

Search of new physics through flavor physics observables

Olcyr Sumensari

Advisor: Damir Bečirević (LPT Orsay/Univ. Paris-Sud)

Journée du laboratoire d'excellence P2IO, November 15, 2017.



UNIVERSITÀ
DEGLI STUDI
DI PADOVA



elusives
neutrinos, dark matter & dark energy physics



This project has received funding from the European Union's Horizon 2020 research and innovation programme under the Marie Skłodowska-Curie grant agreement No 674896.

The Standard Model

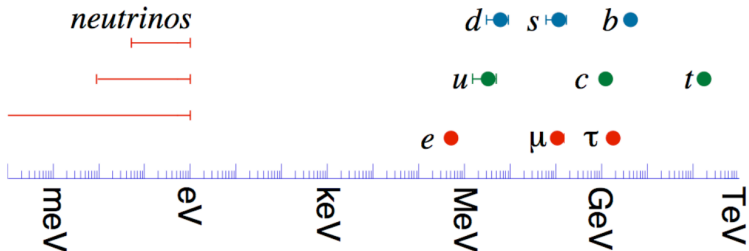
- Gauge sector of the SM entirely **fixed by symmetry**:
 - ⇒ Only a **handful of parameters**.
 - ⇒ Theory renormalizable and **verified** at the **loop-level** (oblique parameters).
- Flavor sector **not fixed by symmetry** :
 - ⇒ 13 free parameters (**masses and quark mixing**) – fixed by data.

$$\mathcal{L}_Y = -Y_\ell \bar{L} \Phi \ell_R - Y_d \bar{Q} \Phi d_R - Y_u \bar{Q} \tilde{\Phi} u_R + \text{h.c.}$$

- ⇒ These (many) parameters exhibit a **hierarchial structure** we do not understand.

The Flavor Problem

- Striking hierarchy \Rightarrow **Flavor theory**?
- Quarks and leptons mix in completely different ways.



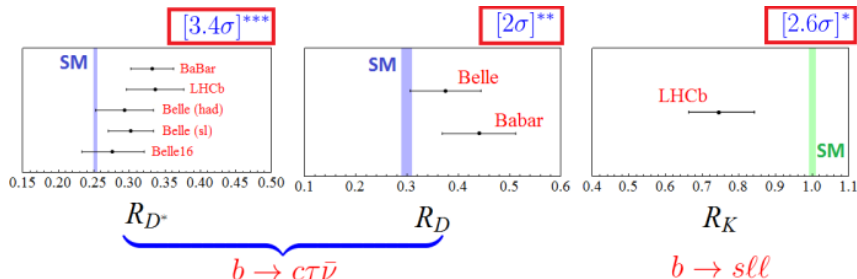
“Who ordered that?”

- Assumption in the SM: the couplings of quarks and leptons with the gauge bosons are **flavor universal** by construction (*up to fermion mass effects*).
- **A few cracks** [$\approx 2 - 3\sigma$] appeared recently in B -meson decays
 - \Rightarrow Violation of **L**epton **F**lavor **U**niversality (LFU)?
 - \Rightarrow **To explain** those observations (*in both tree-level and loop induced decays*), one needs to go beyond the SM.

LFUV in B Decays [pre-2017]

$$R_{D^{(*)}} = \frac{\mathcal{B}(B \rightarrow D^{(*)} \tau \bar{\nu})}{\mathcal{B}(B \rightarrow D^{(*)} \ell \bar{\nu})},$$

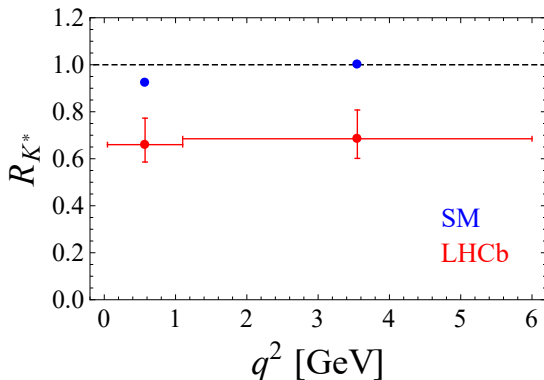
$$R_K = \frac{\mathcal{B}(B^+ \rightarrow K^+ \mu \mu)}{\mathcal{B}(B^+ \rightarrow K^+ e e)} \bigg|_{q^2 \in [1,6] \text{ GeV}^2}$$



- NEW (FPCP17): LHCb, $R_{D^*} = 0.285(35)$.
- NEW: LHCb, $R_{J/\Psi} = 0.71(17)(18)$. Larger than th. predictions (?)

$$R_{K^*} = \frac{\mathcal{B}(B \rightarrow K^* \mu \mu)}{\mathcal{B}(B \rightarrow K^* e e)} \bigg|_{q^2 \in [q_{\min}^2, q_{\max}^2]} \quad [\text{LHCb}, 1705.05802]$$

- **New results** in two bins of q^2 : $[\approx 2.5\sigma]$



TASK:

Using a large ensemble of **flavor physics observables** as constraints (such as $Z \rightarrow \ell\ell$, $W \rightarrow \ell\nu$, $B_s \rightarrow \mu\mu \dots$), we would like to build a **model of New Physics** that can explain the so-called ***B*-physics anomalies**.

\Rightarrow This is a very difficult problem to solve because of the **large amount** of experimental **constraints**!

TASK:

Using a large ensemble of **flavor physics observables** as constraints (such as $Z \rightarrow \ell\ell$, $W \rightarrow \ell\nu$, $B_s \rightarrow \mu\mu \dots$), we would like to build a **model of New Physics** that can explain the so-called **B -physics anomalies**.

\Rightarrow This is a very difficult problem to solve because of the **large amount** of experimental **constraints**!

An important part of my thesis was dedicated to this issue.

\Rightarrow I showed that extending the Higgs sector is not enough.

[P. Arnan, D. Bečirević, F. Mescia, OS. 2017]

\Rightarrow Instead, the scenarios involving a **leptoquark (LQ) boson** can **do the job**. Here I focus on R_K and R_{K^*} .

[D. Bečirević, S. Fajfer, N. Košnik, OS. 2016]

[D. Bečirević, OS. 2017]

The standard strategy of introducing a low-energy LQ boson cannot provide a tree-level solution to R_K and R_{K^*} without causing other problems (*e.g., proton stability and/or flavor constraints*).

c.f. [D. Bečirević, N. Košnik, OS, R. Zukanovich. 2016]

[B. Fornal, B. Grinstein. 2017]

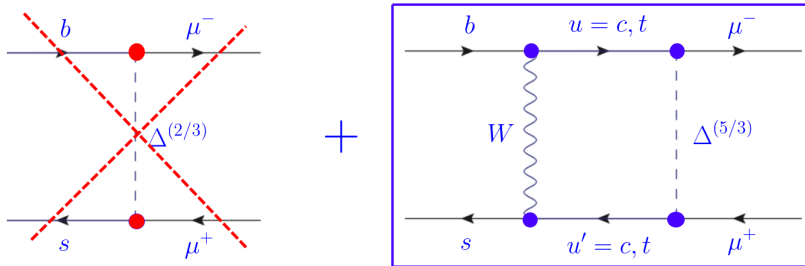
The only scalar LQ which does not disturb the proton stability predicts $R_{K^{(*)}}^{\text{NP}} > R_{K^{(*)}}^{\text{SM}}$, in disagreement with the **LHCb** findings.

I showed that a **peculiar choice** of the **Yukawa couplings** can circumvent this problem and explain R_K and R_{K^*} via loops:

[D. Bečirević, OS. 2017]

$$\mathcal{L}_\Delta = (g_R)_{ij} \bar{Q}_i \Delta \ell_{Rj} + (g_L)_{ij} \bar{u}_{Ri} \tilde{\Delta}^\dagger L_j + \text{h.c.}, \quad \Delta = (\mathbf{3}, \mathbf{2})_{7/6}$$

\Rightarrow We take $g_R = 0$ and $g_L \neq 0$.



\Rightarrow This provides a **viable explanation** which can be **tested experimentally**, e.g. through the direct searches at the LHC.

- We are entering a **precision era** of **flavor physics**: maturity of LQCD and unprecedented precision in flavor experiments.
- **Collective effort** in **flavor experiments** is/will be a **guide to theory**: *NA62, BES-III and LHCb*, and the forthcoming *Belle-II, KOTO, $(g - 2)_\mu$, Mu2E...*
- Interesting **hints of LFU violation** in $R_{K^{(*)}}$ and $R_{D^{(*)}}$:
 \Rightarrow Use the experimental data to build a model of NP and verify its validity in direct searches!
- **Higgs Flavor Era** around the corner?

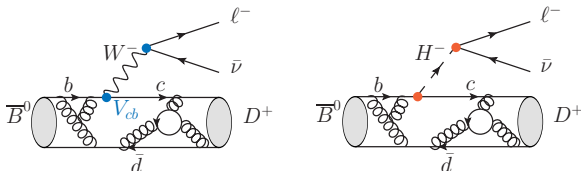
Thank you!

Back-up

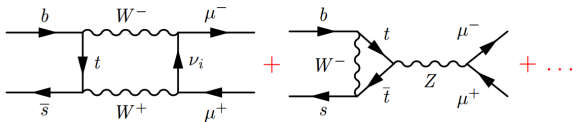
Flavor physics observables

Precision flavor physics: search of deviations w.r.t. the SM predictions

- Flavor changing charged currents: e.g. $b \rightarrow c \tau \nu$



- Flavor changing neutral currents: e.g. $b \rightarrow s \ell \ell$



- Possible mostly due to the maturity of **LQCD** in determining the relevant *hadronic matrix elements* (form factors, decay constants, bag parameters).
- Particularly interesting due to the **deviations** from LFU observed in **B -meson decays**: $B \rightarrow D^{(*)} \ell \bar{\nu}$ ($\ell = e, \mu, \tau$) and $B \rightarrow K^{(*)} \ell \ell$ ($\ell = e, \mu$).

Explaining R_K

Scalar Leptoquark Models

[Becirevic, Kosnik, OS, Zukanovich. 1608.08501]

⇒ Focus on NP couplings to **muons only**

[couplings to electrons are also possible, cf. Hiller, Schmaltz 2014]

$SU(3)_c \times SU(2)_L \times U(1)_Y$:

N.B. $Q = Y + T_3$.

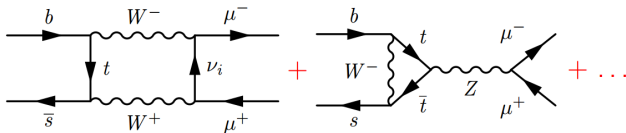
	BNC	Interaction	WC	R_K/R_K^{SM}	$R_{K^*}/R_{K^*}^{\text{SM}}$
$(\bar{3}, 1)_{4/3}$	✗	$\overline{d_R^C} \Delta \ell_R$	$(C_9)' = (C_{10})'$	≈ 1	≈ 1
$(3, 2)_{7/6}$	✓	$\overline{Q} \Delta \ell_R$	$C_9 = C_{10}$	> 1	> 1
$(3, 2)_{1/6}$	✓	$\overline{d_R} \tilde{\Delta}^\dagger L$	$(C_9)' = -(C_{10})'$	< 1	> 1
$(\bar{3}, 3)_{1/3}$	✗	$\overline{Q^C} i\tau_2 \tau \cdot \Delta L$	$C_9 = -C_{10}$	< 1	< 1

⇒ No fully viable model. Triplet can be used, but further symmetries are needed to forbid **proton decay** (see [Dorsner et al. 2017] for a GUT mechanism).

LFU violation

(i) $b \rightarrow s \mu^+ \mu^-$

- FCNC process:



- Form-factor errors cancel out in the ratio \Rightarrow **Extremely clean prediction.**

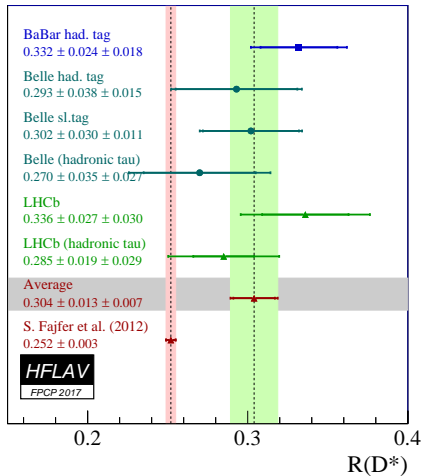
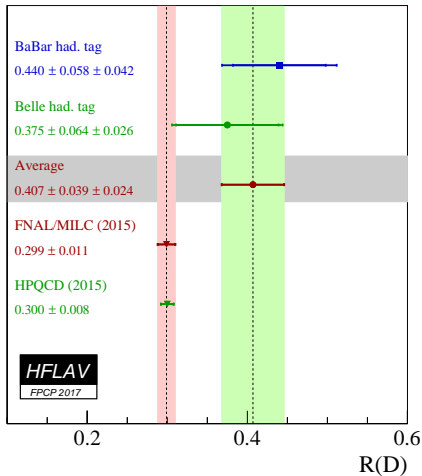
$$R_K \equiv \frac{\mathcal{B}(B^+ \rightarrow K^+ \mu \mu)}{\mathcal{B}(B^+ \rightarrow K^+ e e)} \bigg|_{q^2 \in [1,6] \text{ GeV}^2} \stackrel{\text{SM}}{=} 1.00(1)$$

[Bordone et al. 2016]

- 2.6 σ deviation** observed by LHCb:

$$R_K^{\text{exp}} = 0.745^{+0.090}_{-0.074}(\text{stat}) \pm 0.036(\text{syst})$$

- 2.5 σ deviation** in two bins for $B \rightarrow K^* \mu \mu$: [0.045, 1.1] and [1.1, 6] GeV^2 .



- **3.9 σ combined** deviation from the SM [theory error under control?]
- **2.2 σ** deviation if **only R_D** is considered.
- 2σ deviation in $R_{J/\Psi}$?

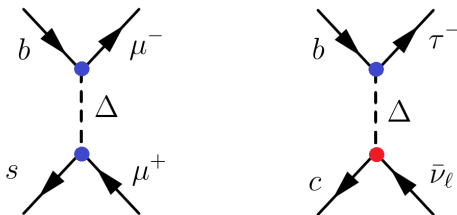
Theory Challenge

A SLQ Model for R_K and R_D

[Becirevic, Fajfer, Kosnik, OS. 1608.08051]

We can also explain R_D if a **new ingredient** is added to the model
 $\Delta^{1/6} = (\mathbf{3}, \mathbf{2})_{1/6}$: three light RH neutrinos ν_R .

$$\mathcal{L}_Y = Y_{ij}^L \bar{L}_i \tilde{\Delta}^{(1/6)} d_{Rj} + Y_{ij}^R \bar{Q}_i \Delta^{(1/6)} \nu_{Rj} + \text{h.c.}$$



For $b \rightarrow c \tau \bar{\nu}$ $\Rightarrow |\mathcal{M}(B \rightarrow D^{(*)} \ell \nu)|^2 = |\mathcal{M}_{\text{SM}}|^2 + |\mathcal{M}_{\text{NP}}|^2$.

Naturally generates $R_{D^{(*)}}^{\text{NP}} > R_{D^{(*)}}^{\text{SM}}$ if $|Y_{b\tau}^L| \gtrsim |Y_{b\mu}^L|$.

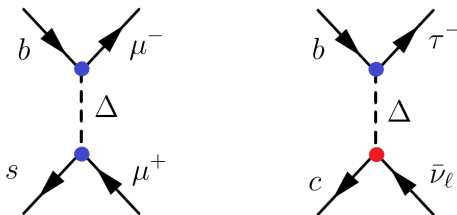
Theory Challenge

A SLQ Model for R_K and R_D

[D. Becirevic, S. Fajfer, N. Kosnik, OS. 1608.08501]

We can also explain R_D if a **new ingredient** is added to the model
 $\Delta^{1/6} = (3, 2)_{1/6}$: three light RH neutrinos ν_R .

$$\mathcal{L}_Y = Y_{ij}^L \bar{L}_i \tilde{\Delta}^{(1/6)} d_{Rj} + Y_{ij}^R \bar{Q}_i \Delta^{(1/6)} \nu_{Rj} + \text{h.c.}$$



- **Passed all flavor tests:** $\mathcal{B}(B_s \rightarrow \mu^+ \mu^-)$, $\mathcal{B}(B \rightarrow K \mu \mu)_{\text{high } q^2}$, Δm_{B_s} , $\mathcal{B}(B \rightarrow \tau \bar{\nu})$, $\mathcal{B}(D_s \rightarrow \tau \bar{\nu})$, $\mathcal{B}(B \rightarrow K \nu \bar{\nu})$, $\mathcal{B}(B \rightarrow K \mu \tau)$ etc.
- Many experimental signatures for LHCb and Belle-2.

Explaining R_K

EFT approach

If the LFUV takes place at scales well above EWSB, then use OPE:

$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \left[\sum_{i=1}^6 C_i(\mu) \mathcal{O}_i(\mu) + \sum_{i=7,8,9,10,P,S,\dots} \left(C_i(\mu) \mathcal{O}_i + C'_i(\mu) \mathcal{O}'_i \right) \right]$$

- Operators relevant to $b \rightarrow s \ell \ell$ are

$$\begin{aligned} \mathcal{O}_9^{(\prime)} &= (\bar{s} \gamma_\mu P_{L(R)} b) (\bar{\ell} \gamma^\mu \ell), & \mathcal{O}_{10}^{(\prime)} &= (\bar{s} \gamma_\mu P_{L(R)} b) (\bar{\ell} \gamma^\mu \gamma^5 \ell), \\ \mathcal{O}_S^{(\prime)} &= (\bar{s} P_{R(L)} b) (\bar{\ell} \ell), & \mathcal{O}_P^{(\prime)} &= (\bar{s} P_{R(L)} b) (\bar{\ell} \gamma_5 \ell), \\ \mathcal{O}_7^{(\prime)} &= m_b (\bar{s} \sigma_{\mu\nu} P_{R(L)} b) F^{\mu\nu} \quad \dots \end{aligned}$$

- To explain $R_{K^{(*)}}^{\text{exp}} < R_{K^{(*)}}^{\text{SM}}$, one needs effective coefficients C_9, C_{10} .