

Status of V_{us}

- 1) V_{us} from K_{l3} decays
- 2) V_{us} from F_K / F_π ratio
- 3) V_{us} from hadronic τ decays
- 4) V_{us} from hyperon decays

Theoretically, the K_{l3} -decay rate is given by:

$$\Gamma[K \rightarrow \pi l \nu_e(\gamma)] = \frac{G_F^2 M_K^5}{192\pi^3} C_K^2 S_{EW} I_K(\{\lambda\}) |V_{us}|^2 f_+^2(0)$$

where with $t = (p-p')^2$:

$$\langle \pi^-(p') | \bar{s} \gamma_\mu u | K^0(p) \rangle = (p+p')_\mu f_+(t) + (p-p')_\mu f_-(t)$$

The largest theoretical uncertainty resides in the hadronic form factor $f_+(0)$. \Rightarrow First determine: $|V_{us}| f_+(0)$

Recent new measurements by E865, KTeV, NA48 and KLOE:

$$\begin{aligned} |V_{us}| f_+(0) &= 0.2169(9) && \text{(PDG 06)} \\ |V_{us}| f_+(0) &= 0.21673(46) && \text{(FLAVIANet } K \text{ WG 06)} \end{aligned}$$

Current data on K_{l3} form-factor slopes

| | Note | $\lambda'_+ \times 10^3$ | $\lambda''_+ \times 10^3$ | $\lambda_0 \times 10^3$ |
|---------------------------------|------------------------|----------------------------------|--------------------------------|-------------------------|
| KTev PRD 70 (2004) | $K_L e3$ $K_L \mu3$ | 21.7 ± 2.0 17.0 ± 3.7 | 2.9 ± 0.8 4.4 ± 1.5 | 12.8 ± 1.8 |
| KLOE PLB 636 (2006) | $K_L e3$ | 25.5 ± 1.8 | 1.4 ± 0.8 | |
| NA48 PLB 604 (2004) | $K_L e3$ | 28.0 ± 2.4 | 0.4 ± 0.9 | |
| NA48 hep-ex/0703002 | $K_L \mu3$ | 20.5 ± 3.3 | 2.6 ± 1.4 | 9.5 ± 1.4 |
| ISTRA+ PLB 581 (2004) | $K^- e3$ | 24.9 ± 1.7 | 1.9 ± 0.9 | |
| ISTRA+ PLB 589 (2004) | $K^- \mu3$ | 23.0 ± 6.4 | 2.3 ± 2.3 | 17.1 ± 2.2 |

Kindly provided by **Matthew Moulson**; Update from **CKM06**.

$$f_+(t) = f_+(0) \left[1 + \lambda'_+ \frac{t}{M_\pi^2} + \lambda''_+ \frac{t^2}{2M_\pi^4} + \dots \right]$$

Fit to K_{l3} form-factor slopes

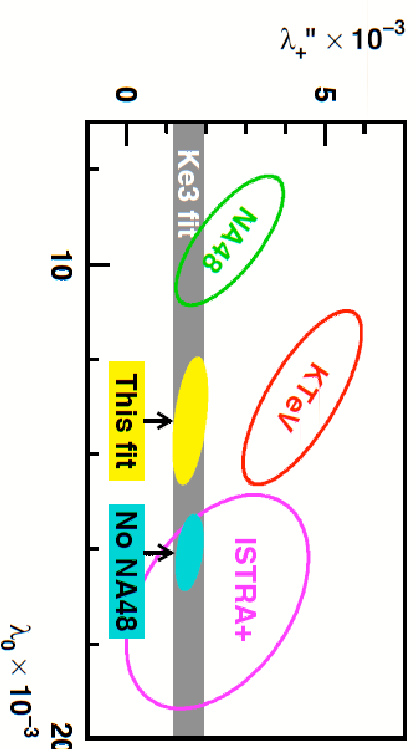
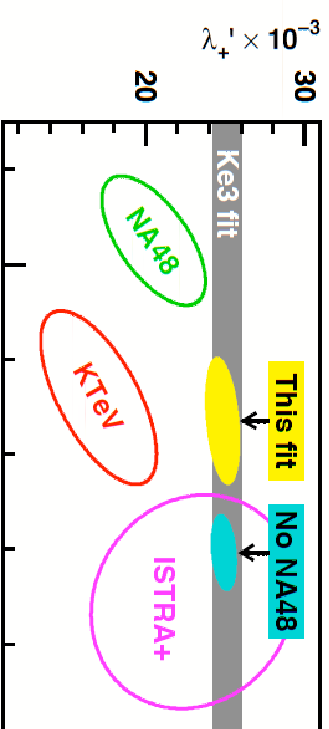
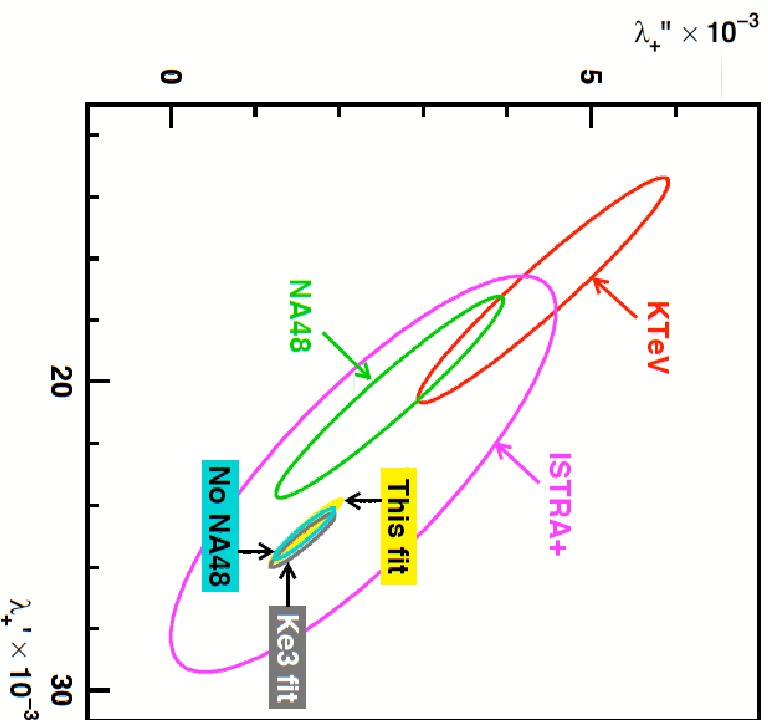
Slopes from

KTev K_{l3}

ISTRA+ K_{l3}

NA48 K_{l3}

Fit to Ke3 data



K_{l3} fit, no NA48 K_{l3} : $\chi^2=11.9/9$ (21.7%)

K_{l3} fit, all data, $\chi^2=58/12$ (10^{-6})

Kindly provided by **Matthew Moulson**; Update from **CKM06**.

K_{l3} form-factor slopes: Fit results

Although compatibility poor, no *a priori* reason to exclude NA48 $K\mu 3$ data
Inconsistency parameterized by scale factors for fit results

Slope parameters $\times 10^3$:

$$\begin{array}{lll} \lambda'_{+} = 24.84 \pm 1.10 & S = 1.4 \\ \lambda''_{+} = 1.61 \pm 0.45 & S = 1.3 \\ \lambda_0 = 13.30 \pm 1.35 & S = 2.1 \\ \chi^2/\text{ndf} = 52/12 \quad (10^{-6}) \end{array}$$

Correlation coefficients:

$$\begin{array}{ll} \lambda'_{+} & \lambda_0 \\ \lambda'_{+} & -0.94 \quad +0.31 \\ \lambda''_{+} & -0.42 \end{array}$$

Kindly provided by **Matthew Moulson**; Update from **CKM06**.

Integrals

$$\begin{array}{ll} I(K^0 e3) & 0.15452(29) \\ I(K^+ e3) & 0.15887(30) \\ I(K^0 \mu 3) & 0.10207(34) \\ I(K^+ \mu 3) & 0.10501(35) \end{array}$$

These results used to evaluate
 $|V_{us}| f_{\pm}(0)$ for all modes

A description of the $K\pi$ vector form factor can be obtained within **chiral** perturbation theory with resonances (**R χ PT**):

(MJ, Pich, Portolés 2006)

$$F_{+}^{K\pi}(s) = \frac{M_{K^*}^2 e_2^3 \text{Re} [\tilde{H}_{K\pi}(s) + \tilde{H}_{K\eta}(s)]}{M_{K^*}^2 - s - iM_{K^*} \Gamma_{K^*}(s)} .$$

where

$$\Gamma_{K^*}(s) = \frac{G_V^2 M_{K^*} s}{64\pi F_\pi^4} [\sigma_{K\pi}^3(s) + \sigma_{K\eta}^3(s)] .$$

The parameters of this model, namely M_{K^*} and G_V , can be fitted from **experimental** data for p -wave $K\pi$ scattering.

Also a **second** resonance contribution can easily be included.

As a prediction of the model, we obtain the slope and the curvature of the vector form factor $F_+^{K\pi}(s)$:

$$\lambda'_+ = 25.6 \cdot 10^{-3}, \quad \lambda''_+ = 1.31 \cdot 10^{-3},$$

to be compared with the experimental result:

$$\lambda'_+ = (24.8 \pm 1.1) \cdot 10^{-3}, \quad \lambda''_+ = (1.61 \pm 0.45) \cdot 10^{-3}.$$

Likewise, the scalar form factor $F_0^{K\pi}(s)$ can be obtained from a dispersion relation analysis of S-wave $K\pi$ scattering data.

(MJ, Oller, Pich 2002/04/06)

Here, the respective result is:

$$\lambda_0 = (14.7 \pm 0.4) \cdot 10^{-3} \Leftrightarrow \lambda_0^{\text{exp}} = (13.3 \pm 1.4) \cdot 10^{-3}.$$

| | Reference | $f_+^{K^0\pi^-}(\mathbf{0})$ |
|---------------------|-----------------------------|------------------------------|
| Quark model | [Leutwyler & Roos, 1984] | 0.961 (8) |
| | [Becirevic et al, 2005] | 0.960 (9) |
| | [MILC Collab., 2005] | 0.962 (11) |
| Lattice | [Dawson et al, 2006] | 0.968 (11) |
| | [UKQCD/RBC Collab., 2006] | 0.961 (5) ! |
| | [Bijnens & Talavera, 2003] | 0.976 (10) |
| K π scalar f.f. | [Jamin, Oller & Pich, 2004] | 0.974 (11) |
| Large- N_c | [Cirigliano et al, 2005] | 0.984 (12) |

Kindly provided by Jorge Portolés.

IFIC - Instituto de Física Corpuscular

$$\Rightarrow f_+(0) = 0.966(5) \Rightarrow$$

$$|V_{us}| = 0.2244(13)$$

Status of V_{us}

M. Jamin, ICREA & IFAE, UA Barcelona

Moriond EW, March 2007

From the leptonic decays $\Gamma[K \rightarrow l\nu_l(\gamma)] / \Gamma[\pi \rightarrow l\nu_l(\gamma)]$,
one can predict: (Marciano 2004)

$$\frac{|V_{us}|}{|V_{ud}|} \frac{F_K}{F_\pi} = 0.27618 \pm 0.00048.$$

Together with the recent lattice result (MILC 2006)

$$\frac{F_K}{F_\pi} = 1.208(2) \left(\begin{smallmatrix} +7 \\ -14 \end{smallmatrix} \right) \quad \text{and} \quad V_{ud} = 0.97377(27) \\ \text{(nuclear } \beta\text{-decay)}$$

this leads to:

$$|V_{us}| = 0.2226 \left(\begin{smallmatrix} +26 \\ -14 \end{smallmatrix} \right)$$

Unitarity and V_{ud} imply: $F_K / F_\pi = 1.182(6)$! LR84: 1.22(1)

Consider the physical quantity R_τ : (Braaten, Narison, Pich (1992))

$$R_\tau \equiv \frac{\Gamma(\tau^- \rightarrow \text{hadrons } \nu_\tau(\gamma))}{\Gamma(\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau(\gamma))} = 3.640 \pm 0.010.$$

R_τ is related to the QCD correlators $\Pi^{T,L}(z)$: ($z \equiv s/M_\tau^2$)

$$R_\tau = 12\pi \int_0^1 dz (1-z)^2 \left[(1+2z) \text{Im} \Pi^T(z) + \text{Im} \Pi^L(z) \right],$$

with the appropriate combinations

$$\Pi^J(z) = |V_{ud}|^2 \left[\Pi_{ud}^{V,J} + \Pi_{ud}^{A,J} \right] + |V_{us}|^2 \left[\Pi_{us}^{V,J} + \Pi_{us}^{A,J} \right].$$

Additional information can be inferred from the moments

$$R_{\tau}^{kl} \equiv \int_0^1 dz (1-z)^k z^l \frac{dR_{\tau}}{dz} = R_{\tau,NS}^{kl} + R_{\tau,S}^{kl}.$$

Theoretically, R_{τ}^{kl} can be expressed as:

$$R_{\tau}^{kl} = N_c S_{EW} \left\{ (|V_{ud}|^2 + |V_{us}|^2) [1 + \delta^{kl(0)}] + \sum_{D \geq 2} \left[|V_{ud}|^2 \delta_{ud}^{kl(D)} + |V_{us}|^2 \delta_{us}^{kl(D)} \right] \right\}.$$

$\delta_{ud}^{kl(D)}$ and $\delta_{us}^{kl(D)}$ are corrections in the Operator Product Expansion, the most important ones being $\sim m_s^2$ and $m_s \langle \bar{q}q \rangle$.

The sensitivity to the strange quark mass can be enhanced by considering the flavour SU(3)-breaking difference:

(Pich, Prades; ALEPH (1998))

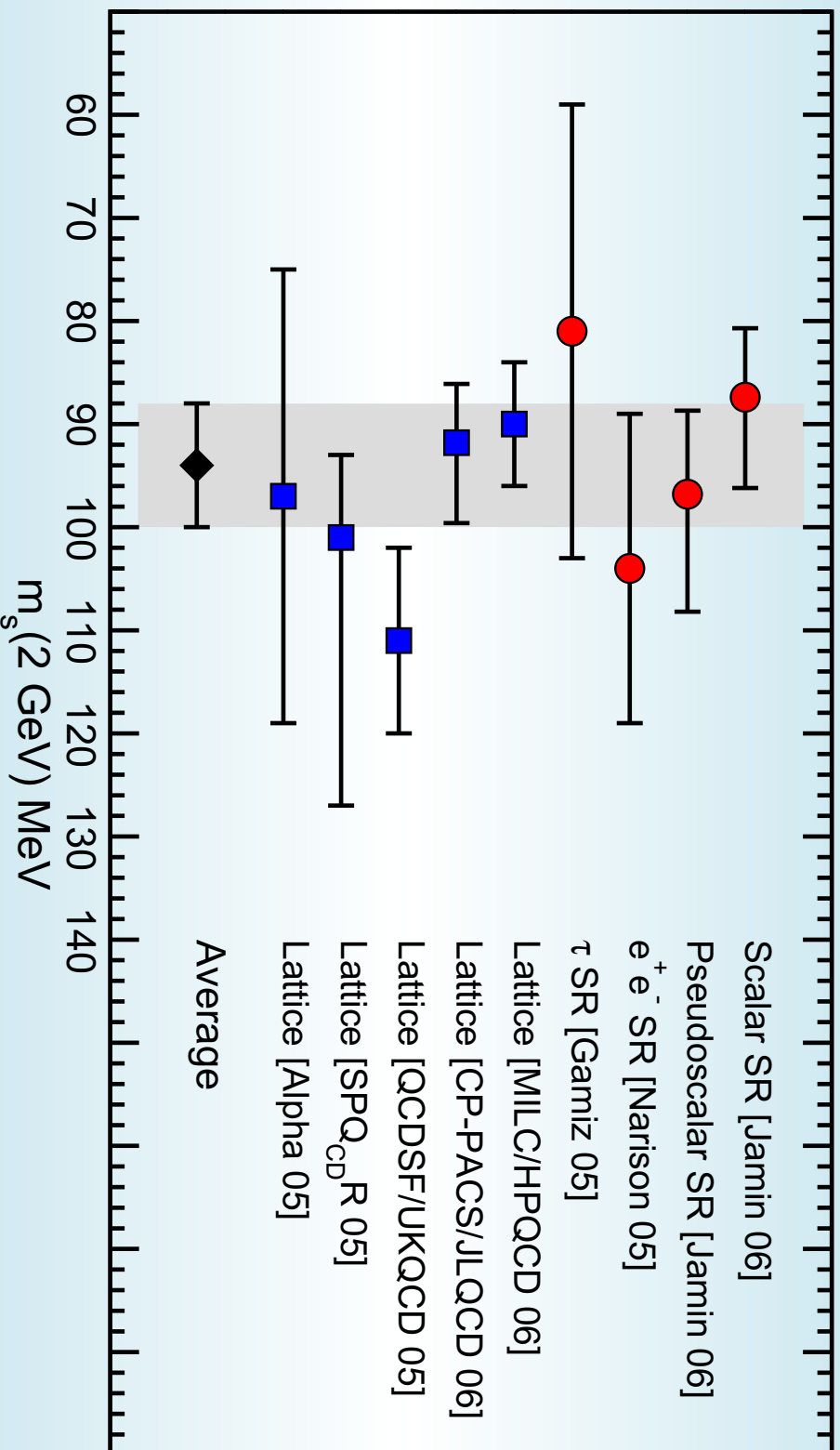
$$\delta R_{\tau}^{kl} \equiv \frac{R_{\tau,NS}^{kl}}{|V_{ud}|^2} - \frac{R_{\tau,S}^{kl}}{|V_{us}|^2} = 3 S_{EW} \sum_{D \geq 2} \left(\delta_{ud}^{kl(D)} - \delta_{us}^{kl(D)} \right).$$

Flavour independent uncertainties drop out in the difference.

In previous analyses a sizeable part of the theoretical error was due to large α_s corrections in the longitudinal contribution.

This uncertainty could be greatly reduced by replacing badly behaved scalar/pseudoscalar correlators with phenomenology.

(Gámiz, M.J., Pich, Prades, Schwab (2003/04))



\Rightarrow Average: $m_s(2 \text{ GeV}) = 94 \pm 6 \text{ MeV}$

Given m_s , we are in a position to predict δR_{τ}^{kl} from theory. Theoretically, the uncertainty is smallest for the $(0,0)$ moment:

$$\delta R_{\tau,th} = \underbrace{0.155}_{\text{Pheno}} + \underbrace{0.078}_{m_s^2} + 0.003 = 0.236 \pm 0.028.$$

Let us now reconsider the equation for δR_{τ} :

$$|V_{us}| = \sqrt{\frac{R_{\tau,S}}{R_{\tau,NS} / |V_{ud}|^2 - \delta R_{\tau,th}}} \approx 3.658$$

Thus the theoretically derived quantity $\delta R_{\tau,th}$ only gives a small correction to experimentally measured quantities.

Together with the experimental results $R_{\tau,NS} = 3.469 \pm 0.014$ as well as $R_{\tau,S} = 0.1677 \pm 0.0050$, V_{us} can be determined:

$$|V_{us}| = 0.2214 \pm 0.0033_{\text{exp}} \pm 0.0010_{\text{th}} = 0.2214 \pm 0.0034$$

The uncertainty on V_{us} is dominated by the experimental error on $R_{\tau,S}$. The theoretical error by the perturbative expansion.

In the near future, it should be possible to reduce the uncertainty with the τ -data sets from BABAR and BELLE.

If the experimental value $B(\tau \rightarrow K\nu_\tau) = (0.686 \pm 0.023)\%$ is replaced by the theoretical prediction $(0.715 \pm 0.004)\%$ based on $K_{\mu 2}$ decays, one finds $|V_{us}| = 0.2225 \pm 0.0034$.

$$|V_{us}|(K_{l3}) = 0.2244(13)$$

$$|V_{us}|(\frac{F_K}{F_\pi}) = 0.2226(^{+26}_{-14})$$

$$|V_{us}|(\tau) = 0.2225(34)$$

$$|V_{us}|(\text{Hyp}) = 0.226(5)$$

\Rightarrow **Average:**

$$|V_{us}| = 0.2240(11)$$

Unitarity relation: $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1 - \delta$.

$$\Rightarrow \delta = (1.58 \pm 0.72) \cdot 10^{-3}$$

$\approx 2.2 \sigma$
difference.

Or vice versa:

$$|V_{us}|_{\text{Unit}} = 0.2275(12)$$