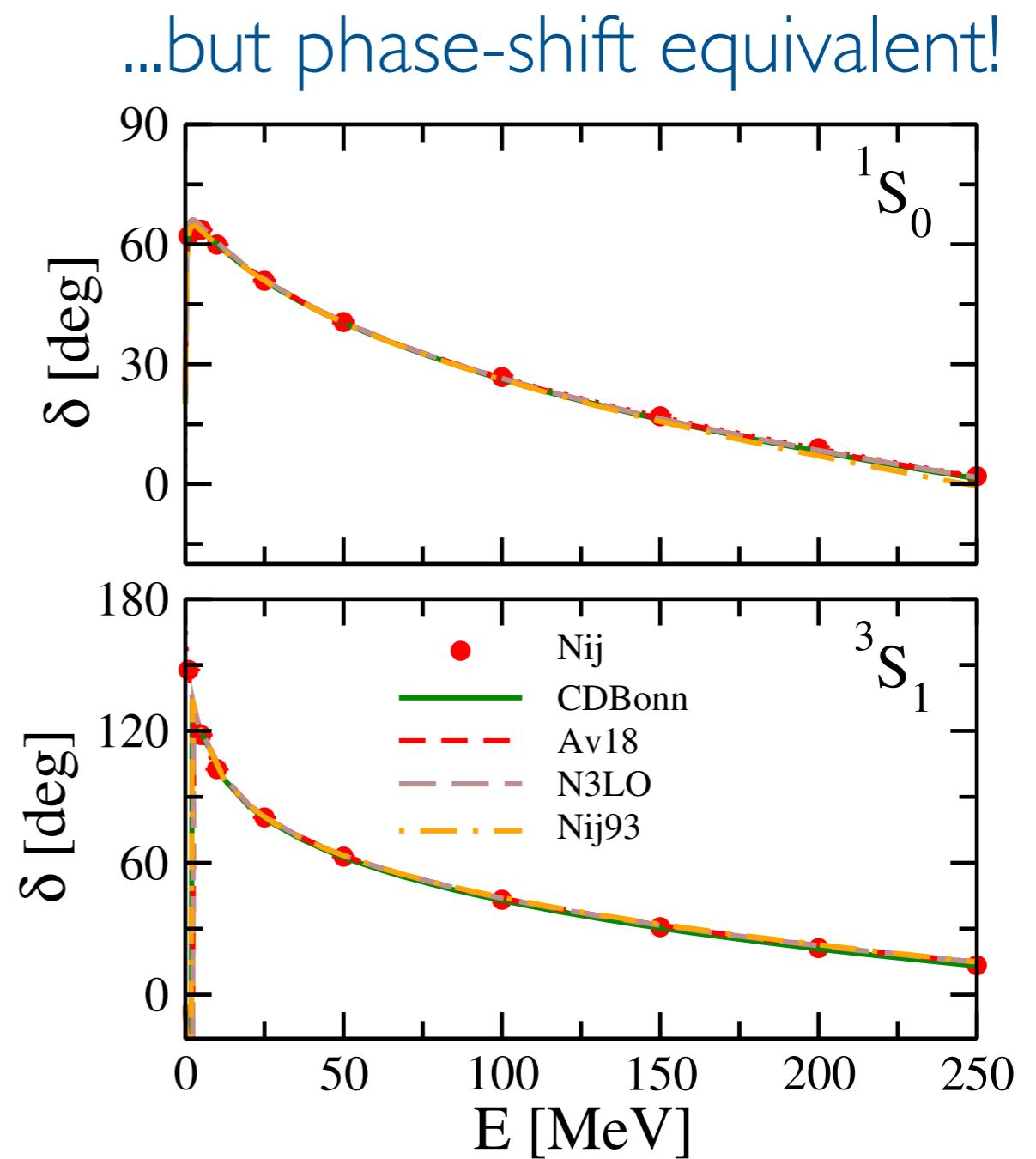
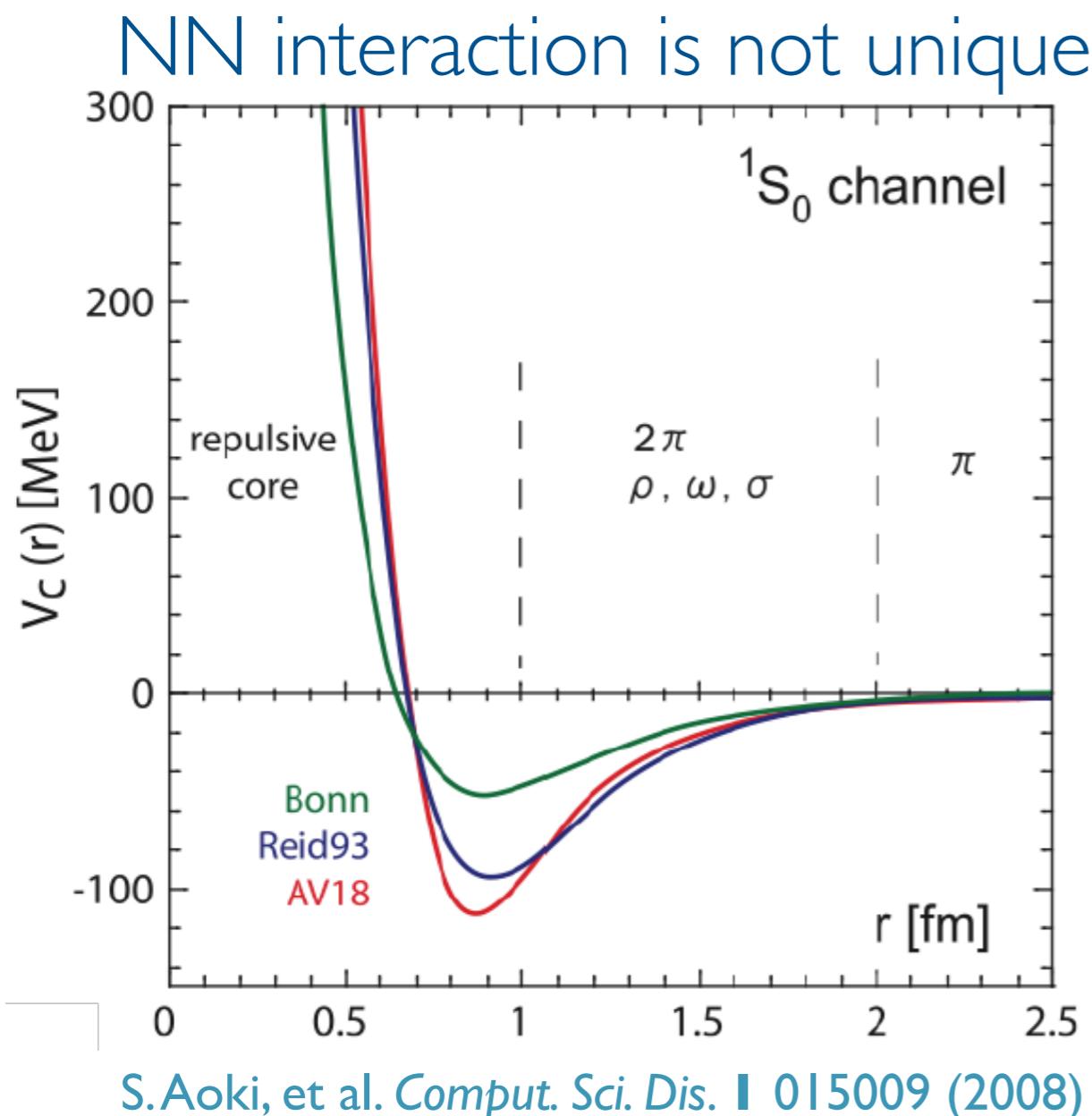


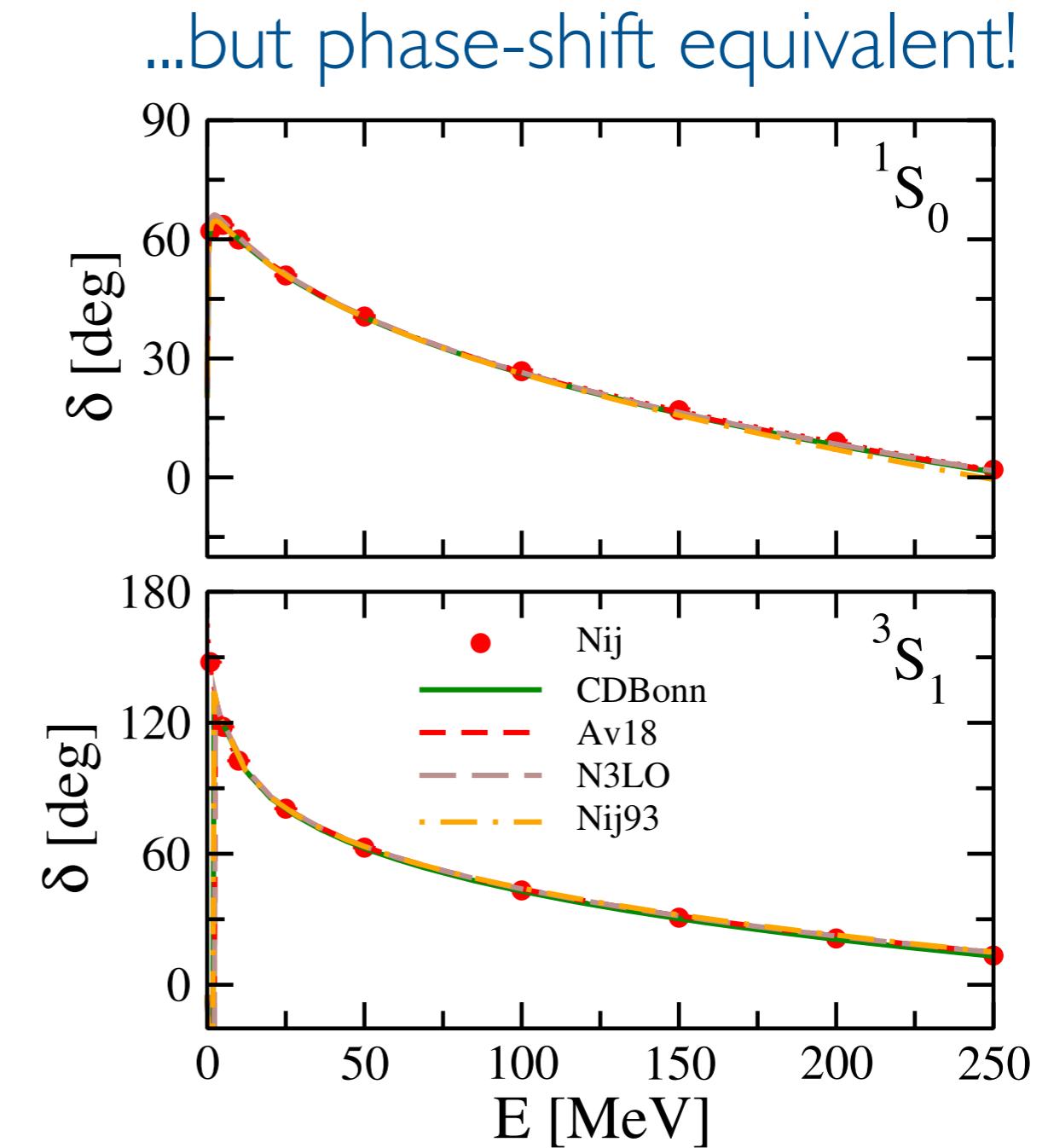
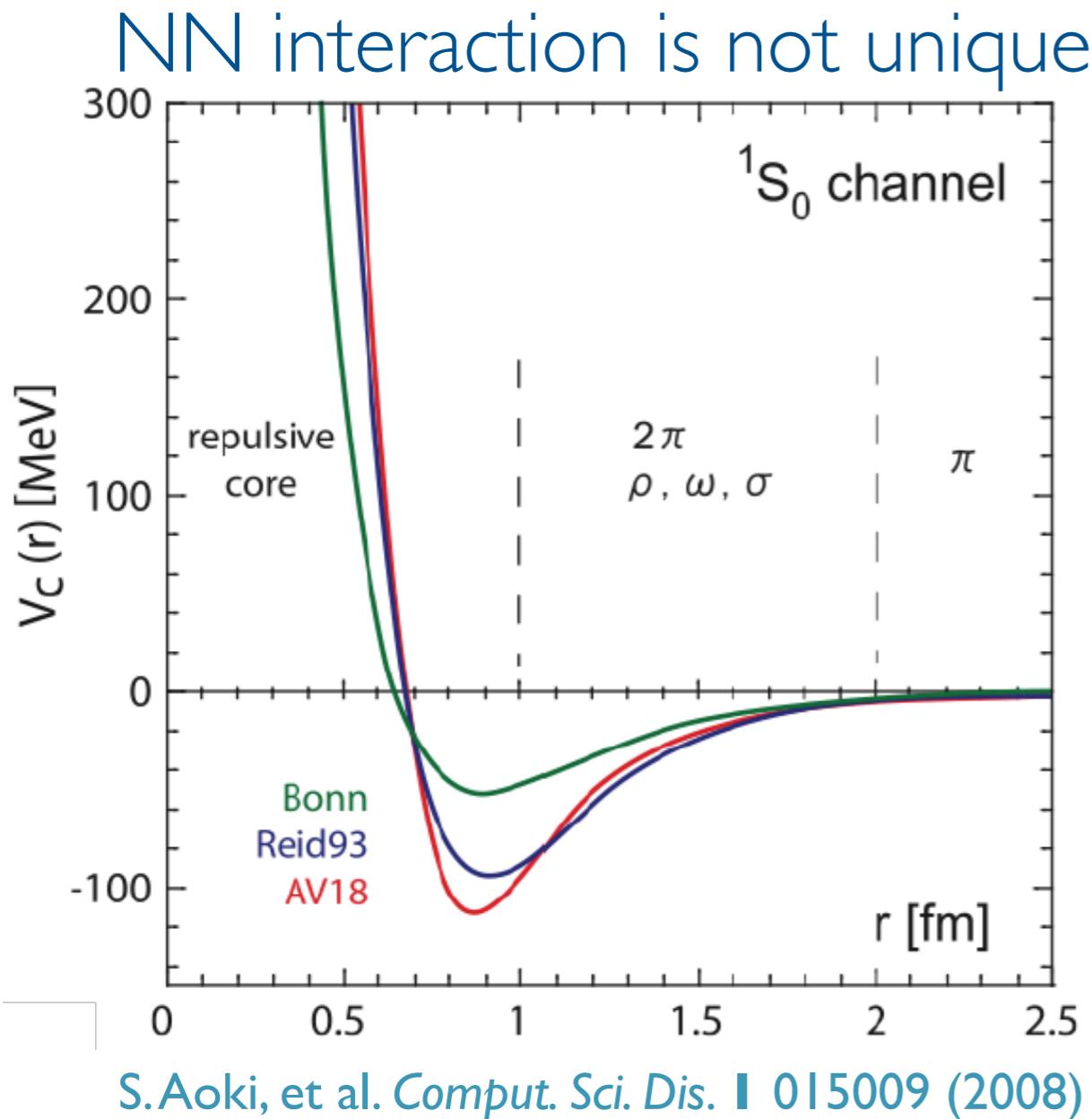
Green's functions for dense matter & mean-field connections

Arnau Rios Huguet
Lecturer in Nuclear Theory
Department of Physics
University of Surrey

Complications

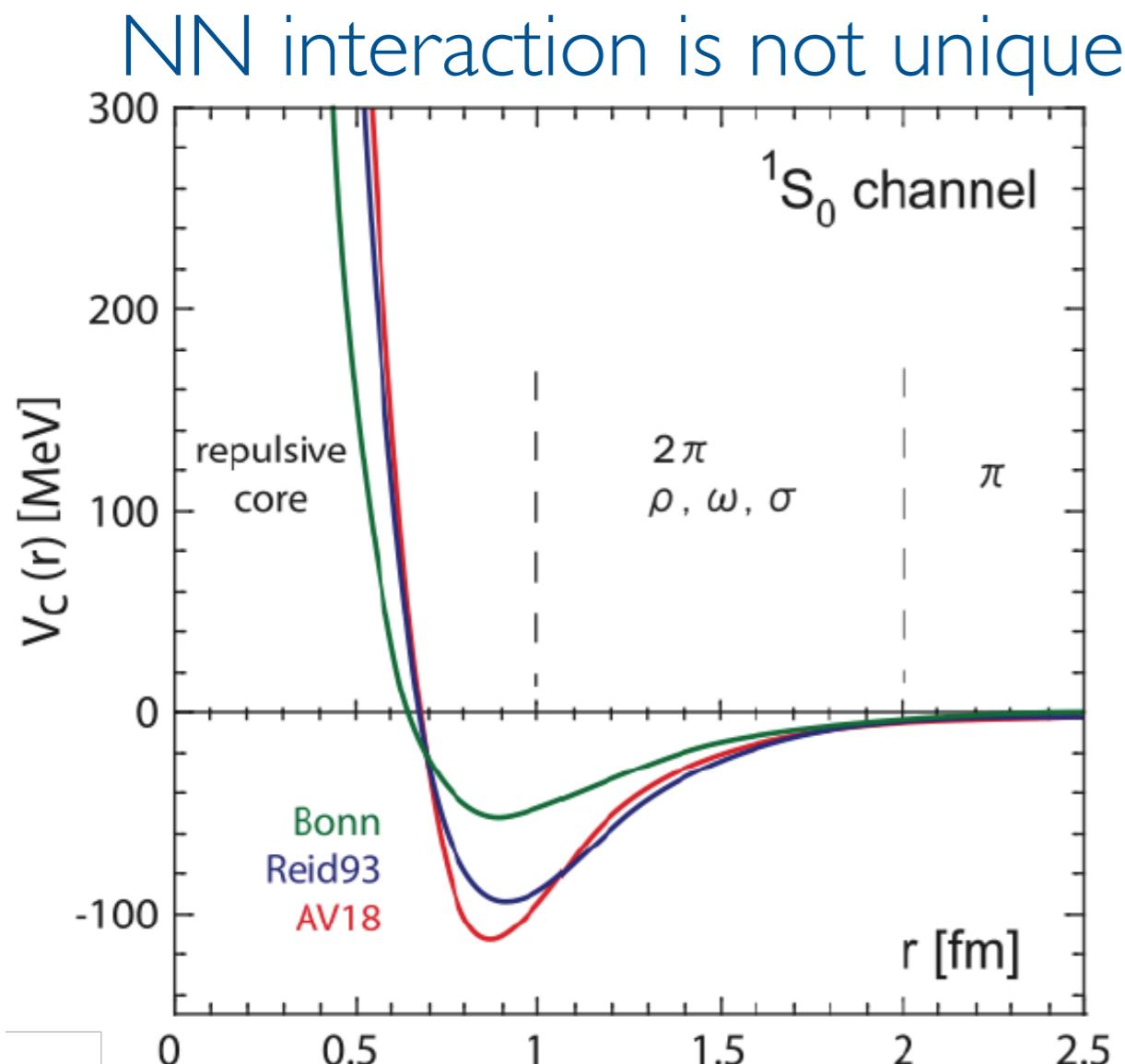


Complications

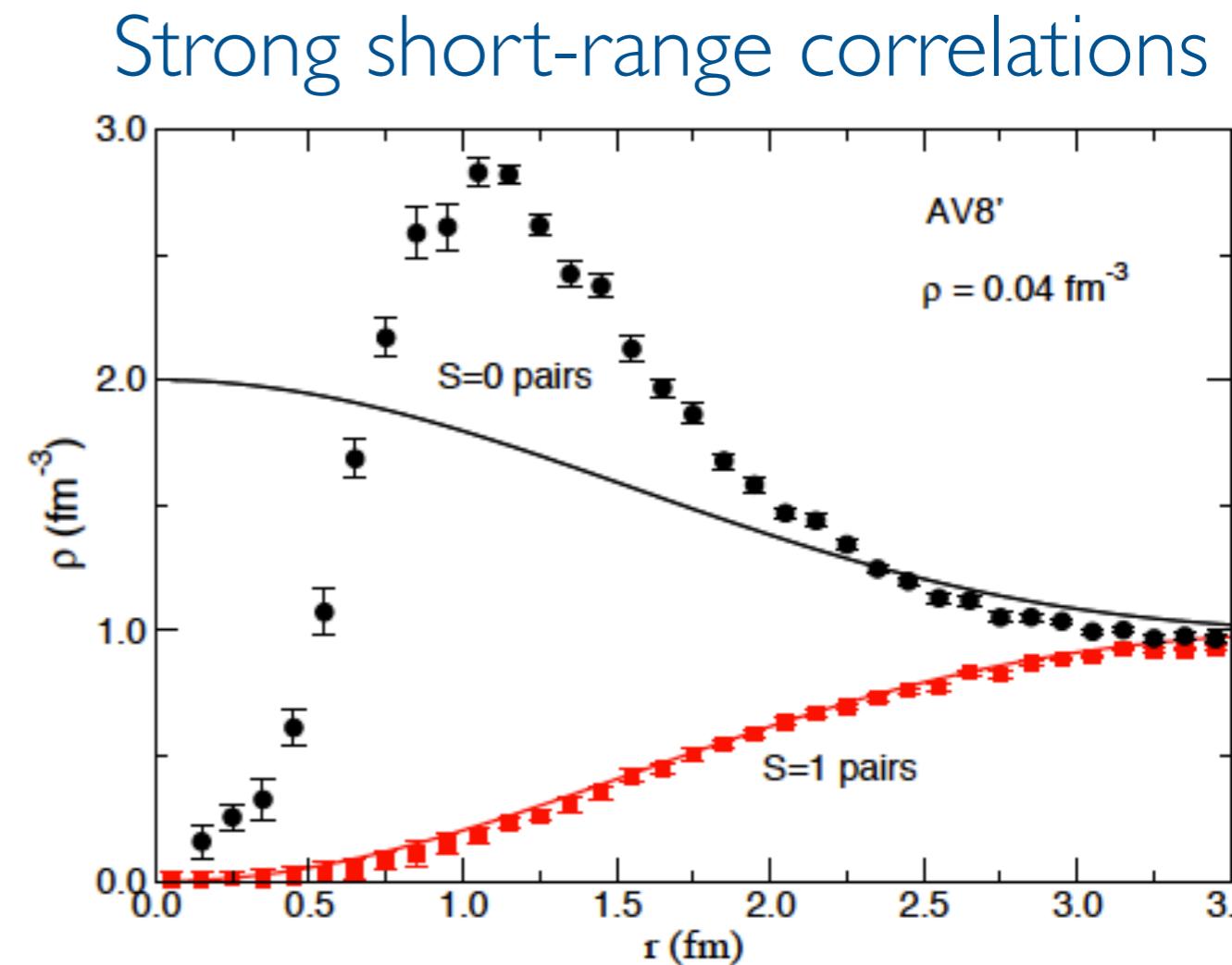


- Non-uniqueness of nucleon forces ✗

Complications



S.Aoki, et al. *Comput. Sci. Dis.* **I** 015009 (2008)

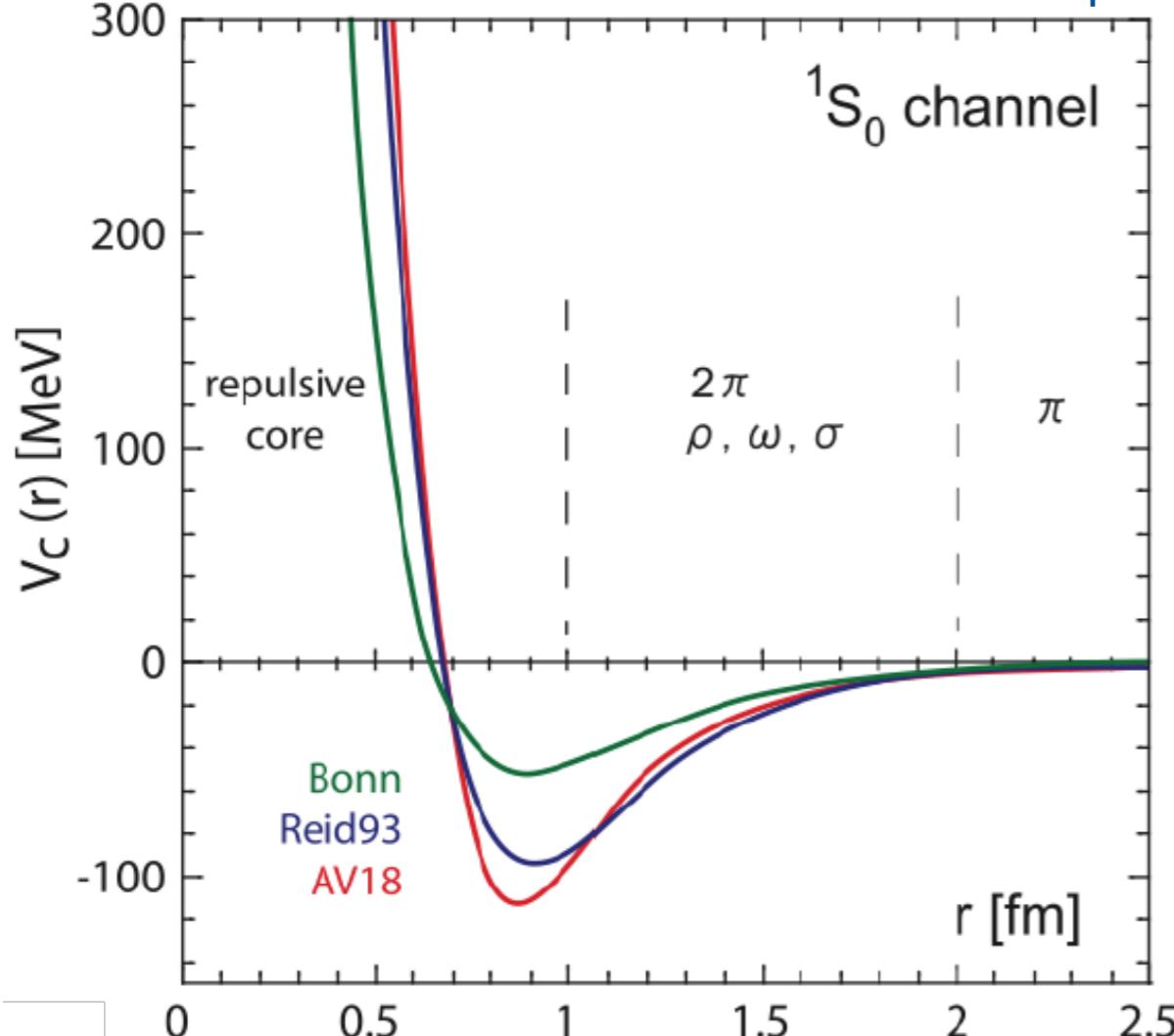


Carlson et al., *Phys. Rev. C* **68** 025802 (2003)

- Non-uniqueness of nucleon forces ✗
- Short-range core needs many-body treatment ✗

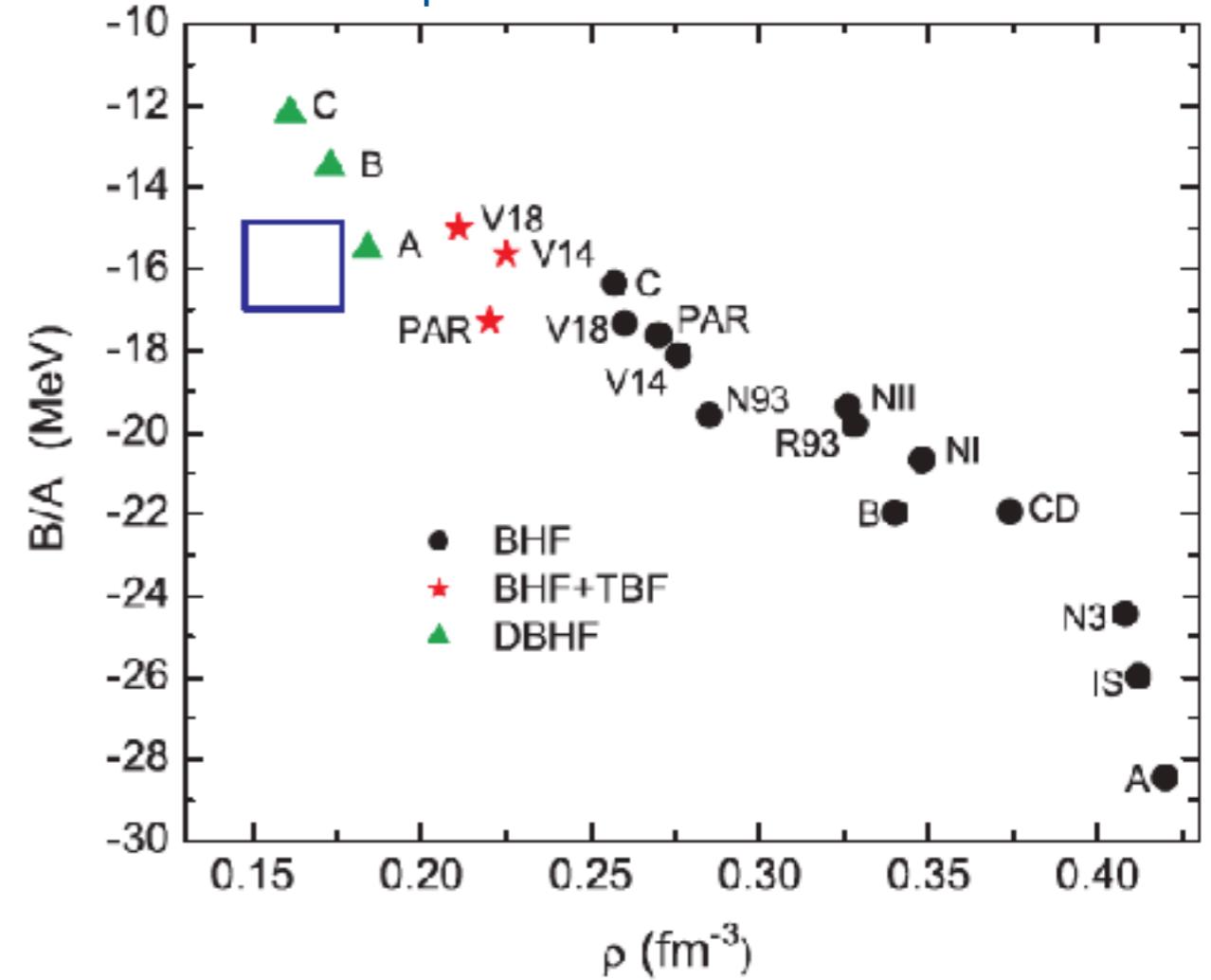
Complications

NN interaction is not unique



S.Aoki, et al. *Comput. Sci. Dis.* **I** 015009 (2008)

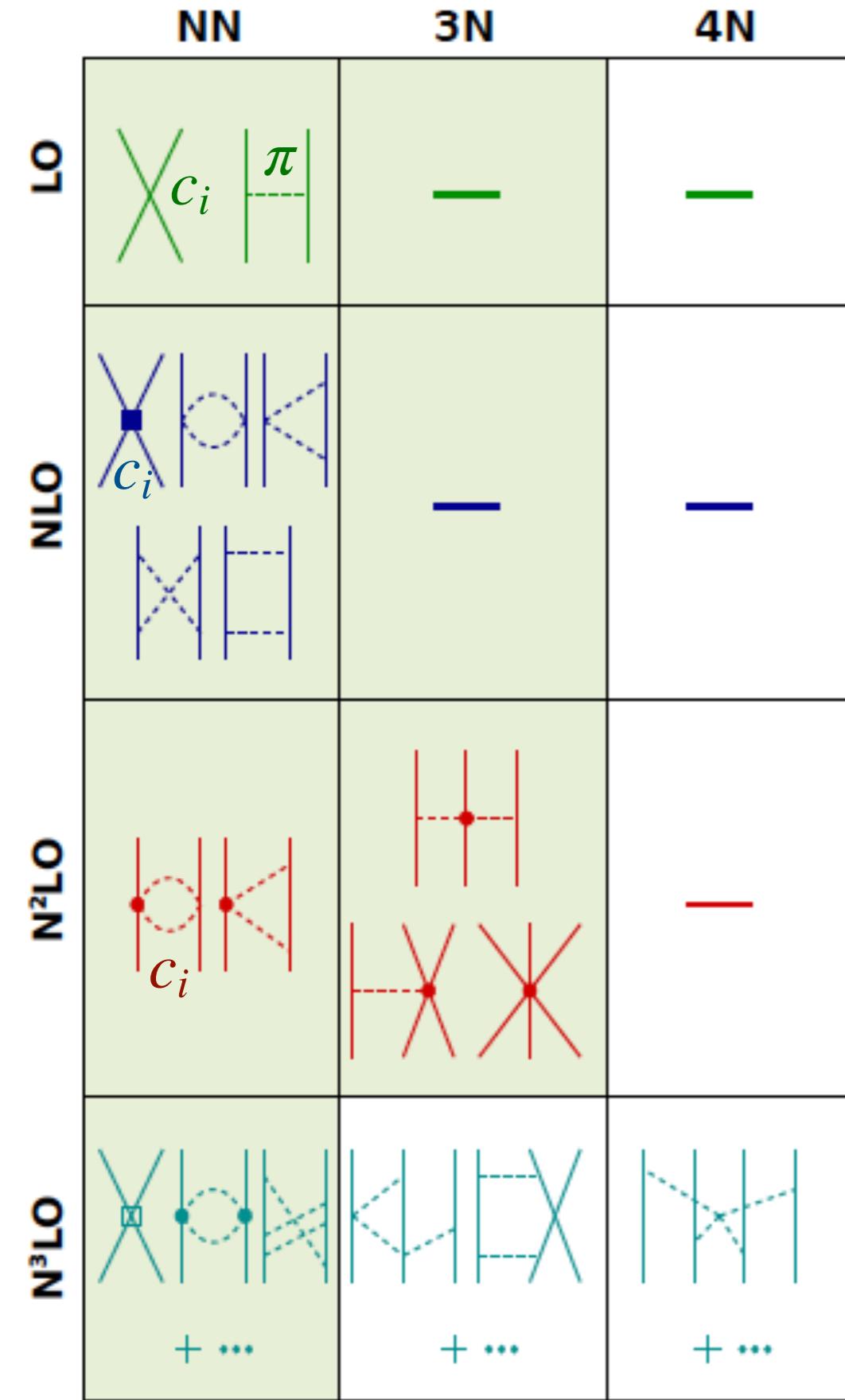
Saturation point of nuclear matter



Li, Lombardo, Schulze et al. *PRC* **74** 047304 (2006)

- Non-uniqueness of nucleon forces ✗
- Short-range core needs many-body treatment ✗
- Three-body forces needed for saturation ✗

NN forces from EFTs of QCD

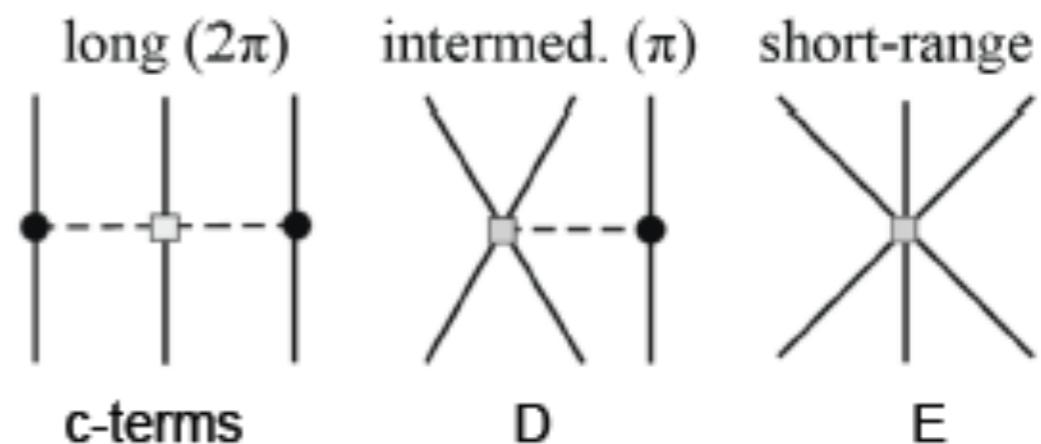


$$\mathcal{O}\left(\frac{Q}{\Lambda}\right)$$

$$\Lambda \sim 1 \text{ GeV}$$

Chiral perturbation theory

- π and N as dof
- Systematic expansion
- 2N at N^3LO - LECs from πN , NN
- 3N at N^2LO - 2 more LECs
- (Often further renormalized)



Weinberg, Phys. Lett. B **251** 288 (1990), NPB **363** 3 (1991)

Entem & Machleidt, PRC **68**, 041001(R) (2003)

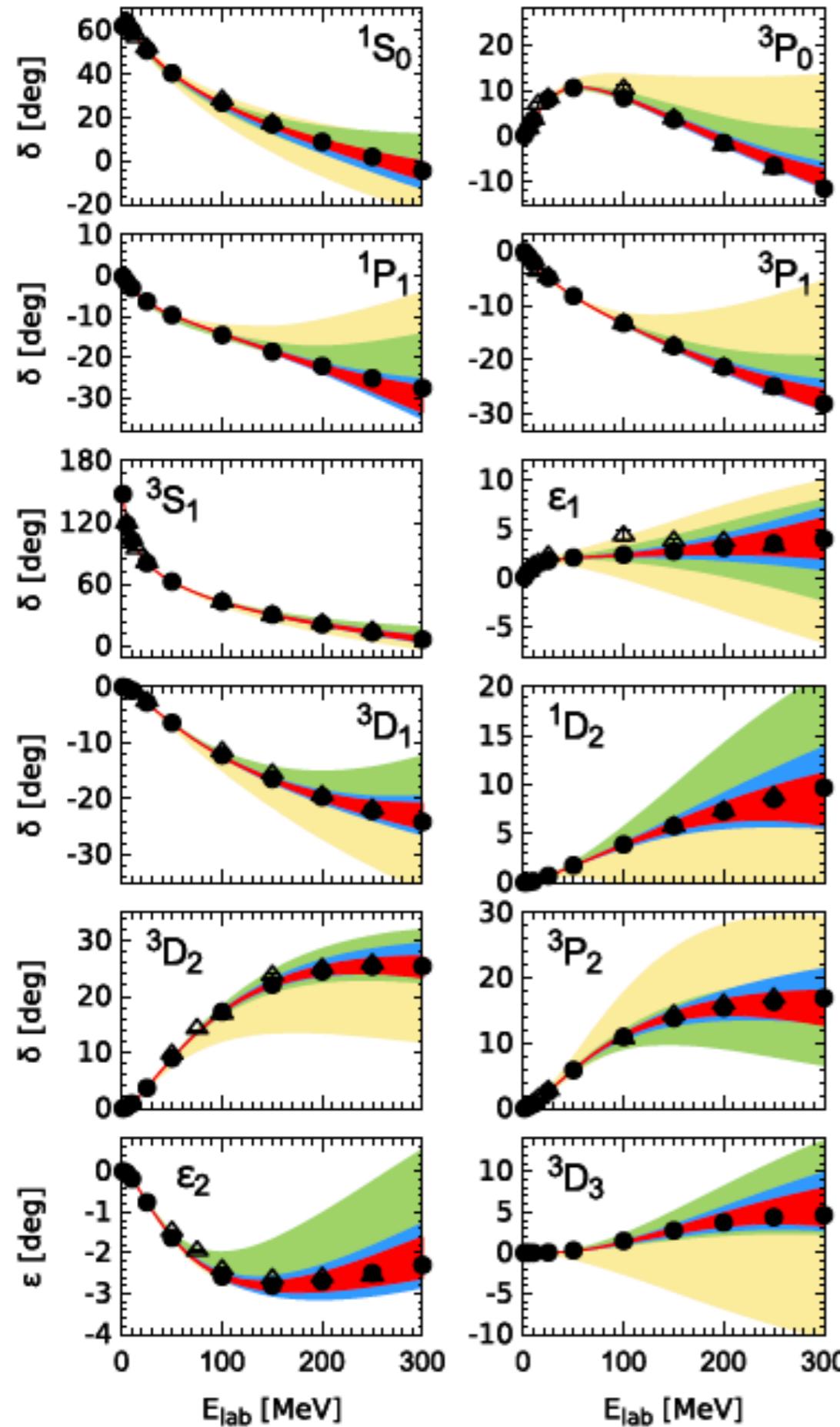
Tews, Schwenk et al., PRL **110**, 032504 (2013)

Epelbaum, Frebs & Meissner, PRL **115**, 122301 (2015)

NN forces from EFTs of QCD



UNIVERSITY OF
SURREY

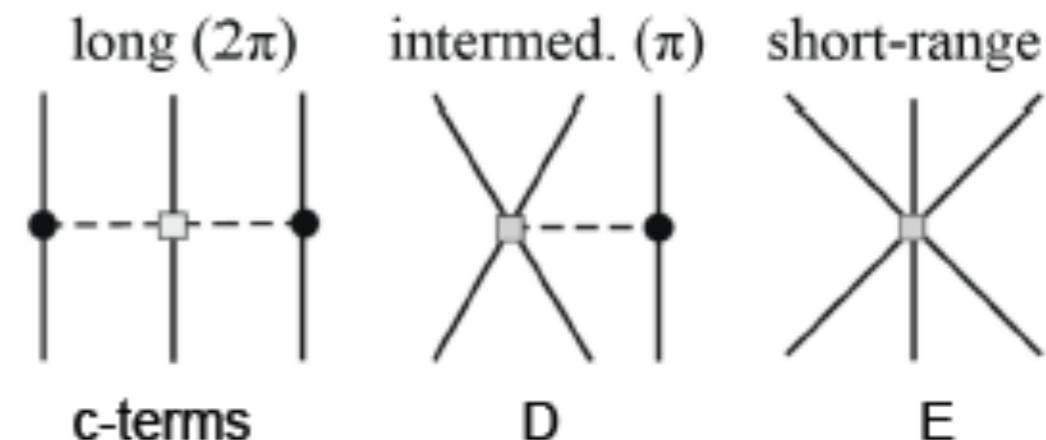


Chiral perturbation theory

- π and N as dof
- Systematic expansion
- 2N at N³LO - LECs from πN , NN
- 3N at N²LO - 2 more LECs
- (Often further renormalized)

$$\mathcal{O}\left(\frac{Q}{\Lambda}\right)$$

$$\Lambda \sim 1 \text{ GeV}$$



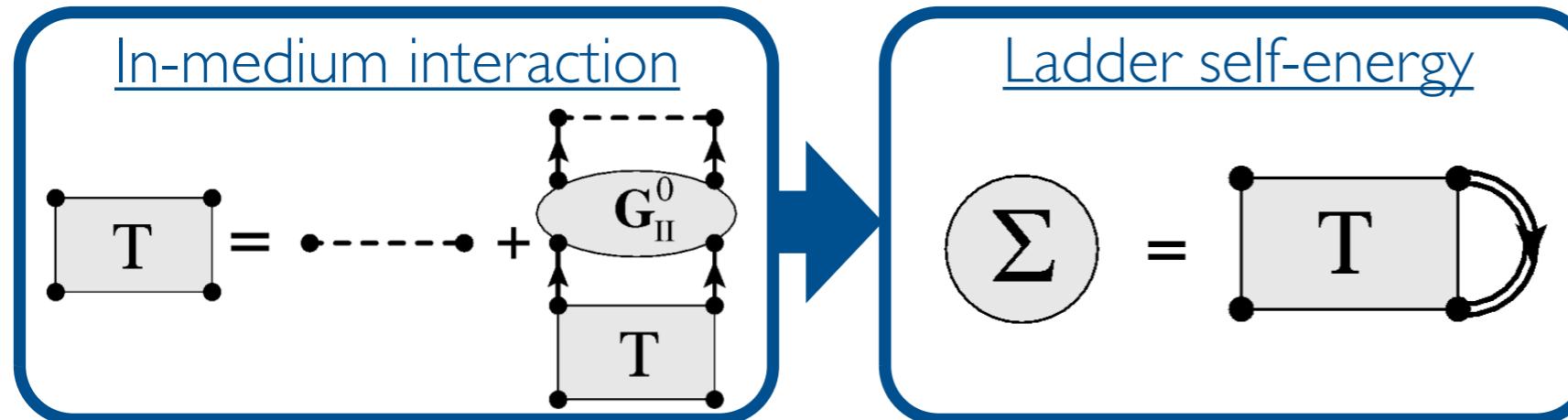
Weinberg, Phys. Lett. B **251** 288 (1990), NPB **363** 3 (1991)

Entem & Machleidt, PRC **68**, 041001(R) (2003)

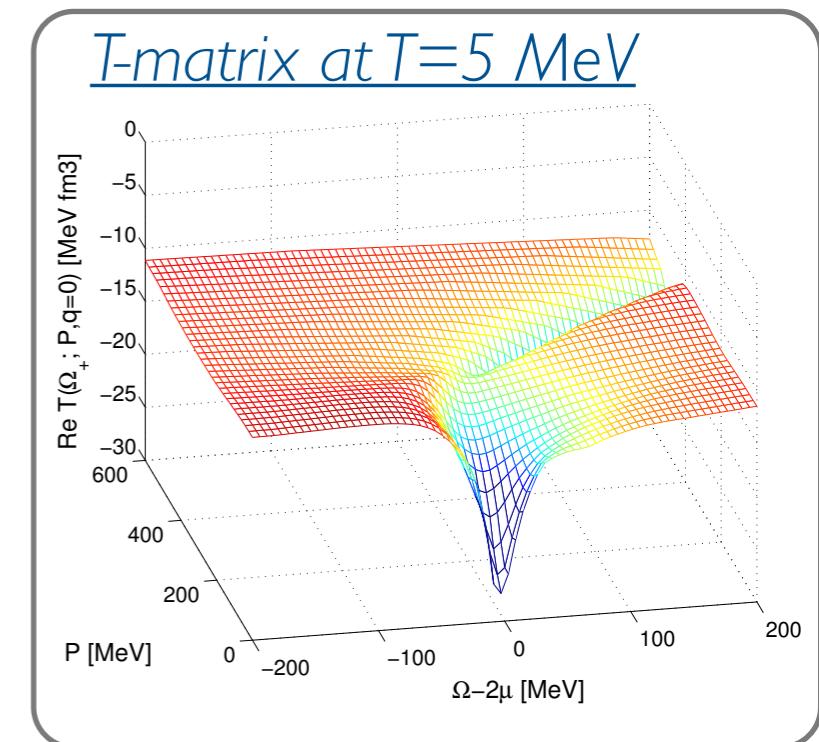
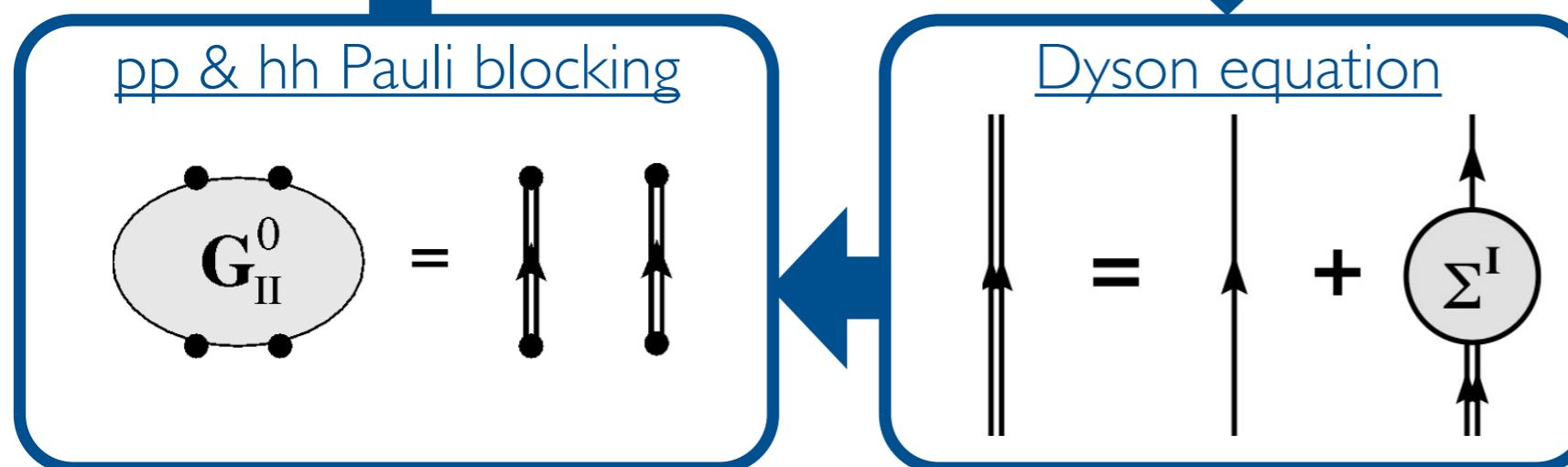
Tews, Schwenk et al., PRL **110**, 032504 (2013)

Epelbaum, Frebs & Meissner, PRL **115**, 122301 (2015)

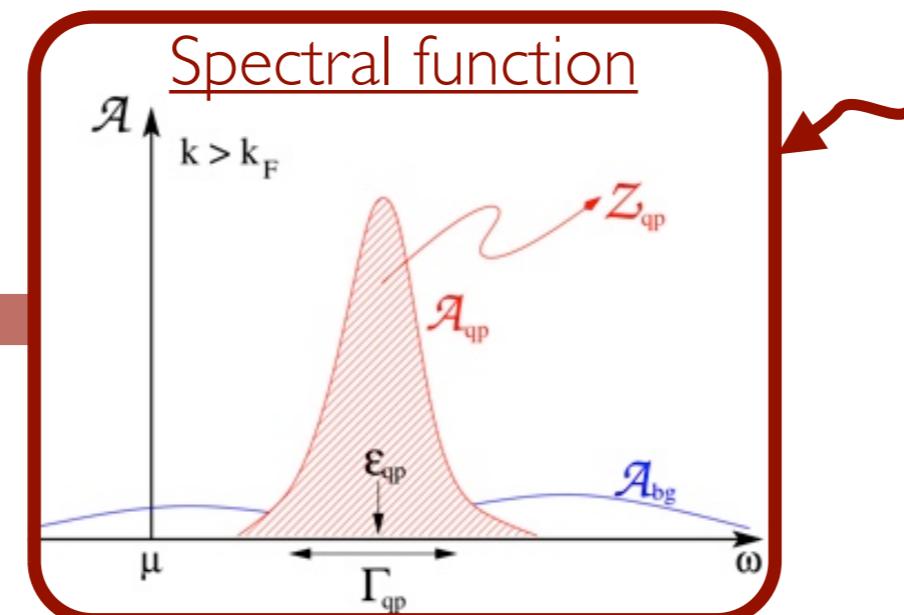
SCGF Ladder approximation



- Self-consistent resummation
- Energy and momentum integral
- @Finite T (Matsubara)



One-body properties
Momentum distribution
Thermodynamics & EoS
Transport



- Ramos, Polls & Dickhoff, NPA **503** 1 (1989)
 Alm et al., PRC **53** 2181 (1996)
 Dewulf et al., PRL **90** 152501 (2003)
 Frick & Muther, PRC **68** 034310 (2003)
 Rios, PhD Thesis, U. Barcelona (2007)
 Soma & Bozek, PRC **78** 054003 (2008)
 Rios & Soma PRL **108** 012501 (2012)

Ladder approximation with 3BF

Two-body interaction



In-medium T-matrix

$$\boxed{T} = \text{---} + \text{---} \quad \text{---} \quad \boxed{T}$$

Self-energy

$$\boxed{\Sigma} = \text{---} \quad \text{---} \quad \boxed{T}$$

Dyson equation

$$\boxed{\Gamma} = \boxed{\Gamma} + \boxed{\Sigma}$$

Effective interactions

Effective one-body force

$$\text{---} = \text{---} + \frac{1}{2} \text{---}$$

Effective two-body force

$$\text{---} = \text{---} + \text{---}$$

In-medium T-matrix

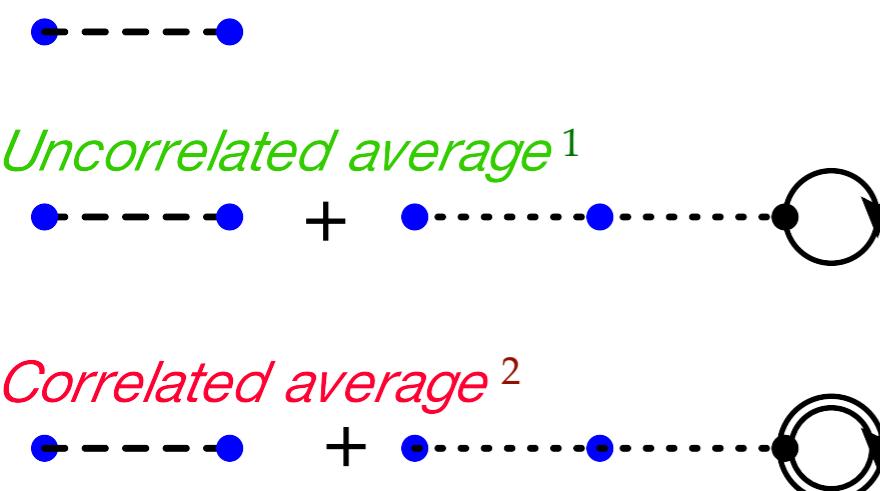
$$\boxed{T} = \text{---} + \text{---} \quad \boxed{T}$$

Self-energy

$$\boxed{\Sigma} = \text{---} + \text{---} \\ + \frac{1}{12} \text{---} + \dots$$

Density-dependent interaction

Two-body N3LO

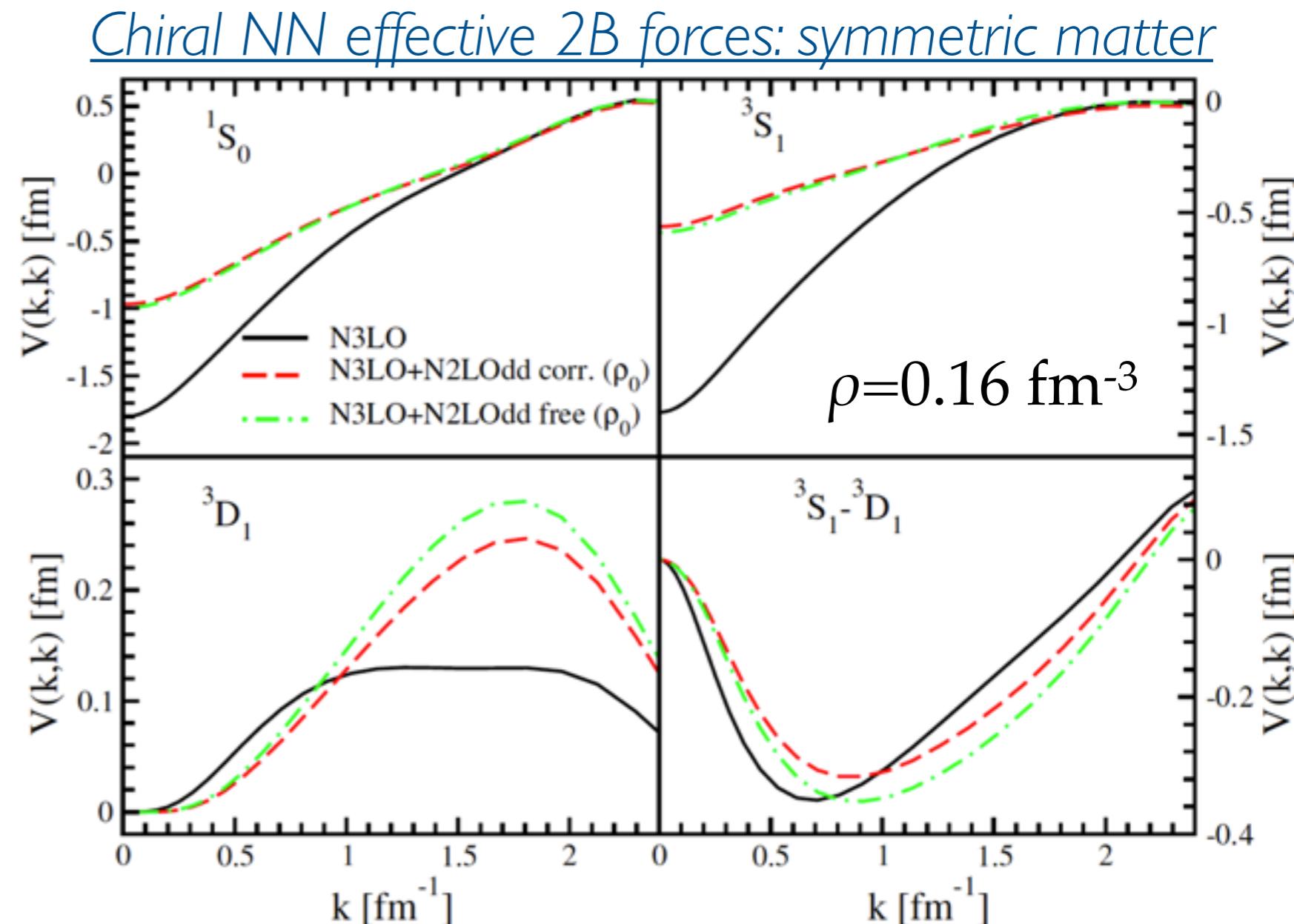


LECs

$$c_D = -1.11$$

$$c_E = -0.66$$

$$k \neq k' \Rightarrow \frac{1}{2}(k + k')$$



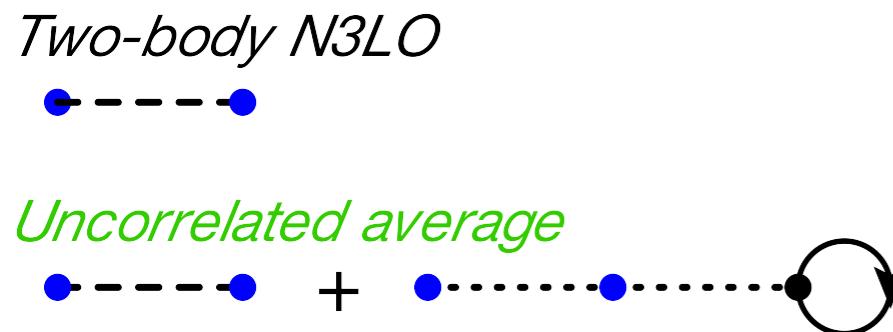
- 3NF bring repulsion: correlated & uncorrelated averages are similar
- Correlated average brings small corrections to 1/2 of terms
- Diagonal $k=k'$ matrix elements computed
- Off-diagonal extrapolated & regulated

¹Holt et al. Phys. Rev. C **81** 024002 (2010)

² Carbone, Polls & Rios, PRC **90**, 054322 (2014); A. Carbone, PhD thesis 6

Symmetric matter

Theoretical uncertainties: Chiral expansion

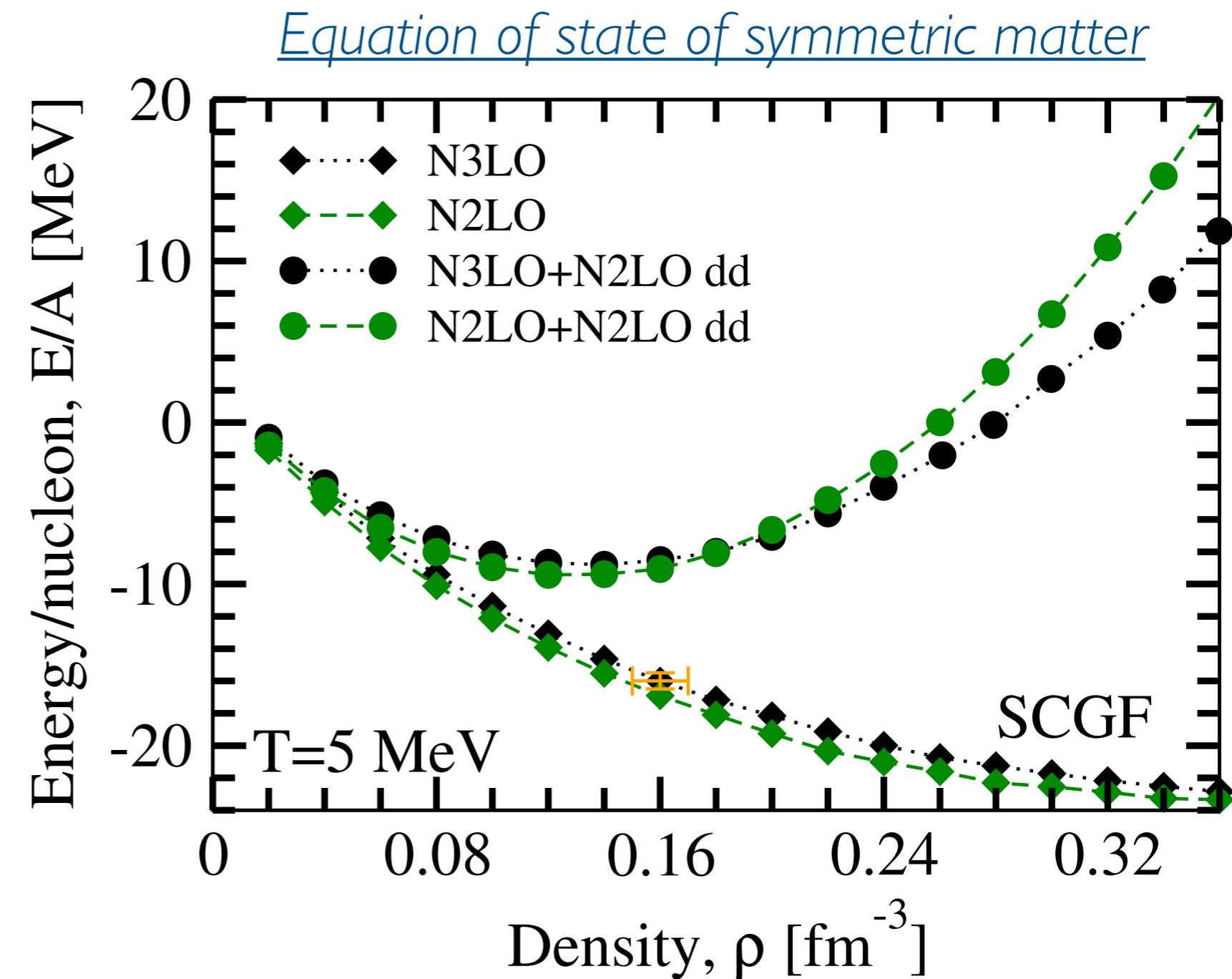


LECs

$$c_D = -1.11$$

$$c_E = -0.66$$

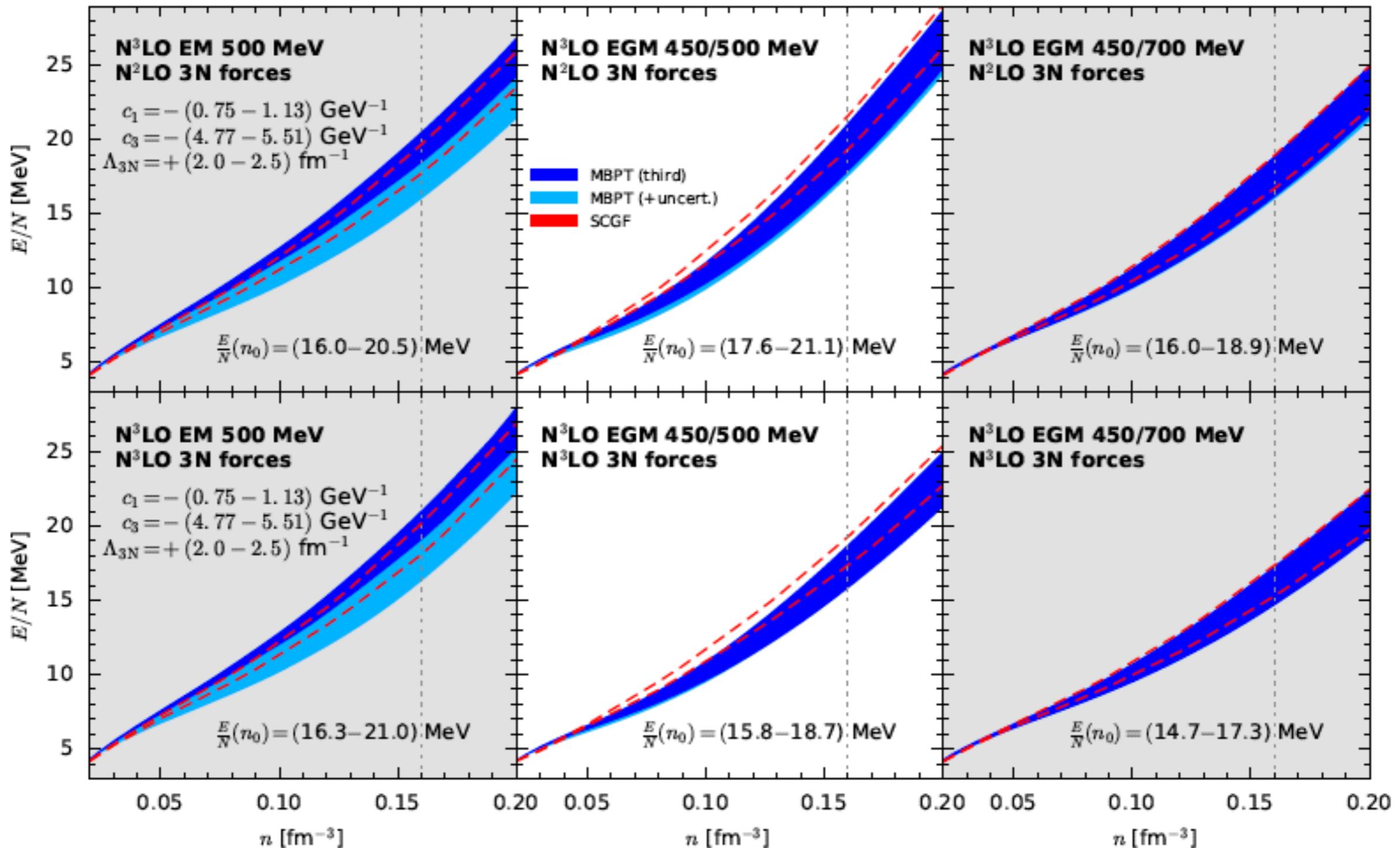
$$K_0 \sim 60 \text{ MeV}$$



Carbone, Polls & Rios, PRC **88** 044302 (2014)

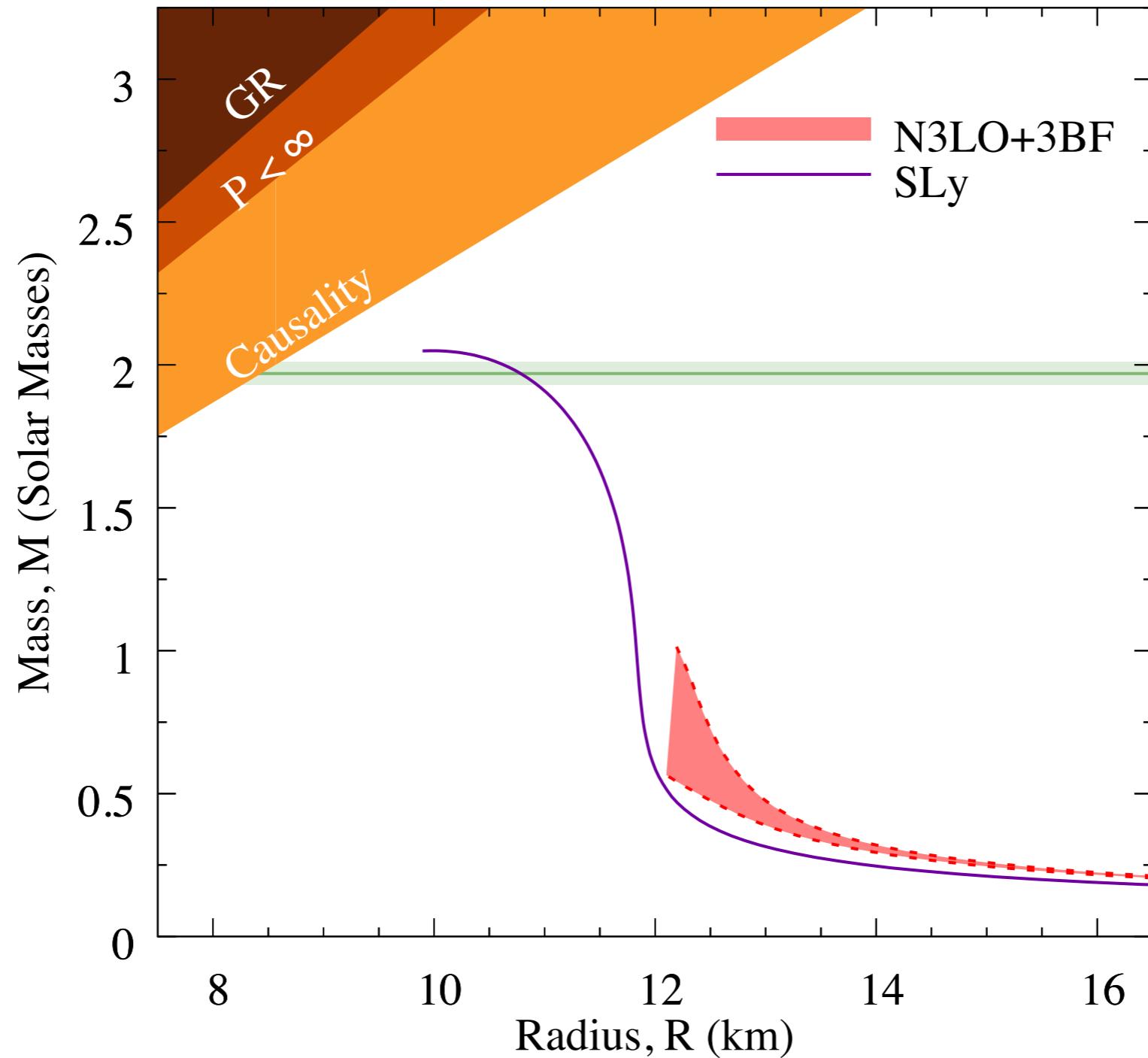
- 3NF result is still underbound
- Small difference in infinite matter for N3LO & N2LO...
- In contrast to finite nuclei!

Neutron matter

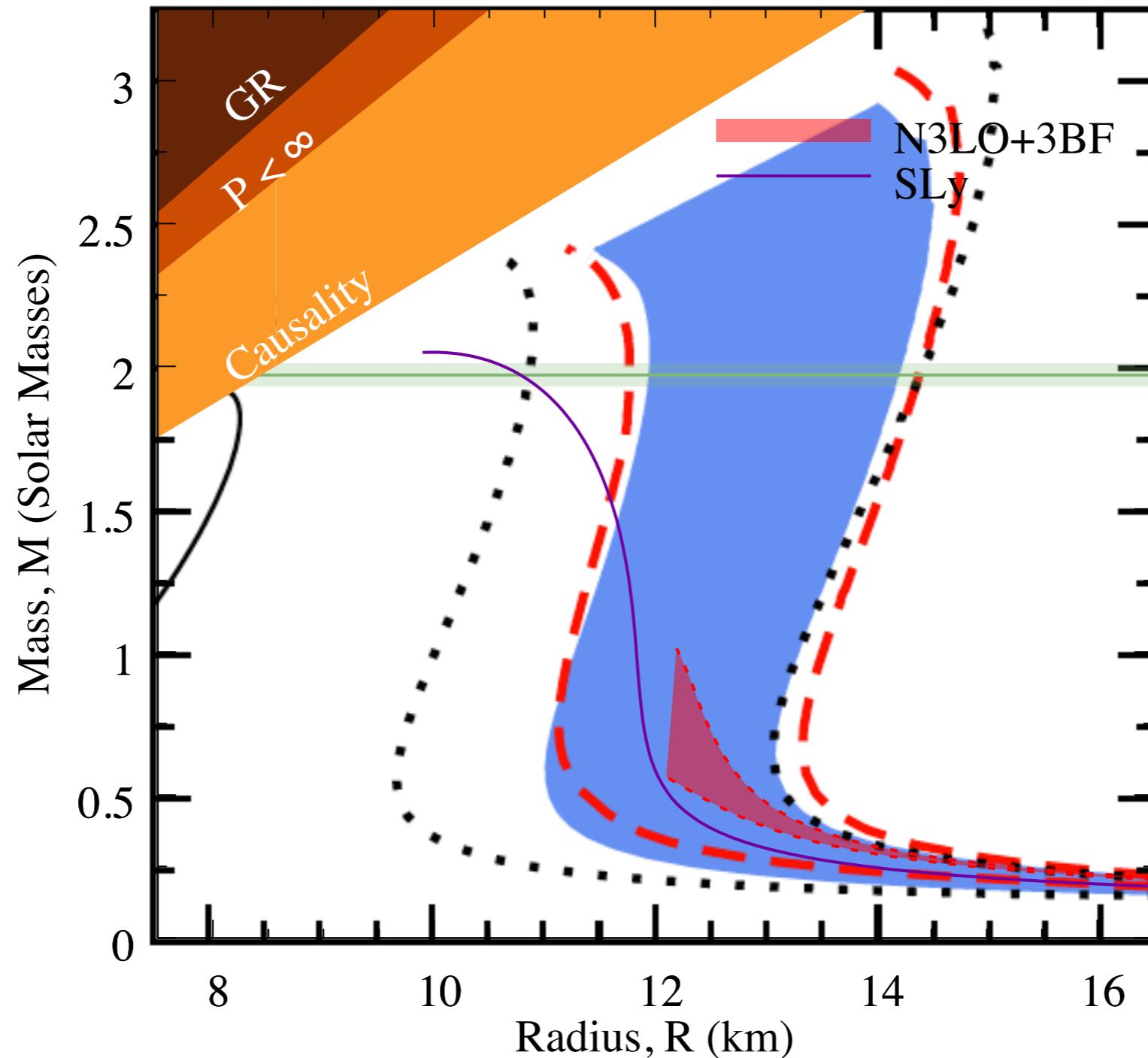


Drischler, Carbone, Hebeler, Schwenk PRC **94** 054307 (2016)

- Uncertainty band from unknown ChPT LECs + cutoff + MBPT
- Finite temperature available too



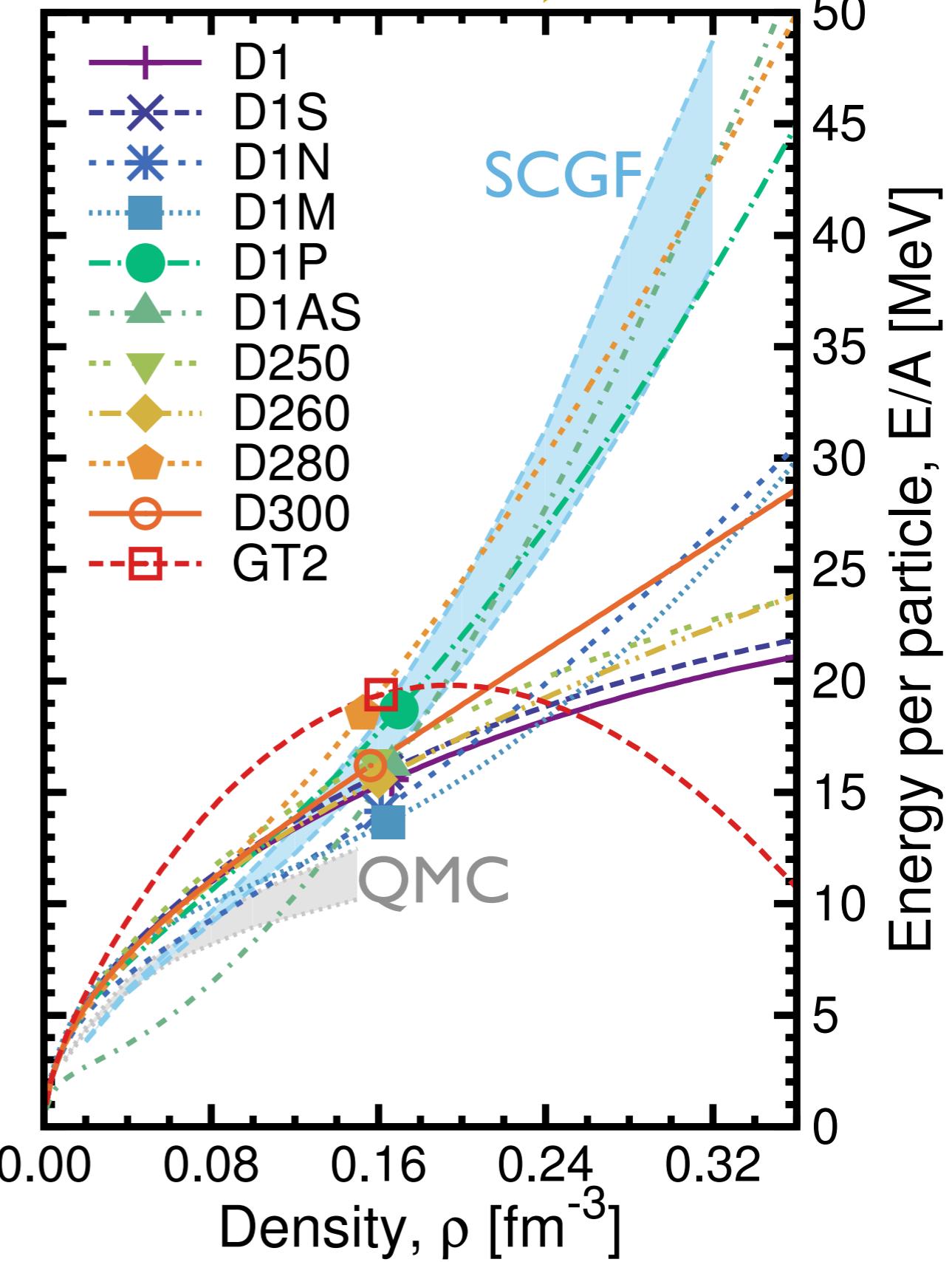
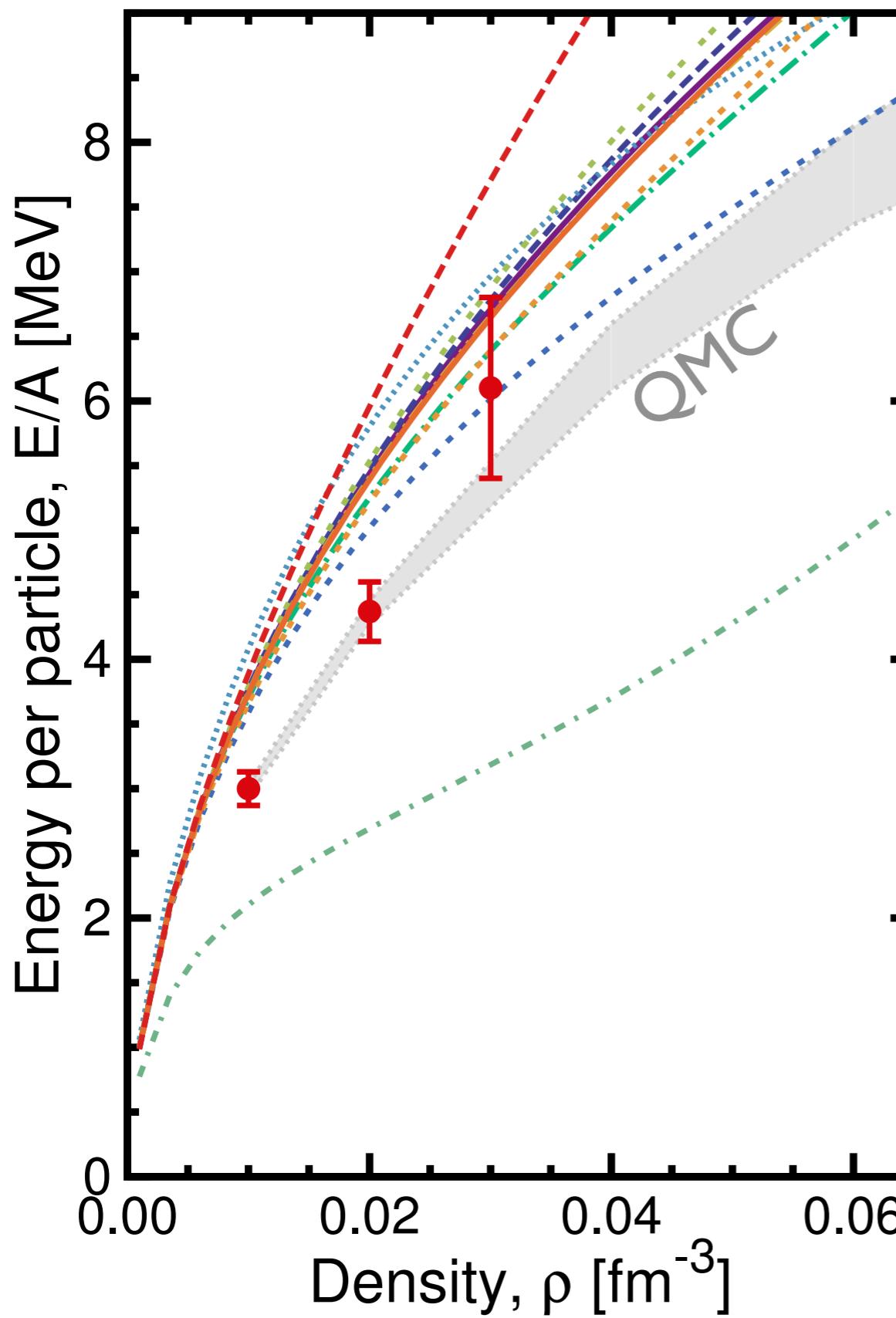
- Mass-Radius relation from SCGF calculations
- Cut-off variation (N3LO) and/or SRG evolution



Hebeler, Lattimer, Pethick, Schwenk ApJ 773 11 (2013)

- Mass-Radius relation from SCGF calculations
- Cut-off variation (N3LO) and/or SRG evolution

Neutron matter



Momentum distribution

Single-particle occupation

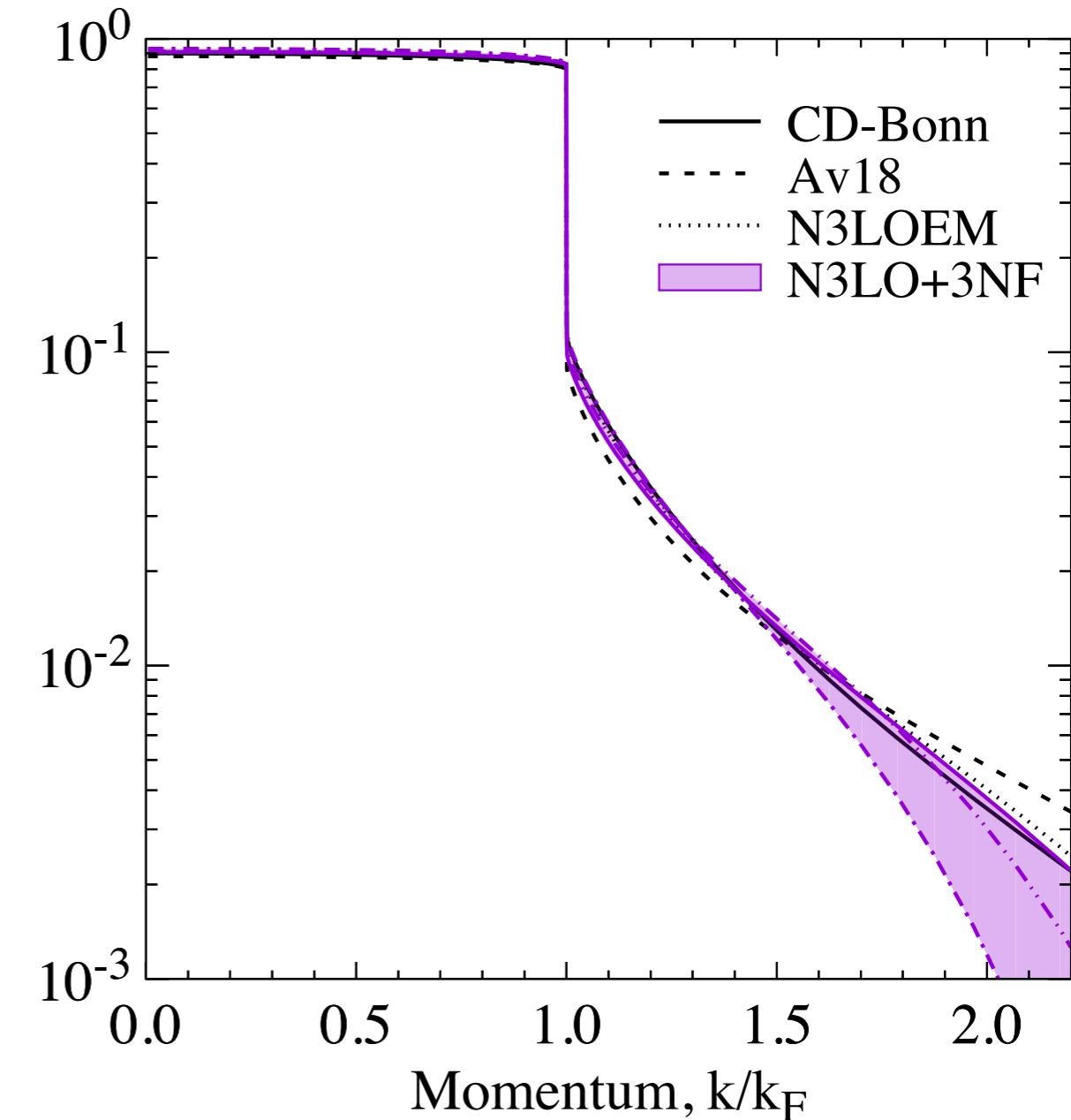
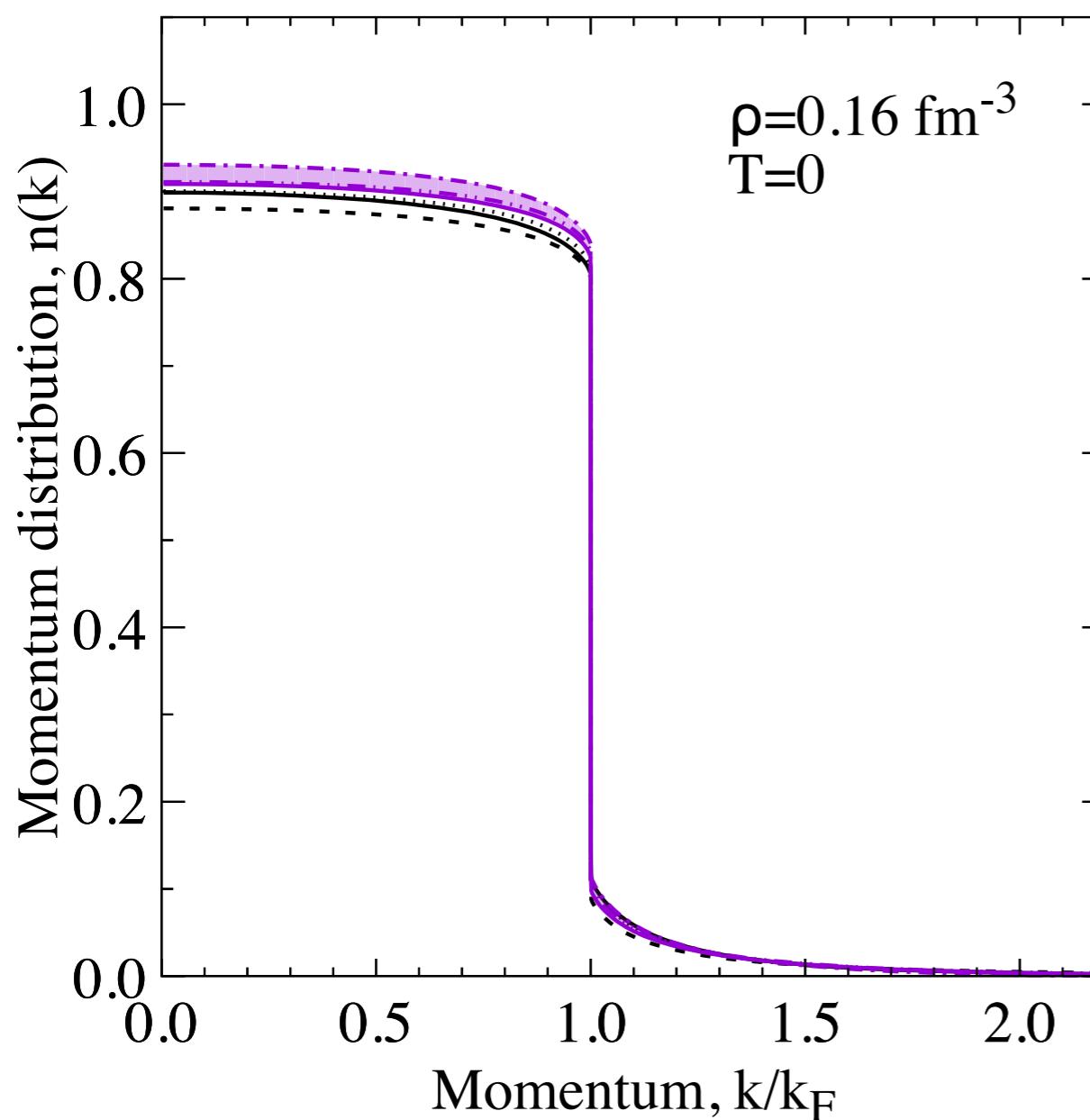
$$n(k) = \langle a_k^\dagger a_k \rangle$$



UNIVERSITY OF
SURREY

$$\nu \int \frac{d^3k}{(2\pi)^3} n(k) = \rho$$

Symmetric matter



- Dependence on NN interaction **understood**
- $\text{N3LO+3NF} = \text{N3LO 2NF} + \text{N2LO 3NF}$ @ $\Lambda = 414\text{-}500 \text{ MeV}$ (cutoff variation only)

Momentum distribution

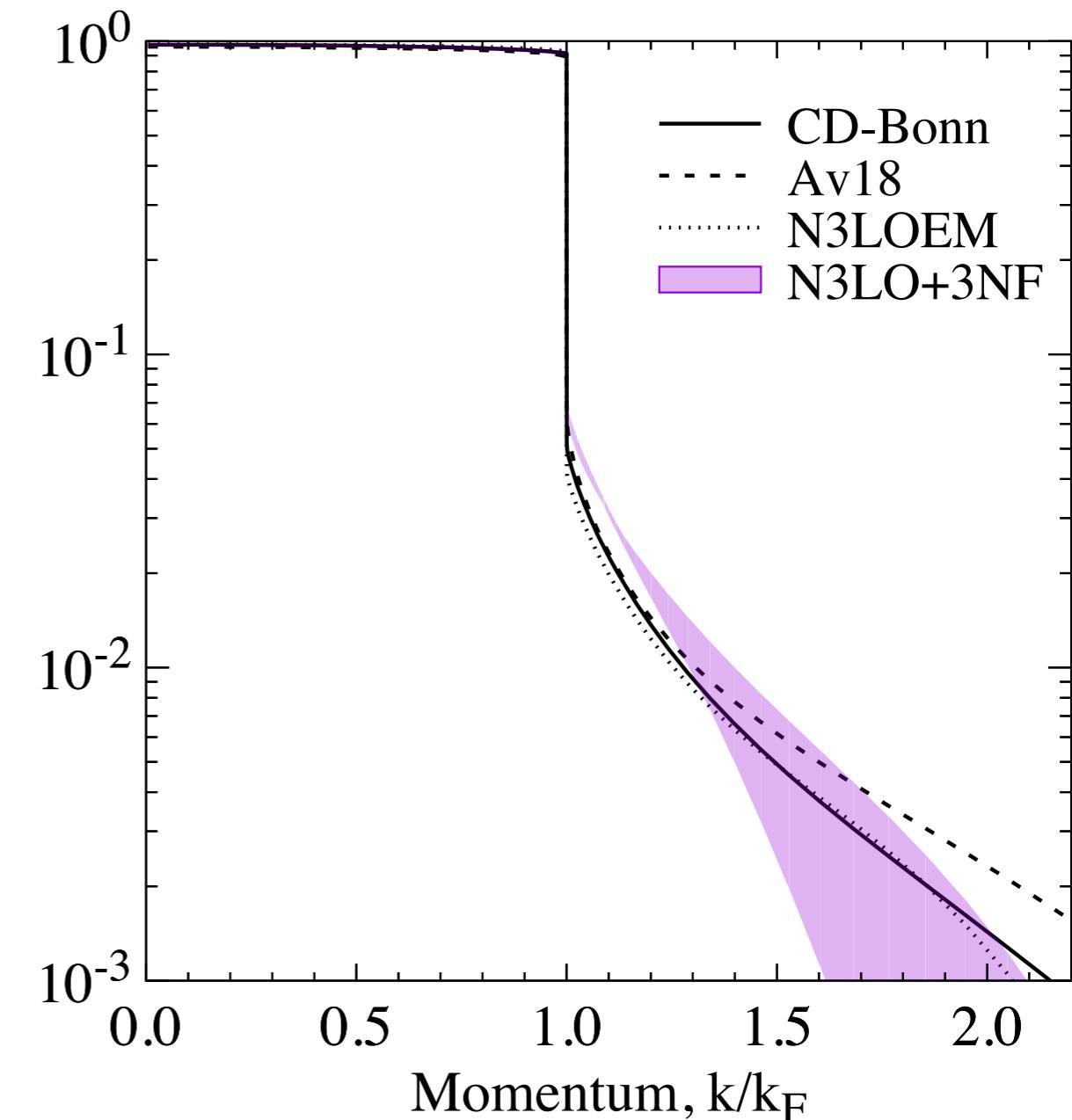
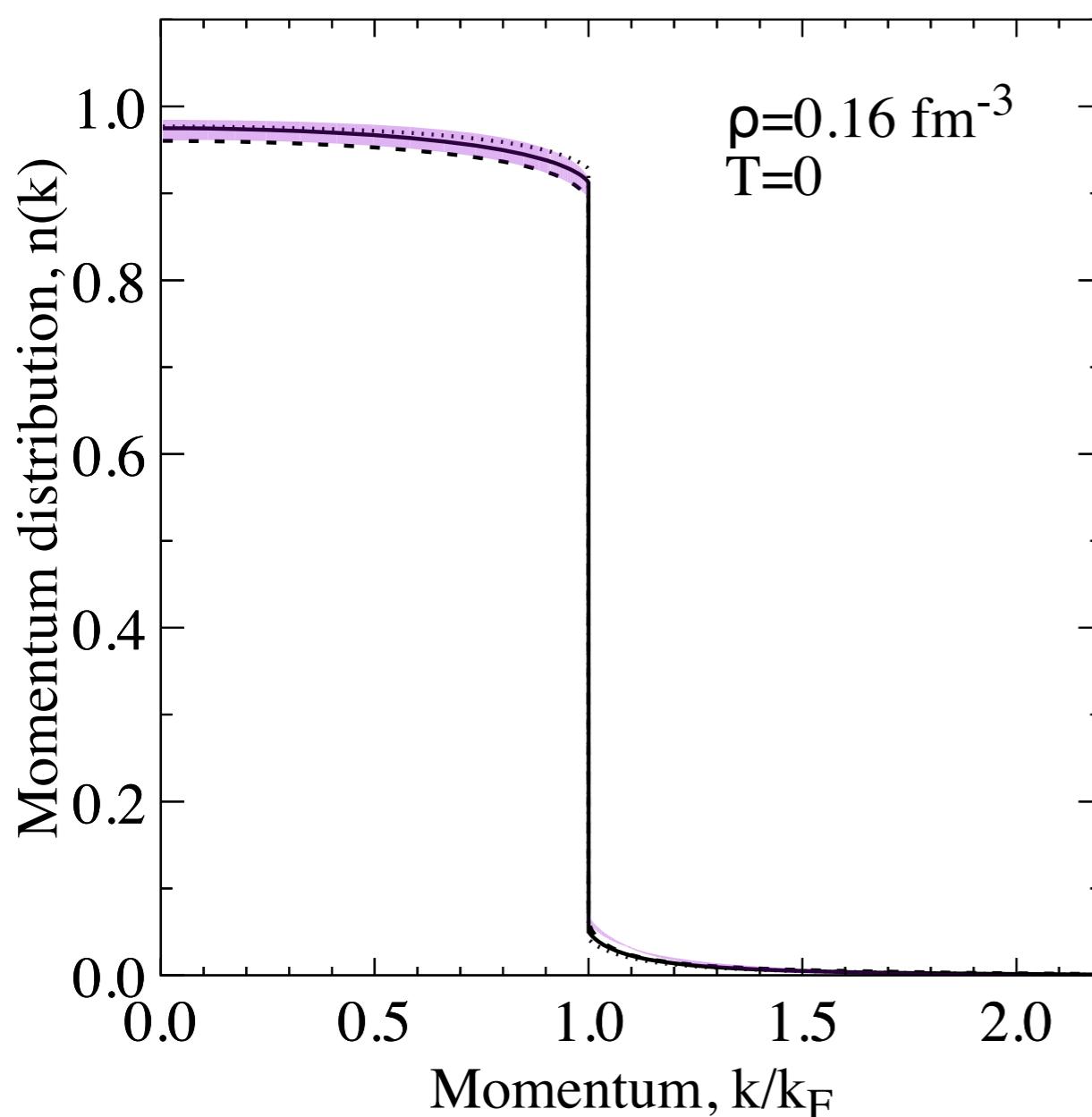
Single-particle occupation

$$n(k) = \langle a_k^\dagger a_k \rangle$$



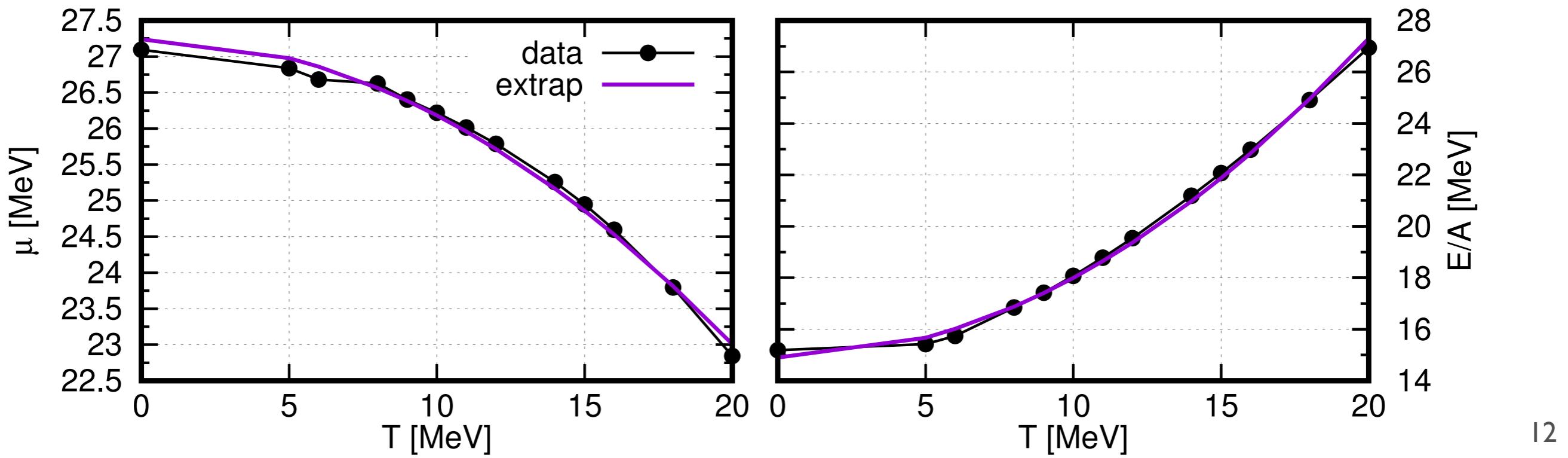
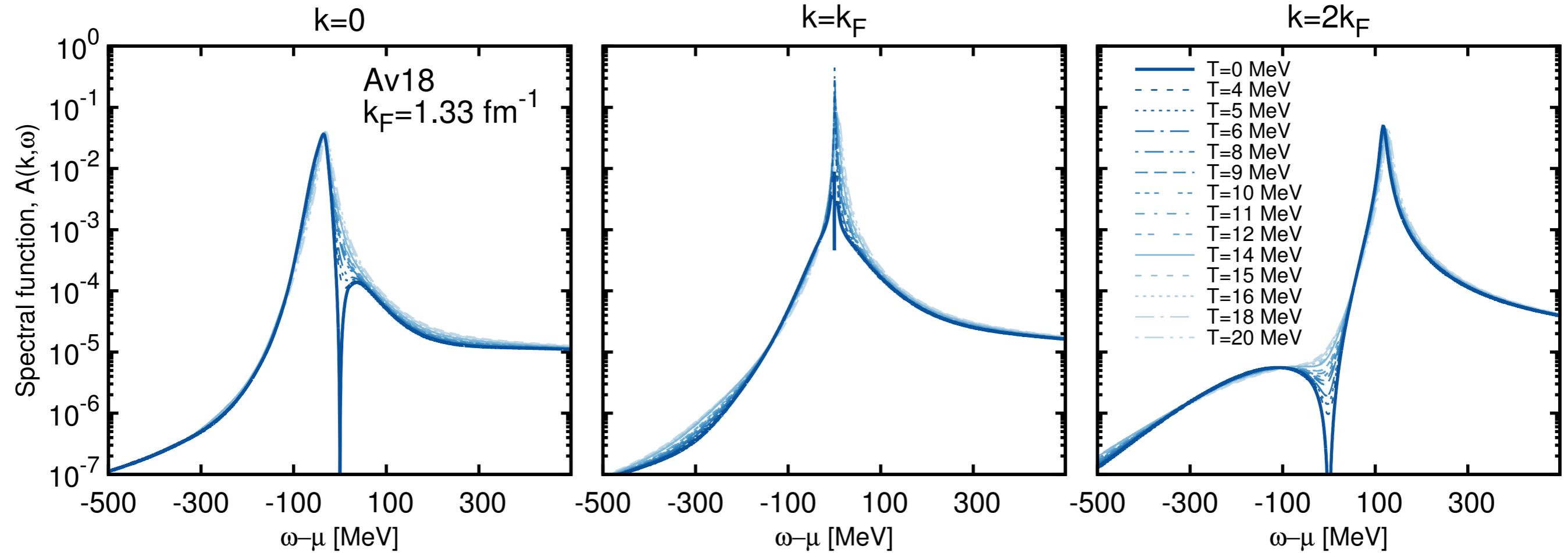
$$\nu \int \frac{d^3 k}{(2\pi)^3} n(k) = \rho$$

Neutron matter



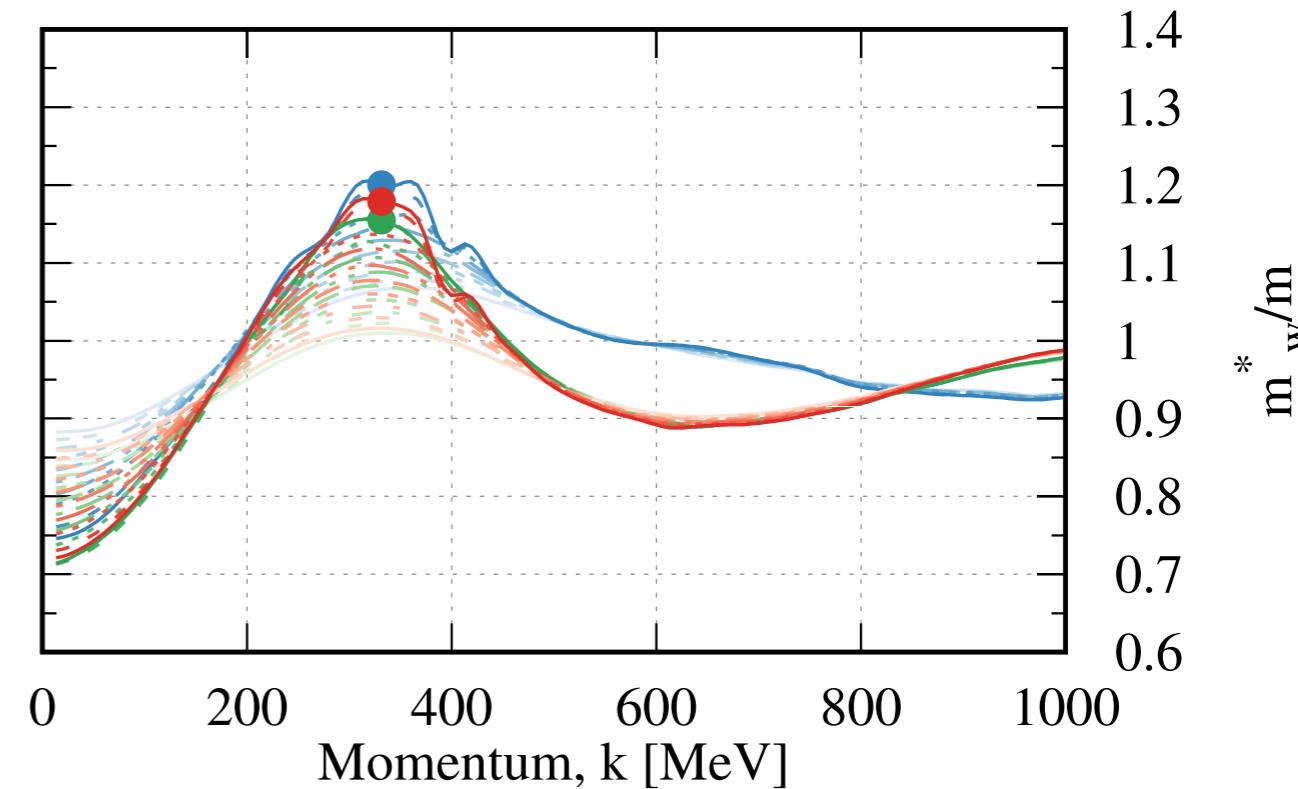
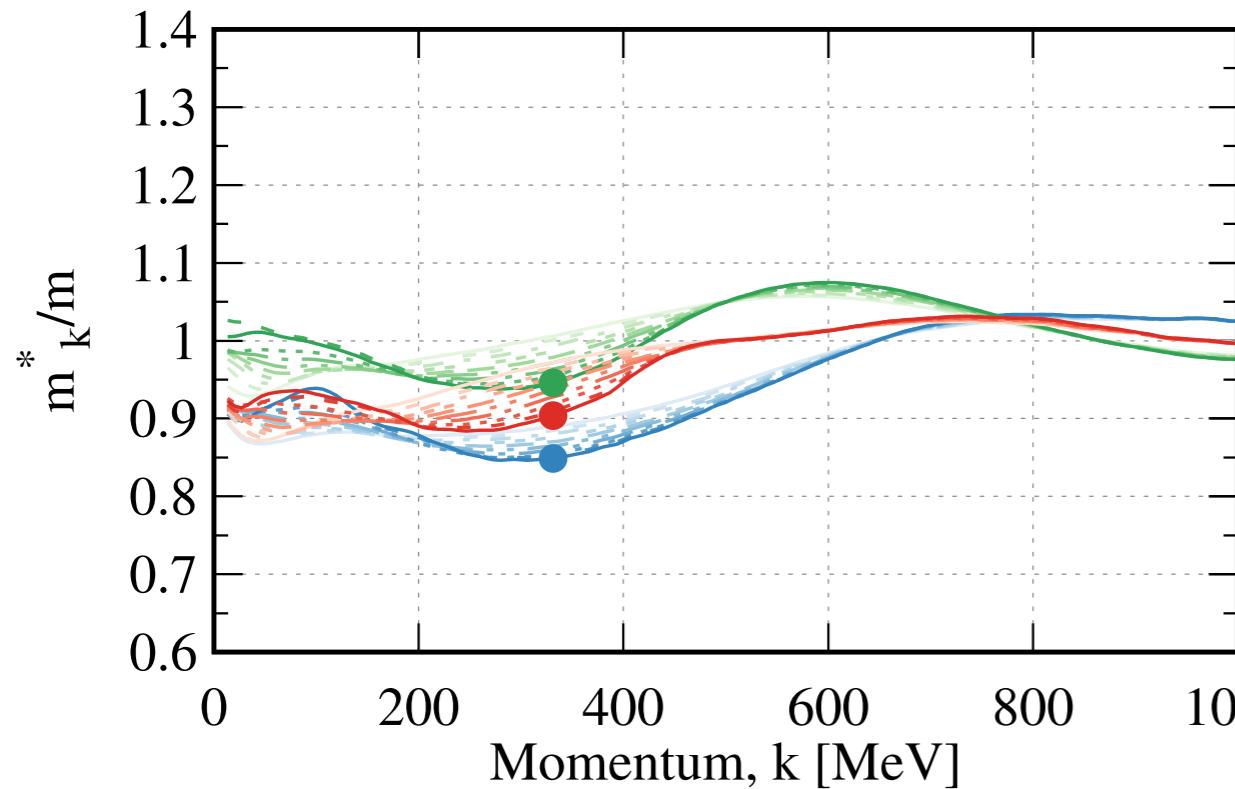
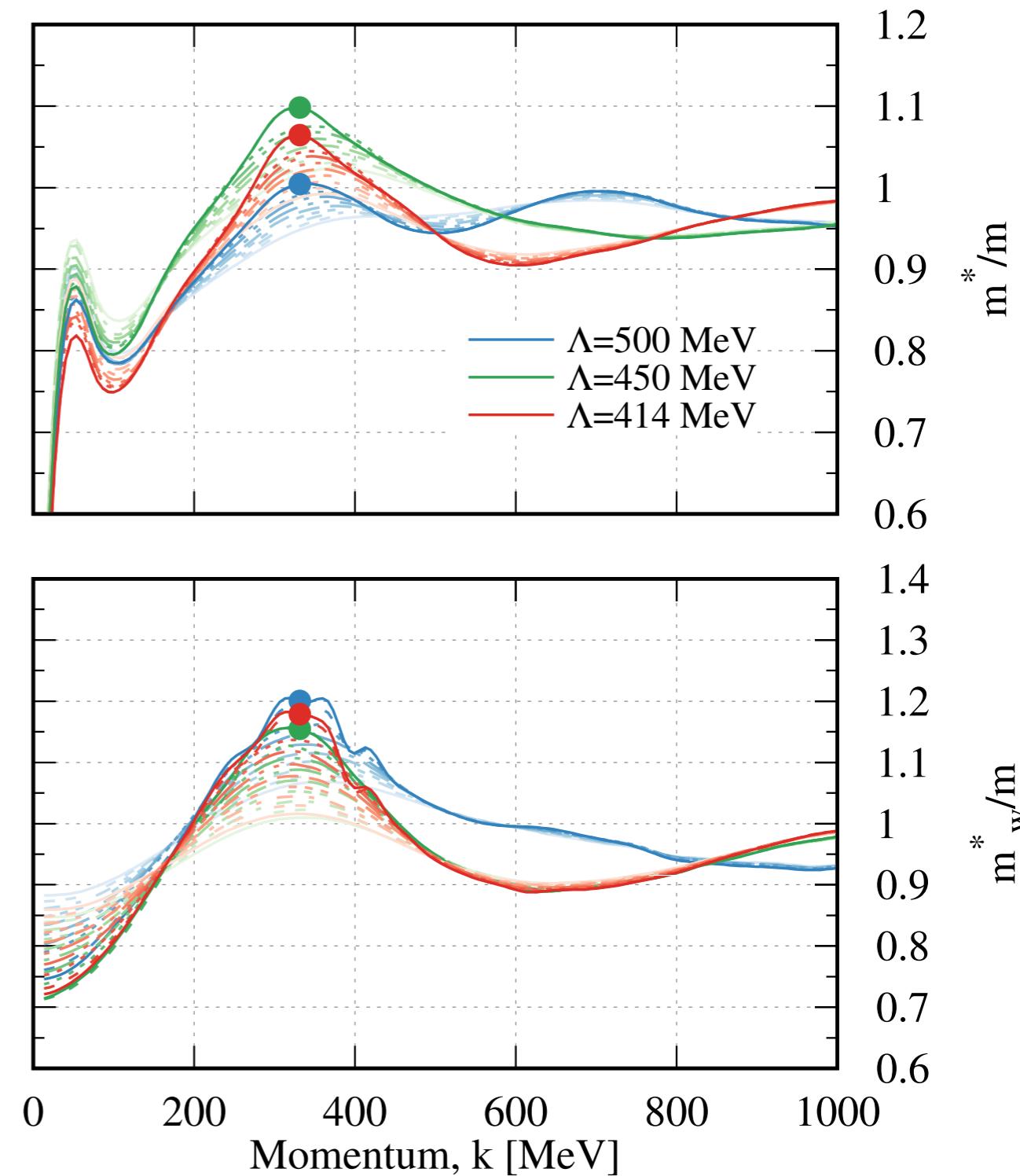
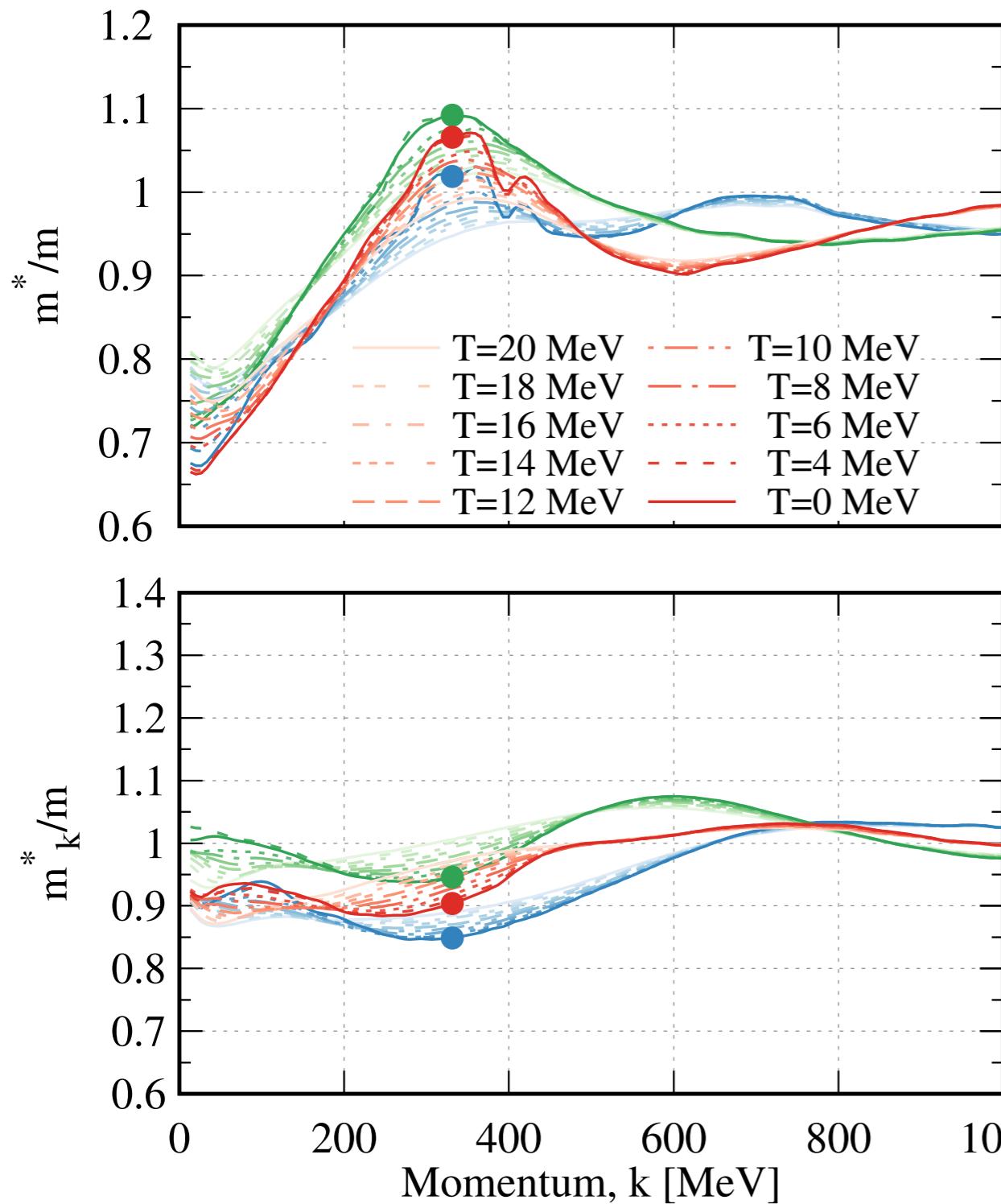
- Dependence on NN interaction **understood**
- N3LO+3NF = N3LO 2NF + N2LO 3NF @ $\Lambda=414-500$ MeV (cutoff variation only)

$T=0$ extrapolations



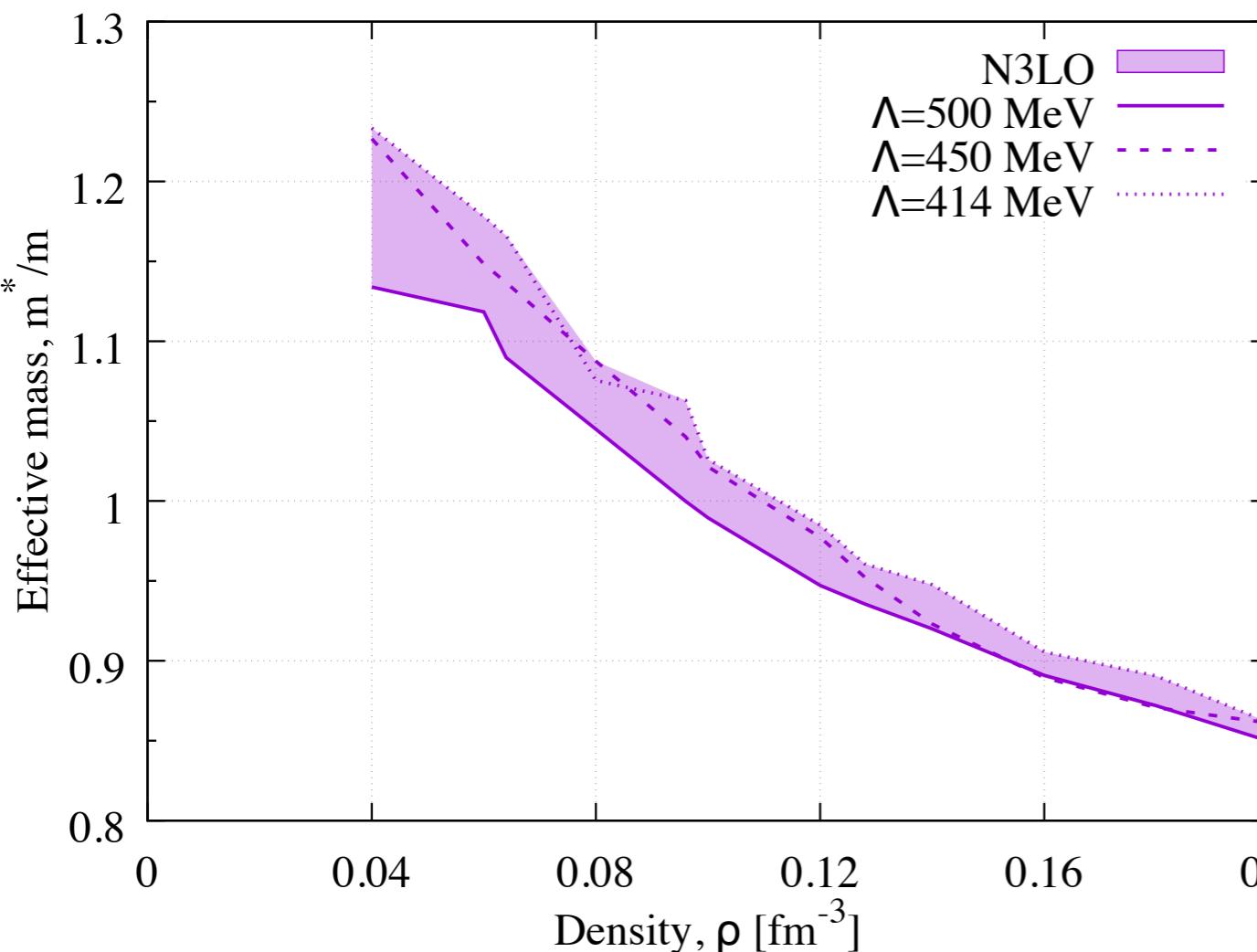
Effective masses

Neutron matter $\rho=0.16 \text{ fm}^{-3}$

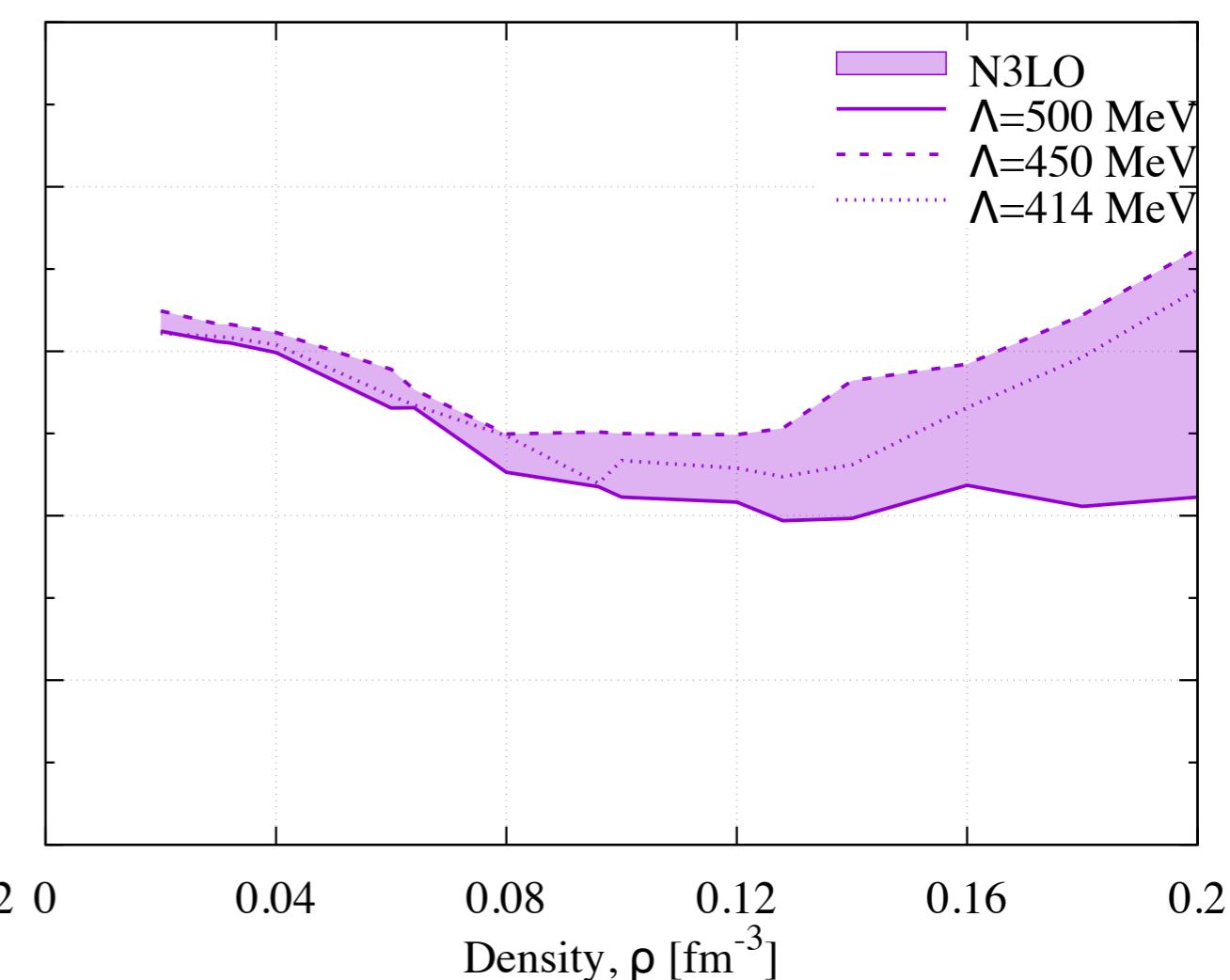


Coraggio, Holt, et al. PRC **89** 044321 (2014)

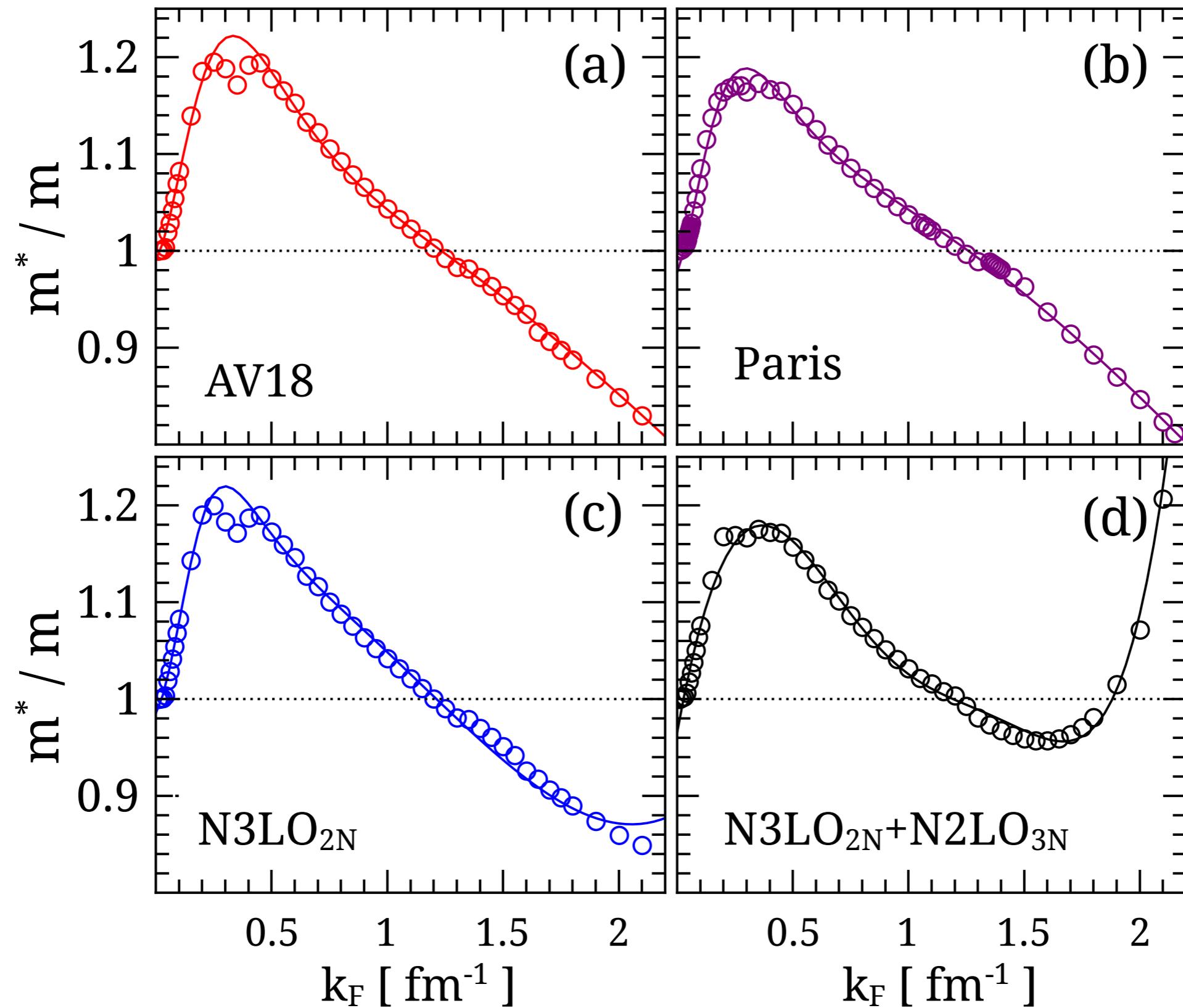
Symmetric matter



Neutron matter

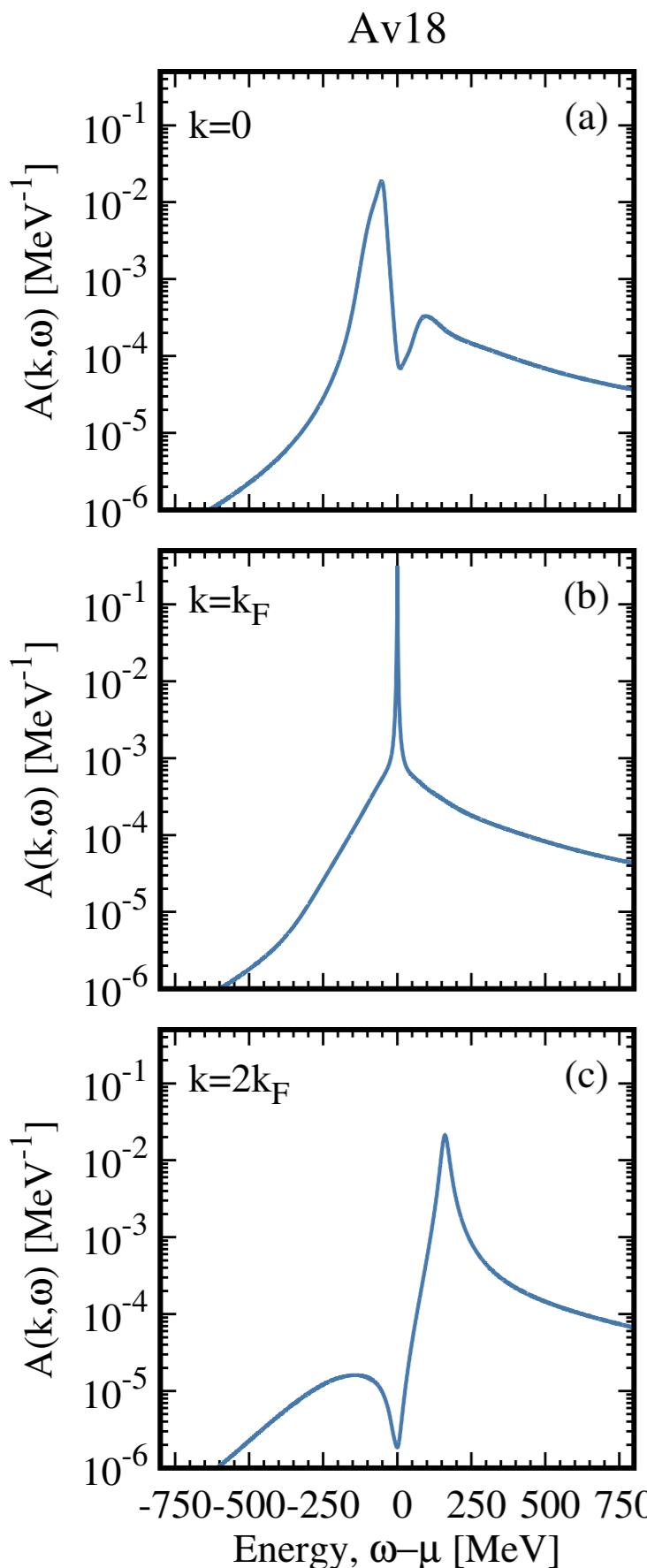


Density dependence of m^*



Spectral function

$T=5 \text{ MeV}$, $\rho=0.20 \text{ fm}^{-3}$



Can we **quantify differences?**

Energy weighted sum rules

$$m_k^{(0)} = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \mathcal{A}_k(\omega) = 1$$

$$m_k^{(1)} = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \omega \mathcal{A}_k(\omega) = \frac{\hbar^2 k^2}{2m} + \Sigma_k^{\infty}$$

$$m_k^{(2)} = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \omega^2 \mathcal{A}_k(\omega) = \left[m_k^{(1)} \right]^2 + \sigma_k^2$$

Ventura, Polls et al, PRC **49** 3050 (1994)

Frick, Mother & Polls, PRC **69** 054305 (2004)

Duguet & Hagen, PRC **85** 034330 (2012)

Duguet, Herbert et al, PRC **92** 034313 (2015)

Rios, Carbone, Polls, PRC **96** 014003 (2017)

Effective sp energies

Centroid energies

$$m_k^{(1)} = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \omega \mathcal{A}_k(\omega) = \frac{\hbar^2 k^2}{2m} + \Sigma_k^{\infty}$$

$$\begin{aligned} \Sigma_k^{\infty} &= \int \frac{d^3 k_1}{(2\pi)^3} \langle \vec{k} \vec{k}_1 | V | \vec{k} \vec{k}_1 \rangle_a n_{k_1} \\ &+ \frac{1}{2} \int \frac{d^3 k_1}{(2\pi)^3} \frac{d^3 k_2}{(2\pi)^3} \langle \vec{k} \vec{k}_1 \vec{k}_2 | W | \vec{k} \vec{k}_1 \vec{k}_2 \rangle_a n_{k_1} n_{k_2} \end{aligned}$$

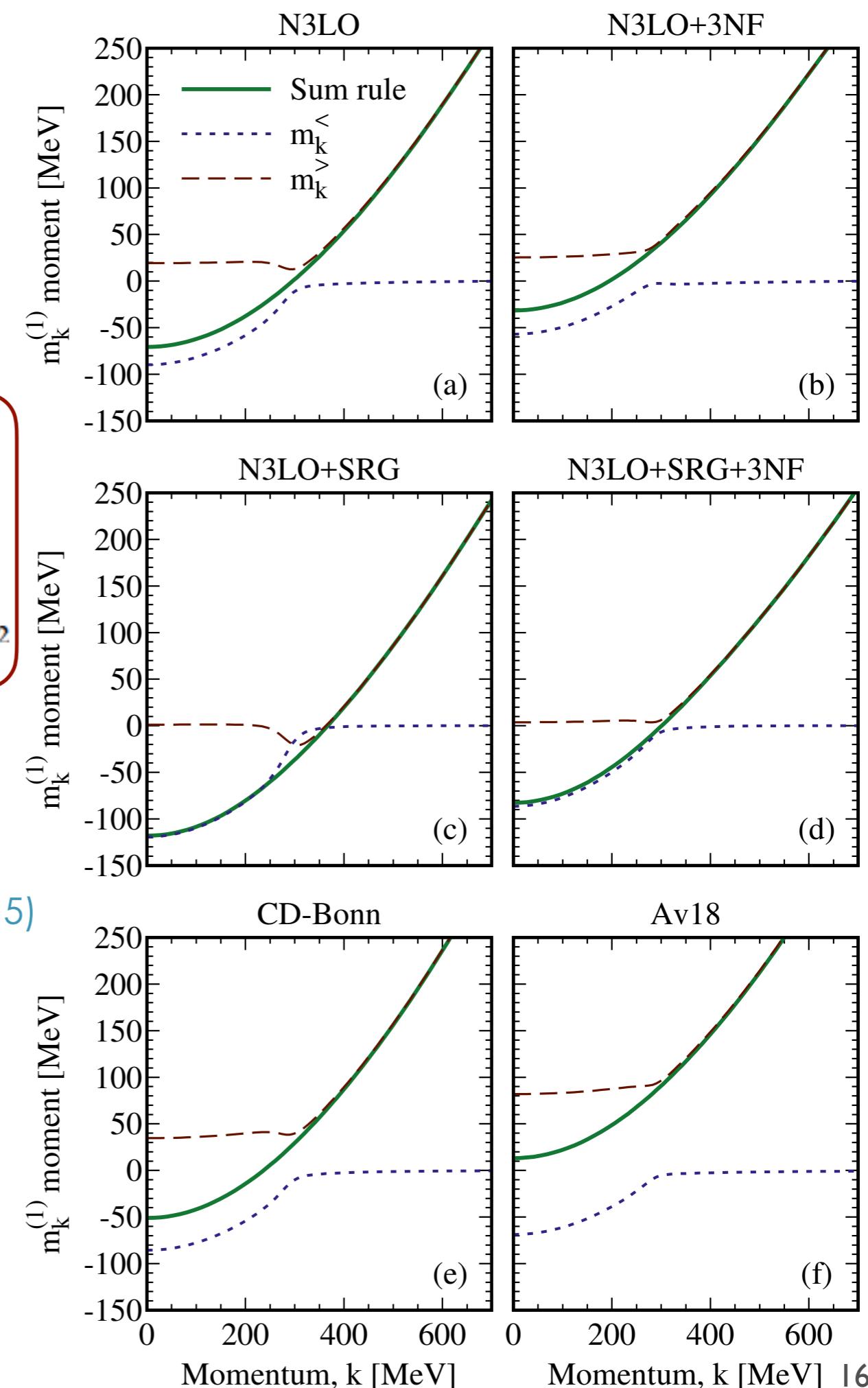
Cutoff evolution

Duguet, Herbert et al, PRC **92** 034313 (2015)

$$\frac{d}{d\lambda} O(\lambda) \equiv [\eta(\lambda), O(\lambda)],$$

$$\frac{d}{d\lambda} M_{pq}^{(0)}(\lambda) = 0,$$

$$\begin{aligned} \frac{d}{d\lambda} M_{pq}^{(1)}(\lambda) &= -\langle \Psi_0^A(\lambda) | \{[[\eta(\lambda), a_p], H(\lambda)], a_q^\dagger \} \\ &+ \{[a_p, H(\lambda)], [\eta(\lambda), a_q^\dagger]\} | \Psi_0^A(\lambda) \rangle. \end{aligned}$$

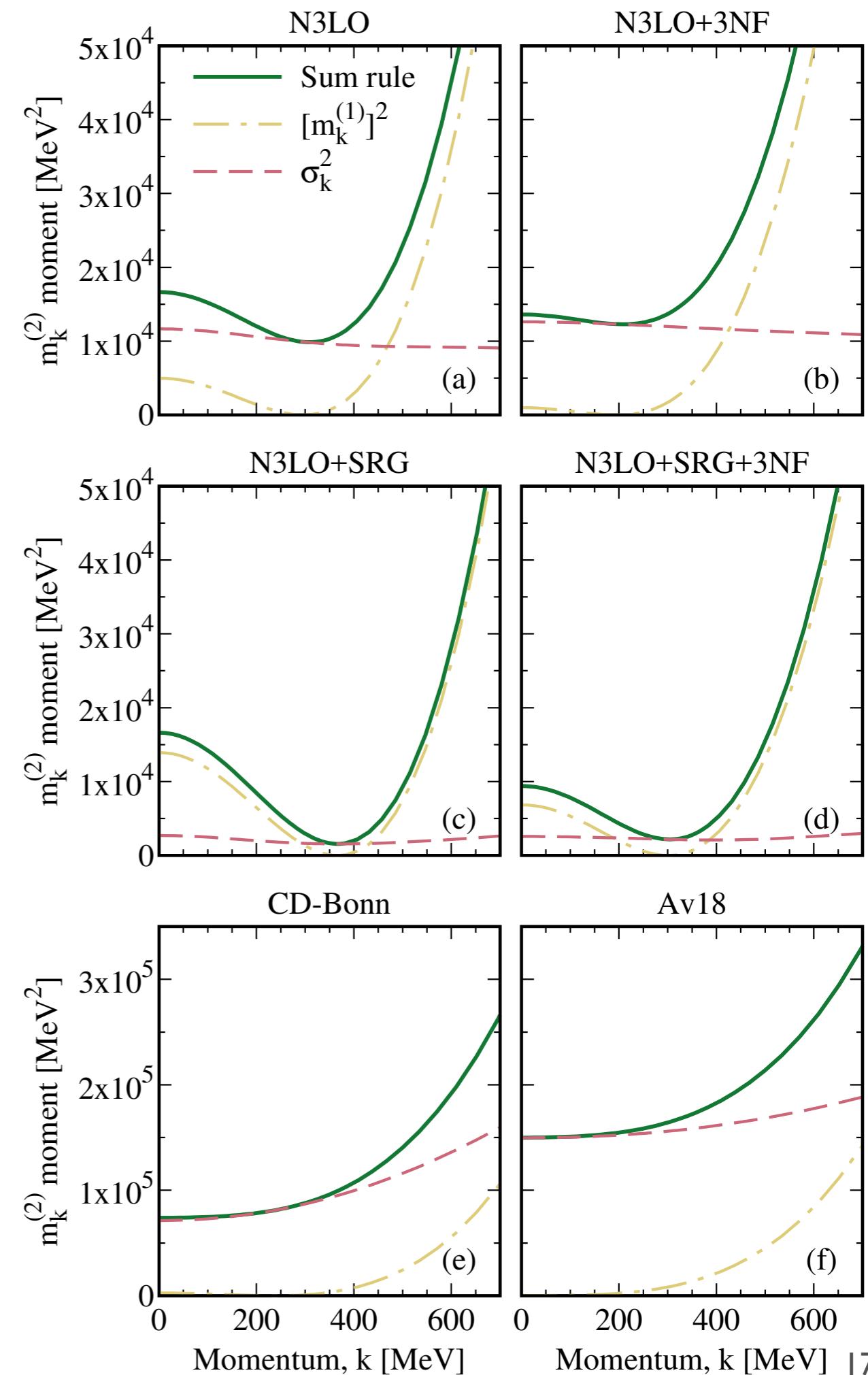


Second sum rule

sf variance

$$\sigma_k^2 = \int_{-\infty}^{+\infty} \frac{d\omega}{2\pi} [\omega - m_k^{(1)}]^2 \mathcal{A}_k(\omega)$$

$$\sigma_k^2 = - \int_{-\infty}^{+\infty} \frac{d\omega}{\pi} \text{Im } \Sigma_k(\omega)$$

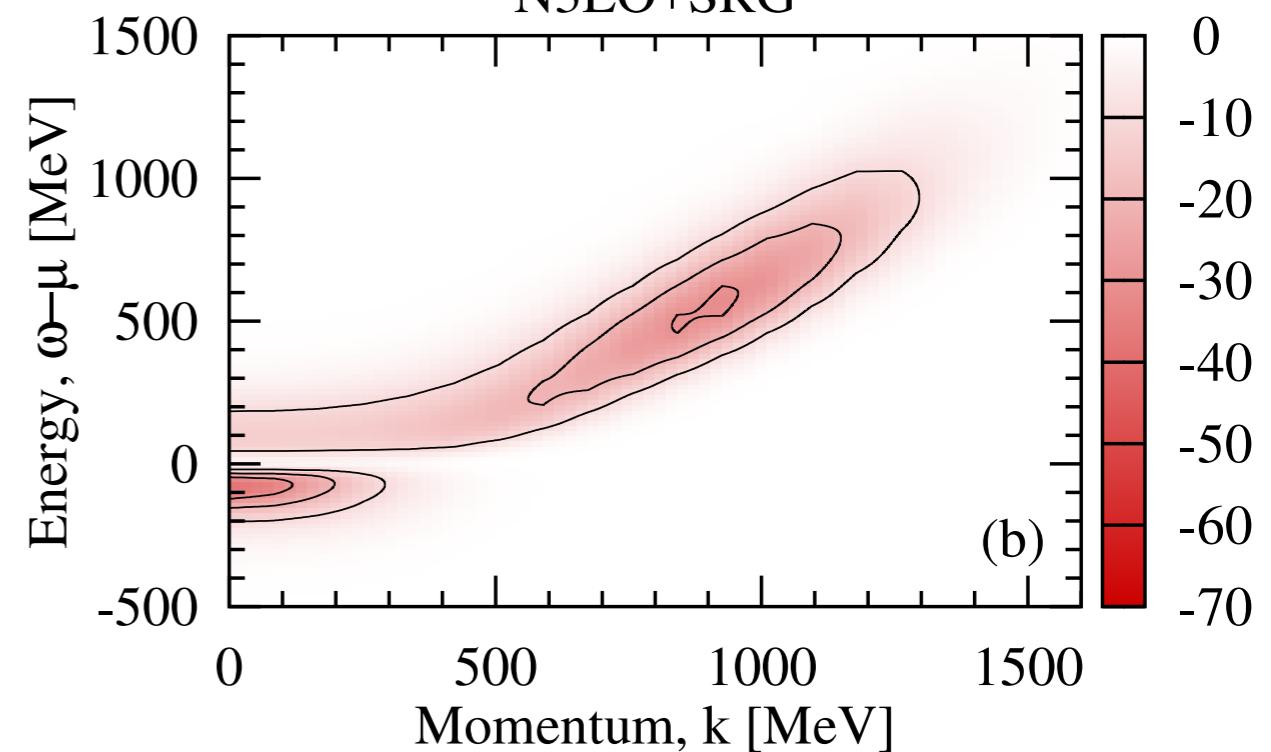
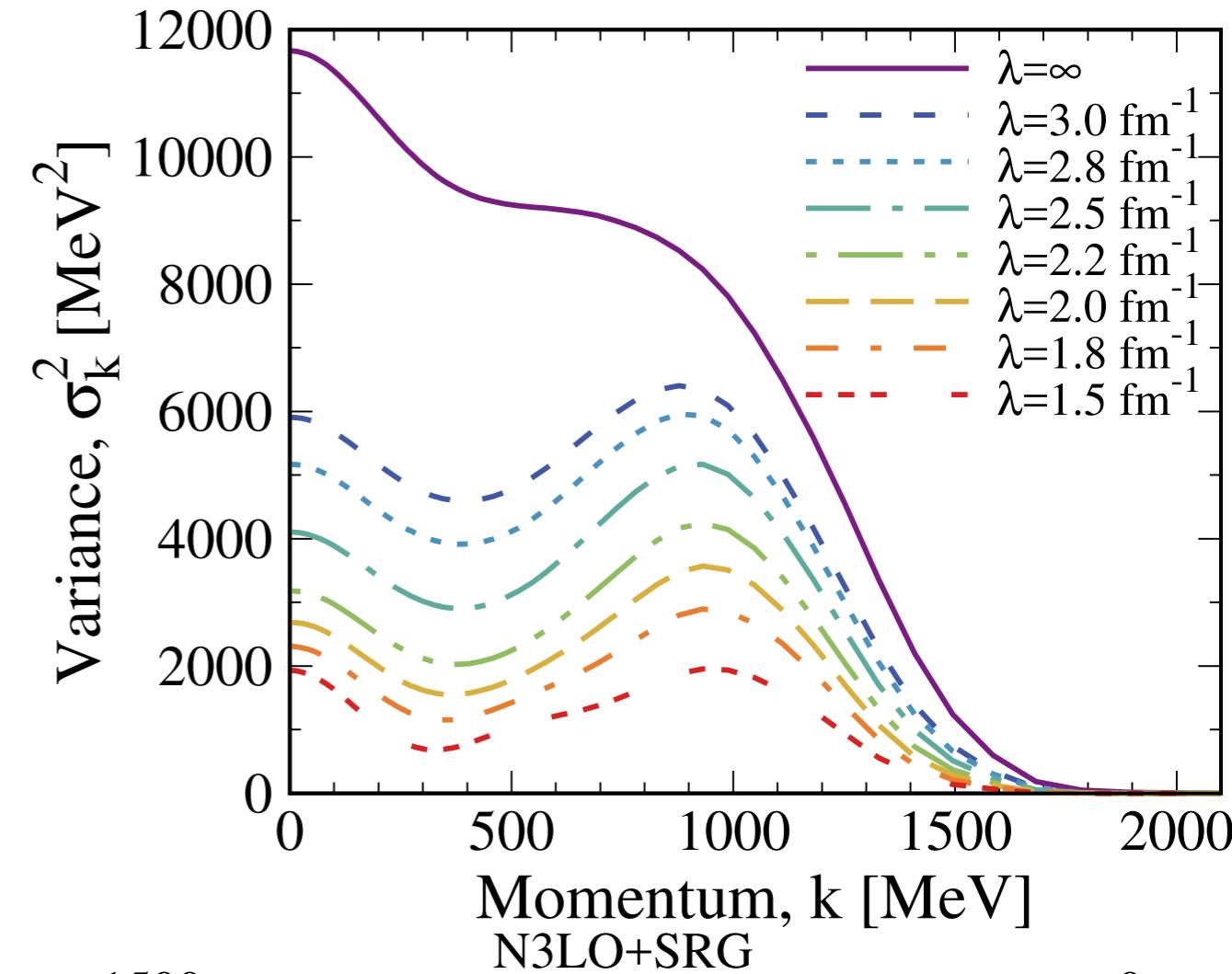
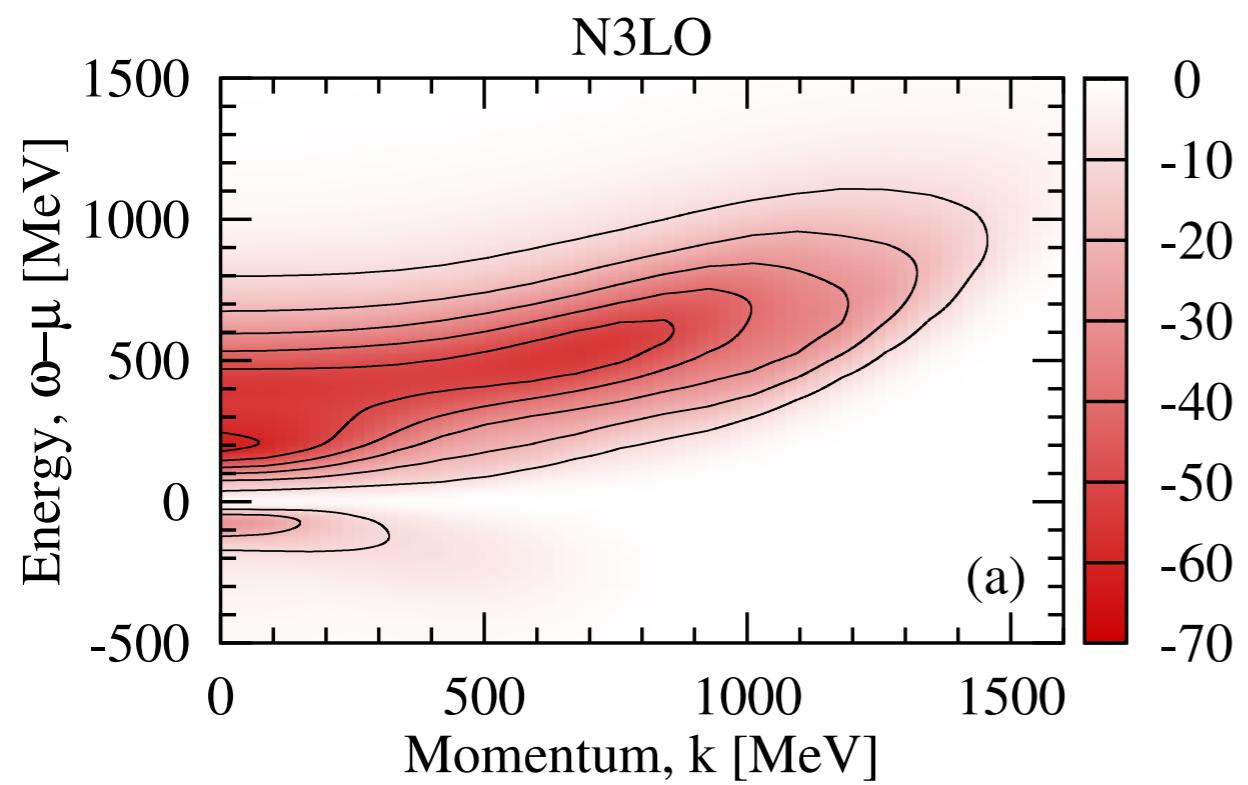


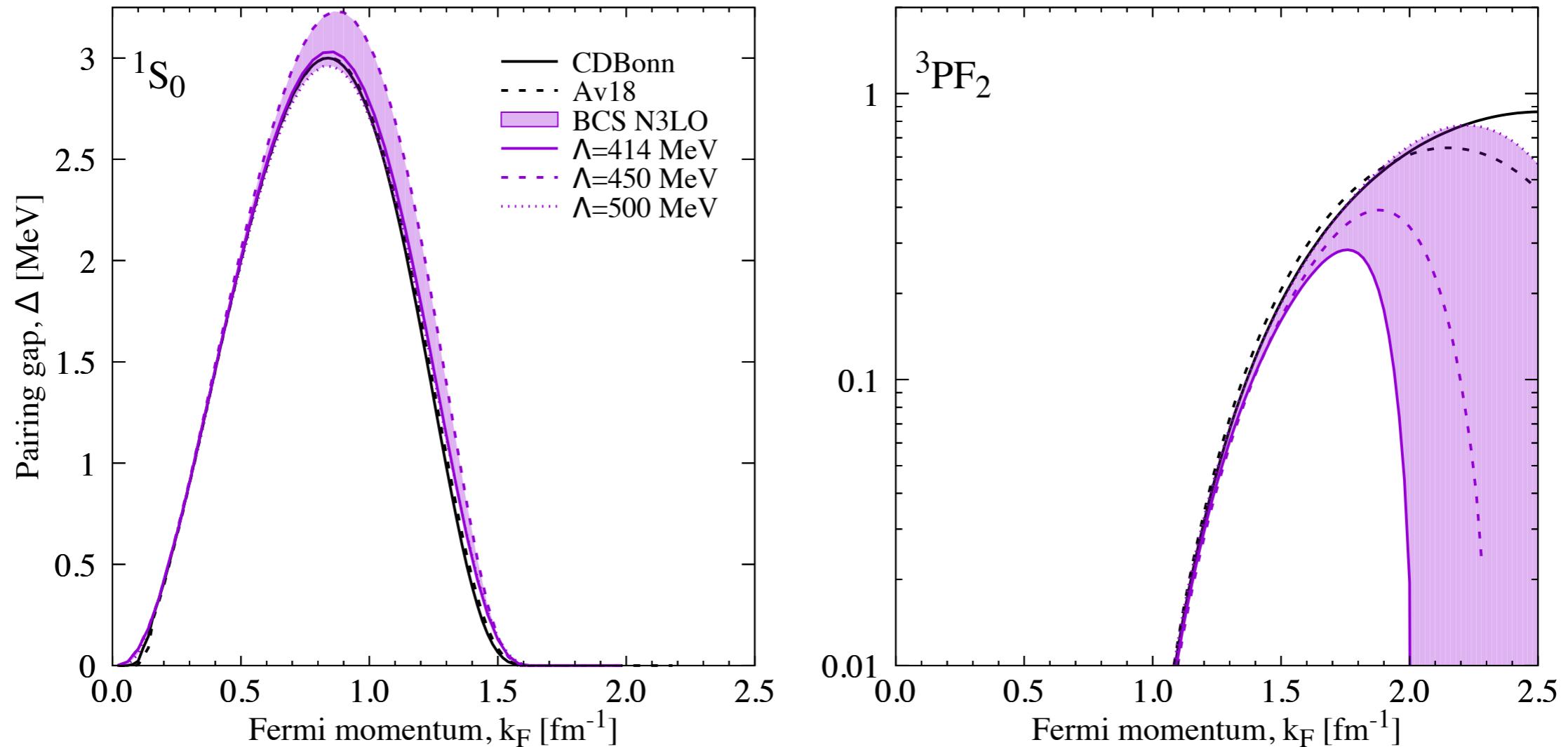
Second sum rule

sf variance

$$\sigma_k^2 = \int_{-\infty}^{+\infty} \frac{d\omega}{2\pi} [\omega - m_k^{(1)}]^2 \mathcal{A}_k(\omega)$$

$$\sigma_k^2 = - \int_{-\infty}^{+\infty} \frac{d\omega}{\pi} \text{Im } \Sigma_k(\omega)$$

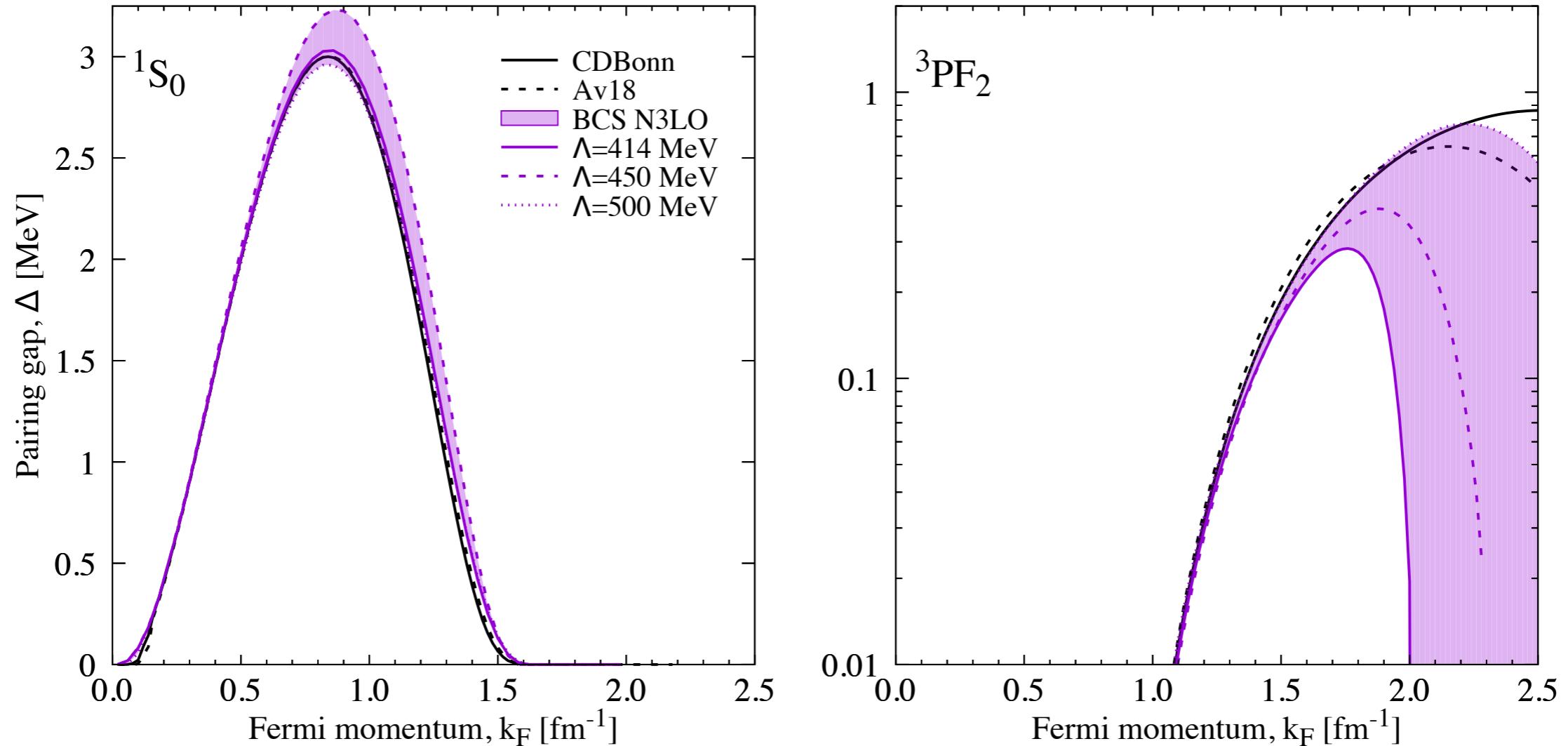




BCS equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'} \quad + \quad \chi_k = \varepsilon_k - \mu$$

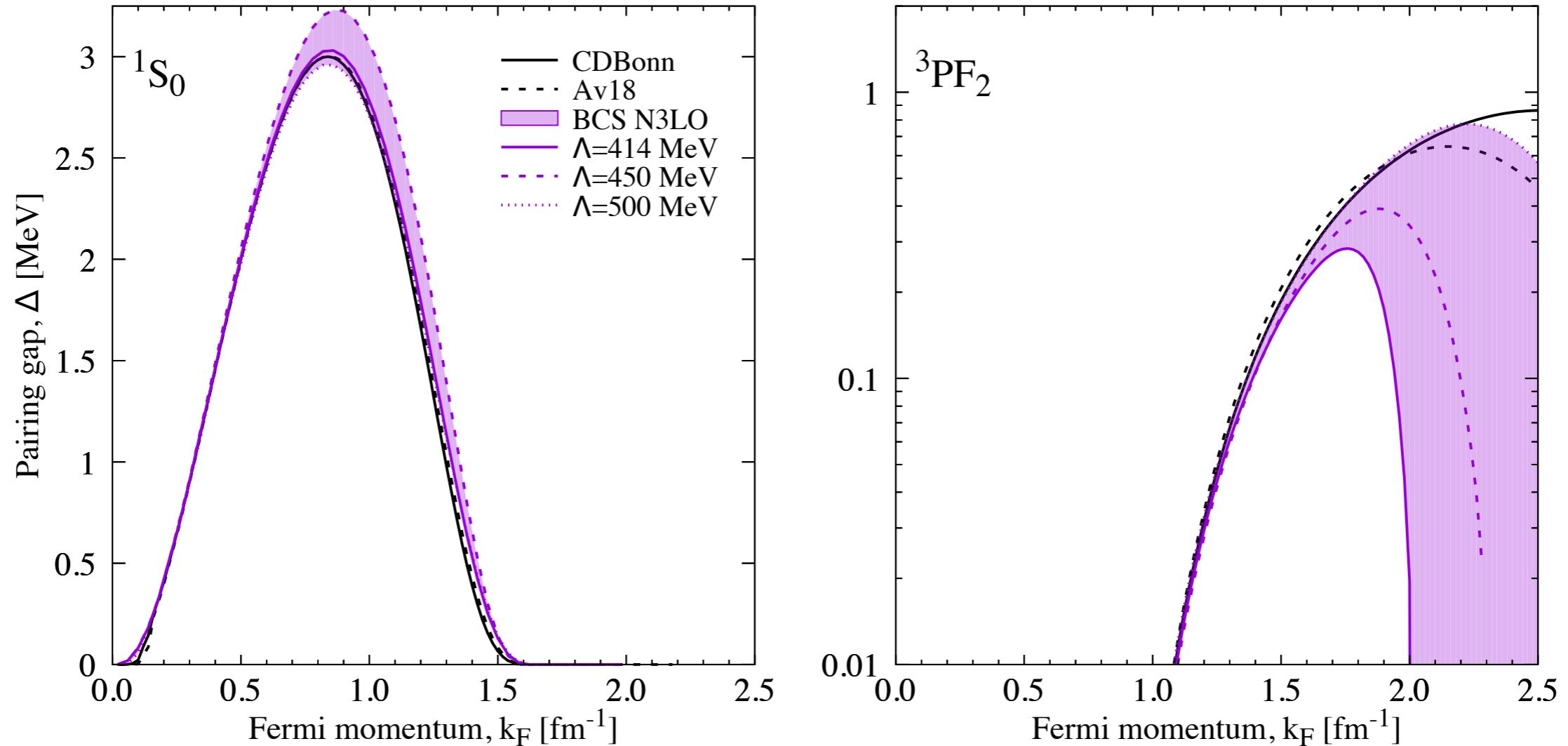
- Single-particle spectrum choice: $\varepsilon_k = \frac{k^2}{2m} + U(k) - \mu$
- Angular gap dependence: $|\Delta_k|^2 = \sum_L |\Delta_k^L|^2$



BCS equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'} \quad + \quad \chi_k = \varepsilon_k - \mu$$

- Single-particle spectrum choice: $\varepsilon_k = \frac{k^2}{2m} + \cancel{U(k)} - \mu$
- Angular gap dependence: $|\Delta_k|^2 = \sum_L |\Delta_k^L|^2$

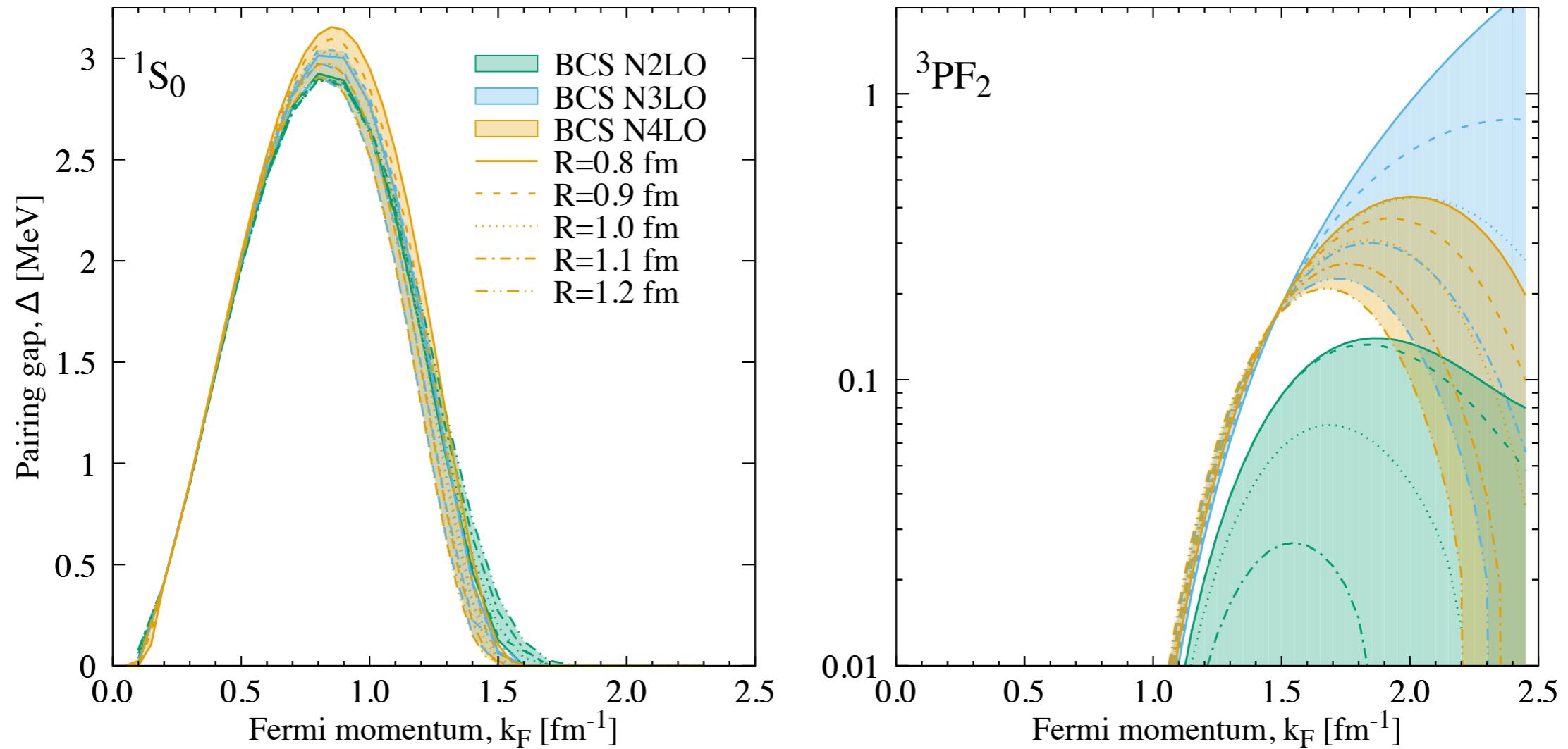


BCS equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'} \quad + \quad \chi_k = \varepsilon_k - \mu$$

- Single-particle spectrum choice: $\varepsilon_k = \frac{k^2}{2m} + \cancel{U(k)} - \mu$

- Angular gap dependence: $|\Delta_k|^2 = \sum_L |\Delta_k^L|^2 \approx |\Delta_k^L|^2$



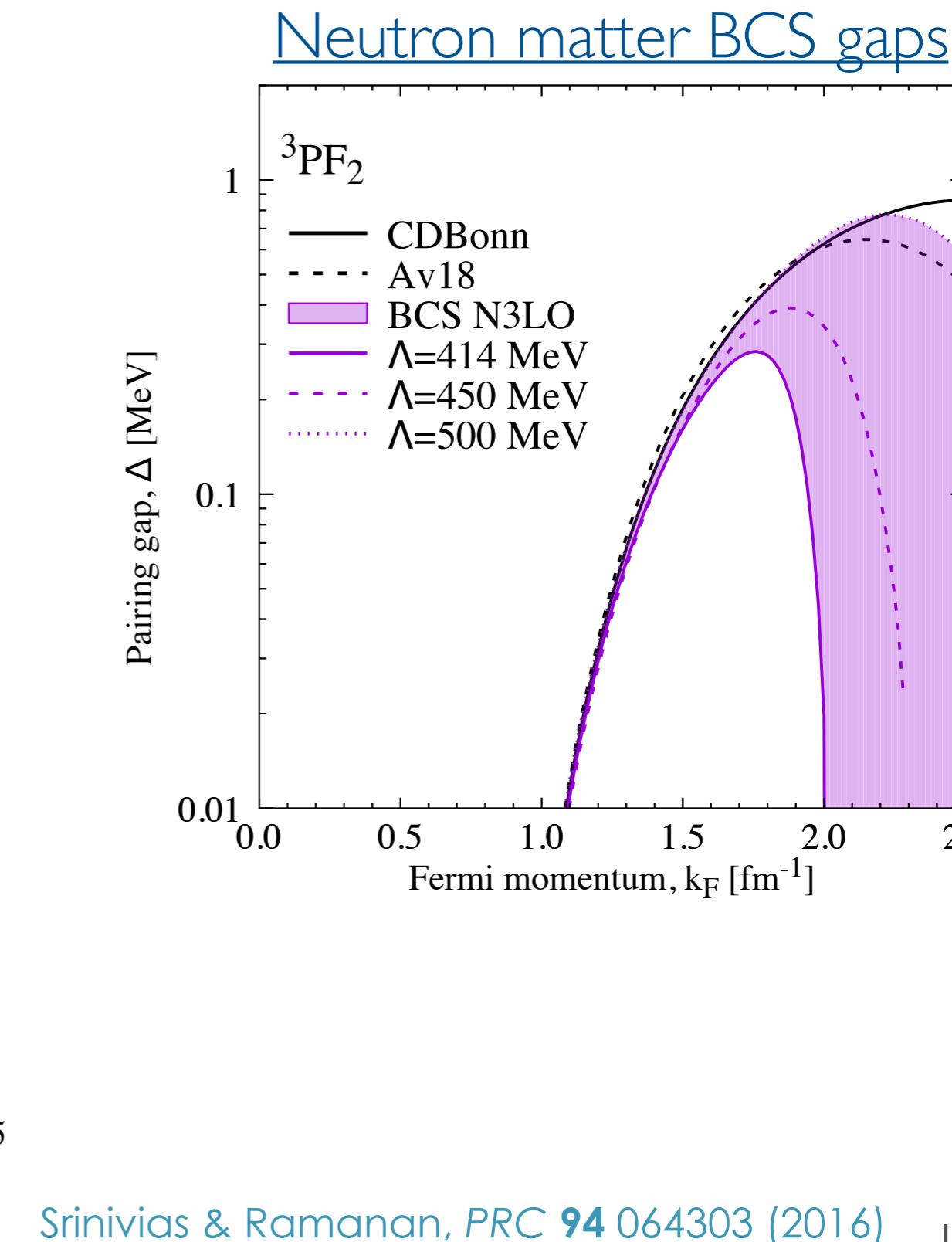
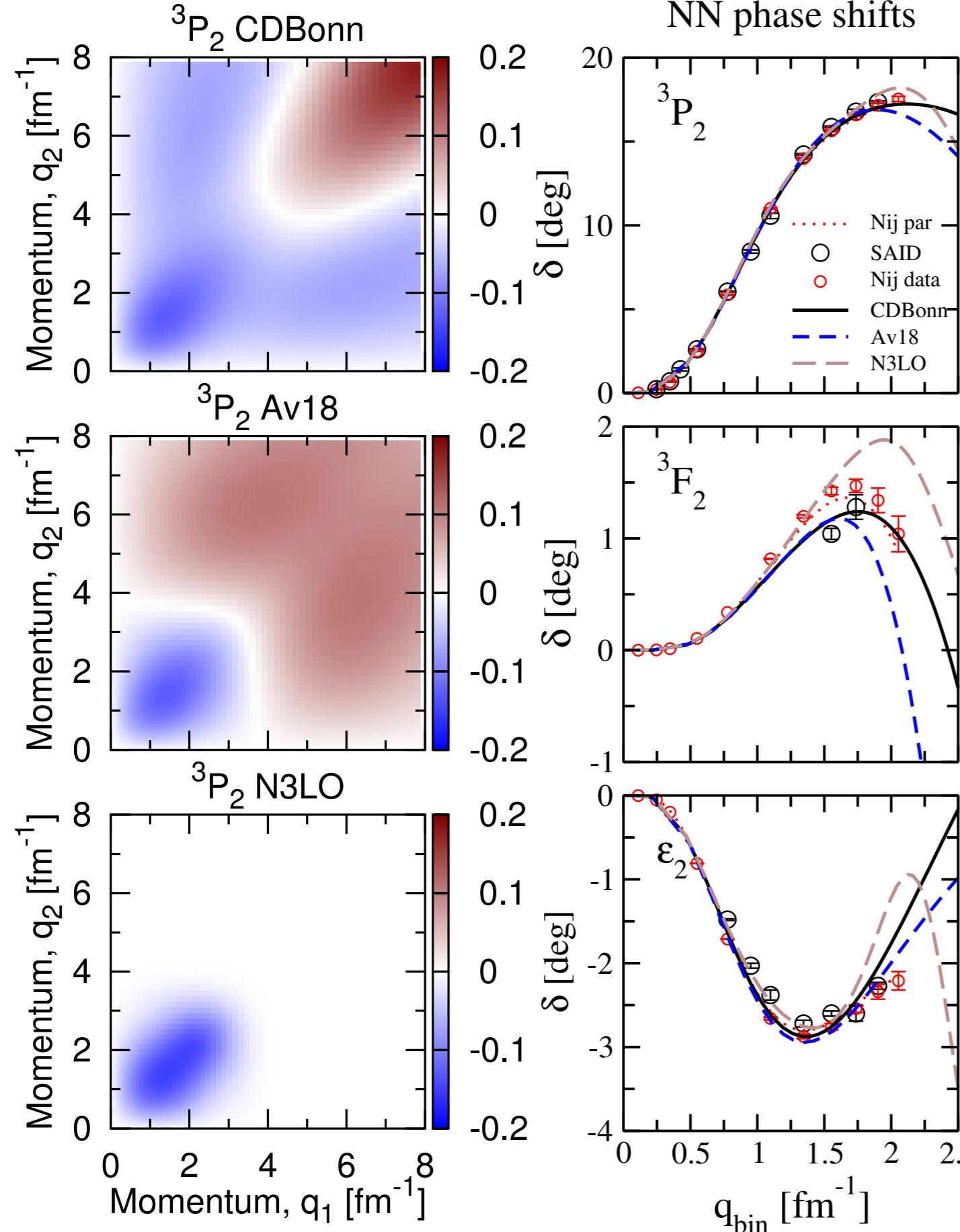
BCS equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'} \quad + \quad \chi_k = \varepsilon_k - \mu$$

- Single-particle spectrum choice: $\varepsilon_k = \frac{k^2}{2m} + \cancel{U(k)} - \mu$

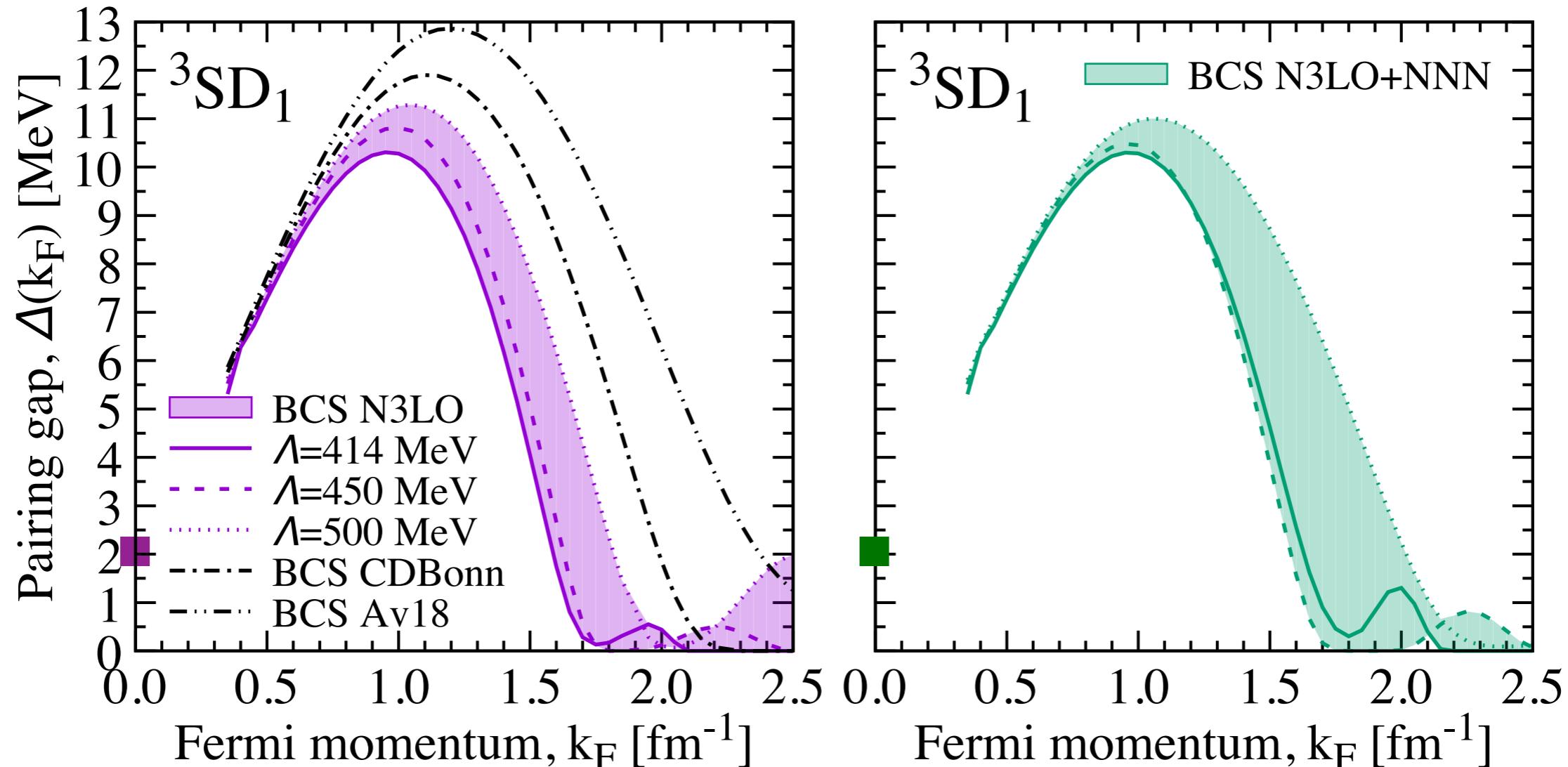
- Angular gap dependence: $|\Delta_k|^2 = \sum_L |\Delta_k^L|^2 \approx |\Delta_k^L|^2$

3PF_2 pairing: phase shift equivalence



Beyond BCS 101: 3SD_1

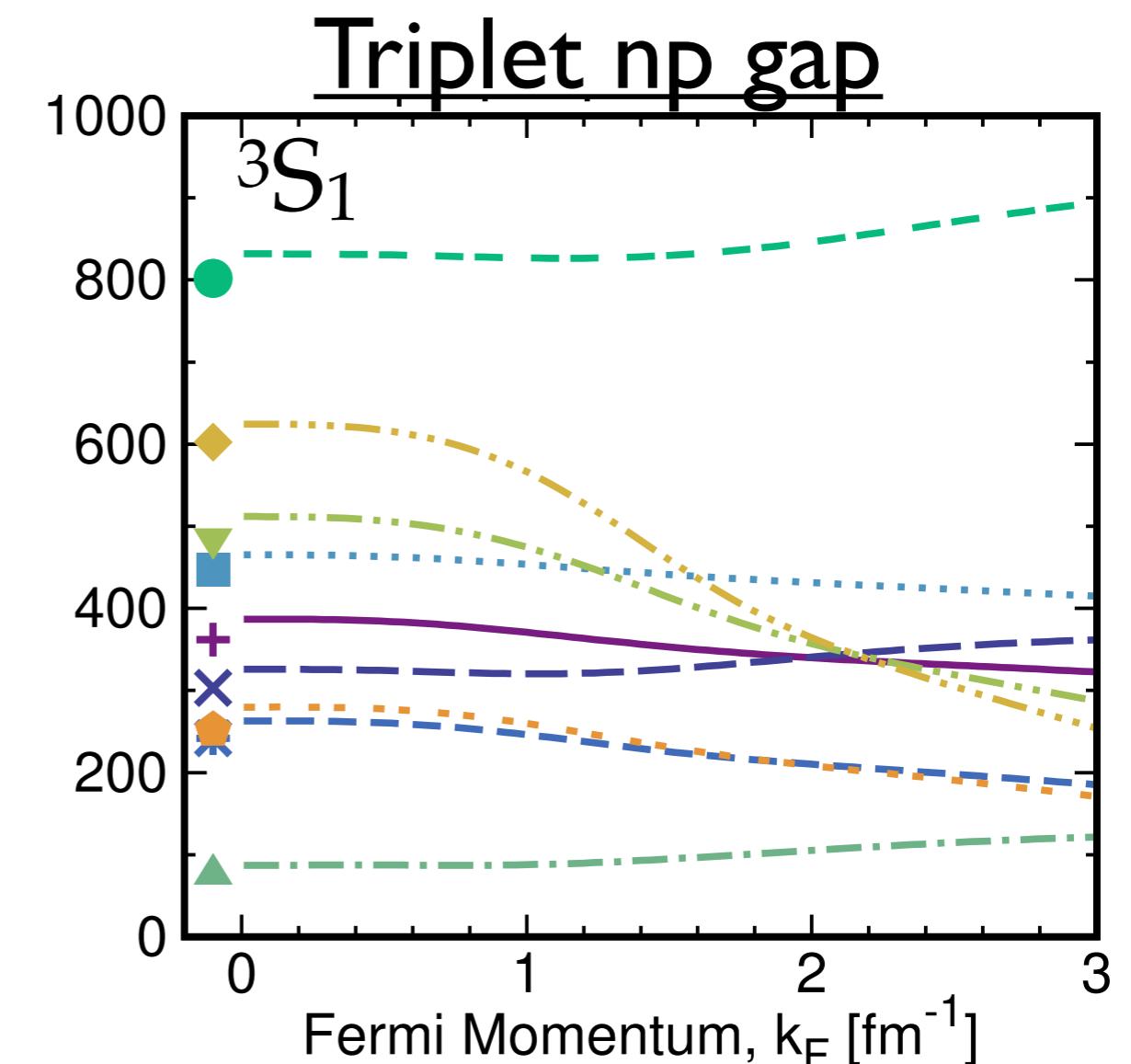
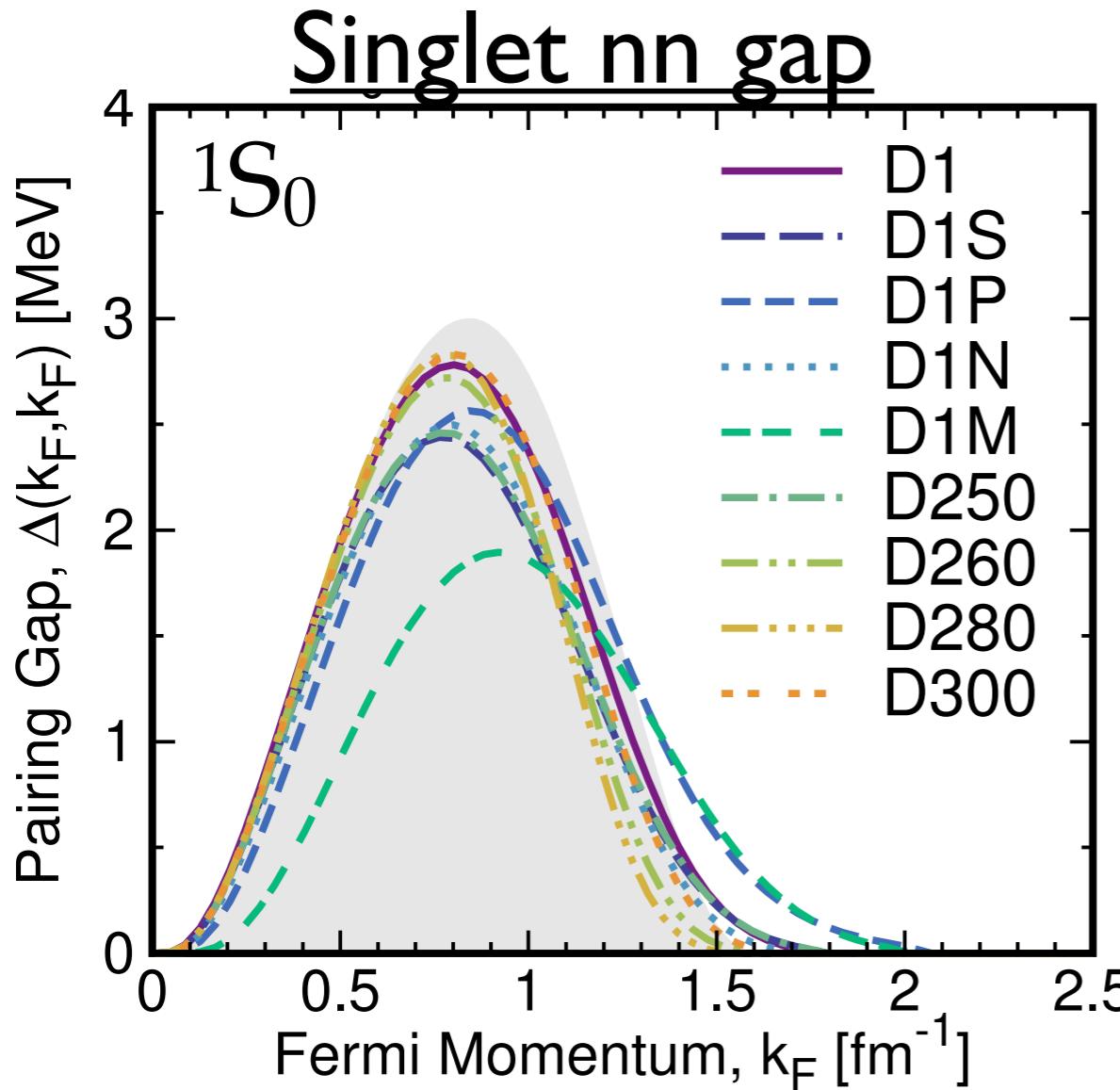
Symmetric matter



Muether & Dickhoff, *PRC* **72** 054313 (2005)

Rios, Polls & Dickhoff, arXiv:1707.04140
Maurizio, Holt & Finelli, *PRC* **90**, 044003 (2014)

- **Massive** gaps 3SD_1 channel but...
- **No** evidence of strong np nuclear pairing
- 3NF do **not** alter picture **significantly**
- Short-range correlations **deplete** gap

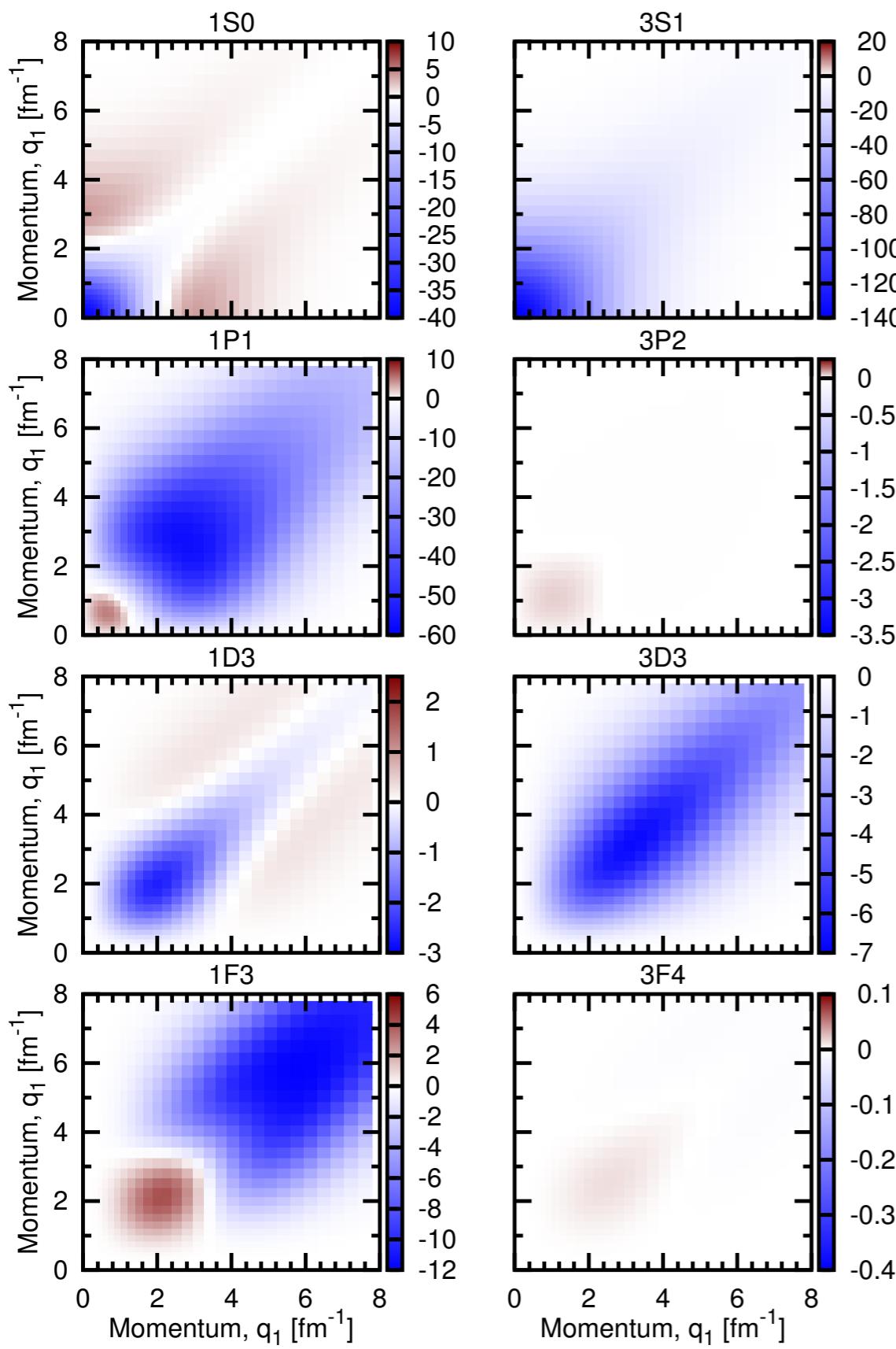


BCS equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'} + \quad \chi_k = \varepsilon_k - \mu$$

Matrix elements

D1S



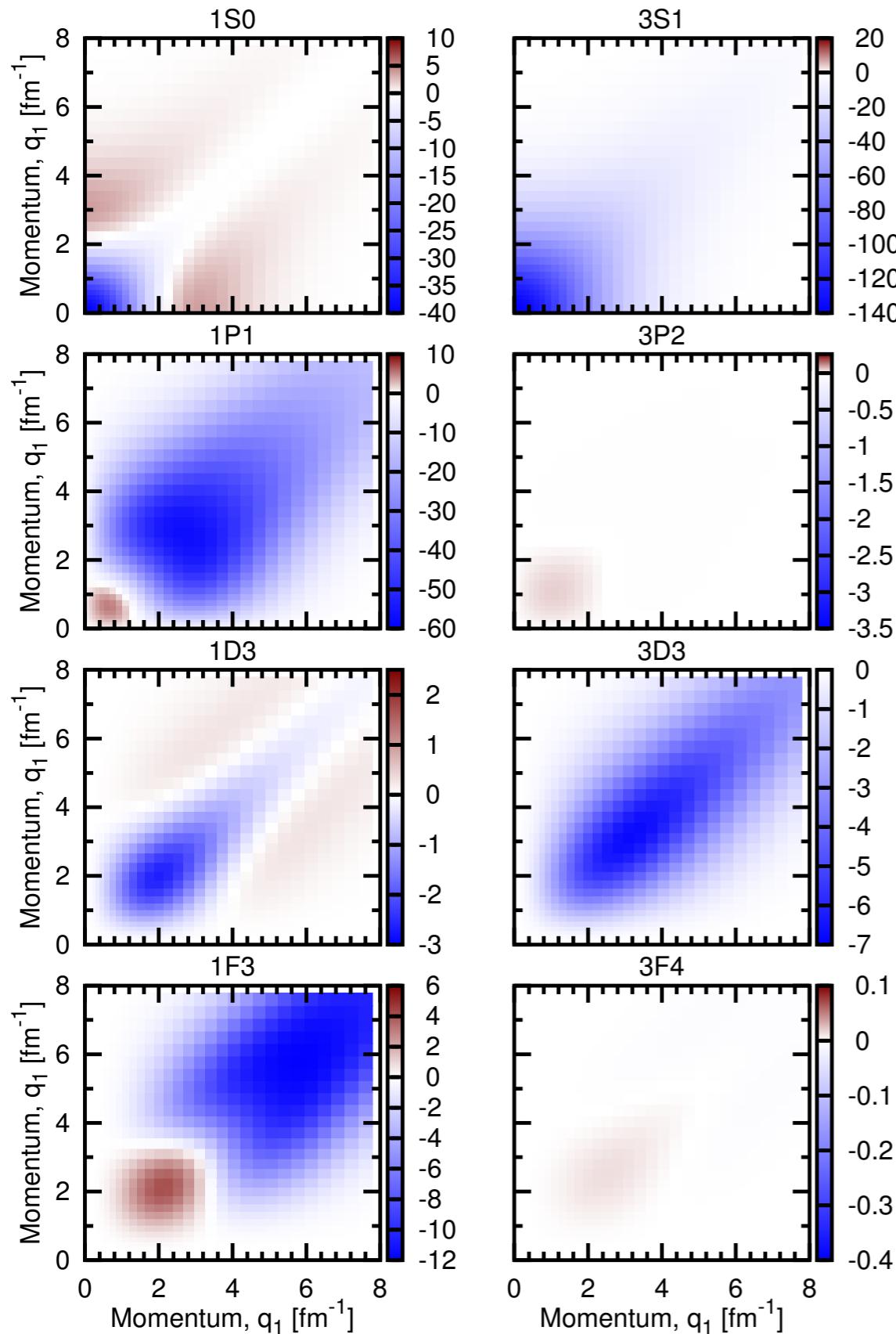
$$V^{LST}(q, q') = \sum_i^2 \left[Z_i^{LS} + F_i^{ST} e^{\frac{-\mu_i^2(k^2 + k'^2)}{4}} I_L \left(\frac{\mu_i^2 k k'}{2} \right) \right]$$

$$Z_i^{LS} = \frac{\delta_{L,0}}{4\pi^3} \left[t_0^i \rho^{\alpha_i} \left(1 - x_0^i (-)^S \right) \right]$$

$$F_i^{ST} = \frac{\mu_i^3}{2\pi^{\frac{1}{2}}} \left[W_i + B_i (-)^S - H_i (-)^T - M_i (-)^{S+T} \right]$$

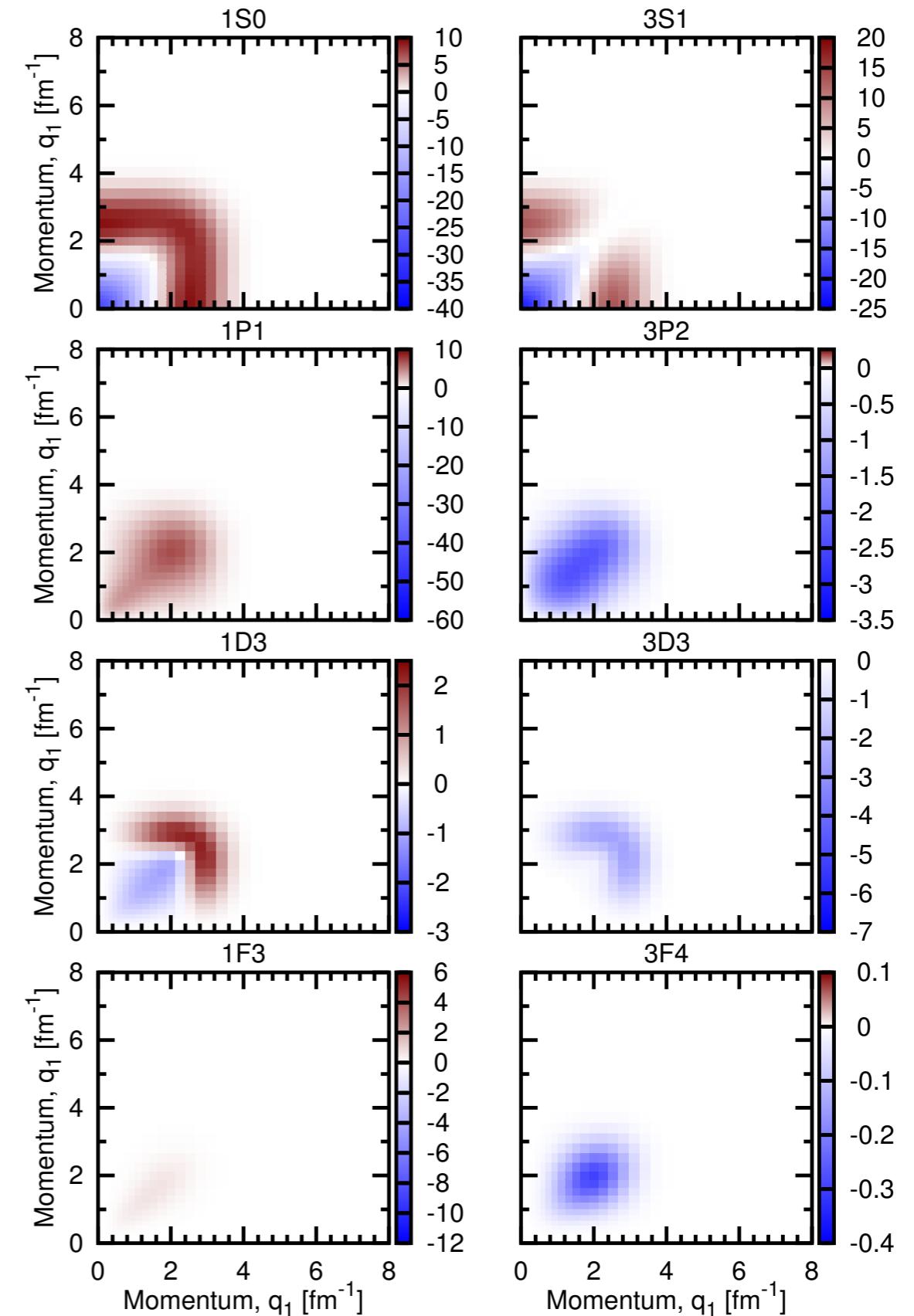
Matrix elements

D1S



UNIVERSITY OF
SURREY

N3LO



Postdoc position opening imminently

+A. Carbone



+D. Ding, W. H. Dickhoff



+A. Polls



+C. Barbieri



+V. Somà



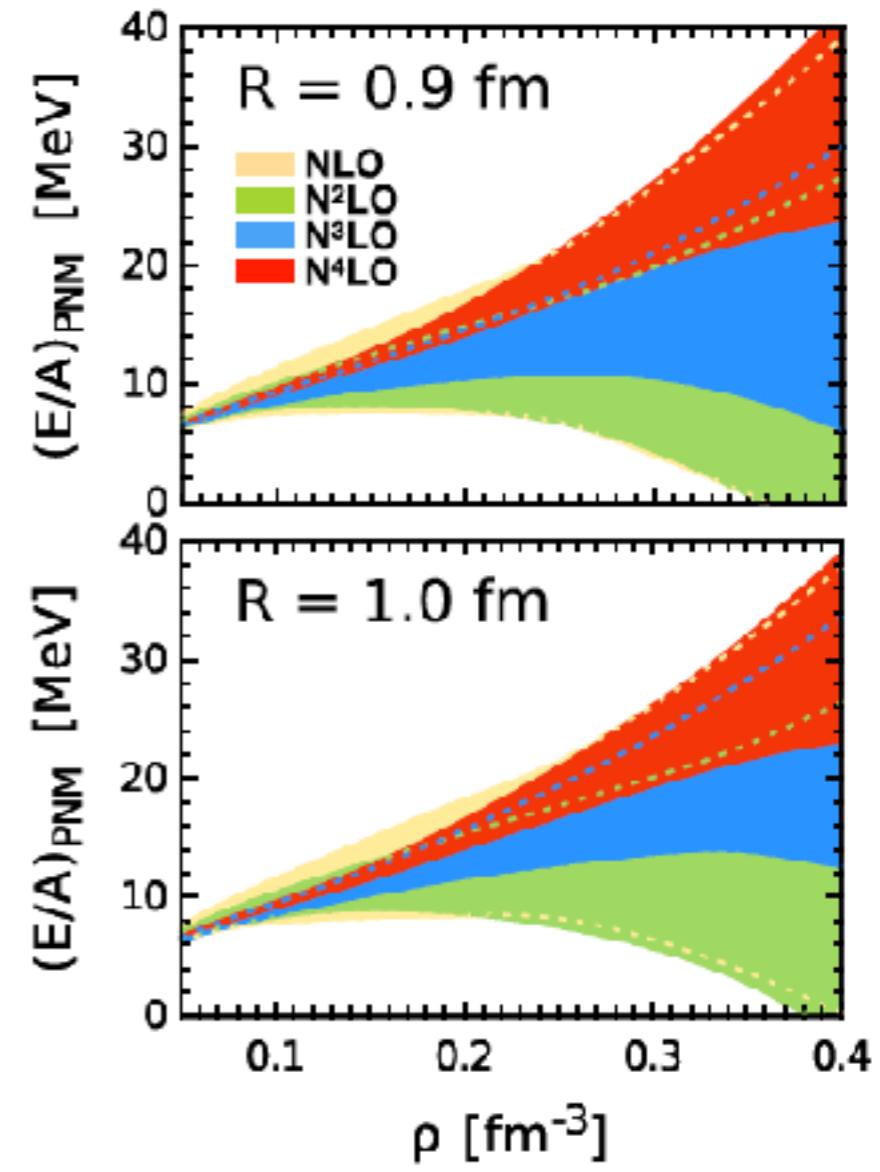
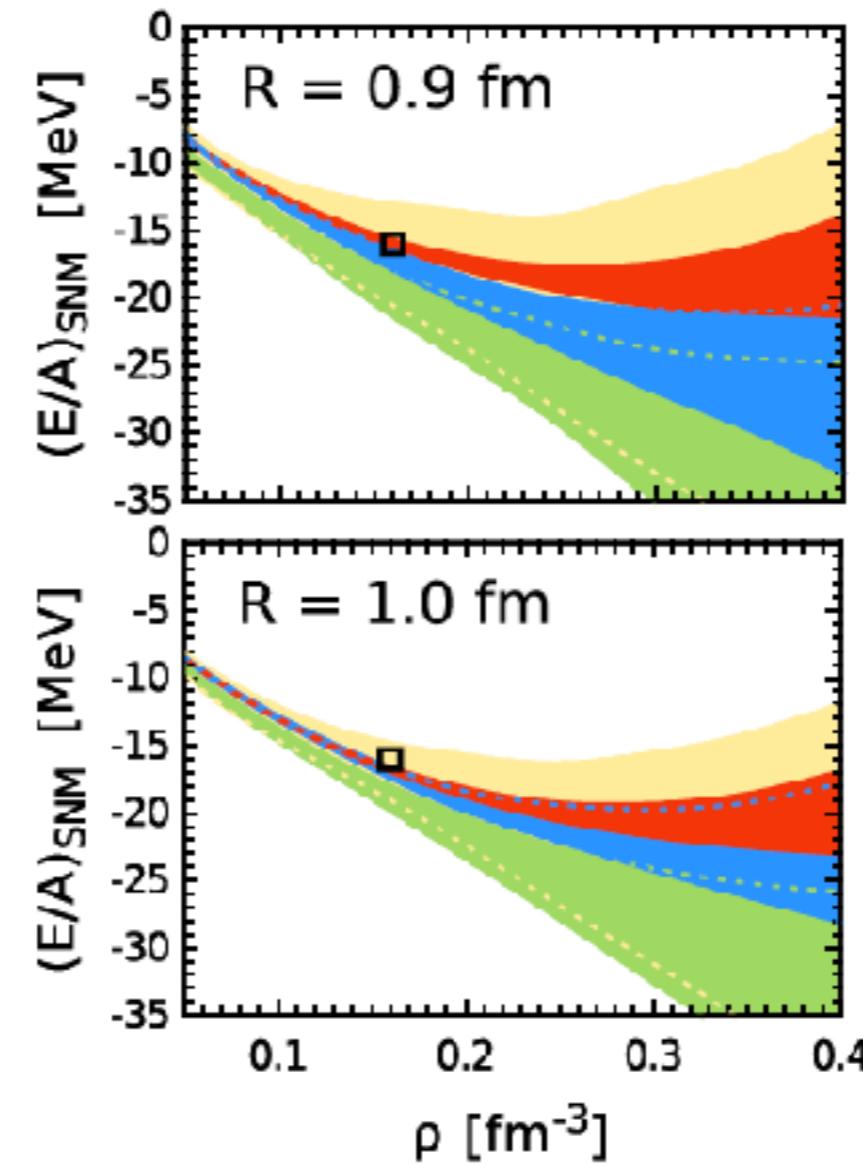
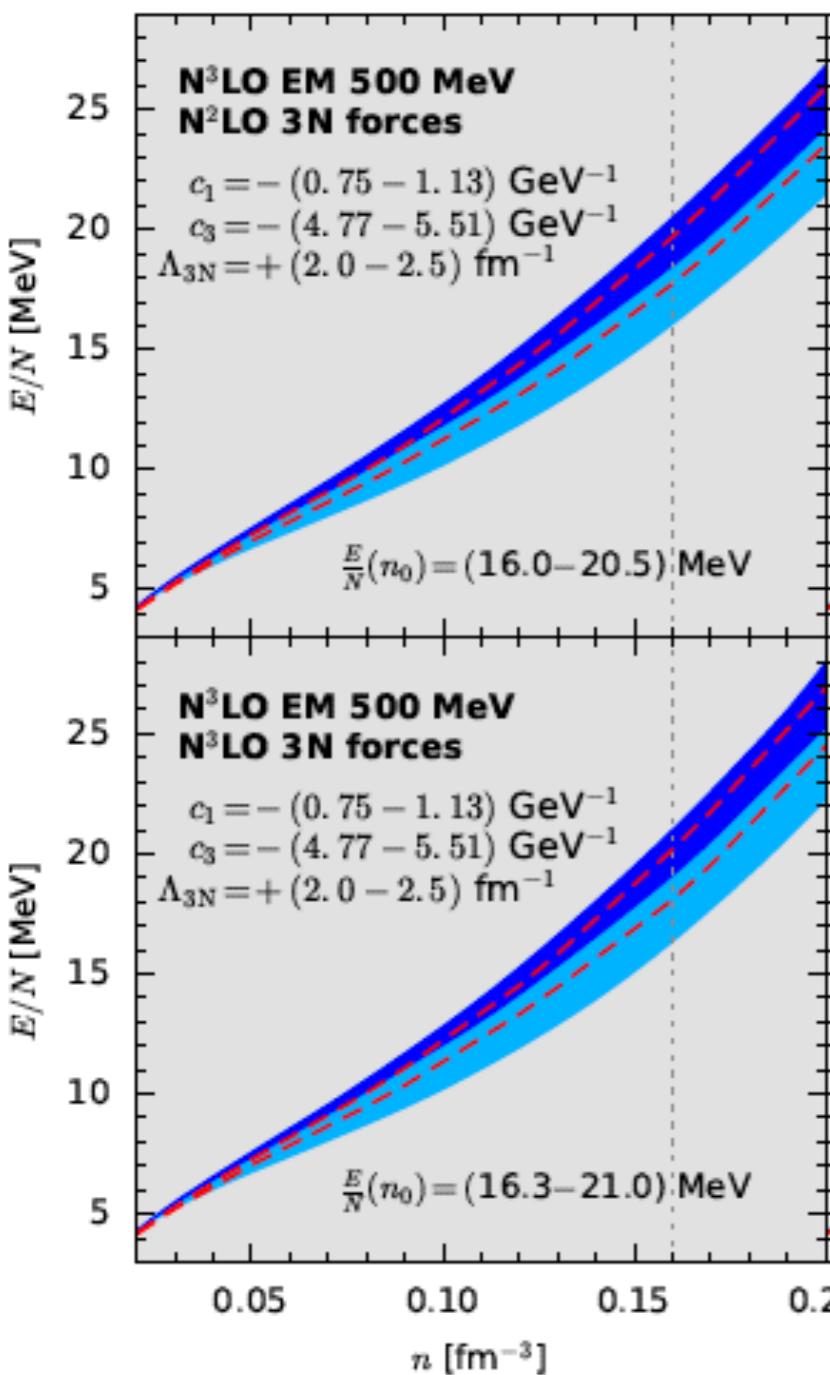
+H. Arellano, F. Isaule



Science & Technology
Facilities Council

a.rios@surrey.ac.uk
[@riosarnau](https://twitter.com/riosarnau)

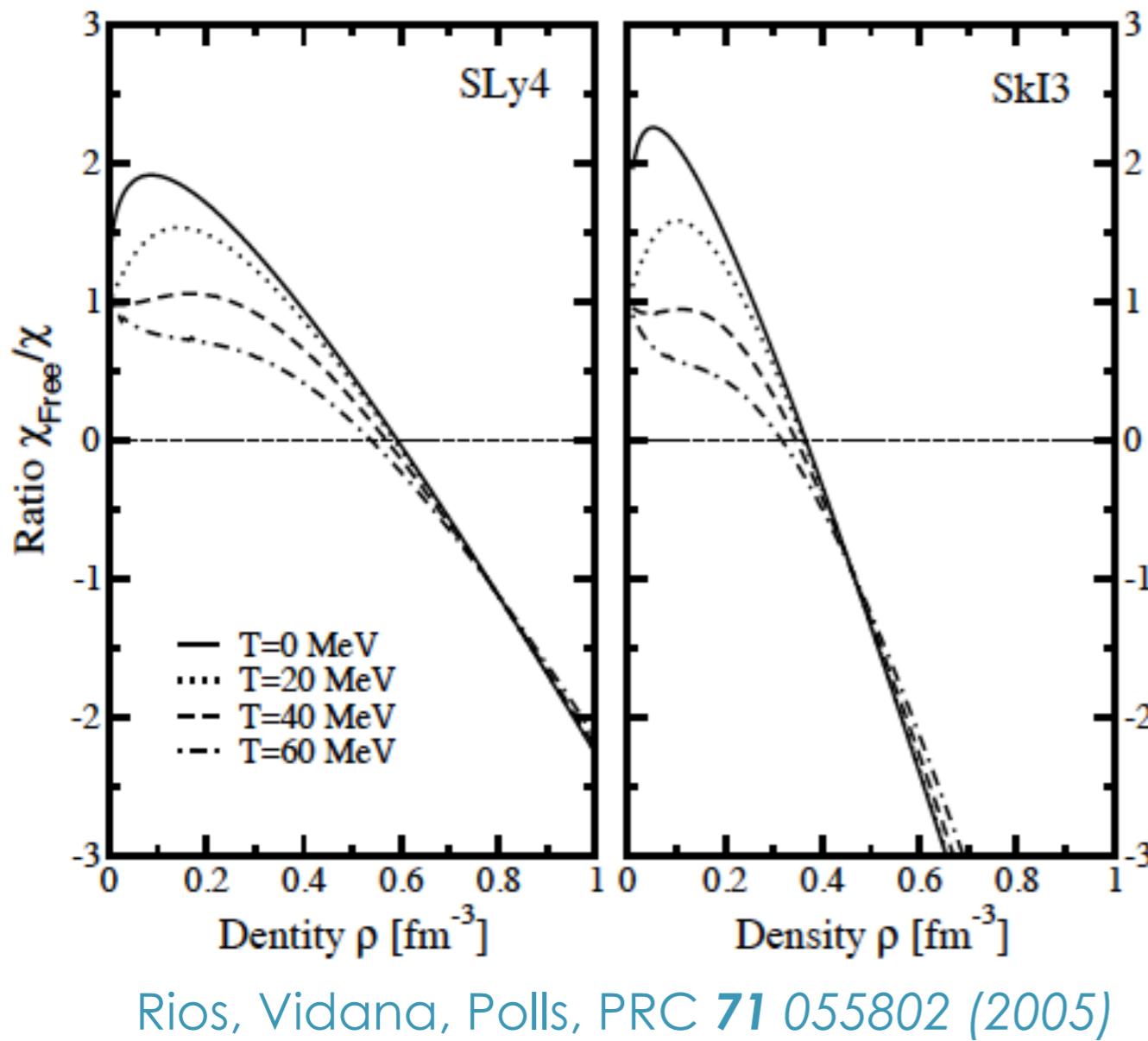
Quantifiable uncertainties should be propagated Not good to fit to a **single line** anymore!



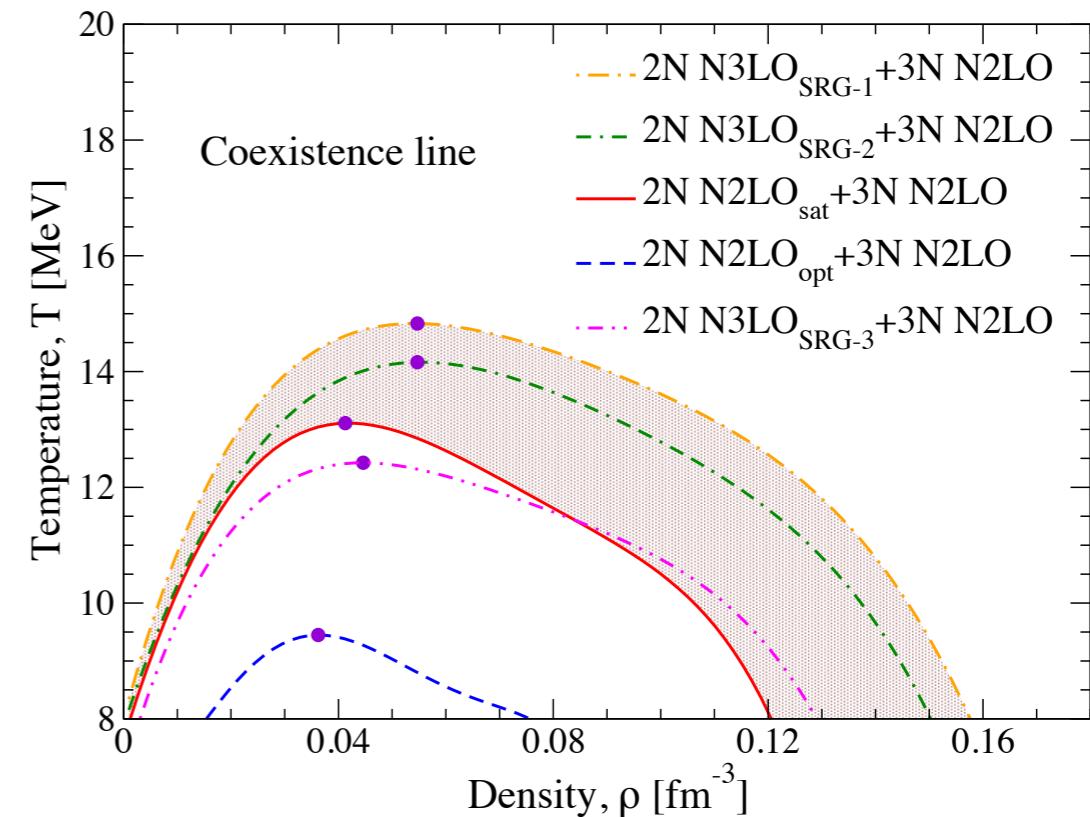
Drischler, Carbone, Hebeler, Schwenk PRC **94** 054307 (2016)
 Hu, Zhang, Epelbaum, Meissner, Meng, PRC **96** 034307 (2017)

Open questions

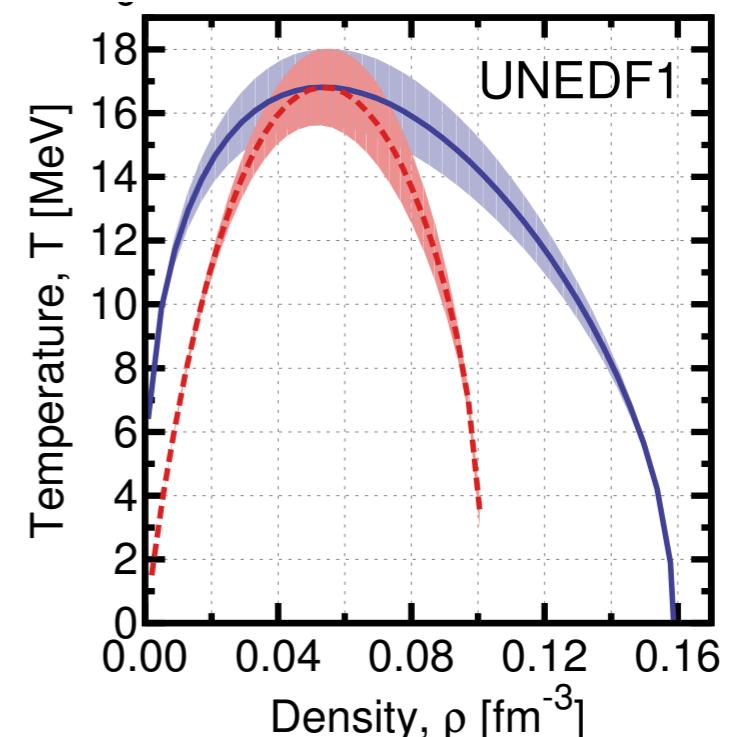
Polarized matter & finite temperature



Rios, Vidana, Polls, PRC **71** 055802 (2005)



Carbone, Rios, Polls, work in progress



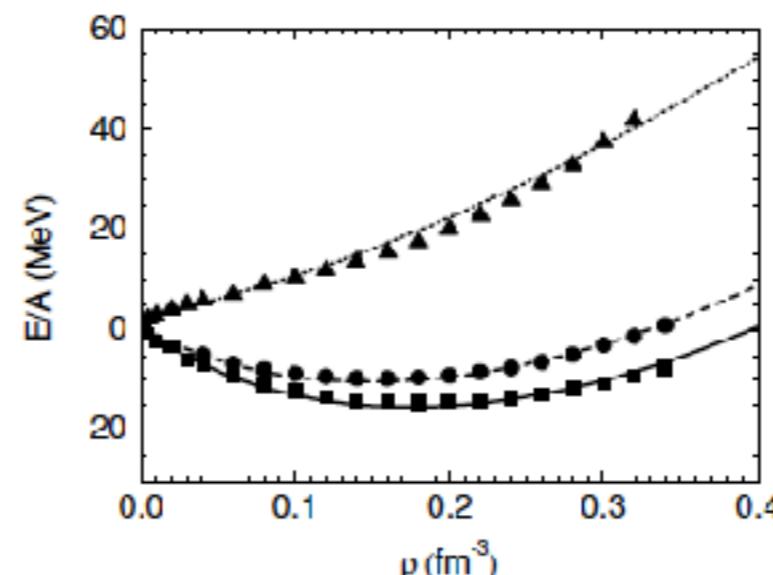
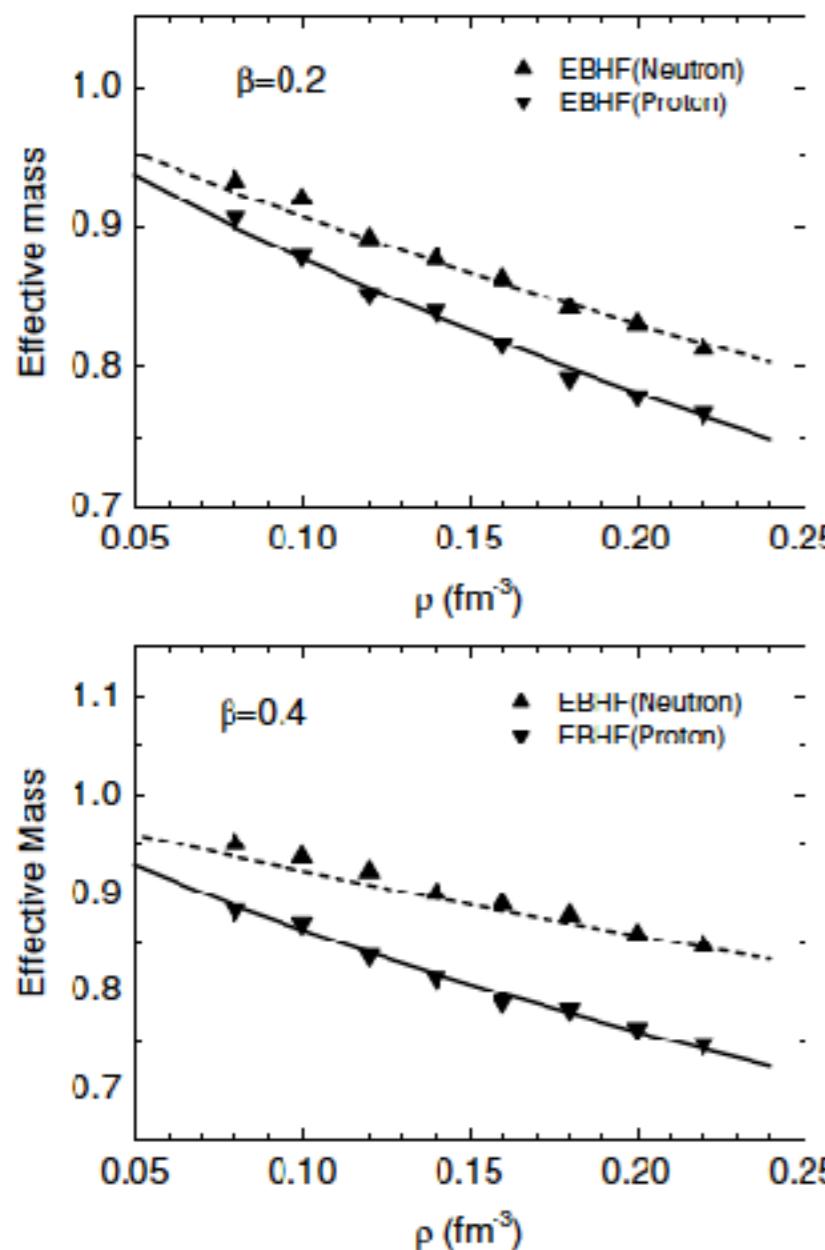
Rios & Roca-Maza, JPG **42** 034005 (2015) 25

Open questions

How can we get more **systematic** links?

Is there a way to go beyond **subjective** choices?

Landau parameters & response?



Skyrme LNS

