

# Probing New Physics Through $B_s$ Mixing

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# Phenomenology of $B_s$ Mixing

- $\bar{B}^0 = (b\bar{s})$  and  $B^0 = (s\bar{b})$  with definite flavour content, but not mass eigenstates  $\rightsquigarrow$  **Mixing!**
- mixing induces mass and width mixing matrices  $M_{ab}^s, \Gamma_{ab}^s$
- $M_{12}^s$  is loop-induced: **sensitive to new physics!**
- observables:
  - $\Delta M_s = 2|M_{12}^s|, \quad \phi_s = \arg(-M_{12}^s/\Gamma_{12}^s) \approx \arg M_{12}^s$
  - CP asymmetry in  $B_s \rightarrow J/\psi\phi$
  - $\Delta\Gamma_s = 2|\Gamma_{12}^s| \cos \phi_s$
  - semileptonic CP asymmetry

$$A_{\text{SL}}^s = \frac{N(\bar{B}_s^0 \rightarrow \ell^+ X) - N(B_s^0 \rightarrow \ell^- X)}{N(\bar{B}_s^0 \rightarrow \ell^+ X) + N(B_s^0 \rightarrow \ell^- X)} = \frac{\Delta\Gamma_s}{\Delta M_s} \tan \phi_s$$

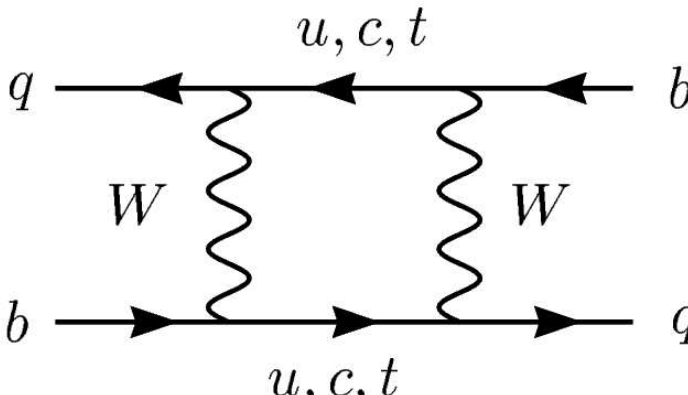
# Experimental status of $B_s$ Observables

- $\Delta M_s = (17.77 \pm 0.10 \pm 0.07) \text{ ps}^{-1}$  (CDF) (2006!)
- $\Delta\Gamma_s = (0.17 \pm 0.09 \pm 0.02) \text{ ps}^{-1}$  (D0) (2007!)
- $A_{SL}^s = 0.001 \pm 0.0090$  (D0) (2007!)
- $\phi_s = -0.70_{-0.39}^{+0.47}$  from  $A_{SL}^s$  and  $B_s \rightarrow J/\psi\phi$  (D0) (2007!)

For LHCb, expect (van Hune CKM06)

- sensitivity  $\sigma_{\phi_s} = 0.04 \text{ rad}$  @  $0.5 \text{ fb}^{-1}$ ,  $\sigma_{\phi_s} = 0.02 \text{ rad}$  @  $2 \text{ fb}^{-1}$
- first results promised for 2008

# $\Delta M_s$ in the SM



$$M_{12}^{\text{SM}} = \frac{G_F^2 M_W^2}{12\pi^2} M_{B_s} \hat{\eta}^B \hat{B}_{B_s} f_{B_s}^2 (V_{ts}^* V_{tb})^2 S_0(x_t)$$

- $S_0(x_t = m_t^2/M_W^2) = 2.35 \pm 0.06$ : Inami-Lim function
- $\hat{\eta}^B = 0.552$ : NLO QCD correction (Buras/Jamin/Weiss '90)
- $\hat{B}_{B_s} f_{B_s}^2 \propto \langle B_s^0 | (\bar{s}b)_{V-A} (\bar{s}b)_{V-A} | \bar{B}_s^0 \rangle$ : hadronic matrix element, from lattice
- $V_{ts}^* V_{tb}$ : from tree-level processes

# Status of Input Parameters

- $\sqrt{\hat{B}_{B_s}} f_{B_s} = (0.281 \pm 0.021) \text{ GeV}$  (HPQCD 06, unquenched,  $N_f = 2 + 1$  staggered light quarks)
- $\sqrt{\hat{B}_{B_s}} f_{B_s} = (0.245 \pm 0.021_{-0.002}^{+0.003}) \text{ GeV}$  (JLQCD 03, unquenched,  $N_f = 2$  Wilson light quarks)
- $|V_{ts}^* V_{tb}| = (41.3 \pm 0.7) \times 10^{-3}$
- SM predictions:

$$\Delta M_s|_{\text{JLQCD}} = (16.1 \pm 2.8) \text{ ps}^{-1}, \quad \Delta M_s|_{\text{HPQCD}} = (21.3 \pm 3.2) \text{ ps}^{-1}$$

- recall CDF result:  $(17.77 \pm 0.12) \text{ ps}^{-1}$

# New Physics Models with impact on $\Delta M_s$

- SUSY: gluino-squark loops (Ball hep-ph/0604249)
- SUSY: Higgs penguins at large  $\tan \beta$  (Freitas hep-ph/0702267)
- SUSY GUT (Dutta hep-ph/0607147)
- SUSY: R-parity violation (Xiang-Dong hep-ph/0609269)
- tree-level:  $Z'$  with flavour-non-diagonal couplings (Baek hep-ph/0607113)
- warped extra dimension (Chang hep-ph/0607313)
- littlest Higgs model (Buras hep-ph/0605214)

# Generic Parametrisation of New Physics

- new physics can significantly affect  $M_{12}^s$ , but not  $\Gamma_{12}^s$ : tree-dominated
- model-independent parametrisation of NP effects in terms of only two real parameters

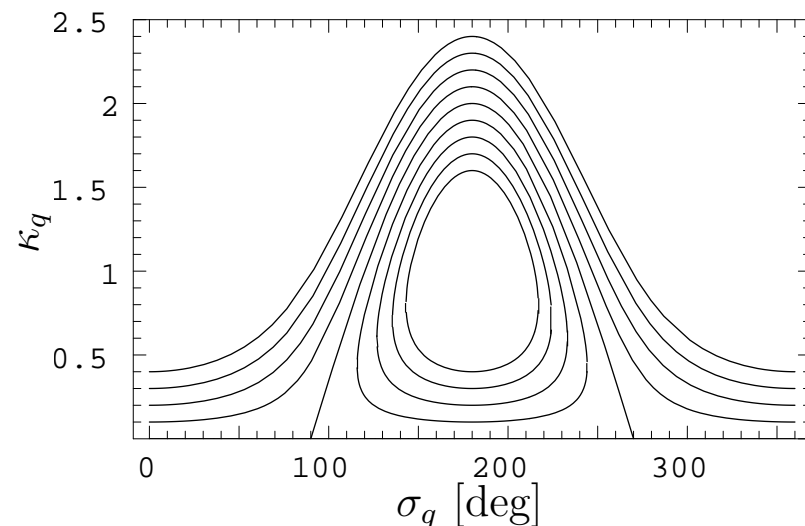
$$M_{12}^s = M_{12}^{s,\text{SM}} (1 + \kappa_s e^{i\sigma_s})$$

- $\kappa_s > 0$ : NP amplitude
- $\sigma_s$ : new CP-violating phase

Deviation from SM measured by

$$\rho_s \equiv \left| \frac{\Delta M_s}{\Delta M_s^{\text{SM}}} \right| = (1 + 2\kappa_s \cos \sigma_s + \kappa_s^2)^{1/2}$$

Lines of  $\rho_s = \text{const.}$ :

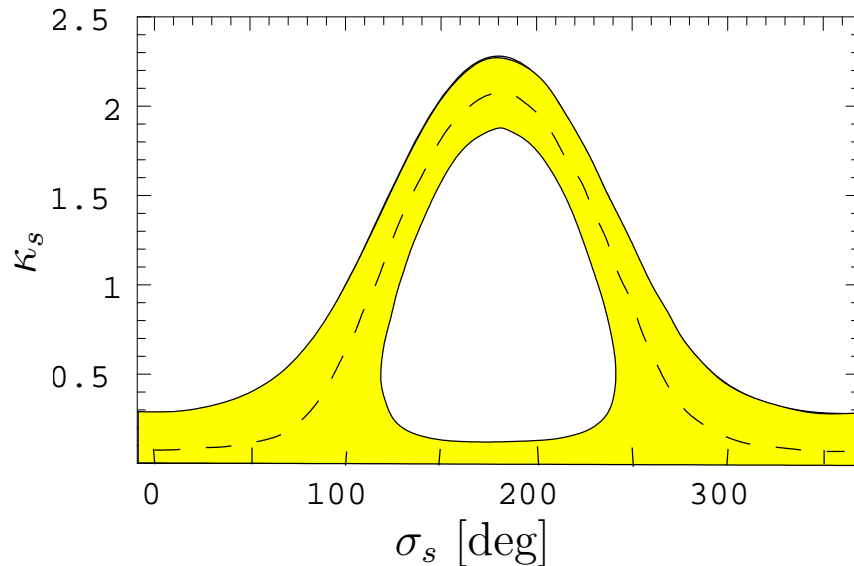


# Constraints from $\Delta M_s^{\text{SM}}$

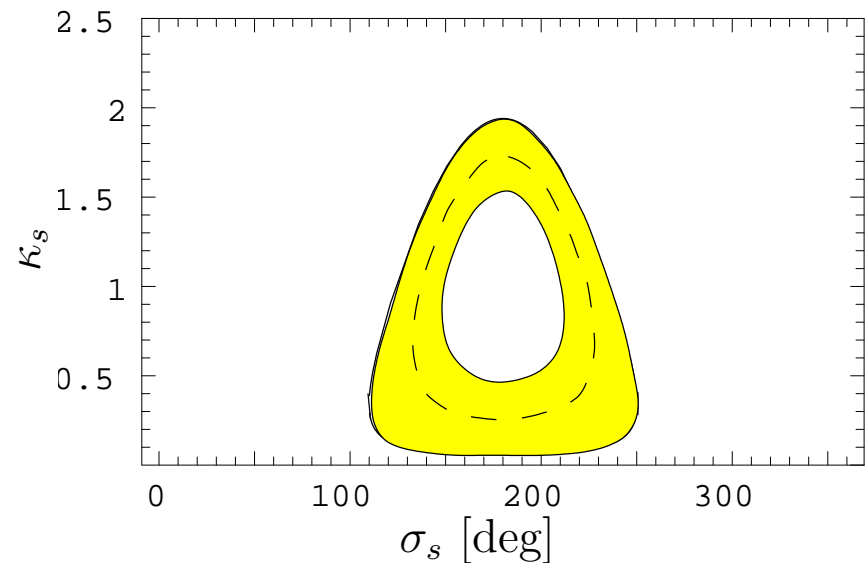
Recall: two unquenched calculations available:

- JLQCD:  $N_f = 2$  Wilson fermions + NRQCD heavy fermions
- HPQCD:  $N_f = 2 + 1$  staggered light + NRQCD heavy fermions

$1\sigma$  constraints from JLQCD:



from HPQCD:



$\Delta M_s|_{\text{exp}}$  compatible...

not compatible with SM (at  $1\sigma$ )



# Constraints from $\Delta M_s^{\text{SM}}$

## Conclusions:

- $\Delta M_s|_{\text{th}}$  **not yet known accurately enough** to exclude even  $|M_{12}^{s,\text{NP}}| \approx |M_{12}^{s,\text{SM}}|$  (i.e.  $\kappa_s < 1$ )
- improved predictions expected in due course thanks to recent **breakthrough in lattice algorithms** to reduce the cost of simulations of light quark masses for Wilson fermions

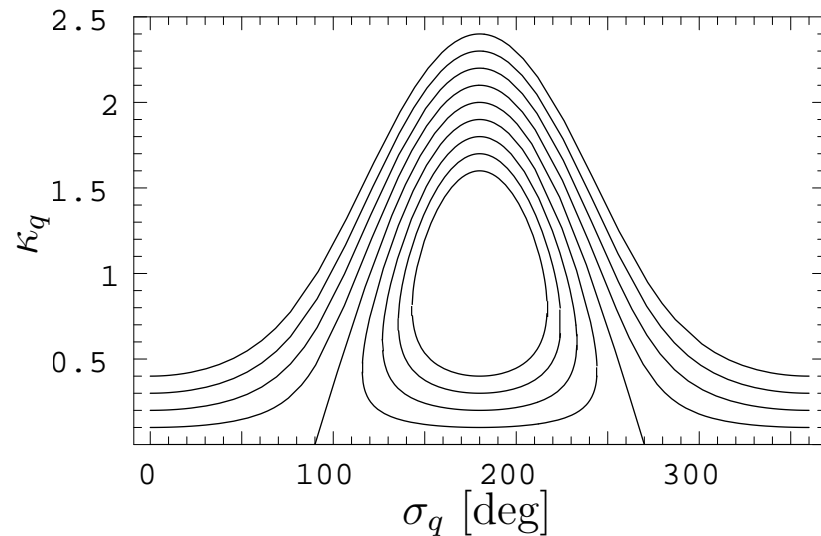
(Del Debbio, Lüscher 06)

- timescale? Current uncertainty of  $\Delta M_s|_{\text{th}} \sim 14\%$ . Sachrajda quoted 1-2% for 2015 (LHCb upgrade workshop Jan 2007)
- also use alternative constraints  $\rightarrow \phi_s$

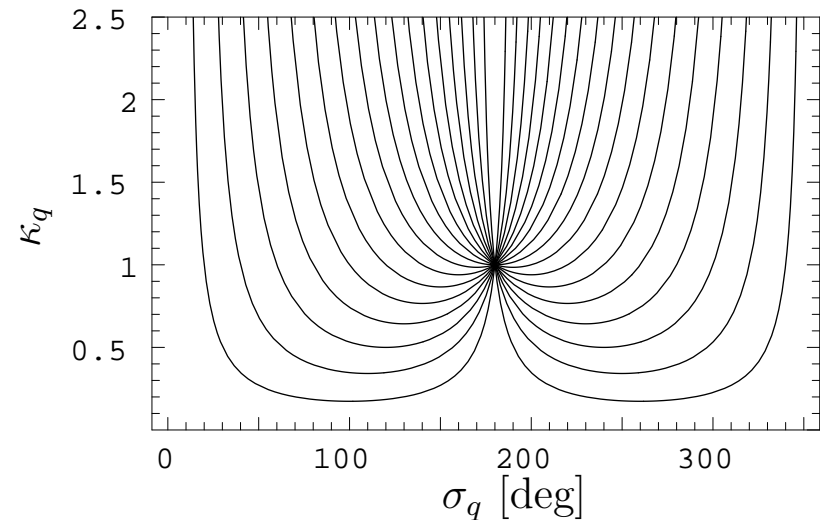
# Constraints from $\phi_s$

$$\phi_s = \arg M_{12}^s = \phi_s^{\text{SM}} + \phi_s^{\text{NP}} \text{ with } \phi_s^{\text{SM}} = -2\lambda^2 R_b \sin \gamma \approx -2^\circ$$

lines of  $\rho_s = \text{const.}$ :



lines of  $\phi_s^{\text{NP}} = \text{const.}$ :



Observables sensitive to  $\phi_s$ :

CP asymmetries in  $B_s \rightarrow J/\psi\phi$  (and related),  $\Delta\Gamma_s$ ,  $A_{\text{SL}}^s$

# Status of $\phi_s$

$\phi_s$  from  $b \rightarrow ccs$  decays, e.g.  $B_s \rightarrow J/\psi\phi$ :

$$\Gamma(t) \propto e^{-\Gamma_s t} \left\{ \cosh \frac{\Delta\Gamma_s t}{2} - \eta_f \cos \phi_s \sinh \frac{\Delta\Gamma_s t}{2} + \eta_f q D \sin \phi_s \sin(\Delta M_s t) \right\}$$

$\eta_f$ : CP eigenvalue of final state ( $\pm 1$ ),  $q = +1(-1)$  if tagged as  $B_s$  ( $\bar{B}_s$ ) at production,  $q = 0$  if untagged,  $D$ : tagging dilution factor

Untagged events:  $\cos \phi_s$  and  $\Delta\Gamma_s$  from 2nd term: Tevatron 2007 (up to discrete ambiguities,  $\phi_s \leftrightarrow -\phi_s$ ,  $\Delta\Gamma_s \leftrightarrow -\Delta\Gamma_s$ ; theory-indep.):

$$\text{D0: } \Delta\Gamma_s = (0.17 \pm 0.09 \pm 0.02) \text{ ps}^{-1}, \quad \phi_s = -0.79 \pm 0.56_{-0.01}^{+0.14}$$

Tagged events:  $\sin \phi_s$  from 3rd term: LHCb 200x (up to dilution factor  $D$ , to be determined from control channels, and discrete ambiguity  $\phi_s \leftrightarrow \pi - \phi_s$ ; theory-indep.)

# Status of $\phi_s$

$\phi_s$  from **flavour-specific** (vulgo: semileptonic) CP-asymmetry  $a_{fs}^s$ , e.g.

$$A_{SL}^s = \frac{N(\bar{B}_s^0 \rightarrow \ell^+ X) - N(B_s^0 \rightarrow \ell^- X)}{N(\bar{B}_s^0 \rightarrow \ell^+ X) + N(B_s^0 \rightarrow \ell^- X)} = \frac{\Delta\Gamma_s}{\Delta M_s} \tan \phi_s$$

Also measurement of **dimuon charge asymmetry**: (D0 06)

$$A_{SL}^{\mu\mu} = \frac{N(b\bar{b} \rightarrow \mu^+ \mu^+ X) - N(b\bar{b} \rightarrow \mu^- \mu^- X)}{N(b\bar{b} \rightarrow \mu^+ \mu^+ X) + N(b\bar{b} \rightarrow \mu^- \mu^- X)} = (-0.92 \pm 0.44 \pm 0.32) \times 10^{-2}$$

• contains contributions from  $B_d$  and  $B_s$ , but not  $B_u$ ,  $\Lambda_b$  etc.

D0 translates both into  $A_{SL}^s = (-0.64 \pm 1.01) \times 10^{-2}$ .

Using the constraint from  $A_{SL}^s$  in the analysis of  $B_s \rightarrow J/\psi\phi$ :

$$\text{D0: } \Delta\Gamma_s = (0.13 \pm 0.09) \text{ps}^{-1}, \quad \phi_s = -0.70_{-0.39}^{+0.47}$$

# Theory Predictions for $\Delta\Gamma_s$ and $A_{\text{SL}}^s$

$\Delta\Gamma = 2|\Gamma_{12}^s| \cos \phi_s$  with

$$\Gamma_{12}^s = -\frac{G_F^2 m_b^2}{24\pi M_{B_s}} (V_{cs}^* V_{cb})^2 [C_1 \langle O_1^s \rangle + C_2 \langle O_2^s \rangle + \delta_{1/m_b}]$$

**Lenz/Nierste** hep-ph/0612167 improve upon Beneke et al. hep-ph/9808385 by switching to a different operator basis (reduced sensitivity to power-suppressed contributions):

$$\Delta\Gamma_s = (0.088 \pm 0.017) \text{ ps}^{-1}, \quad A_{\text{SL}}^s = (2.06 \pm 0.57) \times 10^{-5}$$

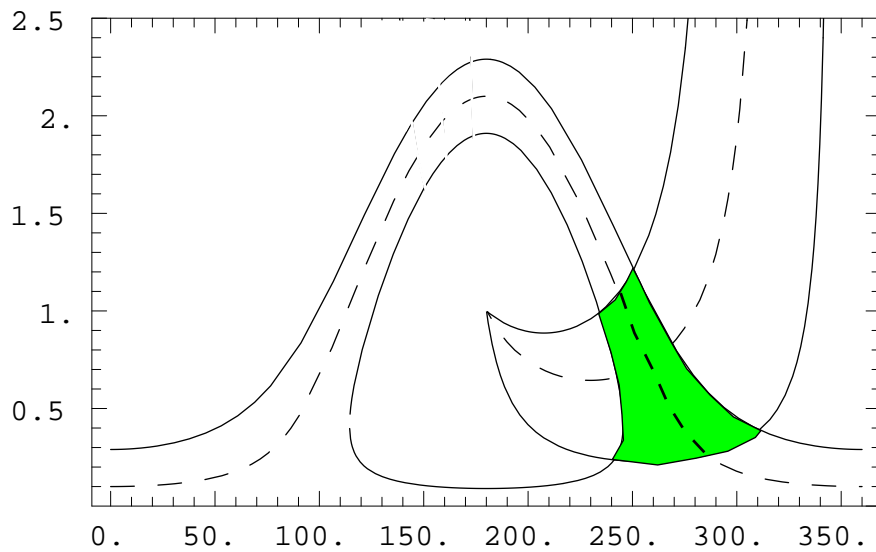
Using theory predictions for the  $B_d$  contribution to  $A_{\text{SL}}^{\mu\mu}$  (instead of exp. data from the B factories), and the D0 data for  $B_s \rightarrow J/\psi\phi$ , **LN** find

$$\sin \phi_s = -0.77 \pm 0.04(\text{th}) \pm 0.34(\text{exp}): \phi_s \neq 0 \text{ at } 2\sigma$$

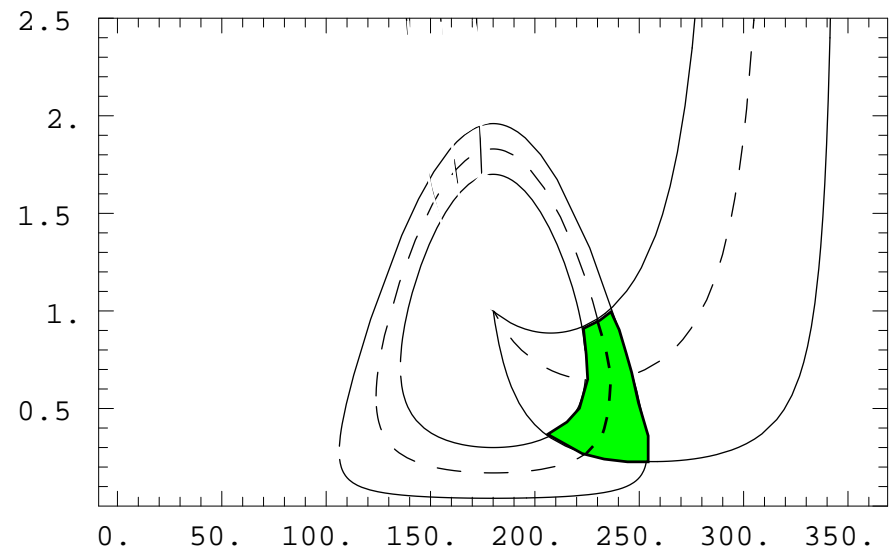
# Combined Constraint on $\kappa_s, \sigma_s$

Using D0 combined result:

$1\sigma$  constraints from JLQCD:



from HPQCD:



•  $|M_{12}^{s,\text{NP}}| \approx |M_{12}^{s,\text{SM}}|$  still allowed

•  $\sigma_s < 180^\circ$  disfavoured

# Constraints on Specific NP Models: $Z'$

- assume absence of  $Z-Z'$  mixing, i.e. flavour-diagonal  $Z$  couplings
- assume flavour non-diagonal  $Z'$  couplings only to  $q_L$
- constrain  $\rho_L \exp(i\phi_L) \equiv (g' M_Z)/(g M_{Z'}) B_{sb}^L$  with  $B_{sb}^L$  being  $\bar{s} Z' b$  coupling
- $\kappa_s < 1.2 \iff \rho_L < 1.3 \cdot 10^{-3}$
- can translate this into bound on  $Z'$  mass:

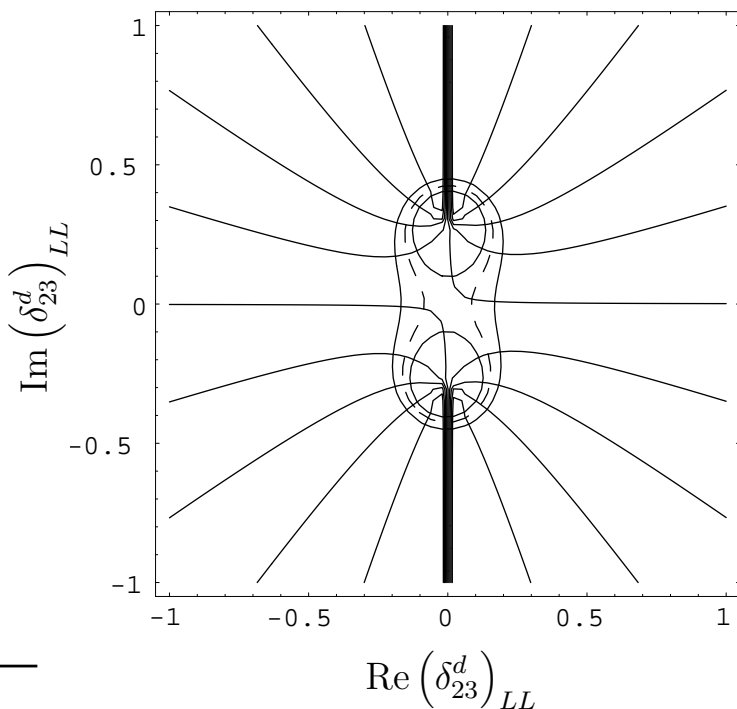
$$3.0 \text{ TeV} \left( \frac{g'}{g} \right) \left| \frac{B_{sb}^L}{V_{ts}} \right| < M_{Z'}$$

- very strong constraint!

# MSSM (in MIA)

- MSSM (box diagram) contributions from charged Higgs, neutralinos, photinos, gluinos and charginos\*
- for  $B_s$  mixing, only gluino contributions relevant
- full NLO Wilson coefficients available (Ciuchini hep-ph/0606197)

\* also from double Higgs penguins, which are however only relevant for large  $\tan\beta$



Constraints on  $(\delta_{23}^d)_{LL}$  insertion using JLQCD lattice data.  
Open lines: constraints from  $\phi_S$ .



# Summary

- new physics in  $B_s$  mixing can be described, in a **model-independent way**, by **2 real parameters**
- constraints from  $\Delta M_s$  not very strong, despite excellent experimental accuracy: large hadronic **(lattice) uncertainties**  $\sim 14\%$
- more decisive constraints from NP mixing phase
- present constraints on  $\phi_s$  from  $A_{\text{SL}}^s$  and untagged  $B_s \rightarrow J/\psi\phi$  (Tevatron) not very strong, although **slight hint at NP** ( $2\sigma$  acc. to Lenz/Nierste)
- expect much more precise constraints on/measurement of  $\phi_s$  from LHCb:  $\sigma_{\phi_s} = 0.04$  with  $0.5 \text{ fb}^{-1}$ : 1st year of running
- together with improved lattice calculations, the  $B_s$  system will prove a powerful tool to constrain/find NP