



Measurement of EPRtype flavour entanglement in $\Upsilon(4S) \rightarrow B^0 \overline{B}^0$ decays

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quant-ph/0702267 submitted to PRL

The EPR argument (1935)

MAY 15, 1935

PHYSICAL REVIEW

Can Quantum-Mechanical Description of Physical Reality Be Considered Complete?

A. EINSTEIN, B. PODOLSKY AND N. ROSEN, Institute for Advanced Study, Princeton, New Jersey (Received March 25, 1935)

In a complete theory there is an element corresponding to each element of reality. A sufficient condition for the reality of a physical quantity is the possibility of predicting it with certainty, without disturbing the system. In quantum mechanics in the case of two physical quantities described by non-commuting operators, the knowledge of one precludes the knowledge of the other. Then either (1) the description of reality given by the wave function in quantum mechanics is not complete or (2) these two quantities cannot have simultaneous reality. Consideration of the problem of making predictions concerning a system on the basis of measurements made on another system that had previously interacted with it leads to the result that if (1) is false then (2) is also false. One is thus led to conclude that the description of reality as given by a wave function is not complete.

- Definition of "element of reality". A "complete theory" must contain a representation for each element of reality;
- EPR consider an "entangled system" of two particles and the measurement of two non-commuting observables (position and momentum);
- Entanglement is used to transport the information from one sub-system to the other;
- EPR identify a contradiction with the QM rule for non-commuting observables.

Bohm (1951), entangled states



* This will happen even if the decision to orient the polarizer for particle **a** is done at the very last moment => no causal connection.

How to cure the problem ?

- with the introduction of a new instantaneous communication channel between the two sub-systems...

- or with the introduction of some new hidden information for particle **b**, so that the particle knows how to behave.

=> QM is incomplete.

Bell (1964), QM vs local models,...

This problem was revitalized in 1964, when Bell suggested a way to distinguish QM from local models featuring hidden variables (J. S. Bell, Physics 1, 195 (1964)).

Several experiments have been done with photon pairs, atoms,...

Tests have been carried out on correlated $K^0 \overline{K}^0$

Apostolakis *et al.*, CPLEAR collab., Phys. Lett. B **422**, 339 (1998) Ambrosino *et al.*, KLOE collab., Phys. Lett. B **642**, 315 (2006).



Study of $\Upsilon(4s) \rightarrow B^0 \overline{B^0}$

The $\Upsilon(4s) \rightarrow B^0 \overline{B}{}^0$ is another case of entangled system: the pair flavour wave function is

$$|\Psi_{\Upsilon(4s)}\rangle = (1/\sqrt{2}) (|\mathbf{B}^0\rangle_a |\overline{\mathbf{B}}^0\rangle_b - |\overline{\mathbf{B}}^0\rangle_a |\mathbf{B}^0\rangle_b)$$

decays occurring at the same proper time are fully correlated: the flavor-specific decay of one meson fixes the (previously undetermined) flavour of the other meson.

⇒ we use the KEKB / BELLE data to explore for the first time this sector

Study of correlated $B^0 \overline{B}{}^0$



N.B. : production vertex position Z_0 not very well known : only ΔZ is available !

Predictions from QM for entangled pairs

Time (Δt)-dependent decay rate into two Opposite Flavour (OF) states

 $R_{OF} \propto 1 + \cos(\Delta m_d \Delta t)$

idem, into two Same Flavour (SF) states

 $R_{SF} \propto 1 - \cos(\Delta m_d \Delta t)$

 Δm_d is the mass difference of the two mass eigenstates



(ignoring CP violation effects O(10⁻⁴), and taking $\Delta \Gamma_d = 0$)

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Local Realism by Pompili & Selleri (PS)

Local Realism, each B has "elements of reality" (hidden variables) $\lambda_1 : CP = +1 \text{ or } -1$ $\lambda_2 : Flavour = +1 \text{ or } -1$ $=> 4 \text{ basic states } B^0_H, \overline{B}^0_H, B^0_L, \overline{B}^0_L$ indexed by i = 1, 2, 3, 4

* Mass states are stable in time, simultaneous anti-correlated flavor jumps.

The model works with probabilities $p_{ij}(t|0) = \text{prob for a B to be}$ in the state j at proper time t=t, conditional of having been in state i at t=0. * p_{ij} set to be consistent with single B⁰ evolution ~ exp{($\Gamma/2 + im$)t}. * PS build a model with a minimal amount of assumptions

- \Rightarrow They only determine upper and lower limits for combined probabilities ...

F. Selleri, Phys. Rev. A 56 (1997) 3493 A. Pompili, and F. Selleri, Eur. Phys. J. C 14 (2000) 469

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Local Realism by Pompili & Selleri (2)

=> analytical expressions for A corresponding to the limits. The Amax is



Spontaneous immediate Disentanglement (SD)

Just after the decay into opposite flavor states, we considers an independent evolution for the B^0 pair

$$A_{SD}(t_1, t_2) = \cos(\Delta m_d t_1) \cos(\Delta m_d t_2) = \frac{1}{2} [\cos(\Delta m_d (t_1 + t_2)) + \cos(\Delta m_d \Delta t)]$$

integrating out $t_1 + t_2$ gives:



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Analysis goals and methods

We want to provide FULLY CORRECTED time-dependent asymmetry.

For this, we will

- -- subtract all backgrounds
- -- correct for events with wrong flavour associations
- -- correct for the detector effects (resolution in Δt) by a deconvolution procedure

=> the result can then be directly compared to the models.

We will use our data to test

- the Pompili and Selleri model,
- the Spontaneous Disentanglement model,
- and we will check for some partial contamination by SD-like events, i.e. we search for decoherence effects from New Physics



152 10⁶ $B^0\overline{B}^0$ pairs

The Belle Detector





Event selection and tagging



* All the remaining tracks are used to guess the flavour of second B, from the standard Belle flavour tagging procedure. We select the best purity subset, from semileptonic decays

 $B \rightarrow X \ell \nu$

Event selection and tagging

After selection, we obtain 6718 OF and 1847 SF events. The Δz is obtained from track fit of the two vertices and converted into a Δt value.

We obtain the OF and SF distributions, with 11 variable-size bins (to account to the fast falling statistics)



$\Delta t \mathrm{bin}$	window [ps]
1	0 - 0.5
2	0.5 - 1
3	1 - 2
4	2 - 3
5	3 - 4
6	4 - 5
7	5 - 6
8	6 - 7
9	7 - 9
10	9 - 13
11	13 - 20

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Background and wrong flavour tags

We correct bin by bin the OF and SF distributions for the following sources of background:

- Fake D*
- Uncorrelated D*-leptons, mainly D* from one B⁰ and the lepton from the other
- $B^{\pm} \rightarrow D^{**} l v$

		OF events	SF events
	Selected	6718	1847
We also correct for a ~1.5% fraction of wrong flavour associations	Fake D^*	-126	- 54
	Uncor. D^*l	-78	-237
	B^{\pm}	-254	-1
	Wrg Flv	-22 + 86	- 86 +22
	N events	6324	1619



Time-dependent asymmetry before and after background subtraction



Data deconvolution

Deconvolution is performed using response matrices for OF and SF distributions. The two 11x11 matrices are build from GEANT MC events. We use a procedure based on singular value decomposition, from H. Höcker and V. Kartvelishvili, NIM A 372 469 (1996).

Toy MC of the 3 models (QM, PS and SD) have been used to study the method and to estimate the associated systematic error.

	Window	Asymmetry				Systematic	errors	
Δt bin	window [ps]	A and total error	stat. error	total	event sel.	bkgd sub.	wrong tags	deconvol.
1	0.0 - 0.5	1.013 ± 0.028	0.020	0.019	0.005	0.006	0.010	0.014
2	0.5 - 1.0	0.916 ± 0.022	0.015	0.016	0.006	0.007	0.010	0.009
3	1.0 - 2.0	0.699 ± 0.038	0.029	0.024	0.013	0.005	0.009	0.017
4	2.0 - 3.0	0.339 ± 0.056	0.047	0.031	0.008	0.005	0.007	0.029
5	3.0 - 4.0	-0.136 ± 0.075	0.060	0.045	0.009	0.009	0.007	0.042
6	4.0 - 5.0	-0.634 ± 0.084	0.062	0.057	0.021	0.014	0.013	0.049
7	5.0 - 6.0	-0.961 ± 0.077	0.060	0.048	0.020	0.017	0.012	0.038
8	6.0 - 7.0	-0.974 ± 0.080	0.060	0.053	0.034	0.025	0.020	0.025
9	7.0 - 9.0	-0.675 ± 0.109	0.092	0.058	0.041	0.027	0.022	0.022
10	9.0 - 13.0	0.089 ± 0.193	0.161	0.107	0.067	0.063	0.038	0.039
11	13.0 - 20.0	0.243 ± 0.435	0.240	0.363	0.145	0.226	0.080	0.231

The result is given here:

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Before to compare with the models, a cross check with the B⁰ lifetime...

Add OF+SF distributions and fit for τ_{B^0}



Comparison with QM

Least-square fits including a term taking the world-average Δm_d into account. To avoid bias we <u>discard BaBar and BELLE</u> <u>measurements</u>, giving $\langle \Delta m_d \rangle = (0.496 \pm 0.013) \text{ ps}^{-1}$



Comparison with PS model

Fit data to PS model, using the closest boundary. We conservatively assign <u>a null deviation when data falls between boundaries</u>



Comparison with SD model



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Search for New Physics: Decoherence

Decoherent fraction into B⁰, \overline{B}^0 by fitting $(1 - \lambda_{B_d})A_{QM} + \lambda_{B_d}A_{SD}$

We obtain $\lambda_{Bd} = 0.029 \pm 0.057$ => consistent with no decoherence

Previous measurements in K⁰ system:

•From CPLEAR measurement: Phys. Lett. B **422**, 339 (1998) Bertlmann *et al.* Phys. Rev. D **60** 114032 (1999) has deduced $\lambda_{K0} = 0.4 \pm 0.7$ •KLOE $\lambda_{K0} = (0.10 \pm 0.22 \pm 0.04) 10^{-5}$

CONCLUSION

We have performed the first experimental test of the EPR-type flavour entanglement in $\Upsilon(4s) \rightarrow B^0 \overline{B}{}^0$ decays.

We have measured the time-dependent asymmetry due to flavour oscillation. The asymmetry has been corrected for the experimental effects and can be used directly to compare with the different theoretical models.

* The asymmetry is consistent with QM predictions * The local realistic model of Pompili and Selleri is disfavoured at the level of 5.1σ .

* A model with immediate disentanglement into flavour eigenstates is excluded by 13σ .

* A decoherent fraction into flavour eigenstates was found to be 0.029 ± 0.057 , consistent with no decoherence.

BACK UP SLIDES

Decoherence

Decoherence in B^0 , \overline{B}^0

 $\begin{array}{lll} A = (1-\lambda)A_{QM} + \lambda A_{SD} & \Longrightarrow & \lambda_{Bd} = 0.029 \pm 0.057 \\ \mbox{Previous mesurements in } K^0 \mbox{ system:} \\ \bullet \mbox{From CPLEAR measurement, PLB 422, 339 (1998), Bertlmann et} \\ al. \mbox{ PRD 60 114032 (1999) has deduced } \lambda_{K0} = 0.4 \pm 0.7 \\ \bullet \mbox{KLOE } \lambda_{K0} = (0.10 \pm 0.22 \pm 0.04) \ 10^{-5} \end{array}$

P. Heberard's ζ parameter for decoherence in B_H , B_L : $A = (1-\zeta)A_{QM} \implies \zeta_{Bd} = 0.004 \pm 0.017 \pm ...$ At present, this result is not considered robust enough.

Previous mesurements in K⁰ system: •CPLEAR $\zeta_{K^0} = 0.13^{+0.16}_{-0.15}$ •KLOE $\zeta_{K^0} = 0.018 \pm 0.040 \pm 0.007$

Event selection

Semileptonic B side:

Variable	Cuts
P_{lepton}^{CMS}	$1.4 GeV/c < P_{lepton}^{CMS} < 2.4 GeV/c$
Slow π vertex constr. to B	$\chi^2/dof < 100$
K/π likelihood for π	$K3\pi$ mode: $Prob_{K/\pi} < 0.5$
	$K\pi, K\pi\pi^0$ mode: $Prob_{K/\pi} < 0.3$
P_{π^0}	$P_{\pi^0} > 0.2 GeV/c$
P_{γ}	$P_{\gamma} > 0.08 \ GeV/c$
$P_{K,\pi}$ (K3 π mode)	$P_{K,\pi} > 0.2 GeV/c$
Impact parameters	$ dr_{IP} < 0.2 \mathrm{cm}$
$\cos(\theta_{B,D^*l})$	$ \cos(\theta_{B,D^*l}) < 1.1$
D^0 mass	$-37MeV/c^2 < M_{K\pi\pi^0} - M_{D^0} < 23MeV/c^2$
	$-13MeV/c^2 < M_{K\pi,K3\pi} - M_{D^0} < 13MeV/c^2$
$M_{D^*} - M_{D^0}$	$144.4 MeV/c^2 < M_{D^*} - M_{D^0} < 146.4 MeV/c^2$
$P_{D^*}^{CMS}$	$P_{D^*}^{CMS} < 2.6 GeV/c$
B^0 vertex	$\chi^2/dof < 75$
B_{tag} vertex	$\chi^2/dof < 75$

All other tracks are used to identify the flavor of the accompanying B.



Background: Wrong D*l combination



Wrong D*1: consistency checks



Background: $B^{\pm} \rightarrow D^{**} \ell \nu$

 $B^{0} \rightarrow D^{**^{-}} \ell^{+} \nu \quad \text{has flavor mixing, signal} \\ B^{+} \rightarrow \overline{D^{**^{0}}} \ell^{+} \nu \quad \text{background} \quad \text{angle}(\vec{p}_{B}^{*}, \vec{p}_{D^{*}\ell}^{*}) \\ (E_{B}^{*} - E_{D^{*}l}^{*})^{2} - |p_{B}^{*}|^{2} - |p_{D^{*}l}^{*}|^{2} + 2|p_{B}^{*}||p_{D^{*}l}^{*}|\cos(\theta_{B,D^{*}l}) = M_{\nu}^{2} \approx 0$

fit $\cos(\theta_{B,D^*\ell})$ distribution using MC shapes for $D^*\ell\nu$ and $D^{**\ell\nu}$



Background: $B^{\pm} \rightarrow D^{**} lv$



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Wrong Flavor

- Use MC to estimate the wrong flavor
- High purity events: $\omega = 0.015 \pm 0.006$
- Expect attenuation on the asymmetry:
- $A(\Delta t) = (1-2\omega) \cos(\Delta m \Delta t)$ $= 0.970 \cos(\Delta m \Delta t)$

 $\Rightarrow \sim 3\%$ attenuation



Toy MC study of deconvolution

- \bullet Toy MC with parametrized resolution in Δz
- Simulate 400 "runs", each consists of
 - ~35000 "MC" events based on QM
 - ~7000 "Data" events based on QM, LR or SD
- Produce 2 unfolding matrices for SF and OF events from "MC"
- Deconvolution performed on "Data" separately for SF and OF.
- Correct for residual systematic effects.



Cross check: Forward Test

At this stage, one can compare data with MC prediction for QM, LR and SD results.

Since our MC is generated with QM correlation, we re-weight each event to produce the prediction of PS and SD models.



 Δm is fixed. The result favors QM

Cross check: extra Δz resolution cut

Select events with better Δz resolution by adding a cut $\sigma(V_z) < 100 \mu m$ cut on both B decay vertices. This discards ~18% of the events.



=> results are consistent

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Sensitivity to λ

Toy MC study of sensitivity to decoherence. Fit of $(1 - \lambda)A_{QM} + \lambda A_{SD}$

 λ from fit vs λ generated



Sensitivity to Δm

