# Asymptotic Scale Invariance, Vacuum Stability and Higgs Inflation

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arXiv: 1807.\*\*\*\* with Mikhail Shaposhnikov

#### Introduction

LHC experiment

- $\cdot$  Higgs boson  $m_h \simeq 125\,{
  m GeV}$
- The SM is consistent so far.

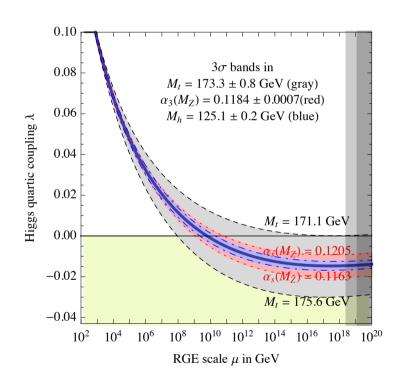
SM is valid up to very high energy scale?

 $ightharpoonup \lambda$  crosses zero, or touches zero.

EW vacuum is meta-stable.

Critical

D.Buttazzo, G.Degrassi, P.P.Giardino, G.F.Giudice, F.Sala, A.Salvio, A.Strumia (2014)



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What does this imply?

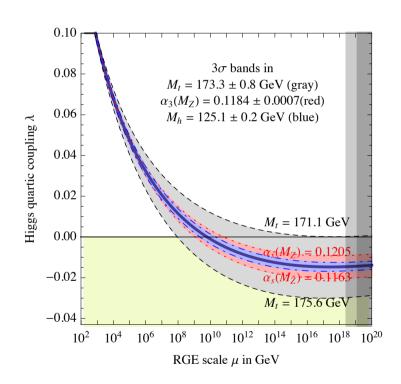


DM, GW from PT, Strong CP...

Asymptotic scale invariance Higgs inflation



D.Buttazzo, G.Degrassi, P.P.Giardino, G.F.Giudice, F.Sala, A.Salvio, A.Strumia (2014)



#### Introduction

Scale Anomaly and Effective Potential

Asymptotic Scale Invariance and Stability

Asymptotic Scale Invariance and Higgs Inflation

Future Directions (DM, GW…??)

Summary

Invariance under 
$$egin{cases} x^{\mu} o \sigma^{-1} x^{\mu} \ \Phi(x) o \sigma^{d_{\Phi}} \Phi(x) \end{cases}$$
  $d_{\Phi}$  : Mass dimension of dynamical fields  $\Phi$ 

\_\_\_\_\_\_

Explicit mass scale breaks SI:

$$V = -\frac{\mu_{\rm EW}^2}{2} h^2 + \frac{\lambda}{4} h^4$$

In the SM of particle physics sector,

SI is broken by the negative Higgs mass term.

Scale invariant for  $h\gg \mu_{\mathrm{EW}}$  . (approximately)

If you want,  $\mu_{\rm EW} \propto \phi$  : dynamical but is NOT crucial here.

SI is anomalous with regularization/renormalization NOT respecting the symmetry.

Dimensional regularization  $n=4-2\varepsilon$ 

$$\frac{\lambda h^4}{4} \Longrightarrow \mu^{\frac{2\varepsilon}{1-\varepsilon}} \times \frac{\lambda h^4}{4}$$

An explicit mass scale is introduced for the divergence and defines coupling "constants".

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One-loop (CW) correction to Higgs potential

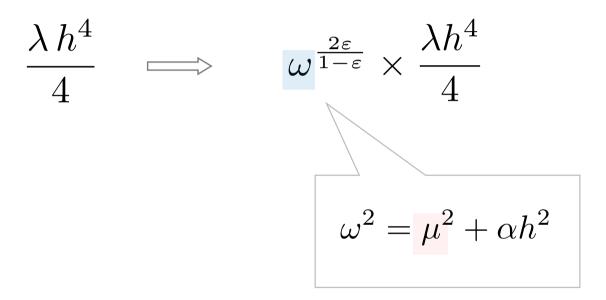
$$\Delta V \sim (-1)^F m^4 \ln \frac{m^2}{\mu^2}$$

Breaks the scale invariance

Effective potential 
$$V_{\rm eff} = \frac{h^4}{4} \left[ \lambda + \frac{B}{2} \ln \frac{h^2}{\mu^2} + \cdots \right]$$
 
$$\equiv \frac{\lambda_{\rm eff}(h)}{4} \, h^4$$

Effective potential 
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$$\equiv \frac{\lambda_{\rm eff}(h)}{4} h^4$$
 Destabilized at 
$$B < 0 \qquad h_{\star} \sim 10^{10} {\rm GeV}$$
 if  $m_t = 173 \, {\rm GeV}$ 

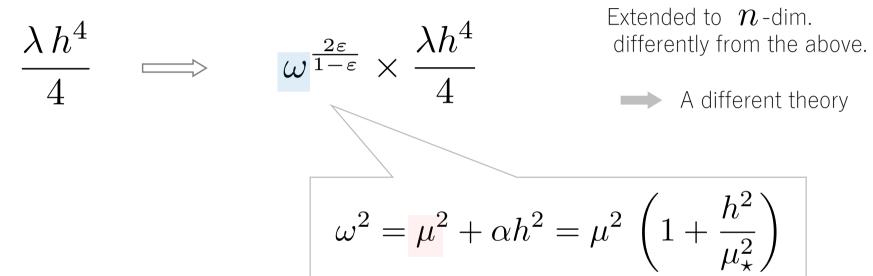
$$V_{\rm eff} = \frac{h^4}{4} \left[ \lambda + \frac{B}{2} \ln \frac{h^2}{\mu^2} + \frac{B'}{8} \left( \ln \frac{h^2}{\mu^2} \right)^2 + \cdots \right]$$
 
$$\equiv \frac{\lambda_{\rm eff}(h)}{4} h^4$$
 
$$h \sim {\rm Planck}$$
 
$$\lambda = B = 0$$
 if  $m_t = 171 \, {\rm GeV}$  "Critical"



Field-dependent

Extended to  $\,n$ -dim. differently from the above.

A different theory

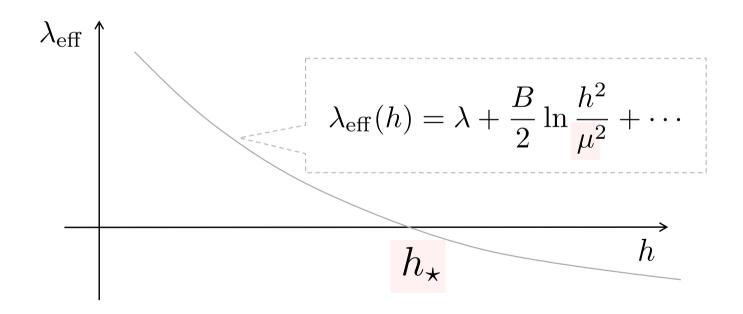


Explicit mass scale is negligible for

$$h \gg \mu_{\star}$$

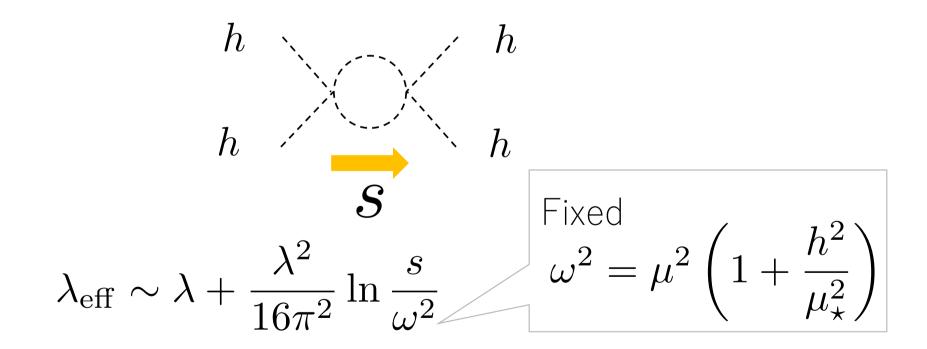
#### **Asymptotic Scale Invariance**

$$\lambda_{\text{eff}}^{\text{aSI}}(h) = \lambda + \frac{B}{2} \ln \frac{h^2}{\omega^2} + \cdots$$

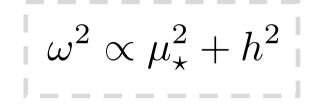


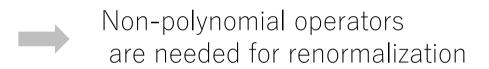
$$\lambda_{\rm eff}^{\rm aSI}(h) = \lambda + \frac{B}{2} \ln \frac{h^2}{\mu^2 (1 + h^2/\mu_\star^2)} + \cdots \qquad \lambda_{\rm eff}(\mu_\star) > 0$$
 EW vacuum is global minimum 
$$h \sim \mu_\star \qquad \text{if } \mu_\star < h_\star$$
 
$$\lambda_{\rm eff}(\mu_\star)$$

Couplings run as energy scale of scattering increases.



$$\omega^{\frac{2\varepsilon}{1-\varepsilon}} \times \frac{\lambda h^4}{4}$$
 : Non-renormalizable





$$\frac{h^{4+2k}}{(\mu_{\star}^2 + h^2)^k} \qquad (k \ge 1)$$

 $\sim h^4$  for  $h \gg \mu_\star$  Asymptotically SI

$$\omega^{\frac{2\varepsilon}{1-\varepsilon}} \times \frac{\lambda h^4}{4}$$
 : Non-renormalizable

$$\omega^2 \propto \mu_\star^2 + h^2$$

Non-polynomial operators are needed for renormalization

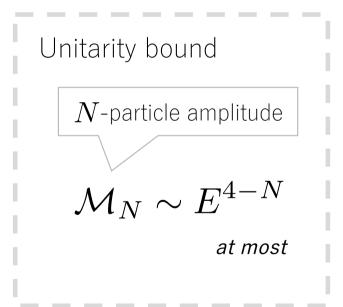
$$\frac{h^{4+2k}}{(\mu_{\star}^2 + h^2)^k} \qquad (k \ge 1)$$

Up to which energy scale is this effective theory valid?

$$\sim h^4$$
 for  $h \gg \mu_\star$  Asymptotically SI

$$\frac{h^{4+2k}}{(\mu_{\star}^2 + h^2)^k}$$

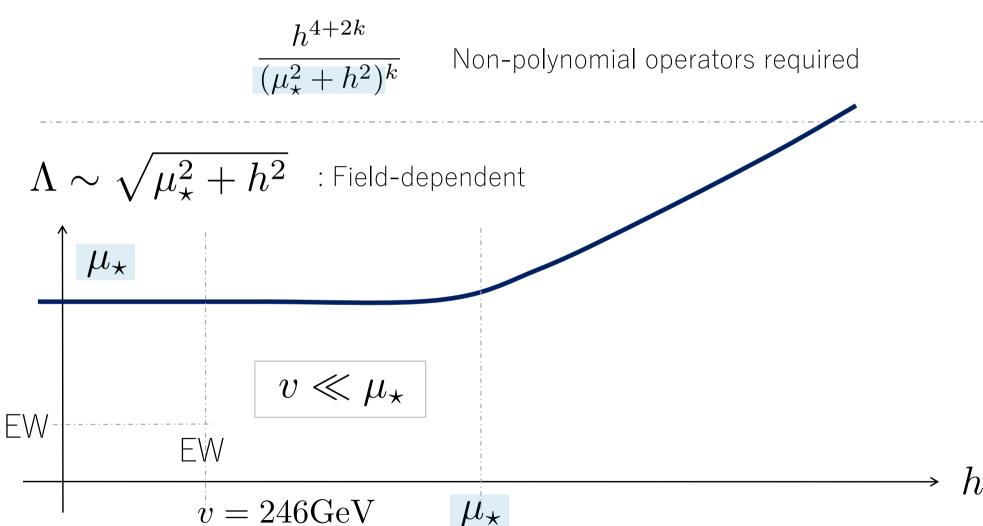
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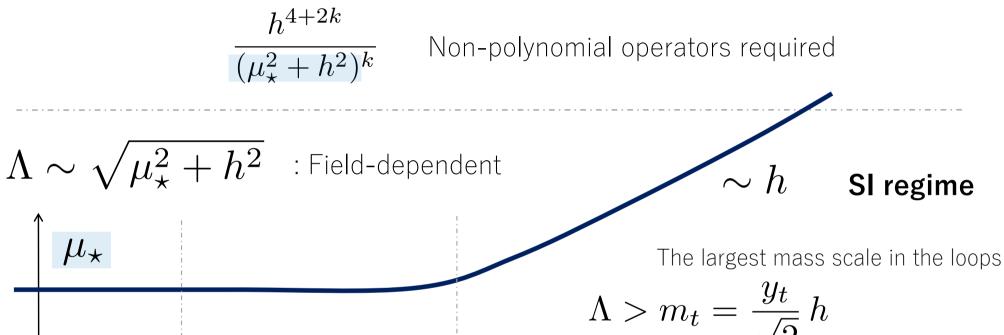


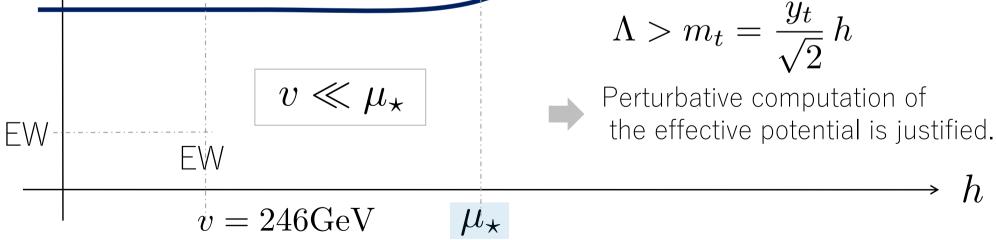
J.M.Cornwall, D.N.Levin, G.Tiktopoulos (1974)

Tree unitarity violation

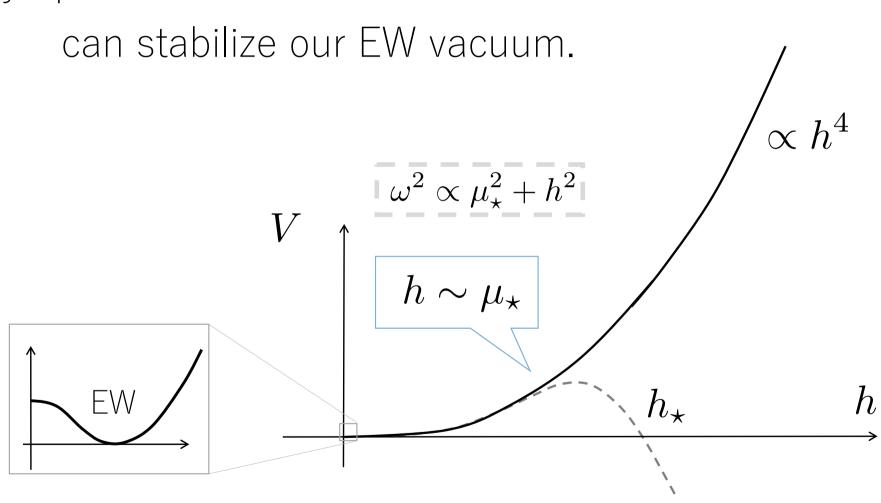
at 
$$\Lambda \sim \sqrt{\mu_{\star}^2 + h^2}$$
 Strong coupling or New physics







# Asymptotic Scale Invariance



$$\frac{\mathcal{L}}{\sqrt{-g}} = \frac{M_{\mathrm{P,eff}}^2}{2} \, R - \frac{g^{\mu\nu}}{2} \partial_\mu h \partial_\nu h - V(h) + \cdots$$
 Effective Planck mass 
$$M_{\mathrm{P,eff}}^2 = M_{\mathrm{P}}^2 + \xi h^2$$

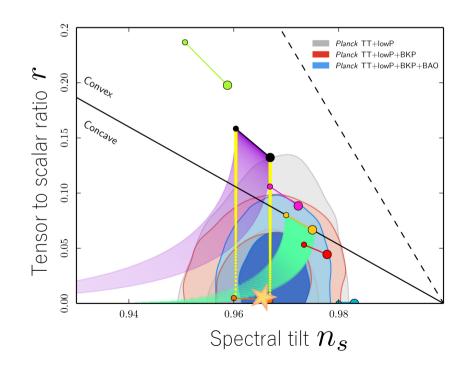
F.Bezrukov, M.Shaposhnikov (2007)

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Effective 
$$M_{
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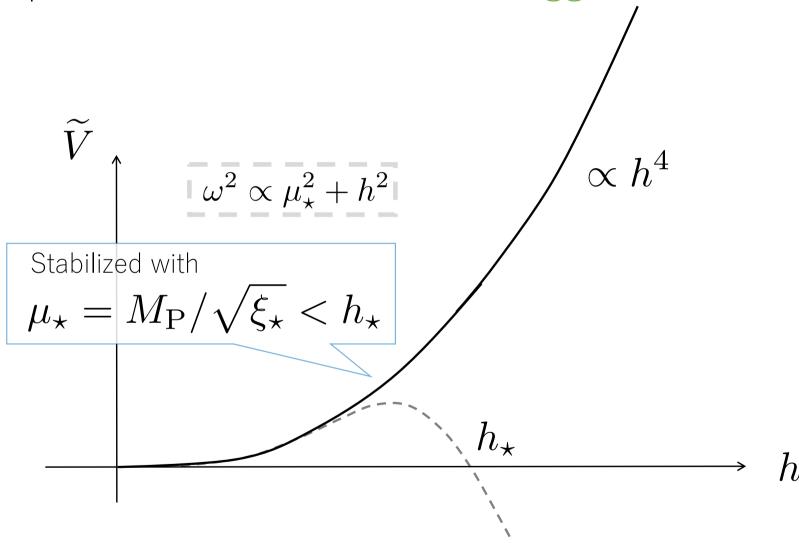
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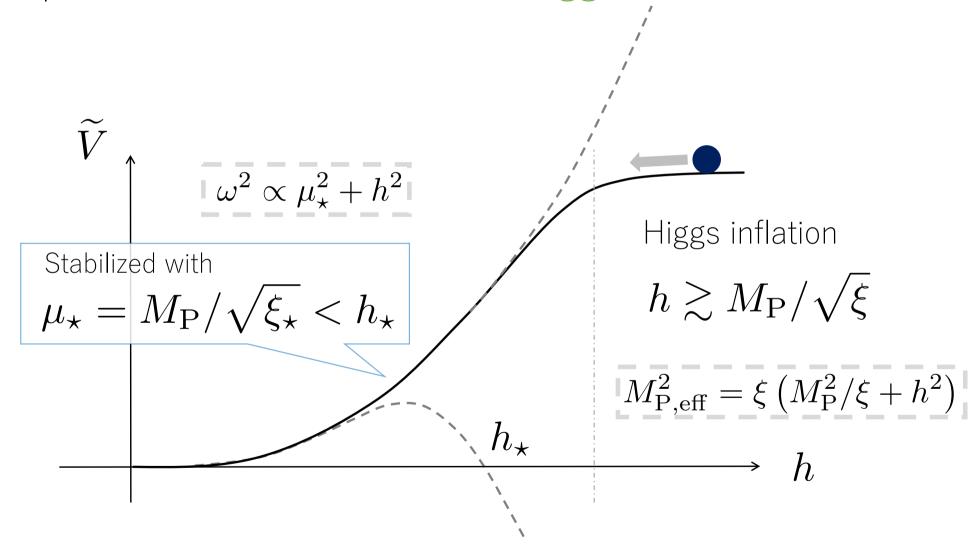
$$\xi \sim 10^4 \sqrt{\lambda}$$
 Large non-minimal coupling 
$$\Rightarrow A_s \simeq 2.2 \times 10^{-9}$$



Renormalization prescription	$\omega^2 \propto$	$\frac{\lambda h^4}{4} \implies \omega^{\frac{2\varepsilon}{1-\varepsilon}} \times \frac{\lambda h^4}{4}$
I	$M_{\rm P}^2 + \xi h^2 = M_{\rm P}^2$	F.Bezrukov, M.Shaposhnikov (2007)
II	$M_{ m P}^2$ (constant)	A.O.Barvinsky, A.Y.Kamenshchik, A.A.Starobinsky (2008)

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$\mu_{\star}^2 + h^2$	$\propto M_P^2 + \xi_{\star} h^2$	
		$\xi_{\star} = M_{\rm P}^2/\mu_{\star}^2 \neq \xi$





✓ Perturbative computation of effective potential is justified.

$$\Lambda > m_t$$
 (the largest mass scale in the loops)

 $\checkmark$  Generation of inflaton (Higgs) fluctuation is also computable.

$$\Lambda > H > k_{
m fluc}$$
 during Higgs inflation

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?? Reheating temperature becomes very high.

$$T_{
m rh} > \Lambda|_{h=0} = \mu_{\star}$$
 ( zero mode vanishes after thermalization )

Thermal history after inflation ( ⇒ Inflationary observable ) cannot be discussed.

- $egin{array}{c} 1, ext{ Theory above $\Lambda$}? \ 2, ext{ When } T_{
  m rh} < \Lambda ? \end{array}$

Critical case

$$V_{\rm eff} = \frac{h^4}{4} \left[ \lambda_c + \frac{B_c}{2} \ln \frac{h^2}{\mu_c^2} + \frac{B_c'}{8} \left( \ln \frac{h^2}{\mu_c^2} \right)^2 + \cdots \right]^{\frac{3\sigma \text{ bands in}}{M_t = 173.3 \pm 0.8 \text{ GeV (gray)}}}_{0.06}$$

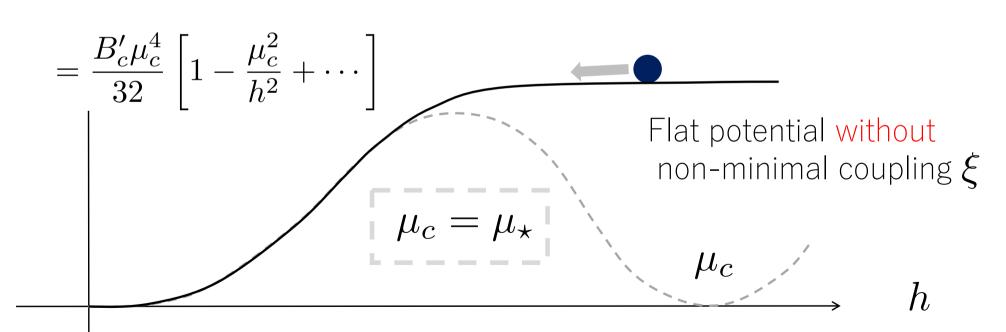
$$V \text{ with the standard renormalization prescription}$$

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 $\mu_c$ 

Critical case

$$V_{\text{eff}} = \frac{h^4}{4} \left[ \lambda_c + \frac{B_c}{2} \ln \frac{h^2}{\mu_c^2 + h^2} + \frac{B_c'}{8} \left( \ln \frac{h^2}{\mu_c^2 + h^2} \right)^2 + \cdots \right]$$

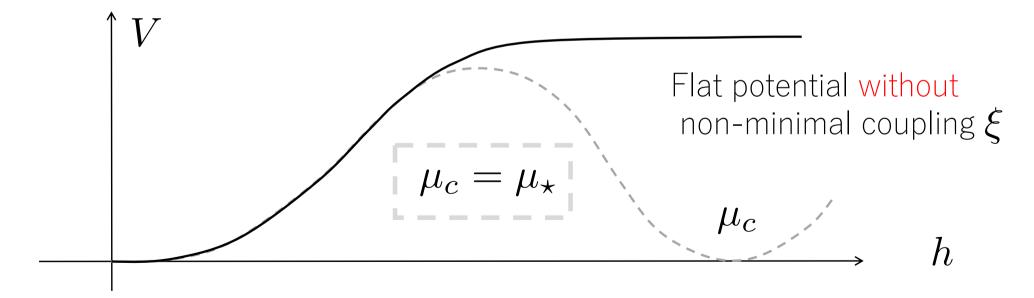


#### Critical case

$$\mu_c \sim 10^{17} \text{GeV} < M_P$$

$$\Rightarrow A_s \simeq 2.2 \times 10^{-9}$$

$$T_{
m rh} \lesssim V^{1/4}$$
 
$$< \mu_c = \Lambda|_{h=0}$$
 Perturbative computation is valid.



Future directions (DM, GW…??)

Critical case (  $T_{
m rh} < \Lambda$  , Asymptotic SI itself needs nothing new for cosmology.)



Any reasonable way to extend the SM??

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$$n_s \approx 0.975 \implies 2\sigma$$
 level

 $n_s \approx 0.975 \implies 2\sigma$  level  $\implies$  Non-standard thermal history is preferred.

$$n_s pprox 1 - rac{3}{2N}$$
 with  $N=60$ 

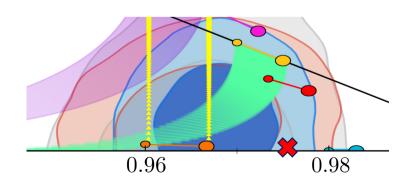
$$\Delta N \sim 15$$

Super-cooling stage before a phase transition?

Gravitational waves, DM production after PT



Typically, in scale invariant models



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Theory above  $\Lambda$ ?

Needs to guarantee the absence of  $\Lambda^2 h^2$ 



DM production with  $T>\Lambda$  Phase transition at  $\,T\sim\Lambda$ 

# Summary

Asymptotic Scale Invariance can be responsible for our EW vacuum stability.

Perturbative computation of the effective potential is valid because tree-unitary violation scale  $\Lambda$  is larger than any others.

Higgs inflation is also possible.

However,  $T_{
m rh} > \Lambda$  requires a theory above  $\Lambda$  .

 $T_{
m rh} < \Lambda$  for the critical case. The theory below  $\Lambda$  is enough.

Thank you