



Introduction To Neutrinos

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The Neutrino Revolution

(1998 – ...)

Neutrinos have nonzero masses!

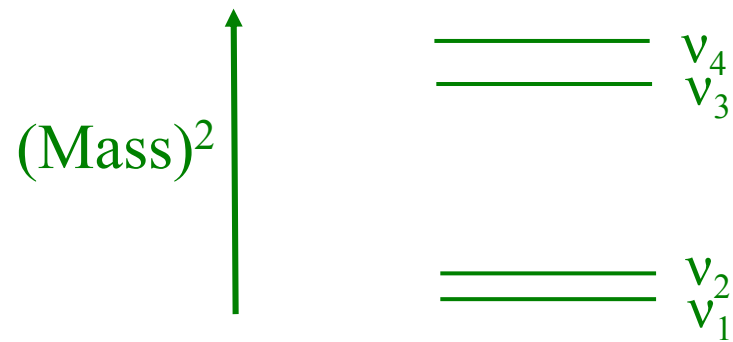
Leptons mix!

These discoveries come from
the observation of
neutrino flavor change
(neutrino oscillation).

The Physics of Neutrino Oscillation

Suppose There Are *Neutrino Masses*

Then there is some spectrum of neutrino mass eigenstates ν_i :



$$\text{Mass}(\nu_i) \equiv m_i$$

Suppose There Is *Leptonic Mixing*

Then the neutrinos $\nu_{e,\mu,\tau}$ of definite flavor

($W \rightarrow e\nu_e$ or $\mu\nu_\mu$ or $\tau\nu_\tau$)

are **superpositions** of the neutrinos of definite mass:

$$|\nu_\alpha\rangle = \sum_i U_{\alpha i}^* |\nu_i\rangle .$$

Neutrino of flavor $\alpha = e, \mu, \text{ or } \tau$

Unitary Leptonic Mixing Matrix

Neutrino of definite mass m_i

Neutrino Flavor Change (“Oscillation”)

Suppose a neutrino is born with flavor α and energy E .

$$\nu_\alpha = \sum_i U_{\alpha i}^* \nu_i \xrightarrow{\text{Distance } L} \sum_i U_{\alpha i}^* \nu_i e^{ip_i(E)L/\hbar} \equiv \nu(L) \neq \nu_\alpha$$

\uparrow
 $\sqrt{E^2 - (m_i c^2)^2} / c$

The neutrino has evolved into a mixture of the flavors.

Often, only 2 neutrinos need to be considered.

Then –

$$U = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix}$$

↙ Mixing Angle

Then–

$$P(\nu_\alpha \rightarrow \nu_{\beta \neq \alpha}) = \sin^2 2\theta \left[\sin^2 1.27 \frac{L(\text{km})}{E(\text{GeV})} \frac{\Delta m^2 (\text{eV}^2)}{\Delta m^2 (\text{eV}^2)} \right]$$

↑ Probability

$$\Delta m^2 \equiv m_2^2 - m_1^2$$

Neutrino oscillation



Neutrino mass & leptonic mixing

Evidence For Flavor Change

Neutrinos

Evidence of Flavor Change

Solar

Compelling

Reactor

Compelling

($L \sim 180$ km)

Atmospheric

Compelling

Accelerator

Compelling

($L = 250$ and 735 km)

Stopped μ^+ Decay

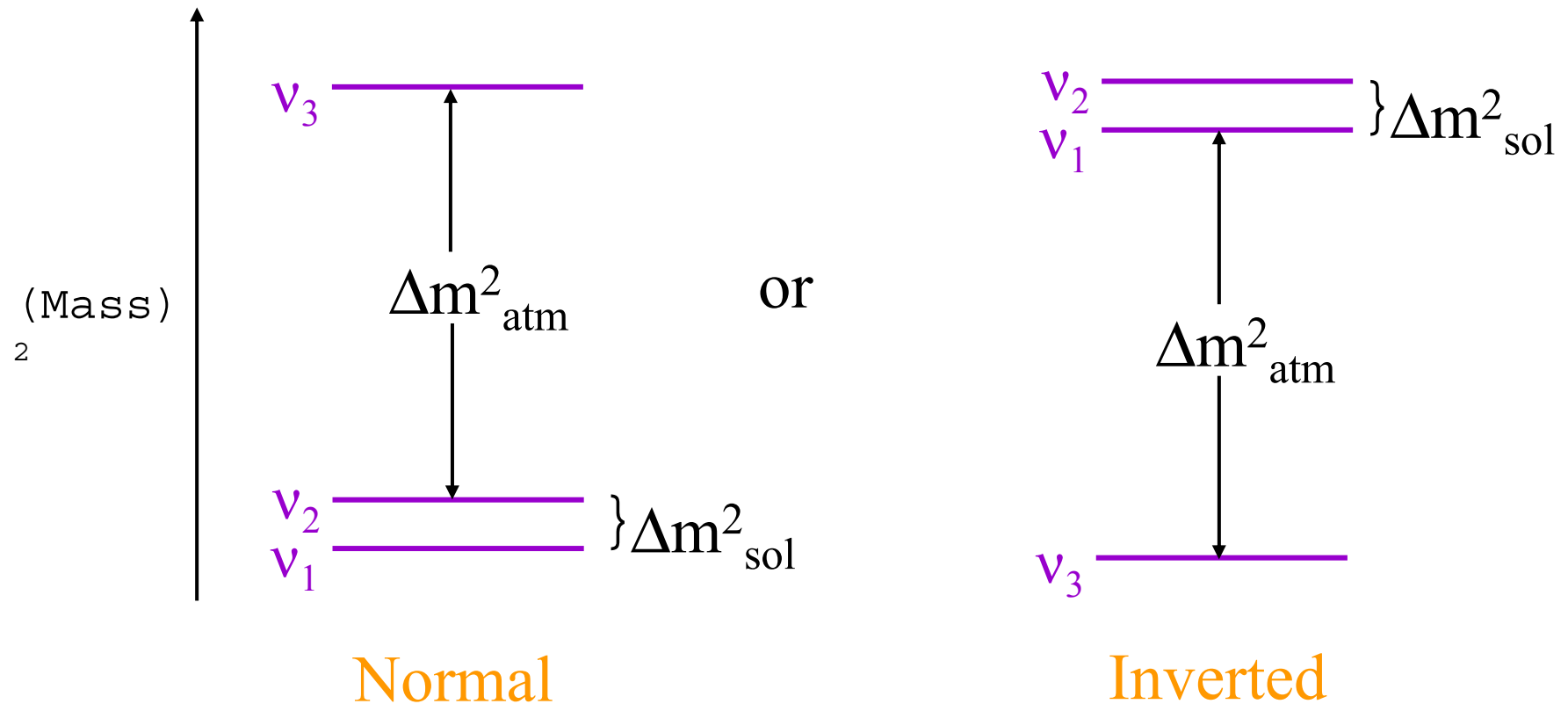
Unconfirmed

(LSND)
($L \approx 30$ m)



**What We
Have Learned**

The (Mass)² Spectrum



$$\Delta m^2_{\text{sol}} \cong 8 \times 10^{-5} \text{ eV}^2, \quad \Delta m^2_{\text{atm}} \cong 2.7 \times 10^{-3} \text{ eV}^2$$

Are there *more* mass eigenstates, as LSND suggests?

The Mixing Matrix

$$U = \begin{matrix} \text{Atmospheric} \\ \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \end{matrix} \times \begin{matrix} \text{Cross-Mixing} \\ \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \end{matrix} \times \begin{matrix} \text{Solar} \\ \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \end{matrix} \\
 \\
 \begin{matrix} c_{ij} \equiv \cos \theta_{ij} \\ s_{ij} \equiv \sin \theta_{ij} \end{matrix} \times \begin{bmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$\theta_{12} \approx \theta_{\text{sol}} \approx 34^\circ$, $\theta_{23} \approx \theta_{\text{atm}} \approx 37\text{-}53^\circ$, $\theta_{13} \lesssim 10^\circ$

Majorana ~~CP~~
phases

δ would lead to $P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta) \neq P(\nu_\alpha \rightarrow \nu_\beta)$. ~~CP~~

But note the crucial role of $s_{13} \equiv \sin \theta_{13}$.

The Open Questions

QuickTime™ and
TIFF (Uncompressed) decompressor
are needed to see this picture.

- What is the absolute scale of neutrino mass?
- Are neutrinos their own antiparticles?
- Are there “sterile” neutrinos?

We must be alert to surprises!

- What is the pattern of mixing among the different types of neutrinos?

What is θ_{13} ?

- Is the spectrum like \equiv or \equiv ?

- Do neutrino – matter interactions violate CP?

Is $P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta) \neq P(\nu_\alpha \rightarrow \nu_\beta)$?

- What can neutrinos and the universe tell us about one another?
- Is CP violation involving neutrinos the key to understanding the matter – antimatter asymmetry of the universe?
- What physics is behind neutrino mass?

The Importance of the Questions, and How They May Be Answered

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are needed to see this picture.

Are Neutrinos Their Own Antiparticles?

For each *mass eigenstate* ν_i , does —

- $\bar{\nu}_i = \nu_i$ (Majorana neutrinos)

or

- $\bar{\nu}_i \neq \nu_i$ (Dirac neutrinos) ?

$e^+ \neq e^-$ since $\text{Charge}(e^+) = -\text{Charge}(e^-)$.

But neutrinos may not carry any conserved charge-like quantum number.

A conserved **Lepton Number** L defined by—

$$L(\nu) = L(l^-) \quad \bar{\nu} \quad -L(\bar{\nu}) = -L(l^+) = 1$$

may not exist. —

If it does not, then nothing

To look for L nonconservation, seek —

Neutrinoless Double Beta Decay



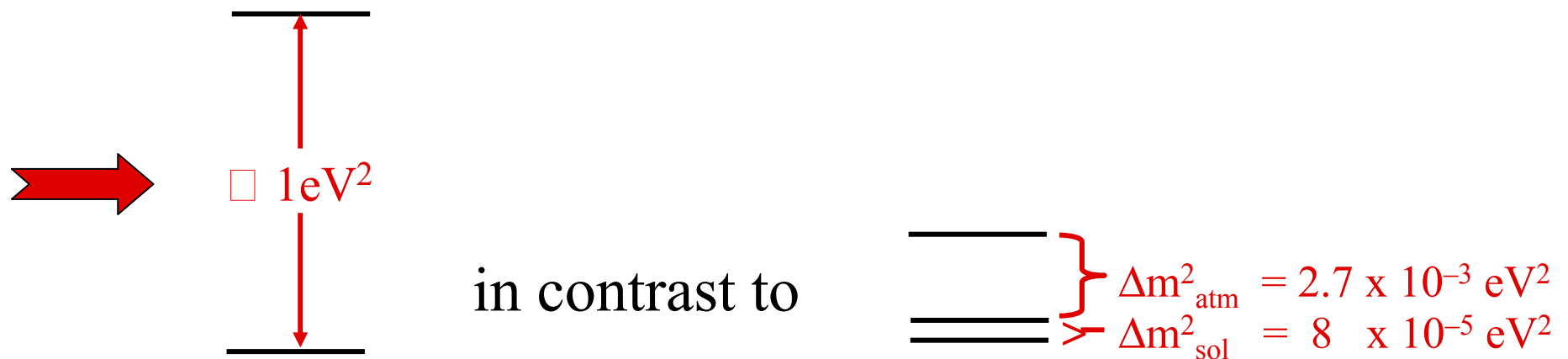
\mathcal{X} will be small, so need a lot (e.g., 1 ton)

Observation would establish that —

- L is not conserved
- $\bar{\nu}_i = \nu_i$ (Neutrinos are Majorana particles)
- Neutrinos have **Majorana masses**, unlike the charged leptons and quarks

Are There Sterile Neutrinos?

Rapid neutrino oscillation reported by the
L(iquid) S(cintillator) N(eutrino) D(etector) —



➡ At least **4** mass eigenstates, hence at least **4** flavors.

Measured $\Gamma(Z \rightarrow \nu\bar{\nu})$ ➡ only **3** different *active* neutrinos.

➡ At least **1** *sterile* neutrino.

Is the so-far unconfirmed oscillation
reported by LSND genuine?

MiniBooNE aims to definitively
answer this question.

The Central Role of θ_{13}

Both CP violation and our ability to tell whether the spectrum is normal or inverted depend on θ_{13} .

If $\sin^2\theta_{13} > (0.0025 - 0.0050)$, we can study both of these issues with intense but conventional accelerator ν and $\bar{\nu}$ beams, produced via $\pi^+ \rightarrow \mu^+ + \nu_{\mu}$ and $\pi^- \rightarrow \mu^- + \bar{\nu}_{\mu}$.

Determining θ_{13} is an important step.

Reactor Experiments

Looking for disappearance of reactor $\bar{\nu}_e$ while they travel $L \sim 1.5$ km with energy $E \sim 3$ MeV is the cleanest way to determine θ_{13} .

$P(\bar{\nu}_e \text{ Disappearance}) =$

$$= \sin^2 2\theta_{13} \sin^2 [1.27 \Delta m_{\text{atm}}^2 (\text{eV}^2) L (\text{km}) / E (\text{GeV})]$$

Accelerator Experiments

Accelerator neutrino experiments can also probe θ_{13} , although now it is entwined with other parameters.

Their *special* capability is to probe *whether the mass spectrum is normal or inverted*, and to look for *CP violation*.

Both are done by studying $\nu_{\mu} \rightarrow \nu_e$ and $\bar{\nu}_{\mu} \rightarrow \bar{\nu}_e$ while the beams travel hundreds of kilometers.

The Mass Spectrum: \equiv or \equiv ?

Generically, grand unified models (GUTS) favor —

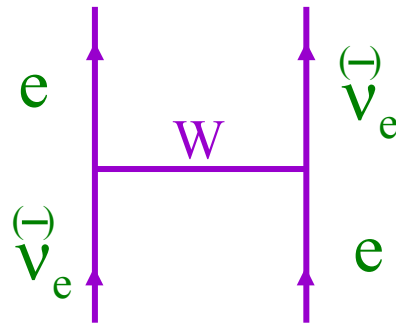
\equiv

GUTS relate the **Leptons** to the **Quarks**.

\equiv is un-quark-like, and would probably involve a lepton symmetry with no quark analogue.

How To Determine If The Spectrum Is Normal Or Inverted

Exploit the fact that, in matter,

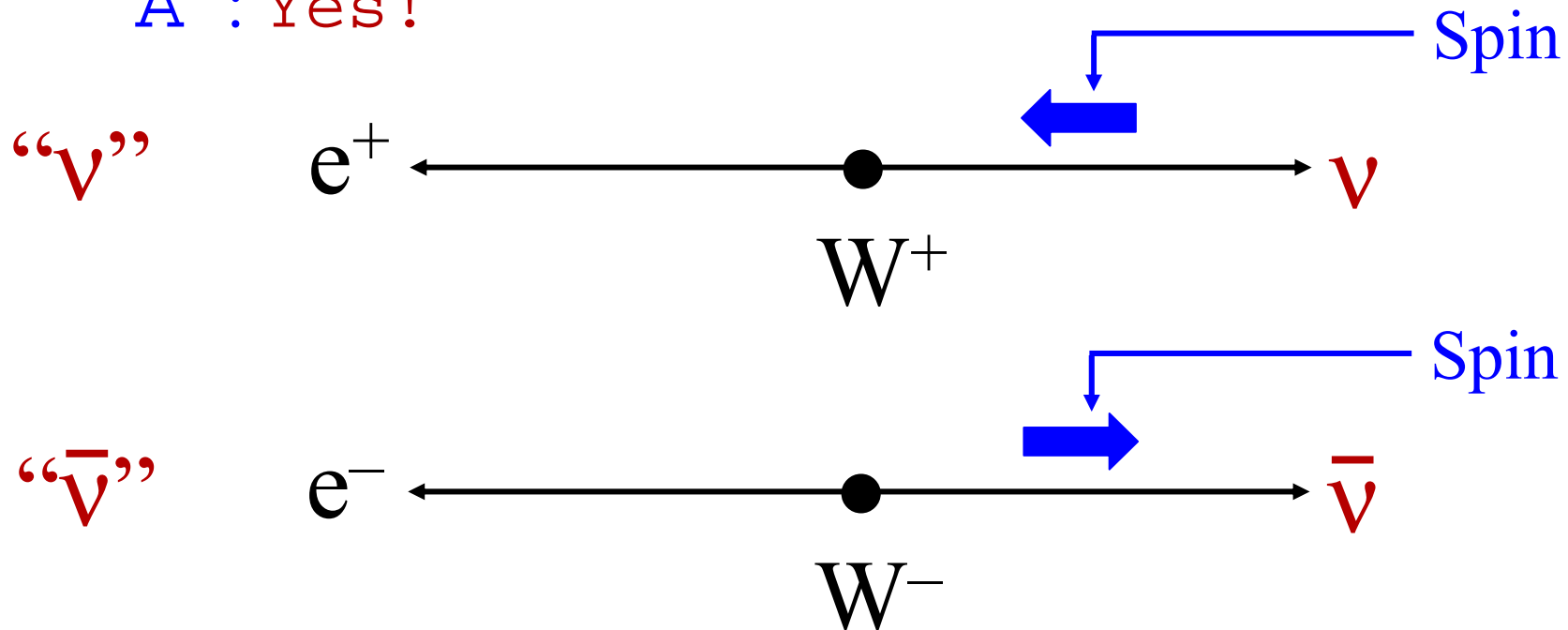


affects ν and $\bar{\nu}$ differently, and leads to —

$$\frac{P(\nu_{\mu} \rightarrow \nu_e)}{P(\bar{\nu}_{\mu} \rightarrow \bar{\nu}_e)} \begin{cases} > 1 ; \text{---} \\ < 1 ; \text{---} \end{cases} \quad \text{Note fake } CP$$

Q : Does matter still affect ν and $\bar{\nu}$ differently when $\bar{\nu} = \nu$?

A : Yes!



The weak interactions violate *parity*. Neutrino – matter interactions depend on the neutrino *polarization*.

Do Neutrino Interactions Violate CP?

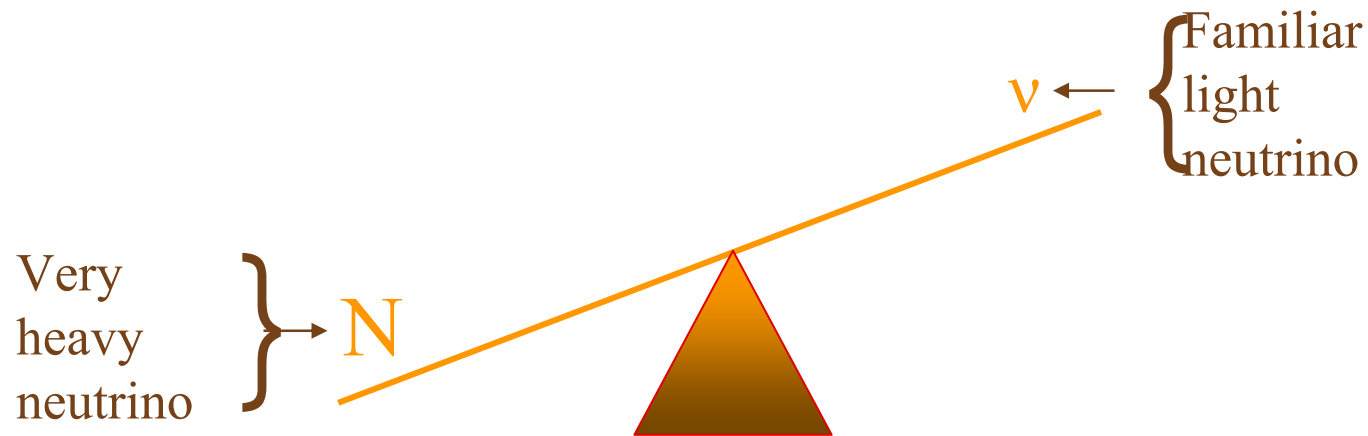
Is leptonic ~~CP~~, through **Leptogenesis**,
the origin of the **Baryon Asymmetry** of
the universe?

Leptogenesis In 60

Seconds

The most popular theory of why neutrinos are so light is the —

See-Saw Mechanism



The heavy neutrinos **N** would have been made in the hot Big Bang.

The heavy neutrinos N , like the light ones ν , are Majorana particles. Thus, an N can decay into l^- or l^+ .

If neutrino oscillation violates CP, then quite likely so does N decay. In the See-Saw, these two CP violations have a common origin.

Then, in the early universe, we would have had different rates for the CP-mirror-image decays –

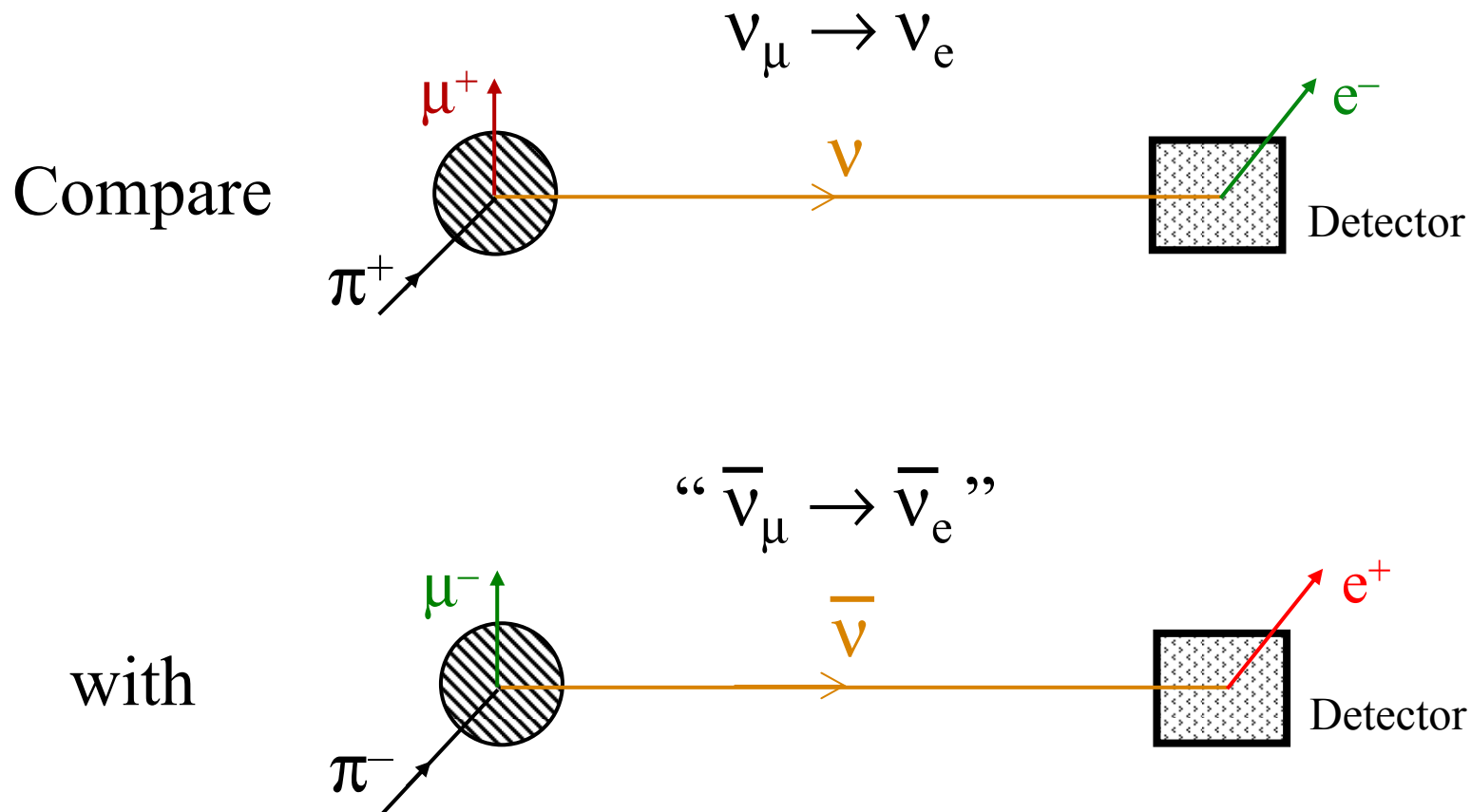
$$N \rightarrow l^- + \dots \quad \text{and} \quad N \rightarrow l^+ + \dots$$

This would have led to unequal numbers of leptons and antileptons (Leptogenesis).

Then, Standard-Model *Sphaleron* processes would have turned $\sim 1/3$ of this leptonic asymmetry into a Baryon Asymmetry.

Q : Can CP violation still lead to
 $P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) \neq P(\nu_\mu \rightarrow \nu_e)$ when $\nu = \bar{\nu}$?

A : Certainly!



Separating \cancel{CP} From the Matter Effect

Genuine \cancel{CP} and the matter effect
both lead to a difference between
“ ν ” and “ $\bar{\nu}$ ” oscillation.

But genuine \cancel{CP} and the matter effect depend
quite differently from each other on L and E.

To disentangle them, one must make oscillation
measurements at different L and/or E.

Summary

We have learned a lot about the neutrinos in the last decade.

What we have learned raises some very interesting questions.

Strategies and experiments for answering them are being developed.