Breaking Supersymmetry Made Easy, Viable, and Generic Hitoshi Murayama (Berkeley) XLIInd Rencontres de Moriond, March 11, 2007

with Yasunori Nomura hep-ph/0612186, 0701231

LHC is coming! Reaching the important energy scale G_F^{-1/2}=300 GeV known since 1933 paper by Fermi. Historic moment!

LHC is coming! Reaching the important energy scale G_F^{-1/2}=300 GeV known since 1933 paper by Fermi. Historic moment!
Growing concern in the community
If there is new physics below TeV, we should have seen its hints by now. Most likely we don't find anything at the LHC.

 LHC is coming! Reaching the important
 energy scale $G_F^{-1/2}$ =300 GeV known since 1933 paper by Fermi. Historic moment! Growing concern in the community If there is new physics below TeV, we should have seen its hints by now. Most likely we don't find anything at the LHC. Now I think It is quite likely to find new physics, especially supersymmetry, at the LHC!













4



4







Why supersymmetry? "hierarchy problem"

- At the end of 19th century: a "crisis" about electron
 - Solution Like charges repel: hard to keep electric charge in a small pack

Selectron is point-like <10⁻¹⁷cm

Need a lot of energy to keep it small!

 $\Delta m_e c^2 \sim \frac{e^2}{r_e} \sim \text{GeV} \frac{10^{-17} \text{cm}}{r_e} \qquad m_e^0 = -3.141082 \text{ GeV}$ $\Delta m_e c^2 \sim \frac{e^2}{r_e} \sim \text{GeV} \frac{10^{-17} \text{cm}}{r_e} \qquad \Delta m_e c^2 = -3.141082 \text{ GeV}$ $\Delta m_e c^2 = -3.141082 \text{ GeV}$ $\Delta m_e c^2 = -3.141082 \text{ GeV}$ $\Delta m_e c^2 = -3.141082 \text{ GeV}$ $m_e c^2 = -3.141082 \text{ GeV}$

Selectron creates a force to repel itself Solution Vacuum bubble of matter anti-matter creation/annihilation Selectron annihilates the positron in the bubble \Rightarrow only 10% of mass even for Planck-size $r_{e} \sim 10^{-33} cm$

Selectron creates a force to repel itself Solution Vacuum bubble of matter anti-matter creation/annihilation Selectron annihilates the positron in the bubble \Rightarrow only 10% of mass even for Planck-size $r_{e} \sim 10^{-33} cm$

Selectron creates a force to repel itself Solution Vacuum bubble of matter anti-matter creation/annihilation Selectron annihilates the positron in the bubble \Rightarrow only 10% of mass even for Planck-size $r_{e} \sim 10^{-33} cm$

Selectron creates a force to repel itself Solution Vacuum bubble of matter anti-matter creation/annihilation Selectron annihilates the positron in the bubble \Rightarrow only 10% of mass even for Planck-size $r_{e} \sim 10^{-33} cm$

 $\Delta m_e \sim m_e \frac{\alpha}{4\pi} \log(m_e r_e)$

Higgs repels itself, too

Just like electron repelling itself becaus of its charge, Higgs boson also repels itself Requires a lot of energy to contain itself in its point-like size! Breakdown of theory of weak force Can't discuss physics
 < TeV⁻¹=10⁻¹⁷cm

History repeats itself?

⊘ Double #particles
 again ⇒

superpartners Vacuum bubbles" of superpartners cancel the energy required to contain Higgs boson in itself Standard Model made consistent with whatever physics at shorter distances

 $\Delta m_{H}^{2} \sim \frac{\alpha}{4\pi} m_{SUSY}^{2} \log(m_{H}r_{H})$ Want m_{susy}<TeV for
>1% tuning

But SUSY is broken

For each term in the superpotential $W_{MSSM} = Y_u^{ij} Q_i u_j^c H_u + Y_d^{ij} Q_i d_j^c H_d + Y_l^{ij} L_i e_j^c H_d + \mu H_u H_d$ we can have the "A-terms" and "B-term" $A_u^{ij}Y_u^{ij}Q_iu_j^cH_u + A_d^{ij}Y_d^{ij}Q_id_j^cH_d + A_l^{ij}Y_l^{ij}L_ie_j^cH_d + B\mu H_uH_d$ Scalar masses for all scalars $m_{Qij}^2 \tilde{Q}_i^* \tilde{Q}_j + m_{uij}^2 \tilde{u}_i^* \tilde{u}_j + m_{dij}^2 \tilde{d}_i^* \tilde{d}_j + m_{Lij}^2 \tilde{L}_i^* \tilde{L}_j + m_{eij}^2 \tilde{e}_i^* \tilde{e}_j + m_{H_d}^2 |H_d|^2 + m_{H_u}^2 |H_u|^2$ gaugino mass for all three gauge factors $M_1 \tilde{B}\tilde{B} + M_2 \tilde{W}^a \tilde{W}^a + M_3 \tilde{g}^a \tilde{g}^c$ $A(18x3)+B(2)+m(9x5+2)+M(2x3)+\mu(2)=111$ $U(1)_{R} \times U(1)_{PQ}$ removes only two phases cf. SM has two params in the Higgs sector 107 more parameters than the SM!

"Typical" parameters excluded "typical" parameters of SUSY breaking too large FCNC and CP violation Made worse after beautiful B data

10

"Typical" parameters excluded "typical" parameters of SUSY breaking \Rightarrow too large FCNC and CP violation Made worse after beautiful B data $(\tilde{d}, \tilde{s}, \tilde{b})^* m_0^2 \begin{pmatrix} 1 & (\delta_{12}^d) & (\delta_{13}^d) \\ (\delta_{21}^d) & 1 & (\delta_{23}^d) \\ (\delta_{31}^d) & (\delta_{22}^d) & 1 \end{pmatrix} \begin{pmatrix} \tilde{d} \\ \tilde{s} \\ \tilde{b} \end{pmatrix}$ $(\delta_{12}^d) < 0.001 \frac{m_{SUSY}}{500 \text{ GeV}}$

10

"Typical" parameters excluded "typical" parameters of SUSY breaking \Rightarrow too large FCNC and CP violation Made worse after beautiful B data \odot EDM $H\propto ec{s}\cdotec{E}$ $(\tilde{d}, \tilde{s}, \tilde{b})^* m_0^2 \begin{pmatrix} 1 & (\delta_{12}^d) & (\delta_{13}^d) \\ (\delta_{21}^d) & 1 & (\delta_{23}^d) \\ (\delta_{31}^d) & (\delta_{22}^d) & 1 \end{pmatrix} \begin{pmatrix} \tilde{d} \\ \tilde{s} \\ \tilde{b} \end{pmatrix}$ stringent limits on electron, neutron, and $(\delta_{12}^d) < 0.001 \frac{\overline{m_{SUSY}}}{500 \text{ GeV}}$ Hg atom @ either msusy>TeV or ď phase<10-2 (δ^d₁₂) **200000000** S

 \overline{K}^0

10

d

 K^0

Accept heavy SUSY > 100 TeV
 the hierarchy problem fine-tuned > 10⁶!

11

Accept heavy SUSY > 100 TeV
 the hierarchy problem fine-tuned > 10⁶!
 Tune SUSY breaking flavor-blind, CP
 probability for viable parameter set 10⁻³_K×10⁻³_B×10⁻³_{µ→ey}×10⁻²_{EDM}×···?

Accept heavy SUSY > 100 TeV \odot the hierarchy problem fine-tuned > 10⁶! Tune SUSY breaking flavor-blind, CP ø probability for viable parameter set 10^{-3} K × 10^{-3} B × 10^{-3} $\mu \rightarrow e_{\rm Y}$ × 10^{-2} EDM × ...? Build an elaborate model to get flavor-blind and CP-conserving SUSY breaking elaborate model = delicate artwork = unlikely choice by Mother Nature (?)

Accept heavy SUSY > 100 TeV Defeatism Tune SUSY breaking flavor-blind, CP ø probability for viable parameter set 10^{-3} K × 10^{-3} B × 10^{-3} $\mu \rightarrow e_{\rm Y}$ × 10^{-2} EDM × ...? Build an elaborate model to get flavor-blind and CP-conserving SUSY breaking elaborate model = delicate artwork = unlikely choice by Mother Nature (?)

Accept heavy SUSY > 100 TeV Defeatism The hierarchy problem fine-tuned > 10⁶! Tune SUSY breaking flavor-blind, CP Selfø probability for viable parameter set righteous 10^{-3} K × 10^{-3} B × 10^{-3} U → eV × 10^{-2} EDM × ···? Build an elaborate model to get flavor-blind and CP-conserving SUSY breaking elaborate model = delicate artwork = unlikely choice by Mother Nature (?)

Accept heavy SUSY > 100 TeV Defeatism The hierarchy problem fine-tuned > 10⁶! Tune SUSY breaking flavor-blind, CP Selfprobability for viable parameter set righteous 10^{-3} K × 10^{-3} B × 10^{-3} $\mu \rightarrow e_{\rm Y}$ × 10^{-2} EDM × ...? Build an elaborate model to get flavor-blind and CP-conserving SUSY breaking Intelligent elaborate model = delicate artwork Design = unlikely choice by Mother Nature (?)

Special Model I SUSY Breaking

Breaking SUSY has been difficult
Nelson-Seiberg: you need either

non-generic superpotential
need exact U(1)_R spontaneously broken

Either way, theory needs to be rather spcial, not a whole lot of models known

Special Model I SUSY Breaking

Breaking SUSY has been difficult
Nelson-Seiberg: you need either
non-generic superpotential
need exact U(1)_R spontaneously broken
Either way, theory needs to be rather spcial, not a whole lot of models known

	SU(6)	U(1)	$U(1)_m$	$U(1)_R$	
A	15	+2	0	$-\frac{18}{7}$	
F	6	-5	0	$-\frac{18}{7}$	
\bar{F}^{\pm}	$\overline{6}$	-1	± 1	$\frac{16}{7}$	$W = A\bar{F}$
$ar{F}^0$	$\overline{6}$	-1	0	$\frac{16}{7}$	
S^{\pm}	1	+6	± 1	$\frac{16}{7}$	
S^0	1	+6	0	$\frac{16}{7}$	

 $W = A\bar{F}^{+}\bar{F}^{-} + \bar{F}^{0}(F^{+}S^{-} + F^{-}S^{+}) + FF^{0}S^{0}$

Dine-Nelson-Nir-Shirman

 $\mu \sim 10^2 - 10^3 \text{ GeV}$

Standard

Model

Mff SUSY SM

16

Our Generic Scheme $\frac{1}{M_{Pl}}\bar{Q}Q\bar{f}f$

SUSY SM

 $M\bar{f}f$

16

 $m_Q \bar{Q} Q$

Our Generic Scheme $\frac{1}{M_{Pl}}\bar{Q}Q\bar{f}f$

SUSY QCD SU(Nc), SO(Nc), Sp(Nc)

$M\bar{f}f$ SUSY SM

no U(1)_R symmetry imposed most general superpotential wide choice of gauge groups, matter content $N_c < N_{f} < \frac{3}{2}N_c$

 $m_Q \bar{Q} Q$

How it works (Most technical slide) • SUSY SU(N_c) QCD N_c<N_f<3N_c/2 $W = m_O^{ij} \bar{Q}_i Q_j$ \odot low-energy free magnetic theory (m_Q< Λ) $W = m_O^{ij} \Lambda M_{ij} + M_{ij} \bar{q}^i q^j$ • SUSY breaking $@M_{ij} = 0, \quad \frac{\partial W}{\partial M_{ij}} = m_Q^{ij} \neq 0$ Socal minimum with long lifetime $W = \frac{1}{M_{Pl}} \bar{Q} Q \bar{f} f$ Generates SUSY breaking in f, fbar their loops⇒gauge mediation

Good news for string theory

- String theory does not predict unique solution
- "Landscape" of possibilities for gauge groups, matter content, number of SUSY
- We at least need SM
- We tend to get extra "junks", i.e. extra gauge groups, extra vector-like matter
 the "junks" are precisely what we need to break SUSY via gauge mediation Easy, Viable, Generic!

Experimental Consequences

Serve wide range of parameters and models SUSY dark matter = gravitino? Iong-lived charged particle (stau)? "LSP" decay into photon+gravitino? Specific mass spectrum of SUSY particles In principle depends on "hidden" sector other dark matter candidates

Conclusion

SUSY breaking happens in a wide class of models expected in the "landscape" of supersymmetric theories (easy, generic)

Most of them provide successful gauge mediation (viable)

flavor-blind and CP-conserving SUSY
It is completely OK to have SUSY<TeV
Looking forward to Tevatron/LHC!