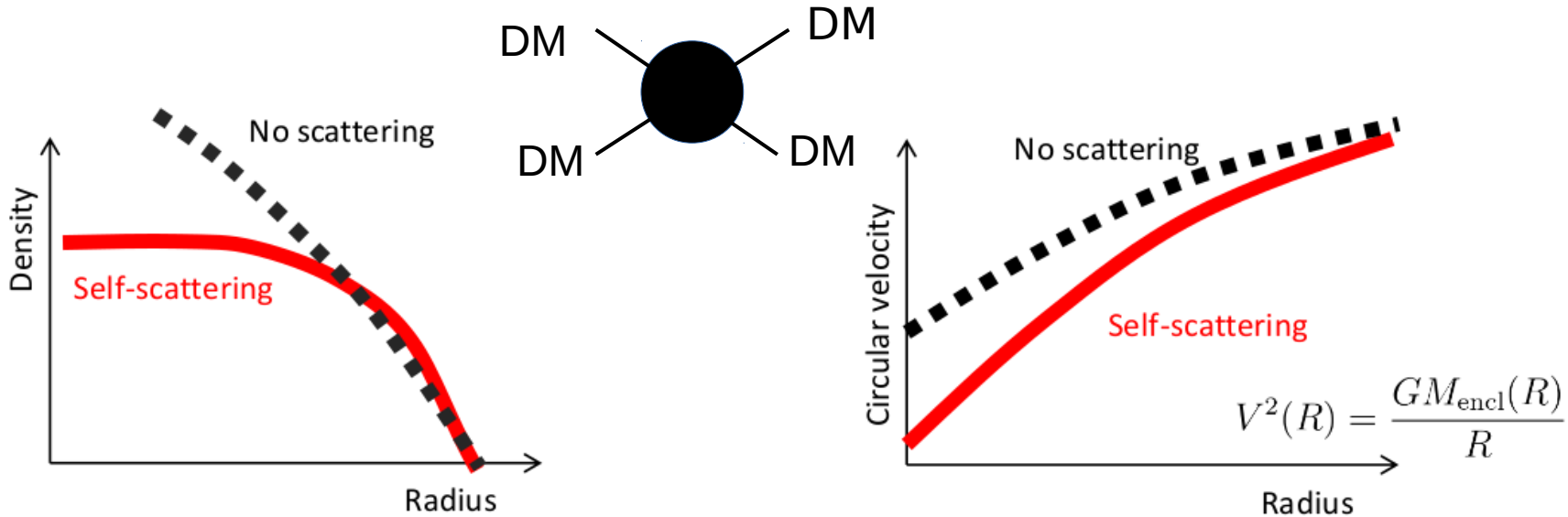


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**Self-interactions solve core-vs-cusp**

*Particles get scattered out of dense halo centers*

**Self-interactions solve too-big-to-fail**

*Rotation curves reduced (less enclosed mass)  
Simulated satellites matched to observations*

# Small-scale problems → Self-interacting DM

## Small-scale problems:

- \* Core-vs-cusp
- \* Too-big-to-fail

## Possible solutions:

- \* Baryonic physics
  - Can't use DM-only simulations to model real DM+baryon Universe
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(Suppressed gas cooling efficiency, low star-formation efficiency, supernova feedback, large velocity anisotropy...)
- \* *Dark Matter*  
*DM may not be collisionless*

$$\left( \frac{\sigma_{\text{scatter}}}{m_{\chi}} \right)_{\text{obs}} = (0.1 - 10) \text{ cm}^2/\text{g} \quad \sim \text{few barns/GeV}$$

$$\frac{\sigma_{\text{scatter}}}{m_{\chi}} \lesssim 1 \text{ cm}^2/\text{g}. \quad \text{From the Bullet Cluster}$$



# **The SIMPlEst DM model ever: Singlet Scalar Dark Matter**

# Singlet Scalar DM

McDonald '07

S is a singlet scalar, **protected by** a  $Z_2$

$$V = \mu_S^2 S^2 + \lambda_S S^4 + \lambda_{HS} |H|^2 S^2$$

3 free parameters:

- \*  $m_S$  DM mass
- \*  $\lambda_{HS}$  Higgs portal
- \*  $\lambda_S$  DM quartic coupling

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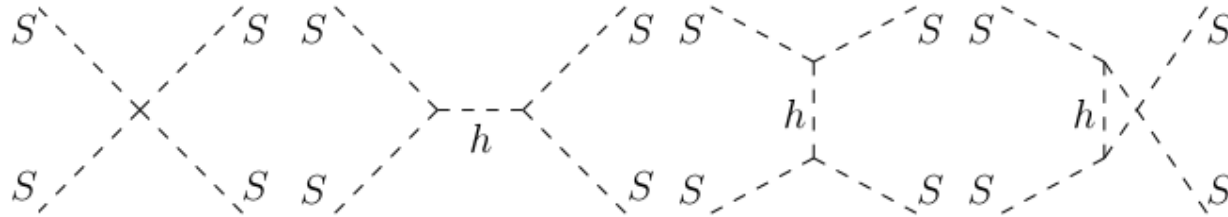
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}

← Concentrated on this

← ~ Ignored!

# Dark Matter Self-Interactions

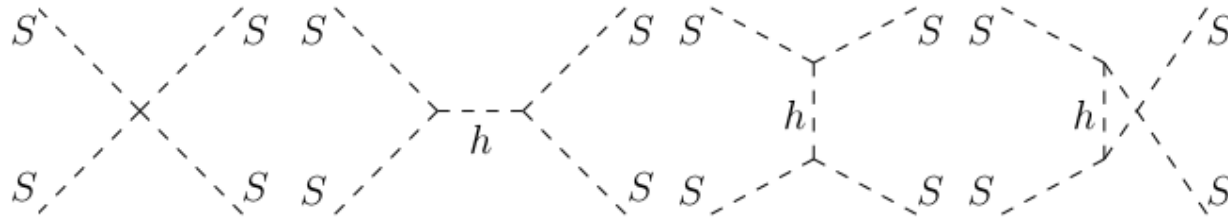


$$\frac{\sigma_{SS}}{m_S} \sim \frac{9}{8\pi} \frac{\lambda_S^2}{m_S^3}$$

$$0.1 \lesssim \frac{\sigma_{SS}}{m_S} \lesssim 10 \text{ cm}^2/\text{g}$$

Implies  $\left\{ \begin{array}{l} * \lambda_S \sim 1 \\ * m_S \sim 100 \text{ MeV} \end{array} \right.$

# Dark Matter Self-Interactions & Invisible Higgs decay



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Implies  $\left\{ \begin{array}{l} * \lambda_S \sim 1 \\ * m_S \sim 100 \text{ MeV} \end{array} \right.$

The Higgs tends to annihilate into DM

$\text{BR}(h \rightarrow \text{inv.}) < 20\%$

\*  $\lambda_{HS} < 7 \times 10^{-3}$

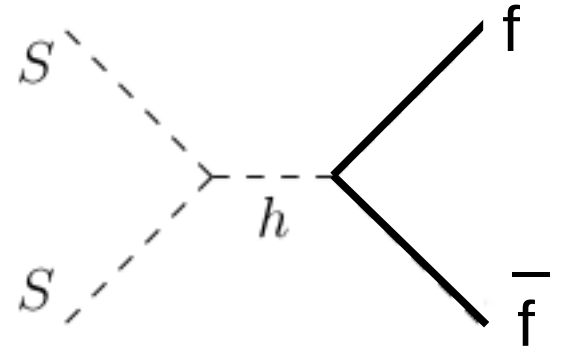
# **How to produce such a Self-Interacting Dark Matter?**

# WIMP DM :-/

DM can (only) annihilate into light fermions  
other annihilation channels kinematically closed!

$$\langle \sigma_{SS \rightarrow f \bar{f} \nu} \rangle \sim \frac{\lambda_{HS}^2}{\pi} \frac{m_f^2}{m_h^4}$$

$$\langle \sigma_{SS \rightarrow f \bar{f} \nu} \rangle \ll 10^{-26} \text{ cm}^3/\text{s}$$



- Universe overclosed
- SSDM with sizable self-interactions can not be a WIMP

**Again:  
How to produce such a  
Self-Interacting Dark Matter?**

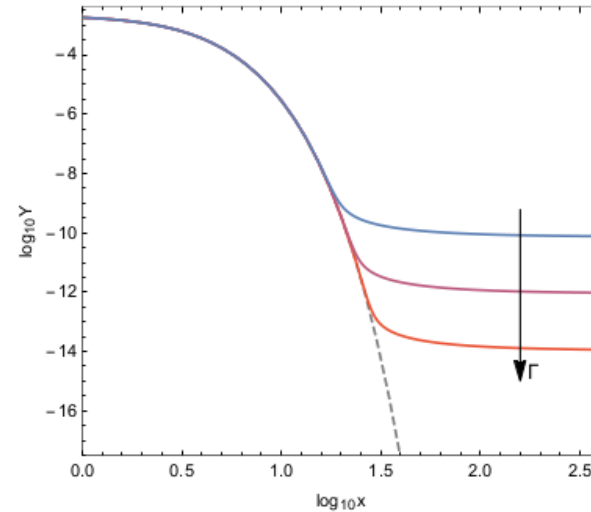
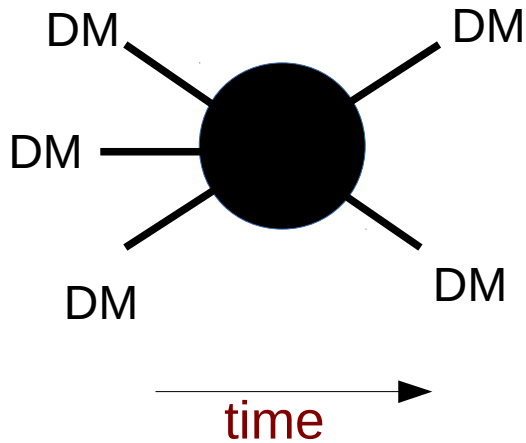


# SIMP DM

## 3 → 2 annihilations

Hochberg, Kuflik, Volansky & Wacker '14

$$\frac{dn}{dt} + 3 H n = -\langle \sigma v^2 \rangle_{3 \rightarrow 2} (n^3 - n^2 n_{\text{eq}})$$

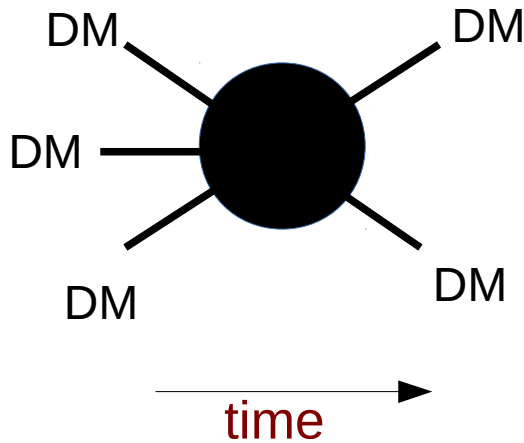


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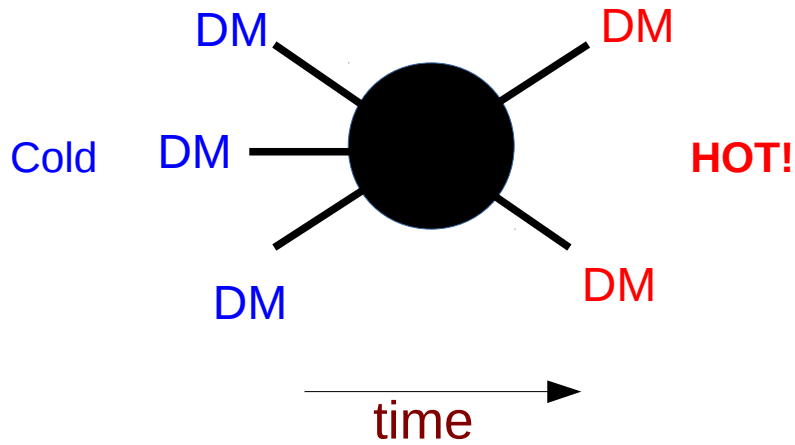


- \* DM in the MeV range
- \* Small DM-SM portal
- \*  $\lambda_S \sim 1$
- \* 'Strong' Self-interactions  
→ SIMP DM

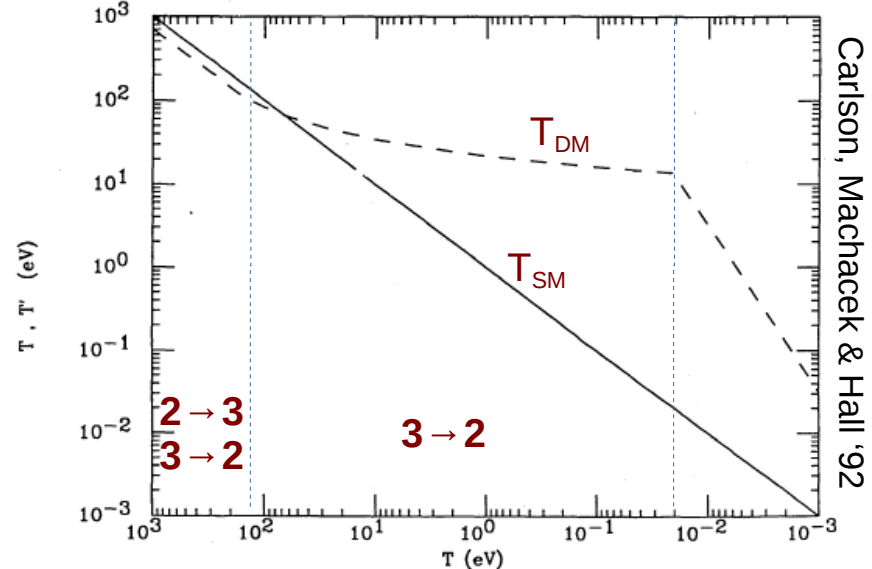
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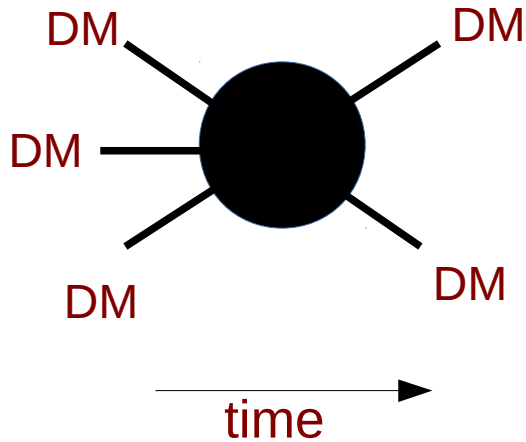
**Caveat:** 3 → 2 annihilations  
pump heat into the dark sector!



# SIMP DM

## 3 → 2 annihilations

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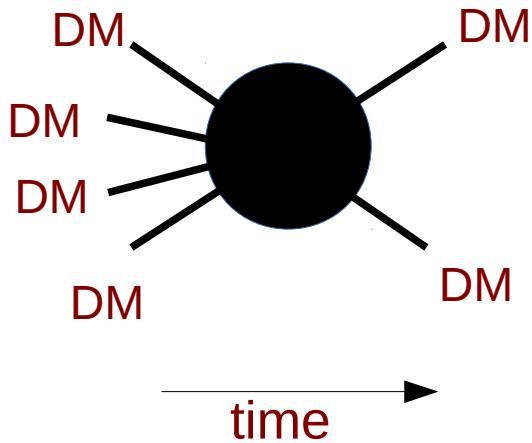


However 3 → 2 reactions are forbidden in most common scenarios where the DM stability is guaranteed by a  $Z_2$  symmetry (R-parity in SUSY, K-parity in Kaluza-Klein...)

# SIMP DM

## 4 → 2 annihilations

$$\frac{dn}{dt} + 3 H n = -\langle \sigma v^3 \rangle_{4 \rightarrow 2} (n^4 - n^2 n_{\text{eq}}^2)$$



However 3 → 2 reactions are forbidden in most common scenarios where the DM stability is guaranteed by a  $Z_2$  symmetry (R-parity in SUSY, K-parity in Kaluza-Klein...)

But  $Z_2$  symmetries allow 4 → 2 annihilations!

# Singlet Scalar DM

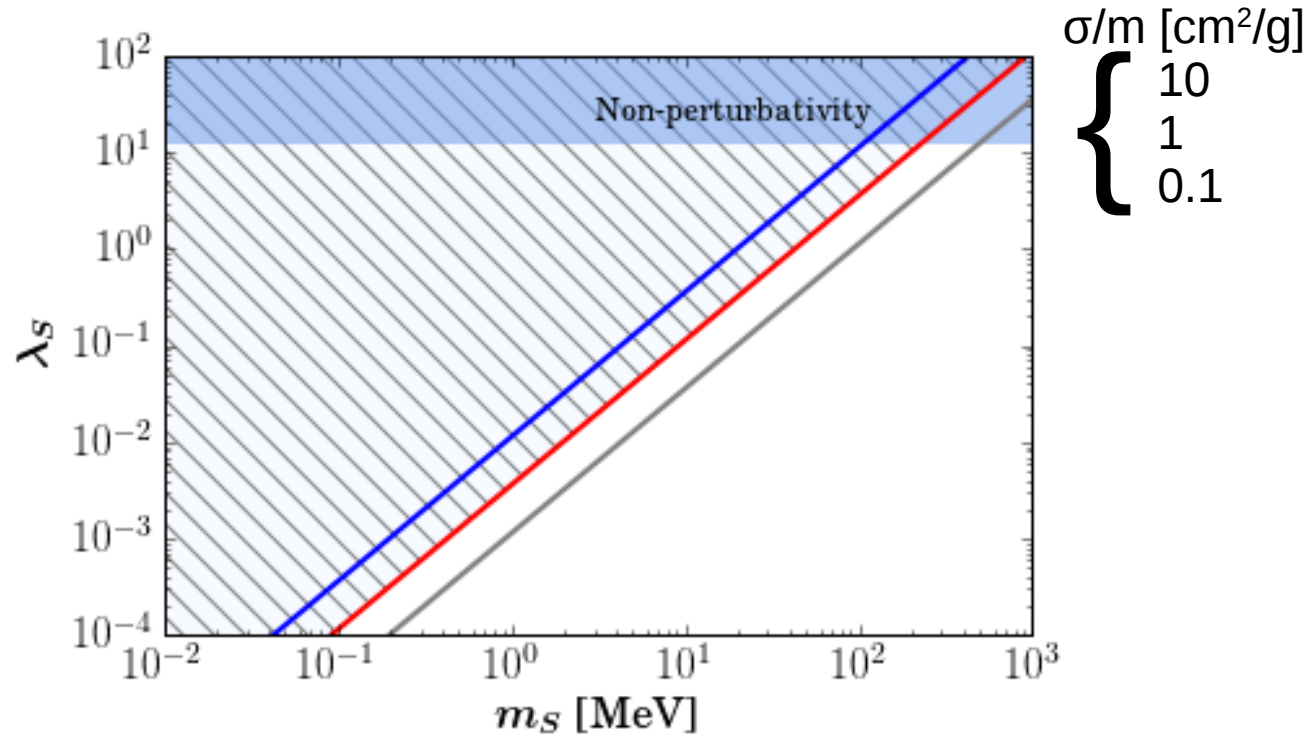
## $4 \rightarrow 2$ annihilations

$$\frac{dn}{dt} + 3 H n = -\langle \sigma v^3 \rangle_{4 \rightarrow 2} (n^4 - n^2 n_{\text{eq}}^2)$$

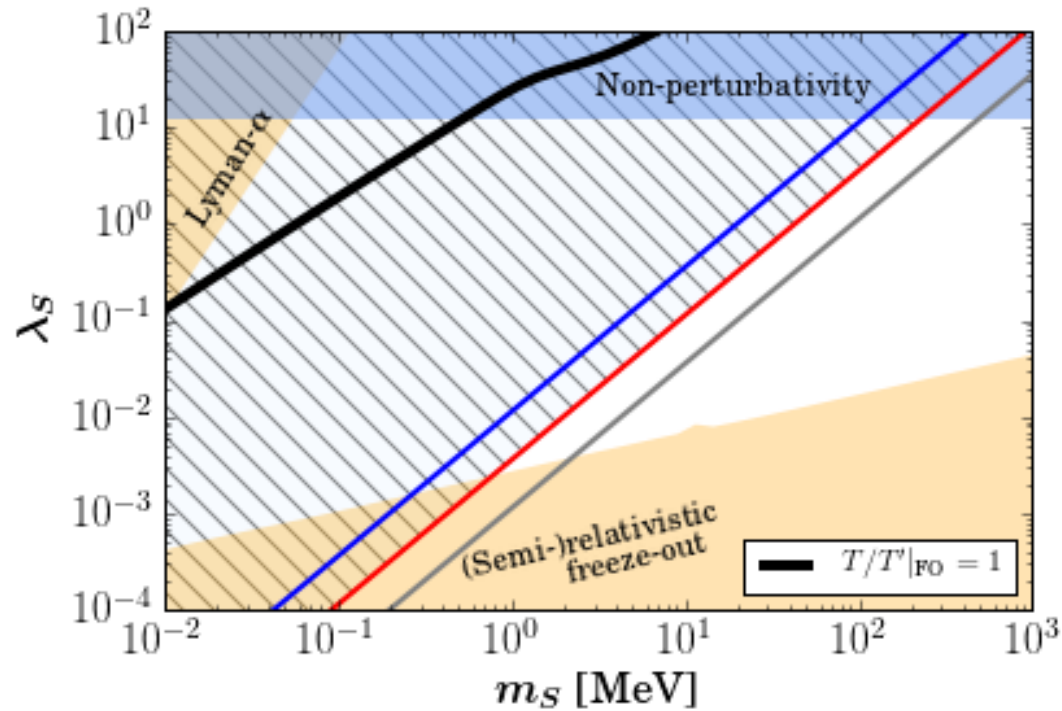


$$\langle \sigma v^3 \rangle_{4 \rightarrow 2} \sim \frac{27\sqrt{3}}{8\pi} \frac{\lambda_S^4}{m_S^8}$$

# Singlet Scalar DM $4 \rightarrow 2$ annihilations



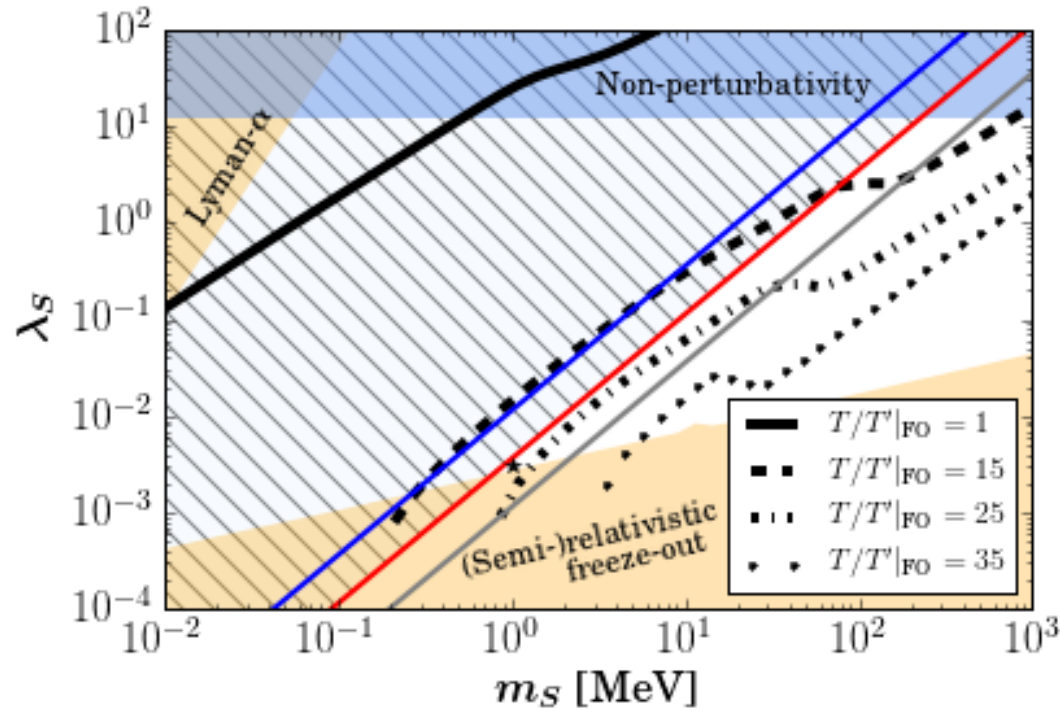
# Singlet Scalar DM $4 \rightarrow 2$ annihilations



$$T_{SM} = T_{DM} \text{ @ DM freeze-out}$$



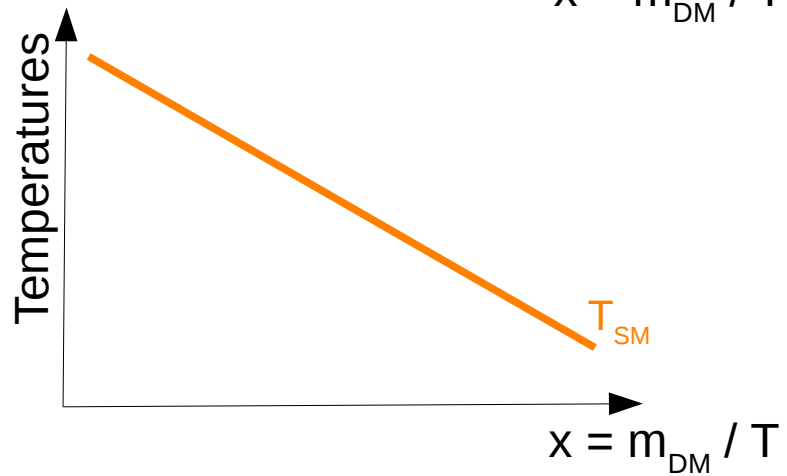
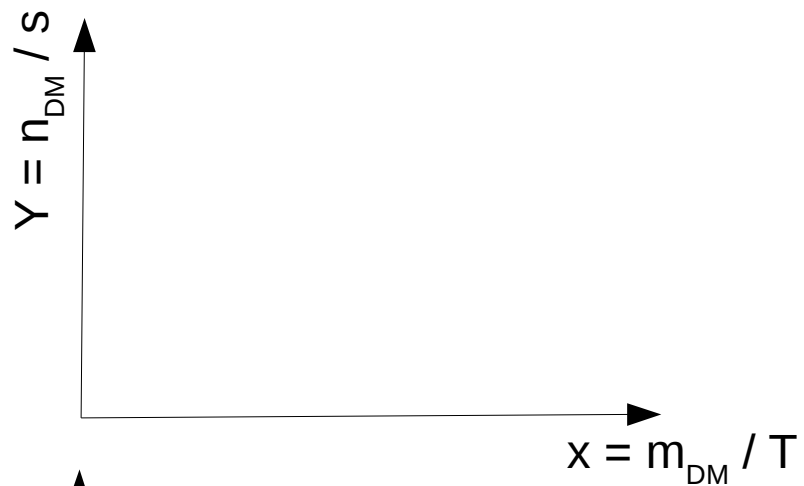
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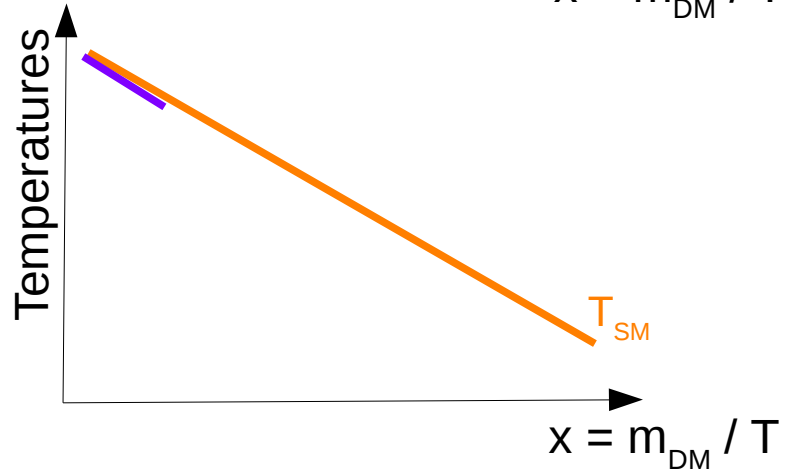
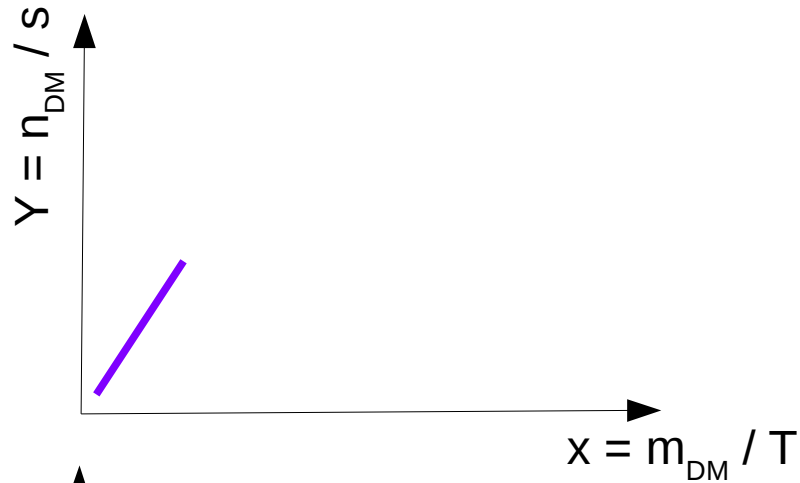
$T_{\text{SM}} = T_{\text{DM}}$   
&  
 $T_{\text{SM}} \neq T_{\text{DM}} @ \text{DM freeze-out}$

**How to produce such a  
difference of temperatures?**

# (Non-) Thermal evolution of DM



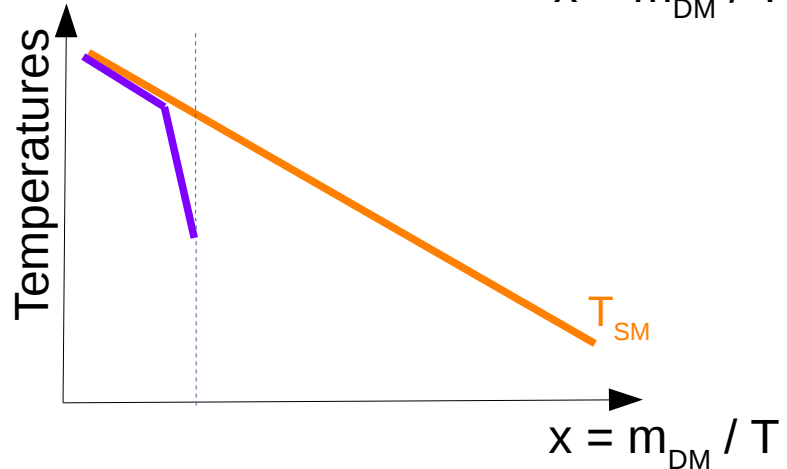
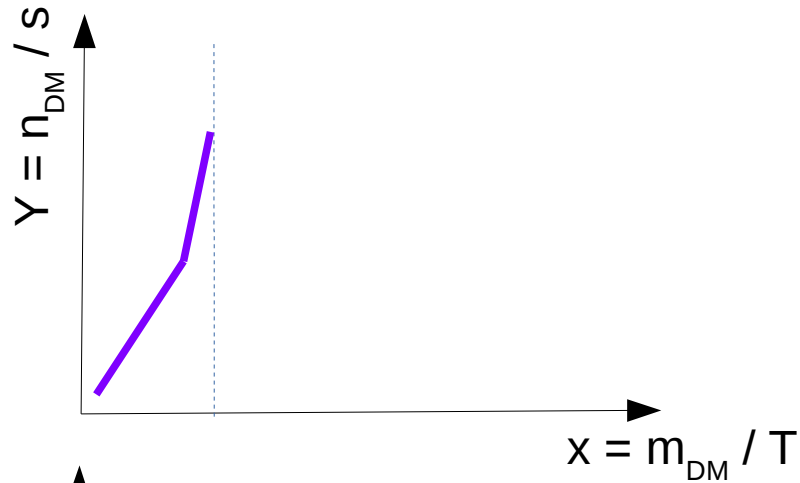
# (Non-) Thermal evolution of DM



## DM Production

- \* Out-of-equilibrium production à la freeze-in:  $\mathbf{h} \rightarrow \mathbf{SS}$
- DM in kinetic equilibrium via  $\mathbf{2} \leftrightarrow \mathbf{2}$
- DM inherits SM temperature

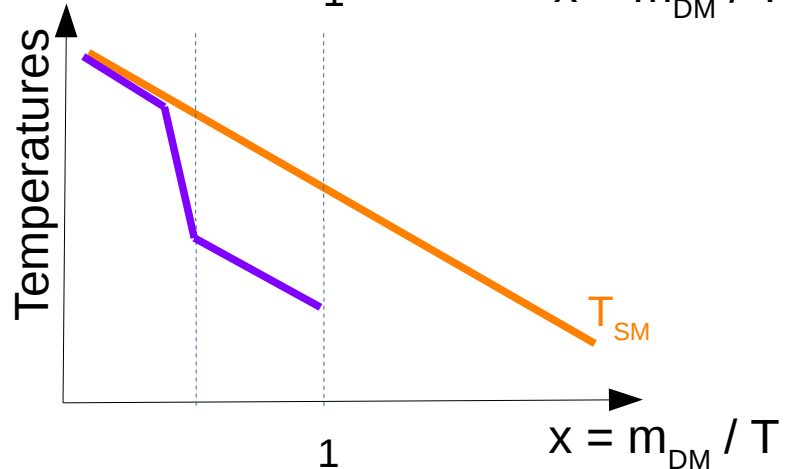
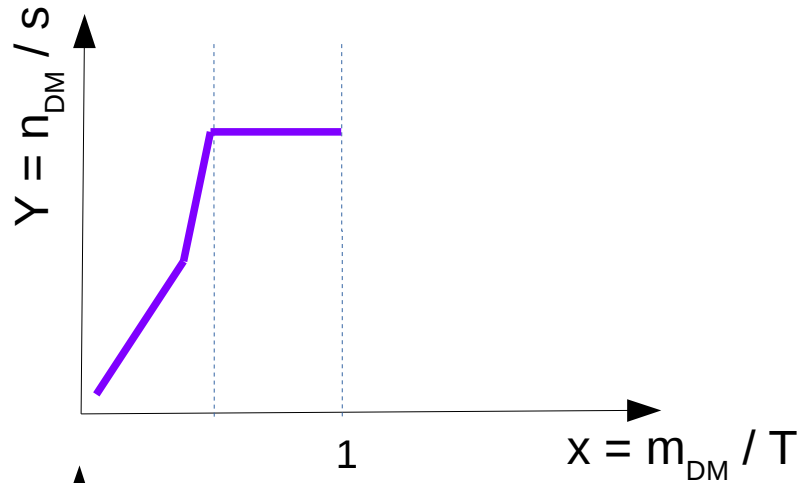
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- \* DM populates rapidly via out-of-equilibrium  $\mathbf{2} \rightarrow \mathbf{4}$ .  
Price to pay: Dramatic decrease of  $T_{\text{DM}}$

# (Non-) Thermal evolution of DM



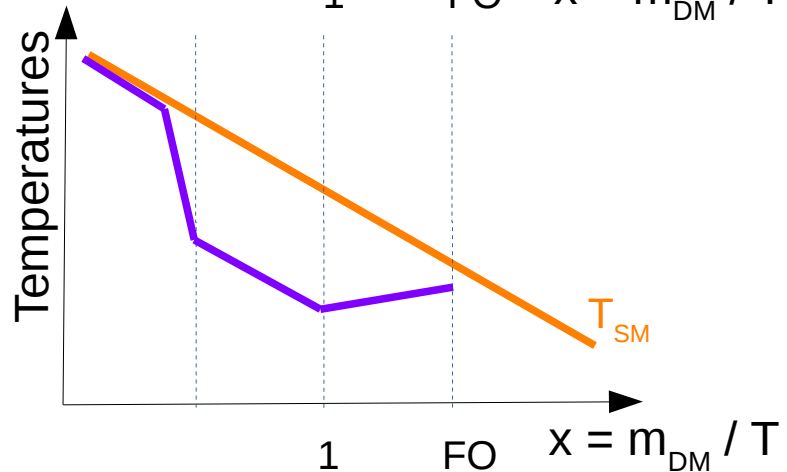
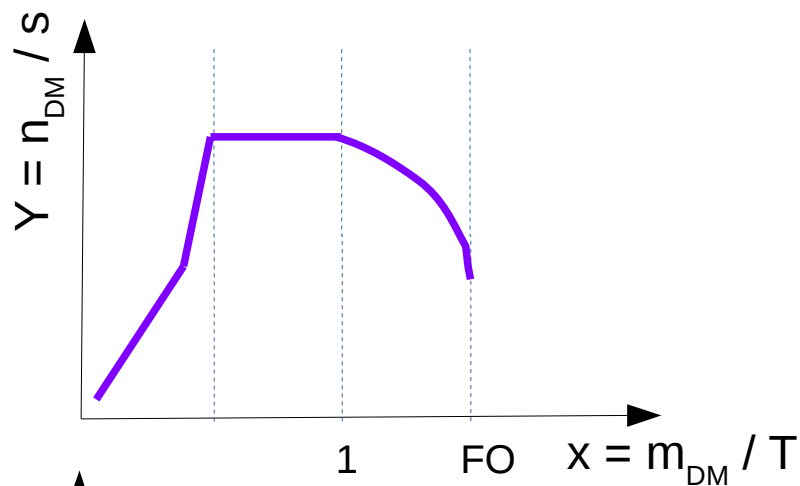
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## Thermal Equilibrium

- \* Chemical equilibrium  $\mathbf{2} \leftrightarrow \mathbf{4}$

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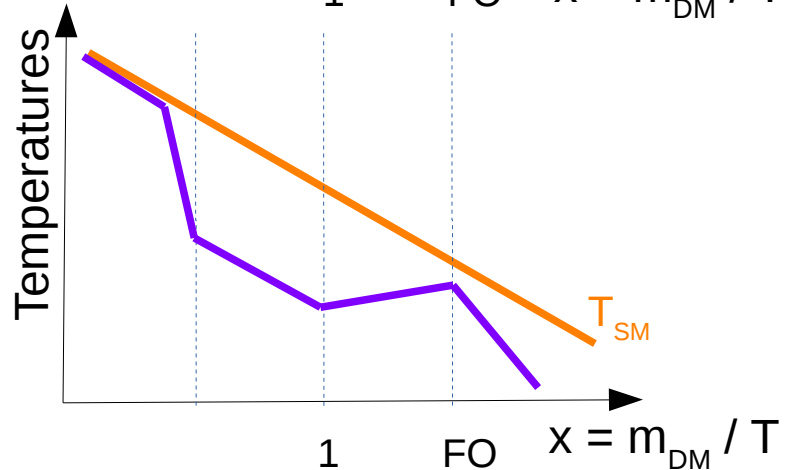
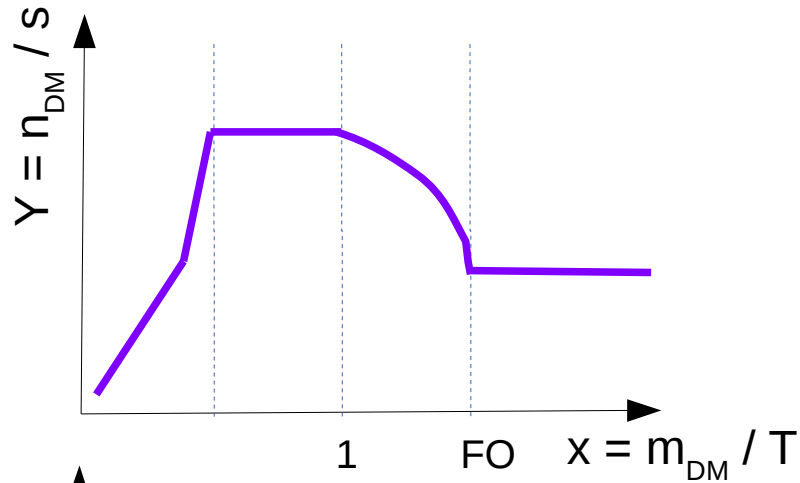
## Thermal Equilibrium

- \* Chemical equilibrium  $\mathbf{2} \leftrightarrow \mathbf{4}$

## DM Annihilation

- \* Freeze-out  $\mathbf{4} \rightarrow \mathbf{2}$

# (Non-) Thermal evolution of DM



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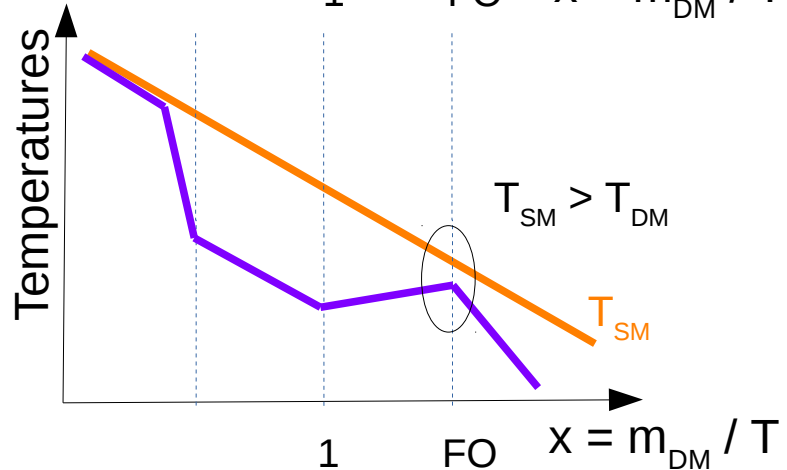
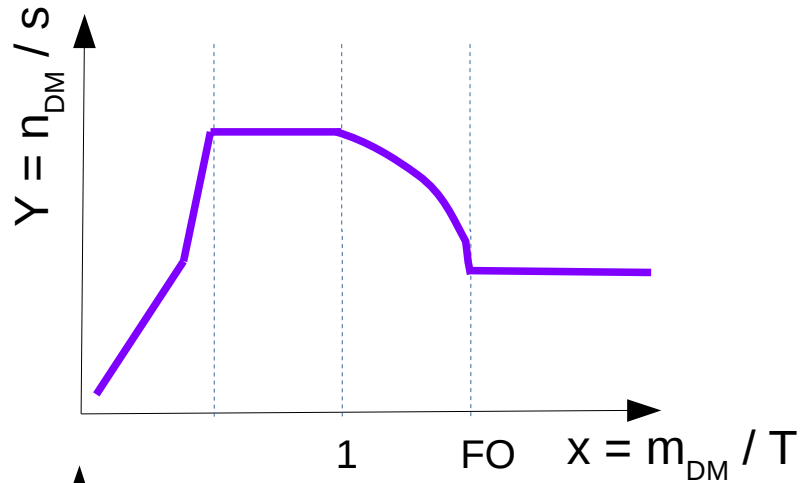
- \* Freeze-out  $\mathbf{4} \rightarrow \mathbf{2}$

## After the Freeze-out

- \* Relic abundance  
Non-relativistic DM cools down faster



# (Non-) Thermal evolution of DM



## DM Production

- \* Out-of-equilibrium production à la freeze-in:  $\mathbf{h} \rightarrow \mathbf{SS}$
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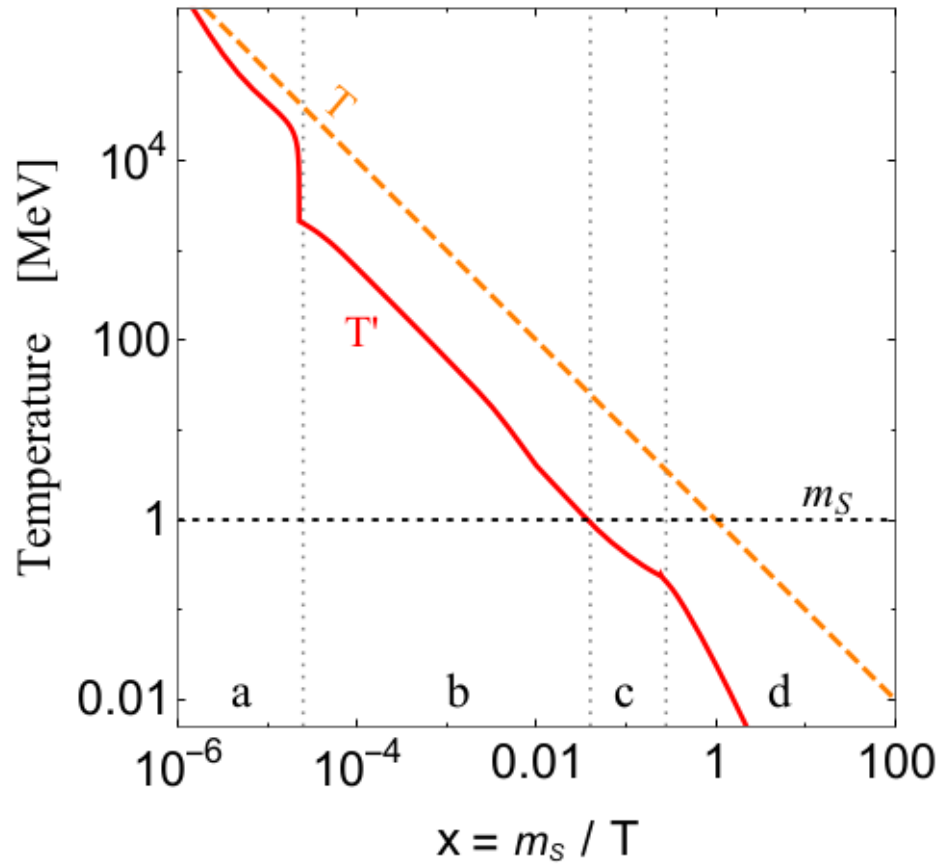
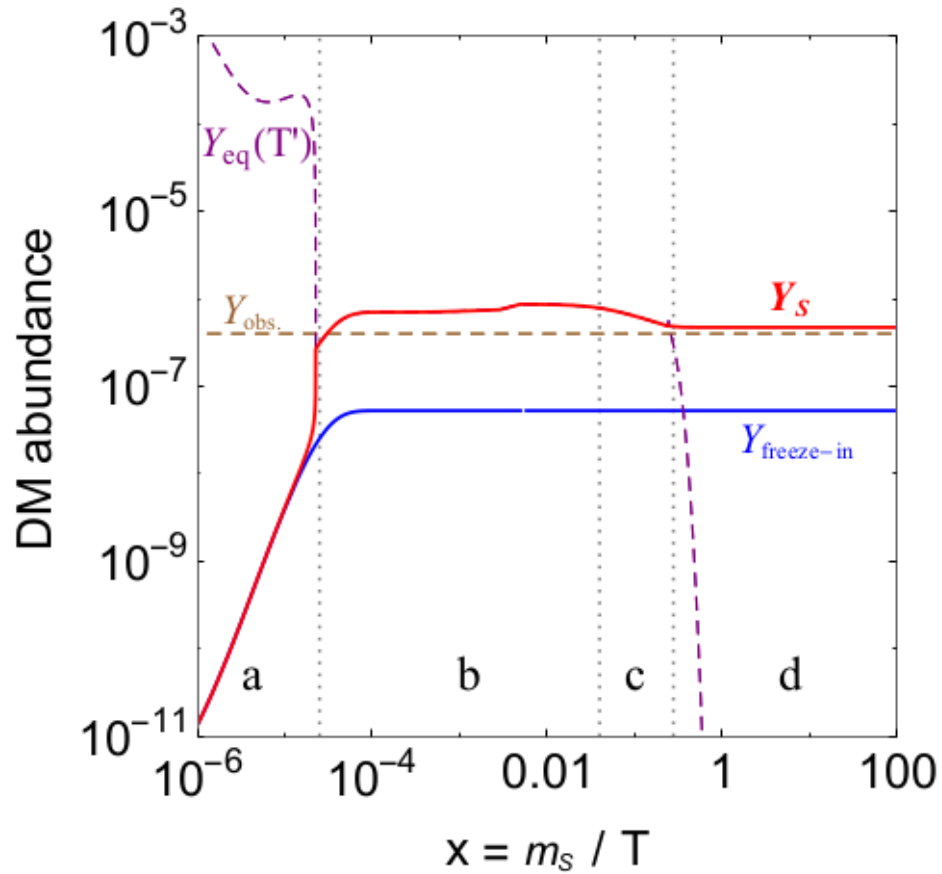
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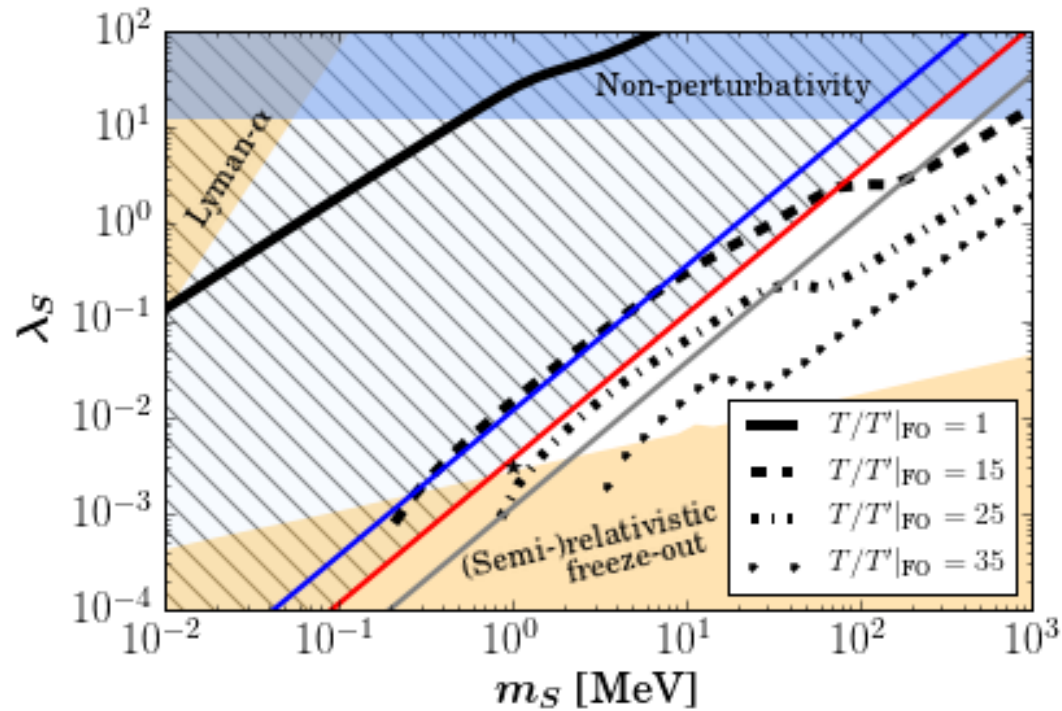
## After the Freeze-out

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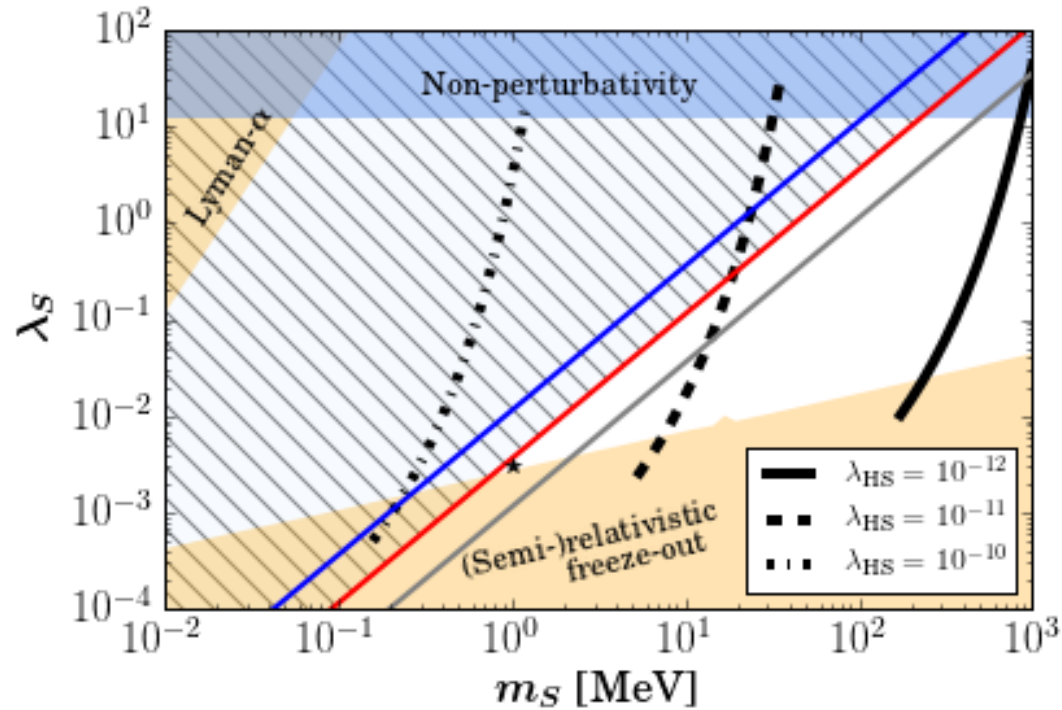
# Generating $T_{\text{DM}} < T_{\text{SM}}$ via the Higgs Portal



# Singlet Scalar DM $4 \rightarrow 2$ annihilations



# Singlet Scalar DM $4 \rightarrow 2$ annihilations



**Now let's Split SIMPs! :-)**

# Splitting SIMPs: fermions

- Fermionic DM:

Dirac fermion  $\Psi$  split by small Majorana masses  $m_L$  and  $m_R$ .

$$\mathcal{L}_\Psi = \bar{\Psi} (i\not{D} - M_D) \Psi - \frac{m_L}{2} (\bar{\Psi}^c P_L \Psi + h.c.) - \frac{m_R}{2} (\bar{\Psi}^c P_R \Psi + h.c.)$$

$$\chi_1 \simeq \frac{i}{\sqrt{2}}(\Psi - \Psi^c), \quad \chi_2 \simeq \frac{1}{\sqrt{2}}(\Psi + \Psi^c)$$

$$\text{Pseudo-Dirac } \chi_{1,2}: \quad m_{1,2} \simeq M_D \mp \frac{m_L + m_R}{2} + O(\delta),$$

# Splitting SIMPs: fermions

- Fermionic DM:

Dirac fermion  $\Psi$  split by small Majorana masses  $m_L$  and  $m_R$ .

$$\mathcal{L}_\Psi = \bar{\Psi} (i\not{D} - M_D) \Psi - \frac{m_L}{2} (\bar{\Psi}^c P_L \Psi + h.c.) - \frac{m_R}{2} (\bar{\Psi}^c P_R \Psi + h.c.)$$

$$\chi_1 \simeq \frac{i}{\sqrt{2}}(\Psi - \Psi^c), \quad \chi_2 \simeq \frac{1}{\sqrt{2}}(\Psi + \Psi^c)$$

$$\text{Pseudo-Dirac } \chi_{1,2}: \quad m_{1,2} \simeq M_D \mp \frac{m_L + m_R}{2} + O(\delta),$$

- The gauged dark U(1) symmetry explicitly broken by  $\Delta m$   
→ the interaction with DM proceeds off-diagonally!

$$\mathcal{L}_{\text{int}, \chi} = i g_V \bar{\chi}_1 \gamma^\mu \chi_2 V_\mu + O(\delta)$$

- The dark gauge boson:

$$\mathcal{L}_V = -\frac{1}{4} V_{\mu\nu} V^{\mu\nu} + \frac{m_V^2}{2} V^2 - \kappa V_\mu J_{\text{SM}}^\mu,$$

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Free parameters:  
 $m, \Delta m, m_V, g_V$  and  $\kappa$ .



# Splitting SIMPs: scalars

- Scalar DM:

Complex scalar  $\Phi$  split by a mass-squared parameter  $m_\phi^2$ .

$$\mathcal{L}_\Phi = |D_\mu \Phi|^2 + M^2 |\Phi|^2 + (m_\phi^2 \Phi^2 + h.c.) - V_\Phi$$

$$V_\Phi = \lambda_\Phi |\Phi|^4 + \frac{\lambda'_\Phi}{2} (\Phi^4 + h.c.) + \lambda_m |\Phi|^2 |H|^2 + (\lambda'_m \Phi^2 + h.c.) |H|^2$$

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- The gauged dark U(1) symmetry explicitly broken by  $\Delta m$   
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$$\mathcal{L}_{\text{int},\phi} = g_V (\phi_1 \partial^\mu \phi_2 - \phi_2 \partial_\mu \phi_1) V_\mu + \frac{1}{2} g_V^2 (\phi_1^2 + \phi_2^2) V^2 - V_\Phi(\phi_1, \phi_2)$$

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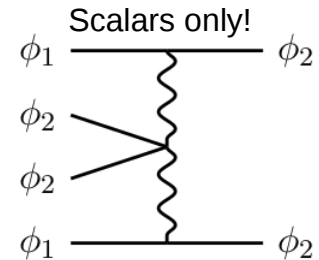
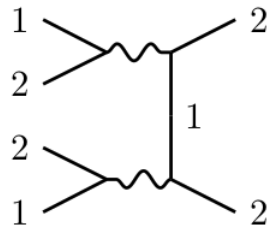
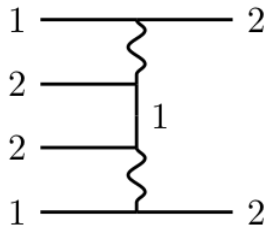
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# Producing Split SIMPs

# Split SIMPs

## 4 → 2 annihilations

$$\frac{dn}{dt} + 3Hn = -\langle \sigma v^3 \rangle_{4 \rightarrow 2} (n^4 - n^2 n_{\text{eq}}^2)$$

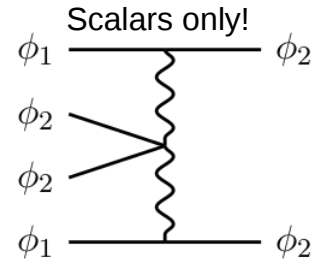
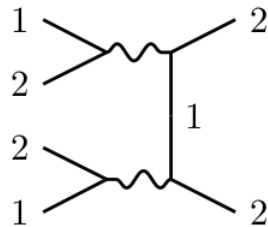
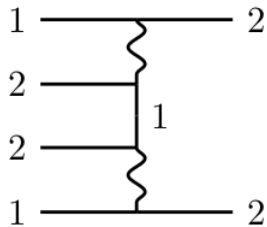


We take  $m_\nu > m_1 + m_2$

# Split SIMPs

## 4 → 2 annihilations

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$$\langle \sigma v^3 \rangle_{4 \rightarrow 2} = [\langle 1122 \rightarrow 22 \rangle + \langle 1122 \rightarrow 11 \rangle] \frac{R^2}{(1+R)^4}$$

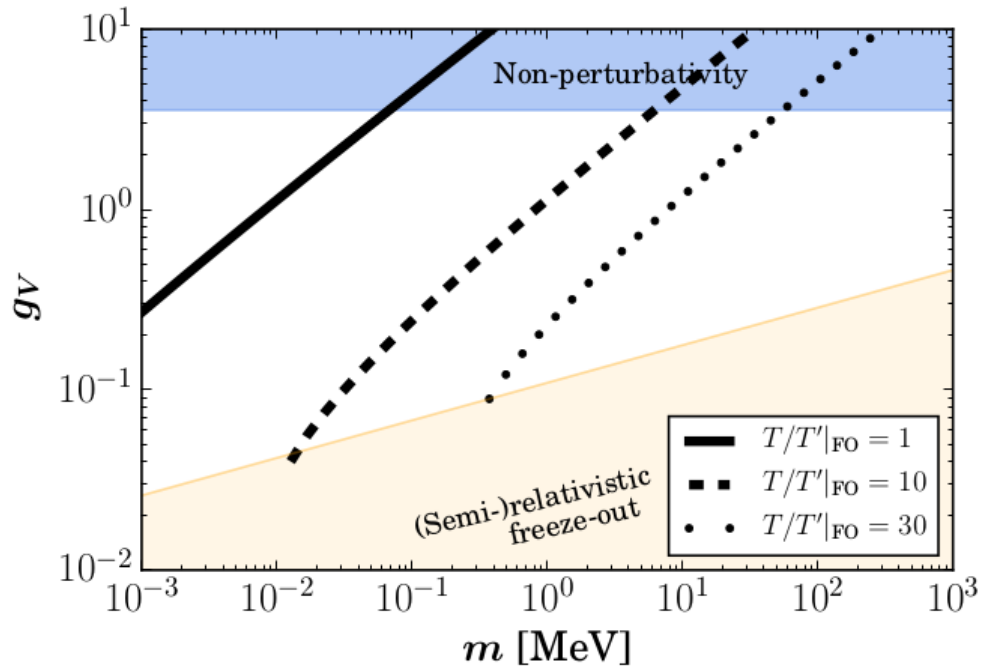
$$\langle 1122 \rightarrow 11 \rangle = \langle 1122 \rightarrow 22 \rangle = \frac{27\sqrt{3} g_V^8 (m_V^4 - 8m^2 m_V^2 - 8m^4)^2}{32\pi (m_V^4 - 2m^2 m_V^2 - 8m^4)^4}$$

We take  $m_V > m_1 + m_2$

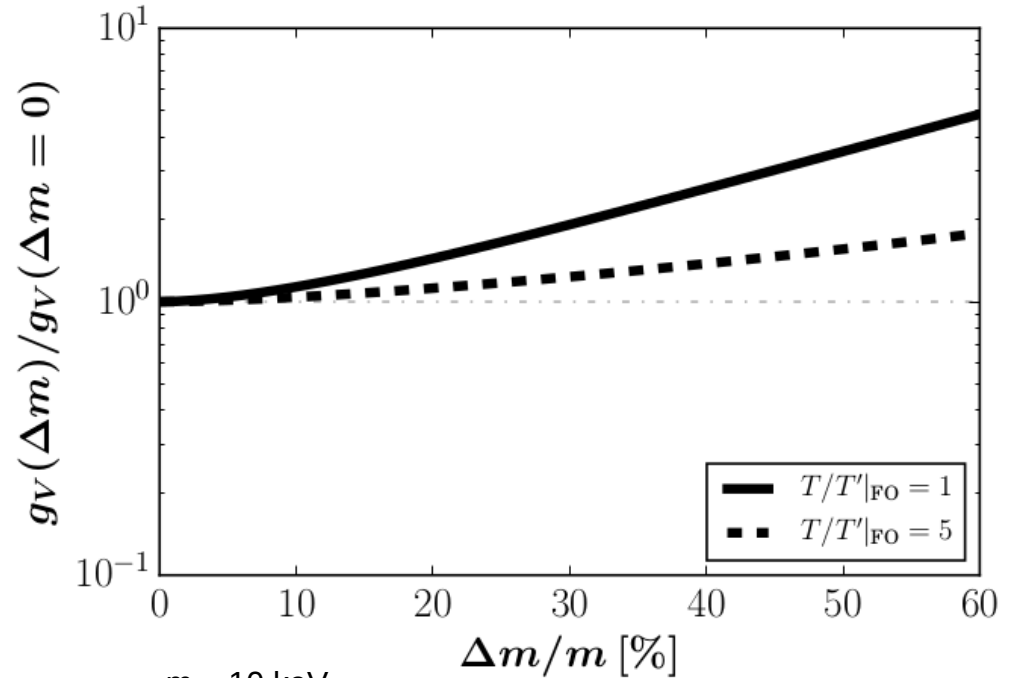
$$R(T') \equiv \frac{n_2(T')}{n_1(T')} = \left(1 + \frac{\Delta m}{m}\right)^{3/2} e^{-\frac{\Delta m}{T'}} \simeq e^{-\frac{\Delta m}{T'}}$$

# Split SIMPs

## $4 \rightarrow 2$ annihilations



$\Delta m = 0$   
 $m_\nu = 2.5 \text{ m}$



$m = 10 \text{ keV}$   
 $m_\nu = 2.5 \text{ m} + \Delta m$

# **Astrophysical Implications of Split SIMPs**

# (Late) Decay of the State 2

The decay of state 2 into state 1 is accompanied by SM radiation, possibly constrained by BBN and CMB.

$$\chi_2 \rightarrow \chi_1 V^* \rightarrow \chi_1 e^+ e^-$$

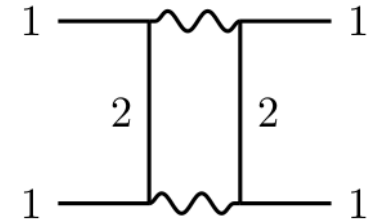
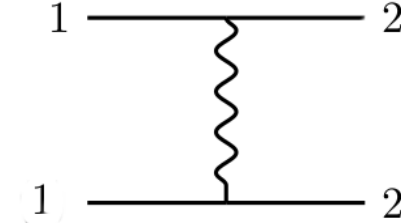
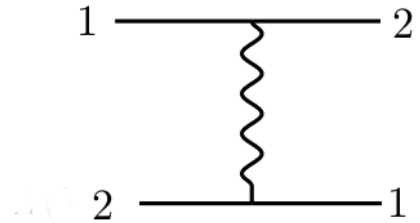
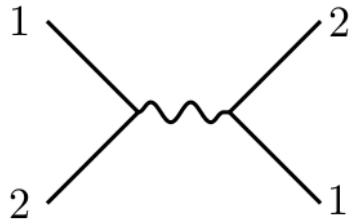
$$\Gamma_{\chi_2 \rightarrow \chi_1 e^+ e^-} \simeq \frac{2\alpha \alpha_V \kappa^2 \Delta m^5}{15\pi m_V^4} \simeq 2 H_0 \times \frac{m}{100 \text{ MeV}} \frac{\alpha_V}{\alpha} \left(\frac{\kappa}{10^{-10}}\right)^2 \left(\frac{\Delta m/m}{10^{-3}}\right)^5 \left(\frac{m}{m_V}\right)^4$$

$$\chi_2 \rightarrow \chi_1 V^* \rightarrow \chi_1 3\gamma;$$

$$\begin{aligned} \Gamma_{\chi_2 \rightarrow \chi_1 3\gamma} &\simeq \Gamma_{\chi_2 \rightarrow \chi_1 \nu \bar{\nu}} \times \left. \frac{\Gamma_{V \rightarrow 3\gamma}}{\Gamma_{V \rightarrow \nu \bar{\nu}}} \right|_{m_V \rightarrow \Delta m} \\ &\simeq H_0 \times \left(\frac{m}{50 \text{ MeV}}\right)^9 \frac{\alpha_V}{\alpha} \left(\frac{\kappa}{10^{-10}}\right)^2 \left(\frac{\Delta m/m}{10^{-2}}\right)^{13} \left(\frac{m}{m_V}\right)^4 \end{aligned}$$



# Self-scatterings



$$\frac{\sigma_{\text{eff}}^{\text{SI}}}{m} \equiv R_0 \frac{\sigma_{12}}{m} + \frac{\langle \sigma_{\text{en}v} \rangle}{m v} + \frac{\sigma_{\text{rad}}}{m} \lesssim 1 \text{ cm}^2/\text{g}$$

# Free-streaming Length

**4-to-2**  
'reheat' DM  
Increasing the FSL

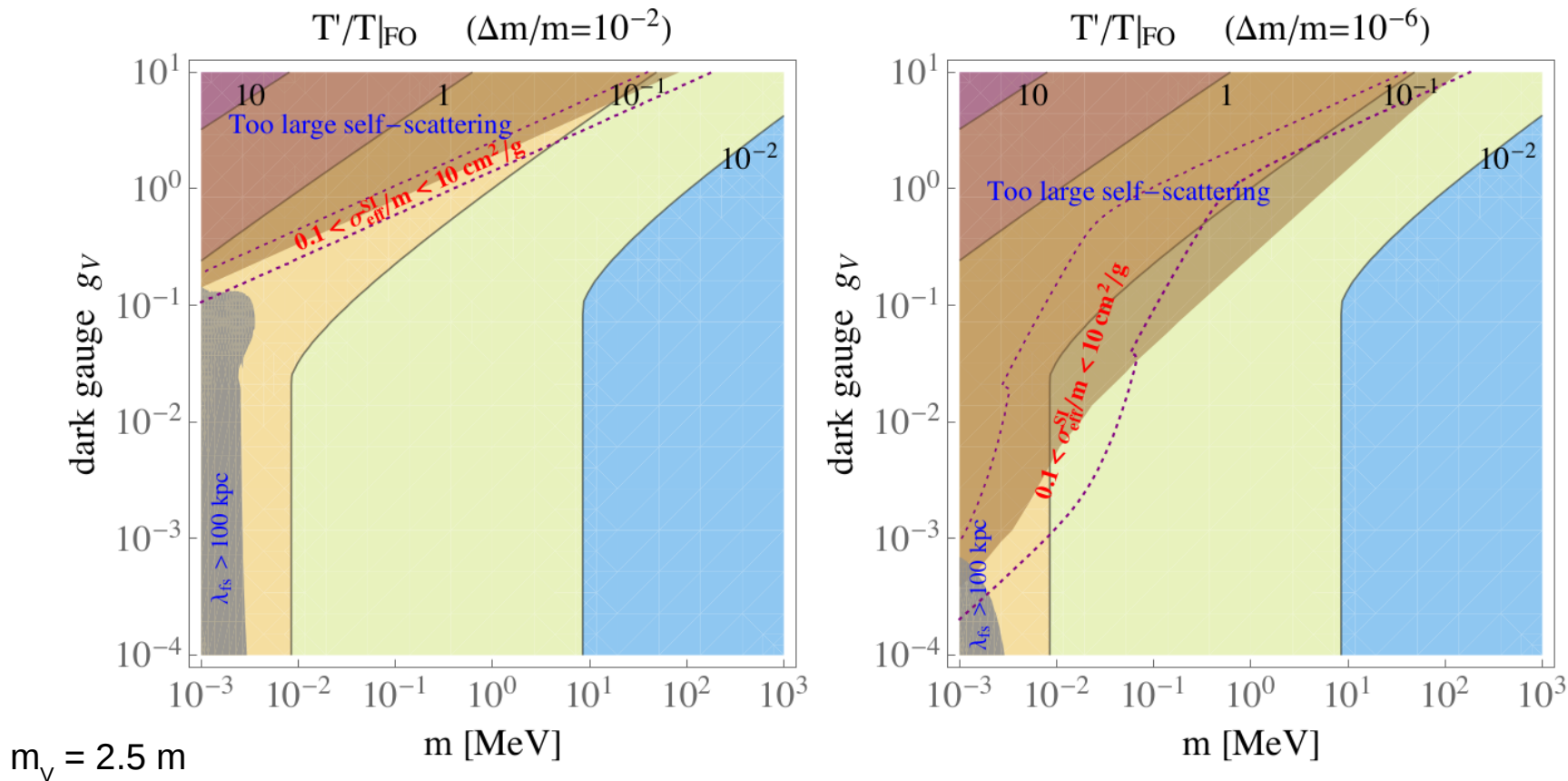
← Versus →

**2-to-2**  
Self-interactions  
Decrease the FSL

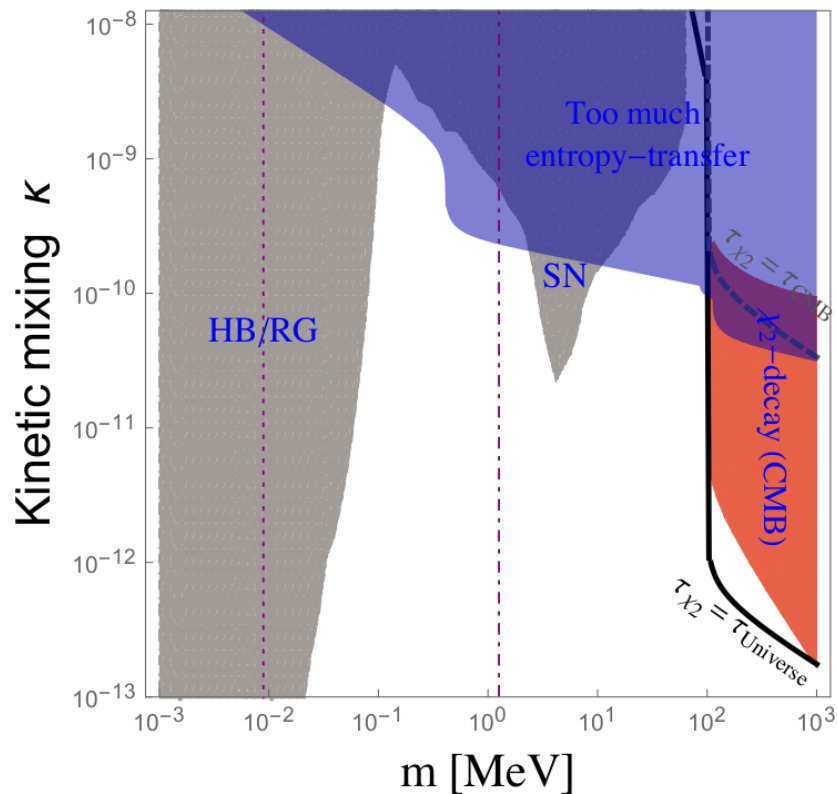
$$\lambda_{\text{fs}} = \int_{t_k}^{t_{\text{eq}}} \frac{v_{\chi}(t)}{a(t)} dt \sim \frac{26 \text{ kpc}}{\sqrt{g_{\star}(T_k)}} \times \frac{10 \text{ keV}}{\sqrt{T_k} m} \left(\frac{T'_k}{T_k}\right)^{1/2} \log_{10} \left(\frac{T_k}{T_{\text{eq}}}\right)$$

$$\lambda_{\text{fs}} \lesssim 100 \text{ kpc} \quad \leftarrow \text{Lyman-}\alpha$$

# Astrophysical implications of Split SIMPs

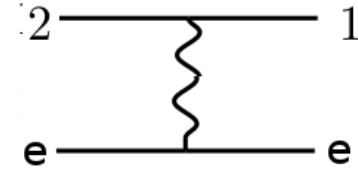
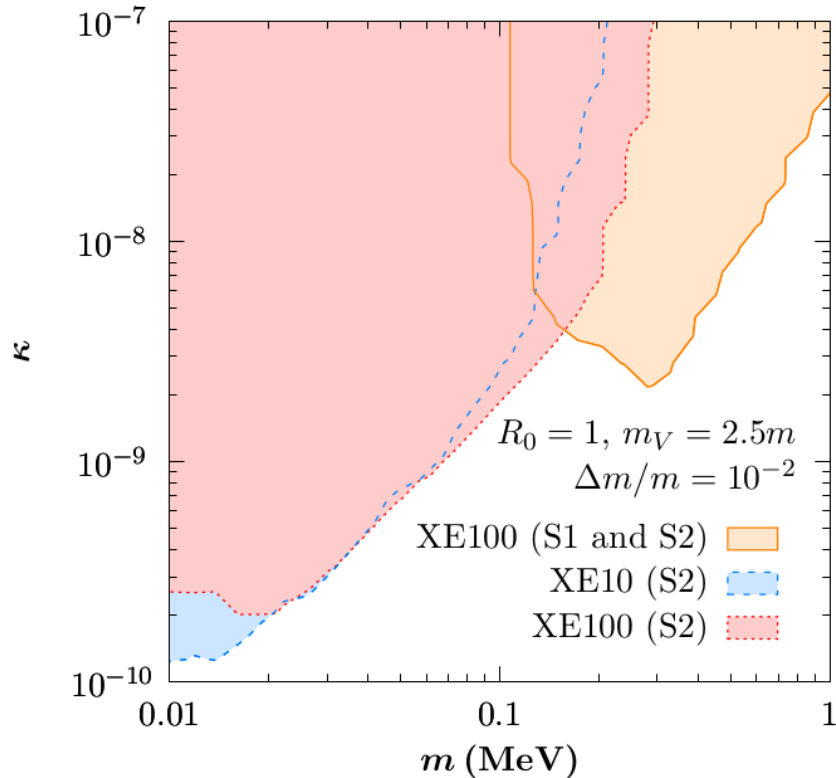


# Constraints on the Kinetic Mixing Portal



$$\Delta m / m = 10^{-2}$$
$$m_{\nu} = 2.5 m$$

# Exothermic DM-Electron Scattering



~ monochromatic signal!

On the verge of being  
probable with reported data  
(S1 and S2) from XENON100!

$$\bar{\sigma}_e = a \frac{16\pi \alpha \alpha_V \kappa^2 \mu_{\chi e}^2}{m_V^4} \simeq 10^{-44} \text{ cm}^2 a \frac{\alpha_V}{\alpha} \left( \frac{\kappa}{10^{-10}} \right)^2 \left( \frac{m}{100 \text{ keV}} \right)^2 \left( \frac{300 \text{ keV}}{m_V} \right)^4$$

# Conclusions

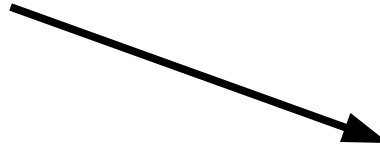
## **Small-scale anomalies**

- \* Cusp-vs-core
- \* Too-big-to-fail

# Conclusions

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- \* Too-big-to-fail



## Self-Interacting Dark Matter

$$\frac{\sigma_{SS \rightarrow SS}}{m_S} \sim 1 \text{ cm}^2/\text{g}$$

# Conclusions

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$$\left. \begin{array}{l} m_S \sim 100 \text{ MeV} \\ \lambda_S \sim 1 \\ \lambda_{HS} < 10^{-3} \end{array} \right\}$$

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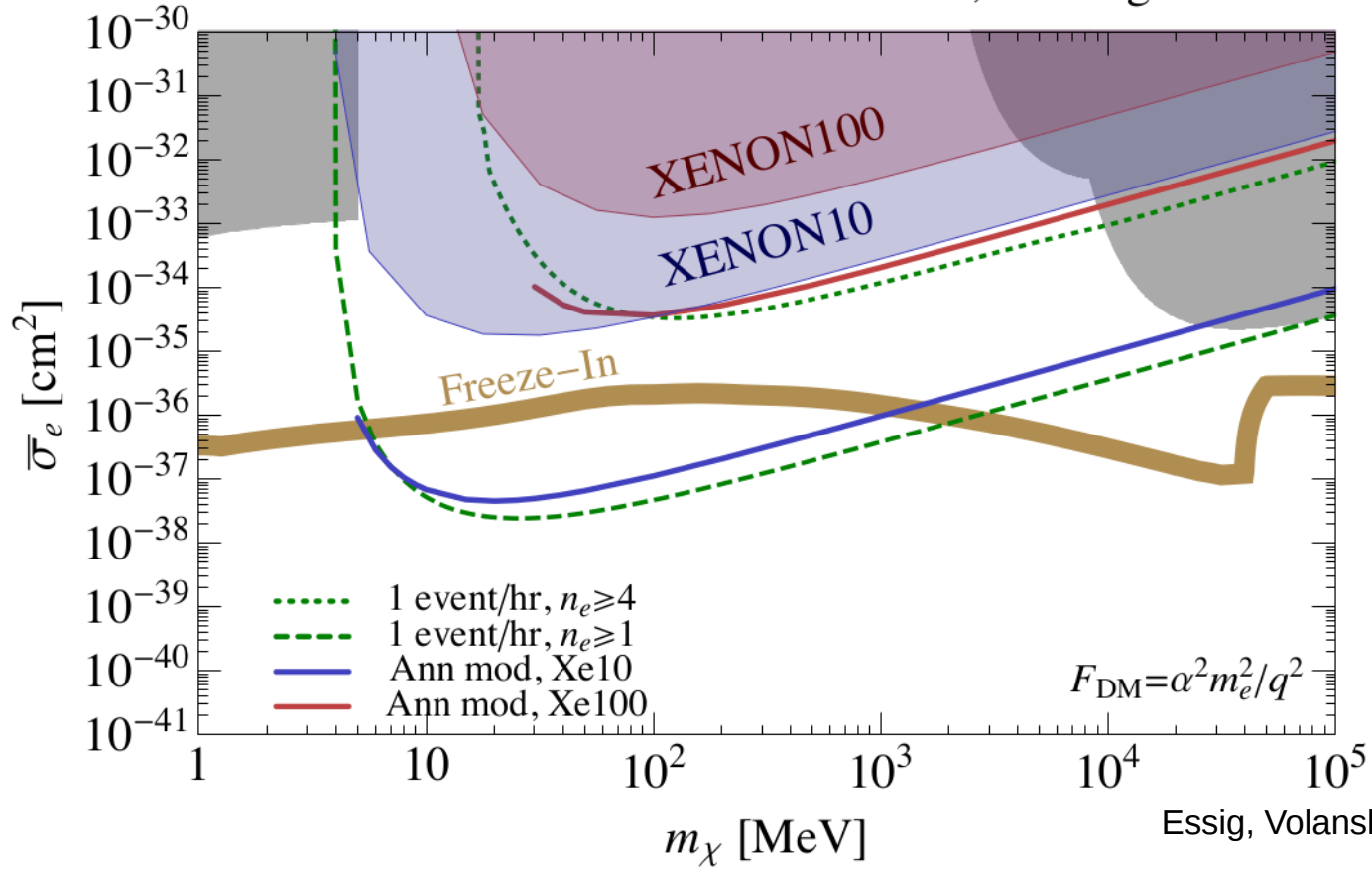
## SIMP DM

- \* dominant  $N \rightarrow n$
- \* need to avoid the 'DM reheating'
- + kinetic equilibrium  $\text{SM} \leftrightarrow \text{DM}$
- + dark sector with relativistic particles @ FO
- + SM and DM never in kinetic equilibrium

# Conclusions

- Self-interacting DM with no light mediators → SIMP DM
- $Z_2$  SIMP DM generated via  $4 \rightarrow 2$  annihilations
- DM: MeV ballpark, ‘large’ self-interactions & ‘small’ portal with the SM
- Difference of temperatures *dynamically* produced via freeze-in!
- Self-interactions: small velocity dependence
  
- SIMPs offer a new window to DM: Points to different physical scales

# Event Rates & Ann. Mod. Search, 1000 kg Xe



Essig, Volansky & Yu '17