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Flavourful Z' portal for vector-like neutrino Dark Matter and $R_{K^{(*)}}$ [arXiv:1803.04430]

Mathias Pierre

In collaboration with **A. Falkowski**, **S. F. King** and **E. Perdomo**

Dark Side of the Universe, LAPTH, June 26 2018

1. Motivation

- The Waning of the WIMP?
- The $R_{K^{(*)}}$ anomalies

2. A flavourful model for $R_{K^{(*)}}$ & Dark Matter

- The model
- Flavour constraints
- Dark Matter searches
- Results

Motivation

The Waning of the WIMP?

- The **WIMP miracle** $\Omega_{\text{DM}} h^2 \propto \langle \sigma v \rangle_{\text{relic}}^{-1} \sim 0.12$

$$\langle \sigma v \rangle_{\text{relic}} \sim 10^{-26} \text{ cm}^3 \text{ s}^{-1} \sim 10^{-9} \text{ GeV}^{-2} \sim \frac{g^4}{m_Z^4} m_{\text{DM}}^2$$

- Higgs and Z -portal are the **simplest** miracles [Arcadi et al. '17]

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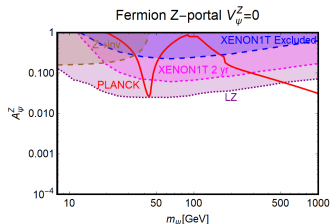
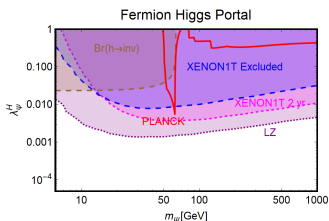
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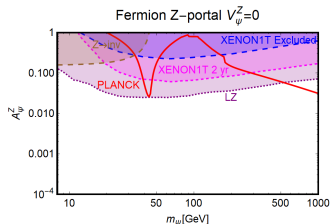
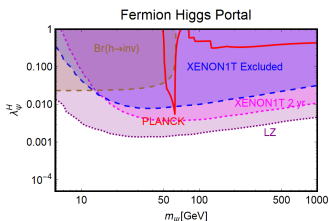
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The $R_{K^{(*)}}$ anomalies

- Test of **Lepton Flavor Universality (LFU) violation** in $b \rightarrow s \ell \ell$

$$R_{K^{(*)}}^{[q_{\min}^2, q_{\max}^2]} \equiv \frac{\mathcal{B}(B \rightarrow K^{(*)} \mu^+ \mu^-)}{\mathcal{B}(B \rightarrow K^{(*)} e^+ e^-)} \Big|_{q^2 \in [q_{\min}^2, q_{\max}^2]}$$

- **Theory under control:** hadronic uncertainties cancels out, below $\bar{c}c$ resonance

$$R_{K^{(*)}}^{\text{SM}} = 1.00(1)$$

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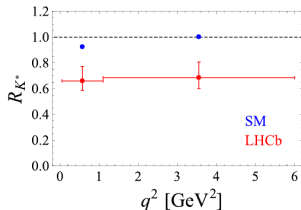
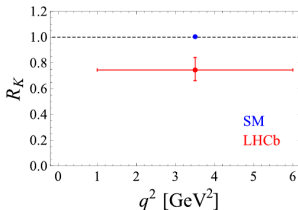
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$$R_{K^*}^{[1.1, 6]} = 0.685_{-0.083}^{+0.122},$$

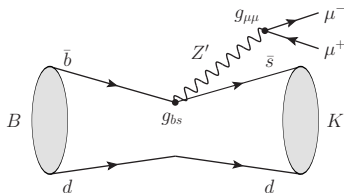
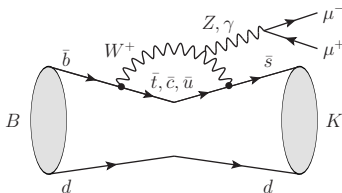
$$R_K^{[1, 6]} = 0.745_{-0.082}^{+0.097}.$$



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- **Flavor Changing Neutral Current (FCNC)** loop generated in SM
 $\rightarrow R_{K^{(*)}}$ sensitive to **New Physics (NP)** featuring FCNC!
- **NP** can be described as

$$\mathcal{H}_{\text{eff}}^{\text{NP}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_{i,\ell} \left(C_{i\ell}^{\text{NP}} \mathcal{O}_i^\ell + C'_{i\ell}{}^{\text{NP}} \mathcal{O}'_{i\ell} \right)$$



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- $R_{K^{(*)}}$ **explained** for $C_{9\mu}^{\text{NP}} = -C_{10\mu}^{\text{NP}}$ [Capdevila et al. '17]

$$\mathcal{O}_9^{(\prime)\ell} = (\bar{s}\gamma_\mu P_{L(R)} b)(\bar{\ell}\gamma^\mu \ell), \quad \mathcal{O}_{10}^{(\prime)\ell} = (\bar{s}\gamma_\mu P_{L(R)} b)(\bar{\ell}\gamma^\mu \gamma_5 \ell) .$$

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$$G_{bs\mu} \sim \frac{1}{(30 \text{ TeV})^2}$$

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A flavourful model for $R_{K^{(*)}}$ & Dark Matter

A flavourful model including a 4th generation

- Consider a 4th **vector-like (V-L) family** charged under extra $U(1)'$
SM particle content not charged under $U(1)'$

Field	Representation/charge			
	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$	$U(1)'$
Q_{Li}	3	2	1/6	0
u_{Ri}	3	1	2/3	0
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L_{Li}	1	2	-1/2	0
e_{Ri}	1	1	-1	0
ν_{Ri}	1	1	0	0
H	1	2	1/2	0
Q_{L4}, \bar{Q}_{R4}	3	2	1/6	q_{Q4}
u_{R4}, \tilde{u}_{L4}	3	1	2/3	q_{u4}
d_{R4}, \tilde{d}_{L4}	3	1	-1/3	q_{d4}
L_{L4}, \tilde{L}_{R4}	1	2	-1/2	q_{L4}
e_{R4}, \tilde{e}_{L4}	1	1	-1	q_{e4}
$\nu_{R4}, \tilde{\nu}_{L4}$	1	1	0	$q_{\nu4}$
$\phi_{Q,u,d,L,e}$	1	1	0	$-q_{Q4,u4,d4,L4,e4}$

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 SM singlets scalars ϕ responsible for **generation mixing**

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The flavourful Lagrangian

- Mass and Yukawa terms in **flavour basis** ($i = 1, 2, 3$)

$$\mathcal{L} \supset x_i^Q \phi_Q \bar{Q}_{Li} \tilde{Q}_{R4} + x_i^u \phi_u \tilde{u}_{L4} u_{Ri} + M_4^Q \bar{Q}_{L4} \tilde{Q}_{R4} + M_4^u \tilde{u}_{L4} u_{R4} + M_4^\nu \tilde{\nu}_{L4} \nu_{R4}$$

- **Non-diagonal** mass terms generated **after the $U(1)'$ breaking** $\langle \phi \rangle \neq 0$
→ diagonalized with 4×4 unitary matrices $V_{\alpha\beta}$ ($\alpha, \beta = 1, \dots, 4$)
- **After $U(1)'$ SSB:** Mixing mostly with **3rd** gen. of Q and **2nd** gen. of L

$$\mathcal{L}_{Z'}^{\text{gauge}} = g' Z'_\mu \left(q_{Q4} (s_{34}^Q)^2 \bar{Q}'_{L3} \gamma^\mu Q'_{L3} + q_{L4} (s_{24}^L)^2 \bar{L}'_{L2} \gamma^\mu L'_{L2} \right)$$

- Expanding the **primed field** in mass eigenstates **after EW SSB**

$$b'_L = (V_{dL}^{\prime\dagger})_{33} b_L + (V_{dL}^{\prime\dagger})_{32} s_L + \dots \quad \simeq b_L + (V_{dL}^{\prime\dagger})_{32} s_L$$

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Assuming $|(V_{dL}^{\prime\dagger})_{32}| \approx |V_{ts}| \approx 0.04$

Generation of **tree-level** $Z'_\mu \bar{b}_L \gamma^\mu s_L$ term

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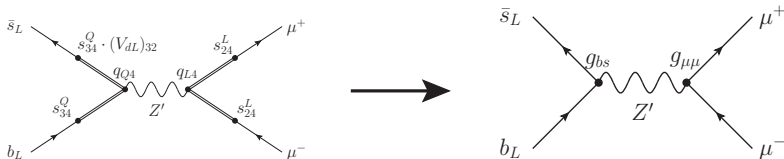
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Generation of **tree-level** $Z'_\mu \bar{b}_L \gamma^\mu s_L$ term

The phenomenological Lagrangian



- The **relevant terms** at the end are:

$$\mathcal{L} \supset Z'_\mu \left(g_{bb} \sum_{q=b,t} \bar{q}_L \gamma^\mu q_L + g_{bs} \bar{b}_L \gamma^\mu s_L + g_{\mu\mu} \sum_{\ell=\mu,\nu_\mu} \bar{\ell}_L \gamma^\mu \ell_L + g_{\nu\nu} \bar{\nu}_4 \gamma^\mu \nu_4 \right),$$

with the relations $g_{bs} = g_{bb} V_{ts}$ and $g_{\nu\nu} = g' q_{\nu 4}$

- A **dark matter** candidate $\nu \equiv \nu_4$ **stable**
- A heavy Z' **mediator** featuring **tree-level FCNC**

5 parameters: $g_{\nu\nu}$, $g_{\mu\mu}$, g_{bb} , $m_{Z'}$ and m_ν

$R_{K^{(*)}}$ relation

$$G_{bs\mu} = -\frac{g_{bs}g_{\mu\mu}}{M_{Z'}^2} = -\frac{V_{ts}g_{bb}g_{\mu\mu}}{M_{Z'}^2} \approx \frac{1}{(31.5 \text{ TeV})^2}$$

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Neutrino trident: $\nu_\mu\gamma^* \rightarrow \nu_\mu\mu^+\mu^-$

$$\Delta\mathcal{L}_{\text{eff}} \supset -\frac{G_\mu}{2}(\bar{\ell}_L\gamma^\mu\ell_L)^2, \quad G_\mu = \frac{g_{\mu\mu}^2}{M_{Z'}^2}$$
$$-\frac{1}{(390 \text{ GeV})^2} \lesssim G_\mu \lesssim \frac{1}{(370 \text{ GeV})^2}$$

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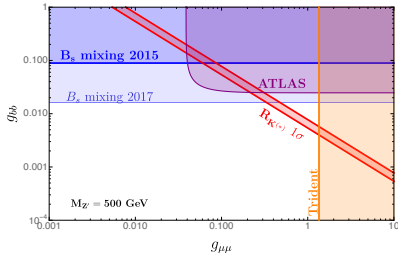
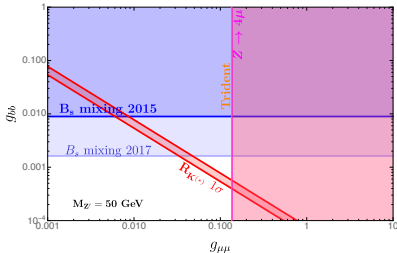
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$B_s - \bar{B}_s$ mixing

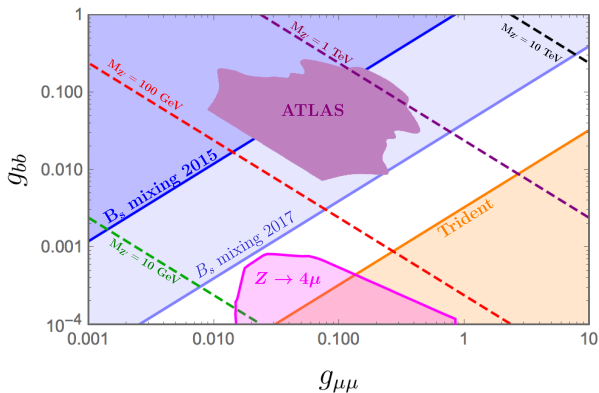
$$\Delta\mathcal{L}_{\text{eff}} \supset -\frac{G_{bs}}{2}(\bar{s}_L\gamma^\mu b_L)^2 + \text{h.c.}, \quad G_{bs} = \frac{g_{bs}^2}{M_{Z'}^2} = \frac{g_{bb}^2 V_{ts}^2}{M_{Z'}^2}$$
$$-\frac{1}{(180 \text{ TeV})^2} \lesssim G_{bs} \lesssim \frac{1}{(770 \text{ TeV})^2}$$

Collider and other constraints

- $pp \rightarrow \gamma^* Z \rightarrow 4\mu$ constrains **low masses** $5 \lesssim M_{Z'} \lesssim 70$ GeV
- $pp \rightarrow Z' \rightarrow \mu^+ \mu^-$ constrains **higher masses** $200 \lesssim M_{Z'} \lesssim 1000$ GeV
- Precise measurements of LFU $R_{1s}^{\tau/\mu}$ of Υ_{1s} and $(g-2)_\mu$ set weaker constraints



Summary of the constraints



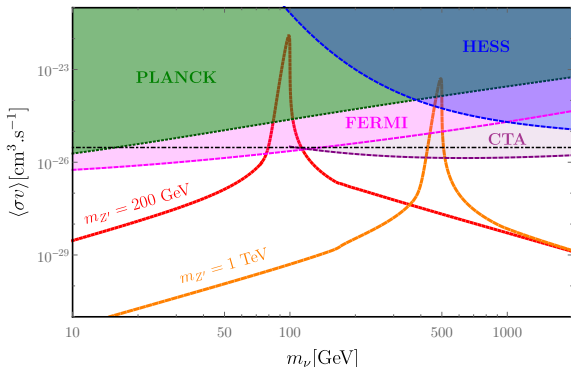
- $M_{Z'} \gtrsim$ few TeV **unnatural**
- $M_{Z'} \lesssim 50 \text{ GeV}$ **strongly constrained**
- Hierarchy $g_{\mu\mu} \gg g_{bb}$ satisfied in the **natural** parameter space

Dark matter relic density

- Dark matter **annihilations** $\bar{\nu}\nu \rightarrow \bar{\psi}\psi, Z'Z'$ with $\psi = b, t, \mu, \nu_\mu$

$$\Omega_\nu h^2 \simeq 0.12 \left(\frac{3 \times 10^{-26} \text{cm}^3 \text{s}^{-1}}{\langle \sigma v \rangle} \right) \quad \text{with} \quad \langle \sigma v \rangle = \sum_\psi \langle \sigma v \rangle_{\bar{\psi}\psi} + \langle \sigma v \rangle_{Z'Z'}$$

- $\langle \sigma v \rangle$ is **s-wave** dominated \rightarrow constraints from **Indirect Searches**:
Fermi and **HESS** ($b\bar{b}$), **Planck** ($\mu^+\mu^-$) and upcoming **CTA** ($\mu^+\mu^-$)



Dark matter direct searches

- Effective operators generated at the scale $\mu \simeq M_{Z'}$

$$\mathcal{L}_{\text{eff}} \supset - \sum_{f=\mu,b} \frac{g_{\nu\nu} g_{ff}}{M_{Z'}^2} \bar{f}_L \gamma^\alpha f_L \bar{\nu} \gamma_\alpha \nu,$$

- Below $\mu < M_{Z'}$, **DM couplings to light quarks** are induced via **RGE**

$$\mathcal{L}_{\text{eff}} \supset \sum_{q=u,d} C_{1,f}^{(6)}(\mu) \bar{q} \gamma^\alpha q \bar{\nu} \gamma_\alpha \nu \quad \text{with} \quad C_{1,f}^{(6)}(\mu) \sim \frac{\alpha}{4\pi} \frac{g_{\nu\nu} g_{ff}}{M_{Z'}^2} \log\left(\frac{M_{Z'}}{\mu}\right)$$

- Effective operator generated at **nuclear scale** $\mu \simeq 2 \text{ GeV}$

$$\mathcal{L}_{\text{eff,NR}} \supset \sum_{N=p,n} c_1^N \bar{\nu} \nu \bar{N} N \rightarrow \sigma_{\text{DD}} \sim \left(\frac{g_{\nu\nu}}{0.2}\right)^2 \left(\frac{g_{\mu\mu}}{0.1}\right)^2 \left(\frac{m_Z}{M_{Z'}}\right)^4 10^{-45} \text{ cm}^2.$$

Dark matter direct searches

- Effective operators generated at the scale $\mu \simeq M_{Z'}$

$$\mathcal{L}_{\text{eff}} \supset - \sum_{f=\mu,b} \frac{g_{\nu\nu} g_{ff}}{M_{Z'}^2} \bar{f}_L \gamma^\alpha f_L \bar{\nu} \gamma_\alpha \nu,$$

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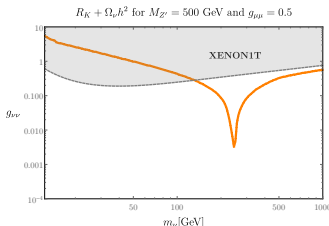
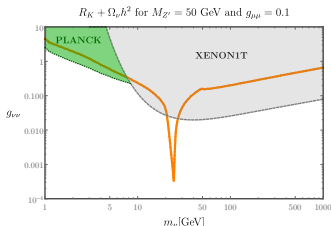
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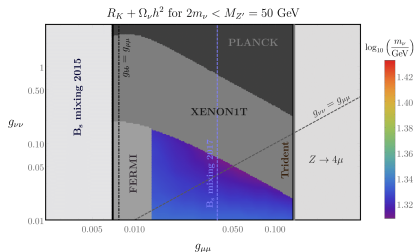
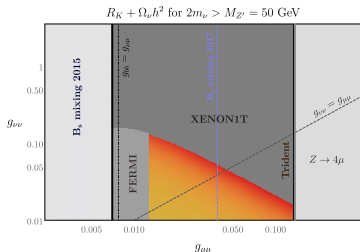
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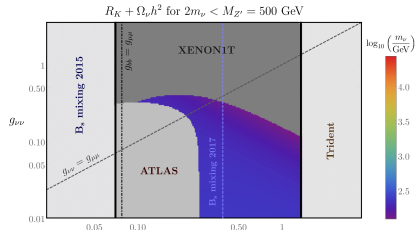
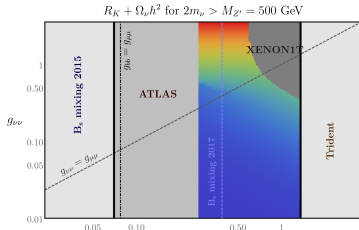
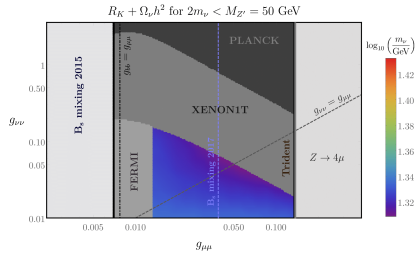
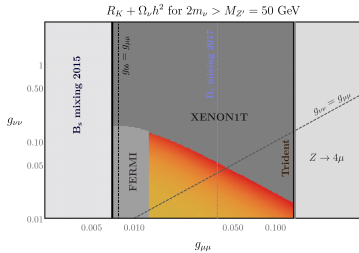
The surviving parameter space

- Each point satisfies $R_{K^{(*)}}$ and $\Omega_{\nu} h^2$, **color code** represents **DM mass**
- All the grey points are excluded



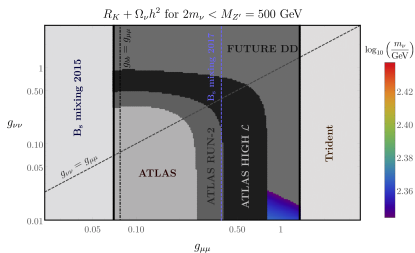
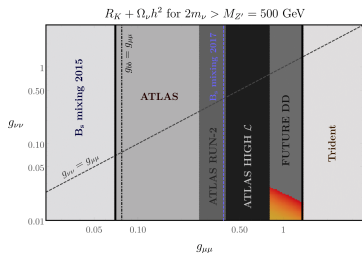
The surviving parameter space

- Each point satisfies $R_{K^{(*)}}$ and $\Omega_\nu h^2$, color code represents DM mass
- All the grey points are excluded



The future parameter space

- Assuming 3000 fb^{-1} and 2 orders of magnitude improvement on σ_{DD}



- The **viable parameter space** should be **probed** in the (near) future!

Summary

- Considered 4th vector-like family charged under $U(1)'$
- Non-diagonal couplings generate mixing between generations
- The 4th right-handed neutrino is a good DM candidate
- $R_{K(*)}$ and $\Omega_{\text{DM}}h^2$ can be explained simultaneously
- The remaining parameter space will be probed

Thank you for your attention!

Back-up slides

The flavourful Lagrangian: Standard Model sector

- Mass and Yukawa terms in **flavour basis** ($i = 1, 2, 3$)

$$\mathcal{L} \supset x_i^Q \phi_Q \bar{Q}_{Li} \tilde{Q}_{R4} + x_i^u \phi_u \tilde{u}_{L4} u_{Ri} + M_4^Q \bar{Q}_{L4} \tilde{Q}_{R4} + M_4^u \tilde{u}_{L4} u_{R4} + M_4^\nu \tilde{\nu}_{L4} \nu_{R4}$$

- **Non-diagonal** mass terms generated **after the $U(1)'$ breaking** $\langle \phi \rangle \neq 0$
→ diagonalized with 4×4 unitary matrices $V_{\alpha\beta}$ ($\alpha, \beta = 1, \dots, 4$)

$$Q_L'^{\alpha} = V_{Q_L}^{\alpha\beta} Q_L^{\beta}, \quad u_R'^{\alpha} = V_{u_R}^{\alpha\beta} u_R^{\beta} \quad \rightarrow \text{mixing between the 4 generations}$$

- $\mathcal{L} \supset y_{ij}^u \bar{Q}_{Li} \tilde{H} u_{Rj}$ becomes $\mathcal{L} \supset y_{ij}'^u \bar{Q}'_{Li} \tilde{H} u'_{Rj}$ with $y_{ij}'^u = (V_{Q_L} y_{ij}^u V_{u_R}^\dagger)_{ij}$
- Diagonalization after EW SSB:

$$V_{uL}^\dagger y'^u V_{uR}'^\dagger = \text{diag}(y_u, y_c, y_t) \rightarrow \boxed{V_{\text{CKM}} \equiv V_{uL}' V_{dL}'^\dagger}$$

- Gauge sectors \mathcal{L}_Z and \mathcal{L}_{W^\pm} unchanged → **GIM mechanism**

The flavourful Lagrangian: The hidden gauge sector

- Before $U(1)'$ SSB:

$$\mathcal{L}^{Z'} \supset g' Z'_\mu \bar{Q}_L D_Q \gamma^\mu Q_L + (Q \rightarrow L) \quad \text{with} \quad D_Q \equiv \text{diag}(0, 0, 0, q_{Q4})$$

- After $U(1)'$ SSB: Mixing mostly with **3rd** gen. of Q and **2nd** gen. of L

$$D'_Q = q_{Q4} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & (s_{34}^Q)^2 & c_{34}^Q s_{34}^Q \\ 0 & 0 & c_{34}^Q s_{34}^Q & (c_{34}^Q)^2 \end{pmatrix}, \quad D'_L = q_{L4} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & (s_{24}^L)^2 & 0 & c_{24}^L s_{24}^L \\ 0 & 0 & 0 & 0 \\ 0 & c_{24}^L s_{24}^L & 0 & (c_{24}^L)^2 \end{pmatrix}$$

$$\mathcal{L}_{Z'}^{\text{gauge}} = g' Z'_\mu \left(q_{Q4} (s_{34}^Q)^2 \bar{Q}'_{L3} \gamma^\mu Q'_{L3} + q_{L4} (s_{24}^L)^2 \bar{L}'_{L2} \gamma^\mu L'_{L2} \right)$$

- Expanding the **primed field** in mass eigenstates **after EW SSB**

$$\begin{aligned} b'_L &= (V'_{dL})_{33} b_L + (V'_{dL})_{32} s_L + \dots && \simeq b_L + (V'_{dL})_{32} s_L \\ \mu'_L &= (V'_{eL})_{22} \mu_L + \dots && \simeq \mu_L \end{aligned}$$

Assuming $(V'_{eL})_{22} \approx (V'_{dL})_{33} \approx 1$ and $|(V'_{dL})_{32}| \approx |V_{ts}| \approx 0.04$