

# Sterile Neutrino

-- neutrino masses, dark matter,  
and baryon asymmetry --

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# Outline

- The  $\nu$ MSM

= the "*neutrino*" Minimal Standard Model

TA, Blanchet, Shaposhnikov

TA, Shaposhnikov

- Dark matter

- Baryogenesis

- Summary

Cf. The "New" MSM

Davoudiasl, Kitano, Li, Murayama

PLB609, 117, 2005

# Motivation

The MSM fails to explain:

- **Neutrino oscillations**

- Neutrinos are exactly massless

- **Dark matter**

- No particle candidate for dark matter

- **Baryon asymmetry**

- No mechanism to generate baryon asymmetry
- No (strong) deviation from thermal equilibrium
- CP violation is too small

→ go beyond the MSM !!

- A simple possibility is the  $\nu$ MSM

What is the vMSM?

# Framework ( $\nu$ MSM)

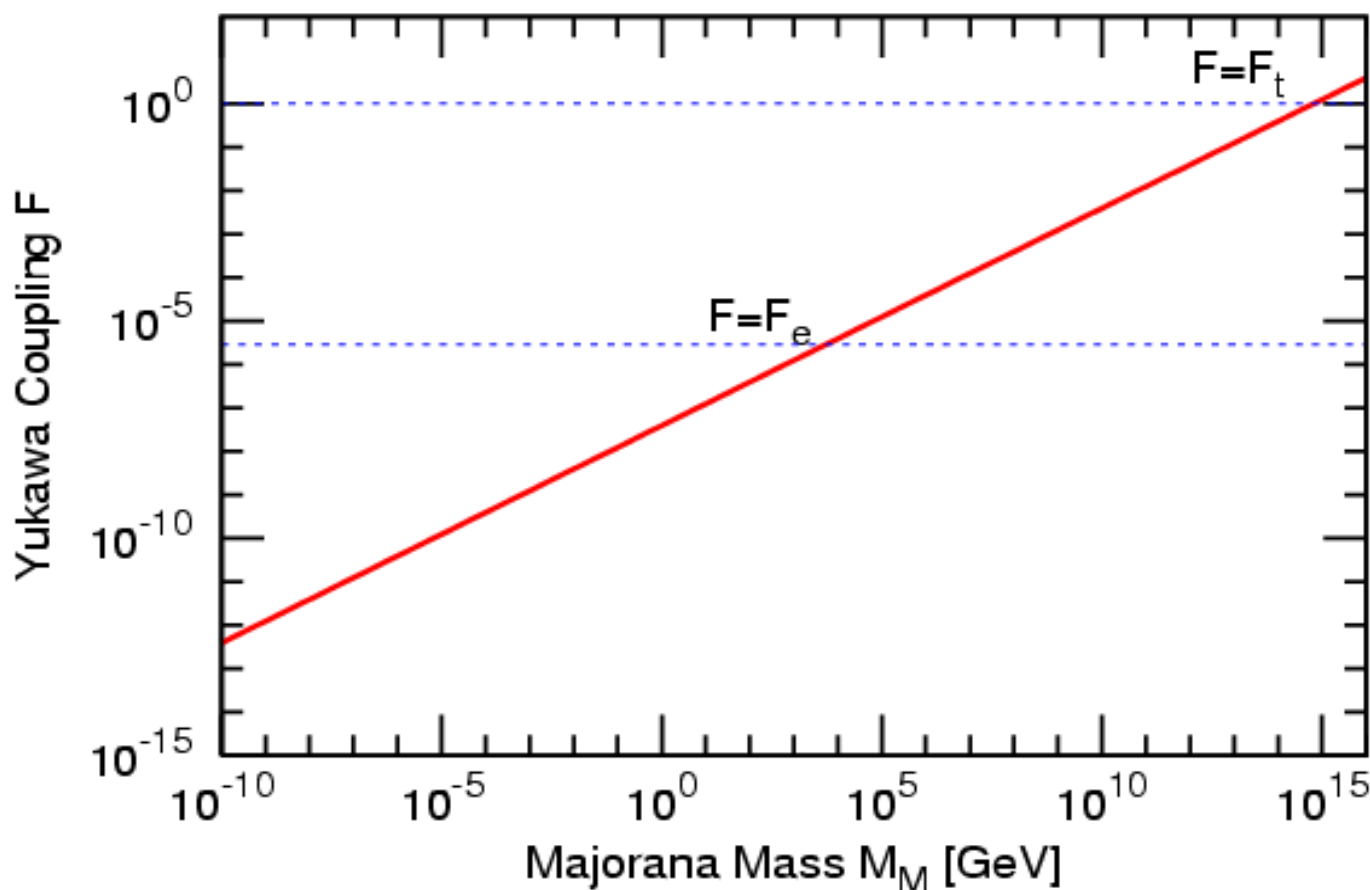
The MSM + three right-handed neutrinos  $N_1, N_2, N_3$

$$\mathcal{L}_{\nu\text{MSM}} = \mathcal{L}_{\text{MSM}} + i\bar{N}_I \not{\partial} N_I - F_{\alpha I} \bar{L}_\alpha \Phi N_I - \frac{M_I}{2} \bar{N}_I^c N_I + h.c.$$

- Dirac masses  $M_D = F\langle\Phi\rangle$  and Majorana masses  $M_I$
- Majorana masses are smaller than  $\sim 100$  GeV
  - All parameters can potentially be determined(?)
  - Seesaw mechanism still works, if  $[M_D]_{\alpha I} \ll M_I$
  - Neutrino oscillations by active neutrino mixing

# Scale of Majorana mass

$$M_\nu = -M_D^T \frac{1}{M_M} M_D \Rightarrow F^2 = M_M M_\nu / \langle \Phi \rangle^2$$



$$M_\nu = \sqrt{\Delta m_{atm}^2}$$

Very small  
Yukawa  
couplings!

# Dark matter in the $\nu$ MSM

# Dark Matter

**Candidate:** lightest sterile neutrino  $N_1$  with  $\sim \text{keV}$  mass

Dodelson, Widrow / Shi, Fuller / Dolgov, Hansen /  
Abazajian, Fuller, Patel

## Requirements:

- Its lifetime is longer than the age of the universe
  - $\tau > t_U \sim 10^{17}$  sec
- It explains dark matter density
  - $\Omega_{\text{dm}} h^2 \simeq 0.1$
- It avoids cosmological constraints
  - Bounds from X-ray and Ly-alpha forest observations

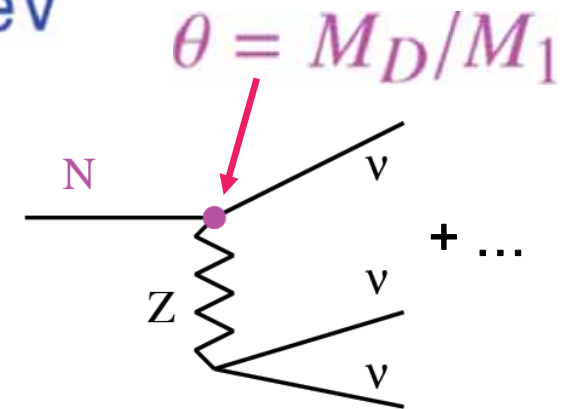


# Decays of sterile neutrino

**N1 is not completely stable particle !**

- Dominant decay:  $N_1 \rightarrow 3\nu$  for  $M_1 \sim \text{keV}$
- Lifetime can be very long

$$\tau_{N_1} \simeq 5 \cdot 10^{26} \text{sec} \left( \frac{\text{keV}}{M_1} \right)^5 \left( \frac{10^{-8}}{\theta^2} \right)$$

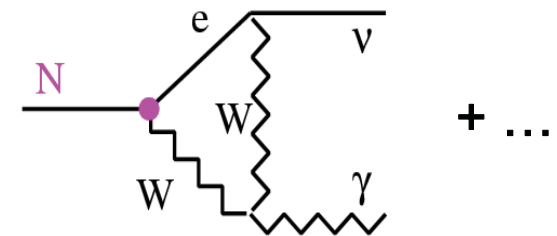


**N1 is not completely dark !**

- Subdominant decay:  $N_1 \rightarrow \nu + \gamma$
- Branching ratio is small

$$\text{Br} = 27\alpha_{\text{em}}/8\pi$$

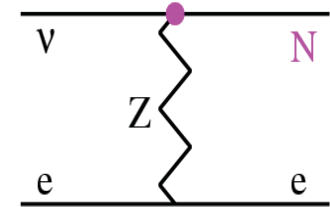
- But, severely restricted from X-ray observations



# Production of sterile neutrino

Dodelson-Widrow scenario:

- N1 is not thermalized
- Production via active-sterile neutrino oscillations
- Dominant production at  $T_* \simeq 100\text{MeV} (M_1/\text{keV})^{1/3}$



Recently, we improve the estimate of the abundance

- Use kinetic equation for density matrix
- Study “hadronic uncertainties” in detail:
  - Hadronic contributions to production rate
  - QCD equation of state (time-temp. relation)

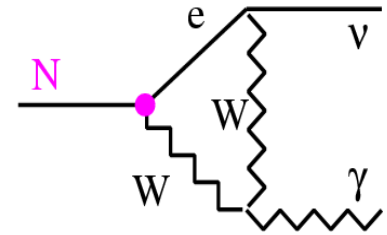
TA, Laine, Shaposhnikov JHEP06 ('06) 053 [hep-ph/0606209]

TA, Laine, Shaposhnikov [hep-ph/0612182]

# Constraints from X-rays

## Radiative decays of sterile neutrino DM

- feature in X-ray background spectrum
- line X-ray from clusters, galaxies, dwarf galaxies...
  - TEST for sterile neutrino DM!
  - No signal → **Upper bound on mixing angle !**



Dolgov, Hansen / Abazajian, Fuller, Tucker  
Boyarsky, Neronov, Ruchayskiy, Shaposhnikov  
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Tkachev  
Riemer-Sorensen, Hansen, Pedersen  
Watson, Beacom, Yuksel, Walker  
Boyarsky, Ruchayskiy, Markevitch  
Riemer-Sorensen, Pedersen, Hansen, Dahle  
Abazajian, Markevitch, Koushiappas, Hickox  
Boyarsky, Herder, Neronov, Ruchayskiy

# Constraints from structure formation

Light sterile neutrino = WDM  $\lambda_{\text{FS}} \sim \text{Mpc} \left( \frac{\text{keV}}{M_1} \right) \frac{\langle |\mathbf{q}_s| \rangle}{\langle |\mathbf{q}_a| \rangle}$

- Erase structures on smaller scales

- Lower bound on mass  $M_1 > \frac{\langle |\mathbf{q}_s| \rangle}{\langle |\mathbf{q}_a| \rangle} M_0$

From Ly- $\alpha$  forest observations

$M_0 \simeq 14.4 \text{keV}$  Seljak, Makarov, McDonald, Trac '06

$M_0 \simeq 10 \text{keV}$  Viel, Lesgourgues, Haehnelt, Matarresse, Riotto '06

We find:  $\frac{\langle |\mathbf{q}_s| \rangle}{\langle |\mathbf{q}_a| \rangle} \simeq 0.8$  for  $M_1 \simeq 10 \text{keV}$

  $M_1 \gtrsim 11.6 \text{keV}$  (SMMT),  $8 \text{keV}$  (VLHMR)

# Parameter space

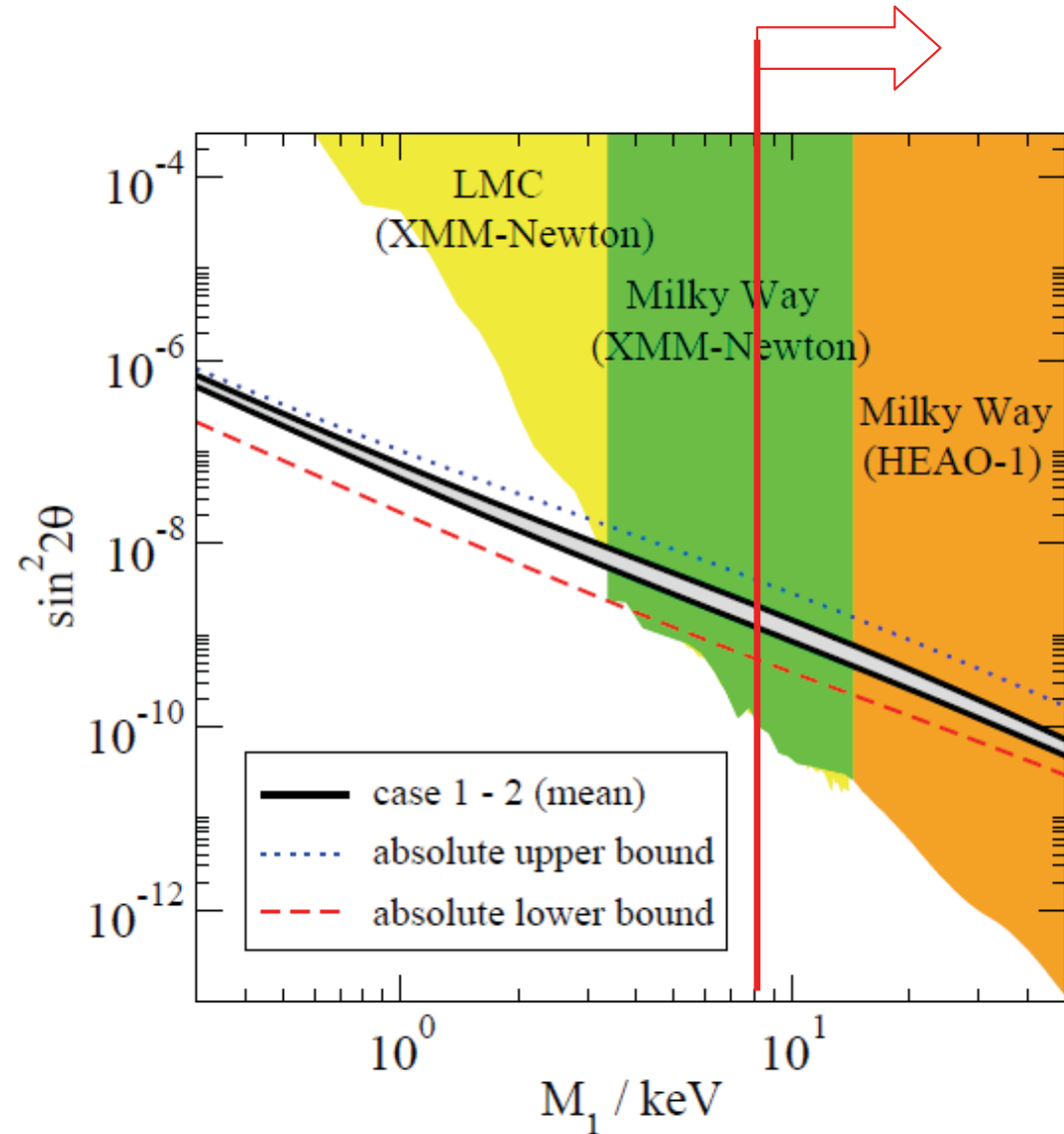
X-ray constraint:

$$M_1 \lesssim (3.5) 6 \text{ keV}$$

Ly-alpha constraint:

$$M_1 \gtrsim 11.6 \text{ keV (SMMT)}$$

$$M_1 \gtrsim 8 \text{ keV (VLHMR)}$$



X-ray constraints from

Boyarsky, Neronov, Ruchayskiy, Shaposhnikov, Tkachev (astro-ph/0603660)

Boyarsky, Nevalainen, Ruchayskiy (astro-ph/0610961)

# Fate of sterile neutrino DM

Dodelson-Widrow scenario assumes:

- No initial abundance at  $T \sim 1\text{GeV}$
- No new interaction at  $E < 1\text{GeV}$
- Charge asymmetries smaller than baryon asymmetry
- No low reheating, i.e. RD universe starts  $T > 1\text{GeV}$

**Sterile neutrino DM is still possible, when**

- Large lepton asymmetry (Shi-Fuller)
  - resonant production via active-sterile oscillations
- Production via decays of Inflaton/Scalar  
(Shaposhnikov, Tkachev / Kusenko)
- Dimension-five operator (Anisimov's talk)
- . . .

# Baryogenesis in the $\nu$ MSM

# Baryon asymmetry

Baryogenesis conditions: (Sakharov)

- Baryon (Lepton) number violation

- B+L violation by sphaleron ( $T > \sim 100 \text{ GeV}$ )
- L violation due to Majorana masses

- C and CP violations

- 1 CP phase in quark sector
- 6 CP phases in neutrino sector

- Out of equilibrium

- Sterile neutrino can be out-of-equilibrium

Scenarios:

- Leptogenesis (Fukugita, Yanagida) ← NO

- Baryogenesis via sterile neutrino oscillations ← YES!



# Baryogenesis via neutrino oscillations

Akhmedov, Rubakov, Smirnov '98

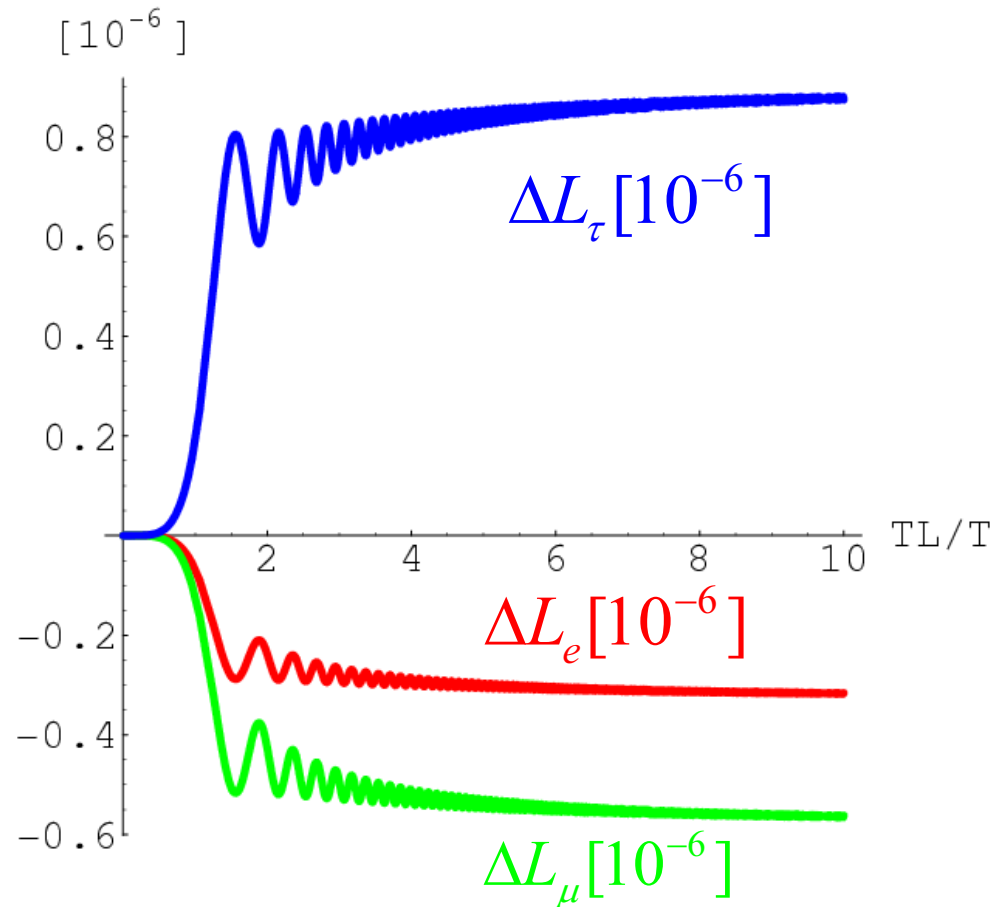
Idea:

- Heavier sterile neutrinos are created and oscillate together with CP violations
- No generation of total lepton asymmetry, since Majorana masses are small ( $M_I \ll T$ )
- Lepton asymmetry is *distributed* between active and sterile sectors
- Asymmetry of active sector is partially transferred into baryon asymmetry

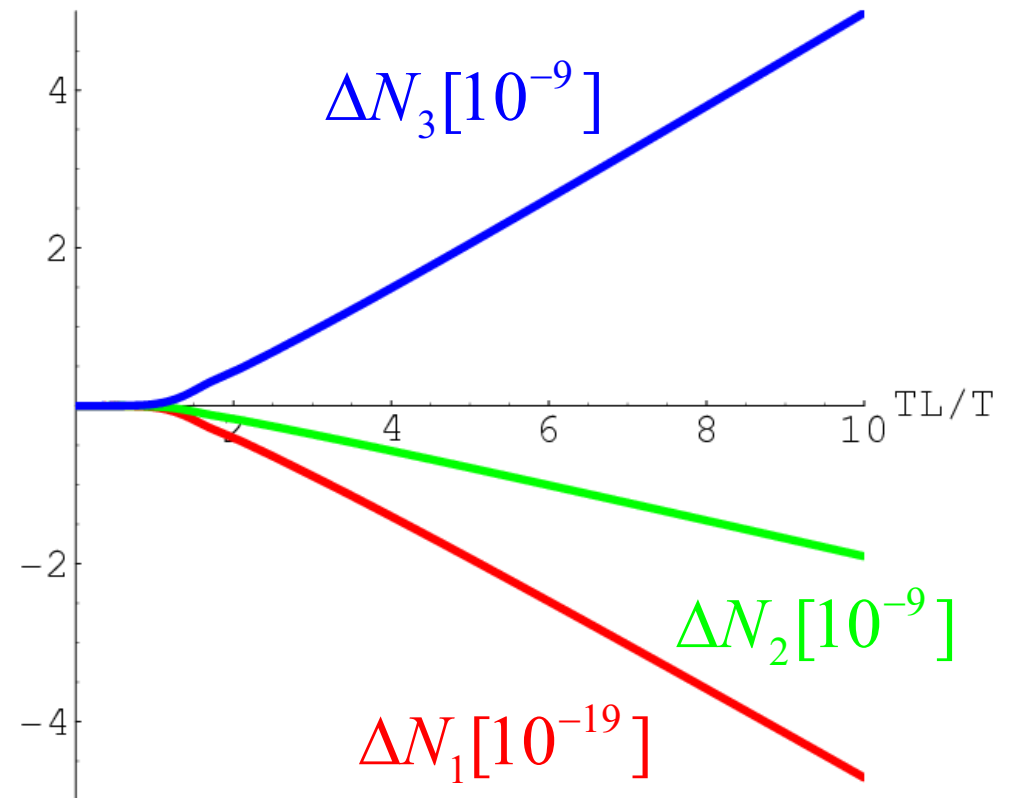


# Evolution of flavor asymmetries

Active sector

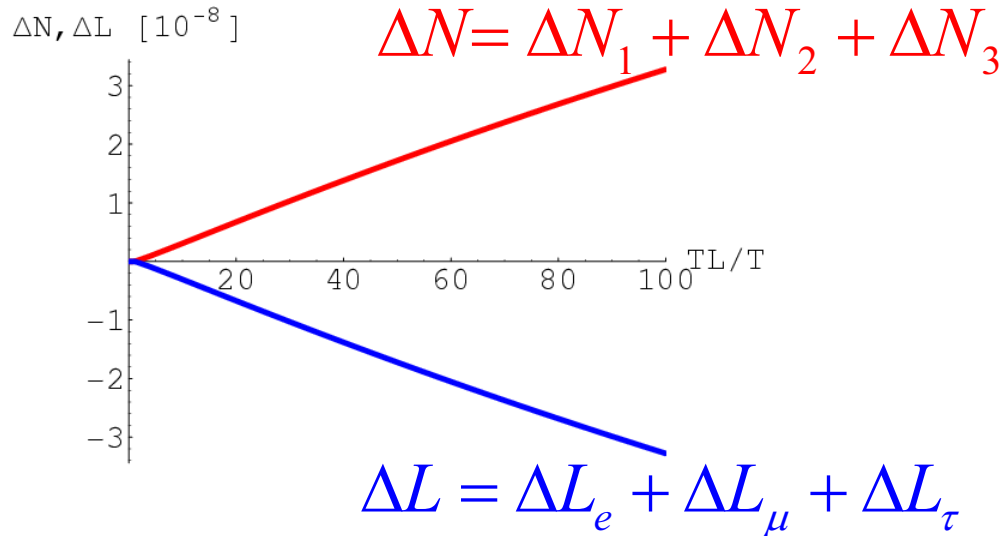


Sterile sector



$T_L \sim 10^4 \text{ GeV}$

# Baryon asymmetry of the universe



Asymmetry in active sector is partially converted into baryon asymmetry

$$\Delta B = -\frac{28}{79} \Delta L \neq 0$$

(Kuzmin, Rubakov, Shaposhnikov)

$$\frac{n_B}{s} \simeq 2 \times 10^{-10} \delta_{CP} \left( \frac{10^{-5}}{\Delta M_{32}^2 / M_3^2} \right)^{2/3} \left( \frac{M_3}{10 \text{ GeV}} \right)^{5/3} \quad (n_B / s)_{\text{OBS}} \sim 10^{-10}$$

- Heavier sterile neutrinos are quasi-degenerate

$$1 \text{ GeV} < M_2, M_3 < 17 \text{ GeV}$$

BBN  $\swarrow$   $\nwarrow$  Out-of-equilibrium

# Summary

The MSM + three right-handed neutrinos ( $\nu$ MSM)

- Lightest sterile neutrino  $\sim 10\text{keV}$  can be dark matter candidate

- The simplest Dodelson-Widrow scenario conflicts with X-ray and Ly-alpha constraints
- Some other production mechanism is needed

- Heavier sterile neutrinos can be responsible to baryon asymmetry of the universe

- Baryogenesis via neutrino oscillations (not leptogenesis)

# Three sterile neutrinos

We may call as "dark", "clear" and "bright"

- Dark,  $N_D$ : Lightest one

- Dark matter (production?)
- $M \sim \text{keV}$ ,  $F < 10^{-12}$ ,  $\theta < 10^{-4}$

- Clear and Bright,  $N_C$  and  $N_B$ : Heavier ones

- Baryon asymmetry
- Neutrino oscillations
- $M \sim 1-10 \text{ GeV}$ ,  $F \sim 10^{-7}$ ,  $\theta \sim 10^{-6}$

# Neutrino masses

**Seesaw mechanism: Dirac masses  $\gg$  Majorana masses**

$$\begin{pmatrix} 0 & M_D \\ M_D^T & M_M \end{pmatrix} \Rightarrow \begin{pmatrix} M_\nu & 0 \\ 0 & M_M \end{pmatrix} \quad M_\nu = -M_D^T \frac{1}{M_M} M_D$$

- **Active neutrinos**

$$\nu_1, \nu_2, \nu_3 \quad (m_1 \leq m_2 \leq m_3)$$

$$U^T M_\nu U = \text{diag}(m_1, m_2, m_3)$$

- **Sterile neutrinos**

$$N_1, N_2, N_3 \quad (M_1 \leq M_2 \leq M_3)$$

$$M_M = \text{diag}(M_1, M_2, M_3)$$

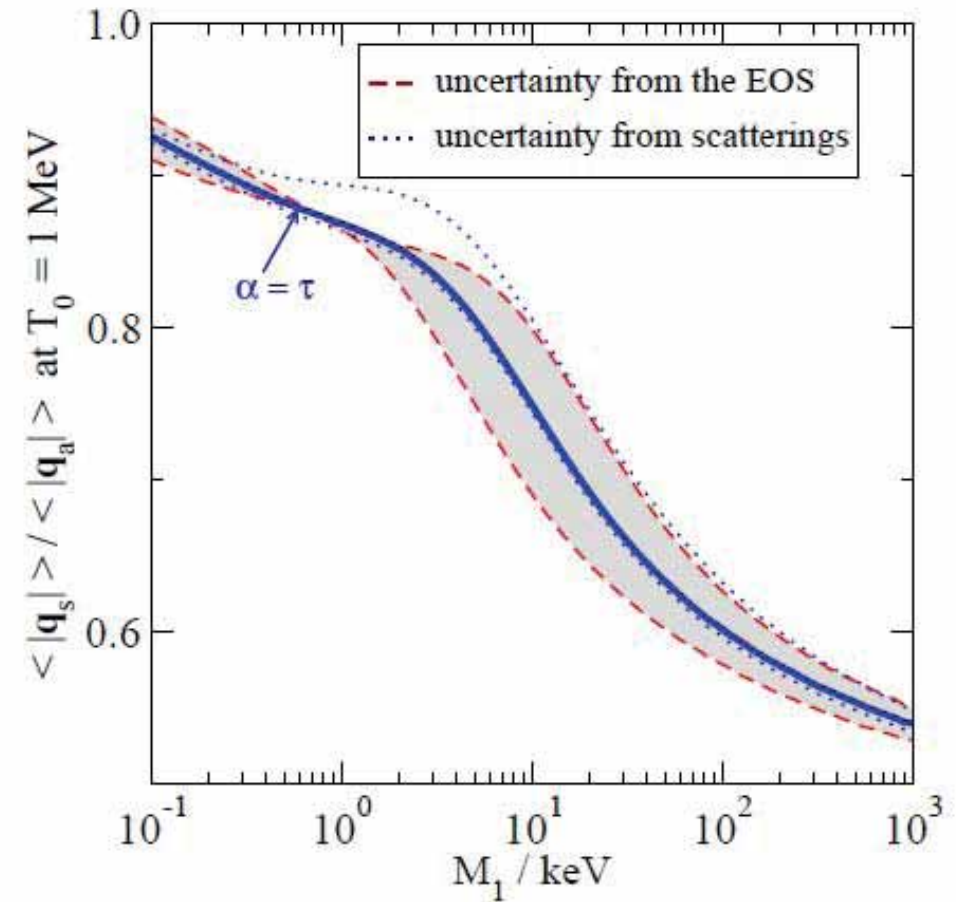
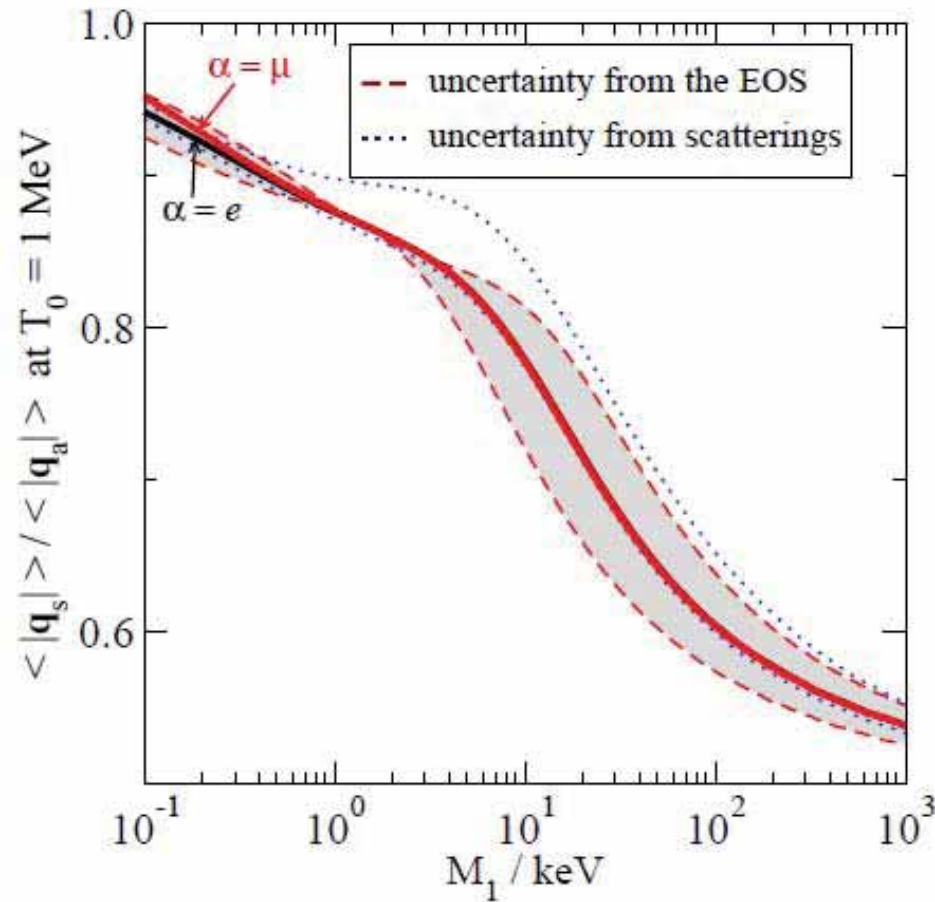
**Neutrino mixing in charged current**

$$\nu_\alpha = U_{\alpha a} \nu_a + \Theta_{\alpha I} N_I^c$$

$(\alpha = e, \mu, \tau) \quad (a = 1, 2, 3)$

$$\Theta_{\alpha I} = (M_D)_{\alpha I} / M_I \ll 1$$

# Averaged momentum



We find:  $\frac{\langle |\mathbf{q}_s| \rangle}{\langle |\mathbf{q}_a| \rangle} \simeq 0.8$  for  $M_1 \simeq 10 \text{ keV}$



$M_1 \gtrsim 11.6 \text{ keV}$  (SMMT)

$M_1 \gtrsim 8 \text{ keV}$  (VLHMR)

# Neutrino Yukawa coupling constants

$$F = K_L \cdot P_\alpha \cdot F_d \cdot K_R^\dagger \cdot P_\beta$$

$$F_d = \text{diag}(f_1, f_2, f_3)$$

$$P_\alpha = \text{diag}(e^{i\alpha_1}, e^{i\alpha_2}, 1), \quad P_\beta = \text{diag}(e^{i\beta_1}, e^{i\beta_2}, 1)$$

$$K_L = \begin{pmatrix} 1 & & \\ & c_{L23} & s_{L23} \\ & -s_{L23} & c_{L23} \end{pmatrix} \begin{pmatrix} c_{L13} & & s_{L13}e^{-i\delta_L} \\ & 1 & \\ -s_{L13}e^{i\delta_L} & & c_{L13} \end{pmatrix} \begin{pmatrix} c_{L12} & s_{L12} & \\ -s_{L12} & c_{L12} & \\ & & 1 \end{pmatrix}$$
$$K_R = \begin{pmatrix} 1 & & \\ & c_{R23} & s_{R23} \\ & -s_{R23} & c_{R23} \end{pmatrix} \begin{pmatrix} c_{R13} & & s_{R13}e^{-i\delta_R} \\ & 1 & \\ -s_{R13}e^{i\delta_R} & & c_{R13} \end{pmatrix} \begin{pmatrix} c_{R12} & s_{R12} & \\ -s_{R12} & c_{R12} & \\ & & 1 \end{pmatrix}$$



# Implication to active neutrino masses

## The lightest active neutrino mass

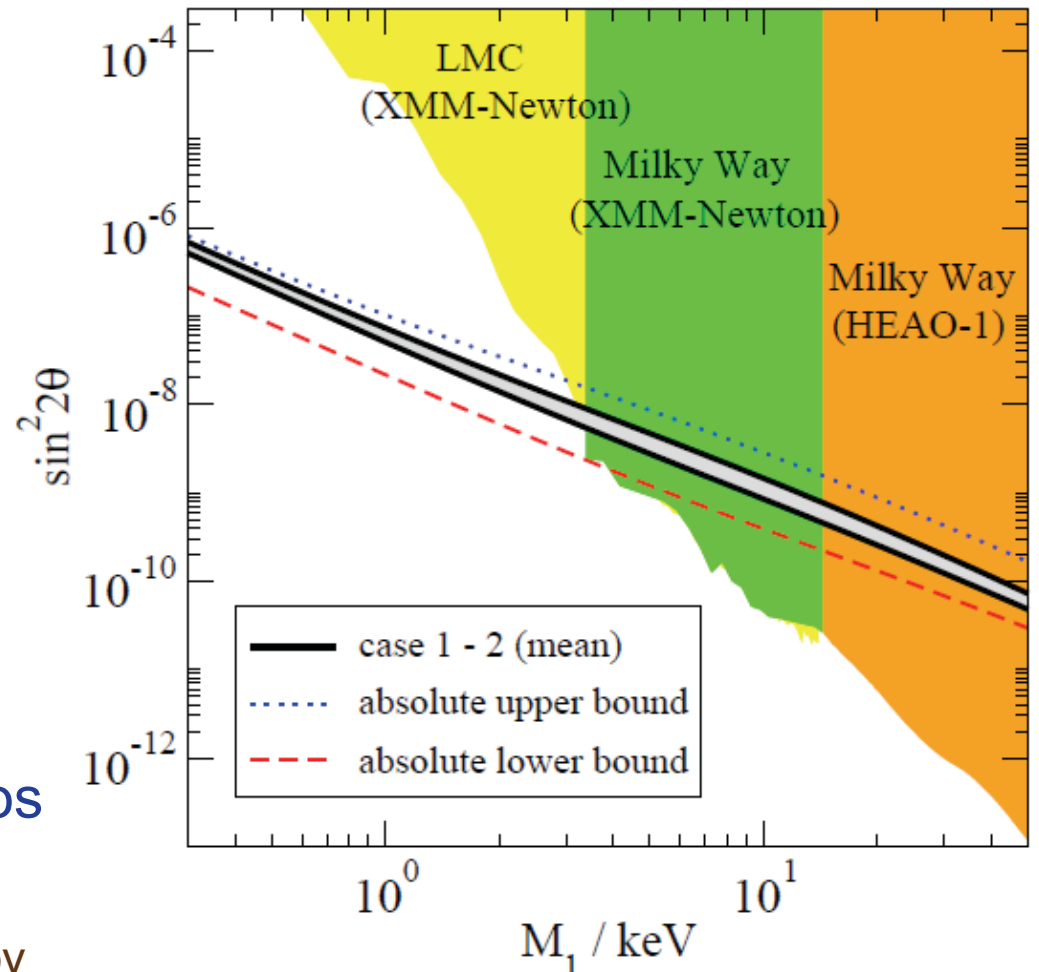
$$m_{\nu_1} \lesssim \sum_{\alpha=e,\mu,\tau} \frac{|M_D|_{\alpha 1}^2}{M_1}$$
$$\approx 2 \cdot 10^{-5} \text{eV} \left( \frac{\text{keV}}{M_1} \right)$$

- Tremaine-Gunn bound:

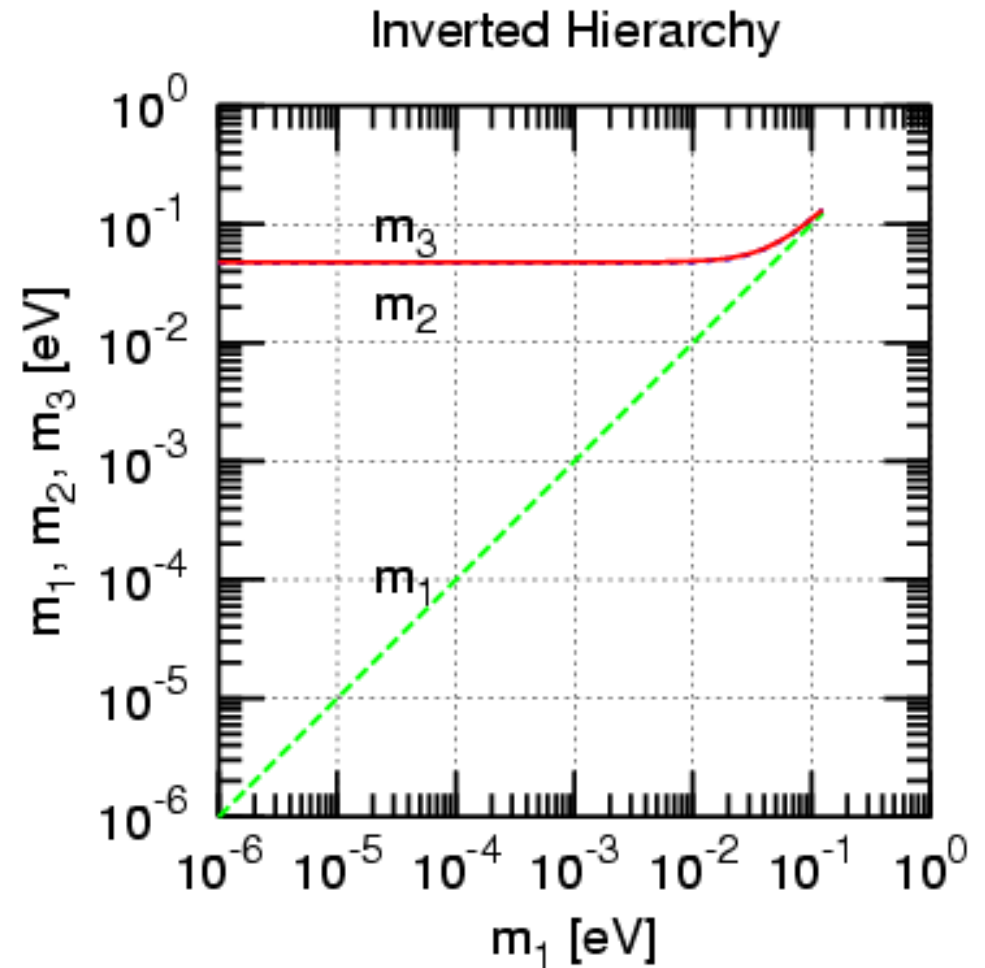
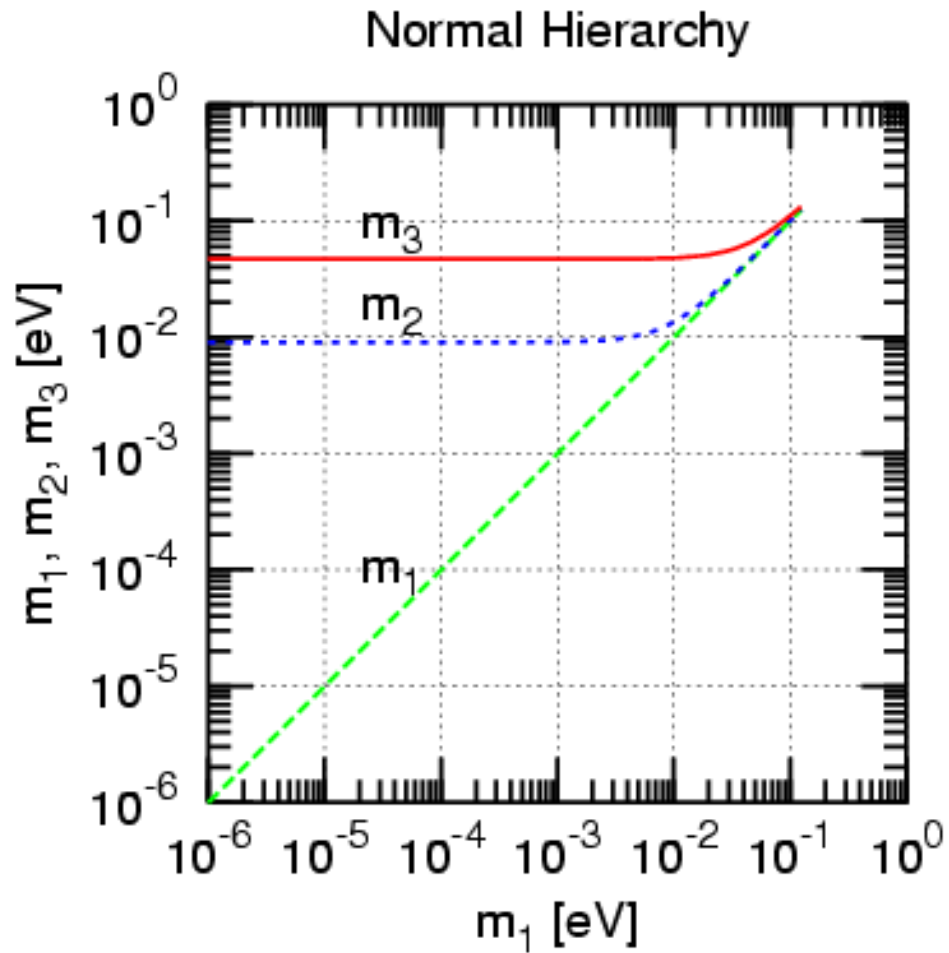
$$M_1 \gtrsim 0.3 \text{keV}$$

➔  $m_{\nu_1} \lesssim 7 \cdot 10^{-5} \text{eV}$

➔ exclude the degenerate masses of active neutrinos



# Active neutrino masses



# Entropy dilution

TA, Kusenko, Shaposhnikov

Decays of heavier sterile neutrinos can produce entropy  
by  $S \lesssim 30$  within the  $\nu$ MSM

- Bound from overclosure:  $\sin^2(2\theta) < f(M_1) \cdot S$
- Bound from X-rays: does not change
- Bound from Ly- $\alpha$ :  $M_1 > \frac{\langle |\mathbf{q}_s| \rangle}{\langle |\mathbf{q}_a| \rangle} M_0 \cdot S^{-1/3}$

## ■ Allowed region opens

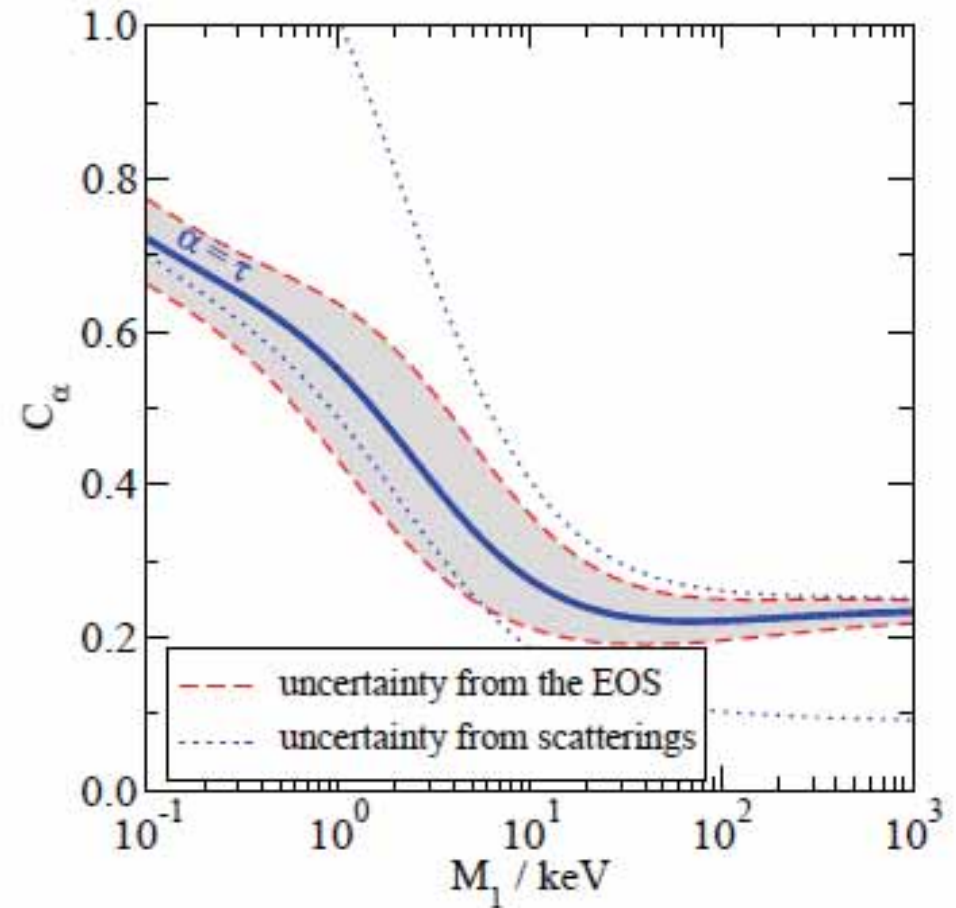
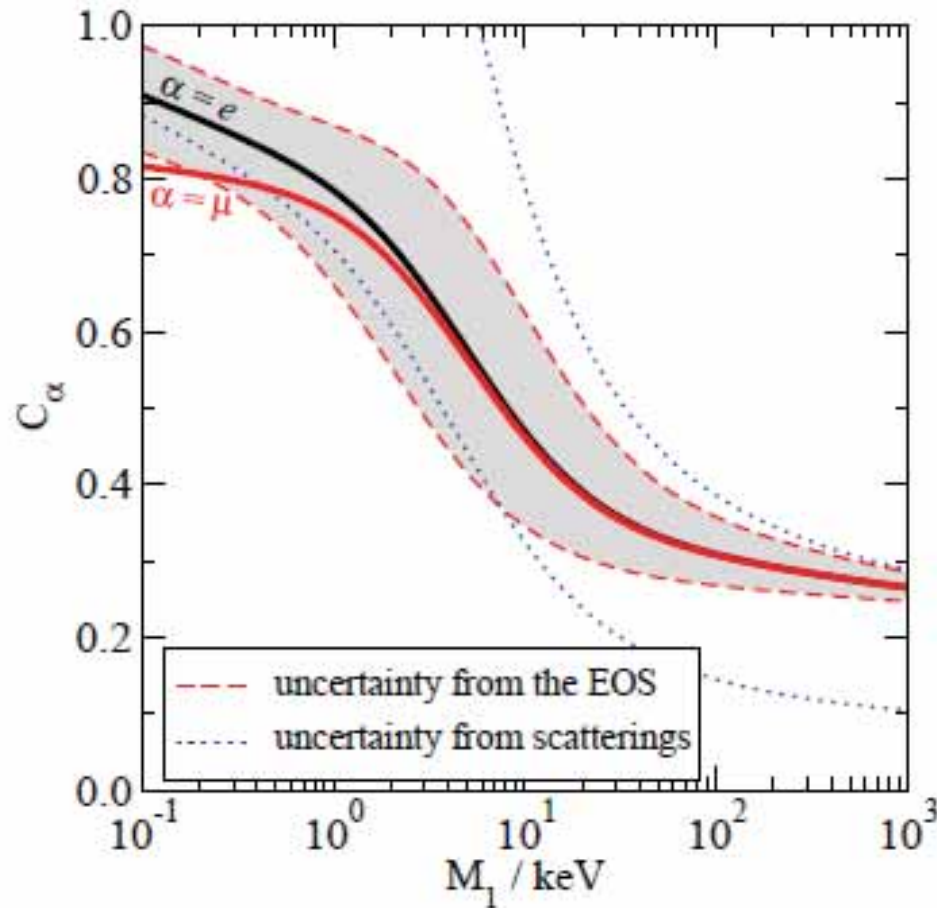
$$S \gtrsim 155, \quad (M_1 \lesssim 1.5\text{keV}, \quad \sin^2(2\theta) < 1.9 \cdot 10^{-6})$$

for  $M_1 \gtrsim 8\text{keV} \cdot S^{-1/3}$  (VLHMR)

- Such rate cannot be obtained within the  $\nu$ MSM
- Sterile neutrino dark-matter cannot be realized within the  $\nu$ MSM, if we apply the current Ly- $\alpha$  constraints !!

# Relic density of sterile neutrino

$$\Omega_{N_1} h^2 = 0.11 \sum_{\alpha=e,\mu,\tau} C_\alpha(M_1) \left( \frac{|M_D|_{\alpha 1}}{0.1 \text{ eV}} \right)^2$$



$$\Omega_{N_1} = \Omega_{\text{DM}} \rightarrow |M_D|_{\alpha I} \sim 0.1 \text{ eV}, \quad |F_{\alpha I}| \sim 10^{-12}$$

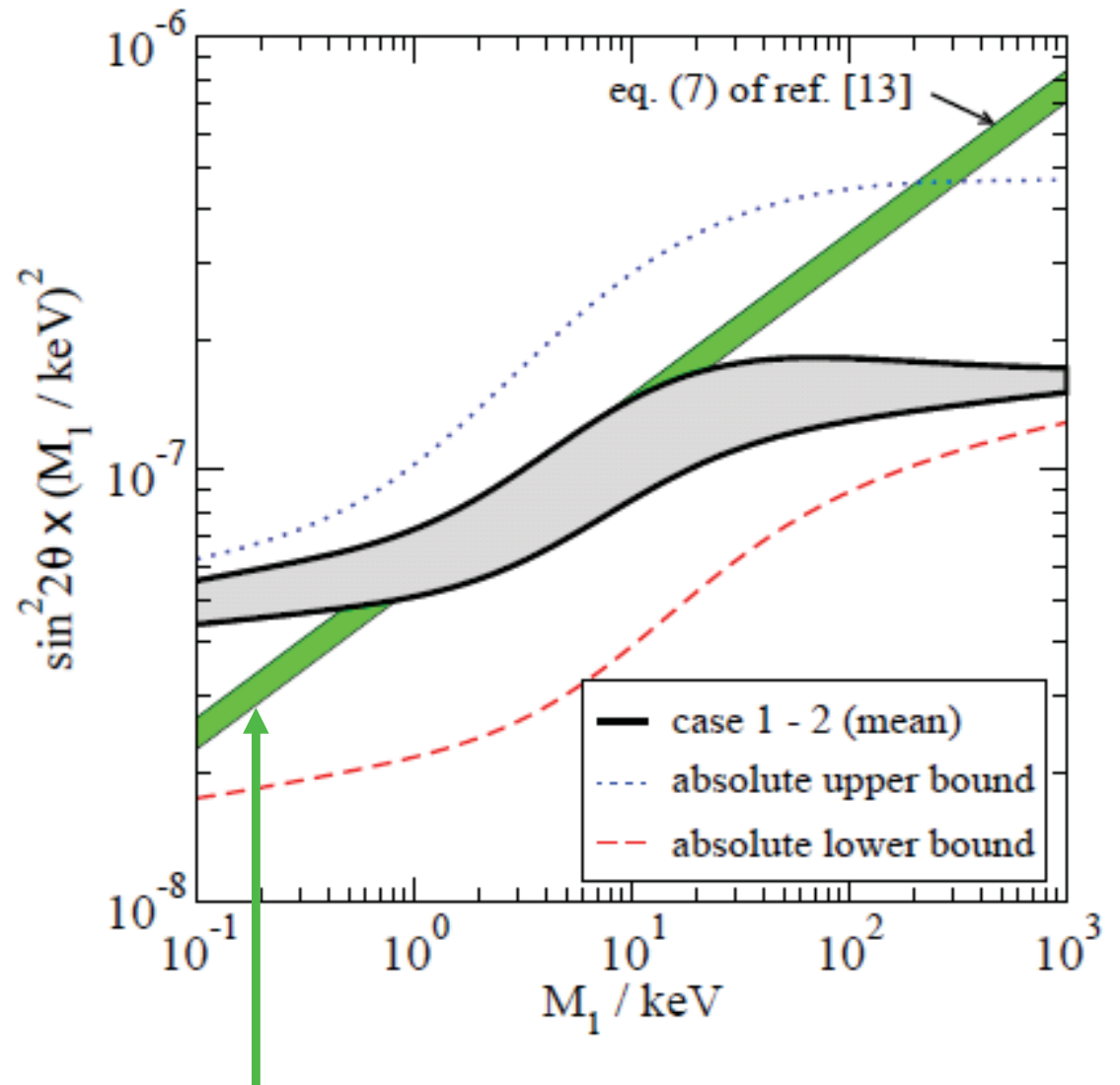
# Upper bound on mixing angle

We find  $C_e > C_\mu > C_\tau$

- The case 1:  
the largest abundance  
the strongest bound  
 $|M_D|_{e1} \neq 0$   
 $|M_D|_{\mu1} = |M_D|_{\tau1} = 0$
- The case 2:  
the smallest abundance  
the weakest bound

$$|M_D|_{\tau1} \neq 0$$

$$|M_D|_{e1} = |M_D|_{\mu1} = 0$$



Abazajian: Phys.Rev. D73 ('06) 063506

$\Omega_{DM} = 0.2$ ,  $T_c = 150 - 200\text{MeV}$