

# Liquid Argon TPCs for T2K experiment and future long baseline neutrino experiments

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# Neutrino oscillation physics

# Neutrino oscillations (I)

- Neutrino 3-flavour oscillation is well established both at the solar and atmospheric scale.
- This means that mass eigenstates and flavour eigenstates are different.
- Neutrinos are produced in weak interactions i.e. as flavour eigenstates and propagate as mass eigenstates (Hamiltonian eigenstates).
- The relationship between the two eigenstate bases can be expressed with a matrix equation:

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} \quad \longrightarrow \quad |\nu_\alpha(t)\rangle = \sum_{i=1}^n U_{\alpha i}^* |\nu_i(t)\rangle$$

- The matrix can be written in terms of 3 mixing angles ( $\theta_{ij}$ ) and 1 complex ( $\delta$ ) phase as:

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad \begin{matrix} s_{ij} = \sin(\theta_{ij}) \\ c_{ij} = \cos(\theta_{ij}) \end{matrix}$$

Atmospheric scale
Interference
Solar scale

# Neutrino oscillations (2)

- Using the standard plane wave approximation we can write the neutrino propagation as:

$$|\nu_i(t)\rangle = e^{-iE_i t} |\nu_i(0)\rangle$$

- Knowing that neutrinos are relativistic we can use the approximation:

$$E_i = \sqrt{p_i^2 + m_i^2} \simeq p + \frac{m_i^2}{2E}$$

- It results that the oscillation probability  $P_{\alpha\beta} = |\langle \nu_\beta | \nu_\alpha(t) \rangle|^2$  has terms proportional to:

$$\sin^2(\Delta m_{ij}^2 L / 4E)$$

→  
where

$$\Delta m_{ij}^2 \equiv m_i^2 - m_j^2$$

L = baseline (normally distance between  
ν source and detector)

E = ν energy

- The probability of oscillation from a flavour to a different one can be expressed as a function of the **3 mixing angles** ( $\theta_{12}, \theta_{13}, \theta_{23}$ ), **the complex phase** ( $\delta$ ) and **2 neutrino mass differences** ( $\Delta m_{21}^2, \Delta m_{31}^2$ ).

**NOTE:** neutrino oscillations are **NOT** sensitive to absolute masses but only to mass differences.

# Neutrino oscillations (3)

- As an example we take the 2-flavour case. The mixing matrix can be written as:

$$U = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix}$$

- The probability of oscillation from neutrino “a” to neutrino “b” is:

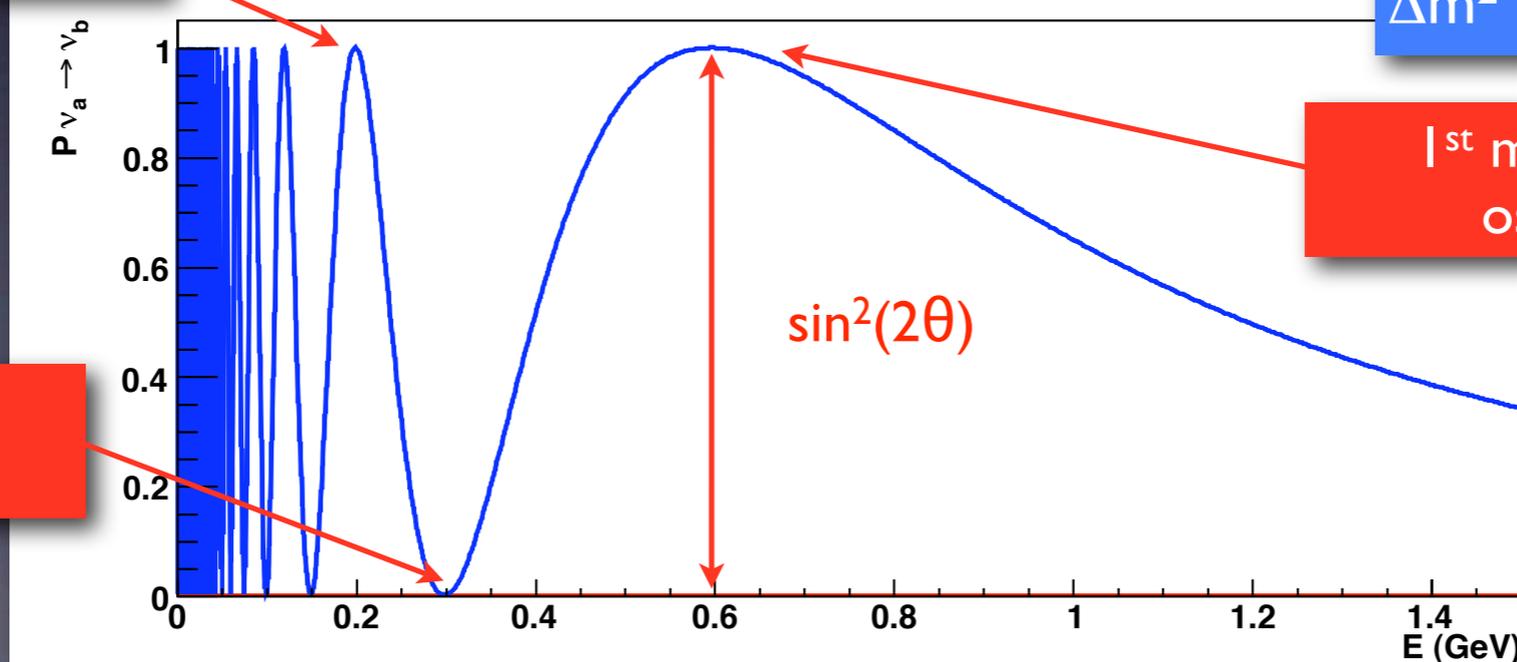
$$P_{ab} = \sin^2 2\theta \sin^2 \left( \frac{\Delta m^2 L}{4E} \right)$$

2<sup>nd</sup> maximum of oscillation

Baseline  $L = 295$  km  
 $\Delta m^2 = 2.5 \times 10^{-3}$  eV<sup>2</sup>

1<sup>st</sup> minimum of oscillation

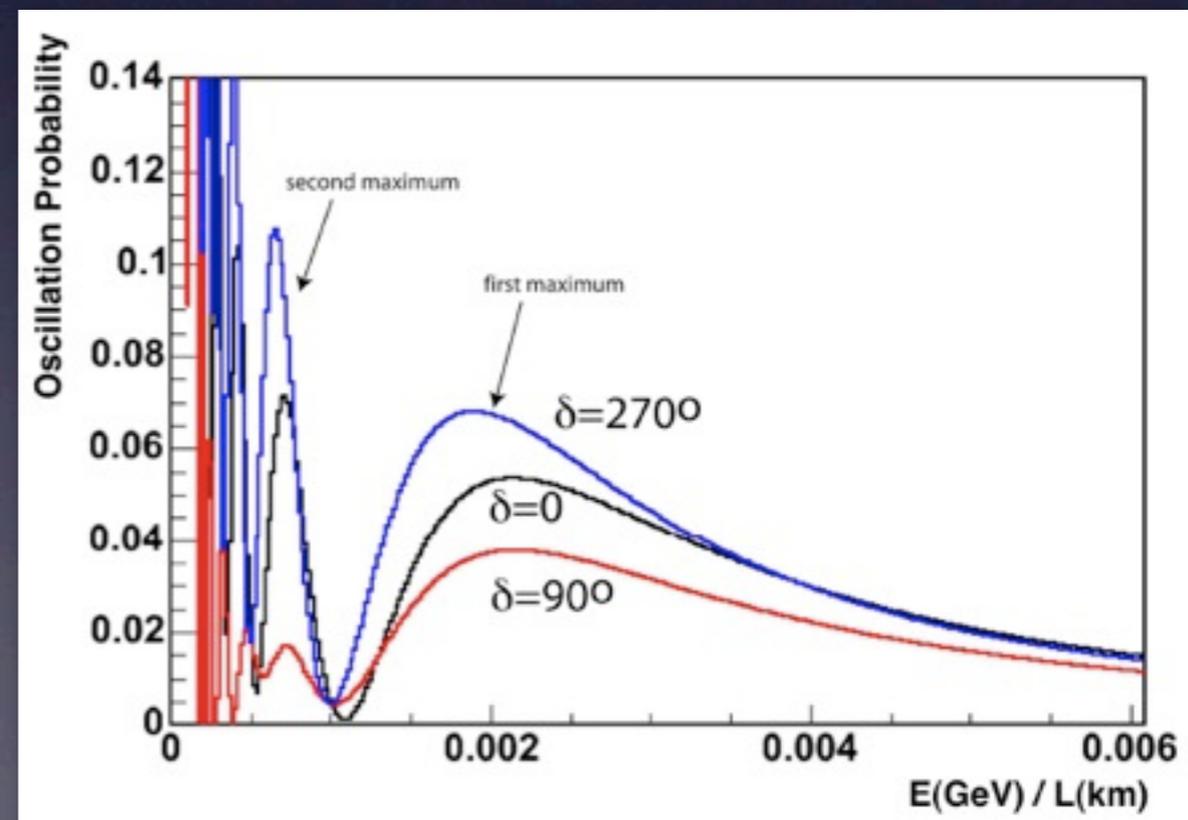
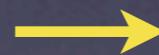
1<sup>st</sup> maximum of oscillation



# Neutrino oscillations (4)

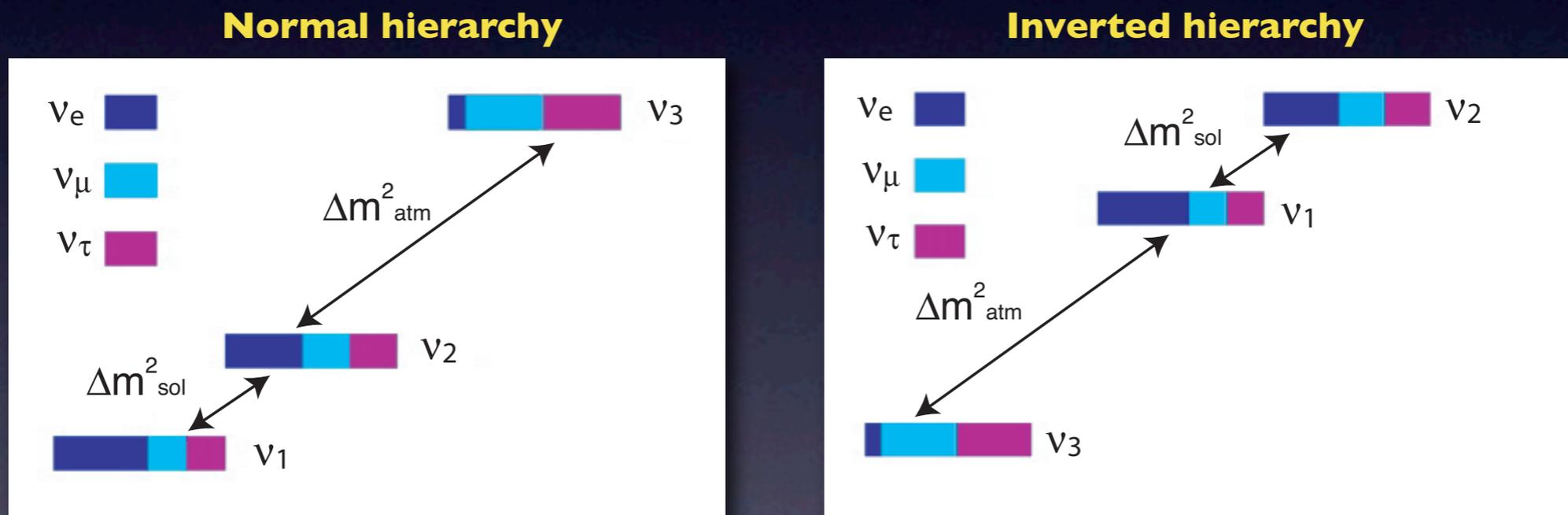
- The presence of matter affects neutrino oscillations since the mass eigenstates (i.e. the ones relevant for the propagation) are different, due to a new potential term present in the system Hamiltonian.
- The probability of oscillation of neutrinos and antineutrinos is affected in a different way: this **“fakes” a CP violation effect**.
- In order to measure intrinsic CP violation due to  $\delta_{CP}$  phase we need to disentangle the two phenomena and this can be done in 2 ways:

1. Measuring the same maximum of oscillation at different baselines.
2. Measuring two maxima of oscillation at the same baseline.



# Neutrino oscillations (5)

- The fact that mass hierarchy (i.e. Sign ( $\Delta m^2_{31}$ )) is unknown makes things harder since changing hierarchy the effects of matter on neutrino and antineutrino are swapped.



- To measure the hierarchy the best option would be to have a very long baseline (a few thousands km) so that matter effect is strong and it is easy to measure if neutrino oscillations are enhanced or reduced.

# Present knowledge

Parameter	Present knowledge (90% C.L.)	Channel	Experiments	Future
$\theta_{23}$	$\sin^2(2\theta_{23}) \geq 0.92$	$P(\nu_\mu \rightarrow \nu_\mu)$	SK, (K2K, MINOS)	T2K
$\theta_{12}$	$0.82 \leq \sin^2(2\theta_{12}) \leq 0.89$	Solar $\nu$ + $P(\text{anti } \nu_e \rightarrow \text{anti } \nu_e)$	SK, SNO, KamLAND	
$\theta_{13}$	$\sin^2(2\theta_{13}) \leq 0.19$	$P(\text{anti } \nu_e \rightarrow \text{anti } \nu_e)$ $P(\nu_\mu \rightarrow \nu_e)$	CHOOZ	T2K, Double CHOOZ Future LBL
$\Delta m_{21}^2$	$7.7 \leq \Delta m_{21}^2 / 10^{-5} \text{ eV}^2 \leq 8.3$	Solar $\nu$ + $P(\text{anti } \nu_e \rightarrow \text{anti } \nu_e)$	SK, SNO, KamLAND	
$ \Delta m_{31}^2 $	$1.9 \leq  \Delta m_{31}^2  / 10^{-3} \text{ eV}^2 \leq 3.0$	$P(\nu_\mu \rightarrow \nu_\mu)$	SK, MINOS	MINOS, T2K
Sign ( $\Delta m_{31}^2$ )	Unknown	$P(\nu_\mu \rightarrow \nu_e)$ Vs $P(\text{anti } \nu_\mu \rightarrow \text{anti } \nu_e)$		Future LBL
$\delta_{CP}$	Unknown	$P(\nu_\mu \rightarrow \nu_e)$ Vs $P(\text{anti } \nu_\mu \rightarrow \text{anti } \nu_e)$		T2K+Reactor Future LBL

# Goals

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- The goal of **long baseline** neutrino oscillation experiments is to precisely measure the relevant parameters and answer to important questions such as:
  - ✓ Is  $\theta_{23}$  mixing maximal?
  - ✓ Is  $\theta_{13}$  different from zero?
  - ✓ Is there CP violation in the leptonic sector? (i.e. is  $\delta \neq 0$ ?)
  - ✓ Is there normal or inverted hierarchy? (i.e. which is the sign of  $\Delta m_{31}^2$ ?)

# Strategy (I)

- The goal of **long baseline** neutrino oscillation experiments is to precisely measure the relevant parameters and answer to important questions such as:

$\nu_\mu$  disappearance

✓ Is  $\theta_{23}$  mixing maximal?

$$P(\nu_\mu \rightarrow \nu_x) \sim \cos^4\theta_{13} \sin^2(2\theta_{23}) \sin^2(\Delta m_{32}^2 L/(4E_\nu))$$

✓ Is  $\theta_{13}$  different from zero?

✓ Is there CP violation in the leptonic sector? (i.e. is  $\delta \neq 0$ ?)

✓ Is there normal or inverted hierarchy? (i.e. which is the sign of  $\Delta m_{31}^2$ ?)

$\nu_e$  appearance

$\nu_\mu \rightarrow \nu_e$  oscillation

# Strategy (2)

- The full 3-flavour neutrino oscillation probability for  $\nu_\mu \rightarrow \nu_e$  is given by:

$$P(\nu_\mu \rightarrow \nu_e) = \sum_{i=1,4} P_i$$

$$P_1 = \sin^2 \theta_{23} \sin^2(2\theta_{13}) \left(\frac{\Delta_{13}}{B_\pm}\right)^2 \sin^2 \frac{B_\pm L}{2}$$

$$P_2 = \cos^2 \theta_{23} \sin^2(2\theta_{12}) \left(\frac{\Delta_{12}}{A}\right)^2 \sin^2 \frac{AL}{2}$$

$$P_3 = J \cos \delta \left(\frac{\Delta_{12}}{A}\right) \left(\frac{\Delta_{13}}{B_\pm}\right) \cos \frac{\Delta_{13}L}{2} \sin \frac{AL}{2} \sin \frac{B_\pm L}{2}$$

$$P_4 = \mp J \sin \delta \left(\frac{\Delta_{12}}{A}\right) \left(\frac{\Delta_{13}}{B_\pm}\right) \sin \frac{\Delta_{13}L}{2} \sin \frac{AL}{2} \sin \frac{B_\pm L}{2}$$

Atmospheric term

Solar term

Interference terms

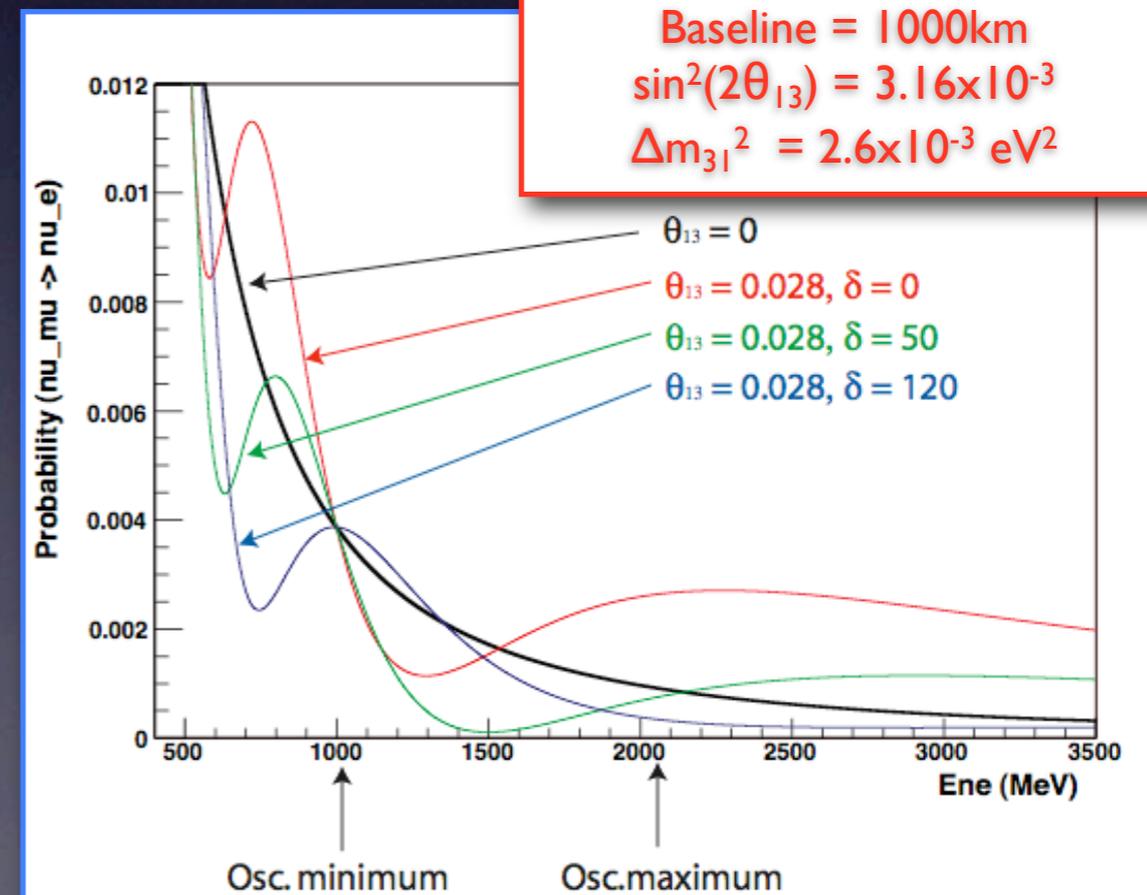
$$\Delta_{ij} = \frac{\Delta m_{ij}^2}{2E_\nu}$$

$$A = \sqrt{2}G_F n_e$$

$$B_\pm = |A \pm \Delta_{13}|$$

$$J = \cos \theta_{13} \sin(2\theta_{12}) \sin(2\theta_{13}) \sin(2\theta_{23})$$

- $\theta_{13}$  is crucial for the atmospheric part of the oscillation, and it must be proved to be non-zero.
- In case of a value of  $\theta_{13}$  different from zero, the oscillation probability depends strongly on the value of  $\delta$ .
- The so far unknown **sign of  $\Delta m_{31}^2$**  also affects the oscillation probability and mass hierarchy must be determined.



# T2K experiment

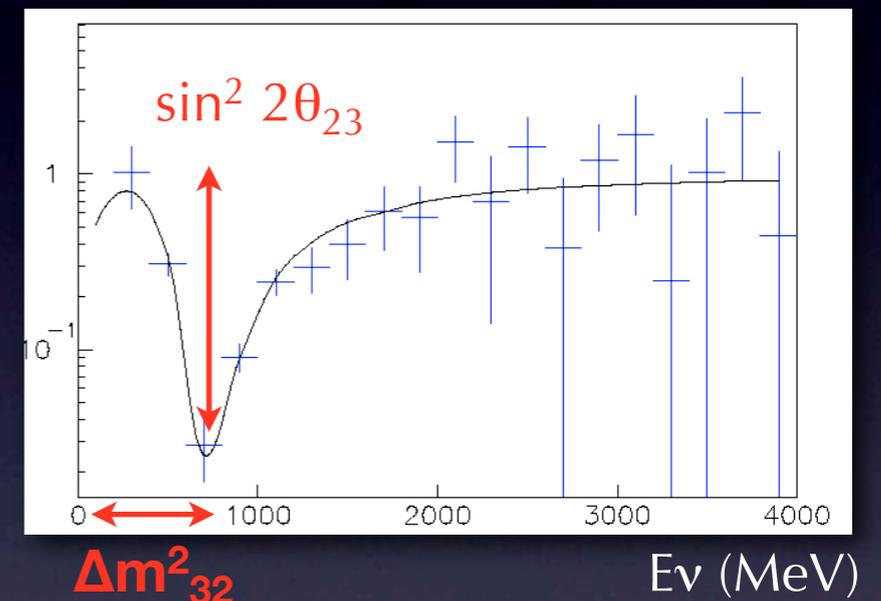
# T2K goals

Using a conventional neutrino beam (mainly  $\nu_\mu$  with a contamination of 0.4%  $\nu_e$  at peak) the goals are:

- $\nu_\mu$  disappearance:

$$P(\nu_\mu \rightarrow \nu_x) \sim \cos^4 \theta_{13} \sin^2(2\theta_{23}) \sin^2(\Delta m_{32}^2 L / (4E_\nu))$$

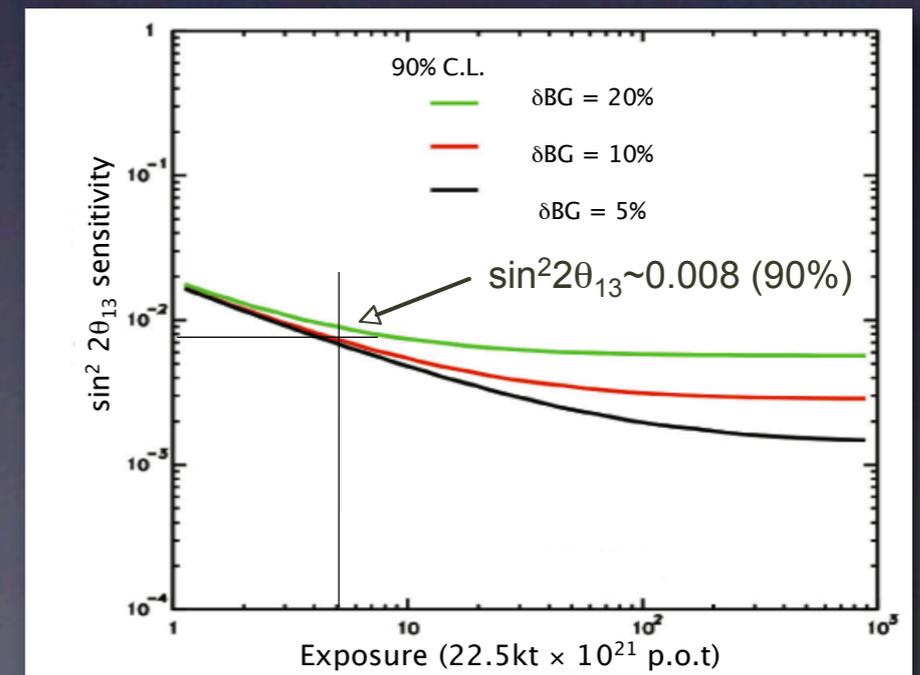
Sensitivity:  $\delta(\sin^2(2\theta_{23})) \sim 0.01$  ( $\approx 1\%$ )  
 $\delta(\Delta m_{32}^2) \leq 3 \times 10^{-5} \text{ eV}^2$  ( $\approx 1\%$ )



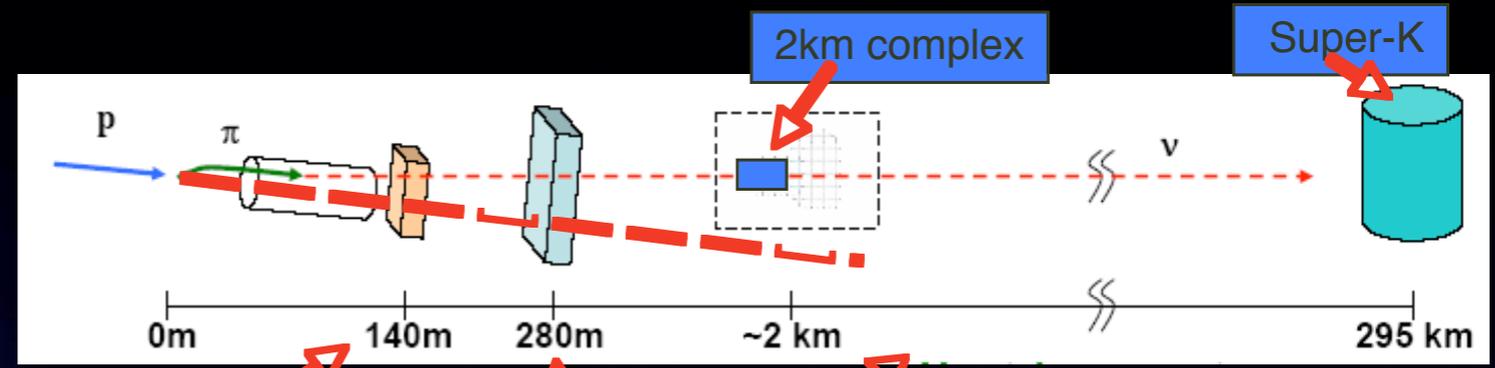
- $\nu_e$  appearance:

This measurement will allow to discover a non-zero value or set a smaller limit on the  $\theta_{13}$  mixing angle.

Sensitivity:  $\sin^2(2\theta_{13}) \leq 8 \times 10^{-3}$  (90% C.L.)



# T2K experiment

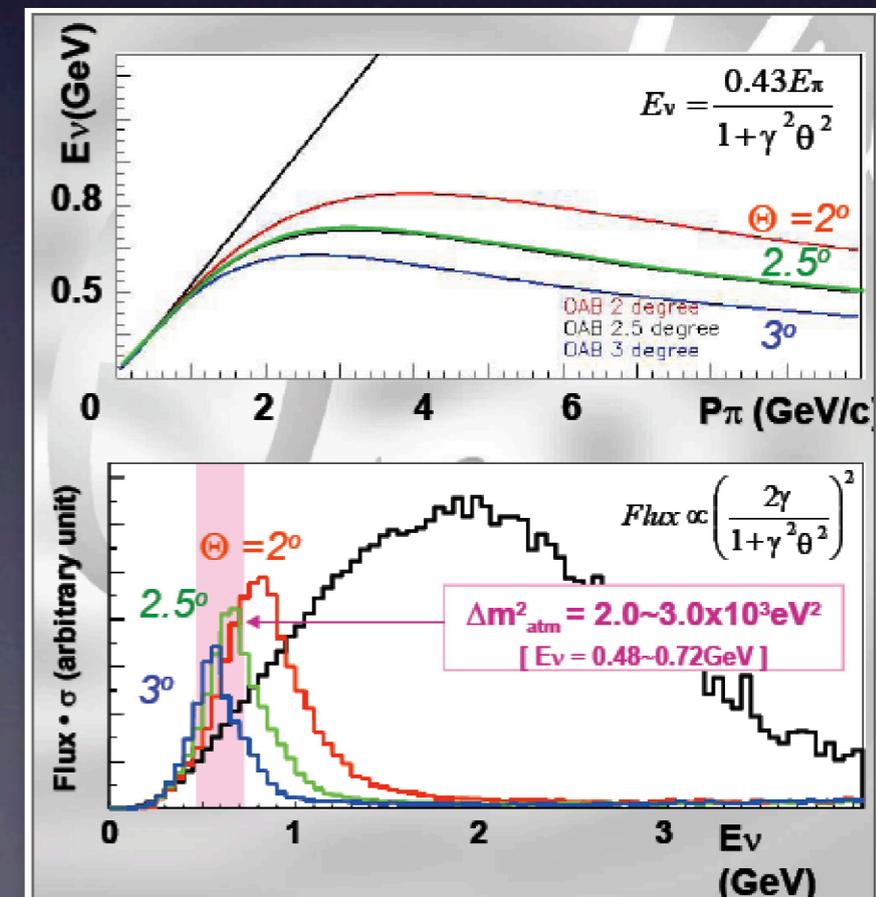


μ monitor (beam direction and intensity)

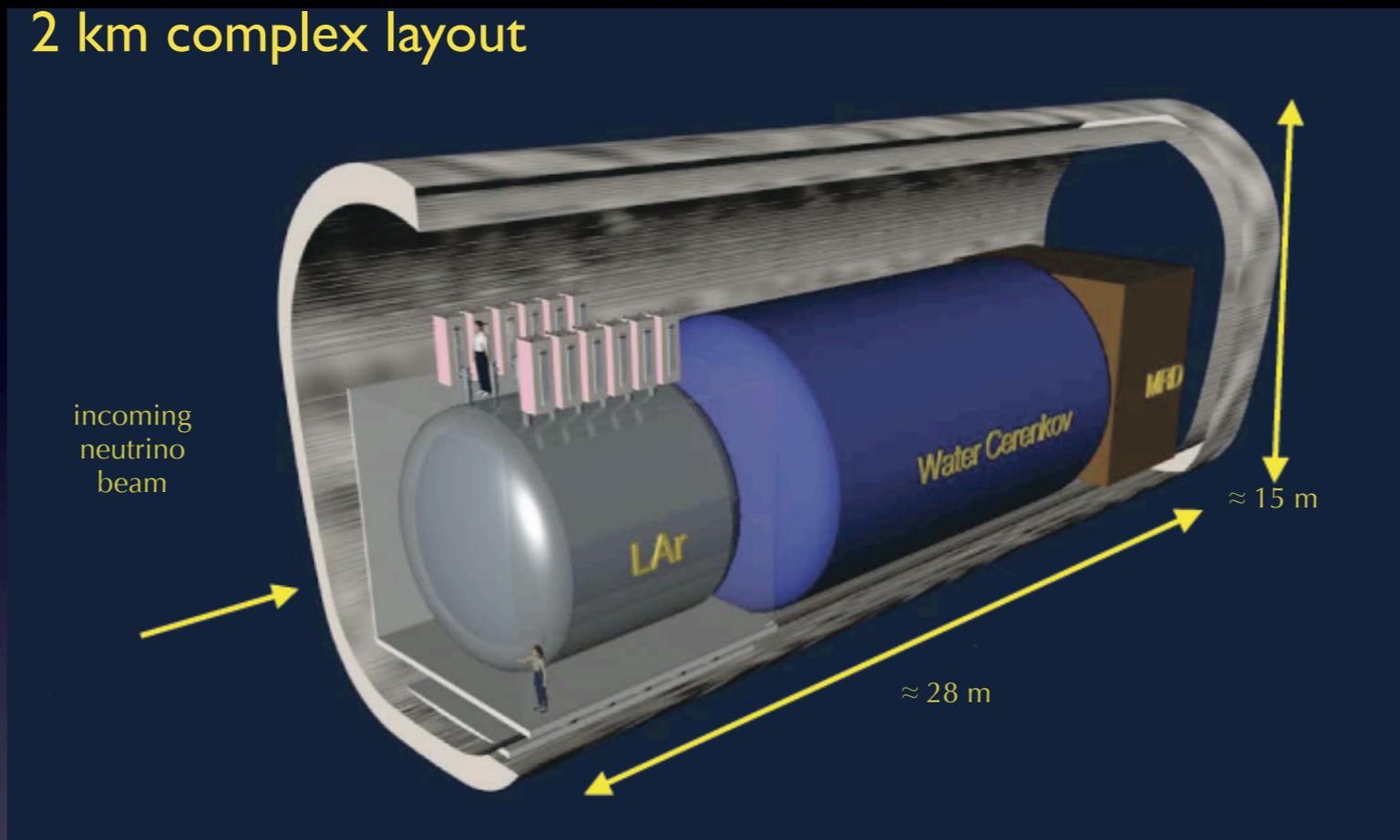
ν energy spectrum and intensity

Same spectrum as SK, BG measurement

- Long baseline (295 km) neutrino oscillation experiment.
- Low E (less than 1 GeV) Super Beam:  $\sim 10^{21}$  p.o.t./year.
- Off axis by 2.5 degrees.
- Start in 2009 aiming to reach a power of 0.75 MW from a 30 GeV proton synchrotron (upgrades to 1.6 MW and 4 MW under study).



# T2K experiment: 2 km complex

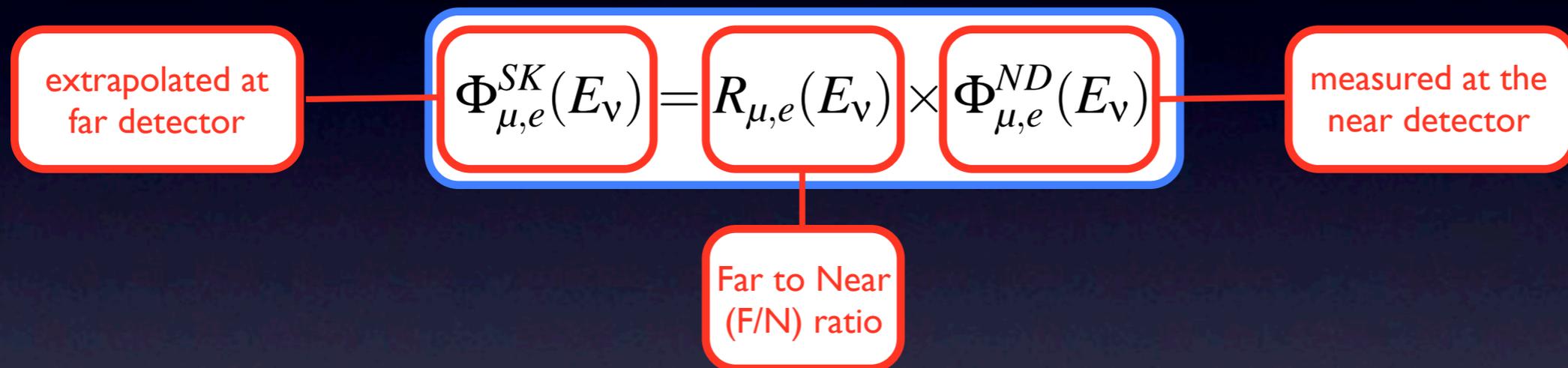


- The 2 km complex has been proposed but not approved yet, it is not foreseen before ~ 2012.
- This detector complex would be located off axis at 2km from the target.
- It is made up of a 100 ton LAr TPC (fine grained detector), a 1 kton Water Cerenkov detector and a muon range detector.

Far/Near ratio

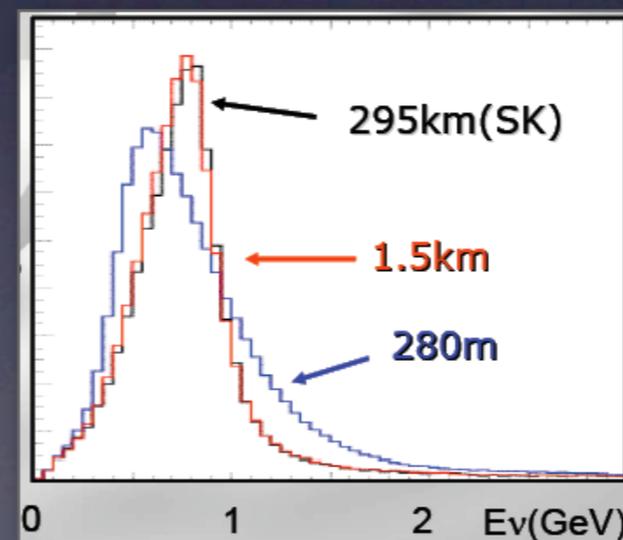
# F/N ratio

- Both appearance and disappearance analyses rely on the neutrino spectra measured at SK (far detector) and the spectra extrapolated at SK from the near detector measurement (ND280):

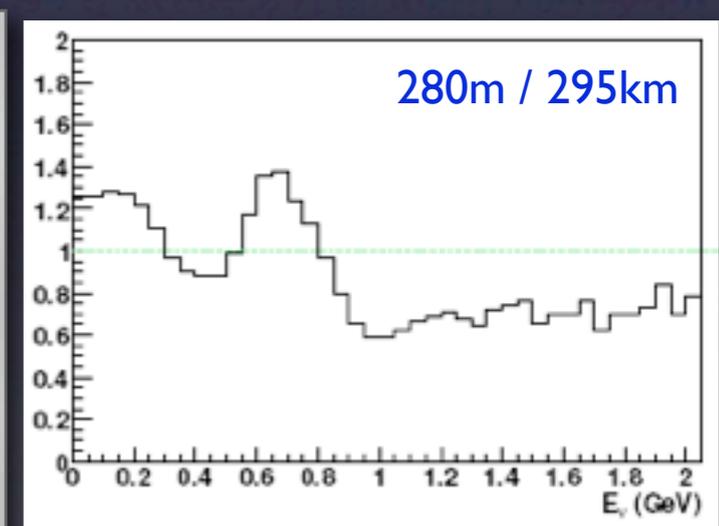


- In case of a point-like and isotropic neutrino source, the ratio is given by the solid angle (i.e. it scales as  $L^{-2}$  where  $L$  is the baseline).
- In practice, due to the finite size of the source, the F/N ratio depends on the neutrino energy.

$\nu_\mu$  Flux (normalised to SK)



F/N ratio ( $\nu_\mu$ )



# Requirements on F/N ratio

- The background in  **$\nu_e$  appearance** is mainly due to  $\nu_\mu$  NC  $\pi^0$  interactions and intrinsic beam  $\nu_e$ :

$$\begin{aligned}
 N_{BG} &= N_{BG}^{\pi^0} + N_{BG}^e \\
 &= \Phi_\mu^{SK} \cdot \sigma_{NC\pi^0} \cdot \epsilon_{SK}^{\pi^0} + \Phi_e^{SK} \cdot \sigma_e \cdot \epsilon_{SK}^e \\
 &= R_\mu \Phi_\mu^{ND} \cdot \sigma_{NC\pi^0} \cdot \epsilon_{SK}^{\pi^0} + R_e \Phi_e^{ND} \cdot \sigma_e \cdot \epsilon_{SK}^e
 \end{aligned}$$

- To achieve the goals of the experiment the systematics should be no more than 10%, therefore  **$\delta(N_{BG}) \leq 10\%$** .

- The  **$\nu_\mu$  disappearance** oscillation is measured as:

$$\begin{aligned}
 N_{sig}(E_\nu) &= P_{osc} \cdot \Phi_\mu^{SK}(E_\nu) \cdot \sigma_{\nu\mu CC}(E_\nu) \cdot \epsilon_{SK}^{\nu\mu CC}(E_\nu) \\
 &= P_{osc} \cdot R_\mu \Phi_\mu^{ND}(E_\nu) \cdot \sigma_{\nu\mu CC}(E_\nu) \cdot \epsilon_{SK}^{\nu\mu CC}(E_\nu)
 \end{aligned}$$

- MC studies on how systematics on F/N ratio affects the precision of T2K measurements showed that if  **$\delta(R_{\mu,e}) \approx 2 - 3\%$**  the contribution to the systematics is negligible compared to other contributions (due to ND280 spectrum measurements, cross sections, efficiencies, etc.) and it results into:

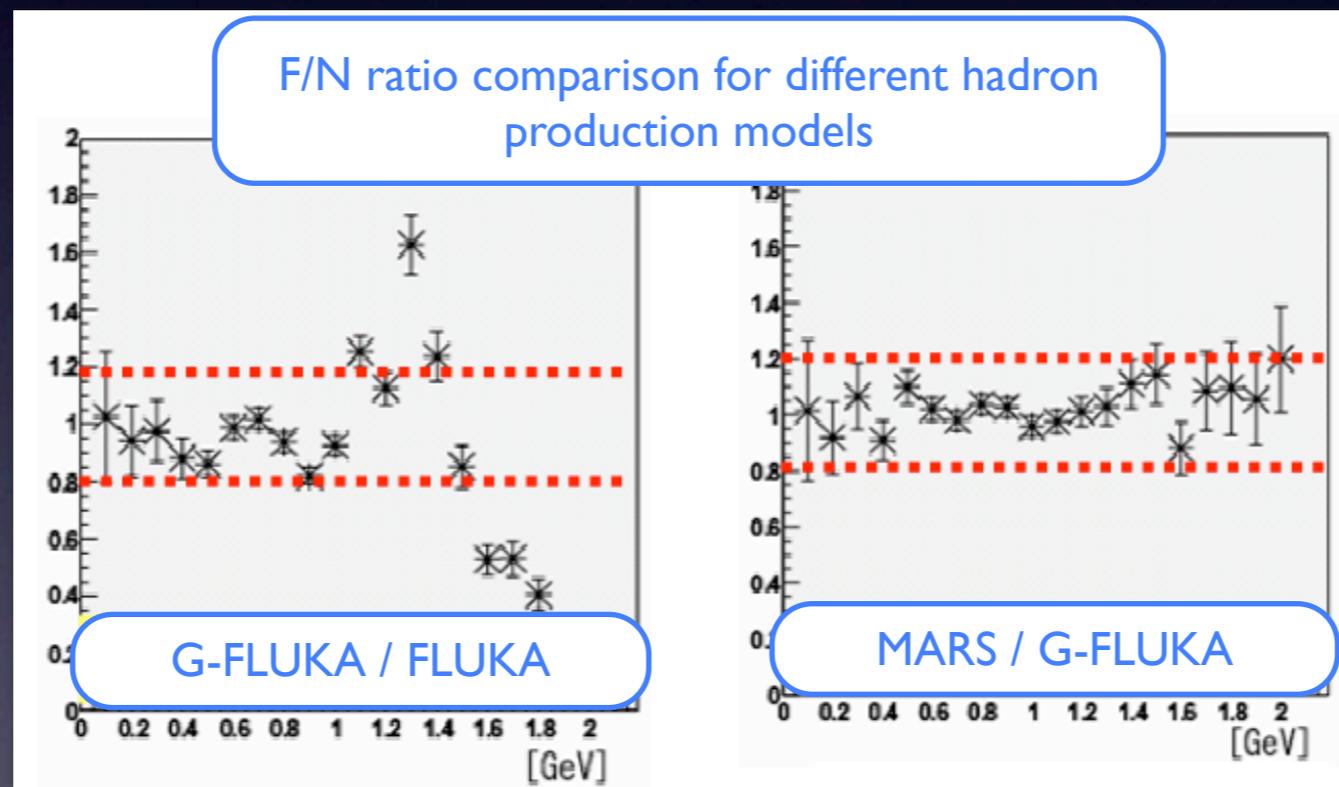
$$\delta(N_{BG}) \approx 2\%$$

$$\delta(\Delta m_{32}^2) \approx \pm 1.5 \times 10^{-5} \text{ eV}^2$$

$$\delta(\sin^2 2\theta_{23}) \approx \pm 0.005$$

# F/N ratio measurement

- No data for p+C interaction at 30 GeV exist, therefore the only way to evaluate F/N ratio is to rely on MC simulation.
- No model has been validated in this range, therefore we can assume as systematic error on F/N ratio the difference between MARS and G-FLUKA models i.e.:



- This corresponds to the following error on the N/F ratio:

$$\delta(R_{\mu}) \approx 20\%$$

# F/N ratio measurement (2)

- The error on F/N ratio of  $\approx 20\%$  corresponds to:

$$\delta(\sin^2 2\theta_{23}) \approx \pm 0.015 - 0.030$$

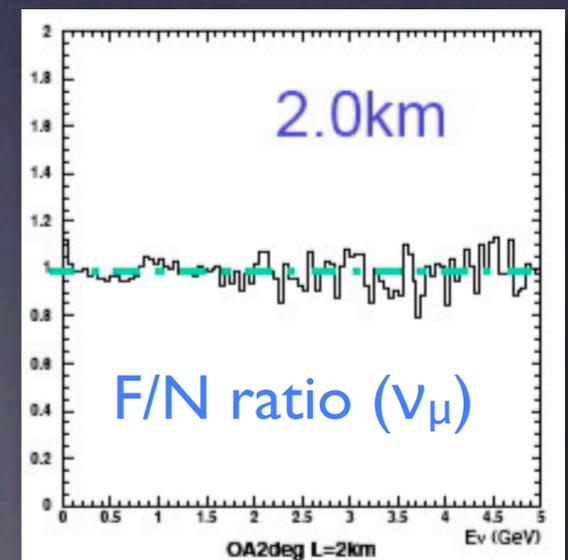
$$\delta(\Delta m^2_{32}) \approx \pm(5 - 10) \times 10^{-5} \text{ eV}^2$$

Larger than T2K goal:  
 $\delta(\sin^2 (2\theta_{23})) \sim 0.01$   
 $\delta(\Delta m^2_{32}) \leq 3 \times 10^{-5} \text{ eV}^2$

$$\delta(N_{\text{BG}}) \approx 15\%$$

Larger than the required  
error on background:  
 $\delta(N_{\text{BG}}) \leq 10\%$ .

- A measurement of the hadron (pions and kaons) production off the T2K target is therefore mandatory to reach the desired sensitivity  $\rightarrow$  **NA61/SHINE experiment**.
- Another option would be the measurement of the neutrino flux at the **2km detector complex** since at this distance the approximation of a point-like source is quite good and a flat F/N ratio within a 5% error can be achieved simply scaling for the distance.

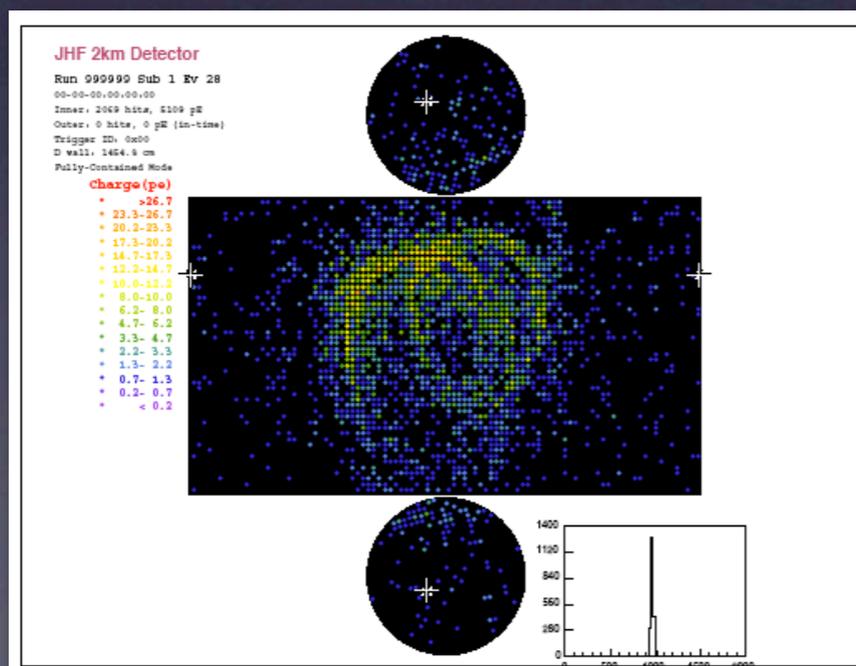


2 km complex

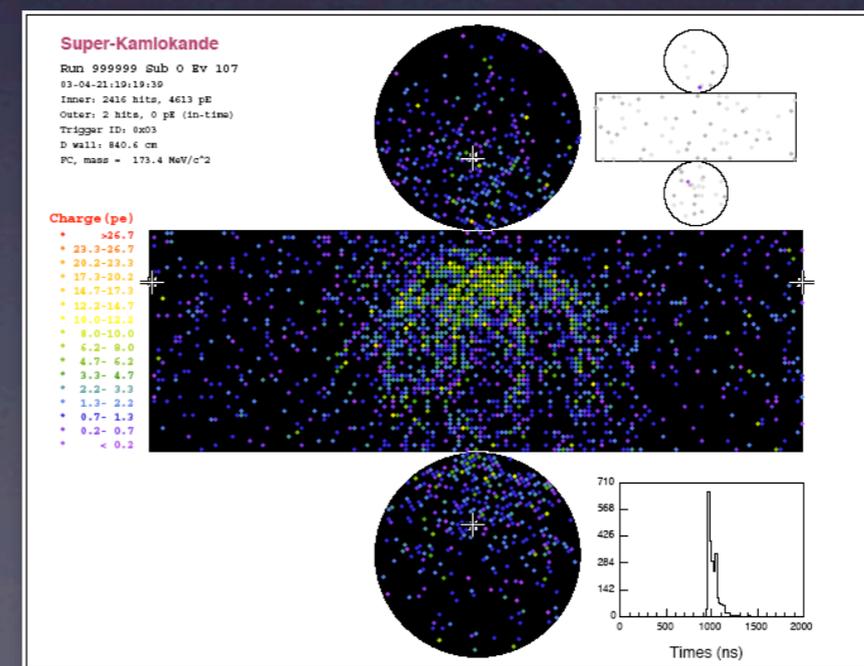
# Further motivations for a 2km complex

- As stated before the flux measured at 2 km will allow to reduce one of the the main sources of systematics: the F/N flux ratio.
- In addition, the 2 km complex would profit from a Water Cerenkov detector crucial to minimise the systematics in prediction at SK, since it has the same target but most important the same events reconstruction procedure.
- A fine grained detector (e.g. LAr TPC) is needed in order to reconstruct recoiling protons, low momentum hadrons, asymmetric decays of  $\pi^0$ , etc., in an unbiased way.

$\pi^0$  - MC 2 km



$\pi^0$  - Real event SK



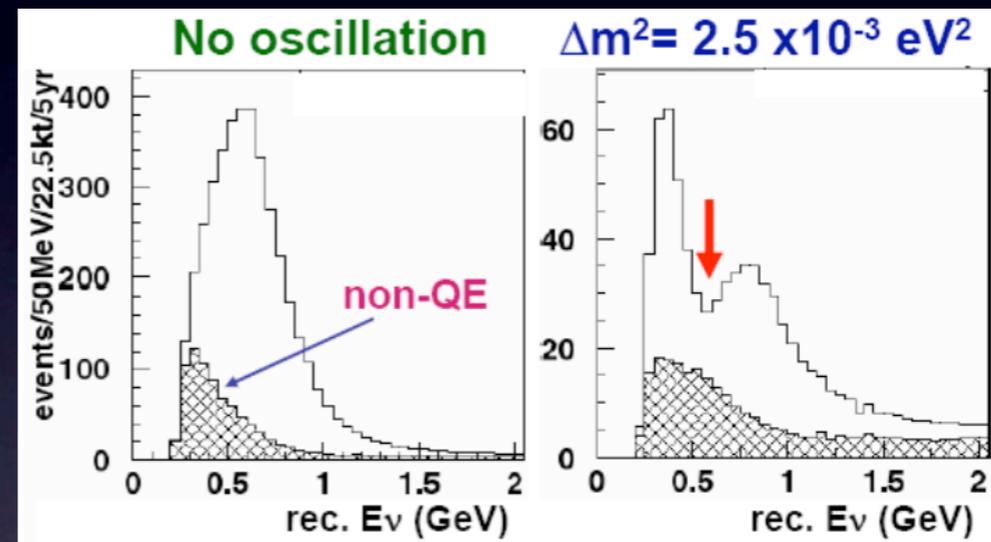
# LAr TPC physics contribution

- According to the different goals and physics scenarios, the LAr TPC will play several important roles for the final T2K measurements.

$\nu_\mu$  disappearance



Measurement of QE/nQE ratio.



$\nu_e$  appearance



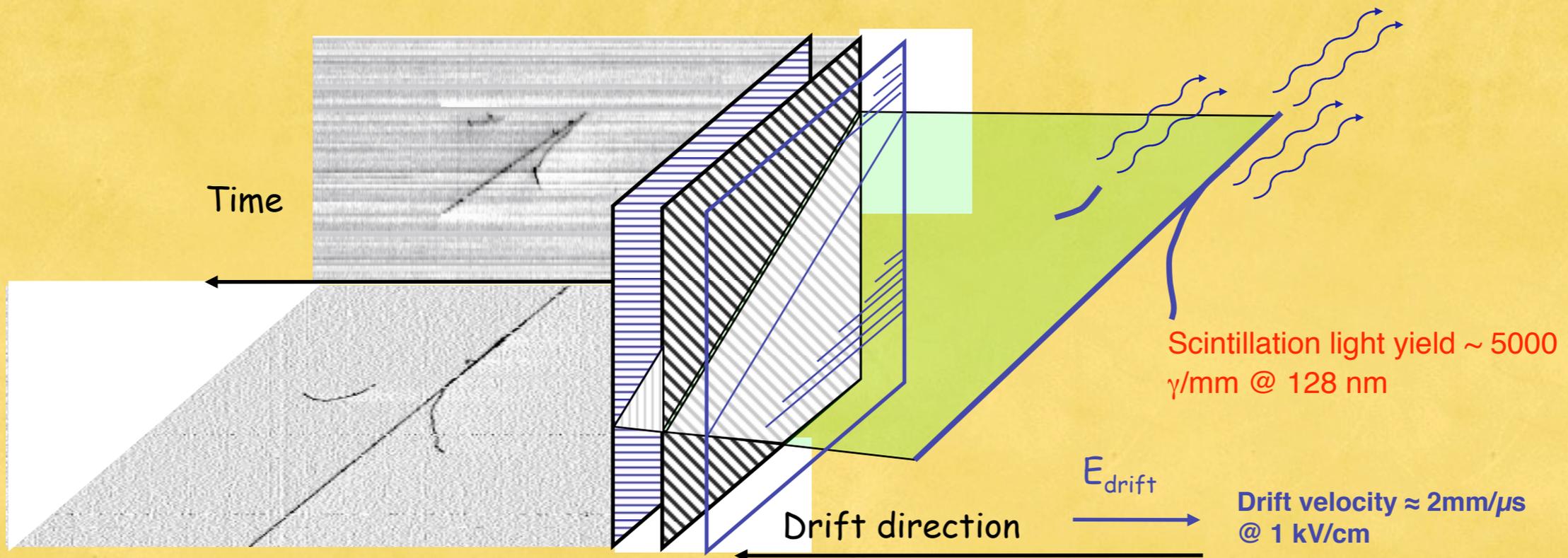
Reduction of the systematics on the different components of the background.

# The LAr TPC principle

Charge yield  $\sim 6000$  electrons/mm  
 ( $\sim 1$  fC/mm)

UV Scintillation Light: L

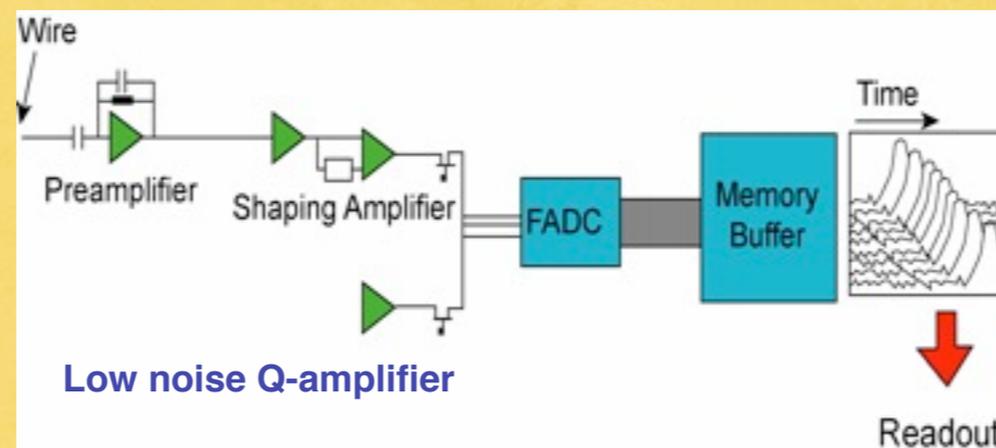
Charge readout planes: Q



Drift electron lifetime:

$$\tau \approx 300\mu\text{s} \times \frac{1\text{ppb}}{N(\text{O}_2)}$$

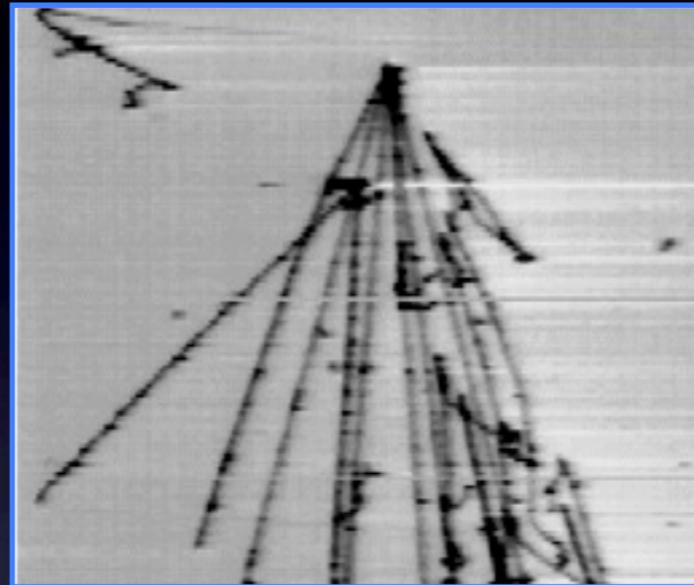
Purity  $< 0.1\text{ppb}$   $\text{O}_2$ -equiv.



# LAr TPC features (I)

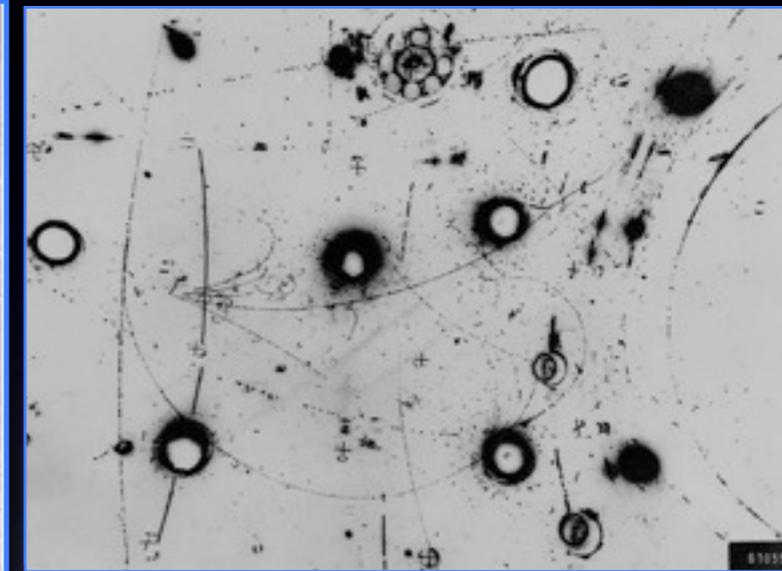
- Fully active, homogeneous, high-resolution device: **high statistics neutrino interaction studies with bubble chamber accuracy.**
- Reconstruction of low momentum hadrons (below Cerenkov threshold), especially recoiling protons.

Real event in ICARUS



High granularity: Sampling =  $0.02 X_0$   
"bubble" size  $\approx 3 \times 3 \times 0.4 \text{ mm}^3$

Gargamelle bubble chamber

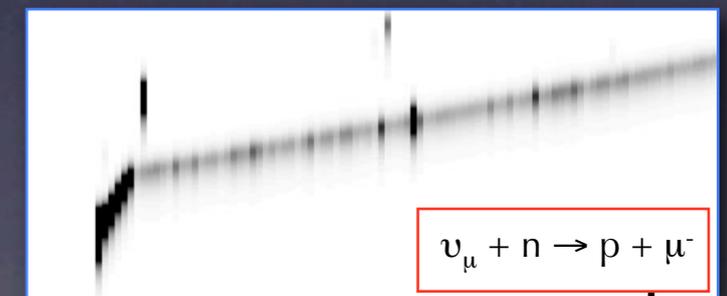


bubble diameter  $\approx 3 \text{ mm}$

## Protons

Kinetic energy T (MeV)	Momentum p (MeV/c)	Range in LAr (cm)
10	43	0.14
40	280	0.93
70	370	4.19
100	446	7.87
300	813	51.9
500	1094	116

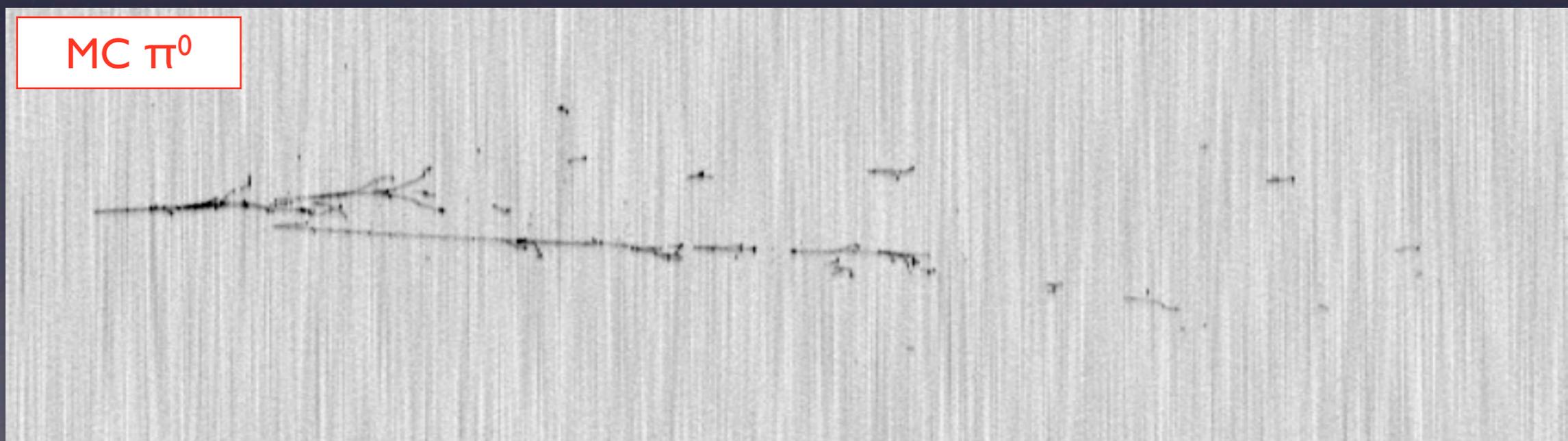
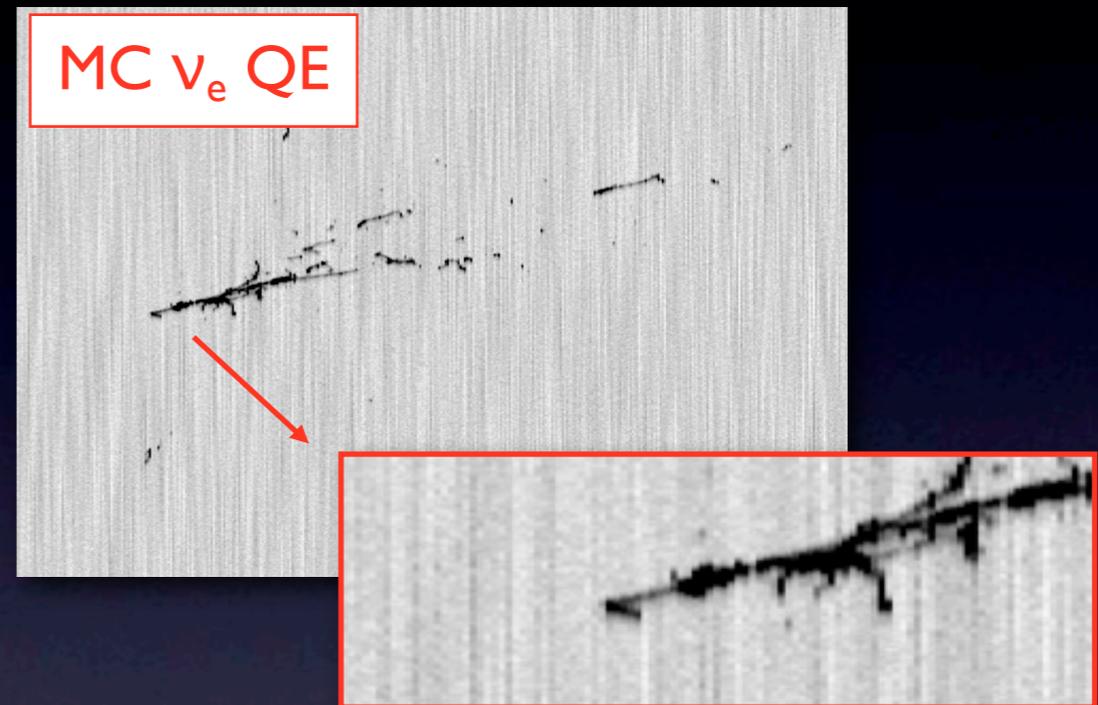
Cherenkov threshold in Water:  
 $p = 1070 \text{ MeV/c}$



MC QE event.  
Proton momentum =  $490 \text{ MeV/c}$

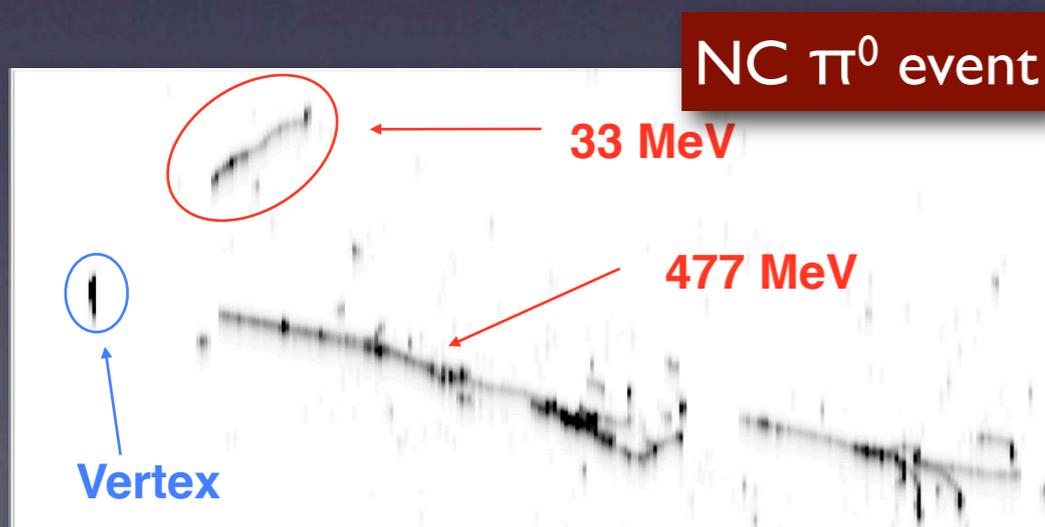
# LAr TPC features (2)

- Exclusive measurement of  $\nu\text{NC}$  events with clean  $\pi^0$  identification for an independent determination of systematic errors on the NC/CC ratio.
- Precise measurement of  $\nu_e\text{CC}$  events.



# LAr TPC performance

- Several studies have been carried out to assess the detector performance:
  - ✓ **Energy reconstruction** finding an overall r.m.s. of 20%.
  - ✓ Use of **inner target** finding that an extrapolation from Argon to Water is possible with an error of the order of 3%.
  - ✓ **QE/nQE event separation** crucial for the  $\nu_\mu$  disappearance measurement, finding a conservative error of 18%, that can be strongly reduced taking into account correctly nuclear reinteractions.
  - ✓ **NC  $\pi^0$  event selection** which is one of the relevant background for appearance measurement: 84% of the events that are misidentified in Water Cerenkov detector are correctly reconstructed.

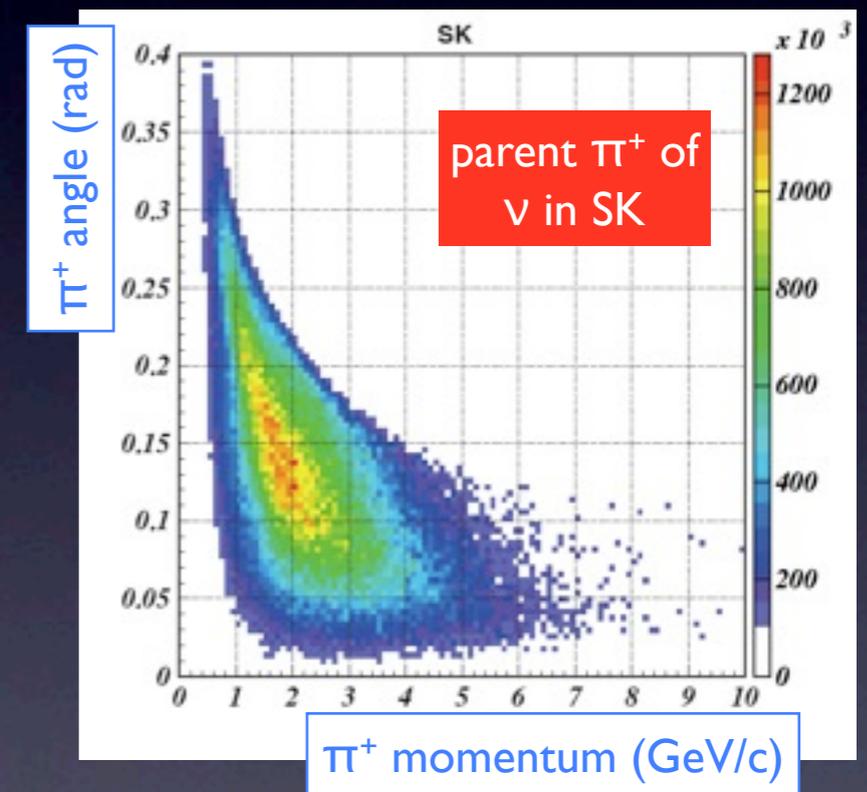
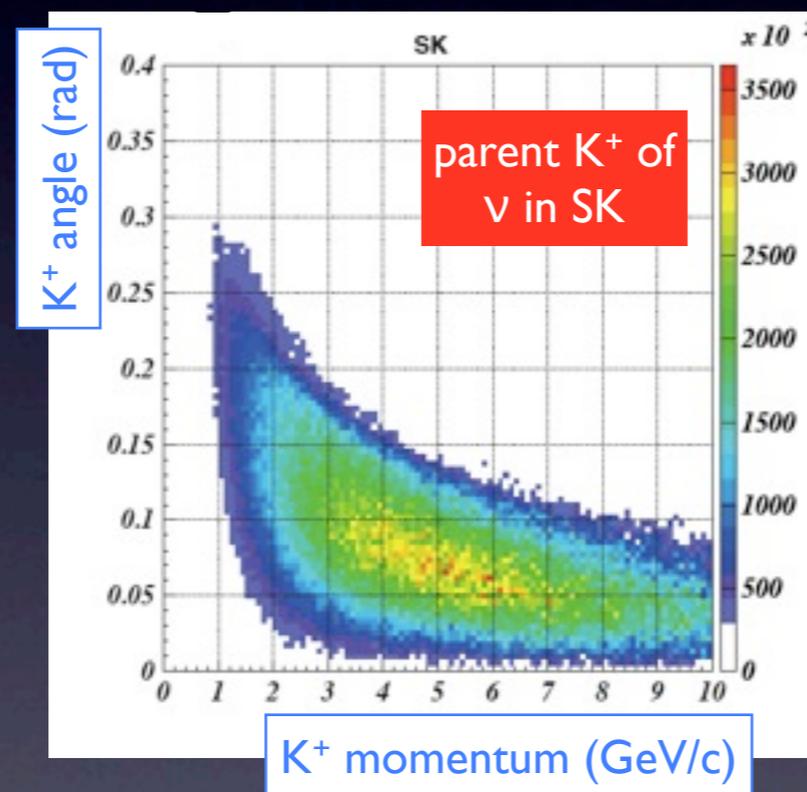


NA61/SHINE

# NA6I: goal and requirements

- The goal of NA6I experiment related to T2K is the **measurement of hadron production** from 30 GeV proton colliding on a replica of the T2K target, in order to correctly extrapolate the neutrino flux at SK from the measurement at the near detector.

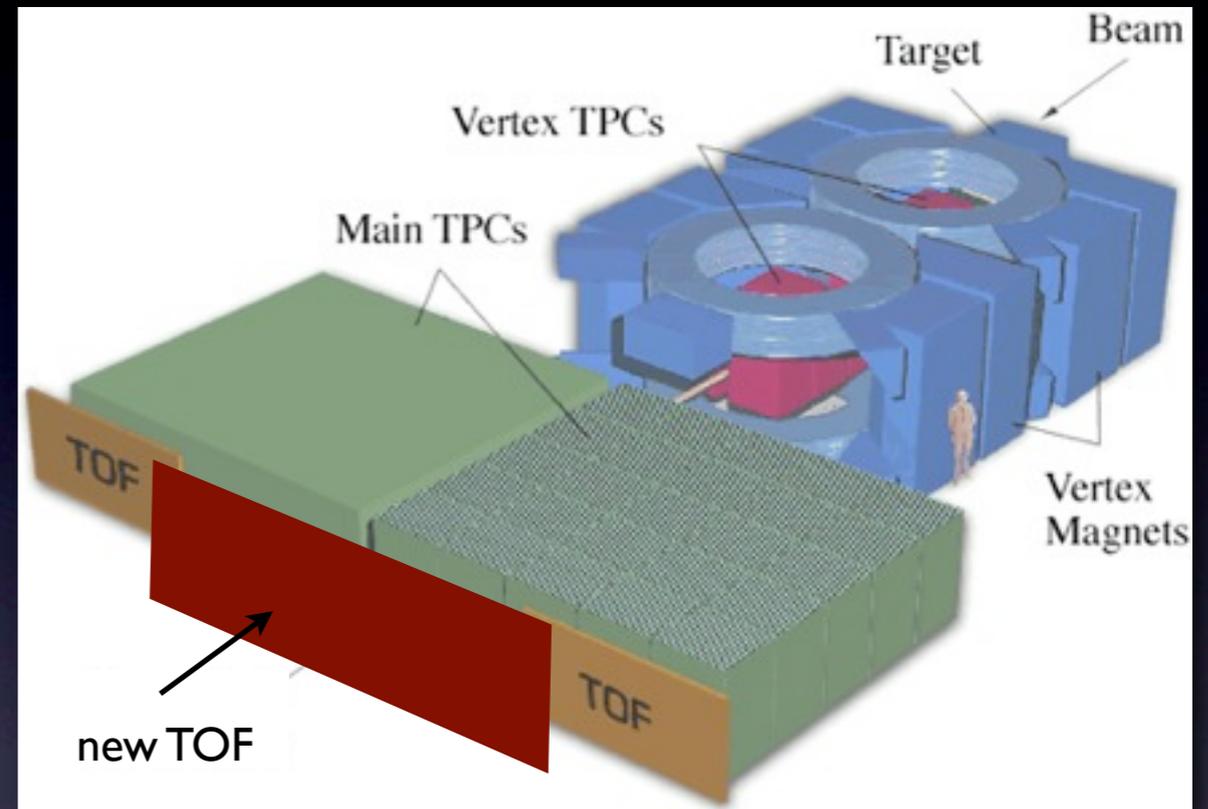
- In particular we need to measure pions and kaons whose daughter neutrinos goes through SK.



- To reach the desired precision on the Far/Near ratio of less than 3% we need to measure  **$\sim 200k \pi^+$  reconstructed tracks**.
- We also need to measure the K/ $\pi$  ratio with an uncertainty  **$\delta(K/\pi) < 10\%$** .

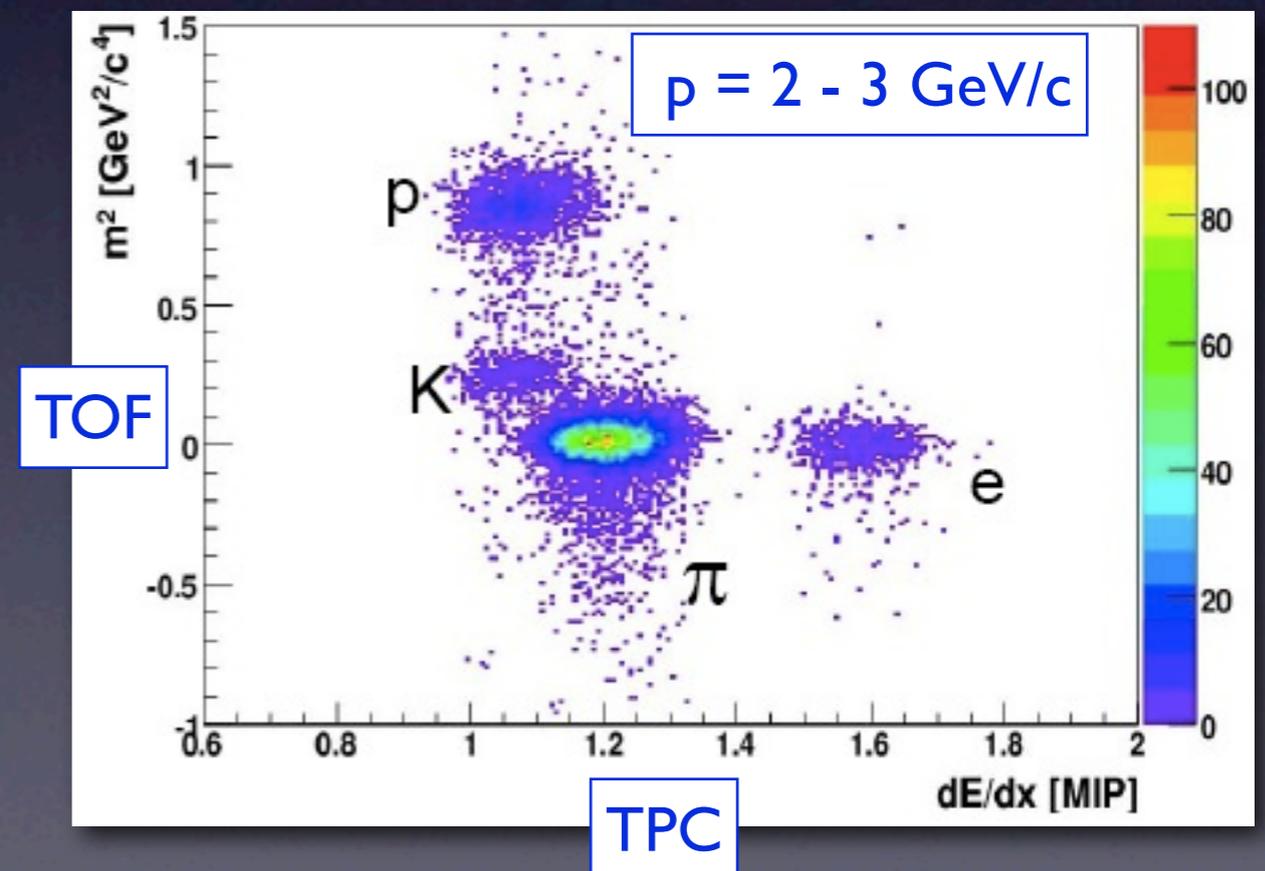
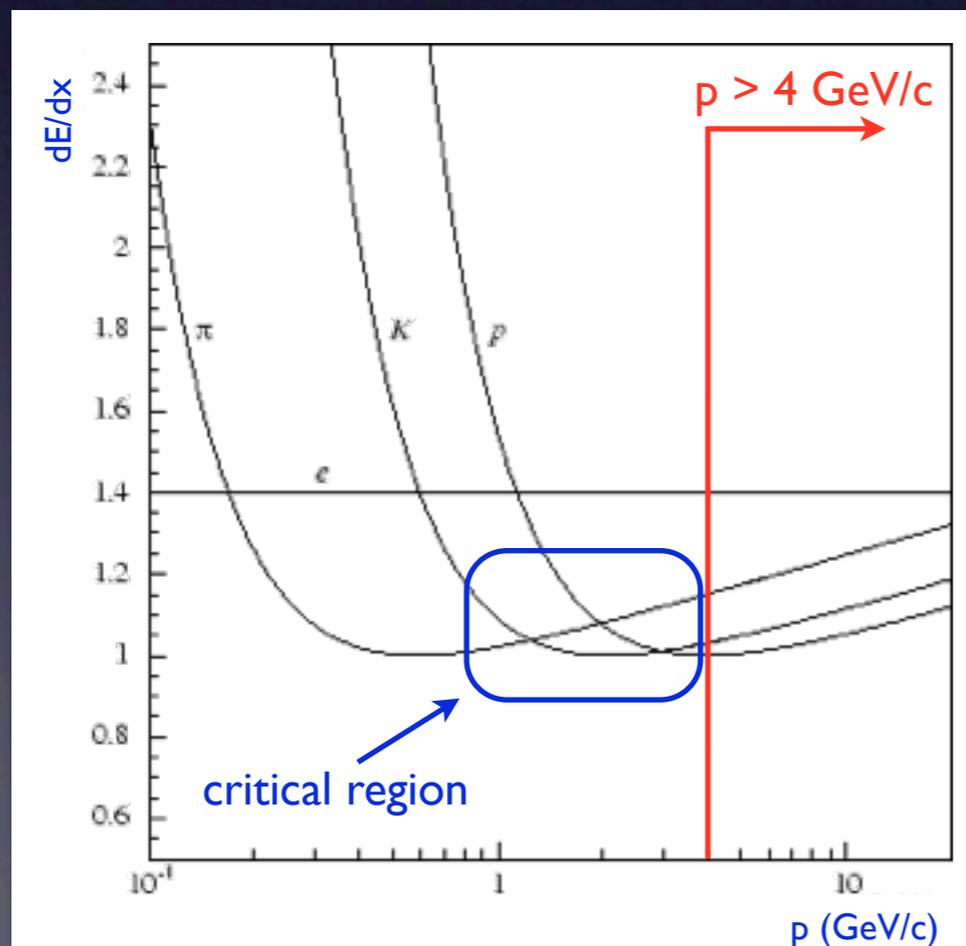
# NA61 experiment

- The NA61 experiment is a fixed target experiment located on the H2 line of SPS at CERN.
- It uses the same detector setup of the NA49 experiment (heavy ions physics) with some modifications to optimise the performance for the T2K measurements.
- The main components are:
  1. 2 dipole magnets with a bending power of about  $1.2 \text{ Tm}$  over 7 m length.
  2. 2 “vertex” TPC located inside the magnet (measurement of tracks momentum).
  3. 2 “main” TPC (measurement of  $dE/dx$ )
  4. 3 TOF , one especially built for the T2K purpose (measurement of the particle mass to be combined to  $dE/dx$  for particle ID).



# Detector sensitivity

- The momentum resolution is  $dp/p^2 \sim 10^{-4} (\text{GeV}/c)^{-1}$ .
- For tracks with momentum larger than about 4 GeV/c the particle ID is performed essentially by dE/dx.
- In the critical region of 1 - 4 GeV/c momentum the information from the TOF is crucial for a correct particle ID.

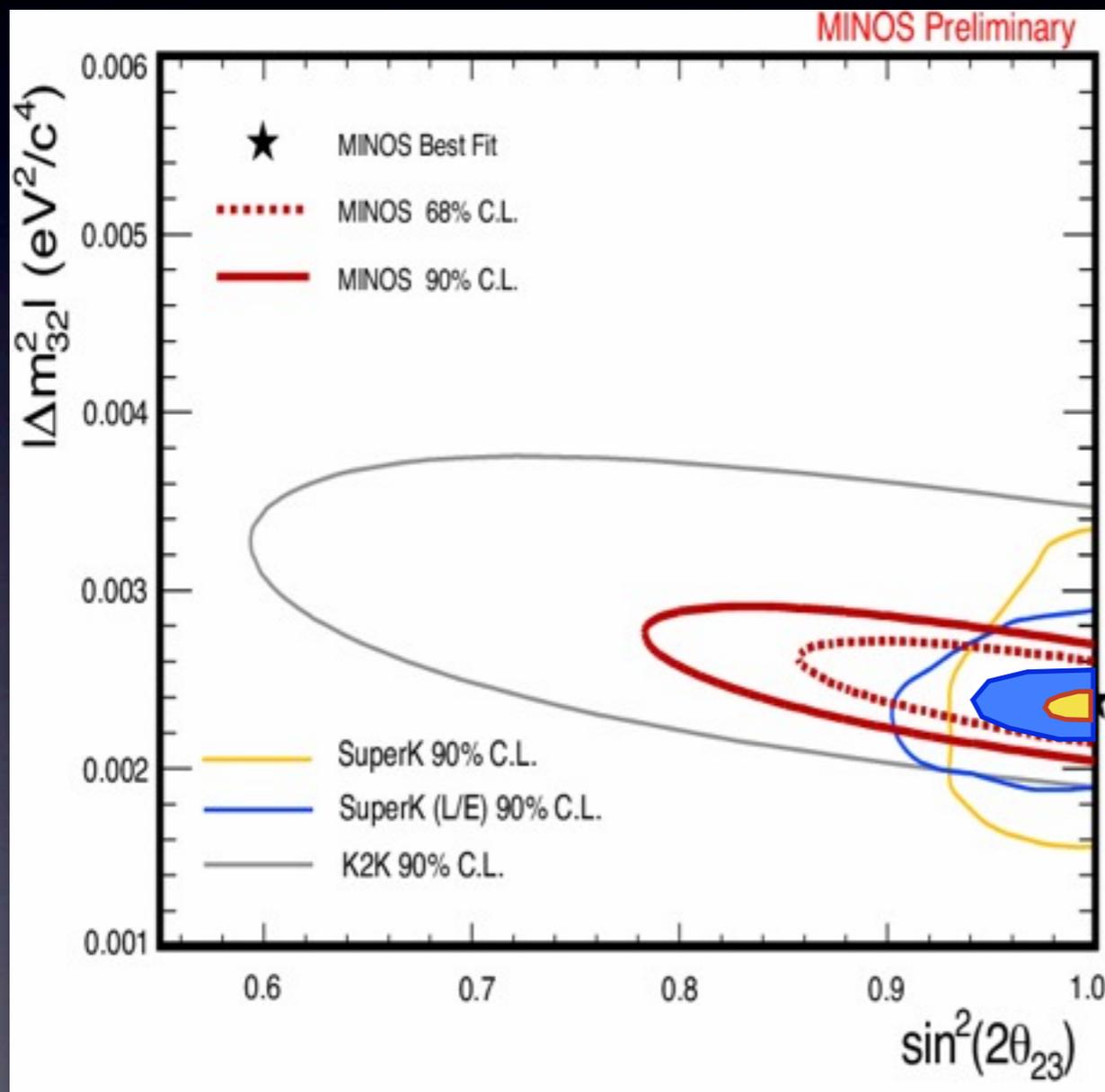


# NA6I strategy and results

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- The hadron production should be measured for **protons at 30, 40 and 50 GeV/c** momentum (as request by present and future plans of the T2K experiment) on 2 **graphite targets**:
  1. 2 cm thick graphite target ( $\sim 4\% \lambda_{\text{int}}$ ) to measure in detail the cross section for  $p+C \rightarrow \pi^++X$  and  $p+C \rightarrow K^++X$ .
  2. 90 cm long T2K replica target ( $\sim 180\% \lambda_{\text{int}}$ ) to study in detail the secondary interactions, quite critical as the beam energy increases.
- In 2007 a pilot run was taken and in 2008 full statistics was expected. However, due to the LHC accident, the goal was not achieved and postponed to 2009.
- In 2008 run we registered enough trigger of protons and pions at 30 GeV and 75 GeV to confirm the preliminary results from 2007 run and complete the proton-carbon cross section measurement.
- Analysis is almost finished, results expected to be made public by summer 2009.

# Visual impact of NA6 I



**T2K without  
NA6 I**

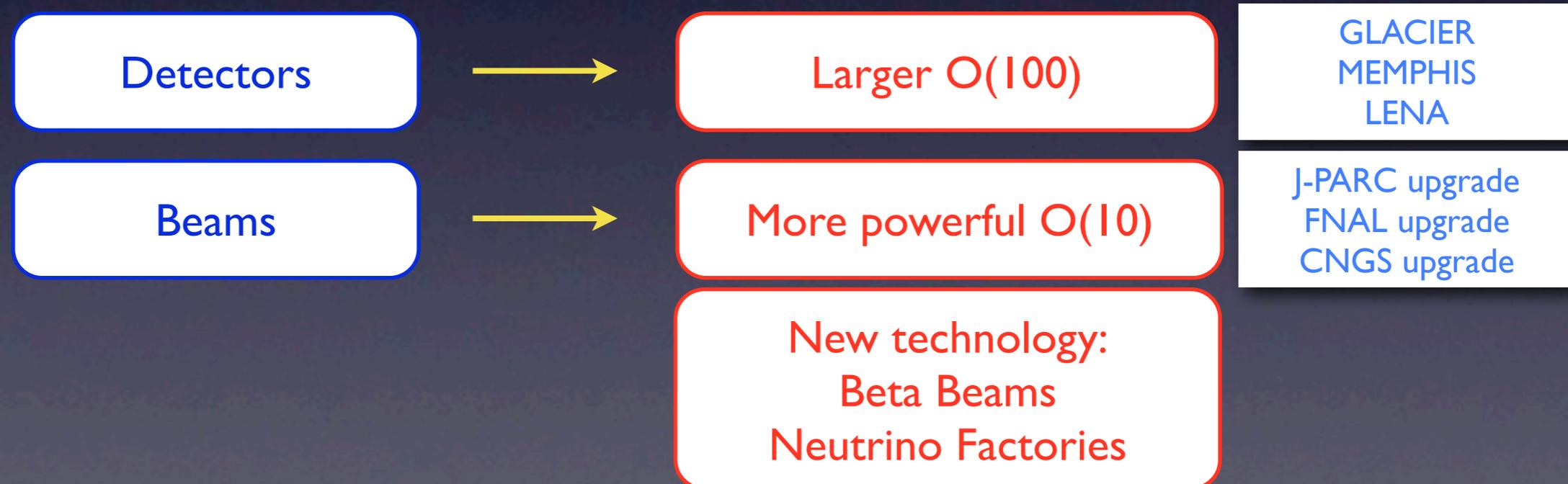


**T2K with  
NA6 I**

# Future long baseline neutrino experiments

# Perspectives

- Thanks to NA61 experiment and potentially with the 2 km detector complex, the T2K experiment will reach a limit on  $\sin^2(2\theta_{13}) \leq 8 \times 10^{-3}$ .
- T2K results will be crucial to determine the future of neutrino physics: in case a signal is measured, it would open the search for **CP violation** in the leptonic sector.
- What will the future be?



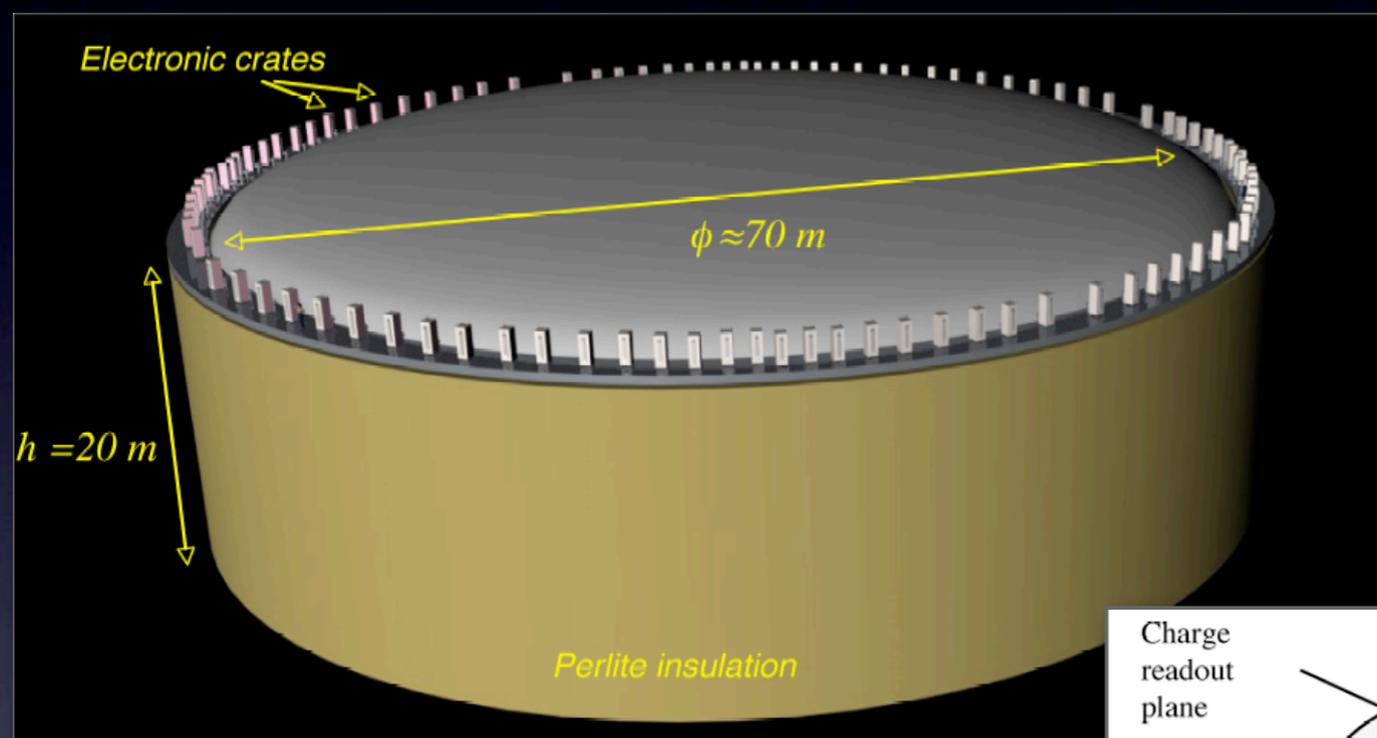
# Perspectives (2)

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- Given the amount of money required for the next generation neutrino long baseline oscillation experiments, people try to find a “common solution” (ISS, LAGUNA, BENE, EUROν).
- Do we really need expensive new beam technologies? The answer lies in the true value of  $\theta_{13}$ .
- In any case a large detector will profit from data taking on a Super Beam while deciding if new technologies are really worth it.
- Several studies have been carried out on the performance of a large (100 kton) LAr TPC on possible upgrades of J-PARC and CNGS neutrino beams.
- I will shortly discuss the general feature of the result found on the J-PARC beam only ([arXiv:0801.4035](#)), since the conclusion for the CNGS beam ([JHEP 0611:032,2006](#)) are quite similar.

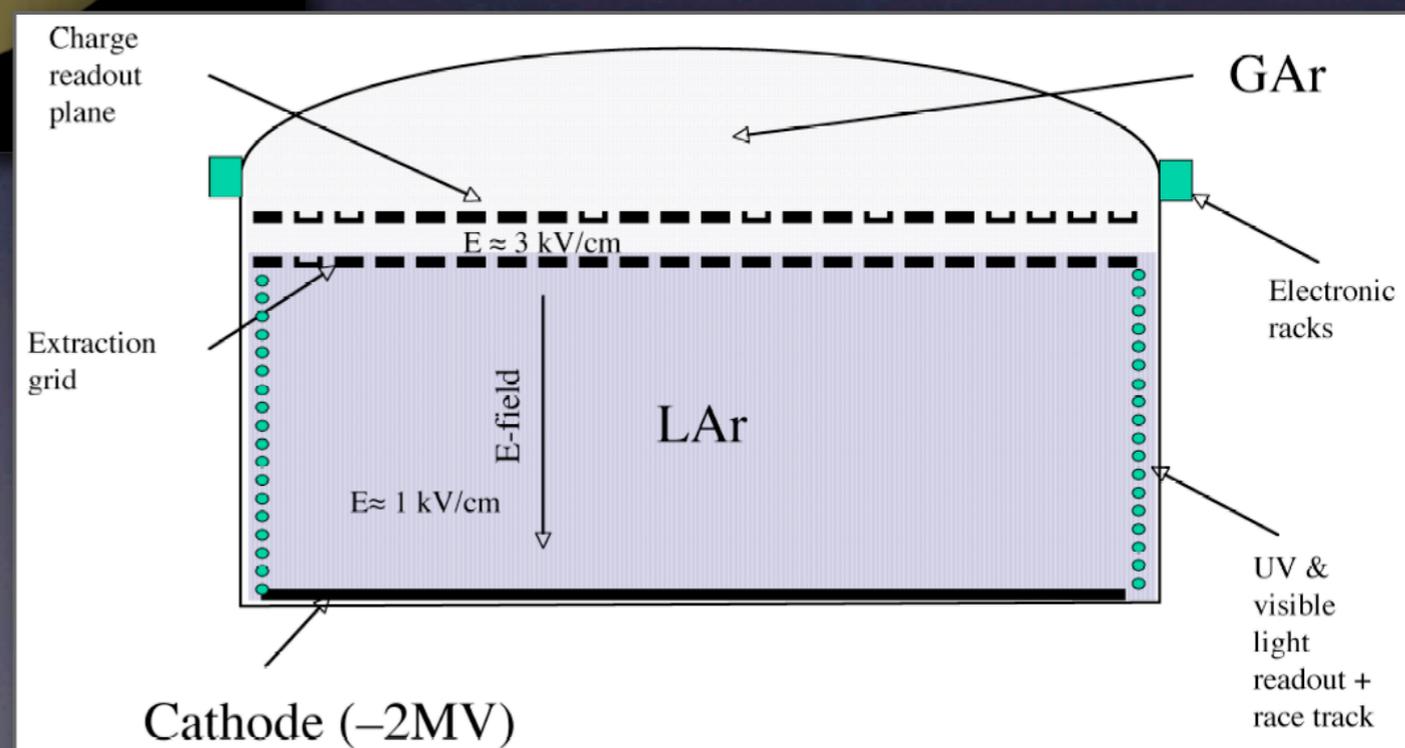
# GLACIER detector

- GLACIER (Giant Liquid Argon Charge Imaging Experiment) is a 100 kton LAr TPC.



- The drift is 20 m long and a new readout other than wires is needed.
- It will be a double phase detector: electrons are drifted in liquid and multiplied in gas.

- New HV supply needed since it is not possible to bring in 2MV with standard feed-through method.
- R&D already ongoing with ArDM, ARGONTUBE, T2K.



# J-PARC neutrino beam

- The same neutrino beam used for the T2K experiment can be measured at different locations:

different off axis angles  
(different energy spectra)



different baselines



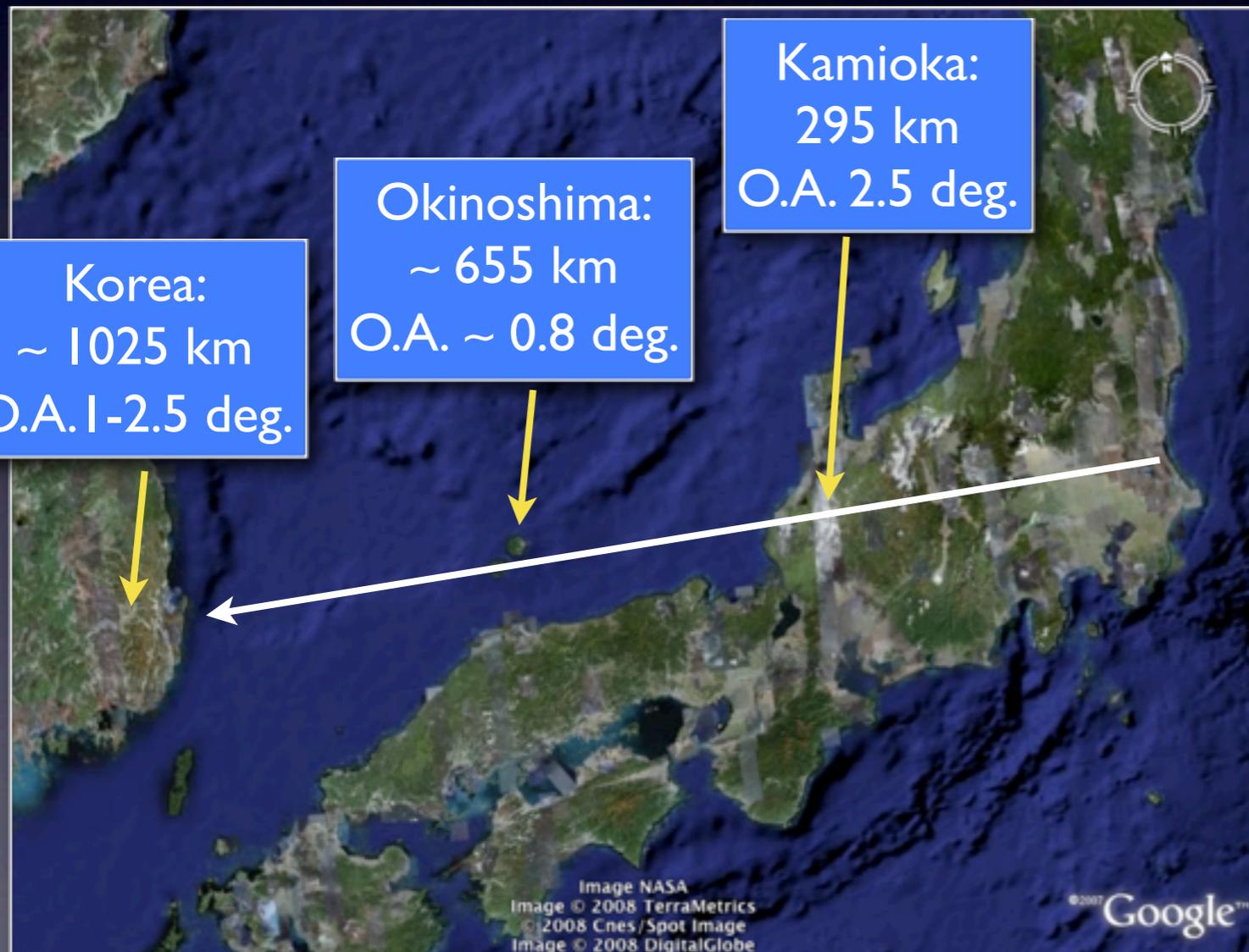
higher maxima of oscillation

different contribution from matter effects



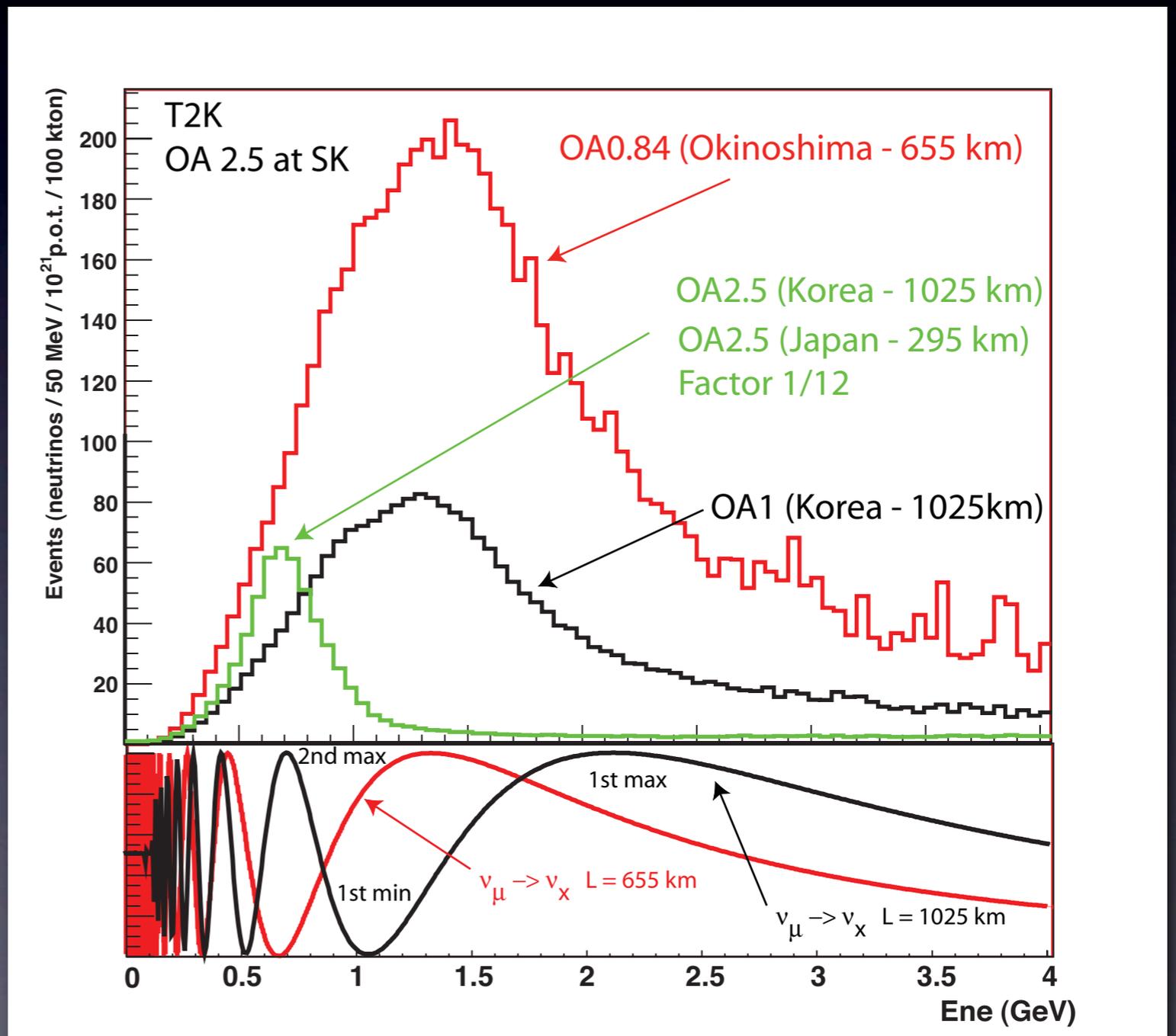
Measurement of **CP violation**  
and **mass hierarchy**

- We located GLACIER in 4 different positions (3 baselines).
- We assumed a beam power upgrade of 4 MW.
- 5 years of neutrino beam and 5 years of antineutrino beam have been assumed in the analysis.



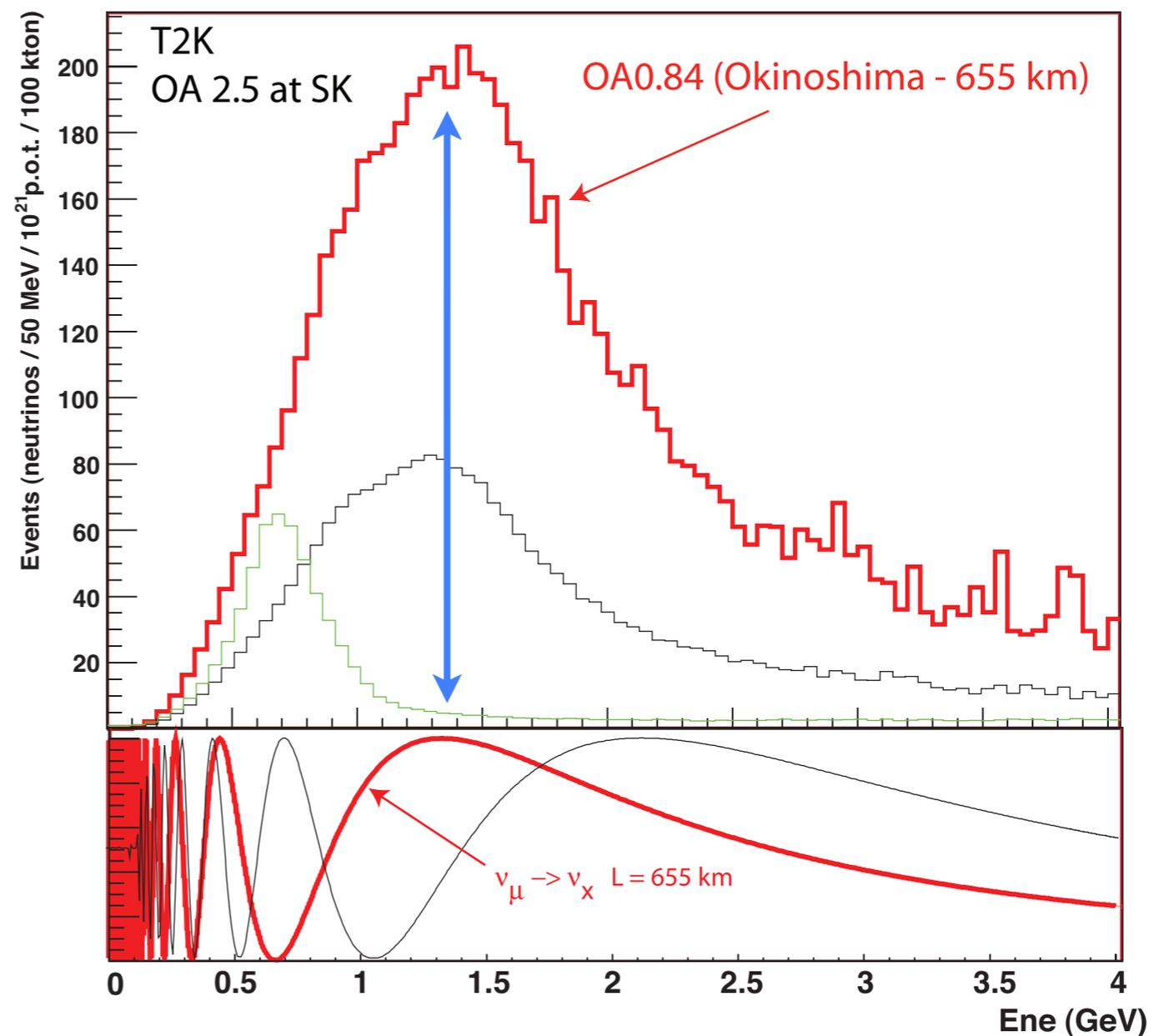
# Spectra

- We calculated the spectra of  $\nu_\mu$  CC at the various locations.



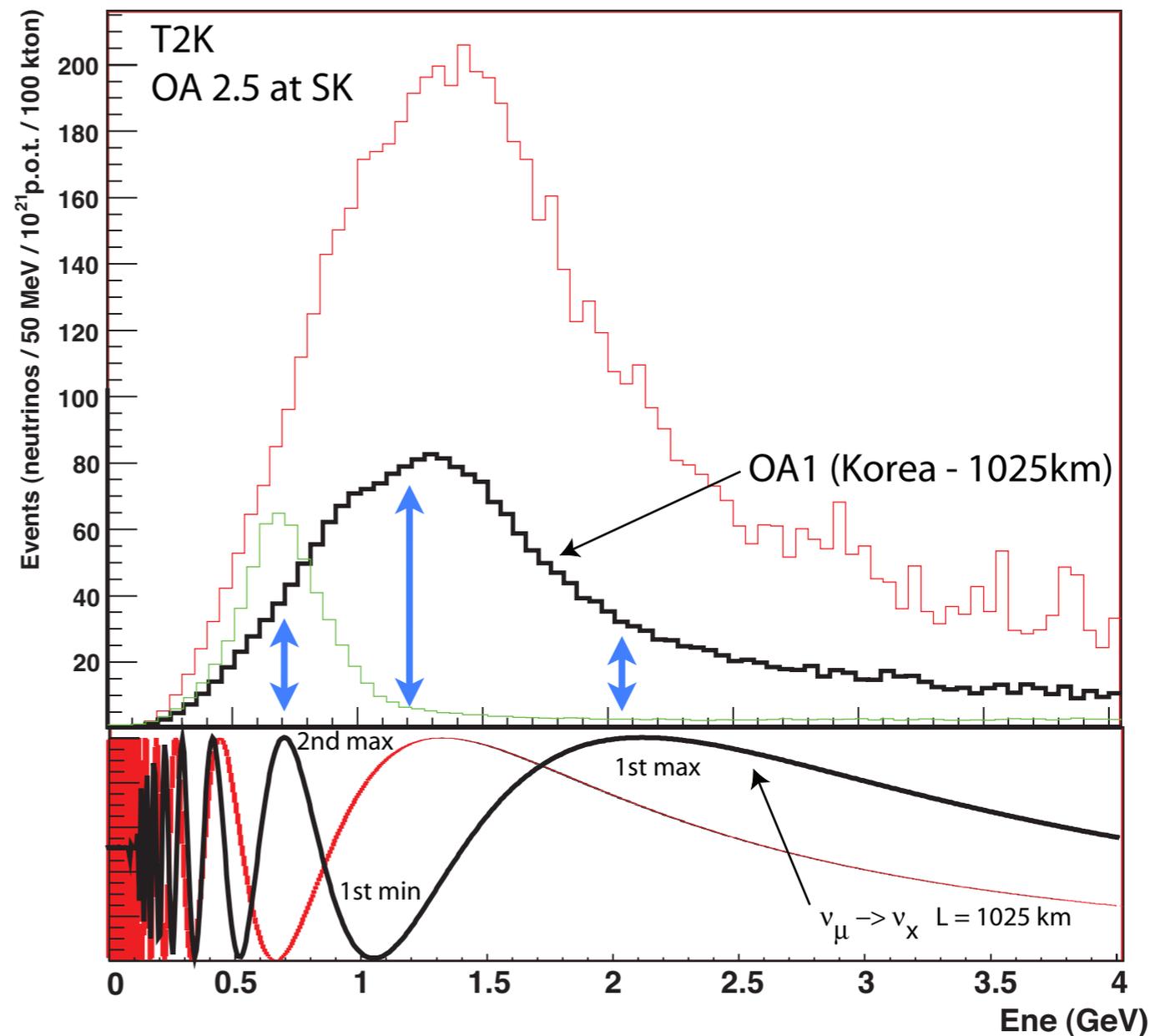
# Spectra

- At Okinoshima, as well as at Kamioka (as we saw for T2K experiment), the peak corresponds to the first maximum of oscillation.



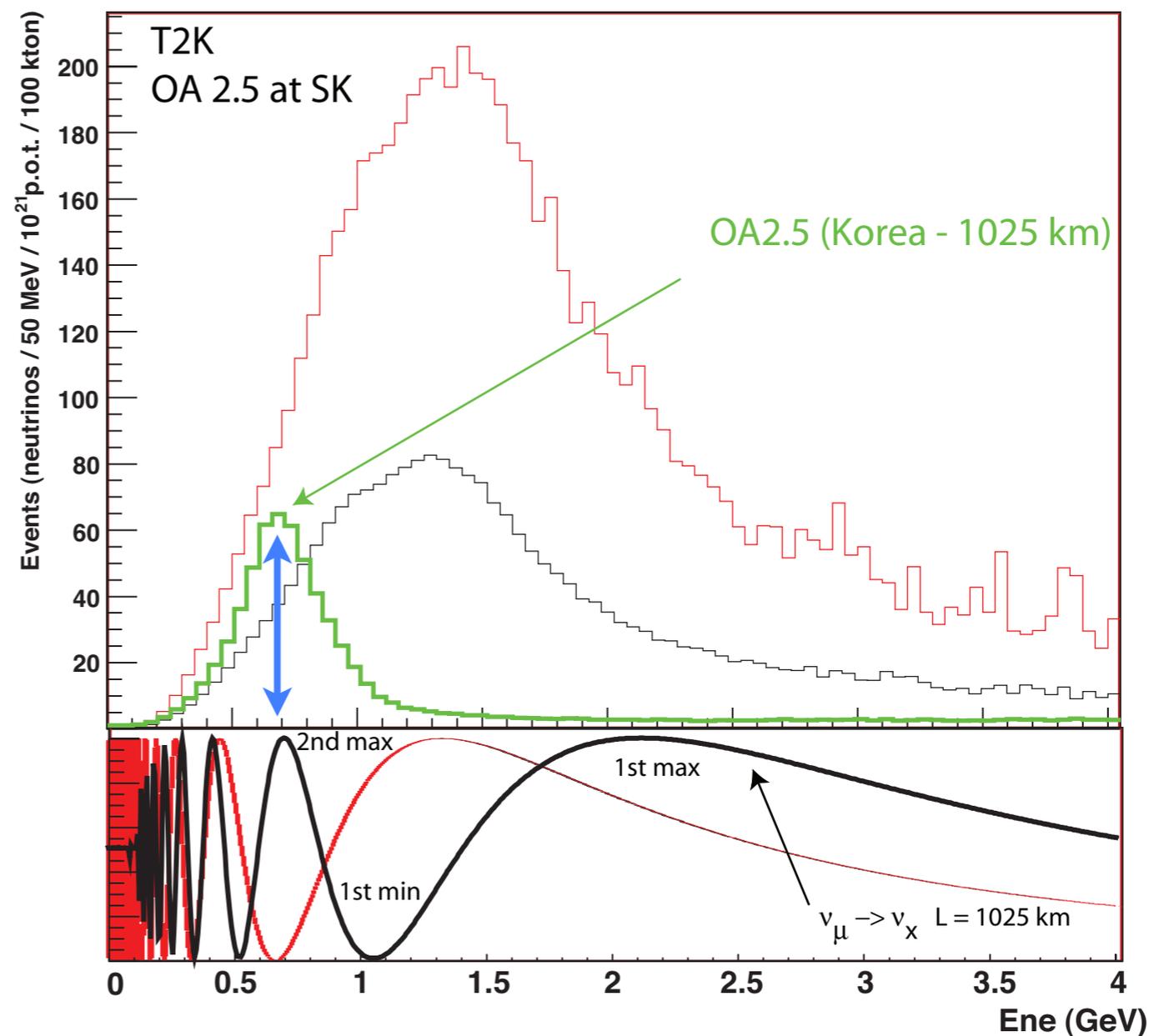
# Spectra

- For the long baseline of 1025 km in **Korea** at **OA 1 deg.** the spectra is almost peaked at the **first minimum** and covers both the **first and second maxima** of oscillations.



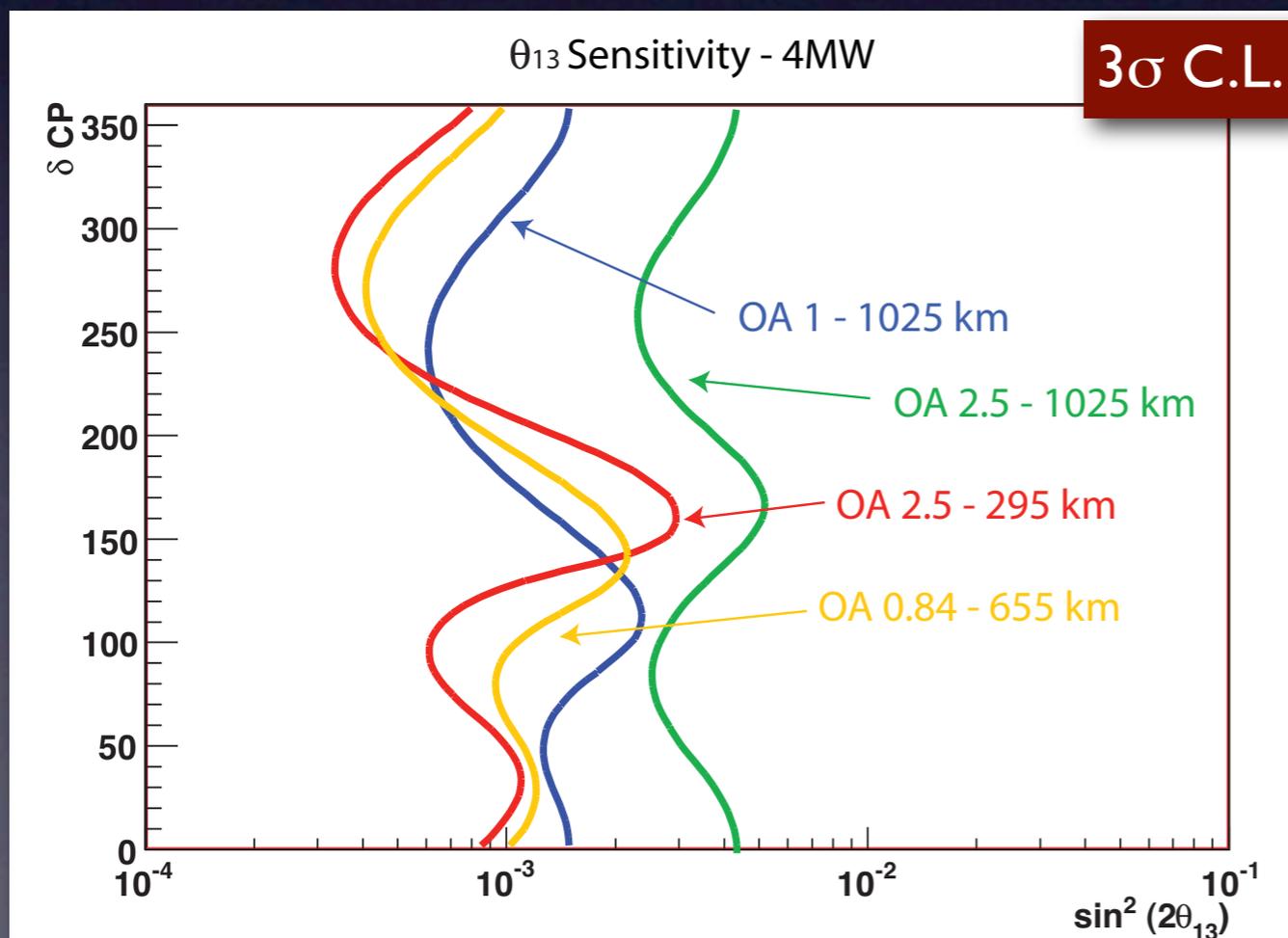
# Spectra

- For the long baseline of 1025 km in **Korea** at **OA 2.5 deg.** the spectra is peaked at **the second maximum** of oscillations.



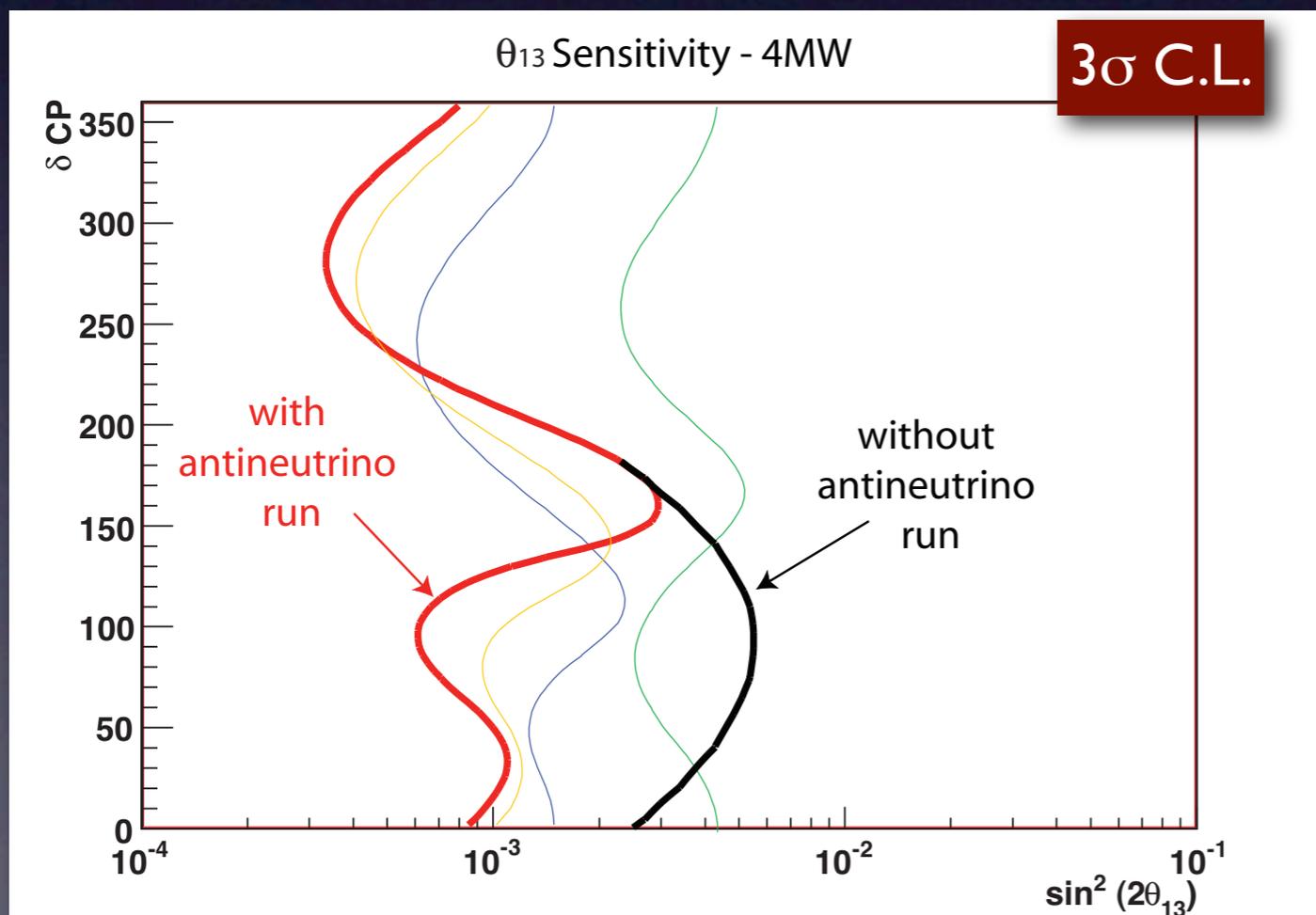
# $\theta_{13}$ sensitivity

- The sensitivity on  $\theta_{13}$  depends simply on the number of oscillated events over the background.
- Therefore, the best option is to peak the spectrum on the first maximum of oscillation and have the largest flux as possible.
- As expected, the **best configuration** for the discovery of a non-zero  $\theta_{13}$ , is the **short baseline** of 295 km at 2.5 deg. OA.



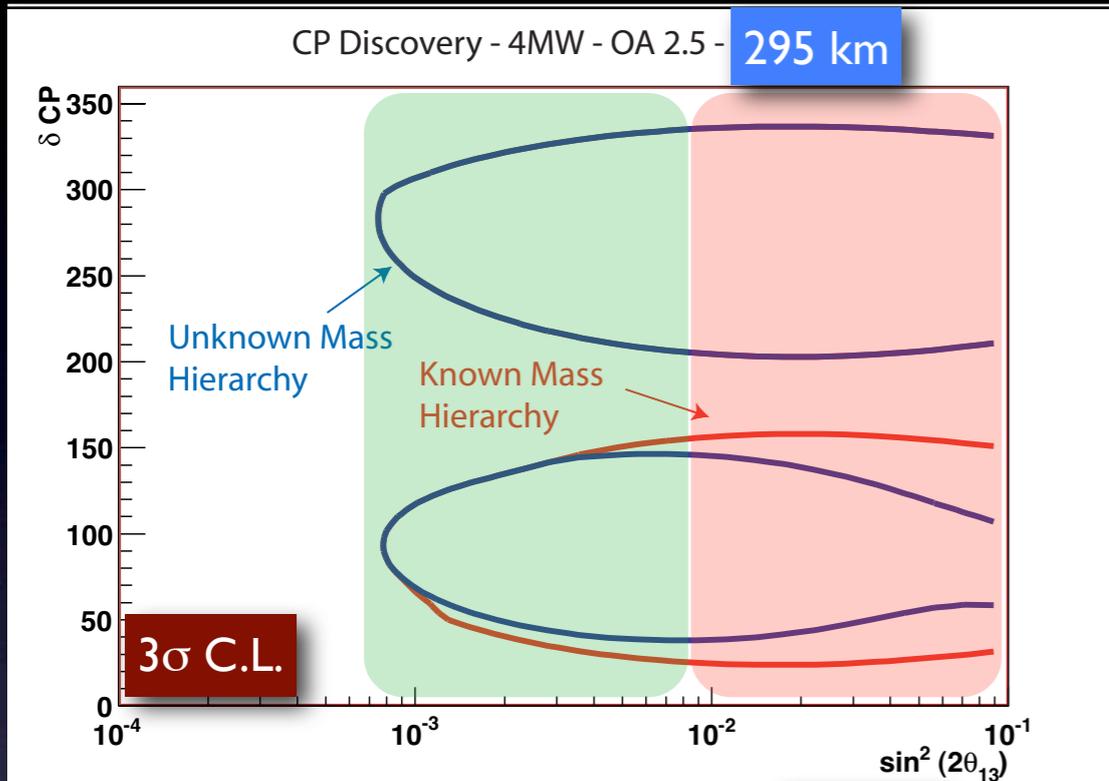
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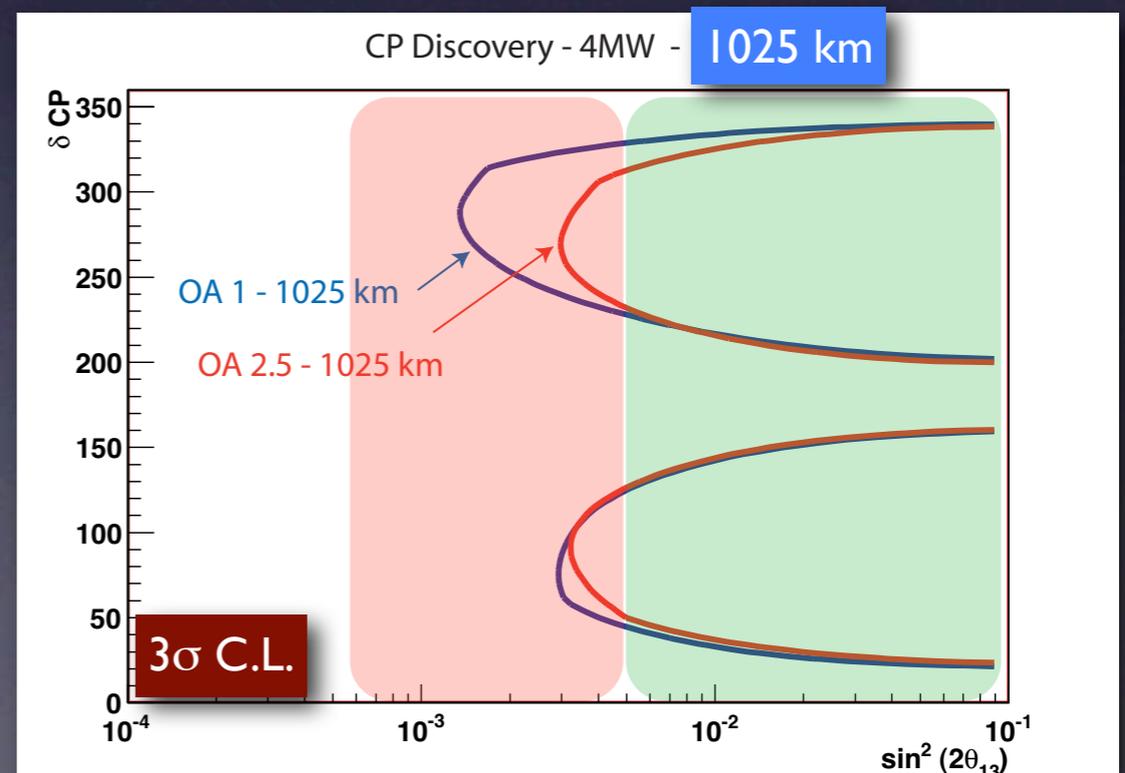
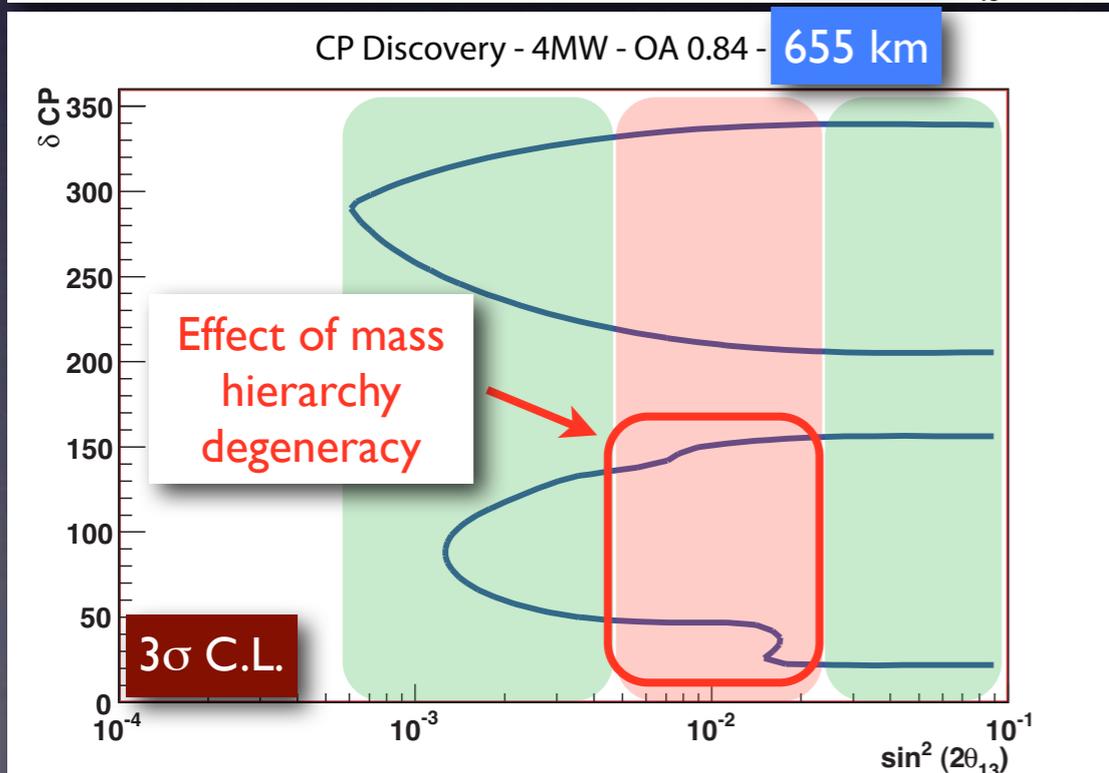


- Note that thanks to the **antineutrino run**, the characteristic “S” shape is not present and the sensitivity is improved in the region of  $\delta_{CP}$  between 0 and 180 degrees.

# CP violation

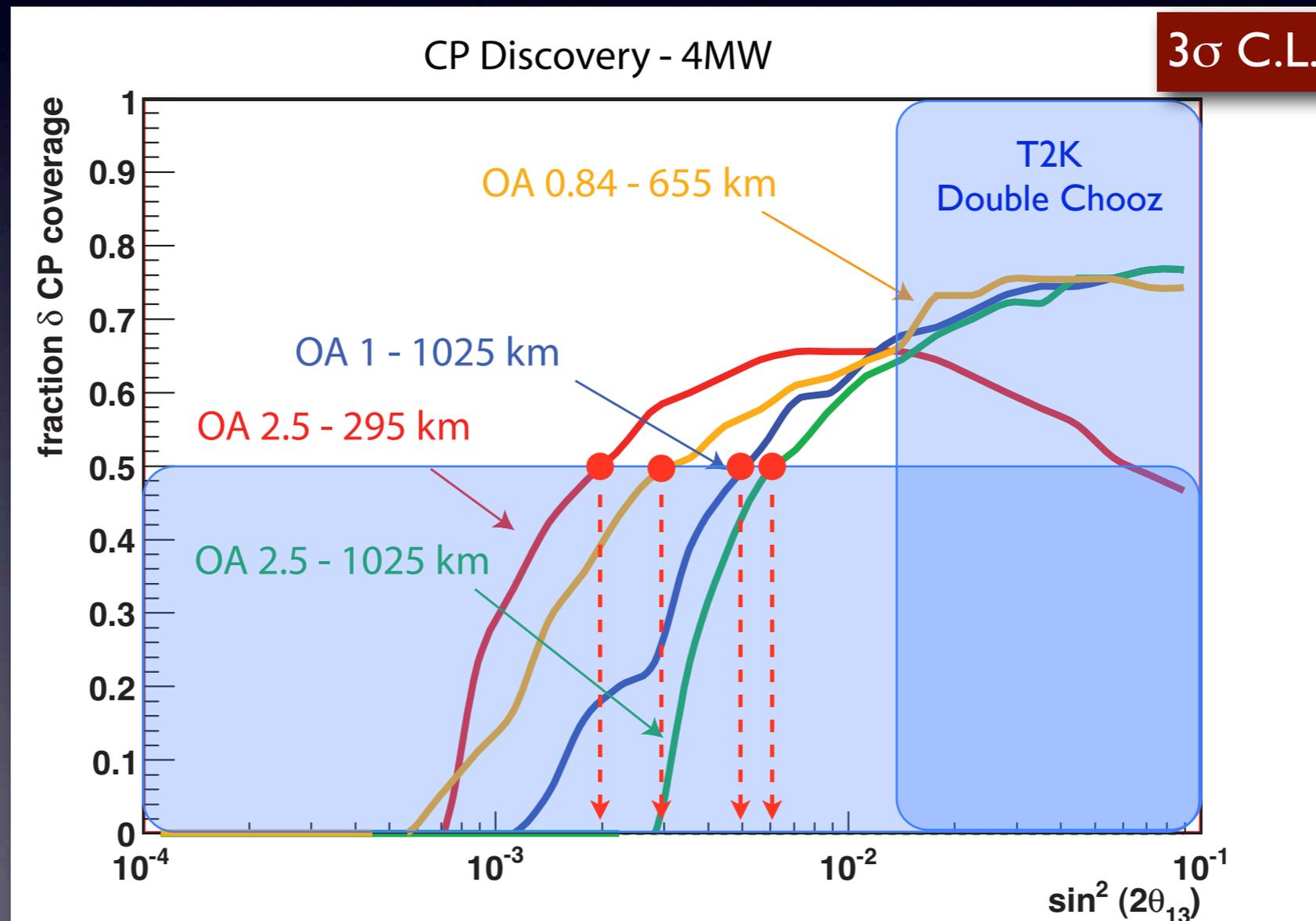


- The **optimal configuration** to observe CP violation (i.e. measure a value of  $\delta_{CP}$  different from 0 or 180 degrees) is less evident and **depends on the value of  $\theta_{13}$** .
- Mass hierarchy degeneracy makes things harder: in some regions (normal hierarchy, CP violation) is equal to (inverted hierarchy, CP conservation).



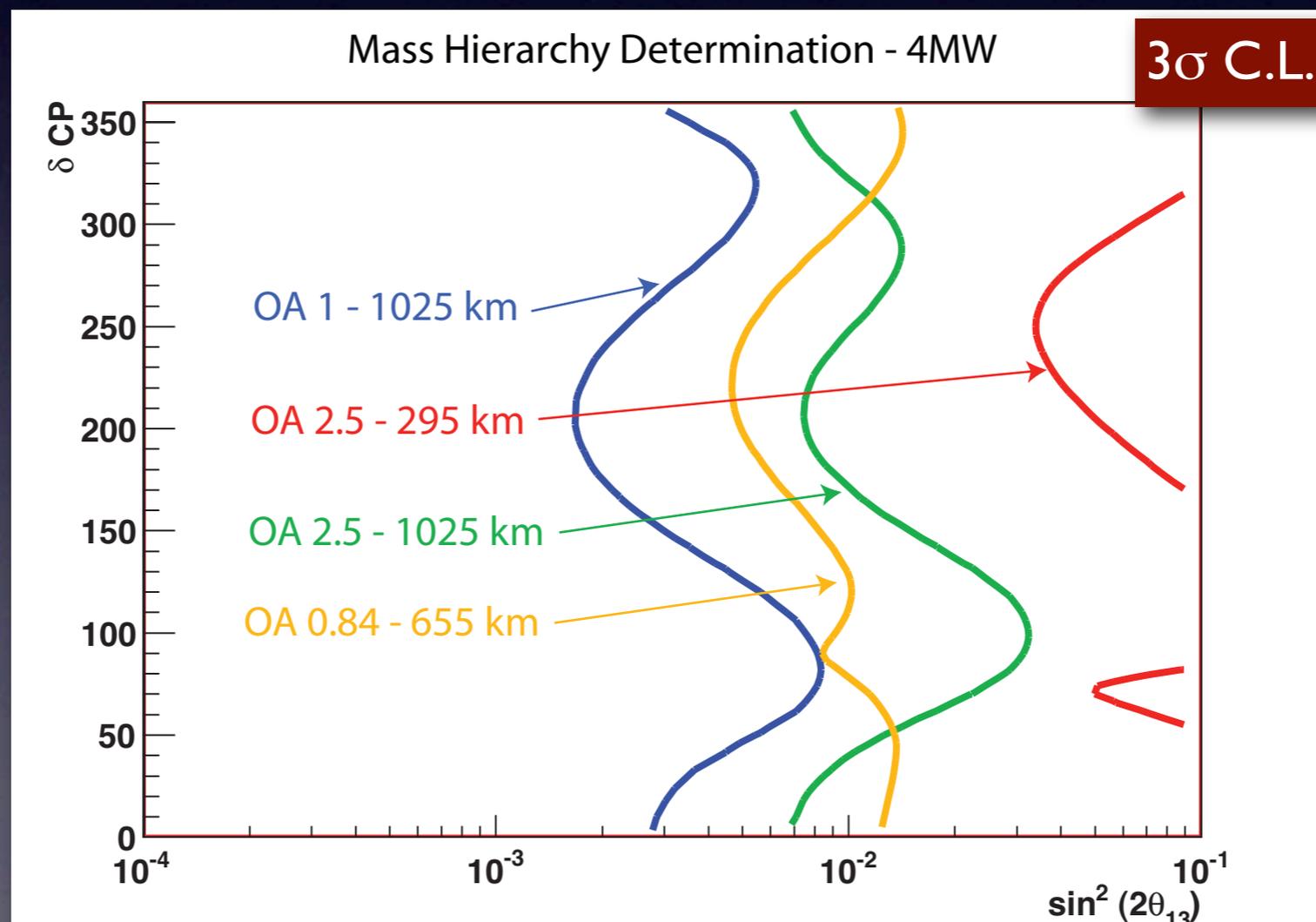
# CP coverage

- According to the value of  $\theta_{13}$  the choice of the “optimal” configuration for  $\delta_{CP}$  discovery is different.
- Results from present experiments (T2K, Double Chooz) are crucial in order to choose which way to follow.



# Mass hierarchy

- To determine the mass hierarchy the **best solution** is represented by the **longest baseline**.
- Of course the higher the flux the better it is, therefore once the baseline is set, the optimal configuration is given by the **smallest off-axis angle** (i.e. between the studied configuration the optimal option is 1025 km at 1 deg. OA).



# Conclusions

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- We studied the performance of a very large LAr TPC on future neutrino beams.
- We found that a “best” location does not exist and different locations are selected according to the measurement to be done (i.e.  $\theta_{13}$ ,  $\delta_{CP}$  and mass hierarchy determination) and for  $\delta_{CP}$  on the value of  $\theta_{13}$  (namely  $\sin^2(2\theta_{13}) < 10^{-2}$  or  $\sin^2(2\theta_{13}) > 10^{-2}$ ).
- Present experiments such as **T2K and Double Chooz will represent a turning point** for choices concerning the future of long baseline neutrino oscillation experiments in particular as far as accelerator technologies are concerned.
- The results presented are independent from the detector technology: similar results have been found with large water Cerenkov detector (about a factor of 3 in mass with respect to LAr). However, the LAr technology presents some advantages in terms of resolution, background reduction and volume.

The end